

Two step photo-ionization of a laser cooled and compressed thermal atomic beam for use in a focused ion beam

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Ultracold Atoms and their Prospects for High Resolution FIB Nanomachining

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Introduction

Goal: Design and fabrication of a focused ion beam (FIB) based on photoionization of a laser cooled and compressed thermal beam of rubidium to reach nm spot sizes

Applications: Imaging and altering structures at the nanometer length scale. By scanning the beam over a substrate material can be removed (milling), a precursor gas can be deposited or the substrate can be imaged by looking at secondary/back scattered electrons/ions.

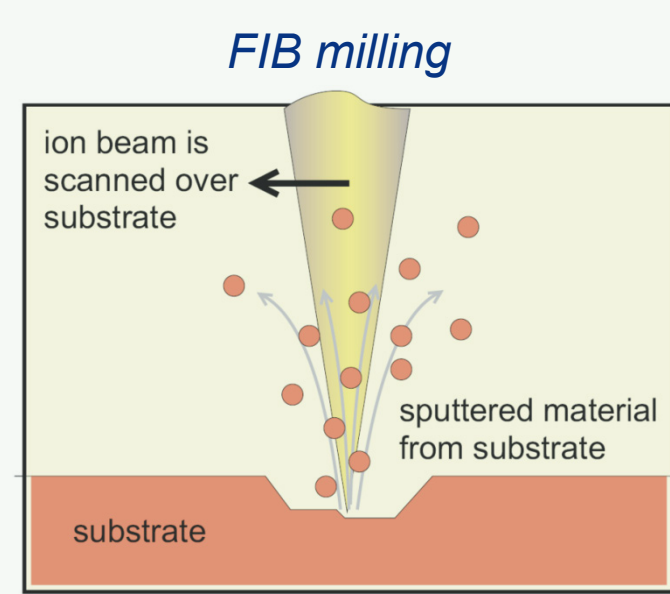
Other FIB sources:

- Liquid metal ion source (LMIS), uses heavy Gallium ions (good for milling), but only has resolution of 5 nm.
- Gas Field Ionization Source (GFIS), uses light helium or neon atoms, but has a high resolution of < 1 nm.

Important beam parameters:

- Reduced brightness B_r : current density per unit of area and solid angle
 - Longitudinal energy spread σ_U
- Higher B_r and lower σ_U give a smaller spot size.

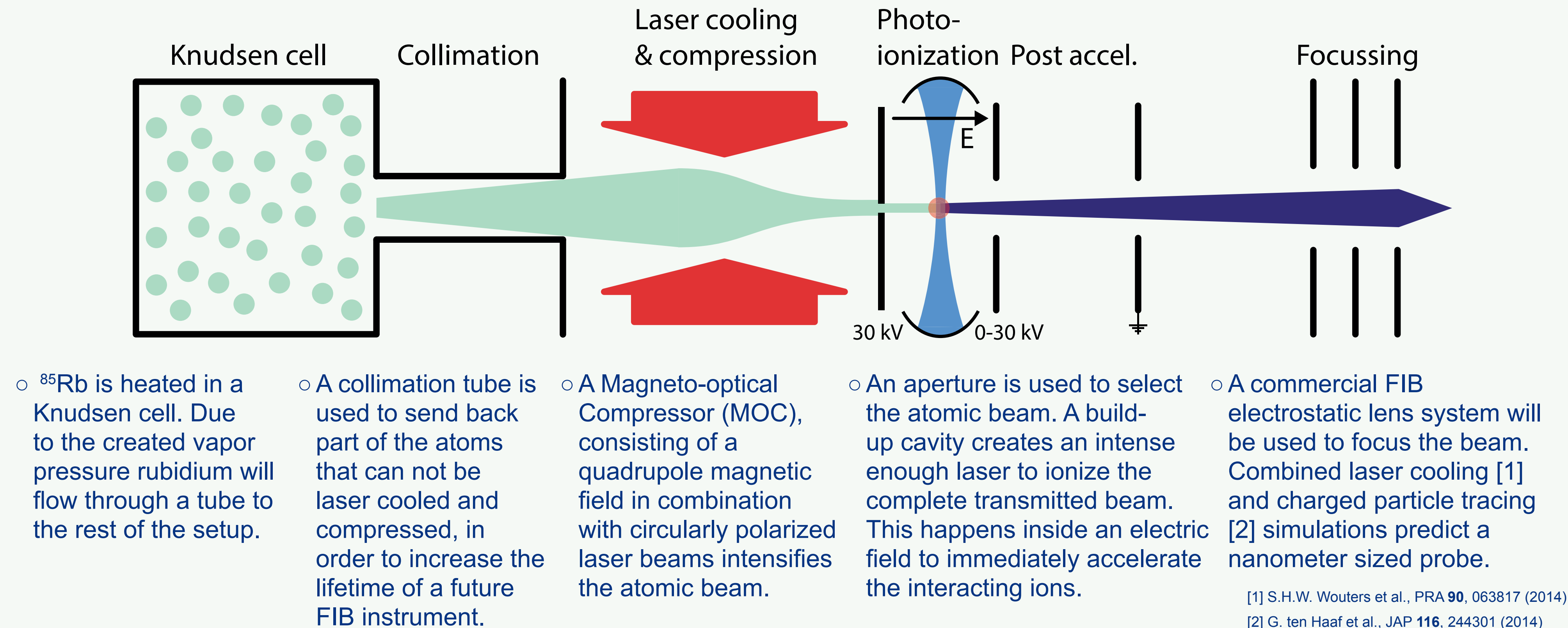
$$B_r = \frac{1}{U} \frac{\partial^2 I}{\partial A \partial \Omega}$$



Commercial (LMIS) FIB system:

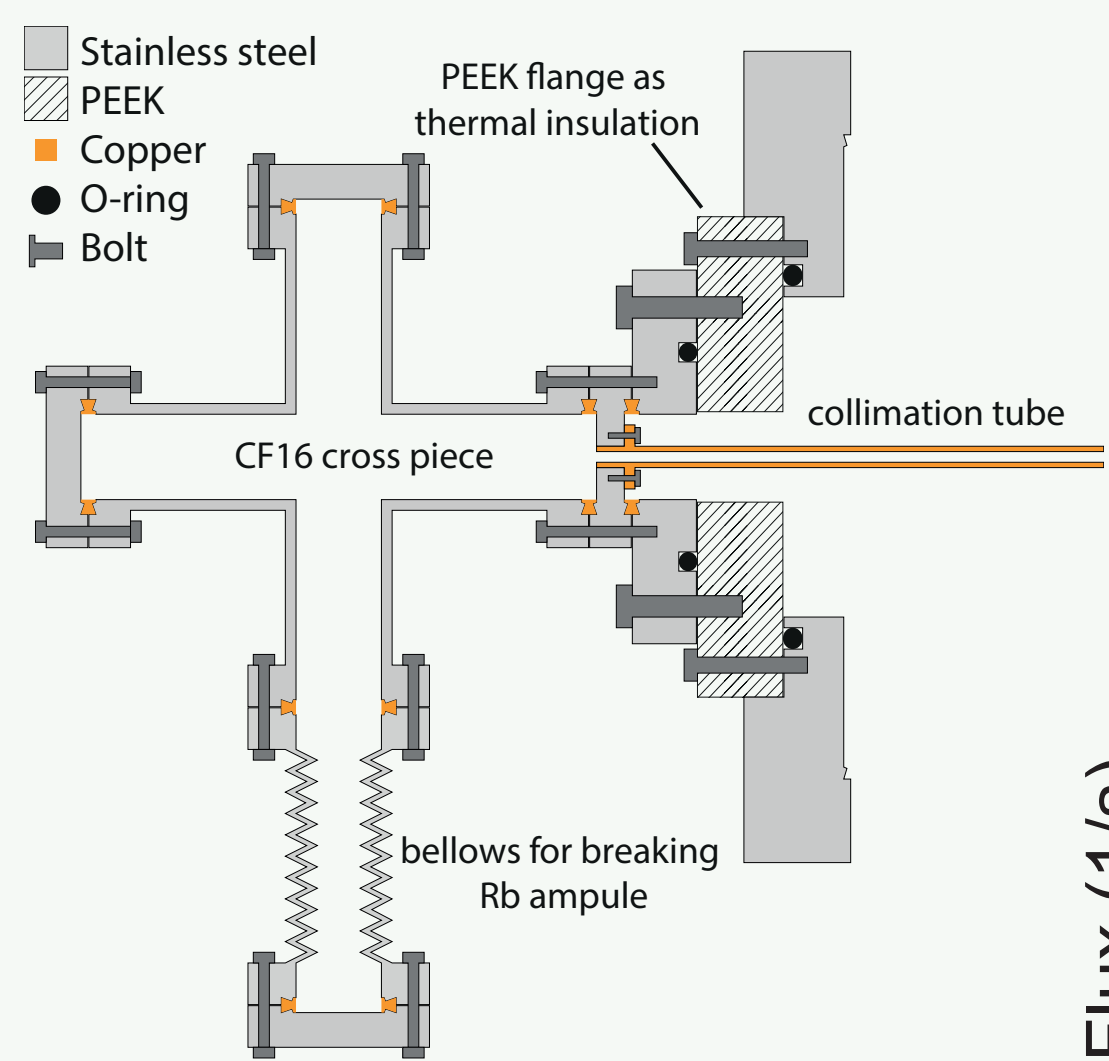


Proposed setup



Knudsen source: design and flux

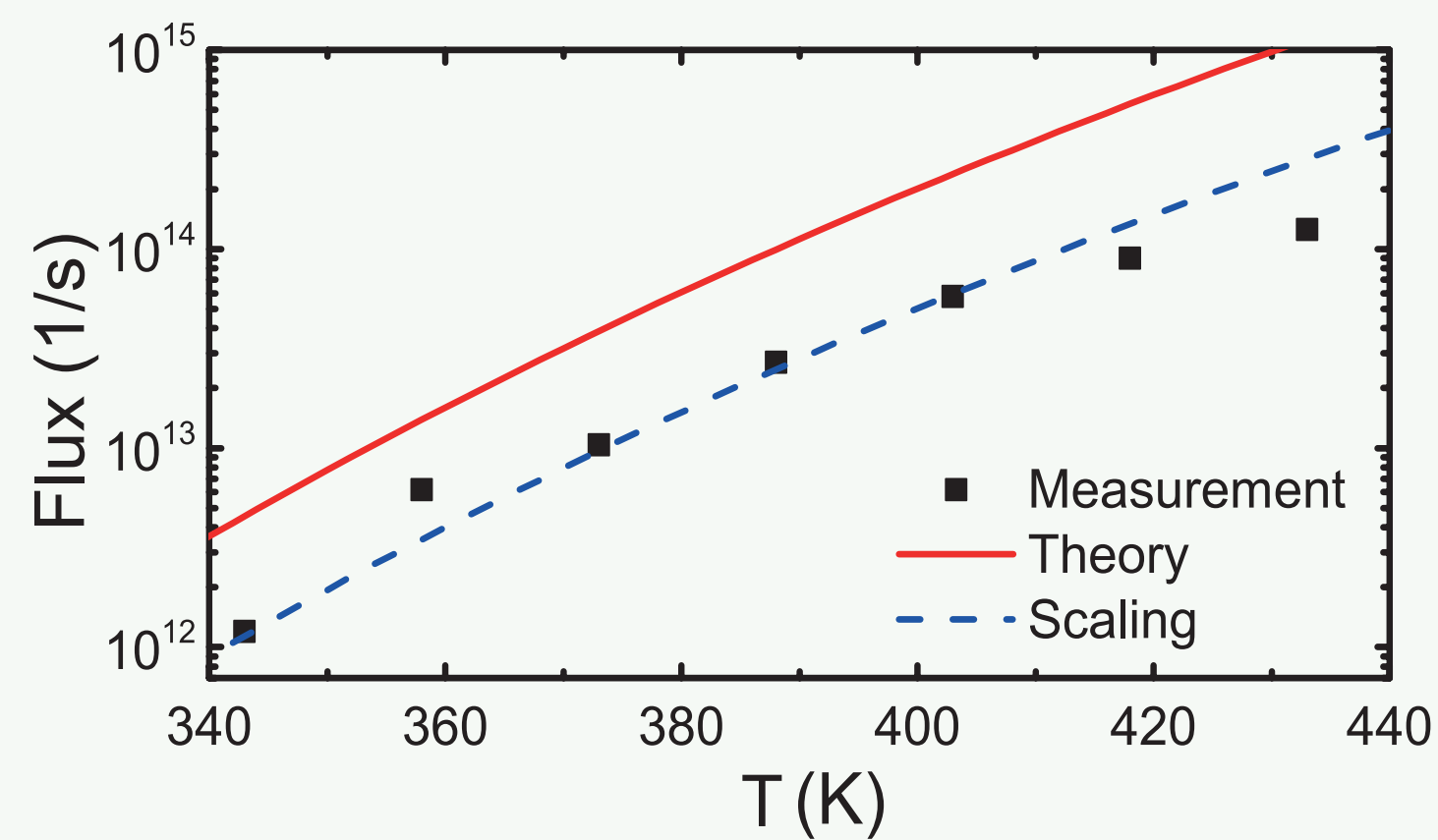
Design



- Bandheaters can heat the cross and bellows up to 160°C and an additional heating wire wrapped around tube heats the tube to 200°C to prevent clogging.
- Collimating tube design chosen over recirculating skimmed oven because of its simple design

Flux measurements

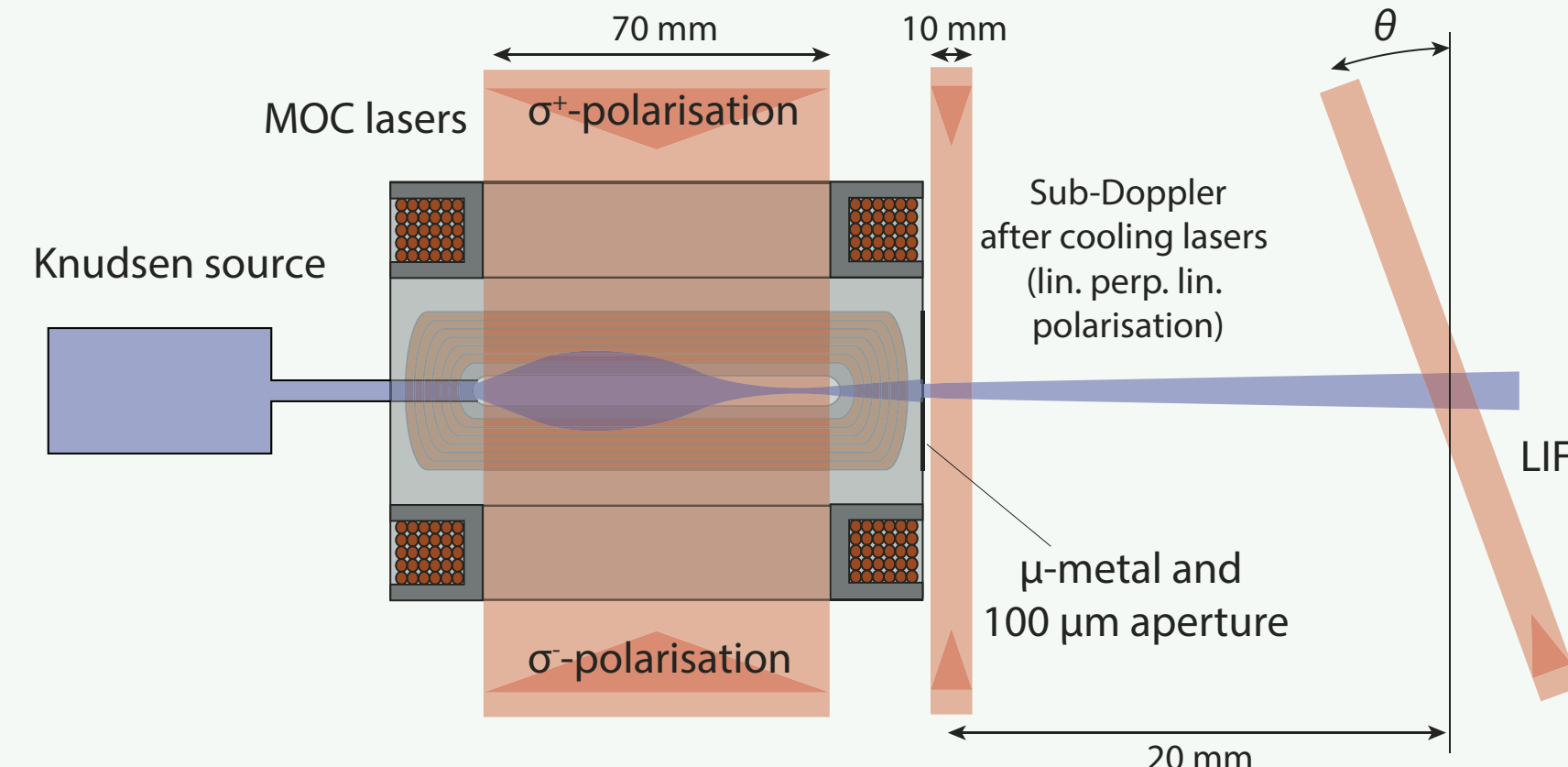
- Measured flux [3] a factor of 4 lower than predicted by molecular flow theory [4]. However, this is not a big problem since it is simply overwon by increasing the Knudsen cell temperature by 20 K.
- At higher temperatures (>400 K) collision in the tube cause the flux to decrease, hence the larger discrepancy with theory.



[3] S.H.W. Wouters et al., Rev. Sci. Instr. 87, 083305 (2016)
[4] D.R. Olander and V. Kruger, Molecular beam sources fabricated from multichannel arrays. iii. the exit density problem. J. Appl. Phys. 41, 2769 (1970)

Magneto-optical Compression (MOC)

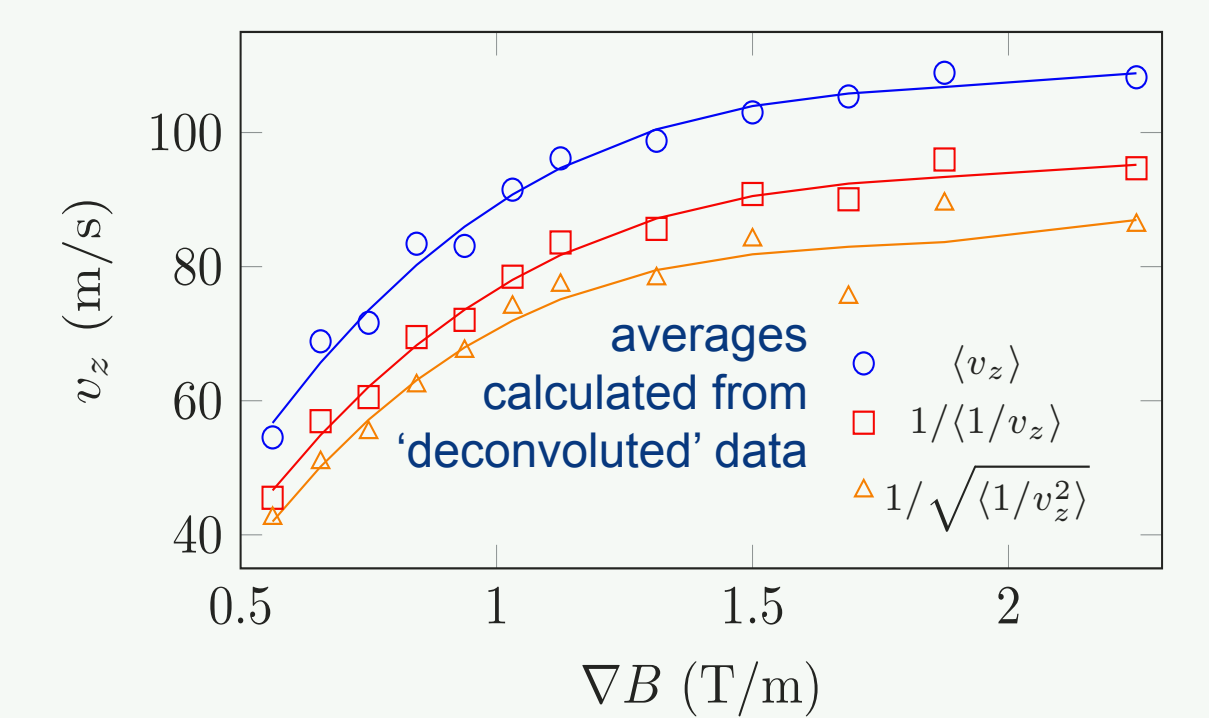
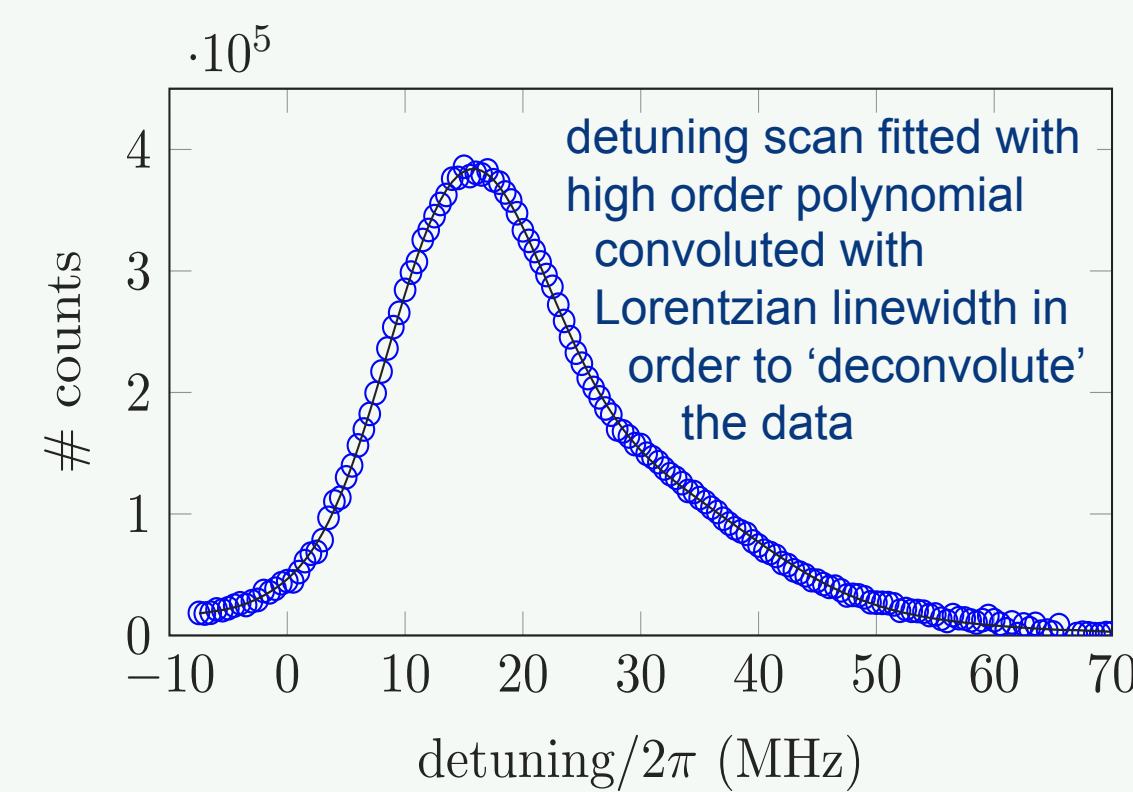
Experimental setup



- 100 μm diameter aperture to select 'useful' central part of the beam
- Imbalance in current through opposing quadrupole coils to overlap magnetic axis with aperture
- Lasers with linear perpendicular linear polarisation configuration were added to improve on the temperature of the beam
- Laser induced fluorescence (LIF) of atoms imaged on camera to measure beam flux and size
- Probe laser perpendicular to beam ($\theta = 0$) for temperature and flux measurements and $\theta = 0.3$ rad for longitudinal velocity distribution measurement.

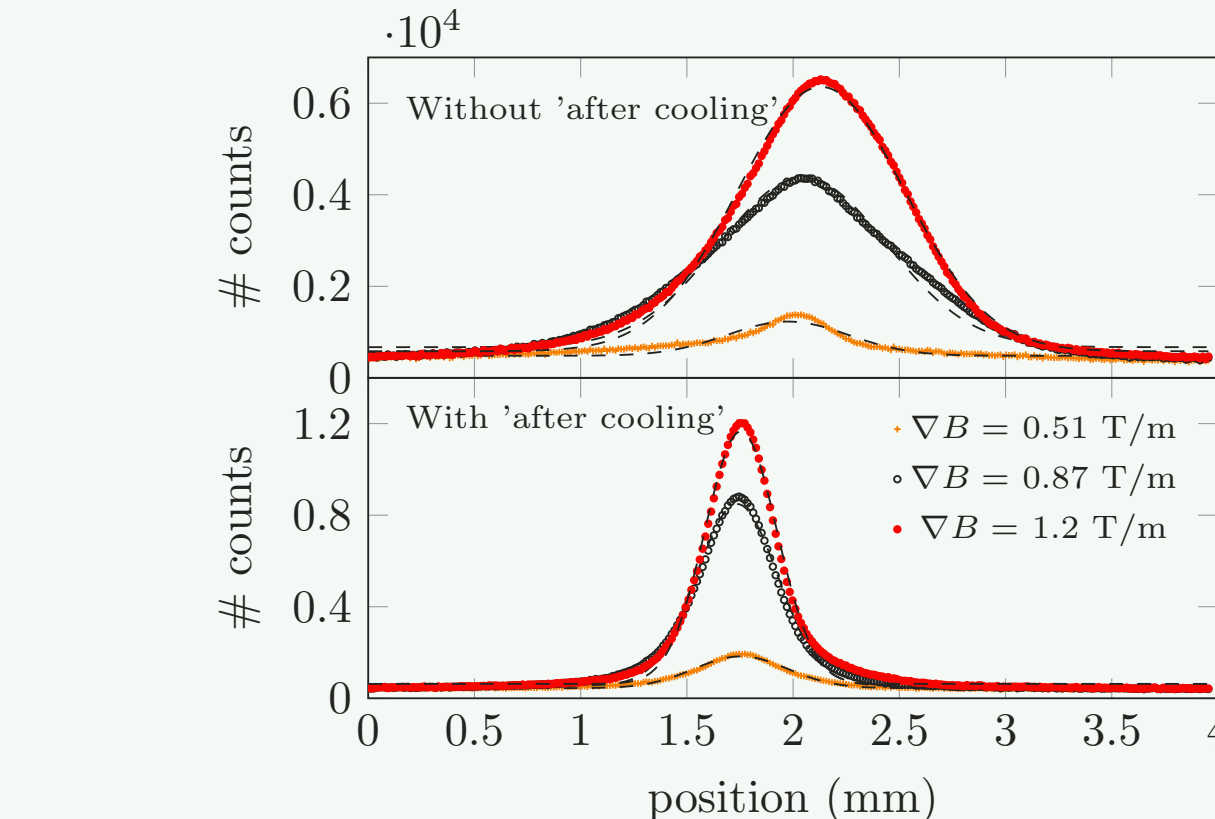
Longitudinal velocity distribution

- Longitudinal velocity required for accurate determination of flux (because of transit time through the probe laser) and temperature (because extracted from divergence)
- Longitudinal velocity of cooled beam lower than input from Knudsen source because slow atoms are captured more easily than fast atoms

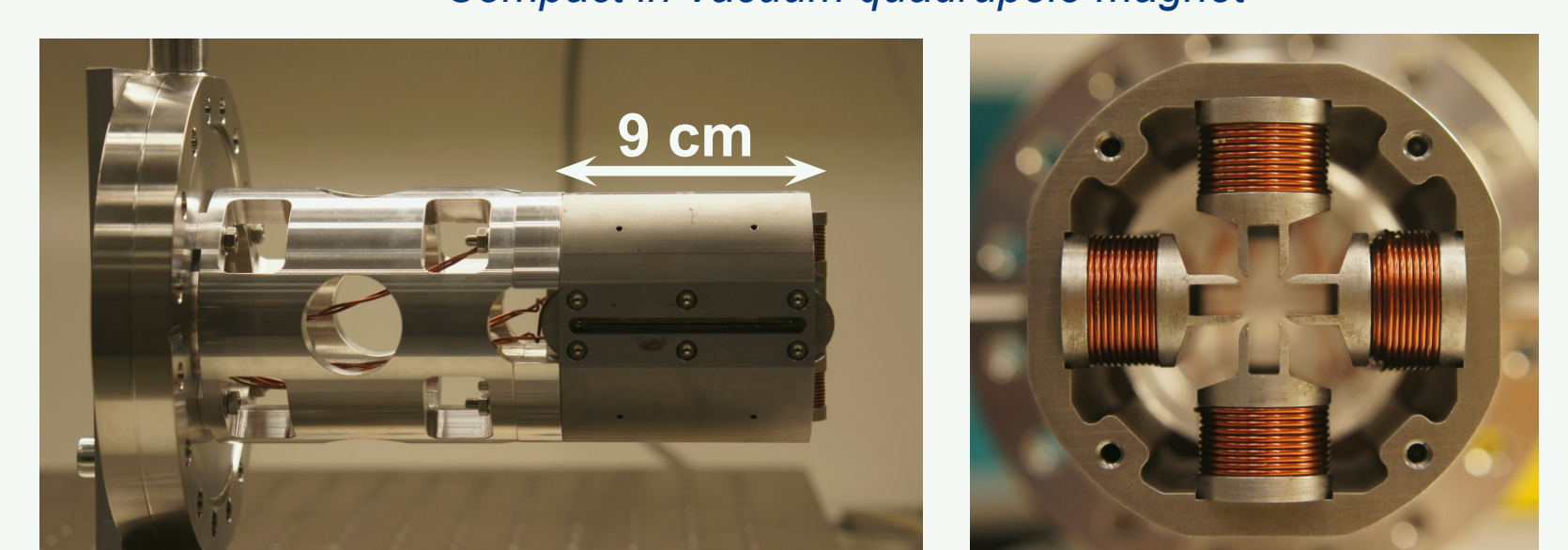
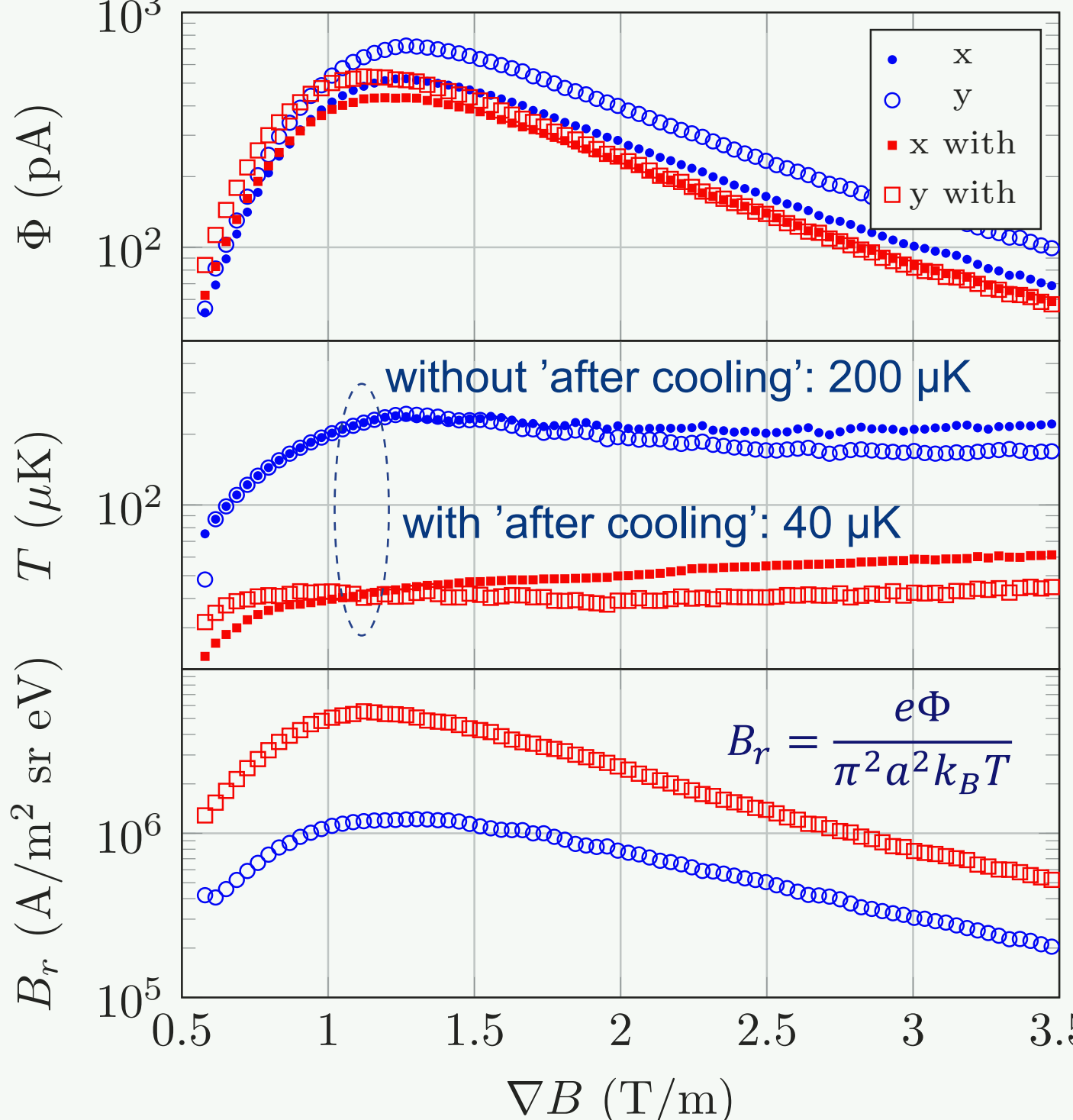


Equivalent reduced brightness

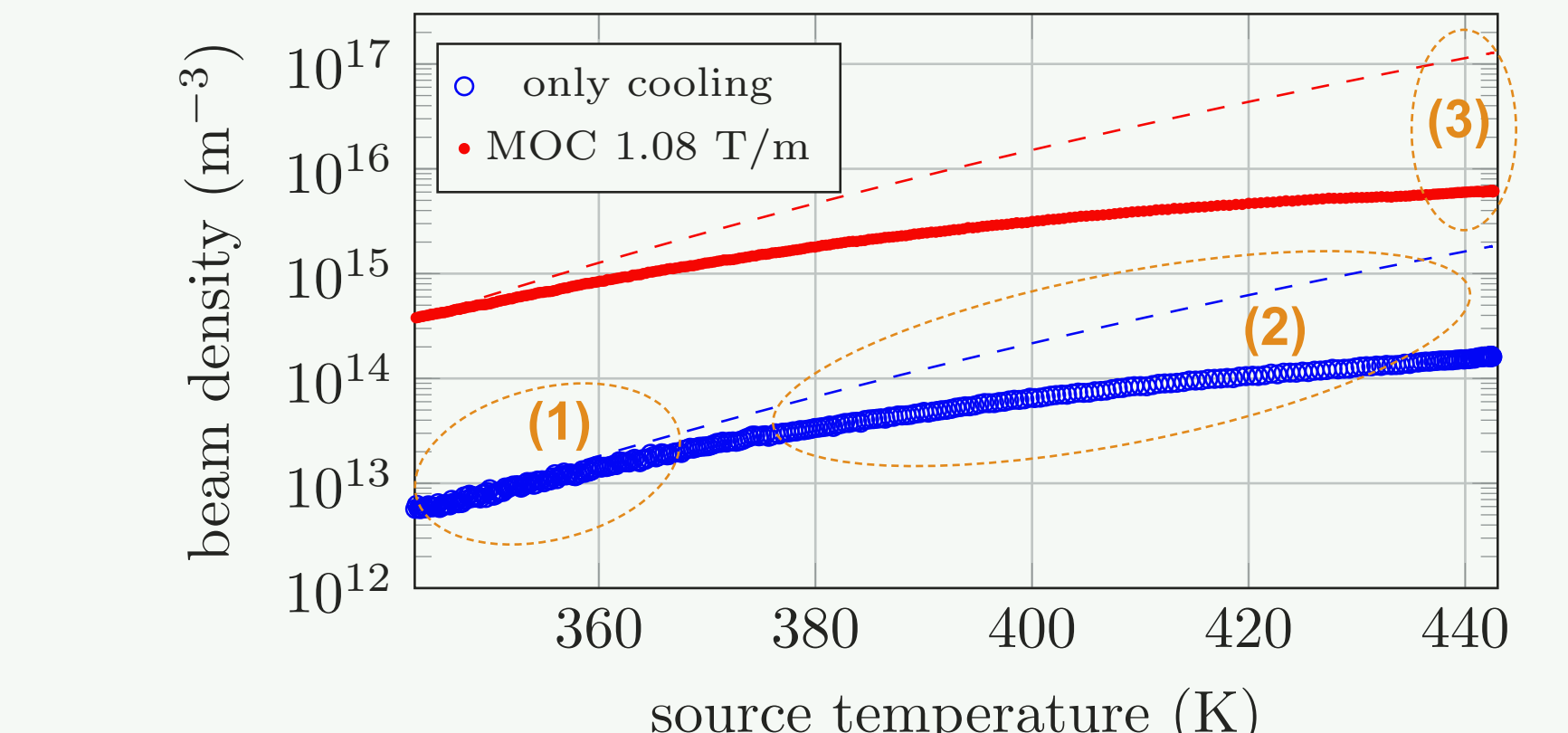
Beam profiles @ different magnetic gradients



Beam parameters vs. magnetic gradient



Beam density vs. source temperature



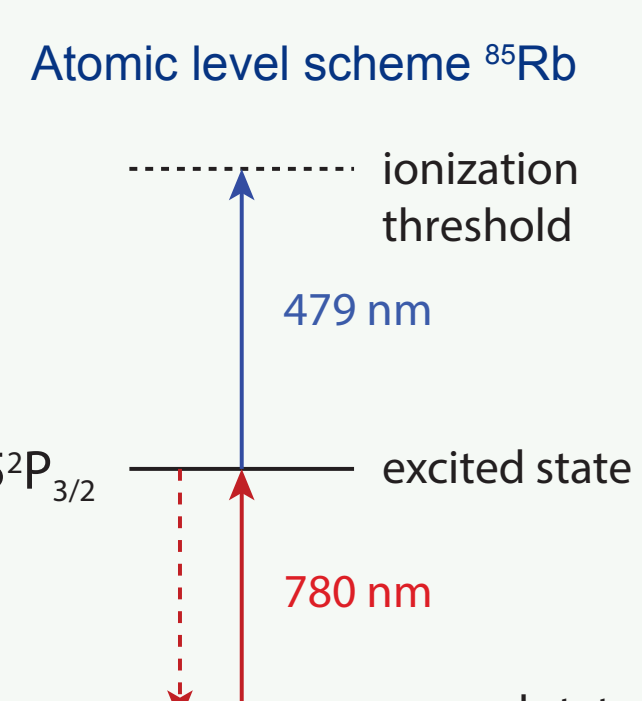
- Without interactions the density in the beam would scale according to the flux from the Knudsen source (dashed lines).
- At low source temperatures (1), scaling agrees with data when only cooling the beam, i.e. without compression.
- At higher temperatures (2), scaling doesn't hold anymore due to atomic collisions inside the collimation tube that lowers the amount of useable flux that can be captured in the cooled beam.
- When also compressing the beam the discrepancy with the scaling is larger (3), possibly due to radiation trapping or inelastic collisions between ground and excited state atoms.

Equivalent brightness of $6 \times 10^6 \text{ A/m}^2/\text{sr/eV}$ with 500 pA equivalent current!

Ionization: process and realization

Two-step ionization

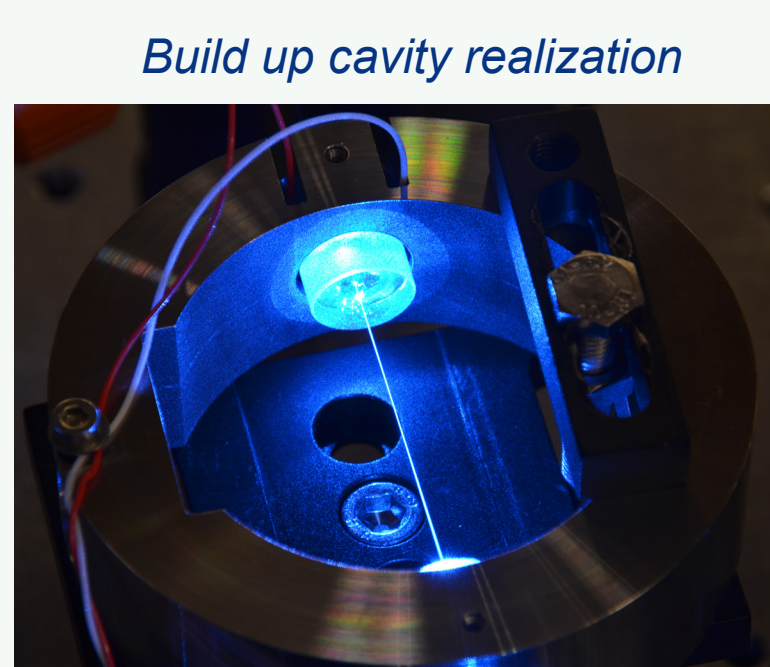
- Two lasers overlapped spatially at the position of the atomic beam.
- For largest brightness nearly complete ionization wanted



- Numerical calculations of the ionization process showed an ionization laser (479 nm) intensity of 10^{10} W/m^2 will ionize $\approx 80\%$ of the beam within approximately 3 μm

Build-up cavity

- Reflectance = 99.7%, max. theor. build up = $\frac{\text{circulating power}}{\text{incident power}} = \frac{1}{1-R} = 333$
- Cavity length chosen to give an 18 μm waist. Two lenses used for mode coupling.
- Pound Drever Hall [5] technique used to lock the length of the cavity to an integer times the half wavelength.
- Build up of (190 \pm 30) measured by locking the cavity and comparing the transmitted power through the cavity to the incident power.



- So 300 mW laser gives $2 \times 10^{10} \text{ W/m}^2$, enough for ionizing $\approx 80\%$.

Conclusion and outlook

- Flux from Knudsen cell scales as expected and is sufficient for our purpose.
- Magneto-optical compression experiments show that an equivalent reduced brightness of $10^6 \text{ A/m}^2/\text{sr/eV}$ can be achieved in a single 70 mm laser cooling stage, additional sub-Doppler cooling improves the equivalent reduced brightness to $6 \times 10^6 \text{ A/m}^2/\text{sr/eV}$.
- Build-up cavity is realized in which an ionization intensity of $2 \times 10^{10} \text{ W/m}^2$ can be reached, which can ionize 80% of the atomic beam.

Next steps:

- Ion beam creation in tailor made vacuum vessel to accommodate accelerator and build up cavity
- Energy spread analysis with retarding field analyzer (0.1 eV FWHM energy spread expected @ 1 pA)
- Mounting source on FIB column (FEI sidewinder) and determine ion beam brightness