# On the equivalence of certain ergodic properties for Gibbs states 

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# On the equivalence of certain ergodic properties for Gibbs states 

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#### Abstract

We extend our previous work by proving that for translation invariant Gibbs states on $\mathbb{Z}^{d}$ with a translation invariant interaction potential $\Psi=\left(\Psi_{A}\right)$ satisfying $\sum_{A \ni 0}|A|^{-1}[\operatorname{diam}(A)]^{d}| | \Psi_{A} \|<\infty$ the following hold: (1) the Kolmogorov-property implies a Trivial Full Tail; (2) the Bernoulli-property implies Følner Independence.

The existence of bilaterally deterministic Bernoulli Shifts tells us that neither (1) nor (2) is true for random fields in general without some further assumption (even when $d=1$ ).


## 1 Introduction

The purpose of this paper is to extend some results for Markov random fields, that were proved in [HS], to a large class of (possibly infinite range) Gibbs states. In §1 we give some notations and definitions. In $\S 2$ we formulate our theorems. In $\S 3$ and $\S 4$ we give proofs.
Notations and definitions. Throughout this paper we consider stationary stochastic processes $X=\left\{X_{x}\right\}_{x \in \mathbb{Z}^{d}}$ taking values in a finite set $F$. We also view $X$ as a probability measure $\mu$ on $\Omega=F^{\mathbb{Z}^{d}}$ that is invariant under the natural $\mathbb{Z}^{d}$-action.

We write $B_{n}=[-n, n]^{d} \cap \mathbb{Z}^{d}$ to denote the $n-$ box in $\mathbb{Z}^{d}$. If $\mu$ is a probability measure on $F^{\mathbb{Z}^{d}}$ and $A \subseteq \mathbb{Z}^{d}$, then we let $\mu_{A}$ denote the probability measure on $F^{A}$ obtained by projecting $\mu$ onto $A$. We also let $X_{A}$ denote the process restricted to $A$, so that $\mu_{A}$ is just the distribution of $X_{A}$.

[^0]In order to save space, rather than repeating verbatum a number of definitions we will frequently refer to [HS]. In particular, the reader can find there the definitions of the $\bar{d}$-distance between two probability measures $\mu_{A}$ and $\nu_{A}$ with finite $A$, entropy, ergodicity, K-automorphism (K), Trivial Full Tail (TFT), Bernoulli (B), Very Weak Bernoulli (VWB), and Følner Independence (FI).

For translation invariant ergodic random fields the following orderings hold (see [HS], $\S 1$ and Theorem 2.4 and references):

$$
\begin{aligned}
& F I \subsetneq V W B ; T F T \subsetneq K \\
& F I \subsetneq T F T ; V W B \subsetneq K \\
& B=V W B .
\end{aligned}
$$

A Gibbs state is defined as follows (see [G], Chapter 2). An interaction potential is a family $\Psi=\left(\Psi_{A}\right)$ of maps $\Psi_{A}: F^{A} \rightarrow \mathbb{R}$ satisfying

$$
\sum_{A: A \cap \Lambda \neq \emptyset}\left\|\Psi_{A}\right\|<\infty \text { for all } \Lambda \subseteq \mathbb{Z}^{d} \text { non-empty and finite }
$$

where $\left\|\Psi_{A}\right\|=\sup _{\eta \in F^{A}}\left|\Psi_{A}(\eta)\right|$ and where $A$ runs over the non-empty finite subsets of $\mathbb{Z}^{d}$. For a given $\Psi$, a Gibbs state for $\Psi$ is any random field $\mu$ whose conditional probabilities on $\Lambda$ given $\sigma$ on $\Lambda^{c}$ are of the form

$$
\mu(\cdot \mid \sigma)=\frac{1}{Z_{\Lambda, \sigma}} \exp \left[-H_{\Lambda}(\cdot \mid \sigma)\right] \text { for all } \Lambda \subseteq \mathbb{Z}^{d} \text { non-empty and finite and } \sigma \in F^{\Lambda^{c}}
$$

where $Z_{\Lambda, \sigma}$ is the normalizing constant (or partition sum),

$$
H_{\Lambda}(\eta \mid \sigma)=-\sum_{A: A \cap A \neq \emptyset} \Psi_{A}\left([\eta \vee \sigma]_{A}\right) \quad\left(\eta \in F^{\Lambda}\right)
$$

is the Hamiltonian on $\Lambda$ given $\sigma$ on $\Lambda^{c}$, and $[\eta \vee \sigma]_{A}$ is the configuration $\eta \vee \sigma$ restricted to $A$.

The class of interaction potentials that we allow in this paper are the ones satisfying

$$
(*) \quad\left\{\begin{array}{l}
\Psi_{A}=\Psi_{A+z} \text { for all } A \text { and all } z \in \mathbb{Z}^{d} \\
\sum_{A \ni 0} \frac{1}{|A|}[\operatorname{diam}(A)]^{d}\left\|\Psi_{A}\right\|<\infty,
\end{array}\right.
$$

where $\operatorname{diam}(A)=\sup _{x, y \in A}|x-y|_{1}$. The second of these conditions means that for large sets the total interaction across the boundary of the set is of the order of the surface of the set.

Despite the fact that the interaction potential is assumed to be translation invariant, there may - and in general will - be Gibbs states that are not translation invariant. In this paper, however, we only consider translation invariant Gibbs states.

## 2 Main theorems

The goal of this paper is to show that the converses of 'FI implies VWB' and 'TFT implies K', though not true in general (see [HS] for a discussion), are true for all $\mathbb{Z}^{d}$-invariant Gibbs states for interactions satisfying $(*)$. That is, we prove the following two theorems.

Theorem 2.1 If $\mu$ is a $\mathbb{Z}^{d}$-invariant Gibbs state for an interaction satisfying (*) and is VWB, then $\mu$ is FI.

Theorem 2.2 If $\mu$ is a $\mathbb{Z}^{d}$-invariant Gibbs state for an interaction satisfying (*) and is $K$, then $\mu$ is TFT.

The proofs of these theorems are given in $\S 3$ and $\S 4$. Thus, for the class $(*)$ we obtain the following ordering:

$$
(* *) \quad F I=V W B \subseteq T F T=K .
$$

## Remarks:

(1) For $d=1,(*)$ precisely coincides with the well-known sufficient condition for uniqueness of the Gibbs state ([G], p. 166). Being the unique Gibbs state, the measure is necessarily TFT ([G], Theorem 7.7(a)). So Theorem 2.2 is of no interest for this case. In fact, for $d=1,(*)$ is known to imply that the unique Gibbs state is Weak Bernoulli ([G], p. 461), which is stronger than FI. Therefore Theorem 2.1 is also of no interest in this case.
(2) Theorem 2.2 is trivial, for any $d \geq 1$, if all (!) Gibbs states for the given interaction are $\mathbb{Z}^{d}$-invariant. In fact, then ergodicity is already enough to imply TFT. The reason for this is that any such ergodic Gibbs state cannot be decomposed as a convex combination of two Gibbs states for the same interaction, since these would necessarily be $\mathbb{Z}^{d}$-invariant and by ergodicity would be identical. Hence, any such ergodic Gibbs state is extremal within the class of all Gibbs states, and therefore must be TFT (again by [G], Theorem 7.7(a)).
(3) In [OW1] it is proved that for the Ising model with ferromagnetic nearestneighbor interaction both the ' + state' and the ' - state' are B. So for this case all four properties in $(* *)$ hold. The proof shows that the same is true for all interactions satisfying the FKG lattice condition ([G], p. 445), the technical reason being that then the conditional measure in a finite set is stochastically increasing as a function of the configuration outside the set.
(4) As will become clear from the proofs, both Theorem 2.1 and Theorem 2.2 are statements of the type: if a certain property holds 'one-sided' then it also holds 'two-sided' (i.e., if the property holds with respect to the lexicographic past of a large box, then it holds with respect to the entire outside of the large box). In
the theory of Gibbs states similar types of statements occur, for instance, for the notions of Markov property ([G], Section 10.1) and entropy [E].
(5) An open question is whether TFT $=$ VWB for the class $(*)$. In [H] an example is constructed of a Markov random field on $\mathbb{Z}^{2}$ that is K but not VWB. Since [HS] shows that $\mathrm{K}=$ TFT for Markov random fields in general, this example violates TFT $=$ VWB. However, it is not Gibbsian (because it is not strictly positive on all cylinder sets). Perhaps a Gibbsian counterexample can be found in the class of nearest-neighbor 'clock models' [FS], where Gibbs states are known to exist that are unique and yet have arbitrarily slow decay of correlations.
(6) Another open question is whether ( $* *$ ) also holds for the larger class of interactions where the second condition in $(*)$ is weakened to $\sum_{A \ni 0}\left\|\Psi_{A}\right\|<\infty$, i.e., the usual summability condition.

## 3 Proof of Theorem 2.1

3.1 Key lemma. We will need the following property of a Gibbs state for an interaction satisfying $(*)$, which plays an important role in the proofs of both Theorem 2.1 and Theorem 2.2.

Lemma 3.1 Fix an interaction satisfying (*) and let $\mu$ be a $\mathbb{Z}^{d}$-invariant Gibbs state for this interaction. Then, given $\ell, m \in \mathbb{N}$ and $\delta>0$, there exists a $C(\ell, m, \delta)$, satisfying

$$
\lim _{\ell \rightarrow \infty} C(\ell, m, \delta)=1 \text { for fixed } m \text { and } \delta
$$

such that for any $k \in[\ell, m \ell] \cap \mathbb{N}$, any $\sigma, \sigma^{\prime} \in F^{B_{k}^{c}}$ that agree on $B_{k+\lfloor\delta \ell]} \backslash B_{k}$, and any $\eta \in F^{B_{k}}$, the following bounds hold a.s.:

$$
\frac{1}{C(\ell, m, \delta)} \leq \frac{\mu_{B_{k}}(\eta \mid \sigma)}{\mu_{B_{k}}\left(\eta \mid \sigma^{\prime}\right)} \leq C(\ell, m, \delta)
$$

Proof: Fix $m \in \mathbb{N}$ and $\delta>0$. For $k, \ell \in \mathbb{N}$, let $\mathcal{A}_{k, \ell, \delta}$ denote the collection of finite sets $A$ satisfying $A \cap B_{k} \neq \emptyset$ and $A \cap B_{k+\lfloor\delta \ell\rfloor}^{c} \neq \emptyset$. Given any finite set $A$, let $T_{A}(k, \ell, \delta)$ denote the number of translates of $A$ that are contained in $\mathcal{A}_{k, \ell, \delta}$. Some elementary combinatorial geometry (left to the reader) shows that there exists a $C_{1}(m, \delta)$ such that

$$
\sup _{A} \sup _{l \in \mathbb{N}} \sup _{k \in[\ell, m \ell] \cap \mathbb{N}} \frac{T_{A}(k, \ell, \delta)}{[\operatorname{diam}(A)]^{d}} \leq C_{1}(m, \delta)
$$

Next, for any $l \in \mathbb{N}$, any $k \in[\ell, m \ell] \cap \mathbb{N}$, any $\sigma, \sigma^{\prime} \in F^{B_{k}^{c}}$ that agree on $B_{k+\lfloor\delta \ell\rfloor} \backslash B_{k}$, and any $\eta \in F^{B_{k}}$, we have

$$
\begin{aligned}
\left|H_{B_{k}}(\eta \mid \sigma)-H_{B_{k}}\left(\eta \mid \sigma^{\prime}\right)\right| & \leq \sum_{A \in \mathcal{A}_{k, \ell, \delta}}\left\|\Psi_{A}\right\| \\
& =\sum_{A \ni 0} \frac{T_{A}(k, \ell, \delta)}{|A|}\left\|\Psi_{A}\right\| \\
& \leq C_{1}(m, \delta) \sum_{A \ni 0, \operatorname{diam}(A) \geq\lfloor\delta \ell\rfloor} \frac{1}{|A|}[\operatorname{diam}(A)]^{d}\left\|\Psi_{A}\right\| .
\end{aligned}
$$

By assumption (*), the sum in the right-hand side tends to zero as $\ell \rightarrow \infty$. Hence there exists a $C_{2}(\ell, m, \delta)$, satisfying $\lim _{\ell \rightarrow \infty} C_{2}(\ell, m, \delta)=1$ for fixed $m$ and $\delta$, such that

$$
\frac{1}{C_{2}(\ell, m, \delta)} \leq \frac{e^{-H_{B_{k}}(\eta \mid \sigma)}}{e^{-H_{B_{k}}\left(\eta \mid \sigma^{\prime}\right)}} \leq C_{2}(\ell, m, \delta)
$$

for any $l, k, \sigma, \sigma^{\prime}, \eta$ as above. These inequalities being true for all $\eta$, the ratio of the corresponding partition functions also satisfies the exact same inequalities. This proves the claim with $C(\ell, m, \delta)=C_{2}(\ell, m, \delta)^{2}$.
3.2 Proof of Theorem 2.1. If a process is VWB, then it is B (see $\S 1$ ). The latter is in turn equivalent to the following condition, called extremality (see [HS], $\S 3$ and references).

Definition 3.2 $A \mathbb{Z}^{d}$-invariant probability measure $\nu$ is called EXTREMAL if for all $\epsilon>0$ there exist an $N \in \mathbb{N}$ and a $\delta>0$ such that: for all $n \geq N$ and for all decompositions of $\nu_{B_{n}}$ of the form

$$
\nu_{B_{n}}=\sum_{i=1}^{M} p_{i} \nu_{i}
$$

with $\left(p_{1}, \ldots, p_{M}\right)$ a probability vector and $M \leq 2^{\delta\left|B_{n}\right|}$, most of the $\nu_{i}$ 's are $\bar{d}$-close to $\nu_{B_{n}}$ in the sense that

$$
\sum_{i: \bar{d}\left(\nu_{B_{n}}, \nu_{i}\right)<\epsilon} p_{i}>1-\epsilon .
$$

In words, any 'not too large' decomposition of the measure on large blocks must have almost all components close to the original measure.

We must contrast extremality with the definition of FI, which reads:
Definition 3.3 $A \mathbb{Z}^{d}$-invariant probability measure $\nu$ is called Følner IndePENDENT if for all $\epsilon>0$ there exists an $N \in \mathbb{N}$ such that: if $n \geq N$ and $S \subseteq B_{n}^{c}$ with $S$ finite, then

$$
\bar{d}\left(\mu_{B_{n}}, \mu_{B_{n}} / \sigma\right)<\epsilon
$$

for all $\sigma \in F^{S}$ except for an $\epsilon-$ portion as measured by $\mu$ (where $\mu_{B_{n}} / \sigma$ denotes $\mu_{B_{n}}$ conditioned on $\sigma$ ).

In words, for large $n$ and for most configurations on $B_{n}^{c}$ the conditional distribution on $B_{n}$ is $\bar{d}$-close to the unconditional distribution.

To show that ' $\mu$ is B ' implies ' $\mu$ is FI ', let $\epsilon>0$ and pick $N_{1}, \delta$ from Definition 3.2. Next, choose $\gamma>0$ sufficiently small and pick $N_{2}$ such that $|F|^{\left|B_{n+\lfloor\gamma n\rfloor} \backslash B_{n}\right|} \leq$ $2^{\delta\left|B_{n}\right|}$ for all $n \geq N_{2}$. Next, pick $N_{3}$ from Lemma 3.1 such that $C(n, 1, \gamma) \leq 1+\epsilon$ for all $n \geq N_{3}$. For such $n$, it follows readily from the bounds in Lemma 3.1 that, for any $\sigma, \sigma^{\prime} \in F^{B_{n}^{c}}$ that agree on $B_{n+\lfloor\gamma n\rfloor} \backslash B_{n}$, the measures $\mu_{B_{n}}(\cdot \mid \sigma)$ and $\mu_{B_{n}}\left(\cdot \mid \sigma^{\prime}\right)$ are within $\epsilon$ in total variation distance.

By Lemma 3.2 in [HS], to verify the FI-condition in Definition 3.3 it suffices to consider $n \geq \max \left\{N_{1}, N_{2}, N_{3}\right\}$ and finite sets $S \subseteq B_{n}^{c}$ that contain $B_{n+\lfloor\sim n\rfloor} \backslash B_{n}$. Since $|F|^{\left.\mid B_{n+\mid \gamma n}\right\rfloor B_{n} \mid} \leq 2^{\delta\left|B_{n}\right|}$, extremality yields that there exist configurations $\eta_{1}, \ldots \eta_{M}$ on $B_{n+\lfloor\gamma n\rfloor} \backslash B_{n}$, with $M \leq\left.|F|\right|^{\left.\left|B_{n+\lfloor }\right| \gamma n\right\rfloor} \backslash B_{n} \mid$, such that their total measure is at least $1-\epsilon$ and such that also $\bar{d}\left(\mu_{B_{n}}, \mu_{B_{n}} / \eta_{i}\right)<\epsilon$ for each $\eta_{i}$.

Now consider all configurations $\sigma$ on $S$ such that the restriction of $\sigma$ to $B_{n+\lfloor\gamma n\rfloor} \backslash B_{n}$ is $\eta_{i}$ for some $i \in\{1, \ldots, M\}$. Clearly, these configurations have total measure at least $1-\epsilon$, and so we need only show that for each such $\sigma$,

$$
\bar{d}\left(\mu_{B_{n}}, \mu_{B_{n}} / \sigma\right)<2 \epsilon .
$$

For this it suffices to show that

$$
\bar{d}\left(\mu_{B_{n}} / \eta, \mu_{B_{n}} / \sigma\right)<\epsilon
$$

whenever $\sigma$ is a configuration on $S$ whose restriction to $B_{n+\left\lfloor\gamma_{n}\right\rfloor} \backslash B_{n}$ is $\eta$. However, $\mu_{B_{n}} / \eta$ and $\mu_{B_{n}} / \sigma$ are each averages of measures that, as we saw earlier, are all within $\epsilon$ in total variation distance of each other. Hence $\mu_{B_{n}} / \eta$ and $\mu_{B_{n}} / \sigma$ are within $\epsilon$ in total variation distance, and therefore also within $\epsilon$ in $\bar{d}$-distance.

## 4 Proof of Theorem 2.2

We will prove the result only for $d=2$, the extension to higher dimensions being straightforward. The proof is a variation on the proof of the analogous statement for Markov random fields given in [HS]. The main point is to implement Lemma 3.1, which requires some estimates.

TFT means that the $\sigma$-algebra $T$ defined by

$$
\begin{aligned}
& T=\cap_{m \geq 1} T_{m} \\
& T_{m}=\sigma\left(X_{x}, x \in B_{m}^{c}\right)
\end{aligned}
$$

is trivial. On the other hand, K is equivalent to the smaller $\sigma$-algebra $T^{\prime}$ defined by

$$
\begin{aligned}
& T^{\prime}=\sigma\left(\cup_{m \geq 1} T_{m}^{\prime}\right) \\
& T_{m}^{\prime}=\cap_{n \geq 1} T_{m, n}^{\prime} \\
& T_{m, n}^{\prime}=\sigma\left(X_{x}, x \in\left\{\left(x_{1}, x_{2}\right): x_{2} \leq-n \text { or }\left(x_{1} \leq-n \text { and } x_{2} \leq m\right)\right\}\right)
\end{aligned}
$$

( $T_{m, n}^{\prime}$ is the lexicographic past of the rectangle $[-n, n] \times[-n, m]$ in $\mathbb{Z}^{2}$ ) being trivial (see [HS], $\S 1$ and references). We will show that $T=T^{\prime}$ a.s, which more than implies the claim that $\mathrm{K}=$ TFT.

In order to do so, we appeal to Lemma 2.10 in [BH] (which is stated there only for $d=1$ but whose proof for higher dimensions is identical). According to this lemma, since $T^{\prime} \subseteq T$ it suffices to show that

$$
h\left(X_{B_{n}} \mid T^{\prime}\right)=h\left(X_{B_{n}} \mid T\right) \quad \text { for all } n \geq 0
$$

where $h(\cdot \mid \cdot)$ denotes conditional entropy.
Fix $n \geq 0$. Since $T_{n}^{\prime} \subseteq T^{\prime} \subseteq T$, it suffices to show that

$$
\begin{equation*}
h\left(X_{B_{n}} \mid T_{n}^{\prime}\right) \leq h\left(X_{B_{n}} \mid T\right) . \tag{1}
\end{equation*}
$$

To achieve this, we will show that there exists a function $\Delta(k, \ell, \delta) \geq 0$, defined for $k, \ell \in \mathbb{N}$ with $k>2 n$ and for $\delta>0$, satisfying

$$
\begin{equation*}
\lim _{\ell \rightarrow \infty} \frac{\Delta(k, \ell, \delta)}{(2 \ell+1)^{2}}=0 \text { for fixed } k \text { and } \delta \tag{2}
\end{equation*}
$$

such that

$$
\begin{equation*}
h\left(X_{B_{n}} \mid T_{n, k-n}^{\prime}\right) \leq h\left(X_{B_{n}} \mid T_{k(2 \ell+1)-n}\right)+\alpha_{k, \ell, \delta} h\left(X_{0}\right)+\frac{\Delta(k, \ell, \delta)}{(2 \ell+1)^{2}}, \tag{3}
\end{equation*}
$$

where $h(\cdot)$ denotes entropy and

$$
\alpha_{k, \ell, \delta}=\frac{\lfloor\delta \ell\rfloor(6 r-1)+\lfloor\delta \ell\rfloor(\lfloor\delta \ell\rfloor+1)}{(2 \ell+1)^{2}} \text { with } r=k(\ell+1)-n .
$$

Assuming the latter, we can let $\ell \rightarrow \infty, \delta \rightarrow 0, k \rightarrow \infty$ (in this order) in (3) and use (2) to obtain (1). Note that $\alpha_{k, \ell, \delta}$ vanishes in this limit and that, by the backwards martingale convergence theorem, the two entropies in (3) converge to the two entropies in (1).

To construct $\Delta(k, \ell, \delta)$, we define

$$
C_{k, \ell}=\cup_{x, y:}|x| \leq \ell,|y| \leq \ell\left\{B_{n}+(k x, k y)\right\}
$$

and note that the $(2 \ell+1)^{2}$ translates of $B_{n}$ comprising $C_{k, \ell}$ are disjoint and have distance at least $k-2 n$ between them. Let $r=k(\ell+1)-n$ as above and define

$$
\begin{aligned}
& E_{r}=\{(i, j): j<-r\} \\
& D_{r, \delta}=B_{r+\lfloor\delta \ell\rfloor} \backslash\left(B_{r} \cup E_{r}\right) .
\end{aligned}
$$

In words, $E_{r}$ is the lower half plane adjacent to the bottom segment of the boundary of $B_{r}$, while $D_{r, \delta}$ consists of $\lfloor\delta \ell\rfloor$ layers adjacent to the left, right and top
segments of the boundary of $B_{r}$. Note that the boundary of $B_{r}$ encloses $C_{k, \ell}$ and is a distance $k-2 n$ away from it.

We next order the $(2 \ell+1)^{2}$ translates of $B_{n}$ in $C_{k, \ell}$ lexicographically. Namely, we say that $B_{n}+(x, y)$ precedes $B_{n}+\left(x^{\prime}, y^{\prime}\right)$ if $y<y^{\prime}$ or $\left(y=y^{\prime}\right.$ and $\left.x<x^{\prime}\right)$. In this way, we get an ordering of the translates of $B_{n}$, which we enumerate as $B^{1}, B^{2}, \ldots B^{(2 \ell+1)^{2}}$. The idea of the proof is to compute the conditional entropy

$$
(\dagger)=h\left(X_{D_{r, \delta}} \vee X_{C_{k, \ell}} \mid X_{E_{r}}\right)
$$

in two different ways, to derive an upper, respectively, lower bound for the two resulting expressions, and in this way obtain an inequality between these bounds. This inequality will then be exploited to complete the proof.

For the lower bound, we estimate

$$
\begin{aligned}
(\dagger) \geq h\left(X_{C_{k, \ell}} \mid X_{E_{r}}\right) & =h\left(\vee_{i=1}^{(2 \ell+1)^{2}} X_{B^{i}} \mid X_{E_{r}}\right) \\
& =\sum_{i=1}^{(2 \ell+1)^{2}} h\left(X_{B^{i}} \mid X_{E_{r}} \vee X_{B^{1} \cup \ldots \cup B^{i-1}}\right) .
\end{aligned}
$$

Clearly, each of the terms in the sum is bounded below by $h\left(X_{B_{n}} \mid T_{n, k-n}^{\prime}\right)$, because the distance between the translates $B^{i}$ is $k-2 n$ and so is the distance between $U_{i} B_{i}$ and $E_{r}$. Hence

$$
\begin{equation*}
(\dagger) \geq(2 \ell+1)^{2} h\left(X_{B_{n}} \mid T_{n, k-n}^{\prime}\right) . \tag{4}
\end{equation*}
$$

For the upper bound, we write

$$
(\dagger)=h\left(X_{D_{r, \delta}} \mid X_{E_{r}}\right)+h\left(X_{C_{k, \ell}} \mid X_{E_{r}} \vee X_{D_{r, \delta}}\right)
$$

The first term is at most $\left|D_{r, \delta}\right| h\left(X_{0}\right)$, where $\left|D_{r, \delta}\right|=\sum_{i=1}^{\lfloor\delta\rfloor}(6 r-1+2 i)=\lfloor\delta \ell\rfloor(6 r-$ $1)+\lfloor\delta \ell\rfloor(\lfloor\delta \ell\rfloor+1)$. We express the second term as $h\left(X_{C_{k, \ell}} \mid X_{B_{r}^{c}}\right)+\Delta(k, \ell, \delta)$ with

$$
\Delta(k, \ell, \delta)=h\left(X_{C_{k, \ell}} \mid X_{E_{r}} \vee X_{D_{r, \delta}}\right)-h\left(X_{C_{k, \ell}} \mid X_{B_{r}^{c}}\right) \geq 0
$$

(the inequality coming from $\left.E_{r} \cup D_{r, \delta} \subseteq B_{r}^{c}\right)$. We develop $h\left(X_{C_{k, \ell}} \mid X_{B_{r}^{c}}\right)$ as

$$
\begin{aligned}
h\left(X_{C_{k, \ell}} \mid X_{B_{r}^{c}}\right) & =h\left(V_{i=1}^{(2 \ell+1)^{2}} X_{B^{i}} \mid X_{B_{r}^{c}}\right) \\
& \leq \sum_{i=1}^{(2 \ell+1)^{2}} h\left(X_{B^{i}} \mid X_{B_{r}^{c}}\right) \\
& \leq(2 \ell+1)^{2} h\left(X_{B_{n}} \mid T_{2 r-(k-n)}\right)
\end{aligned}
$$

using the fact that the largest distance between the boundary of $B_{r}$ and the center of a translate $B^{i}$ is $2 r-(k-n)$. Thus

$$
\begin{equation*}
(\dagger) \leq(2 \ell+1)^{2} h\left(X_{B_{n}} \mid T_{2 r-(k-n)}\right)+\left|D_{r, \delta}\right| h\left(X_{0}\right)+\Delta(k, \ell, \delta) . \tag{5}
\end{equation*}
$$

Comparing (4) and (5), noting that $2 r-(k-n)=k(2 \ell+1)-n$ and dividing by $(2 \ell+1)^{2}$, we obtain (3). Hence we need only verify (2) with the above definition of $\Delta(k, \ell, \delta)$.

To achieve the latter, we need the following trivial lemma.
Lemma 4.1 Let $p=\left\{p_{i}\right\}_{i \in I}$ and $q=\left\{q_{i}\right\}_{i \in I}$ be two finite probability vectors satisfying

$$
\frac{1}{C} \leq \frac{p_{i}}{q_{i}} \leq C \text { for all } i \in I
$$

Then $h(q) \geq-\log C+\frac{1}{C} h(p)$, where $h(\cdot)$ denotes entropy.
Proof: Write

$$
h(p)=\sum_{i} p_{i} \log \left(\frac{1}{p_{i}}\right) \leq \sum_{i} C q_{i} \log \left(\frac{C}{q_{i}}\right)=C \log C+C h(q) .
$$

We want to apply Lemma 4.1 when $p$ is the conditional law of $X_{C_{k, \ell}}$ given $X_{E_{r}} \vee X_{D_{r, \delta}}$ and $q$ is the conditional law of $X_{C_{k, \ell}}$ given $X_{B_{r}^{c}}$. Fix $k$ and $\delta$. Applying Lemma 3.1 and averaging over the configuration in $B_{r} \backslash C_{k, \ell}$, we find that there exists a $C(\ell)$ (namely, $C(\ell)=C(\ell, 2 k, \delta)$ in the notation of Lemma 3.1 because $k l \leq r \leq 2 k l$ ), satisfying

$$
\lim _{\ell \rightarrow \infty} C(\ell)=1
$$

such that for any $\ell \in \mathbb{N}$, any $\eta \in F^{C_{k, \ell}}$, any $\sigma \in F^{E_{r} \cup D_{r, \delta}}$ and any $\sigma^{\prime} \in F^{B_{r}^{c}}$ whose restriction to $E_{r} \cup D_{r, \delta}$ is $\sigma$, the following bounds hold a.s.:

$$
\frac{1}{C(\ell)} \leq \frac{\mu\left(X_{C_{k, \ell}}=\eta \mid X_{E_{r}} \vee X_{D_{r, \delta}}=\sigma\right)}{\mu\left(X_{C_{k, \ell}}=\eta \mid X_{B_{r}^{c}}=\sigma^{\prime}\right)} \leq C(\ell)
$$

(use that $B_{r+\lfloor\delta \ell\rfloor} \backslash B_{r} \subseteq E_{r} \cup D_{r, \delta} \subseteq B_{r}^{c}$ ). Using Lemma 4.1, we now obtain (integrate over $\eta, \sigma, \sigma^{\prime}$ )

$$
h\left(X_{C_{k, \ell}} \mid X_{B_{r}^{c}}\right) \geq-\log C(\ell)+\frac{1}{C(\ell)} h\left(X_{C_{k, \ell} \mid} \mid X_{E_{r}} \vee X_{D_{r, \delta}}\right)
$$

and so

$$
\begin{aligned}
0 \leq \Delta(k, \ell, \delta) & =h\left(X_{C_{k, \ell}} \mid X_{E_{r}} \vee X_{D_{r, \delta}}\right)-h\left(X_{C_{k, \ell}} \mid X_{B_{r}^{c}}\right) \\
& \leq \log C(\ell)+\left(1-\frac{1}{C(\ell)}\right) h\left(X_{C_{k, \ell} \mid} X_{E_{r}} \vee X_{D_{r, \delta}}\right) .
\end{aligned}
$$

But $h\left(X_{C_{k, \ell}} \mid X_{E_{r}} \vee X_{D_{r, \delta}}\right)$ can be bounded above by $(2 \ell+1)^{2} h\left(X_{B_{n}}\right)$. Hence (2) follows because $\lim _{\ell \rightarrow \infty} C(\ell)=1$.
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## References

[BH] Berbee, H.C.P. and den Hollander, W.Th.F., Tail triviality for sums of stationary random variables, Ann. Probab. 17 (1989) 1635-1645.
[E] van Enter, A., On a question of Bratteli and Robinson, Lett. Math. Phys. 6 (1982) 289-291.
[FS] Fröhlich, J. and Spencer, T., The Kosterlitz-Thouless transition in twodimensional Abelian spin systems and the Coulomb gas, Commun. Math. Phys. 81 (1981) 527-602.
[G] Georgii, H.-O., Gibbs Measures and Phase Transitions, de Gruyter 1988, New York.
[H] Hoffman, C., A Markov random field which is $K$ but not Bernoulli, Israel J. Math. (1998), to appear.
[HS] den Hollander, F. and Steif, J.E., On K-automorphisms, Bernoulli shifts and Markov random fields, Ergod. Th. Dynam. Sys. 17 (1997) 405-415.
[OW1] Ornstein, D.S. and Weiss, B., $\mathbb{Z}^{d}$-actions and the Ising model, unpublished manuscript (1973).
[OW2] Ornstein, D.S. and Weiss, B., Every transformation is bilaterally deterministic, Israel J. Math. 24 (1975) 154-158.
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