# How to extend Roberts' Law for eccentrically driven, inverted slider-cranks 

Citation for published version (APA):

Dijksman, E. A., \& Smals, A. T. J. M. (1996). How to extend Roberts' Law for eccentrically driven, inverted slidercranks. In J. Lenarcic, \& V. Parenti-Castelli (Eds.), Recent advances in robot kinematics (pp. 325-336). Kluwer Academic Publishers.

## Document status and date:

Published: 01/01/1996

## Document Version:

Publisher's PDF, also known as Version of Record (includes final page, issue and volume numbers)

## Please check the document version of this publication:

- A submitted manuscript is the version of the article upon submission and before peer-review. There can be important differences between the submitted version and the official published version of record. People interested in the research are advised to contact the author for the final version of the publication, or visit the DOI to the publisher's website.
- The final author version and the galley proof are versions of the publication after peer review.
- The final published version features the final layout of the paper including the volume, issue and page numbers.
Link to publication


## General rights

Copyright and moral rights for the publications made accessible in the public portal are retained by the authors and/or other copyright owners and it is a condition of accessing publications that users recognise and abide by the legal requirements associated with these rights.

- Users may download and print one copy of any publication from the public portal for the purpose of private study or research.
- You may not further distribute the material or use it for any profit-making activity or commercial gain
- You may freely distribute the URL identifying the publication in the public portal.

If the publication is distributed under the terms of Article 25 fa of the Dutch Copyright Act, indicated by the "Taverne" license above, please follow below link for the End User Agreement:
www.tue.nl/taverne

## Take down policy

If you believe that this document breaches copyright please contact us at:
openaccess@tue.nl
providing details and we will investigate your claim.

# HOW TO EXTEND ROBERTS' LAW FOR ECCENTRICALLY DRIVEN, INVERTED SLIDER-CRANKS 

E.A.DIJKSMAN (author) \& A.T.J.M.SMALS (software)<br>Faculty of Mechanical Engineering<br>Eindhoven University of Technology<br>The Netherlands


#### Abstract

The extension of Roberts' Law concerns inverted slider-cranks with an eccentricity and a general tracing-point attached to the moving plane of the slider. As both curve-cognates degenerate in this case, the infinite turning-joint has first been replaced by a finite far-away joint, but such that the coupler-plane also meets two accuracy-positions of the erstwhile slider. One of the curve-cognates may then be adapted or modified until up to six accuracypositions are met. They seem to be sufficient to attain an extraordinary good approximation of the entire branch of the original curve. Besides, even the transmission-angle attains a value about twice as large as the one obtained with the azoxiliary four-bar representing the first approximation. Cases, where the inverted slider-crank turns into a crank-and-rocker, into a double crank or even into a stretchable four-bar, are shown. The latter corresponds with one of Grashof's border cases for the inverted slider-crank. Roberts' Law then ensures that each inverted slider-crank, having a fully revolving crank, possesses three four-bar branchcognates this way.


## 1. Introduction

Though Roberts' Law (ref. [1]) describes the general existence of three four-bar curve cognates, all able to produce the same four-bar coupler curve, its application seems to collapse when applied on a degenerated four-bar such as the inverted slider-crank.
However, a previous paper [4] showed a practical circumvention in the special case where the inverted slider-cranks produced symmetrical curves. Then, a very good approximation was obtained based on stretch-rotation and symmetrization. However, for the more general type of the (eccentric) inverted slider-crank, symmetrization isn't applicable and has to be replaced by another procedure, apparently leading to a similar concurrence between the curves as in the symmetrical case.
In the eccentric case, we are going to start from in this paper, also non-Grashof types occur. Then, there is no crank making a complete revolution, whereas the curve produced will be singular branched as then only one branch occurs as the gradual merging result of two different branches originally produced by Grashof types.
In the symmetrical case the replacement four-bar reproduced only one branch generated by the then centrically driven inverted slider-crank. Naturally, something alike is to be expected for the eccentric case. Thus, when an eccentrically driven inverted slider-crank is of the Grasfof-type, two branches appear, each possibly replaceable by one couplerbranch of a Grashof four-bar. For the Non-Grashof type only part of the curve may be reproduced. This part then corresponds with either a forward - or, otherwise, a backward stroke of the "input-rocker" between its extreme positions (Fig. 9).


Fig. 1: Initial coupler braach produced by point $K$ attached to the slider of an inverted slidercrank (4 accuracy positions chosen to design a 4-bar branch-cognate with)


Fig. 2: Design of 4-bar branch cognate (nearly) producing the same coupler branch as the one produced by the eccentric inverted slider-crank (4 position synthesis with position reduction)

## 2. Investigated Types

It is possible to distinguish between four types of eccentric inverted slider cranks. Two of them are of the Grashof type, whereas two others are Non-Grashof. For the Grashof-type
either $a \cdot e_{l}<d$ (Fig. 1), or $d \cdot e_{l}<a$ (Figrs. 3, 10).
For the Non-Grashof type, the crank-circle about A ${ }_{0}$ has to intersect the eccentricity-circle (about $B_{0}$ ) at real points R and $\bar{R}$, leading to the very existence of $\Delta \mathrm{A}_{0} \mathrm{RB}_{0}$ meeting the three Non-Grashof conditions:

```
\(a<d+a_{l}\)
\(d<e_{1} \cdot a\)
\(a_{1}<a+d\)
```

Of these, the last one has to be met anyway, otherwise no real mechanism is to be drawn with real tangents to the eccentricity-circle. Thus, as $\quad \boldsymbol{e}_{\boldsymbol{f}}<\boldsymbol{a} \cdot \boldsymbol{d}$ represents an almost trivial condition, true also for Grashof linkages, only two cases are left for the NonGrashof type. They are:
case I $\quad d-\varepsilon_{I}<a<d^{\prime} \quad$ (Fig. 9)
case II $\quad d<a<d \cdot e_{1}$
Note that for the centric case, for which $e_{1}=0$, the Non-Grashof types disappear. Further, border-cases appear when either $a \cdot e_{l}=d$ (Fig. 6) or $d \cdot e_{l}=a$

Of course, in order to be complete, all occurring cases are of interest when looking for the possible existence of four-bar branch cognates when Grashof linkages are at hand, or "half curve cognates" when Non-Grashof inverted slider cranks are presented (Fig. 9).

## 3. Design of the 4-bar branch-cognate

The actual design of the four-bar cognate replacing the inverted slider-crank occurs in two stages: (Figrs. 2, 4, 7, (9) and 10)
In the $1^{\text {st }}$ stage a first approximation is adopted leading to an auxiliary four-bar, based on two specifically chosen accuracy positions of the coupler-point and on a simultaneous replacement of the slider.
In the $2^{\text {nd }}$ stage, part of Roberts' Configuration is used to find another crank-circle, simultaneously giving the designer the opportunity to improve his first approximation of the branch through two other pairs of specifically chosen accuracy positions.
As a result, only one branch of the 4-bar coupler curve approximates the chosen branch, initially produced by the inverted slider-crank. For the other branch one finds another branch-cognate somewhat different from the first.
Accuracy positions, as will be proved, occur in pairs. That is to say, for each accuracyposition another exists leading to the same result. So, basically, the design uses only three accuracy positions, in reality being six.
The first pair is to be allocated for the auxiliary four-bar, whereas the remaining pairs are to be used for the final mechanism found in the $2{ }^{\text {nd }}$ stage.


Fig. 3: Original coupler-branch produced by a coupler point $K$ of an inverted slider-crank with eccentricity $e_{1}$ (4 accuracy positions are used to obtain the 4-bar branch cognate)


Fig. 4: The ecc. inverted slider-crank as well as the double crank (being the branch cognate) producing the same coupler branch

### 3.1 THE AUXILLARY FOUR-BAR

In order to simplify matters and to attain a unified method for all cases, the first pair of accuracy-positions are taken to correspond with the perpendicular positions of the inputcrank with respect to the frame. (Figrs. 1, 3, 6, 10). Then, $\mathrm{P}_{12}$, the virtual rotation center of the positions, $\mathrm{A}_{1} \mathrm{~K}_{1}$ and $\mathrm{A}_{2} \mathrm{~K}_{2}$, coincides at the center $\mathrm{B}_{0}$ of the eccentricity-circle. (The mid-normal of $\mathrm{B}_{1}{ }^{*} \mathrm{~B}_{2}{ }^{\text {" }}$ always joins $\mathrm{B}_{0}$.)
One further replaces the slider $A_{1} R_{1}$, touching the eccentricity-circle at $R_{1}$, by a basically required large "circle" with center $\mathrm{B}_{1}$. Then, the last point, becoming a finite turningjoint, has to comply with the first two accuracy-positions, whence $B_{1}$ joins the normal of $A_{1} K_{1}$, simultaneously meeting point $P_{12}=B_{0}$.

In order to have a large length for $c=\overline{\boldsymbol{B}_{i} \boldsymbol{B}_{o}}=h+e_{1}$
one chooses the angle $\alpha=\varangle B_{1} A_{1} R_{1}=80^{\circ}$. (In case $\alpha=90^{\circ}$, one reattains the inverted slider-crank.)
The dimensions of the auxiliary mechanism ( $\mathrm{A}_{0}-\mathrm{A}_{1} \mathrm{~K}_{1} \mathrm{~B}_{1}-\mathrm{B}_{0}$ ) having the same crank, the same coupler-point, and also the same center - $\mathrm{B}_{0}-$ as the inverted slider-crank we started from, is then to be established from the initial dimensions: $\overline{A_{0} B_{0}}=d ; \overline{A_{F} A_{\theta}}=a$;
$\overline{B_{0} R_{1}}=e_{1} ; \overline{A_{1} K_{1}}=p$ and $e_{1}$, the latter being the shortest distance from $K_{1}$ to $\mathrm{A}_{1} \mathrm{R}_{1}$. We so find that:

$$
\begin{align*}
& \overline{A_{1} R_{1}}=\sqrt{d^{2} \cdot a^{2}-a_{l}^{2}}  \tag{2}\\
& \overline{B_{1} R_{1}}=h \cdot \overline{A_{1} R_{1}} \tan a  \tag{3}\\
& \overline{A_{1} B_{1}}=b-\overline{A_{1} R_{1}} \cos ^{-1} a  \tag{4}\\
& \overline{B_{1} K_{1}}=q=\sqrt{\left(h \pm e_{l}^{2}\right)^{2} \cdot\left(\sqrt{p^{2}-e_{2}^{2}}-\overline{\left.A_{1} R_{1}\right)^{2}}\right.}  \tag{5}\\
& \varangle K_{P} B_{2} A_{l}=\beta=\arcsin \left(\frac{p}{q} \sin \left(\alpha-\operatorname{arcstn} \frac{e_{2}}{p}\right)\right] \tag{6}
\end{align*}
$$

Clearly, the design of the auxiliary mechanism is the same for all Grashof-types having a fully revolving crank. Thus,

$$
\begin{array}{rlllll}
\text { for } & a \leq d-e_{1} & \text { with } & d>e_{1} & \text { (Figrs. 1, 6) } \\
\text { for } & a \geq d \cdot e_{1} & \text { with } & d>e_{1} & \text { (Fig. 3) } \\
\text { and for } & a \geq d \cdot e_{1} & \text { with } & d<\varepsilon_{1} & \text { (Fig. 10) }
\end{array}
$$

Of course, as $\alpha<90^{\circ}$, the curve produced by the auxiliary mechanism will deviate a bit from the initial curve as produced by the eccentric inverted slider-crank. However, the deviation is to be corrected by application of the $2{ }^{\text {nd }}$ stage leading to the final replacement mechanism.


Fig. 5: The double crank as a four-bar branch cognate of an eccentrically driven inverted slidercrank (the 2 mechanisms produce about the same branch with nearly the same input velocity)


Fig. 6: Eccentrically driven, inverted slider crank (4 accuracy positions to be used for the design of a 4-bar branch cognate)

### 3.2 FINAL FOUR-BAR BRANCH COGNATE

Naturally, further enlargement of the angle $\propto$ leads to lower values for the transmission angle $\mu_{1}=\varangle \mathrm{B}_{0} \mathrm{~B}_{1} \mathrm{~A}_{1}$. (For instance, if $\alpha=90^{\circ}$, then $\mu_{1}=0^{\circ}$ )
As this is undesired with regard to force-transmission, other ways have to be found. Whence, we abide by $\alpha=80^{\circ}$, but look at Roberts' Law and ditto Configuration instead. Though Roberts' Law may then be applied caused by the now finite location of point B ${ }_{1}$, the transformation does not change the curve itself.
However, if we consider a particular curve-cognate of Roberts', namely the one having a crank rotating with the same speed as the angular velocity of the initial crank, it becomes feasible to change three of its cognate-dimensions, in order to meet three different accuracy positions. To carry this out in detail, the following procedure is adopted:

1. First, change the auxiliary four-bar into its Roberts'curve-cognate having a crank $A_{0}{ }^{\prime \prime} A^{\prime \prime}$ rotating with the same speed as the initial crank $A_{0} A$,
 three (or six) accuracy-positions being chosen.

In more detail, this would lead to the stretch-rotation of $\Delta A_{1} A_{0} B_{0}$ about $B_{0}$ with the complex multiplication-factor $\frac{q}{b} e^{\ell \beta}$, yielding

$$
\begin{equation*}
\Delta A_{1}{ }^{\prime \prime} A_{0}{ }^{\prime \prime} B_{0}=\frac{q}{b} e^{\ell \beta}\left(\Delta A_{1} A_{0} B_{0}\right) \tag{7}
\end{equation*}
$$

For the $2^{\text {nd }}$ accuracy-position, a similar equation holds:

$$
\begin{equation*}
\Delta A_{2}^{*} A_{0}^{*} B_{0}=\frac{g}{b} e^{t \beta}\left(\Delta A_{2} A_{0} B_{0}\right) \tag{8}
\end{equation*}
$$

Hence, we so obtain the simple formulas:

$$
\begin{align*}
& \overline{A_{0}^{\prime \prime} B_{0}}=d^{\prime}=\frac{q}{b} d  \tag{9}\\
& \overline{A_{1}^{\prime \prime} A_{0}^{\prime \prime}}=a^{\prime}=\frac{q}{b} a \tag{10}
\end{align*}
$$

Further, $\varangle B_{0} A_{0}^{\prime \prime} A_{1}^{\prime \prime}=90^{\circ}=\varangle B_{0} A_{0}^{\prime \prime} A_{2}^{\prime \prime}$
Whereas, similarly according to Roberts' Configuration, also

$$
\begin{equation*}
\overline{A_{2}^{\prime \prime} K_{2}}=\overline{A_{1}^{\prime \prime} K_{1}}=p^{\prime}-\frac{c}{b} p \tag{12}
\end{equation*}
$$

Naturally, the perpendicular bisectors of $\overline{A_{1}^{\prime \prime} A_{2}^{\prime \prime}}$ and of $\overline{K_{1} K_{2}}$ intersect at $\mathrm{B}_{0}$. Thus,

$$
\begin{equation*}
\Delta A_{2}^{\prime \prime} B_{0} K_{2}=\Delta A_{1}^{\prime \prime} B_{02} K_{1} \tag{13}
\end{equation*}
$$

yielding $\quad B_{01}=B_{0}=B_{02}$.
Then, with two other chosen accuracy positions from the initially given mechanism, say $\mathrm{A}_{3} \mathrm{~K}_{3}$ and $\mathrm{A}_{4} \mathrm{~K}_{4}$, it becomes quite possible to determine accurate corresponding positions,


Fig. 7. Design of four-bar Branch-Cognate of inverted slider-crank (Grashof's Border case $a+e_{1}=d$ leading to $a^{\prime}+b^{\prime}>d^{\prime}+c^{\prime}$ )


Fig. 8: Four-bar Branch Cognate of inverted slider-crank
$\boldsymbol{A}_{3}^{\prime \prime}$ and $\boldsymbol{A}_{4}^{\prime \prime}$, at the cognated crank-circle about $\boldsymbol{A}_{\theta}^{\prime \prime}$. This may be realized by intersection of circles having radius $p^{\prime}$ respectively about $K_{3}$ and $K_{4}$, with the cognated crank-circle about $A_{o}^{\prime \prime}$

When, for instance, $\mathrm{A}_{3}$ as well as $\mathrm{A}_{4}$ join the initial fixed link $\mathrm{A}_{0} \mathrm{~B}_{0}$, we find that both, $A_{3}^{\prime \prime} \quad$ as well as $A_{4}^{\prime \prime}$, slightly deviates from the line containing the cognated fixed link, that is to say from $A_{0}^{H} B_{0}$. Anyway, having determined their exact locations at the cognated crank-circle, we may establish the locations of the points, $B_{03}$ and $B_{a r}$. The way to find these points, will be based on the so-called attachment-method, hence, on the equations:

$$
\begin{align*}
& \Delta A_{3}^{\prime \prime} B_{0} K_{3}=\Delta A_{1}^{\prime \prime} B_{03} K_{1}  \tag{15}\\
& \Delta A_{4}^{\prime \prime} B_{0} K_{4}=\Delta A_{1}^{\prime \prime} B_{00} K_{1} \tag{16}
\end{align*}
$$

The circle joining the three points $\mathrm{B}_{01}=\mathrm{B}_{02}, \mathrm{~B}_{03}$ and $\mathrm{B}_{04}$ then has the required turningjoint $B_{I}^{\prime \prime}$ for its center.

As a result, one obtains the new dimensions:

$$
\begin{equation*}
\overline{A_{1}^{\prime \prime} B_{1}^{\prime \prime}}-b^{\prime} \quad, \overline{B_{0} B_{1}^{\mu}}=c^{\prime} \text { and } \overline{K_{1} B_{1}^{\prime \prime}}=q^{\prime} \tag{17}
\end{equation*}
$$

In the particular case for which the initial branch is symmetrical, which occurs when the two eccentricities are zero, $B_{I}^{\prime \prime}$ coincides at the center of a circle joining $K_{1}, B_{0}$ and $A_{\boldsymbol{I}}^{a}$. In that case $b^{\prime}=c^{\prime}=\boldsymbol{q}^{\prime}$, meaning that the four-bar is of the $\lambda$-type (ref.[4]).

The accuracy-positions chosen, corresponded with the initial crank-positions dividing $360^{\circ}$ into equal parts of $90^{\circ}$.
A measure for the accuracy of the cognate curve may be found by comparing the Grashof-distances of the two mechanisms. Then, the best approximation will be the one having its Grashof-distance-ratio much nearer the value 1 than with other approximations. For the mechanisms of Fig. 2, for instance, the auxiliary four-bar yields the Grashof-distance-ratio $\frac{a \cdot c-d-b}{a-d \cdot e_{1}}=1.57$, whereas the final branch-cognate results into the ratio $\frac{d\left(a^{\prime} \cdot b^{\prime}-c^{\prime}-d^{\prime}\right)}{d^{\prime}\left(a-d \cdot e_{I}\right)}=0.924$. Thus, the latter represents the better mechanism. (Note that the Grashof-distance-ratio equals 1 when two Roberts' 4-bar curve cognates are compared.)
For the auxiliary mechanism of Fig. 4, we obtain the Grashof-distance-ratio


Fig. 9: Non-Grashof Curve-Cognate


Fig. 10: Iverted slider crank and its 4-bar branch-cognate producing the same branch (4 position-synthesis using a point position reduction method: $B_{01}=B_{02}$ )
$\frac{d \cdot c-a-b}{a-d-e_{1}}=2.34$, whereas the branch-cognate gives rise to the ratio

$$
\frac{d\left(d^{\prime}+b^{\prime}-a^{\prime}-c^{\prime}\right)}{d^{\prime}\left(a-d-e_{1}\right)}=0.81
$$

Even for Non-Grashof linkages, namely for those without revolving bars, the method remains applicable. Fig. 9, for instance, yields for the auxiliary mechanism the Grashof-distance-ratio $\frac{a+b-c-d}{a-d+a_{I}}=1.65$, whereas the final curve-cognate gives the ratio

$$
\frac{d\left(a^{\prime}+c^{\prime}-d^{\prime}-b^{\prime}\right)}{d^{\prime}\left(a-d+a_{1}\right)}=1.35 .
$$

It is quite possible that an other distribution of our accuracy-positions, gives a better Grashof-ratio. Remarkable better results though, are not to be expected. For instance, if we choose point $A_{4}$ of Fig. 9 at the other intersection of the crank - and the eccentricitycircle, we obtain the ratio $\frac{d\left(a^{\prime}+c^{\prime}-d^{\prime}-b^{\prime}\right)}{d^{\prime}\left(a-d+a_{l}\right)}=0.73$. Now, its distance to 1 is only slightly less than with the 1.35 -value belonging to the mechanism demonstrated in Fig. 9. In Grashof's border case, such as the particular one demonstrated in the Figrs. 6 and 7, we rather observe the difference between these Grashof-distances.
Then, for the auxiliary mechanism we precisely observe the value

$$
\begin{aligned}
(a \cdot c-b-d)-\left(a \cdot e_{1}-d\right)= & \sqrt{d^{2} \cdot a^{2}-e_{1}^{2}} \cdot(\sin \alpha-1) \cos ^{-1} \alpha- \\
& =-40 \sqrt{3} \tan 1 / 2[(\pi / 2)-\alpha]=-40 \sqrt{3} \tan 5^{\circ}=-6.0614
\end{aligned}
$$

As in this case each Grashof-distance should tend to zero, a considerable improvement is obtained with the final branch-cognate as demonstrated in Fig. 8. Then, namely

$$
\frac{d}{d^{\prime}} \cdot\left(a^{\prime}+b^{\prime}-c^{\prime}-d^{\prime}\right)-\left(a+e_{1}-d\right)=+0.037
$$

A random circle about $\mathrm{B}_{0}$ intersecting the cognated crank-circle at possible accuracypositions $A_{m}^{\prime \prime}$ and $A_{n}^{\prime \prime}$, yields a virtual rotation-centre $\mathrm{P}_{\mathrm{mn}}=\mathrm{B}_{0}$, giving $\mathrm{B}_{0 \mathrm{~m}}=\mathrm{B}_{0 \mathrm{n}}$. Indeed, accuracy-positions only occur in pairs. Thus, each accuracy-position has to be counted twice, unless of course the cognated crank-positions already are at a circle about $\mathrm{B}_{0}$, as will be the case for the accuracy-positions, $\boldsymbol{A}_{1}^{\prime \prime}$ and $\boldsymbol{A}_{2}^{\prime \prime}$.
Generally, the transmission angle $\mu_{1}=\varangle A_{1}^{\prime \prime} B_{1}^{\prime \prime} B_{0} \quad$ appears to be about twice as large as the one obtained with the auxiliary four-bar or with its Roberts' curve-cognate. Complete application of Roberts' Law would have lead to a point $\left.\quad\left(B_{1}^{u}\right)^{\prime}\right)$ as the fourth vertex of a linkage parallelogram $B_{B_{1}} B_{1} K_{l}\left(B_{i}^{\prime \prime}\right)_{\text {aques }}$. However, such a point $\quad\left(\mathcal{B}_{1}^{\prime \prime}\right)_{\text {cases }}$ would have been much farther away, giving about half the transmission-angle. All examples demonstrated show this same phenomenon. We conclude, that the $2^{\text {nd }}$ stage leads to a better approximation of the generated branch, simultaneously giving about twice the transmission angle in comparision to the one obtained in the $1{ }^{\text {st }}$ stage.

Theoretically, the circular curve described by point $\mathrm{B}_{0}$ of the Stephenson-2 six-bar

$$
\left(A_{0}-A-B^{-}-B_{0}-A_{0}^{\prime \prime}-A_{1}^{\prime \prime}-K_{1}\right) \quad \text { with respect to the link } A_{1}^{\prime \prime} K_{1} \text { will be a }
$$

Stephenson-2 six-bar curve of order 16, (ref.[5]). Clearly, this curve approximates a circle very neatly, although in reality not more than 16 intersections exist. (We count 2 times 16 , minus the ( 8 times 2 ) intersections of the circle at the two circular or isotropic points of the curve.)

## 4. Conclusions

Branches of curves produced by inverted slider-cranks having an eccentricity and/or an eccentrically located coupler point, are to be reproduced by even three different four-bar branch-cogn ates. In the two cases the curves are singular branched, the three four-bar curve-cognates reproduce only part of the curve, namely the part corresponding to either a forward - or, otherwise a backward stroke of the input-rocker. Better transmission angles may be attained at the cost of the accuracy of the reproduction. In most cases though a high accuracy is obtained with acceptable transmission angles. (The accuracy of the replacement-method has been measured with a newly introduced Grashof-distanceratio.)

## References

1. Roberts, S. (1875) Three-bar motion in plane space, Proc.London Math.Soc.7, pp.14-23
2. Stoimenov, M., Panteli \&, T. (1979) Possibility of the development of a curved translation and for a member of an extended kinematic chain of a curved sliding mechanism, Proc. of the $5^{\text {th }}$ World Congress on the Theory of Machines and Mechanisms, pp 1436-1439, published by the ASME.
3. Dijksman, E.A. (1976) Motion Geometry of Mechanisms, (Chapter 5: Cognate Linkages), Cambridge University Press.
(1981) Cinemática De Mecanismos. (Cap.5: Mecanismos Cognados) Editorial LIMUSA, S.A., Mexico
4. Dijksman, E.A. (1996) How to exchange centric inverted slider cranks with $\lambda$-formed fourbar linkages, Mechanism and Machine Theory (in press)
5. Primrose, E.J.F., Freudenstein, F., Roth, B. (1967) Six-Bar Motion Il The Stephenson-1 and Stephenson-2 Mechanisms, Archive for Rational Mechanics and Analysis, Volume 24,Nr.1, pp.55-72.
