

Thickness dependence of unidirectional spin-Hall magnetoresistance in metallic bilayers

Citation for published version (APA):

Yin, Y., Han, D. S., de Jong, M. C. H., Lavrijsen, R., Duine, R. A., Swagten, H. J. M., & Koopmans, B. (2017). Thickness dependence of unidirectional spin-Hall magnetoresistance in metallic bilayers. *Applied Physics Letters*, 111(23), Article 232405. <https://doi.org/10.1063/1.5003725>

DOI:

[10.1063/1.5003725](https://doi.org/10.1063/1.5003725)

Document status and date:

Published: 04/12/2017

Document Version:

Publisher's PDF, also known as Version of Record (includes final page, issue and volume numbers)

Please check the document version of this publication:

- A submitted manuscript is the version of the article upon submission and before peer-review. There can be important differences between the submitted version and the official published version of record. People interested in the research are advised to contact the author for the final version of the publication, or visit the DOI to the publisher's website.
- The final author version and the galley proof are versions of the publication after peer review.
- The final published version features the final layout of the paper including the volume, issue and page numbers.

[Link to publication](#)

General rights

Copyright and moral rights for the publications made accessible in the public portal are retained by the authors and/or other copyright owners and it is a condition of accessing publications that users recognise and abide by the legal requirements associated with these rights.

- Users may download and print one copy of any publication from the public portal for the purpose of private study or research.
- You may not further distribute the material or use it for any profit-making activity or commercial gain
- You may freely distribute the URL identifying the publication in the public portal.

If the publication is distributed under the terms of Article 25fa of the Dutch Copyright Act, indicated by the "Taverne" license above, please follow below link for the End User Agreement:

www.tue.nl/taverne

Take down policy

If you believe that this document breaches copyright please contact us at:

openaccess@tue.nl

providing details and we will investigate your claim.

Thickness dependence of unidirectional spin-Hall magnetoresistance in metallic bilayers

Yuxiang Yin, Dong-Soo Han, Mark C. H. de Jong, Reinoud Lavrijsen, Rembert A. Duine, Henk J. M. Swagten, and Bert Koopmans

Citation: *Appl. Phys. Lett.* **111**, 232405 (2017);

View online: <https://doi.org/10.1063/1.5003725>

View Table of Contents: <http://aip.scitation.org/toc/apl/111/23>

Published by the [American Institute of Physics](#)

Articles you may be interested in

[Antidamping spin-orbit torques in epitaxial-Py\(100\)/ \$\beta\$ -Ta](#)

Applied Physics Letters **111**, 232407 (2017); 10.1063/1.5007202

[Anomalous Nernst effect in Ir₂₂Mn₇₈/Co₂₀Fe₆₀B₂₀/MgO layers with perpendicular magnetic anisotropy](#)

Applied Physics Letters **111**, 222401 (2017); 10.1063/1.4996399

[Tuning the metamagnetism in a metallic helical antiferromagnet](#)

Applied Physics Letters **111**, 232404 (2017); 10.1063/1.5004667

[Layer-sensitive magneto-optical Kerr effect study of magnetization reversal in Fe/MnAs/GaAs\(001\)](#)

Applied Physics Letters **111**, 232403 (2017); 10.1063/1.5004248

[Tailoring magnetic skyrmions by geometric confinement of magnetic structures](#)

Applied Physics Letters **111**, 242405 (2017); 10.1063/1.5005904

[Spin-phonon coupling in antiferromagnetic nickel oxide](#)

Applied Physics Letters **111**, 252402 (2017); 10.1063/1.5009598

Scilight

Sharp, quick summaries **illuminating**
the latest physics research

Sign up for **FREE!**



Thickness dependence of unidirectional spin-Hall magnetoresistance in metallic bilayers

Yuxiang Yin,^{1,a)} Dong-Soo Han,¹ Mark C. H. de Jong,¹ Reinoud Lavrijsen,¹ Rembert A. Duine,^{1,2} Henk J. M. Swagten,¹ and Bert Koopmans¹

¹Department of Applied Physics, Eindhoven University of Technology, PO Box 513, 5600 MB Eindhoven, The Netherlands

²Institute for Theoretical Physics and Center for Extreme Matter and Emergent Phenomena, Utrecht University, Leuvenlaan 4, 3584 CE Utrecht, The Netherlands

(Received 7 September 2017; accepted 17 November 2017; published online 5 December 2017)

A nonlinear magnetoresistance—called unidirectional spin-Hall magnetoresistance—is recently experimentally discovered in metallic bilayers consisting of a heavy metal and a ferromagnetic metal. To study the fundamental mechanism of unidirectional spin-Hall magnetoresistance (USMR), both ferromagnetic and heavy metallic layer thickness dependence of the USMR are presented in a Pt/Co/AlOx trilayer at room temperature. To avoid ambiguities, second harmonic Hall measurements are used for separating spin-Hall and thermal contributions to the non-linear magnetoresistance. The experimental results are fitted by using a drift-diffusion theory, with parameters extracted from an analysis of longitudinal resistivity of the Co layer within the framework of the Fuchs-Sondheimer model. A good agreement with the theory is found, demonstrating that the USMR is governed by both the spin-Hall effect in the heavy metallic layer and the metallic diffusion process in the ferromagnetic layer. *Published by AIP Publishing.* <https://doi.org/10.1063/1.5003725>

In the field of Spintronics, a new way of spin control based on the spin-Hall effect (SHE) recently has attracted a great deal of attention. It originates from the spin-orbit (SO) interaction which converts a charge current into a net flow of spin angular momentum, exerting a SO torque on the magnetization. This leads to an energy-efficient way of writing information to magnetic memories by switching a magnetic entity via sending a charge current through a nearby nonmagnetic metal.^{1,2} Apart from writing, a possible way of reading the memory could be achieved by measuring a so-called spin-Hall magnetoresistance (SMR),^{3,4} i.e., a change in electrical resistance due to the spin Hall effect when the spin-accumulation is perpendicular or parallel to the magnetization. Although SMR provides a promising way towards reading memory devices using a two-terminal architecture, the fact that it can only distinguish between the perpendicular and parallel states limits its application.

Very recently, a unidirectional contribution to magnetoresistance—called unidirectional spin-Hall magnetoresistance (USMR)—has been reported in a ferromagnetic/heavy metallic (FM/HM) bilayer structure.^{5–7} Being different from the ordinary SMR, the resistance changes by reversing the magnetization or the current direction, which could be potentially utilized for reading operation. Based on a drift-diffusion-relaxation theory,⁸ this nonlinear behavior is attributed to the dependence of electron mobility on spin-polarization, which is tuned by the spin-Hall effect induced spin accumulation. This spin accumulation is limited to a thin region at the FM/HM interface due to a finite spin diffusion length in both layers, leading to a non-trivial FM and HM thickness dependence of the USMR. So far, this particular dependence on thickness is not evidenced by any experiments. Thus, a systematic investigation of how USMR depends on the layer

thickness is urgently needed not only for a better understanding of the origin of USMR but also for the enhancement of USMR in practical applications.

In this paper, we present the FM and HM layer thickness dependence of USMR in Pt/Co/AlOx trilayers at room temperature. The experimental results are fitted by using the aforementioned drift-diffusion-relaxation theory, with parameters extracted from an analysis of the longitudinal resistivity of the Co layer within the framework of the Fuchs-Sondheimer model.⁹ Furthermore, second harmonic Hall measurements enable us to disentangle spin-Hall and thermal gradient contributions to the non-linear magnetoresistance, allowing for a more precise fitting. Good agreement with the theory is found, demonstrating that the USMR depends on both the spin-Hall effect in the HM layer and the electron spin diffusion and relaxation in the FM layer.

For the measurement of magnetoresistance, the multilayer structures are Pt(1–8 nm)/Co(4 nm)/AlOx(1.15 nm) (6 samples) and Pt(4 nm)/Co(1–50 nm)/AlOx(1.15 nm) (11 samples), where we vary either the Pt or Co thickness as indicated by the thickness range in the parentheses. These multilayers are then patterned in the form of a Hall bar, shown in Fig. 1(a), by using electron-beam lithography and lift-off. The length of the Hall bar is 100 μm , the lateral width is 5 μm , and the spacing between two Hall bars is 20 μm . The samples are deposited on Si/SiO₂ substrates by DC magnetron sputtering. Pt was deposited at a rate of 0.08 nm/s, and Co was sputtered at a rate of 0.05 nm/s. After deposition, a 1.15 nm thick Al capping layer was finally deposited and further oxidized (by using plasma oxidation during 90 s at 1×10^{-1} mbar) on top of the Pt/Co stack, to prevent oxidation of the Co layer in air.

The magnetoresistance measurements presented in this work were performed at room temperature by using an AC current source with a current density of 1×10^7 A/cm²

^{a)}E-mail: y.yin@tue.nl

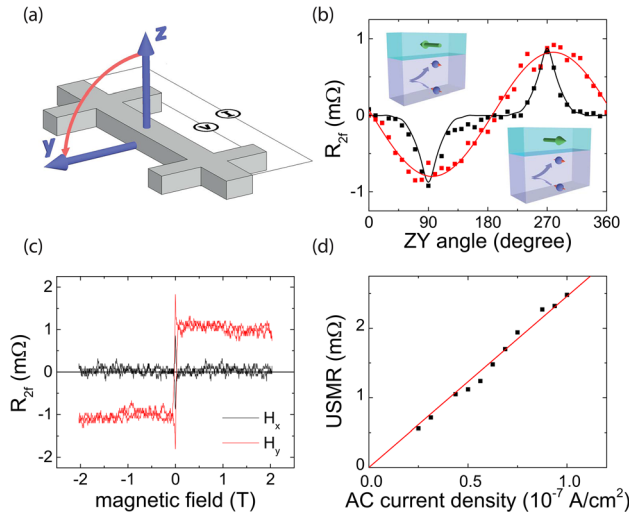


FIG. 1. (a) A schematic illustration of the Hall bar used in the experiment. The measurement scheme of longitudinal resistances is shown. (b) Angle dependence measurements of R_{2f} in the sample Pt(4 nm)/Co(1 nm)/AlOx(1.15 nm) (black dots) and Pt(4 nm)/Co(2 nm)/AlOx(1.15 nm) (red dots). Solid lines represent the fitting results. (c) y (red line) and x (black line) field dependence measurements of R_{2f} in the sample Pt(4 nm)/Co(2 nm)/AlOx(1.15 nm). (d) USMR as a function of current density.

modulated at $f = 801$ Hz. The second harmonics component of the longitudinal resistance R_{2f} is recorded during sample rotation or sweeps of an external magnetic field.

For the measurement of USMR, first an angle dependence measurement of R_{2f} is performed. An external magnetic field $B_{\text{ext}} = 2$ T, which is high enough to saturate the magnetization in every configuration, rotates in the yz plane while R_{2f} is recorded. The results for two types of samples are plotted in Fig. 1(b). The black dots show the results for a Pt(4 nm)/Co(1 nm)/AlOx(1.15 nm) stack with perpendicular anisotropy (PMA), while the red dots show the results for a Pt(4 nm)/Co(2 nm)/AlOx(1.15 nm) stack with in-plane anisotropy. Starting from the $+z$ direction, R_{2f} gradually decreases and reaches a minimum when the field is along the $+y$ direction. After that, R_{2f} starts to rise and reaches a peak at $-y$. Finally, it returns to its original value after a full rotation. The measurement evidences a resistance contribution that depends on the sign of M_y , and the difference of R_{2f} between the $+y$ and $-y$ directions is defined as the USMR. This is further confirmed by a good agreement between the data and a fitted line with respect to M_y (obtained by an anomalous Hall effect measurement), plotted as a solid line in Fig. 1(b). Note that the transition for the PMA stack is significantly sharper than that for the in-plane sample, since a high field is needed to pull the magnetization in the plane for a stack with PMA.

To further investigate the USMR, we have measured R_{2f} while sweeping the external magnetic field along transverse (y) and longitudinal (x) directions. Figure 1(c) shows that R_{2f} is constant as a function of y field and reverses sign upon sweeping the field from the y direction to the $-y$ direction. Two spikes are observed near zero field due to the formation of magnetic domains during magnetization switching. In contrast to a field in the y direction, no difference is measured between the x and $-x$ directions, indicating that, as expected, the USMR only exists in the transverse direction. Compared with the angle dependent measurement, where

one has to ensure that the field is strong enough to saturate the sample in the z direction, the field dependent measurement serves as a more efficient way of quantifying the USMR. Thus, in the following, USMR will be obtained by sweeping the field. As a further test, Fig. 1(d) shows the current dependence of USMR measured in this way, which is linear with the injected current density and converges to zero for decreasing current, since the spin accumulation at the interface scales with the current density.

To verify the role of the interfacial spin accumulation due to the SHE, we examined the dependence of the USMR on the thickness of the HM and FM layers. Figures 2(a) and 2(b) show the absolute change of second harmonic resistance $\text{USMR} = R_{2f}(+M) - R_{2f}(-M)$ measured at constant current density as a function of the Co and Pt thickness. Both curves exhibit qualitatively similar behavior: An initial sharp increase and a gradual decrease as the layer becomes thicker. Apart from the USMR, thermal effects could also contribute to the R_{2f} . Thus, exclusion of these thermal effects is required before further analysis, which will be discussed in the following part.

Figure 2(c) shows the temperature profile $T(z)$ in a line along the thickness direction of the nanowire by using simulation software suites Comsol multiphysics 5.2. The maximum temperature is found at the top owing to the fact that the heat dissipation is faster through the bottom substrate than through the top ambient air. A temperature gradient in the z direction will generate an electric current in the same direction, which interacts with the FM layer through the anomalous Hall effect and generates an electric field ($\propto \Delta T \times M_{\text{sat}}$) in the x direction.¹⁰ This will cause a resistance change in the x direction and possesses the same symmetry as that of the USMR. In order to separate the thermal contribution from the USMR, we measure the second harmonic

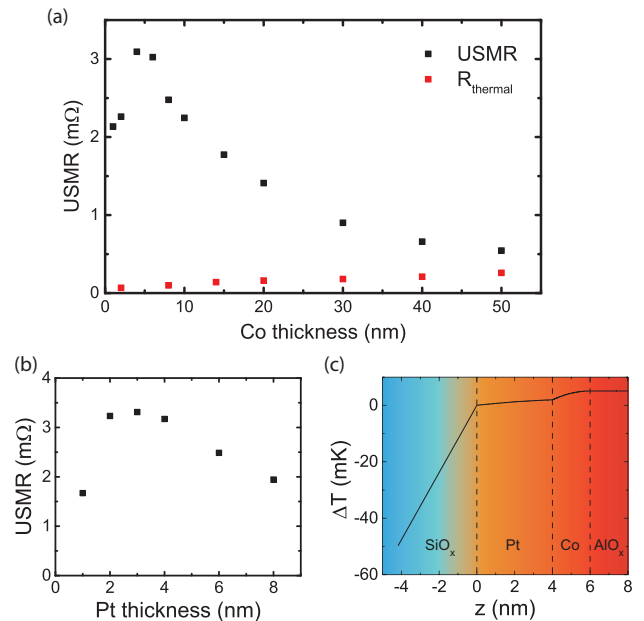


FIG. 2. (a) USMR as a function of the Co thickness and a resistance contribution originated from the thermal effect for Pt(4 nm)/Co(1–50 nm)/AlOx(1.15 nm). (b) USMR as a function of the Pt thickness for Pt(1–8 nm)/Co(4 nm)/AlOx(1.15 nm). (c) Temperature profile in the sample Pt(4 nm)/Co(2 nm)/AlOx(1.15 nm) simulated by using the finite element method.

Hall resistance to quantify the thermal resistance¹¹ (see [Supplementary Material](#) for details). As plotted in the same figure of USMR, see the red dots in Fig. 2(a), the thermal contributions are found to be increased with the thickness. The maximal thermal resistance is observed at a Co thickness of 50 nm, which accounts for about 50% of the USMR and the ratio is smaller for thinner Co. In the following part, the thermal resistance will be subtracted from the USMR to achieve an accurate analysis.

In order to compare the experimental measurement of USMR with the model, we first convert the absolute USMR into normalized USMR, i.e., USMR divided by normal longitudinal resistance. For this purpose, the longitudinal resistance R_{xx} is measured and plotted versus Pt and Co thickness in Figs. 3(a) and 3(b). The plot reveals that the resistance monotonically decreases with thicknesses. The solid line represents the fit which utilizes the Fuchs-Sondheimer approach to extend the conventional t^{-1} resistance model by considering the scattering at the two Co/Pt interfaces.⁹ The fit describes

$$\frac{\rho(E) - \rho(-E)}{\rho(E)} = \frac{6\theta L_F L_H (p_\sigma - p_N) \sigma_F \epsilon \tanh\left(\frac{d_F}{L_F}\right) \tanh\left(\frac{d_H}{2L_H}\right)}{\epsilon_F (d_F \sigma_F + d_H \sigma_H) \left(\frac{L_H (1 - p_\sigma^2) \sigma_F \tanh\left(\frac{d_F}{L_F}\right) \coth\left(\frac{d_H}{L_H}\right)}{L_F \sigma_H} + 1 \right)}, \quad (1)$$

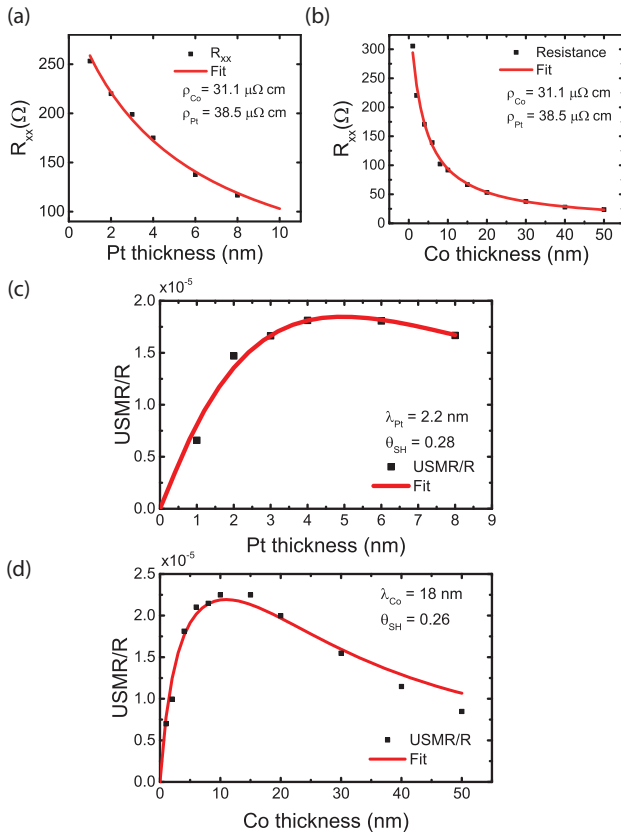


FIG. 3. (a) and (b) Longitudinal resistance as a function of the Co and Pt thickness together with a theoretical fitting curve. (c) and (d) Normalized USMR as a function of the Co and Pt thickness together with a theoretical fitting curve.

the experimental data well and gives the bulk resistivities of $\rho_{Pt} = 37.5 \mu\Omega \text{ cm}$ and $\rho_{Co} = 31.1 \mu\Omega \text{ cm}$, which are comparable to the values in the literature.^{12,13}

Next, we examine the dependence of USMR on the Pt and Co thickness in Pt/Co/AlOx samples. As shown in Fig. 3(c), the normalized USMR is the largest for a Pt thickness of about 5 nm and is reduced for a thicker or thinner Pt layer. Its strong thickness dependence shows that USMR in the structures is mainly influenced by the SHE in Pt. USMR decreases when the Pt layer is thinner than the spin diffusion length due to the reduced spin current caused by back reflection at the interface. On the other hand, for a thicker Pt layer, USMR is also reduced by current shunting. A similar behavior is found for USMR upon varying the Co thickness, as shown in Fig. 3(d), although the maximal USMR is now reached at a Co thickness of 10 nm. Above all, the qualitative behavior matches the prediction of the drift-diffusion-relaxation model. In addition, we also make a quantitative comparison of the experimentally observed values with the model, which describes the USMR as⁸

where d_F (d_H) is the thickness of the FM (HM), L_F (L_H) is the spin diffusion length of the FM (HM), σ_F (σ_H) is the conductivity of the FM (HM), θ is the spin Hall angle of the HM, ϵ is the electric field in FM, ϵ_F is the Fermi energy, p_σ is the conductivity spin asymmetry, and p_N is the difference of density of states at Fermi energy. In our sample, the Fermi energy $\epsilon_F = 5 \text{ eV}$ and spin asymmetry $p_\sigma - p_N = 0.5$.⁸

By fitting the thickness dependence of the normalized USMR to Eq. (1) (red line in Fig. 3), it can be seen that the specific behavior of the data is in line with the drift-diffusion-relaxation model based on the spin-Hall effect, and a spin Hall angle of 0.3 for Pt and spin diffusion lengths of 18 nm and 2.2 nm for Co and Pt, respectively, can be extracted.^{14,15} This implies that the USMR in the Pt/Co system is governed by (1) spin-Hall effect in the Pt layer and (2) electron spin diffusion and relaxation in the Co layer. In a recent work¹⁶ which uses a similar structure (Py/Co), a unidirectional contribution was found in the first harmonic resistance by applying a high current density ($J \sim 10^8 \text{ A/cm}^2$). In this experiment, magnon excitation, instead of an electronic diffusion-relaxation process, is claimed to attribute to this first harmonic USMR. It is also found that the magnon-induced USMR increases with increasing temperature. We do not intend to conduct a temperature dependence measurement here, due to the complex variation of all transport parameters with temperature¹⁷⁻¹⁹ (polarization, spin diffusion length, spin Hall angle, and conductivity) in the drift-diffusion-relaxation model [Eq. (1)], which will make it extremely difficult to draw pertinent conclusions. Moreover, in the magnon experiment,¹⁶ a much

higher current density is used compared to ours, and only the first harmonic in the resistance is addressed, which further complicates a meaningful comparison.

To disentangle electronic and magnonic contribution, a measurement of the temperature dependence of the USMR needs to be performed, which is beyond the scope of this paper. For the electron contribution described before [Eq. (1)], taking into account the temperature variation of all transport parameters (polarization, spin diffusion length, spin Hall angle, and conductivity) predicts that the USMR decreases with increasing temperature,^{17–19} whereas the magnon-induced USMR behaves oppositely.¹⁶

The extracted room-temperature spin Hall angle in Pt appears to be higher than the value ~ 0.1 measured in other work.^{20,21} We still think that this model captures the essential physics of the observed effect, although full quantitative agreement cannot be reached due to various reasons. One reason is the simplifications of the model by assuming spherical Fermi surfaces and constant density of states at the Fermi energy,⁸ which underestimates the magnitude of USMR. Moreover, the model⁸ includes only spin-dependent scattering in the bulk of the ferromagnetic layer. Like in the giant magnetoresistance effect, however, both bulk and interface scattering can contribute to the USMR.⁵ The underestimation would be more if the spin-mixing conductance is incorporated, since the Pt/Co interface is regarded as fully transparent in the model, i.e., the spin-mixing conductance is infinite. Finally, additional charge-spin conversion can take place at either the Pt/Co or Co/AIOx interface, which may lead to a larger spin-Hall effect.^{22,23}

In conclusion, USMR is observed in Pt/Co/AIOx systems and we have shown that the dependence of the USMR on the thickness of both the HM and FM layers agrees qualitatively with the theory based on the electron spin drift-diffusion-relaxation model. We believe that this result provides a better understanding of the physical origin of the USMR and is of importance for its possible applications in spintronic devices.

See [supplementary material](#) for the quantification of thermal contributions to the unidirectional spin-Hall magnetoresistance by measuring the second harmonic Hall resistance.

This work is part of the research program electrical control of spin dynamics and nanomagnetic devices with Project No. FOM-11PR2906, which was (partly) financed by the Netherlands Organisation for Scientific Research (NWO).

- ¹I. M. Miron, K. Garello, G. Gaudin, P.-J. Zermatten, M. V. Costache, S. Auffret, S. Bandiera, B. Rodmacq, A. Schuhl, and P. Gambardella, *Nature* **476**, 189 (2011).
- ²L. Liu, C.-F. Pai, Y. Li, H. Tseng, D. Ralph, and R. Buhrman, *Science* **336**, 555 (2012).
- ³A. Kobs, S. Heße, W. Kreuzpaintner, G. Winkler, D. Lott, P. Weinberger, A. Schreyer, and H. Oepen, *Phys. Rev. Lett.* **106**, 217207 (2011).
- ⁴H. Nakayama, M. Althammer, Y.-T. Chen, K. Uchida, Y. Kajiwara, D. Kikuchi, T. Ohtani, S. Geprägs, M. Opel, S. Takahashi *et al.*, *Phys. Rev. Lett.* **110**, 206601 (2013).
- ⁵C. O. Avci, K. Garello, A. Ghosh, M. Gabureac, S. F. Alvarado, and P. Gambardella, *Nat. Phys.* **11**, 570 (2015).
- ⁶K. Olejnik, V. Novák, J. Wunderlich, and T. Jungwirth, *Phys. Rev. B* **91**, 180402 (2015).
- ⁷K. Yasuda, A. Tsukazaki, R. Yoshimi, K. Takahashi, M. Kawasaki, and Y. Tokura, *Phys. Rev. Lett.* **117**, 127202 (2016).
- ⁸S. S.-L. Zhang and G. Vignale, *Phys. Rev. B* **94**, 140411 (2016).
- ⁹A. Kobs and H. P. Oepen, *Phys. Rev. B* **93**, 014426 (2016).
- ¹⁰W. Nernst, *Ann. Phys.* **267**, 760 (1887).
- ¹¹C. O. Avci, K. Garello, M. Gabureac, A. Ghosh, A. Fuhrer, S. F. Alvarado, and P. Gambardella, *Phys. Rev. B* **90**, 224427 (2014).
- ¹²J. Agustsson, U. Arnalds, A. Ingason, K. B. Gylfason, K. Johnsen, S. Olafsson, and J. T. Gudmundsson, *J. Phys.: Conf. Ser.* **100**, 082006 (2008).
- ¹³P. Carcia, *J. Appl. Phys.* **63**, 5066 (1988).
- ¹⁴W. Zhang, V. Vlaminc, J. E. Pearson, R. Divan, S. D. Bader, and A. Hoffmann, *Appl. Phys. Lett.* **103**, 242414 (2013).
- ¹⁵L. Piraux, S. Dubois, A. Fert, and L. Belliard, *Eur. Phys. J. B-Condens. Matter Complex Syst.* **4**, 413 (1998).
- ¹⁶K. Kim, T. Moriyama, T. Koyama, D. Chiba, S. Lee, S. Lee, K. Lee, H. Lee, and T. Ono, preprint [arXiv:1603.08746](#) (2016).
- ¹⁷M. Isasa, E. Villamor, L. E. Hueso, M. Gradhand, and F. Casanova, *Phys. Rev. B* **91**, 024402 (2015).
- ¹⁸K. Ueda, T. Koyama, R. Hiramatsu, D. Chiba, S. Fukami, H. Tanigawa, T. Suzuki, N. Ohshima, N. Ishiwata, Y. Nakatani *et al.*, *Appl. Phys. Lett.* **100**, 202407 (2012).
- ¹⁹J. Kim, P. Sheng, S. Takahashi, S. Mitani, and M. Hayashi, *Phys. Rev. Lett.* **116**, 097201 (2016).
- ²⁰L. Liu, O. J. Lee, T. J. Gudmundsen, D. C. Ralph, and R. A. Buhrman, *Phys. Rev. Lett.* **109**, 096602 (2012).
- ²¹W. Zhang, W. Han, X. Jiang, S.-H. Yang, and S. S. Parkin, *Nat. Phys.* **11**, 496 (2015).
- ²²V. P. Amin and M. D. Stiles, *Phys. Rev. B* **94**, 104419 (2016).
- ²³L. Wang, R. Wesselink, Y. Liu, Z. Yuan, K. Xia, and P. J. Kelly, *Phys. Rev. Lett.* **116**, 196602 (2016).