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Distributed Bragg grating frequency control in metallic nano lasers

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Abstract - We show that Bragg gratings can be readily incorporated into metallic nano-lasers which exploit waveguides with semiconductor cores, via modulation of the waveguide width. This provides a simple way to implement laser wavelength control.

Sub-wavelength confinement of light in passive plasmonic devices, such as waveguides, has received considerable attention in the last couple of years [1,2,3]. For a long time it was believed that, due to the high optical loss, it was impossible to achieve lasing in a plasmonic cavity. However, Hill and colleagues have shown that it is possible to reproducibly fabricate plasmonic cavities with moderate Q-factors and that it is also possible to sustain a lasing mode inside such a cavity [4,5].

The devices made by Hill are circular and rectangular shaped, with dimensions down to 80 nm, and are fully enclosed in silver. The devices are characterized by collecting light scattered through the substrate on which they are fabricated. To open the way for integration of these lasers in optical systems, side-emission is required. We believe that distributed feedback is suitable to enable side-emission in plasmonic cavities and propose to implement this via incorporation of vertical groove gratings inside the cavity.

To ensure that the propagating mode in the laser cavity is indeed the fundamental plasmon mode of the structure, the width of the waveguide has to be such that it is below cut-off of the 1st order plasmon mode and below the cut-off of the fundamental TE mode of the waveguide. For operation at 1550 nm, the waveguide widths that satisfy these conditions were determined from 2-D FDTD simulations [6] and were found to be ± 140 nm, for a waveguide with a 20 nm thick Si₃N₄ insulation layer and ± 150 nm for a waveguide with a 10 nm thick insulation, figures 1 and 2.

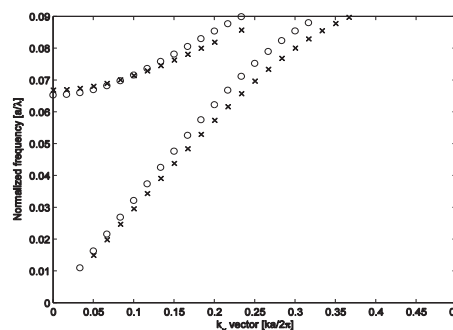


Fig. 1. Geometric dispersion for TM polarization. X & O indicate a 10 nm & 20 nm Si₃N₄ layer respectively, a = 100 nm

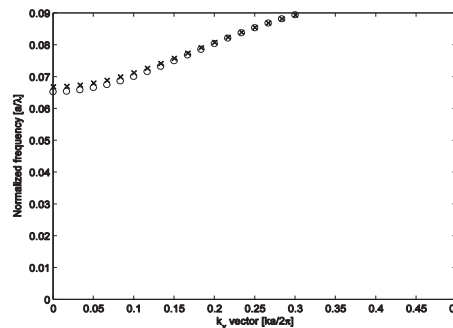


Fig. 2. Geometric dispersion for TE polarization. X & O indicate a 10 nm & 20 nm Si₃N₄ layer respectively, a = 100 nm

Distributed feedback (DFB) lasers are easier to fabricate than laser in which distributed Bragg reflectors (DBR) are incorporated. They are less sensitive to reflections caused by termination of the laser, provided that $\kappa L > 1$ for the grating [7]. Normal DFB lasers support two modes (1st order operation), one on either side of the stop-band. Single mode operation, at the Bragg wavelength, can be forced by either applying a $\lambda/4$ -shift inside the laser cavity or making the laser asymmetric (e.g. by covering one facet with silver and removing it from the other facet), figure 3. We have

determined the optimum period for operation at 1550 nm to be ± 220 nm for a core waveguide width of 80 nm and the total width of the grating is chosen to be 130 nm (=80 nm core + 50 nm grating).

Definition of the devices by means of EBL in combination with lift-off is impossible for devices with these dimensions. The feature size is in the same order of magnitude as the thickness of the metal layer normally used to perform lift-off with. Instead we use a bilayer resist, consisting of HSQ/HPR504, to obtain the required dimensions.

First a 450 nm thick layer of HPR504 photoresist is spun on top of a SiO₂ hardmask layer. The HPR504 is then baked at a very high temperature, in order to harden it. A 80 nm thick HSQ is then spun on top of the hardbaked HPR504.

Hydroxy Silsesquioxane (HSQ) is an electron beam resist with excellent lateral contrast, but with low sensitivity [8,9]. We use this resist to define devices with the desired dimensions. A top view of such a resist pattern, after litho and development is shown in figure 4. The dose used to obtain these dimensions is around 3000 $\mu\text{C}/\text{cm}^2$.

The HSQ is used to etch the underlying HPR504, which in its turn is used to transfer the pattern to a SiO₂ hardmask layer, by means of a pure CHF₃ process in a dedicated reactive ion etcher (RIE). This hardmask can then be used to etch the semiconductor material in a CH₄/H₂ ICP-RIE, figure 5.

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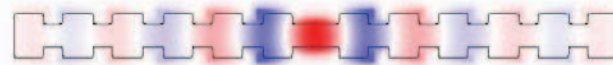


Fig. 3. Topview of the H_z-field in a 2D-FDTD simulation of a 'cold'-cavity incorporating $\lambda/4$ -shift

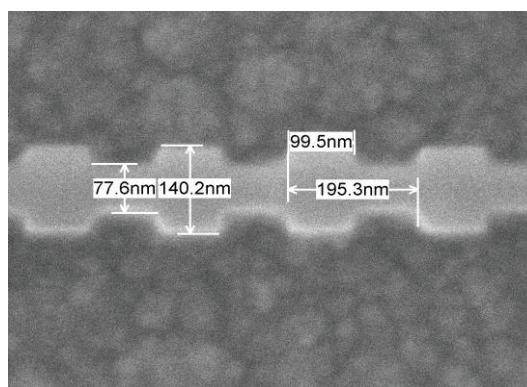


Fig. 4. Topview of HSQ resist pattern on SiO₂ after exposure and development

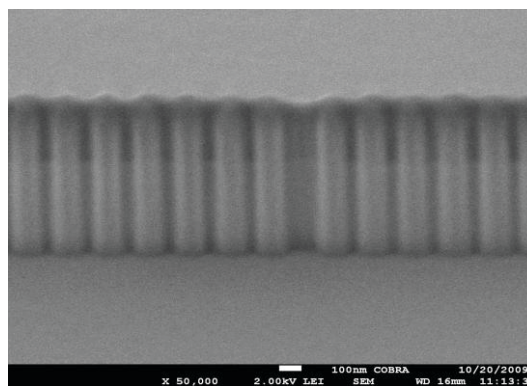


Fig. 5. Sideview of a grating in InP, the SiO₂ hardmask is still on top