

Departures from local thermodynamic equilibrium in HID lamps

Citation for published version (APA):

Mullen, van der, J. J. A. M., Nimalasuriya, T., Flikweert, A. J., Harskamp, van, W. E. N., Zhu, X-Y., Vries, de, N., Bekk, M. L., Haverlag, M., & Stoffels, W. W. (2008). Departures from local thermodynamic equilibrium in HID lamps. In *Proceedings of the 10th biennial European Plasma Conference 2008, 7-11 July 2008, Patras, Greece* (pp. HTPP10-1/110).

Document status and date:

Published: 01/01/2008

Document Version:

Publisher's PDF, also known as Version of Record (includes final page, issue and volume numbers)

Please check the document version of this publication:

- A submitted manuscript is the version of the article upon submission and before peer-review. There can be important differences between the submitted version and the official published version of record. People interested in the research are advised to contact the author for the final version of the publication, or visit the DOI to the publisher's website.
- The final author version and the galley proof are versions of the publication after peer review.
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Departures from Local Thermodynamic Equilibrium in HID lamps

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Technische Universiteit Eindhoven

Universidad de Cordoba

HTPP10
Patras
July 07, 2008

TU/e

Outline

1) Introduction: Plasma type: High Intensity Discharge lamps

2) (non-) Equilibrium aspects;

3) Polydiagnostics

4) Results

Thermal equilibrium T_e versus T_a

Ionisation equilibrium

Time dependence

5) Discussion

General aspects of (non) equilibrium in High pressure Plasmas

Applications of HID lamps I



Applications of HID lamps II



Examples of HID lamps



HID

HID lamps: High Intensity Discharge Lamps

Lighting Streets, Stadiums, Sport fields
 Car head lights

Mercury; Sodium, Xenon, ...Molecules

UHP	200 bar	100 Watt
Metal Halide	10-50 bar	100 Watt
Sulfur lamp	10 bar	1000 Watt

Typical parameters

$R < 1 \text{ cm}$

$L < 20 \text{ cm}$

$P \sim 100 \text{ W}$

Power density 10^9 W/m^3

$n_e \sim 10^{22} \text{ m}^{-3}$

$T_e \sim 6000 \text{ K}$

The plasma in MH lamps



10 bar Hg

10 mbar addition
color rendering
Efficiency:

DyI₃
better
higher

Color non-uniformity

Segregation

Metal Halide Lamps

COST project 529 Efficient lighting of the 21-th century

Many lamp types; electrons and materials

Low- & high pressure lamps

COST standard lamp
Arges project

Philips + TUE (Eindhoven)
Used in space ISS April 2004 !!

Special interest

Competition convection <-> diffusion

Example of the COST lamp



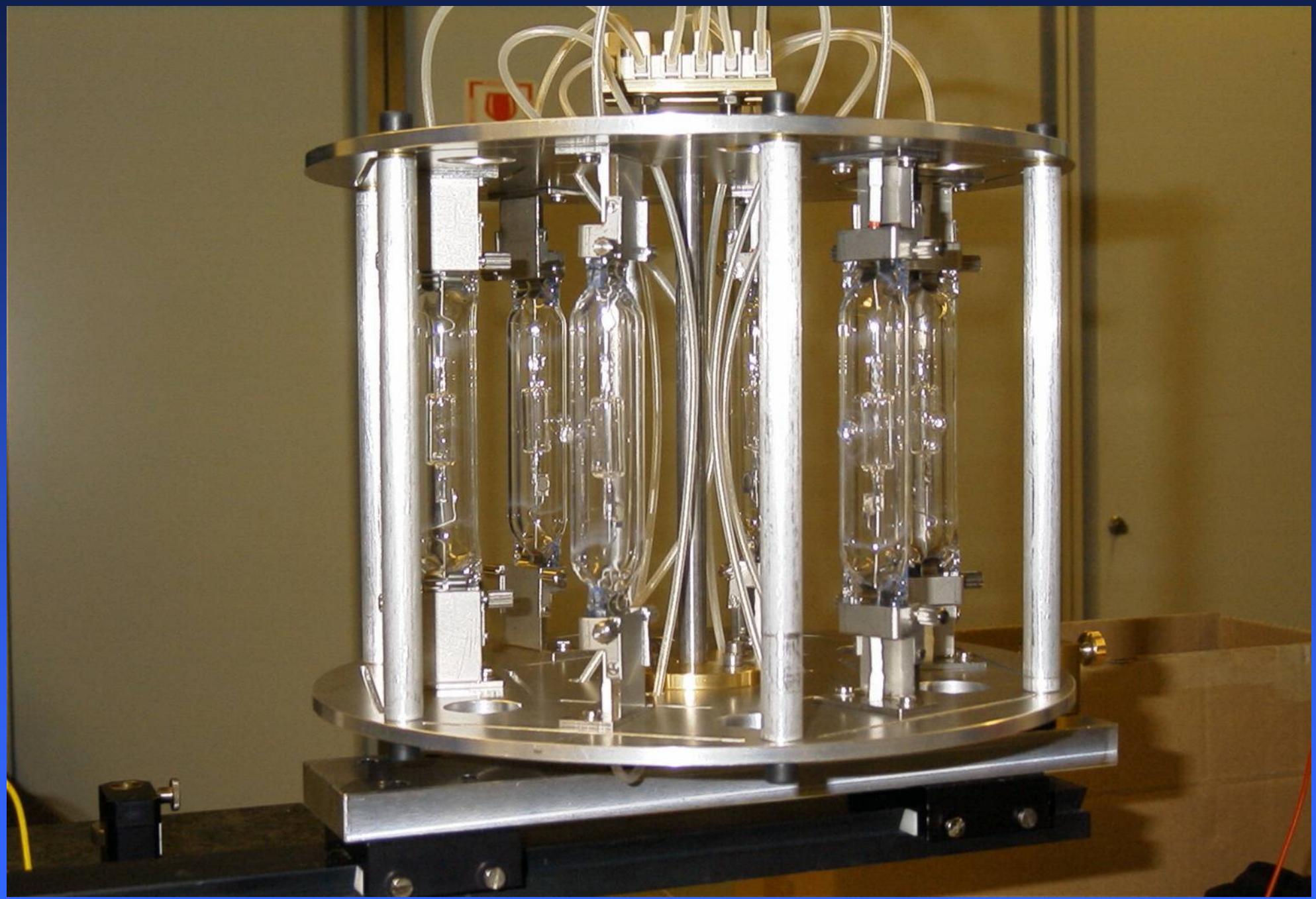
Productivity

Xiaoyan Zhu 2005

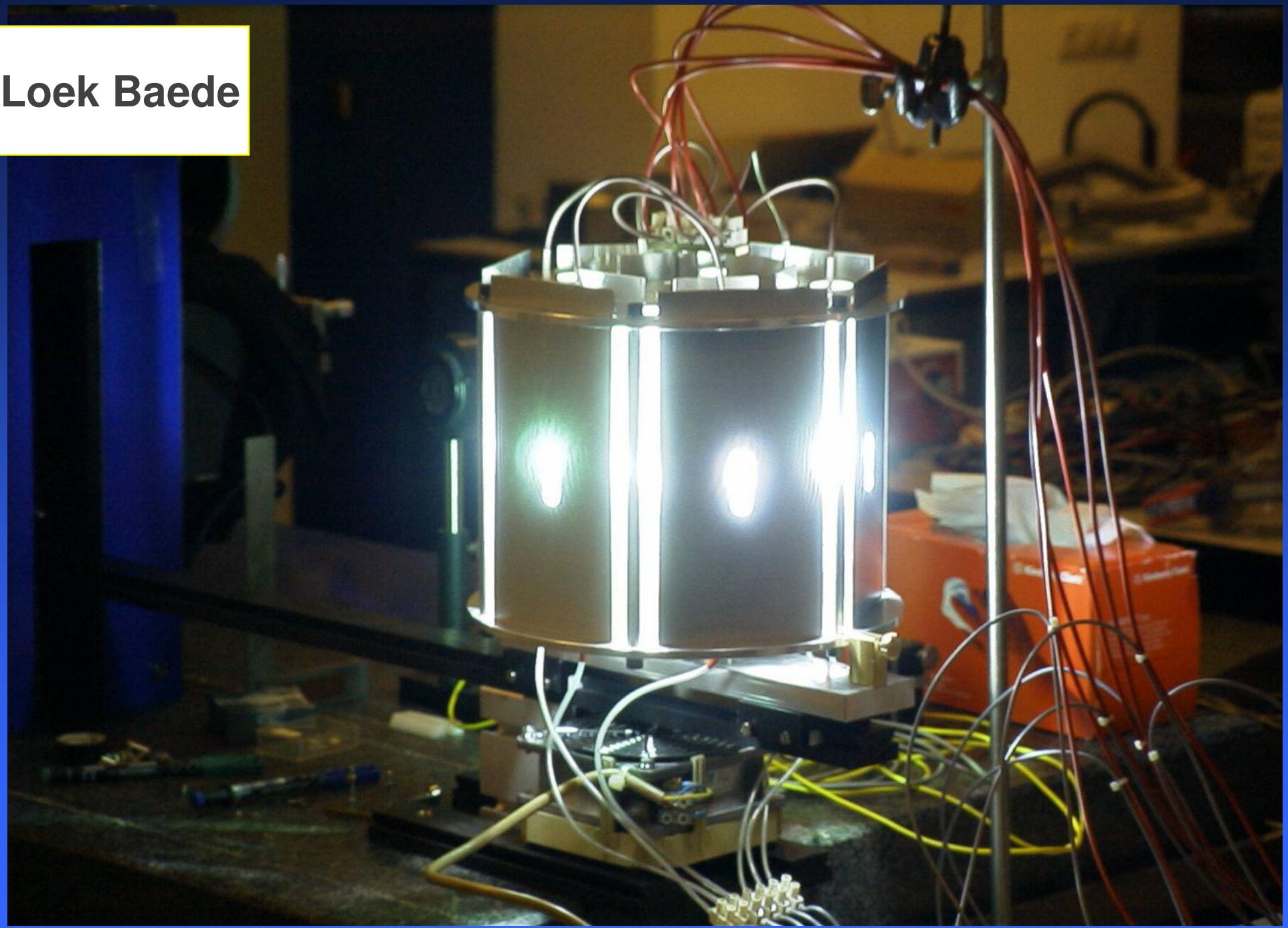
Tanya Nimalasuriya

Mark Bekk

Arjan Flikweert



Loek Baede



Zero g: methods limited



TUE (EPG)



technische universiteit eindhoven

Gerrit Kroesen, Mark Bax,
Danny van den Akker, Guido
Schiffelers, Pim Kemps, Frank
van den Hout, Marc van
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N.N.2, ... , N.N. 15

Philips (CDL)



Marco Haverlag, Rob
Keijser, Jos Eijsermans,
Jacques Claassens, Paul
Huibregts, Wally Dekkers,
Jacques Heuts, Jan Peeraer,
John Etman, Joop
Geijtenbeek, Folke
Nörtemann, Cees
Reynhout, Bruno Smets,
Hans Wernars

Astronaut: André Kuipers (ESA)



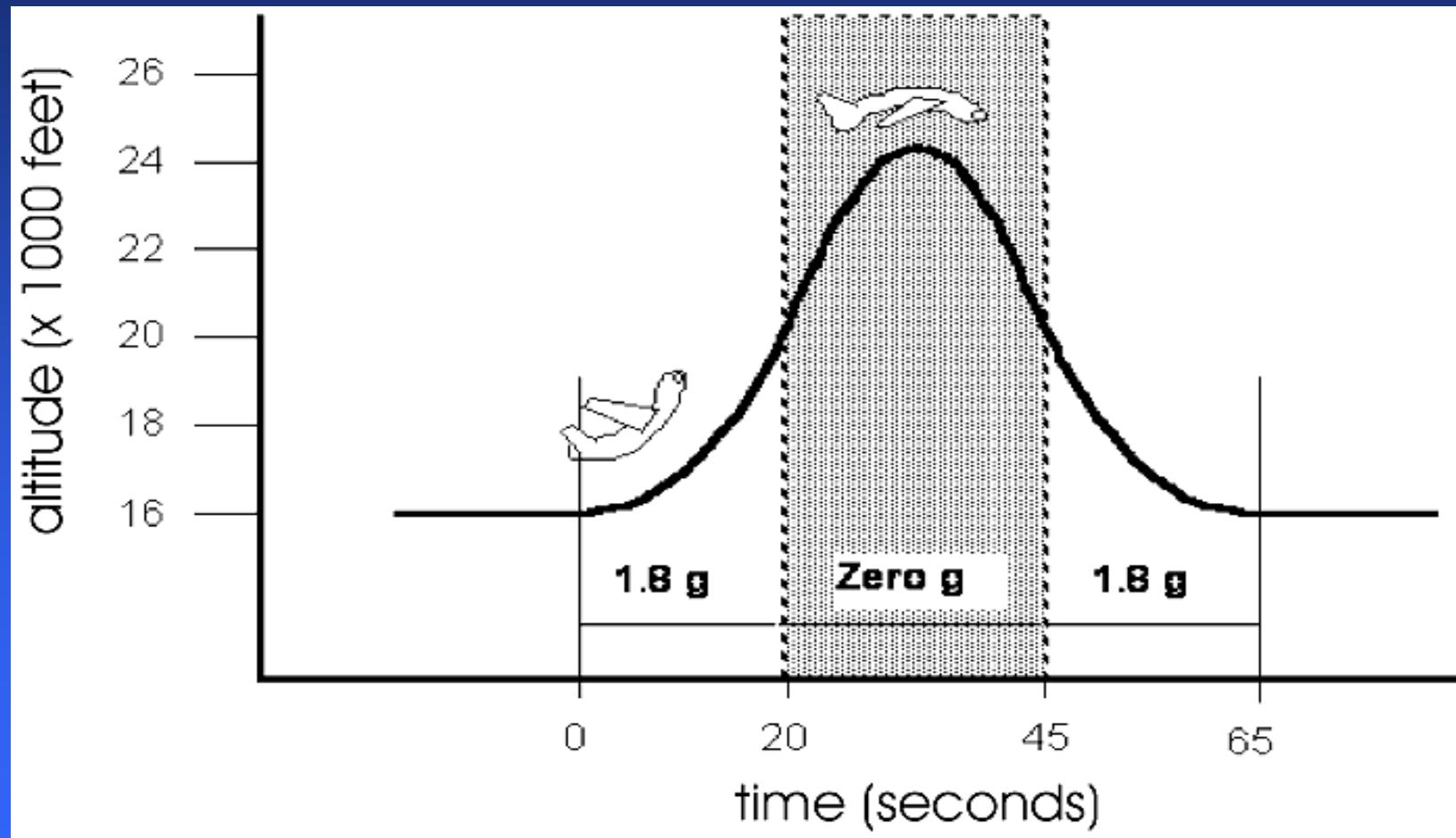
External consultants

Dutch Space: Ron Huijser, Jan Doornink, Geert Brouwer,
Fons van Wijk, King Lam, Luc van den Bergh
Kayser-Threde: Roland Seurig, Andreas Kellig
Verhaert: Piet Rosiers Bradford: Gerard Maas



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Parabolic Flights



D092T12-16-47



1 g

Equilibrium concept

LTE: Local Thermal Equilibrium

or better ?

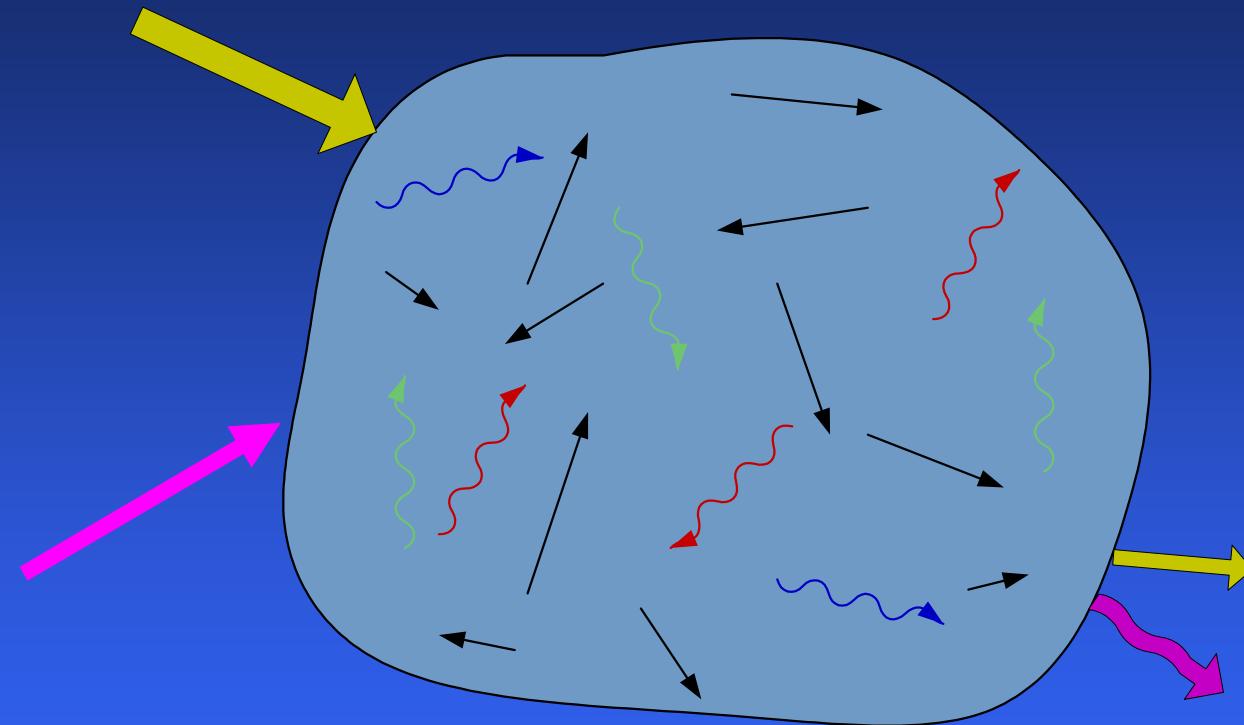
LTE: Local ThermoDynamics Equilibrium

or

LTCE: Local Thermal and Chemical Equilibrium

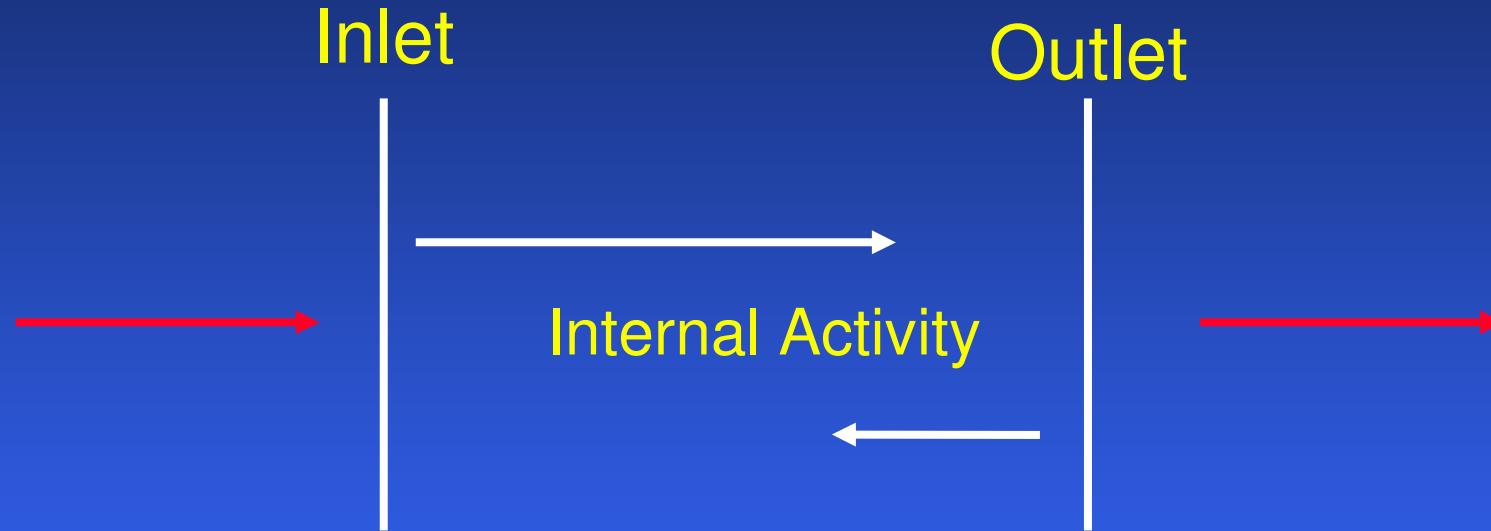
$$T_i \stackrel{?}{=} T_j$$

Plasma Artist Impression



Input and Output
Intermediated by
Vivid Internal Activity

Global Structure

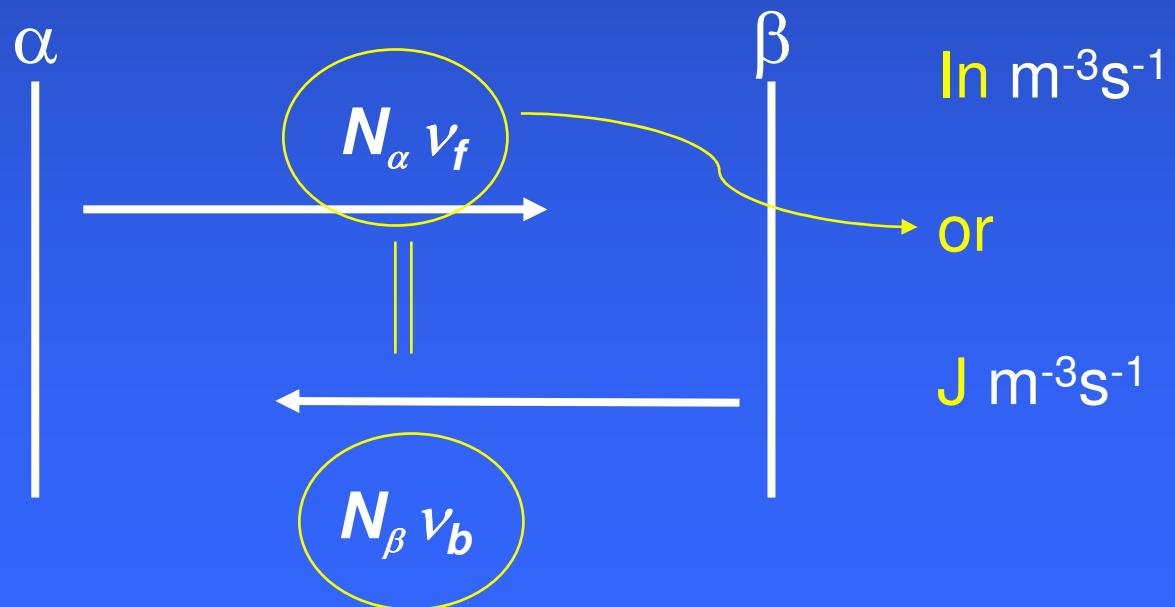


The In/Efflux couple will disturb internal Equilibrium

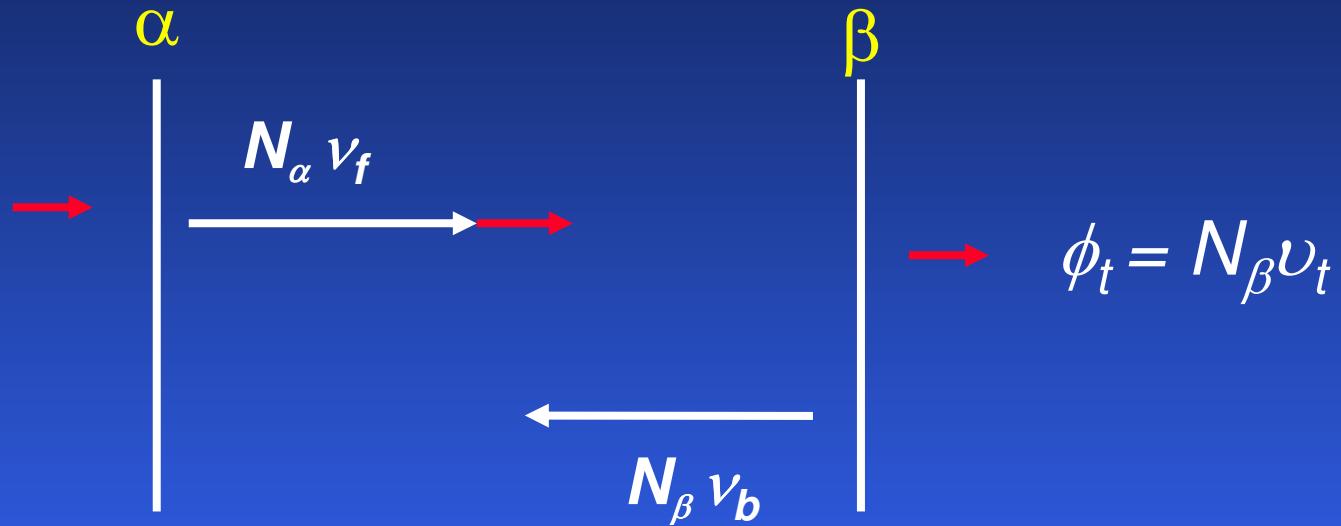
Thermodynamic Equilibrium: Collection of Bilateral Relations

TE Equilibrium in (violet) thermal dynamics

DB Equilibrium on any **level**
for any process-couple along the same route



Disturbance of BR by an Efflux



Equilibrium Condition: $v_t/v_b \ll 1$ or $v_t \tau_b \ll 1$

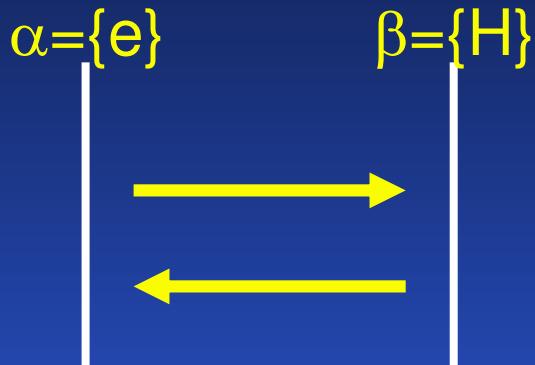
The escape per balance time must be small

EEK

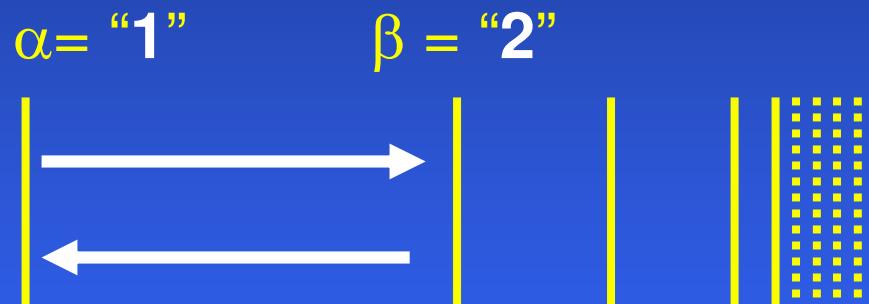
Are electrons the primary agents ??!!

EEK electron excitation kinetics dominant !?

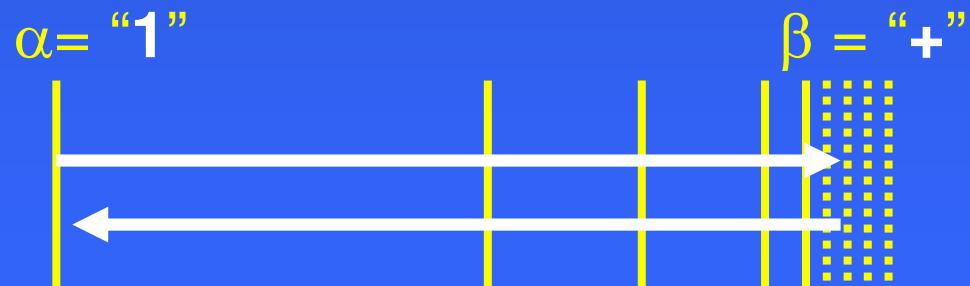
Underlying PROPER Balances



Kinetic Energy Transfer
Maxwell
 $T_e = T_h$

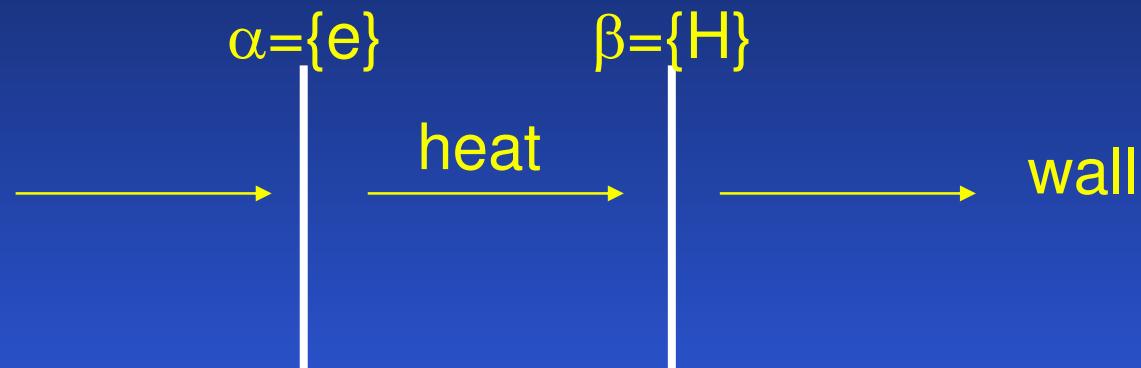


Excitation = Deexcitation
Boltzmann



Ionization = Recombin
Saha

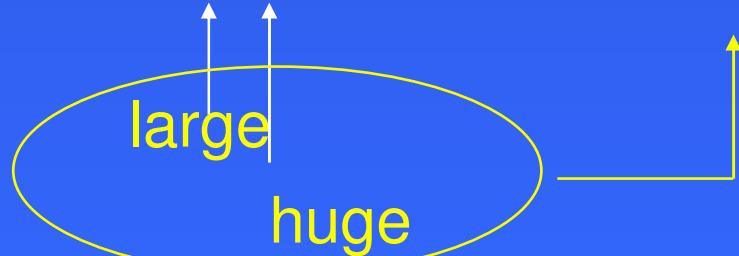
Equilibrium Disturbance



$$n_e n_1 S_{\text{heat}} (kT_e - kT_h) = \lambda / R^2 T_h$$

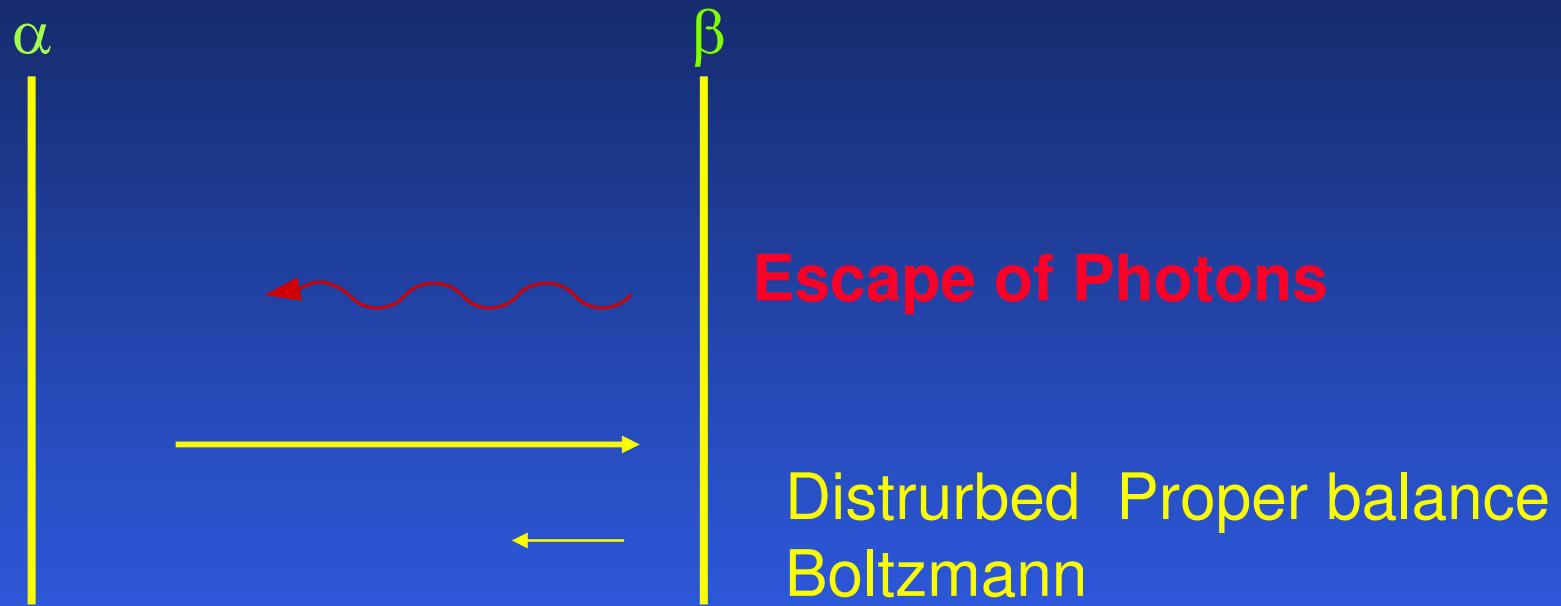
$$v_t \tau_b = \lambda / (R^2 n_e n_1 k S_{\text{heat}})^{-1} \ll 1$$

$$T_e = T_h ?$$



TU/e

Radiation escape Griem's criterion

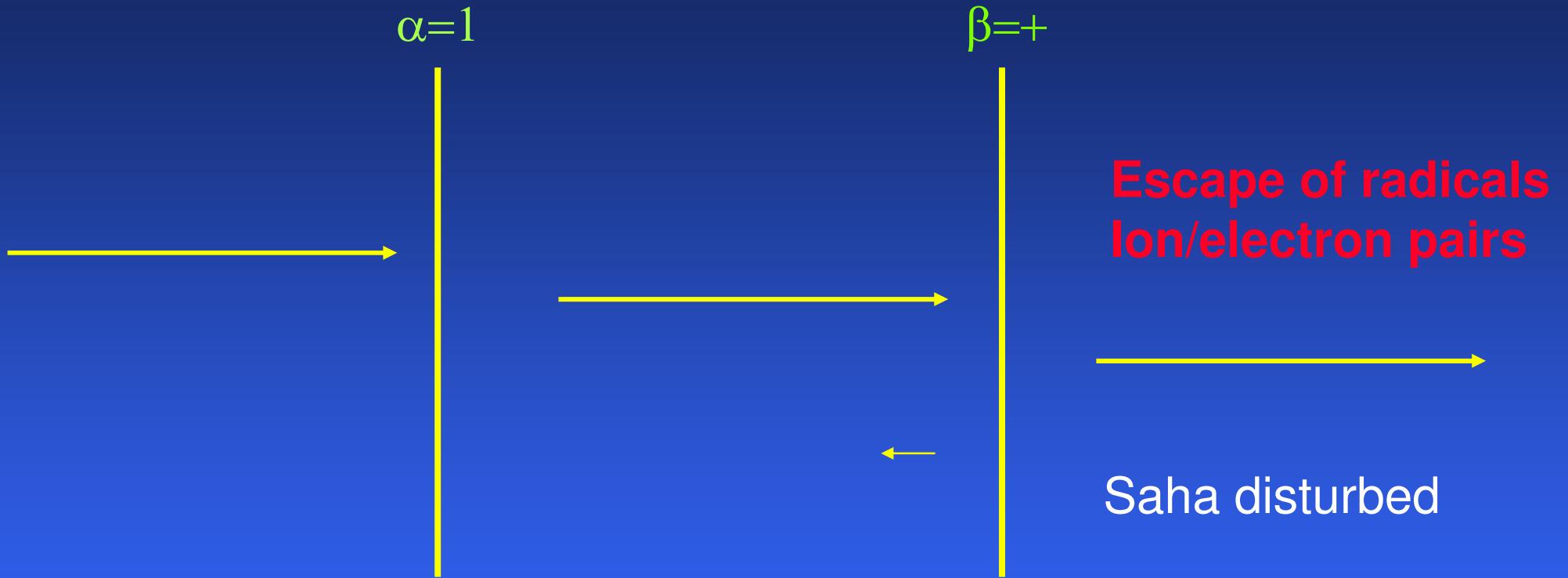


$$y(\alpha) = y(\beta)[1 + (\nu_t \tau_b)^B] \quad \text{with} \quad (\nu_t \tau_b)^B = A/n_e K(2,1)$$

Large n_e : small departure

Boltzmann present

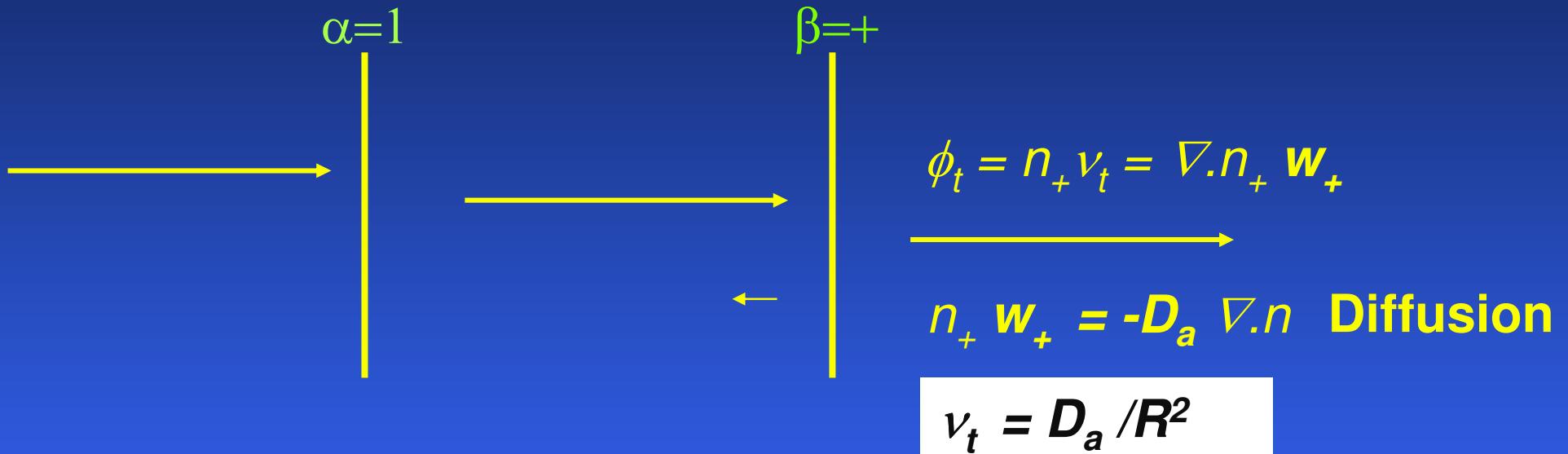
Distortion of Saha



$$y(1) = y(+) (1 + \nu_t \tau_b) \quad \text{or} \quad \delta b(1) \equiv b(1) - 1 = (\nu_t \tau_b)^S$$

$$y(1)/y(+) \equiv b(1)$$

If Ambipolar Diffusion Dominates

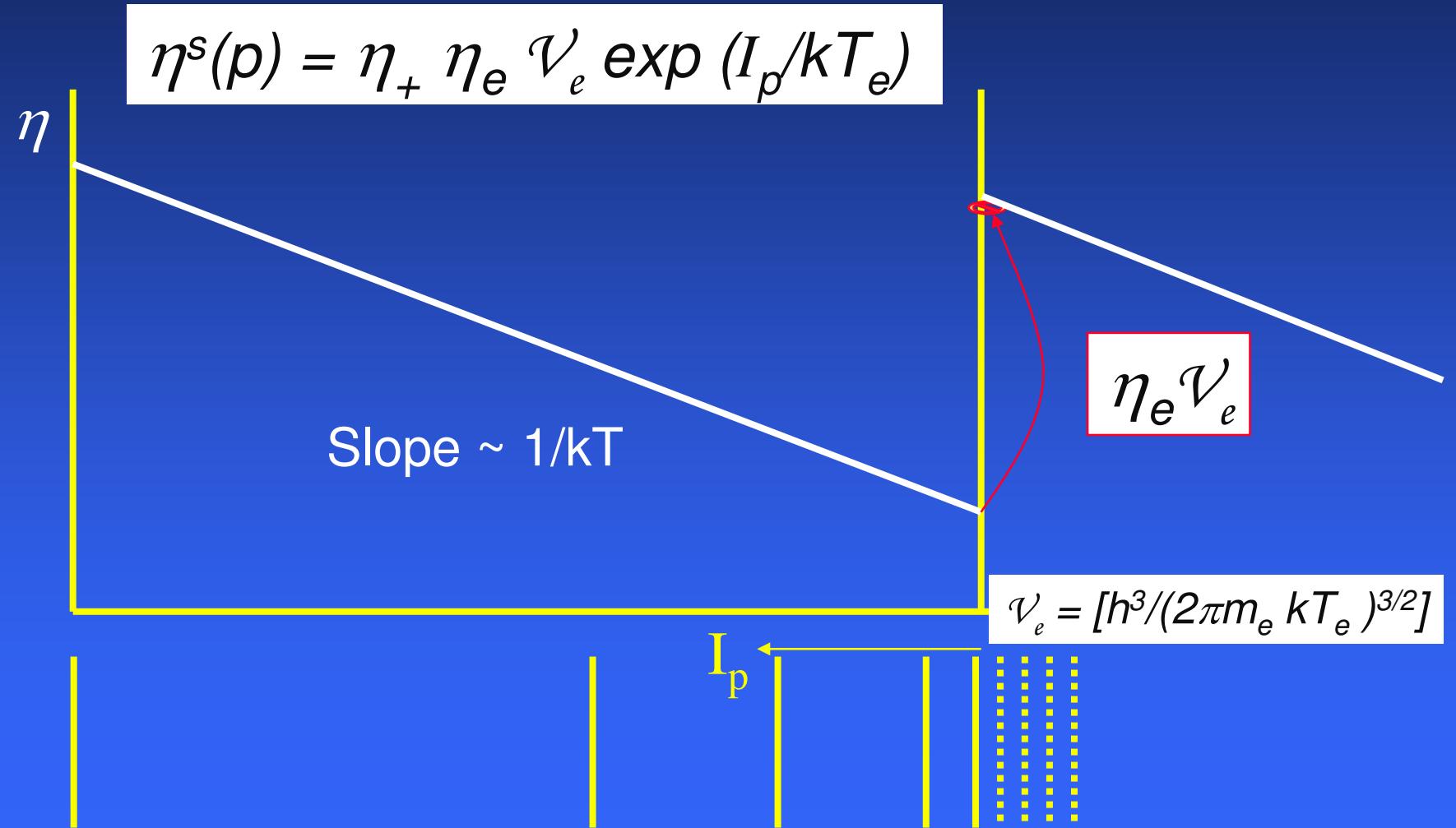


$$\delta b(1) = (v_t/v_b)^s = v_t / (n^s(1) S_{ion}) \approx C_b (A) \times 10^8 D_a (n_e R)^{-2}$$

Large n_e , very small D_a

Saha established !?

General Structure ASDF in full LSE



Polydiagnostics

Passive

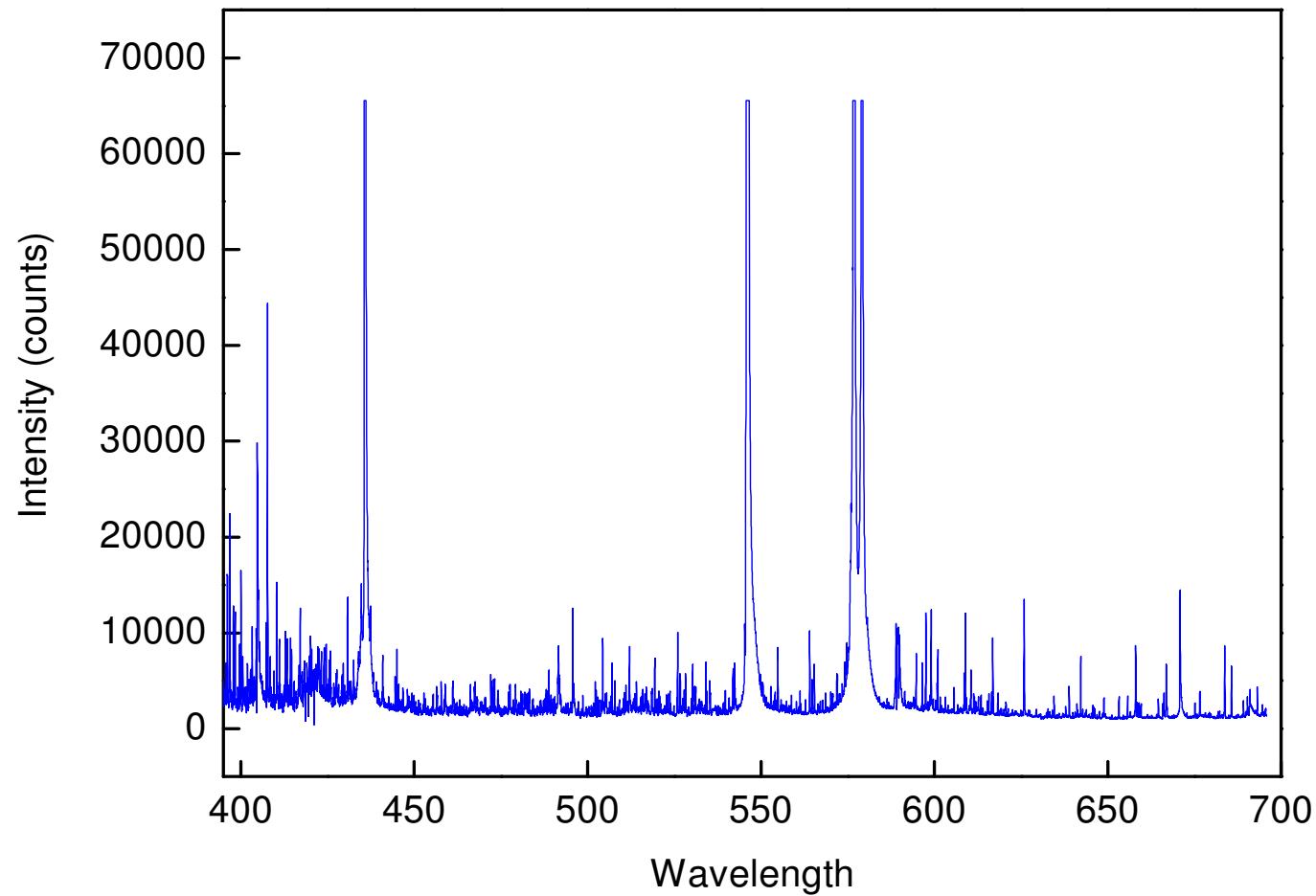
Absolute Line Intensity

Active

Thomson scattering
Xray absorption
Xray fluorescencence

Spectral Impression: grass field

Line identification: not trivial



Results Construction ASDF

ASDF constructed in

Hg

Dy I

DyII

Plasma parameters

Slopes give T
Saha jump n_e

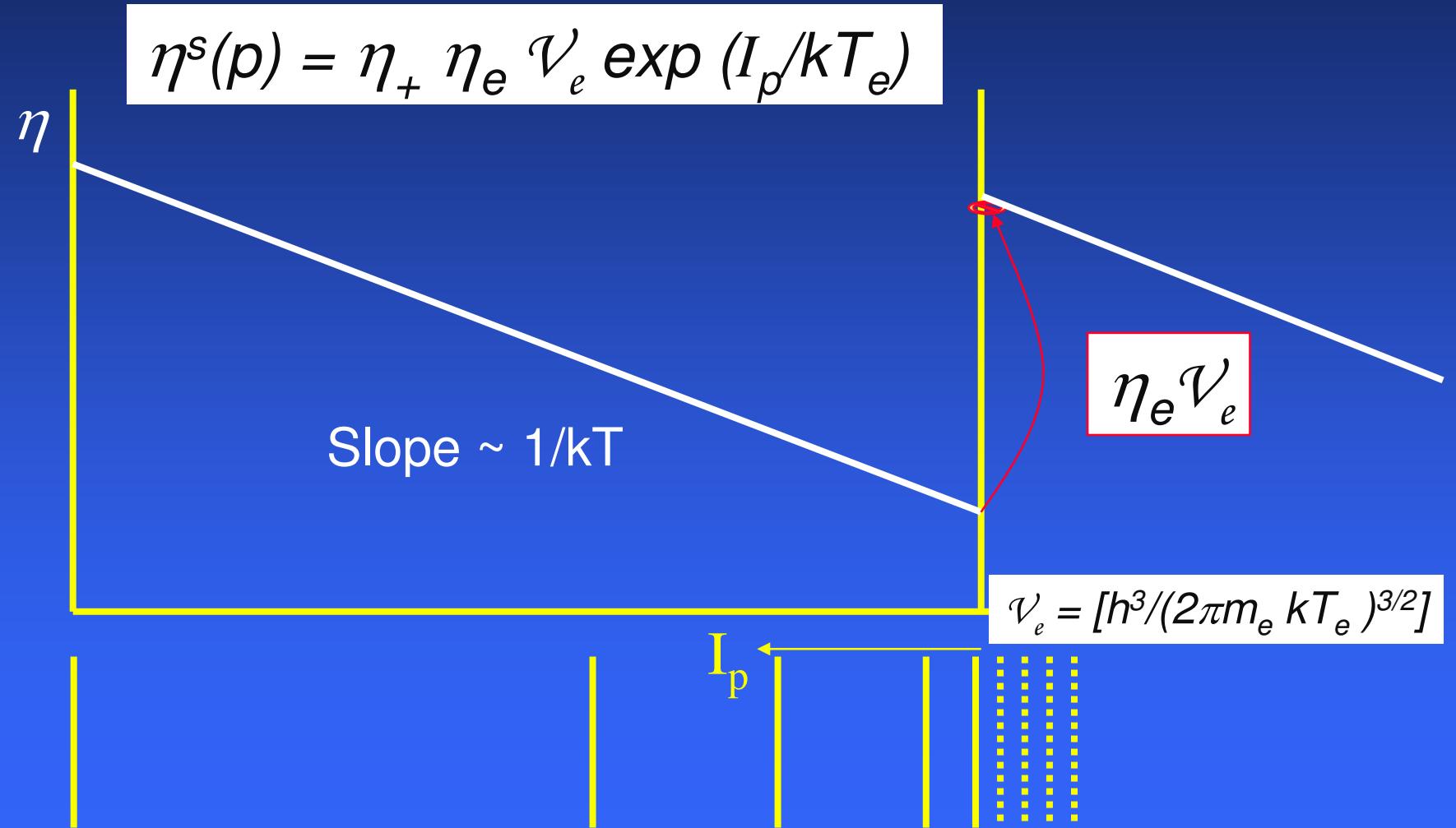
Results

in general agreement
with other methods

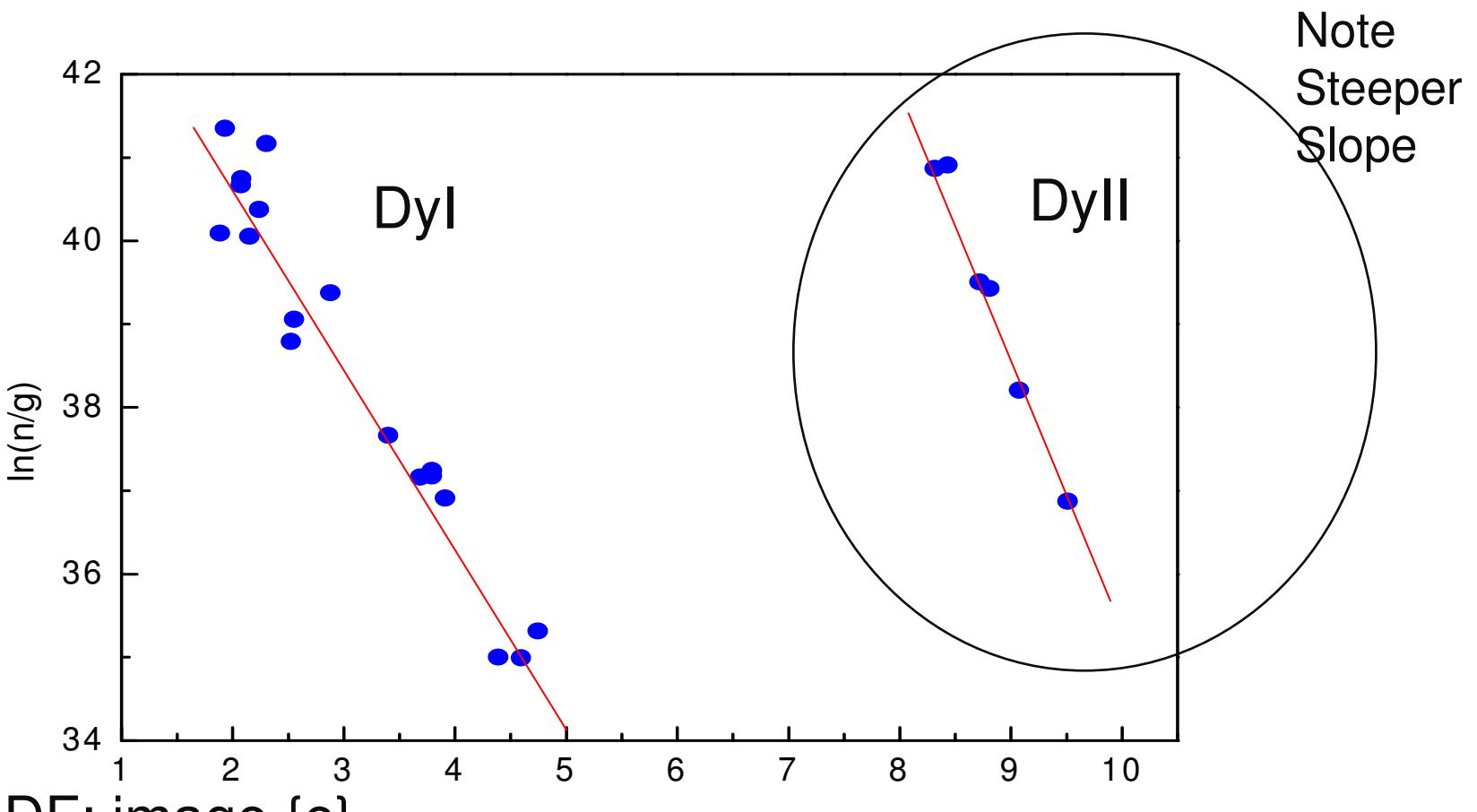
However

deviations in slope Dy II
for low n_e situations

General Structure ASDF in full LSE



ASDF in Dysprosium



ASDF: image {e}
but blurred

T=5524 K

$$\frac{n_{ground}}{g_{ground}} = 3 \cdot 10^{19} m^{-3}$$

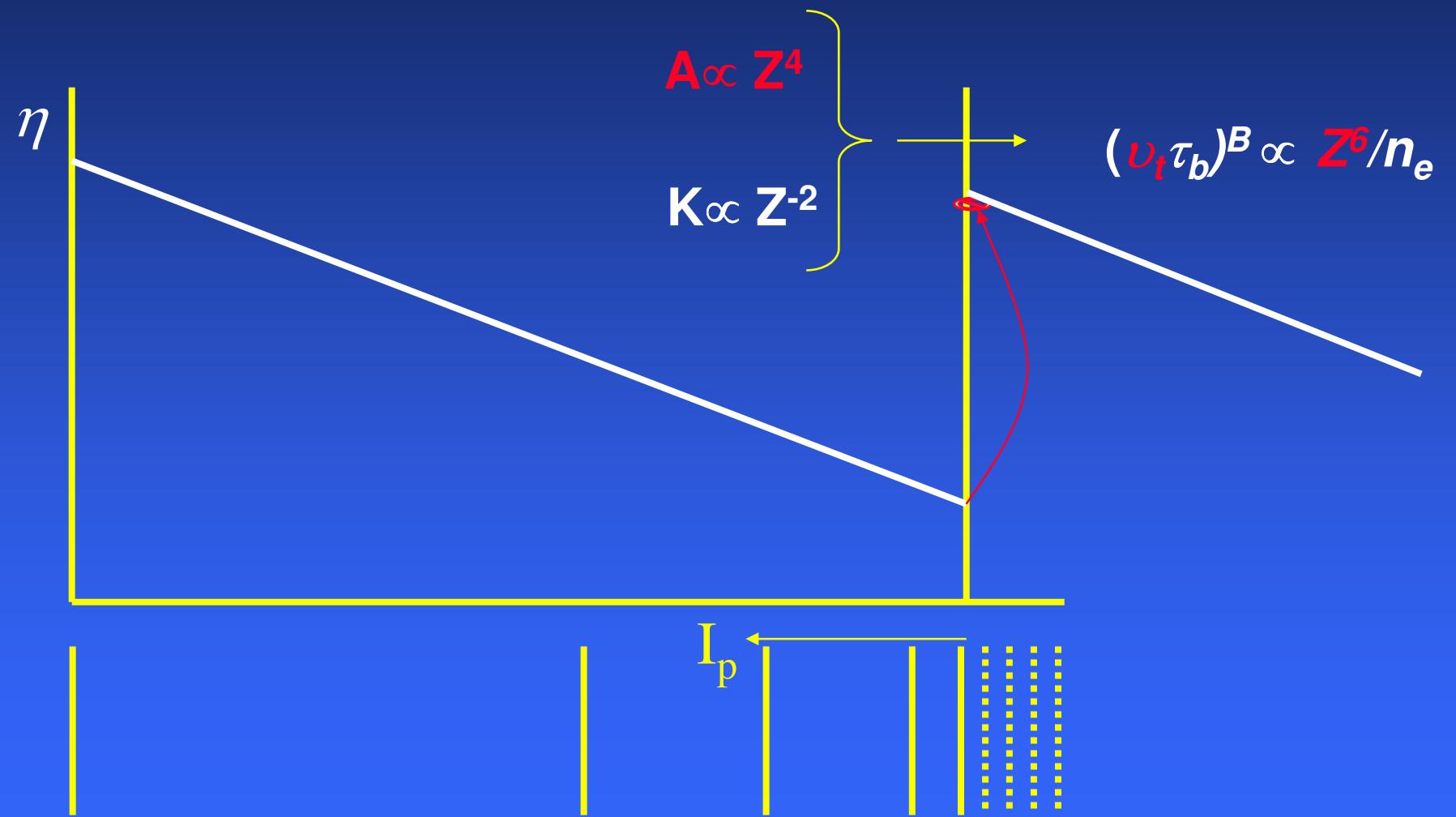
The influence of charge number Z ?

$$y(\alpha) = y(\beta)[1 + (\nu_t \tau_b)^B] \quad \text{with} \quad (\nu_t \tau_b)^B = A/n_e K(2,1)$$

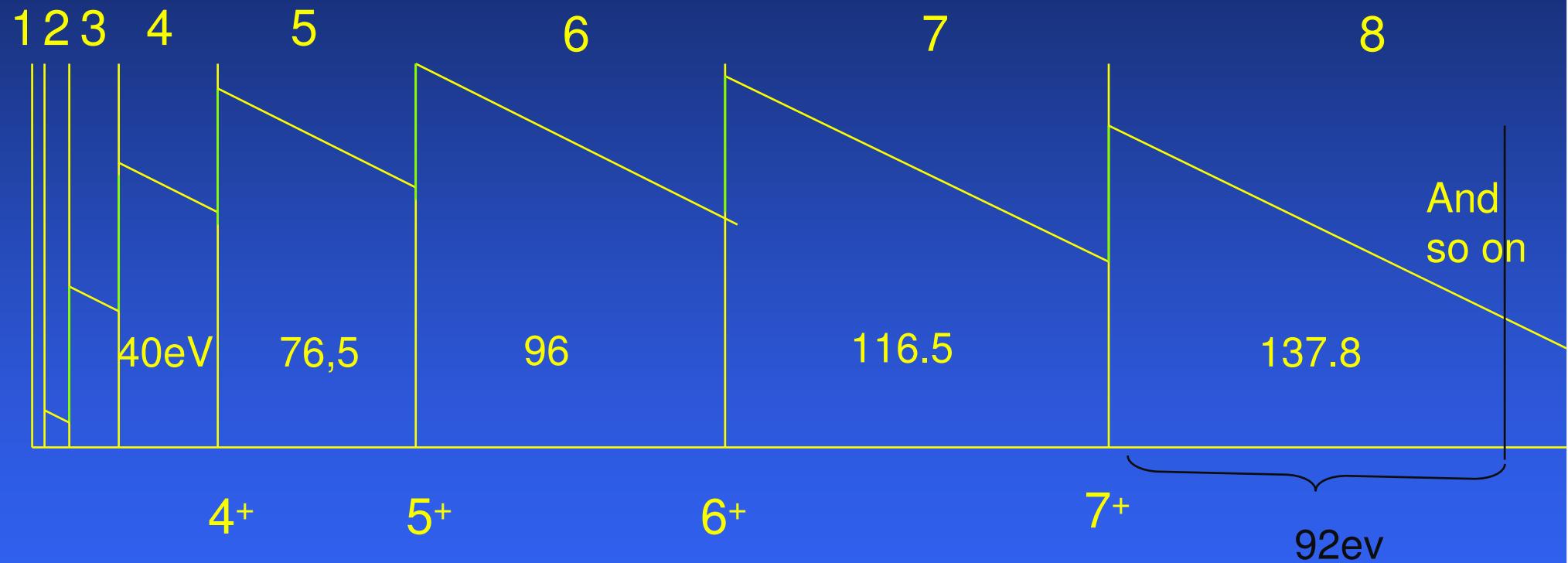
$$\left. \begin{array}{l} A \propto Z^4 \\ K \propto Z^{-2} \end{array} \right\} \longrightarrow (\nu_t \tau_b)^B \propto Z^6/n_e$$

Increasing Z → increase importance radiative escape!!

General Structure ASDF in full LSE



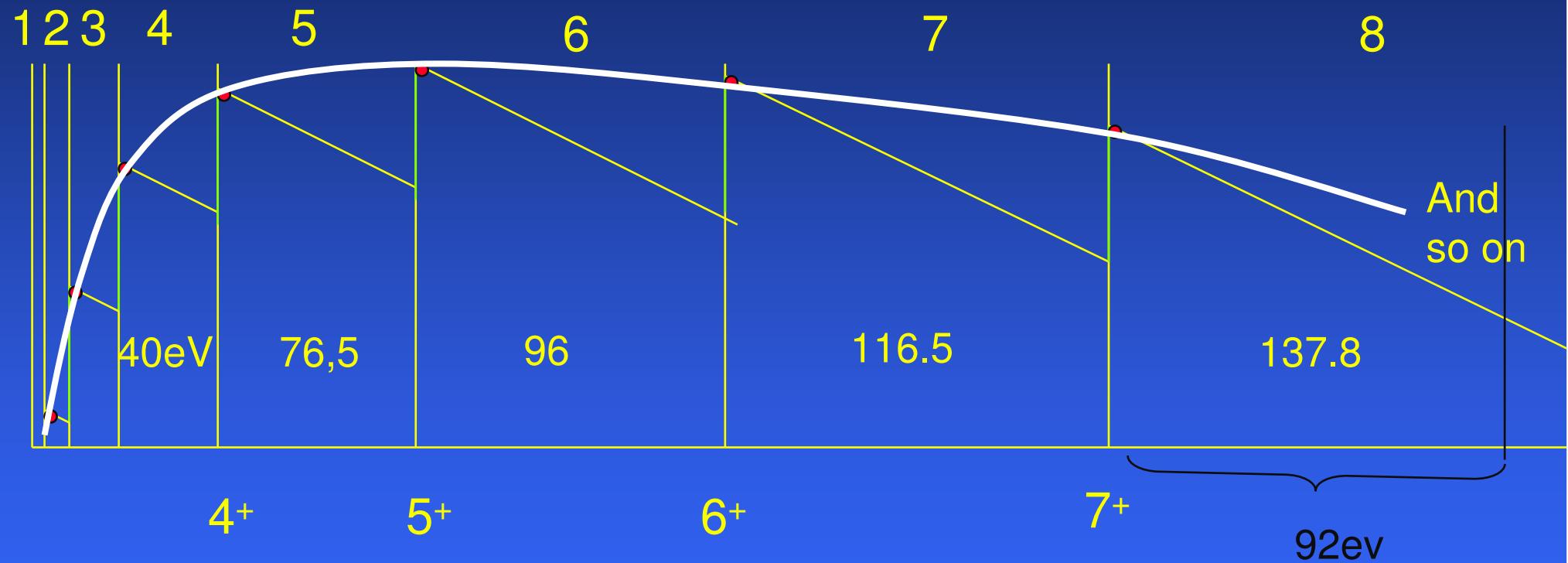
Comparison with EUV plasmas in Sn wanted 92eV radiation



$$= (n_e/2) [h^3/(2\pi m_e k T_e)^{3/2}]$$

TU/e

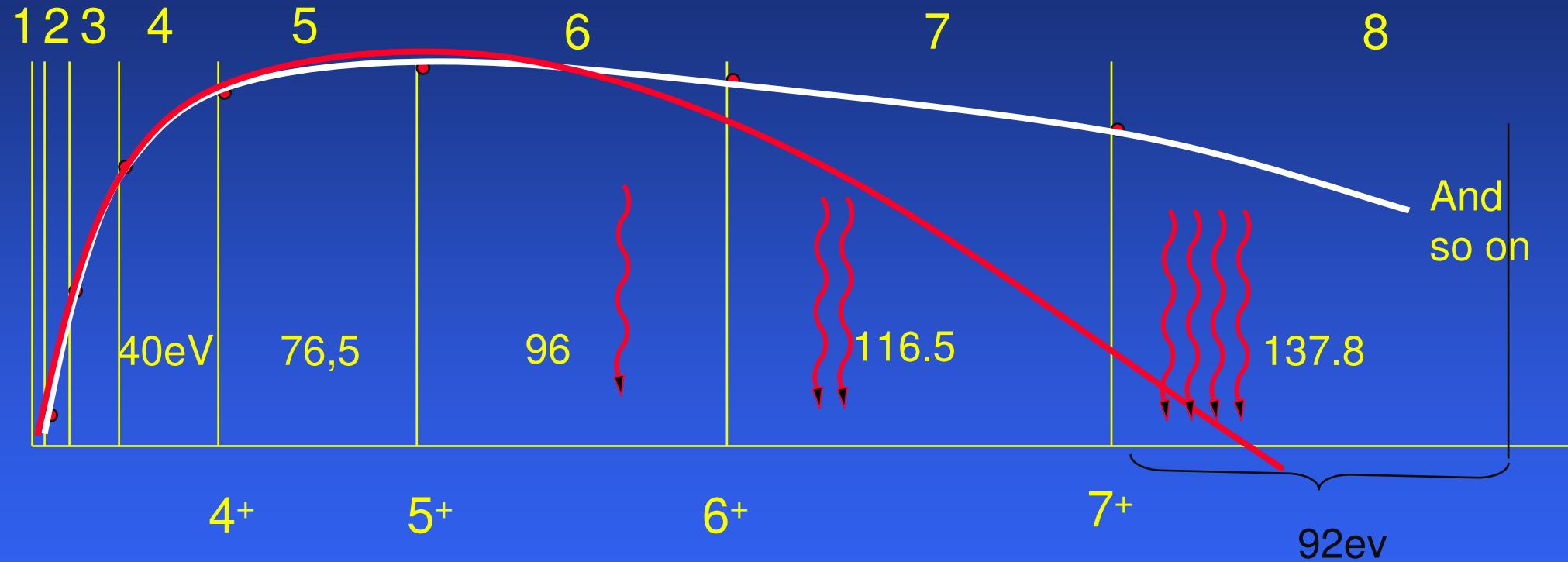
If Saha remains present



$$= (n_e/2) [h^3/(2\pi m_e k T_e)^{3/2}]$$

TU/e

Influence radiation



$$= (n_e/2) [h^3/(2\pi m_e k T_e)^{3/2}]$$

TU/e

For increasing Z
there will be a stage for which radiation escape
will disturb the ASDF

However these states are already low populated
No change in light generating properties

Purely Academic

Presence of $T_e = T_g$

Relevant question for modeling

One- or two-Temperature plasma

Wanted two methods:

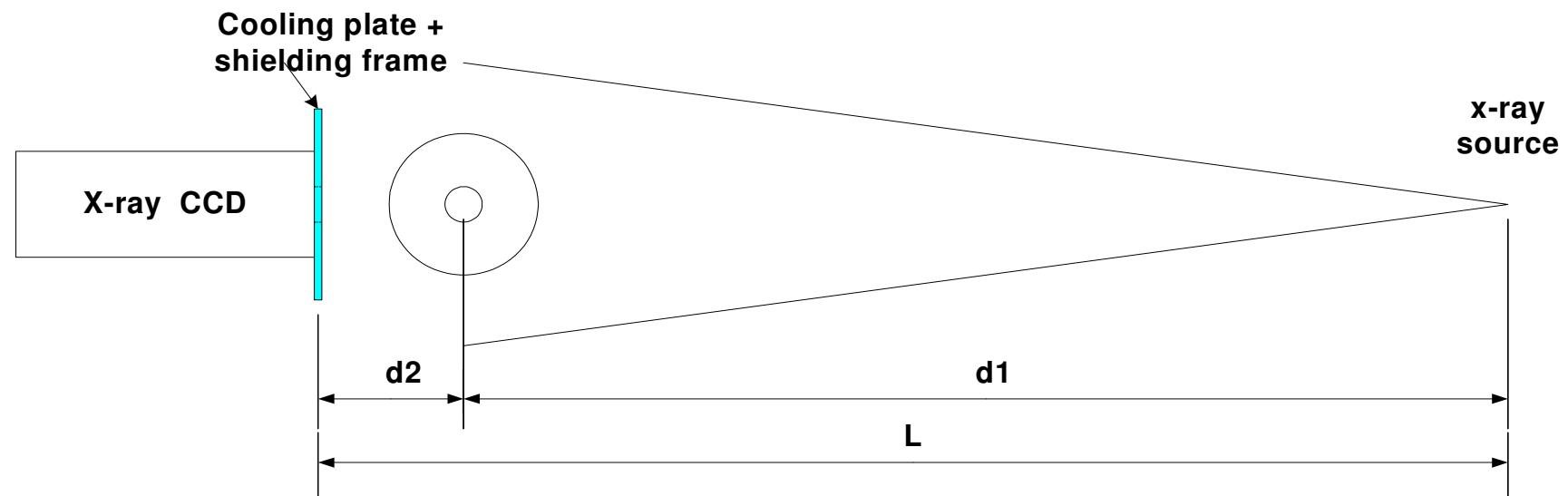
T_e

T_g

Thomson Sc
Xray absorption

X-ray absorption

Xiaoyan Zhu, Tanya Nimalasuriya
Marco Haverlag; Evert Ridderhof



Procedure

Hg is dominant



$$(n T)_{\text{any pos}} = (n T)_{\text{wall}}$$



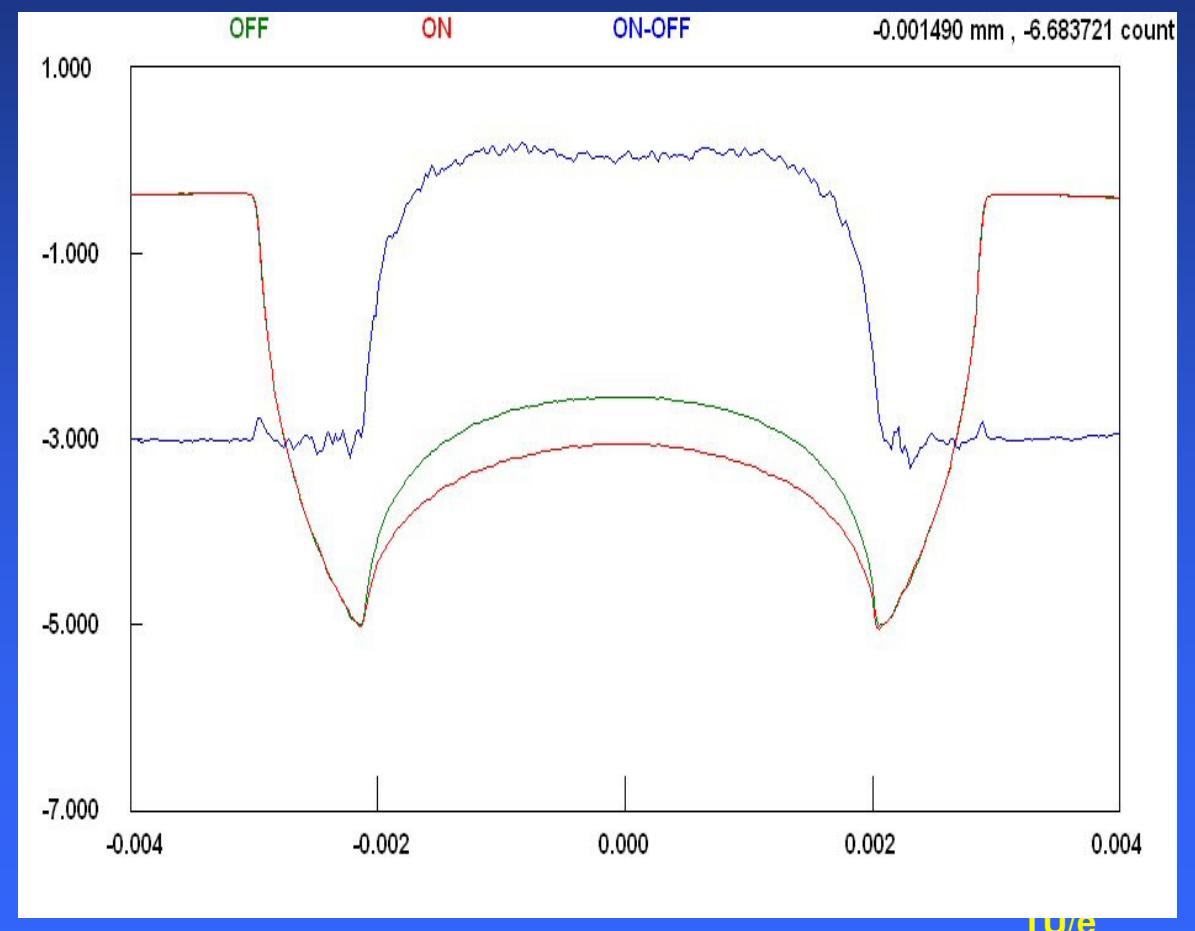
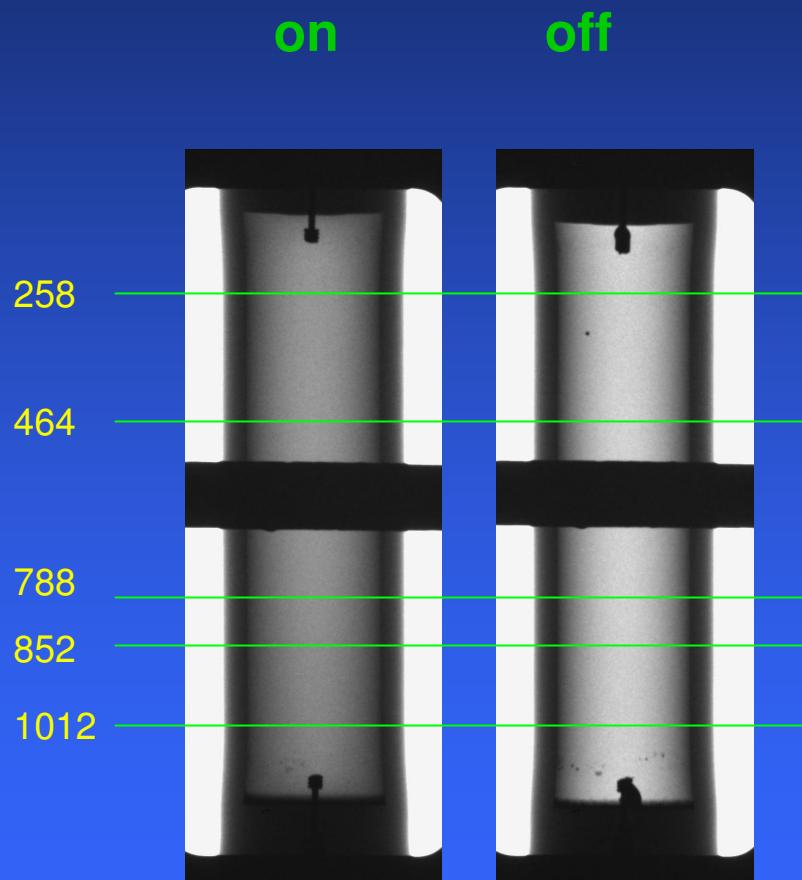
Xray

Pyrometer

T_g on any position

XRA on Helios lamp

- Exposure time: 200s.



Thomson scattering: expensive but always surprise surprise

Expensive: Laser system + Spectrograph+ ICCD

Results: Interpretation independent non-equilibrium model

Scattering on electron gas: Real n_e and T_e

Always a surprise

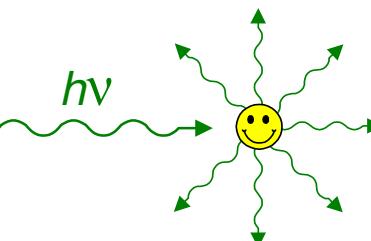
General structure set-up

plasma			
	Laser	Spect	Detector
Academic Industrial			
1972	Ruby	Mono	PMT-array (7)
Xxx	xxx	xxxx	xxxx
1994	YAG	Mono	IPDA (1064)
2000	YAG	TGS (1eV)	ICCD (500x700)
2005	YAG (200 ps)	TGS (30eV)	ICCD ()

Thomson Scattering

⇒ scattering of photons on
free electrons in a plasma

More direct
measurement



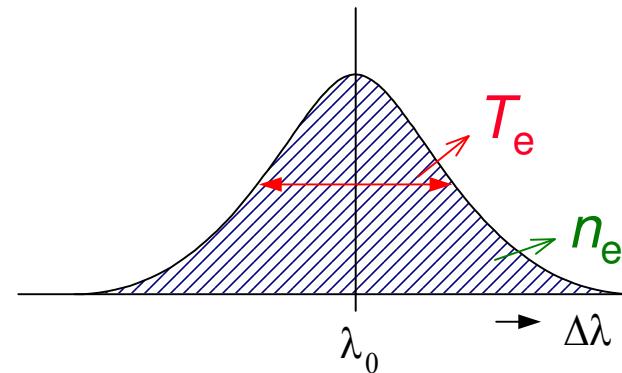
Scattering Doppler
intensity broadening



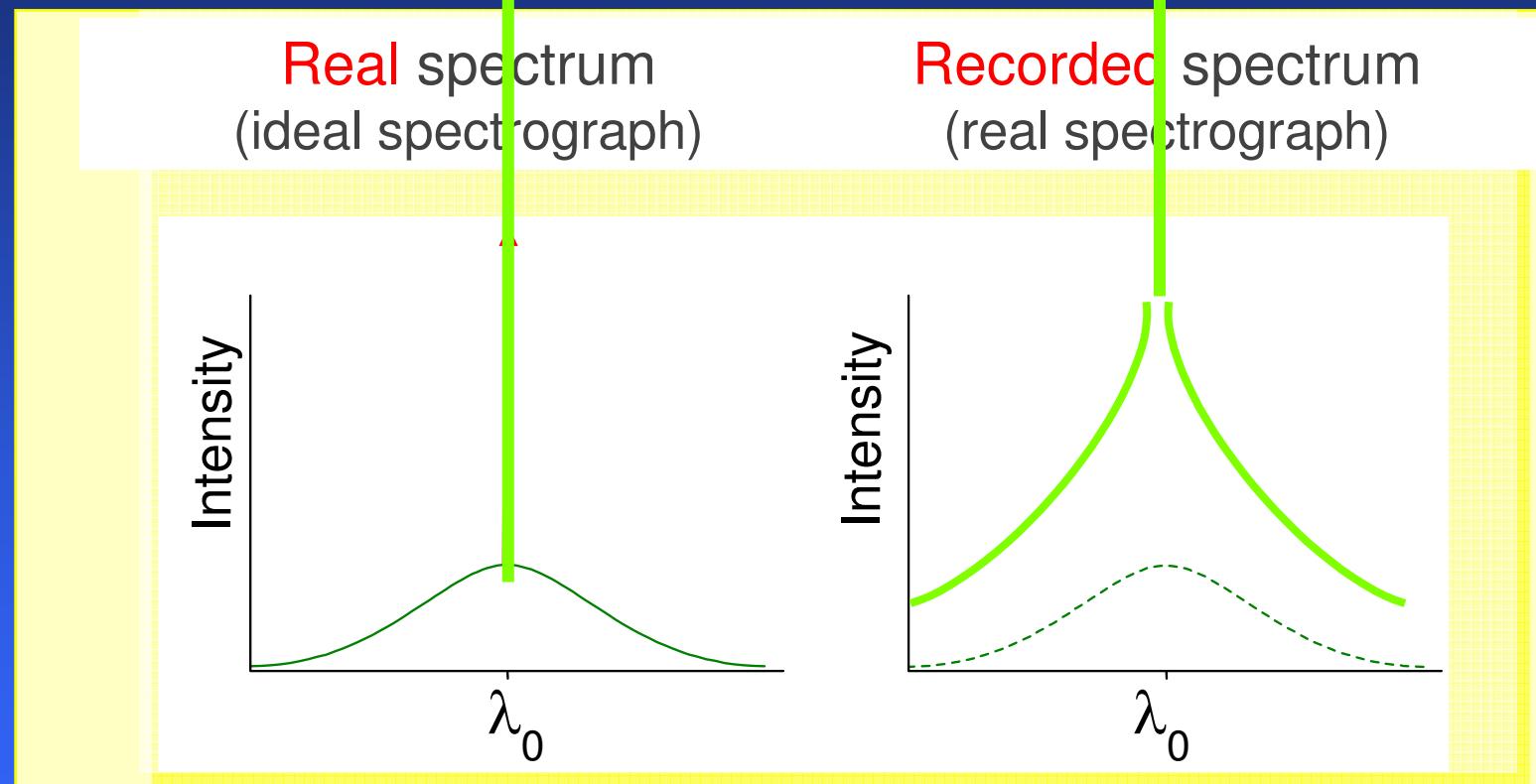
n_e



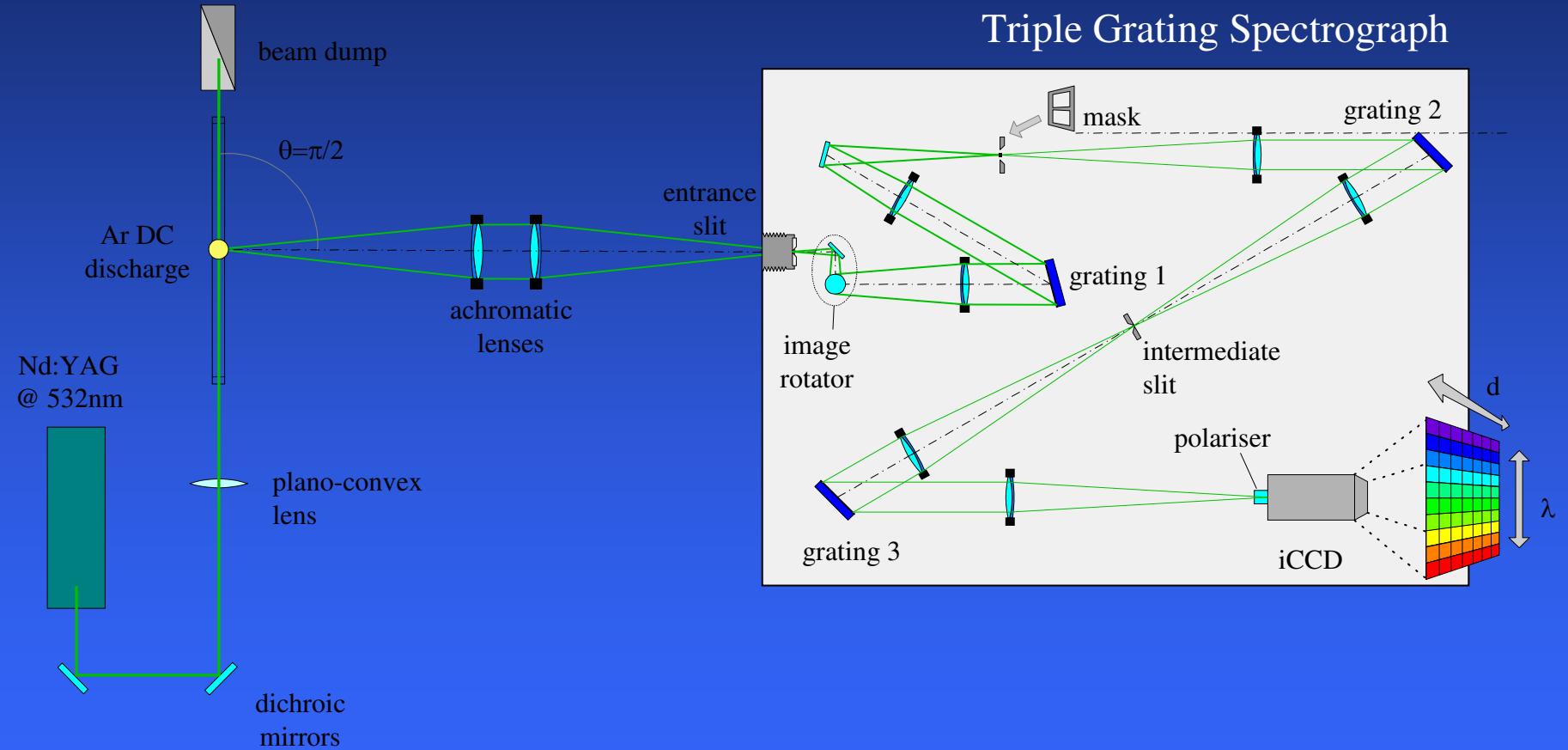
T_e



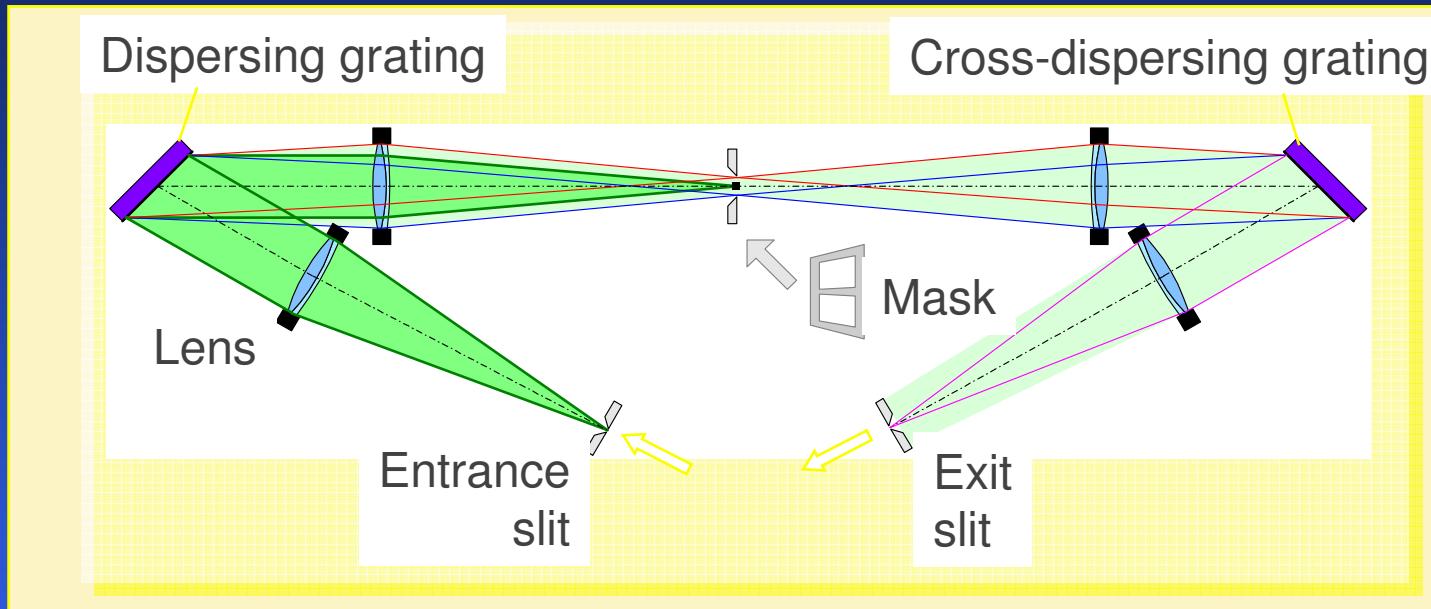
Redistribution of monochromatic light: Instrumental profile



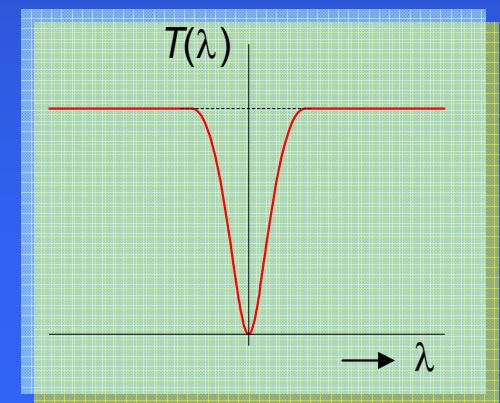
Triple Grating Spectrograph



Home-made Filter



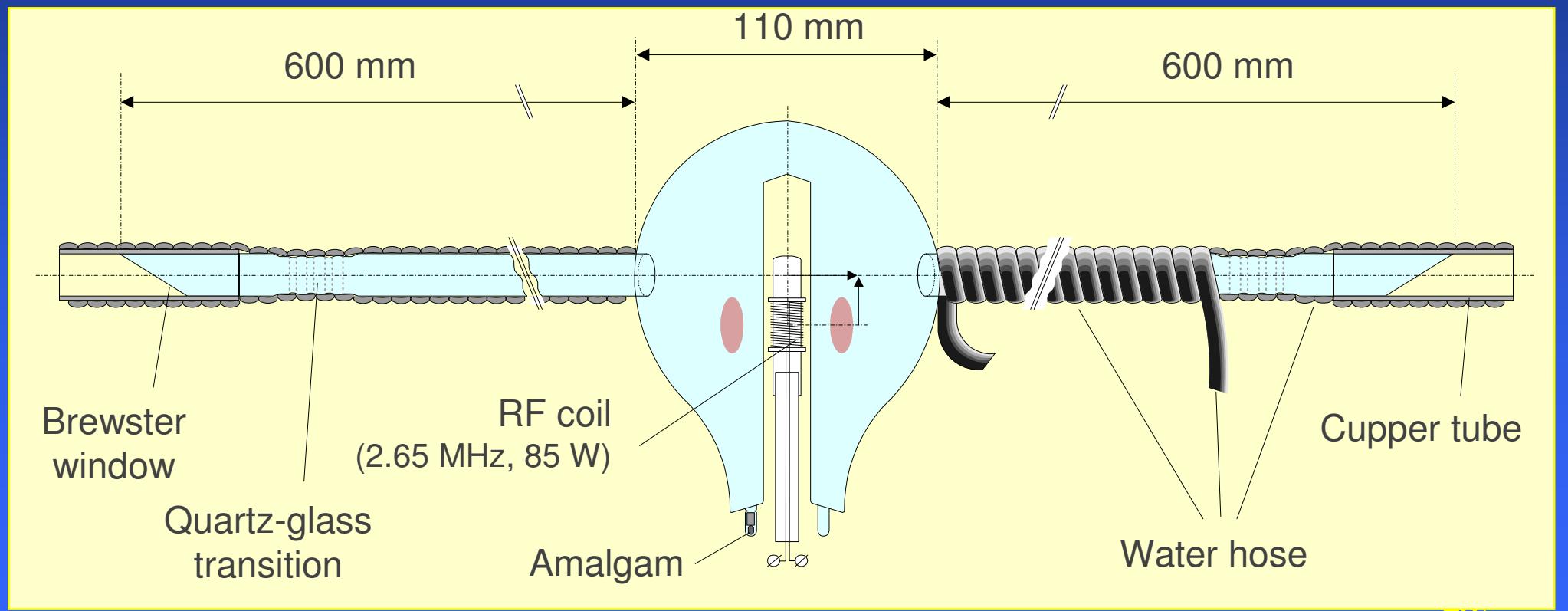
Transmission function



TU/e

Experimental Version QL

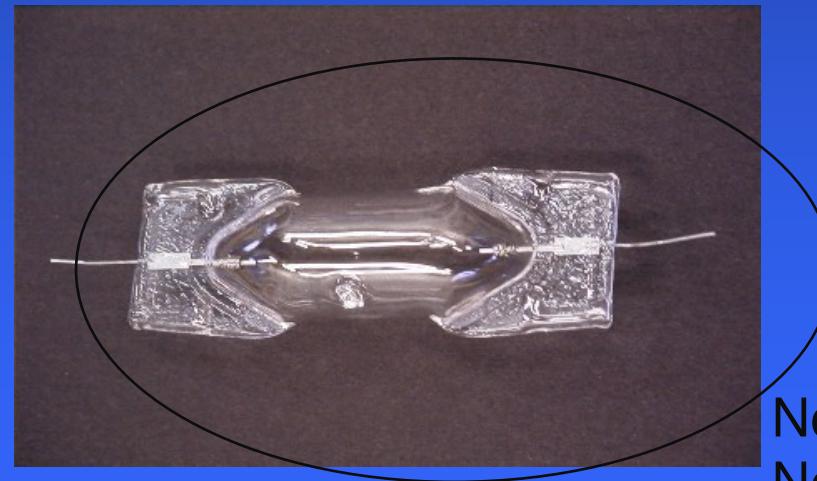
- Extension tubes
- Quartz, Brewster angle windows
- Heating
- Pressure correction



“Real” Lamp

400W high pressure Hg lamp

Inner diameter is 20 mm.

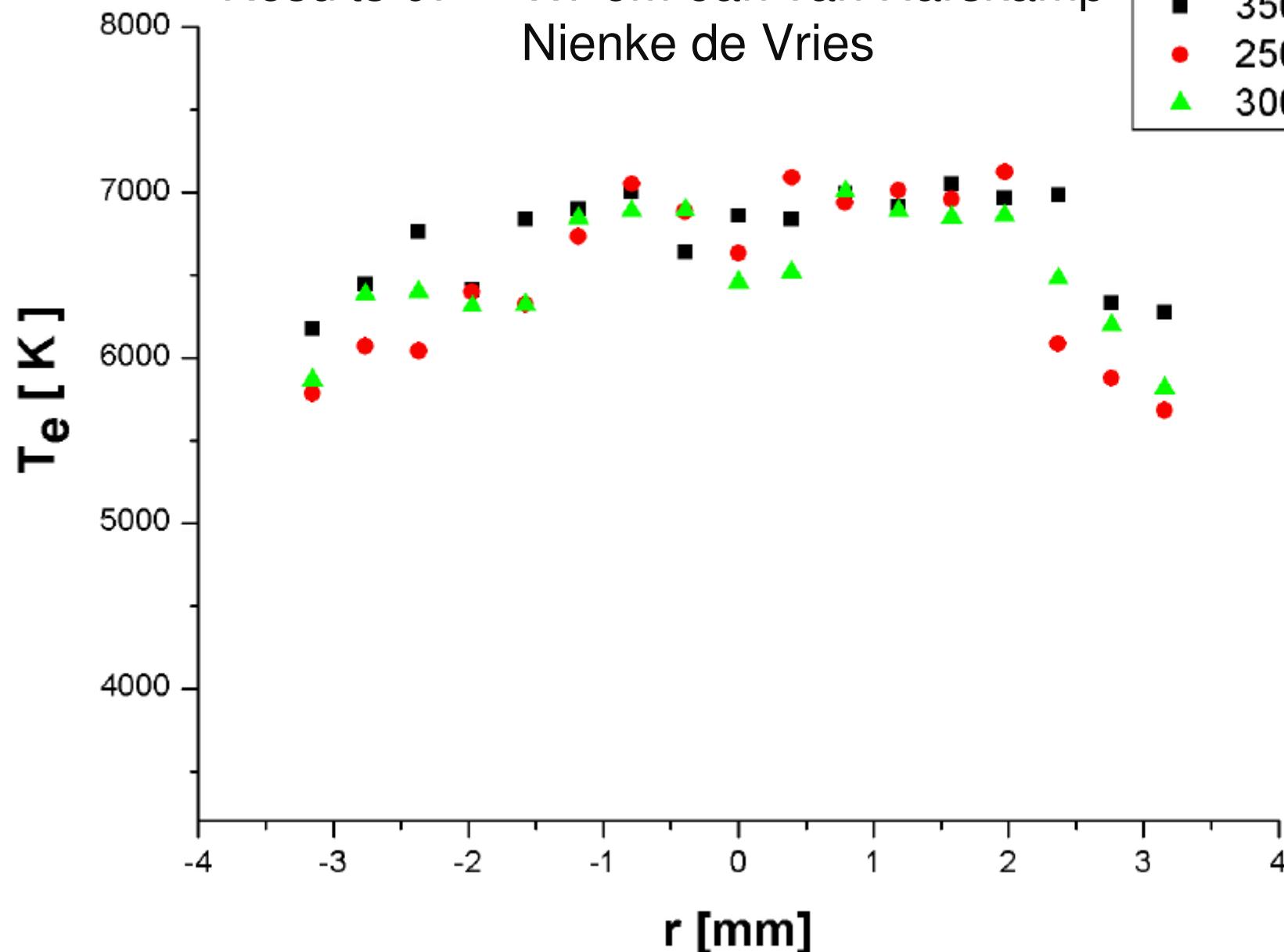


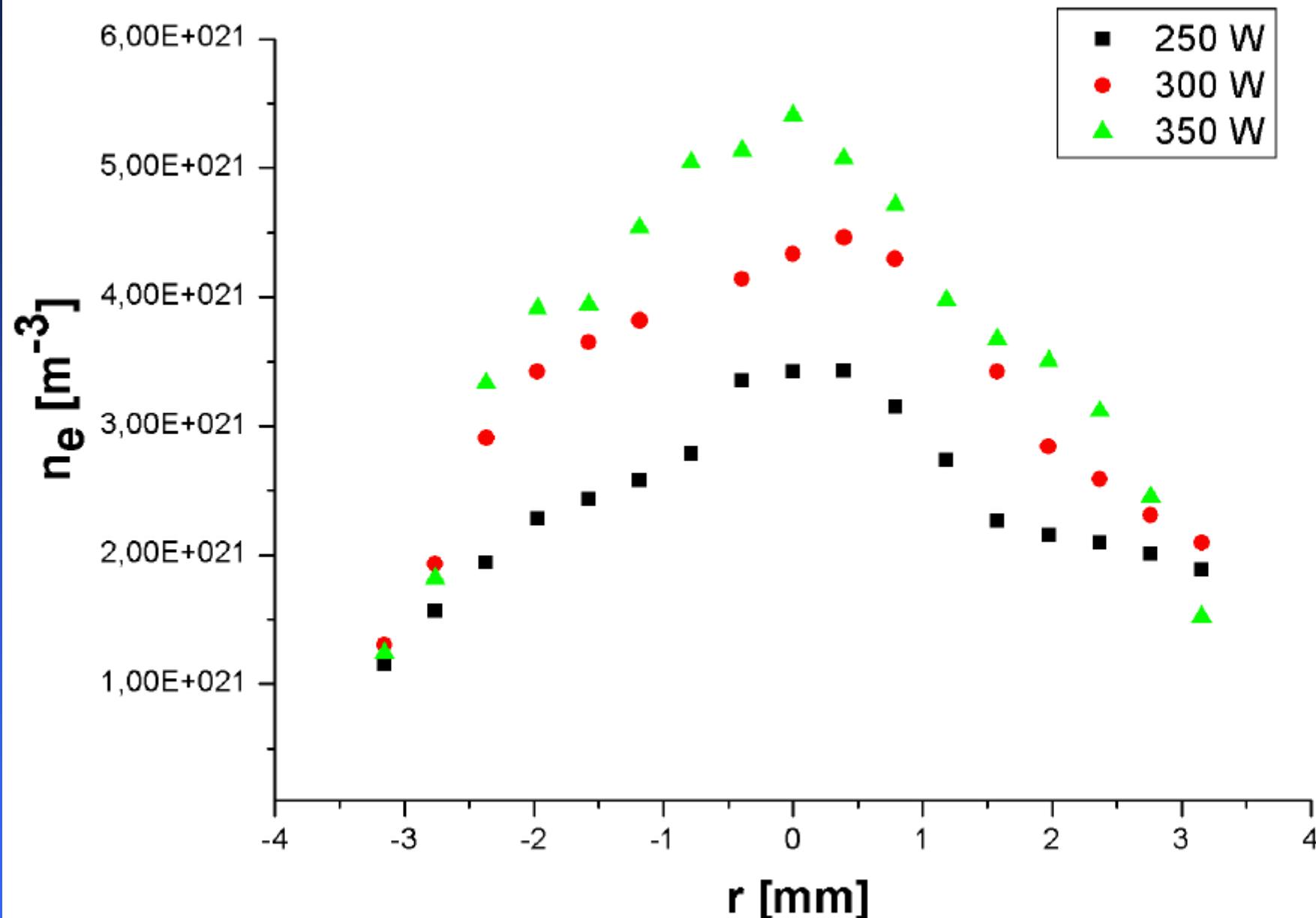
Note
Not a MH Lamp

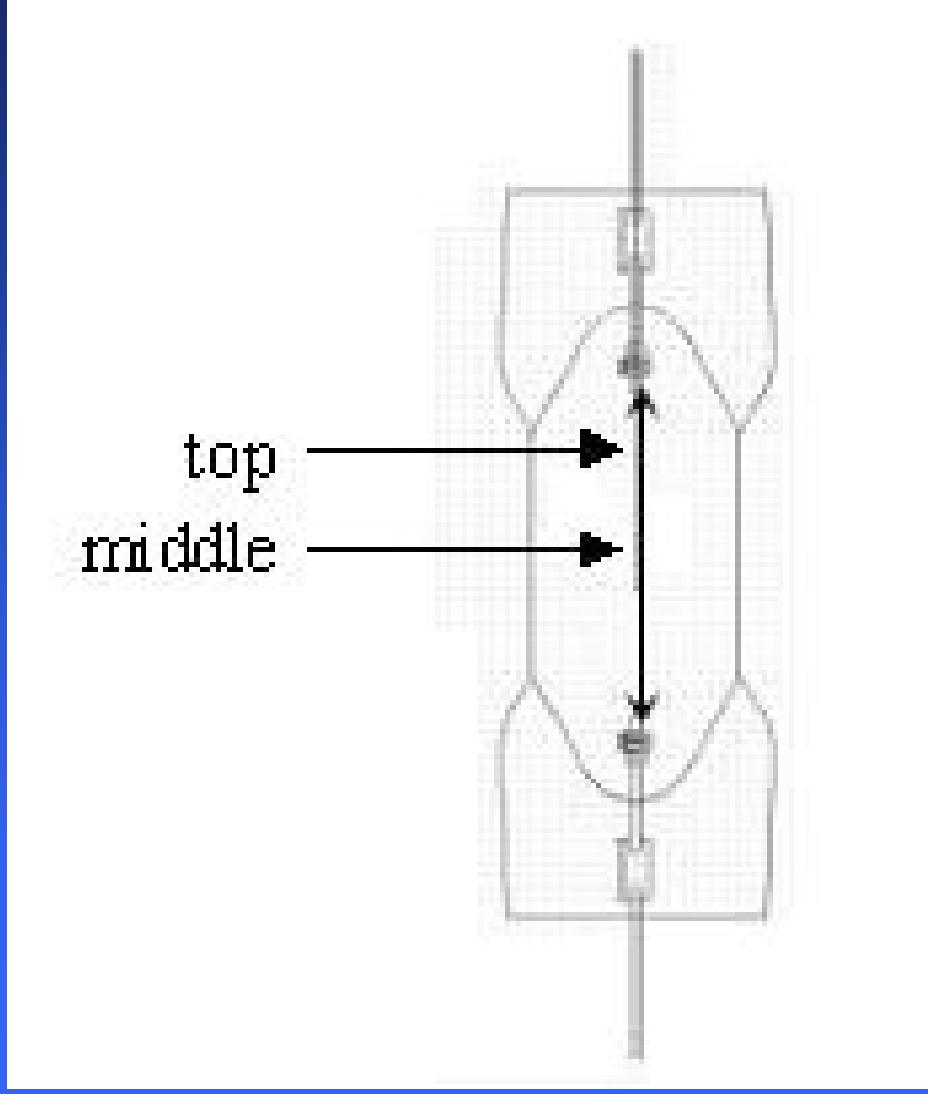
Results 07

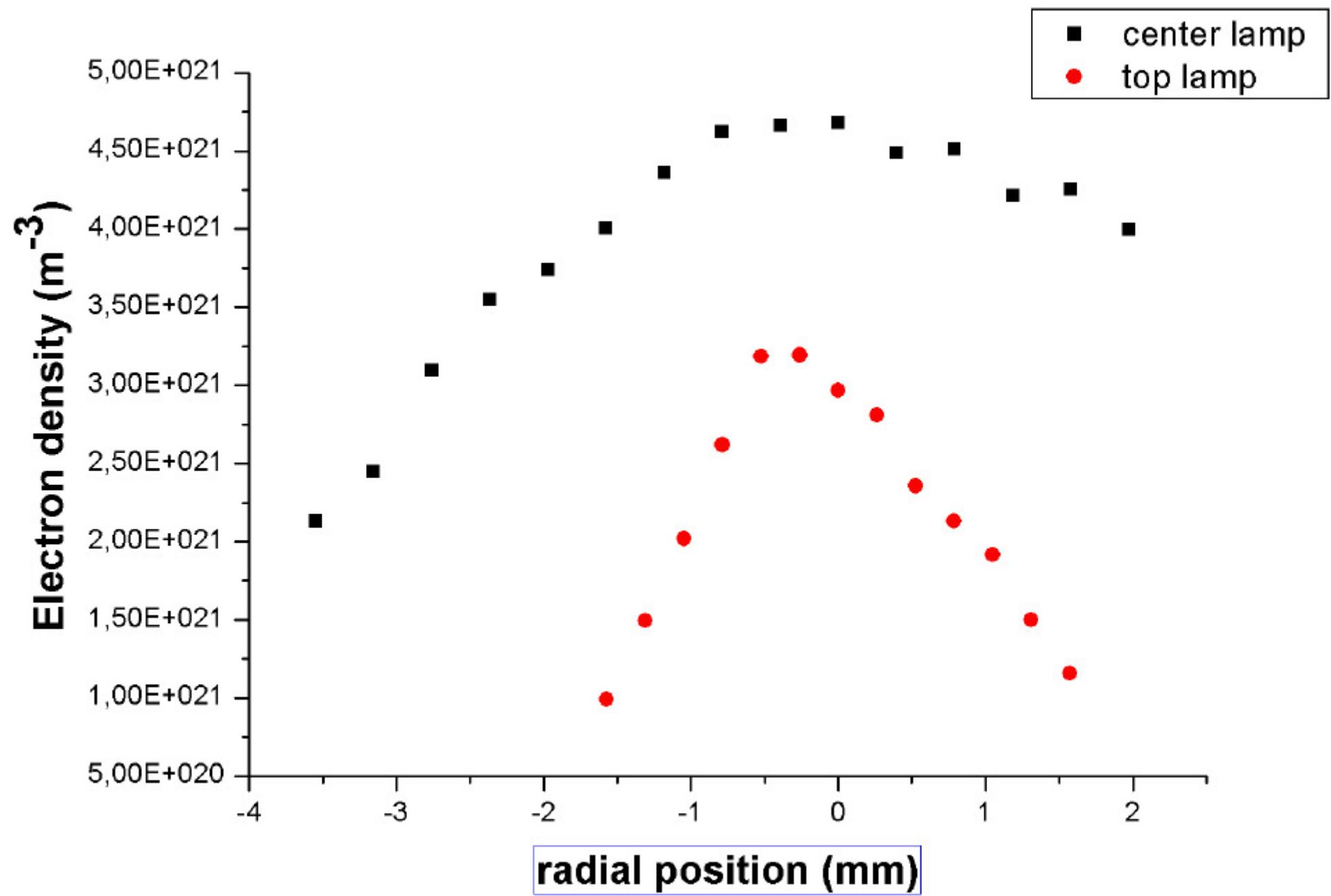
Willem-Jan van Harskamp
Nienke de Vries

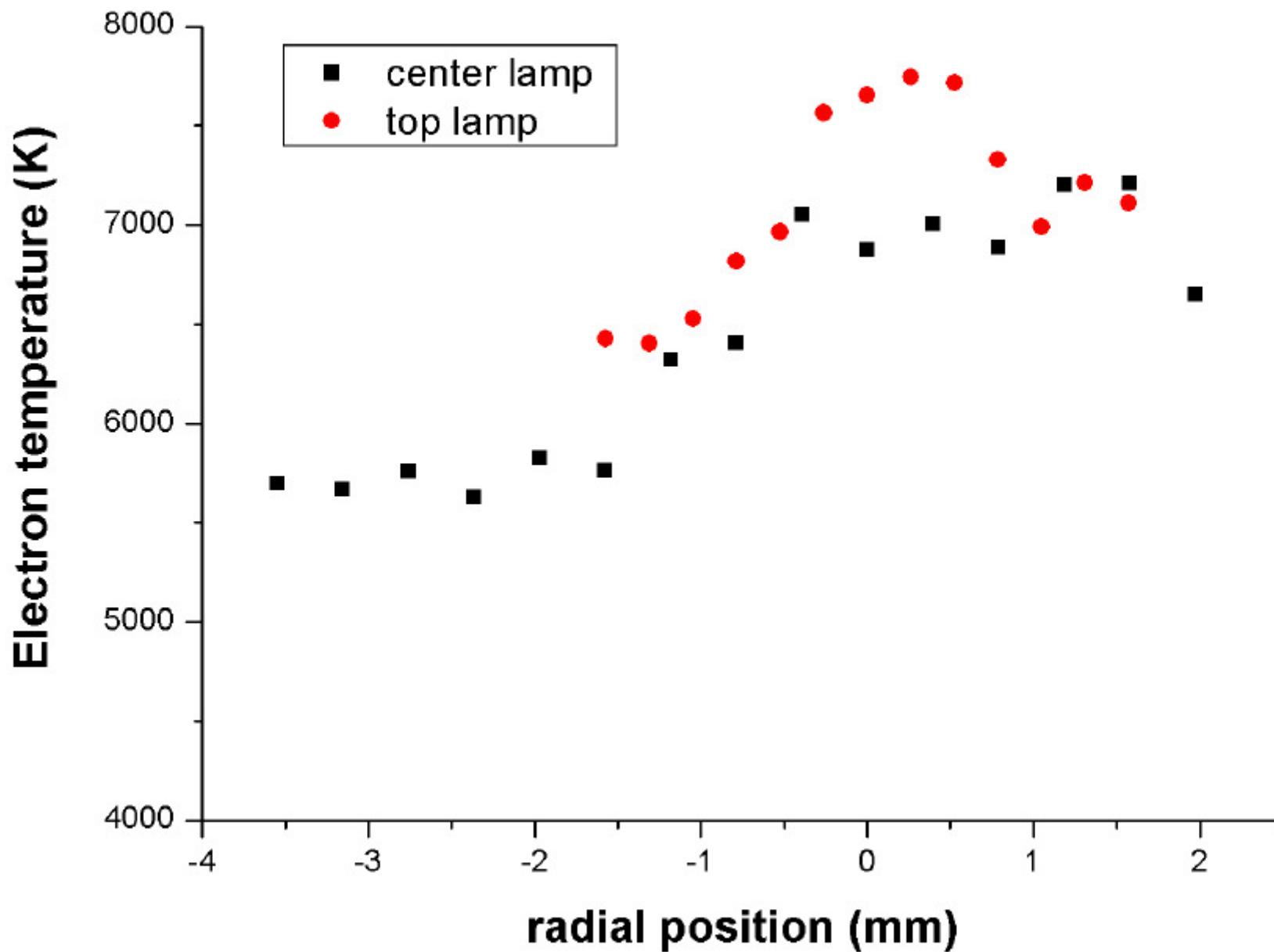
- 350 W
- 250 W
- ▲ 300 W











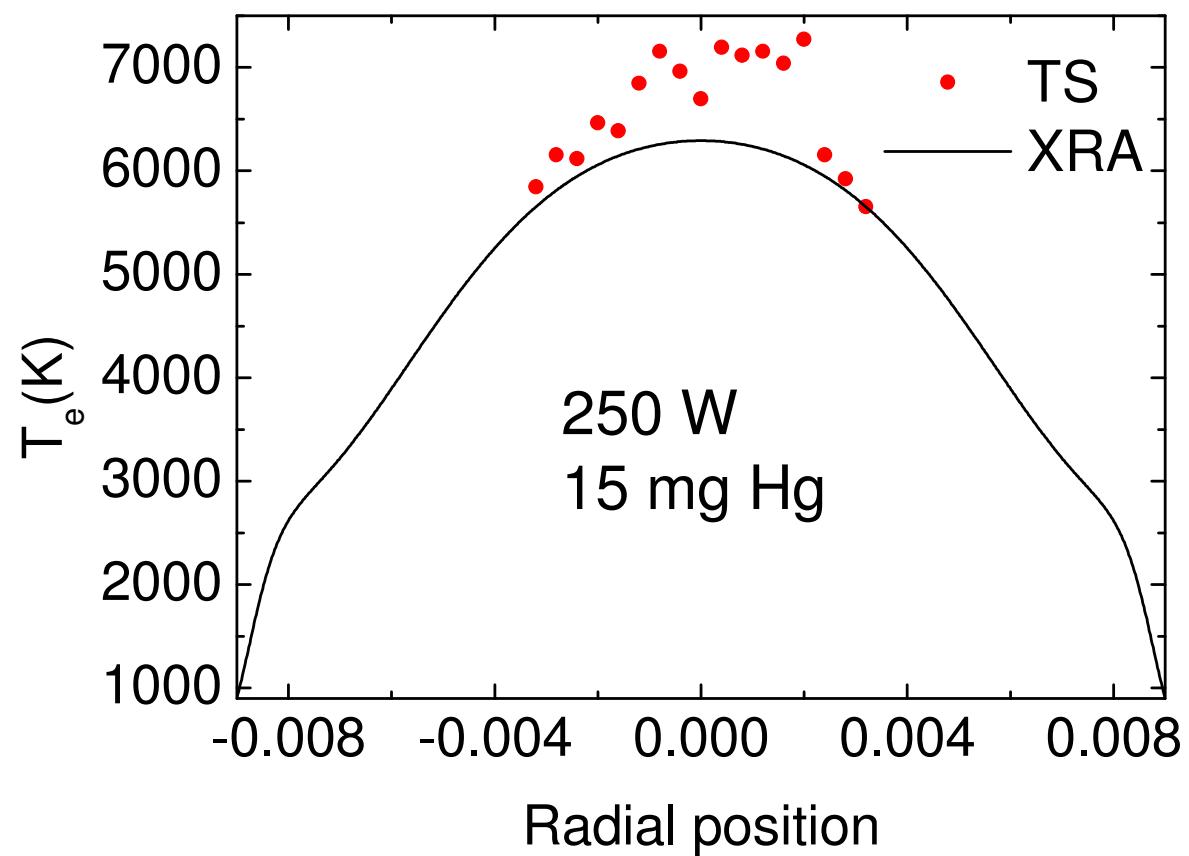
Deviation from equilibrium

Usual philosophy: high pressures
high reaction frequencies; forward/backward
low diffusion velocities

so LTE present.

Type of departure: Thermal: $T_e \neq T_g$
Chemical: ionization degree

Comparison TS and XRA



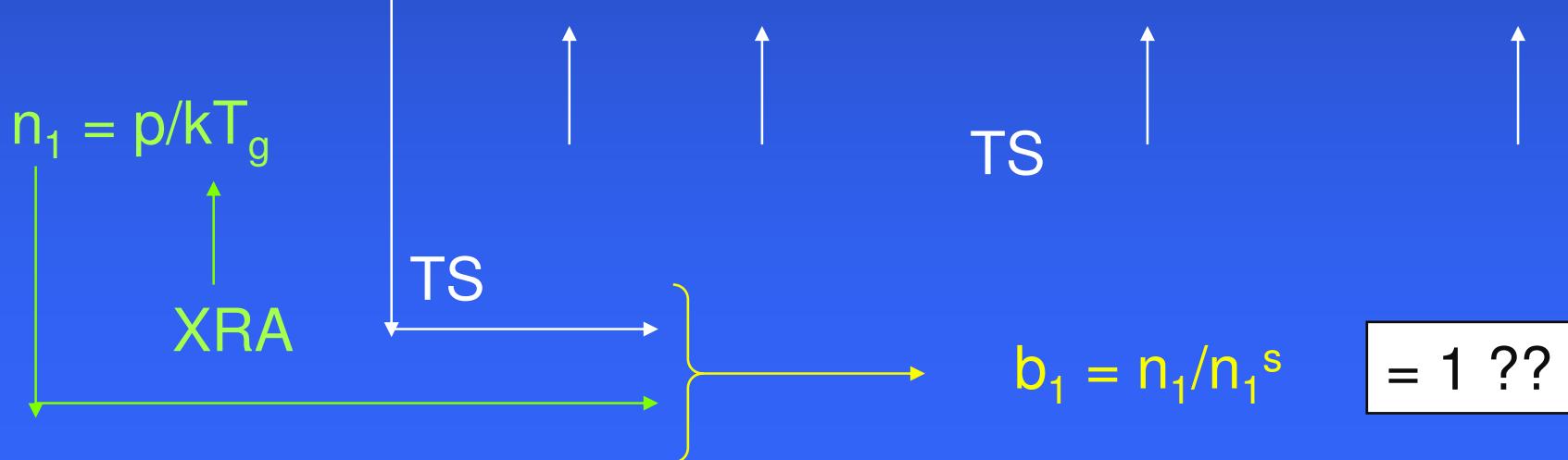
Agreement
within error bars?

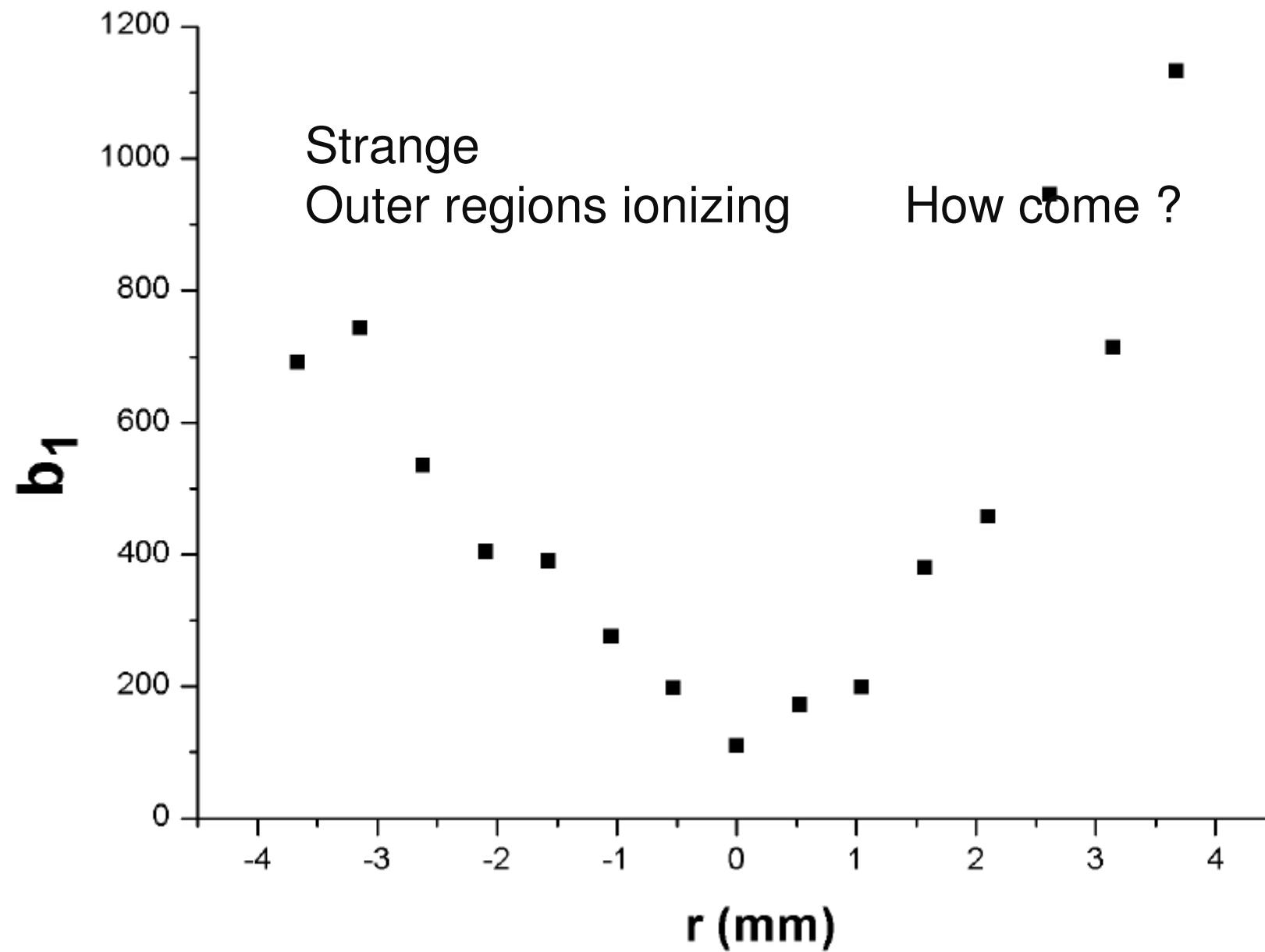
Saha equilibrium



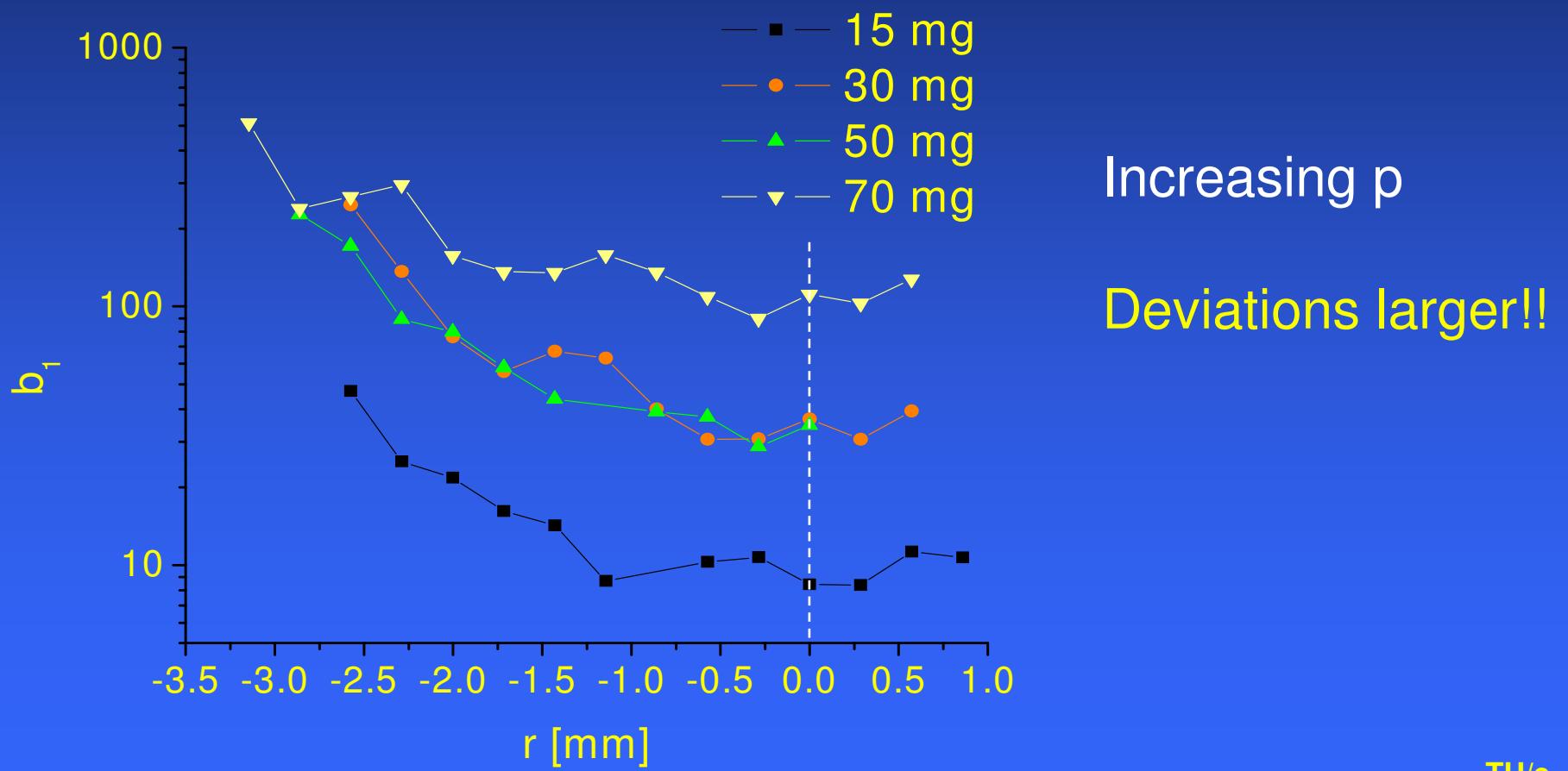
If Saha equilibrium present

$$n_1 = n_1^s = [n_e/2] [n_+/g_+] \{h^3/(2\pi m_e k T_e)^{3/2}\} \exp(I_1/kT)$$





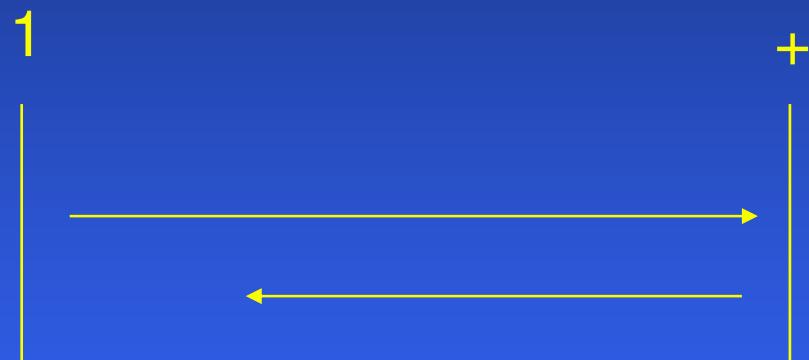
Chemical Equilibrium: b_1 -factor different gas fillings



Meaning $b_1 > 1$

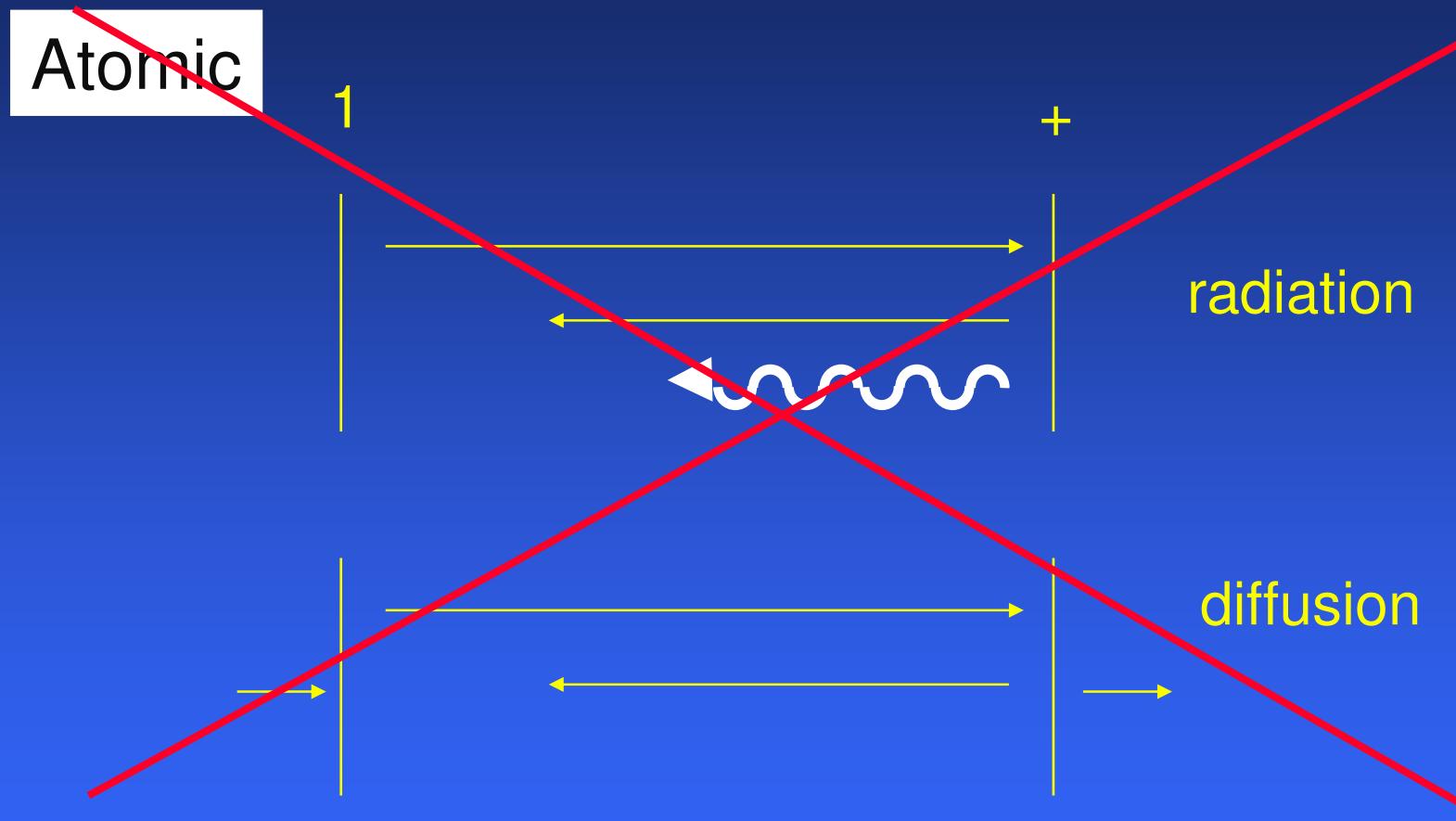
$b_1 > 1$ implies $n_1 > n_1^s \rightarrow$ ionization > 2e-recombination

Not : ionization > recombination



The ground state is relatively over-populated
The continuum is relatively under-populated

Possible Non- 2e-recombination

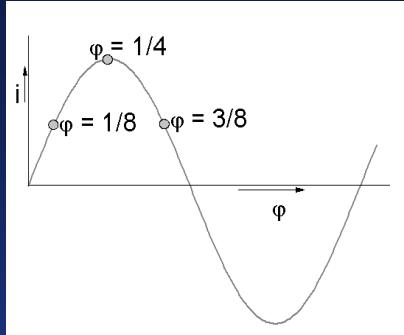


The Role of Molecules



Followed by diss recom DR

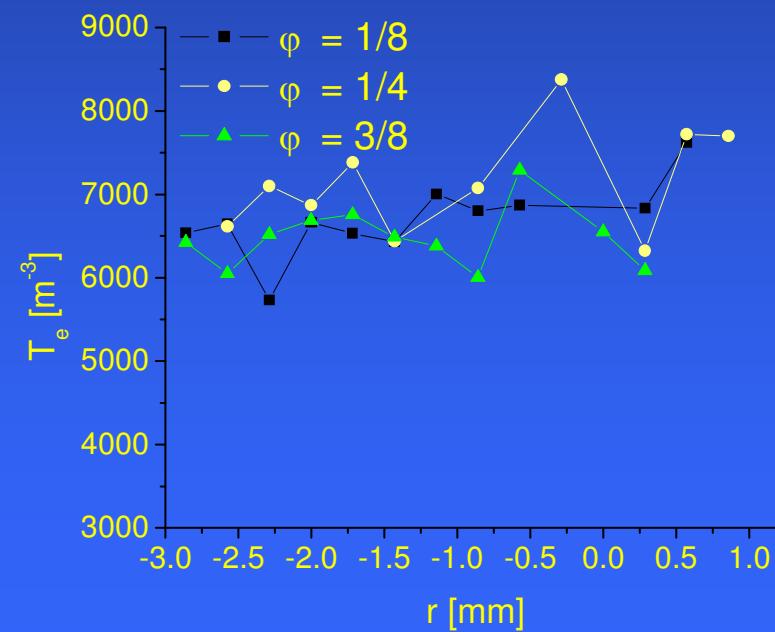
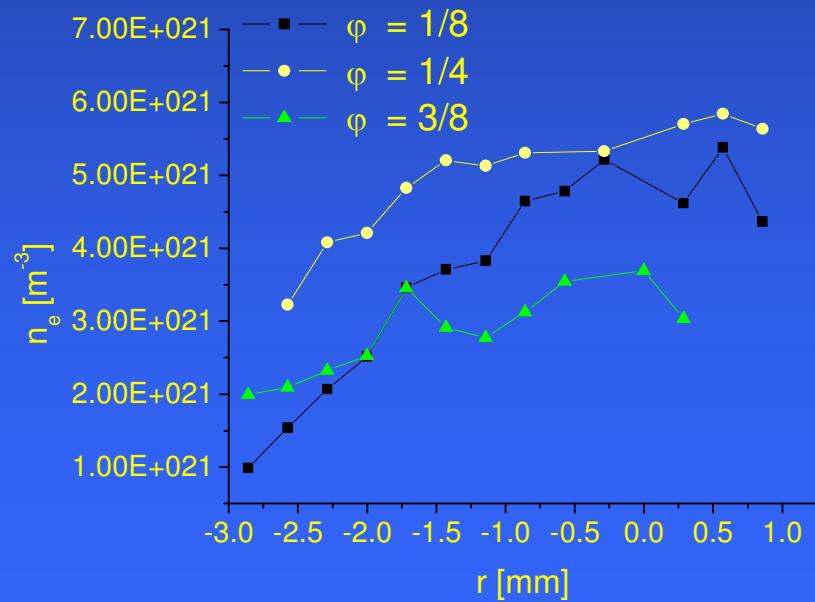




Time dependence TS

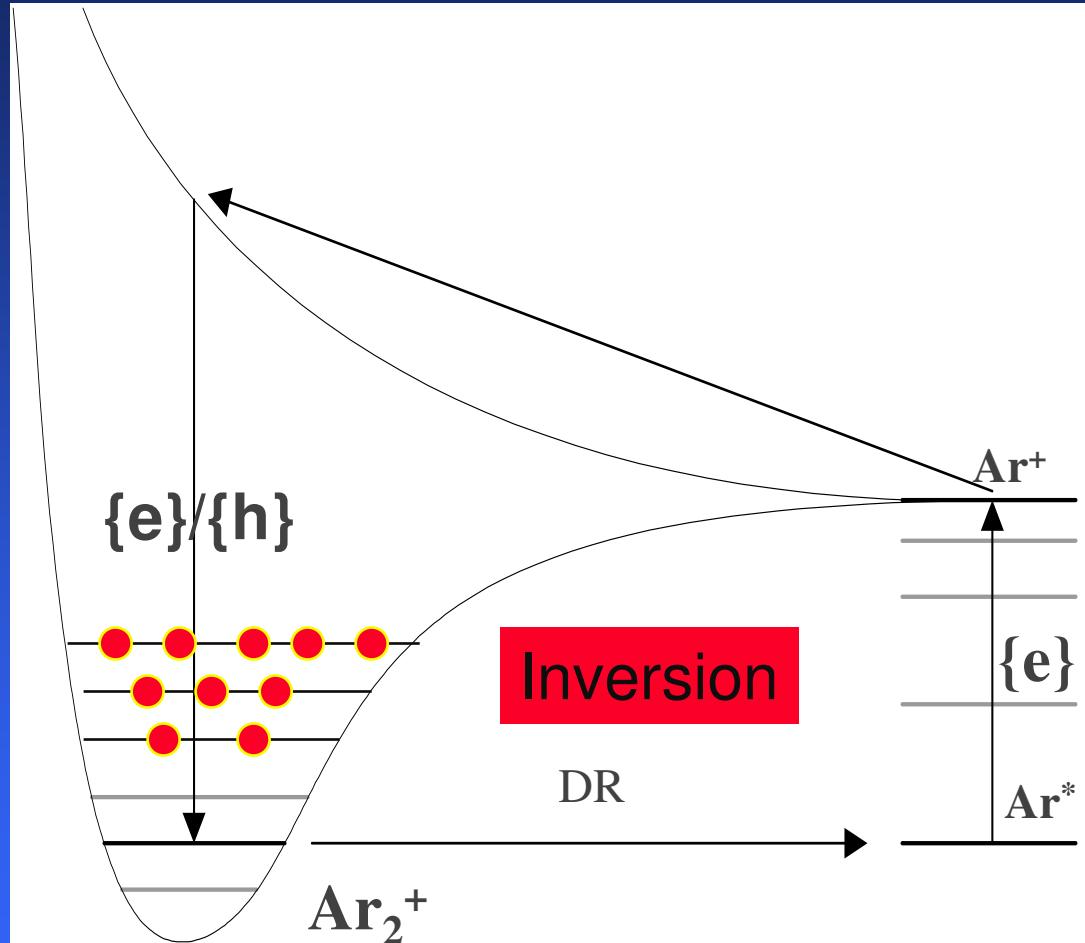
- Alternating Current: sine wave
- Radial profiles of n_e and T_e
 - different phases of the current

Variation in n_e
not in T_e



TU/e

What keeps the electrons hot??



Demand: High $\{e\}$ heating

Known
DR fast

**Solution:
Cyclic Process**

Heating during
1) Recombination
2) Super-elastics

Considerations

Although electrons {e} are primary agents
They form a minority at boundaries and afterglows

The {h} heat reservoir is much larger $3/2N kT$

During association inversion possible

{f} heat reservoir is large as well

Transport

By radiation

In space

In conversion space

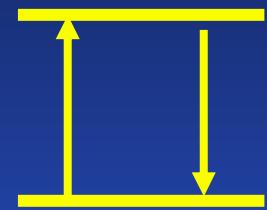
EEK

Are electrons the primary agents ??!!

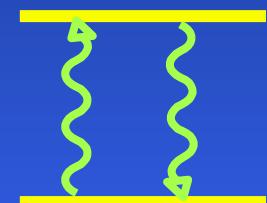
EEK electron excitation kinetics dominant !?

Competing Agents

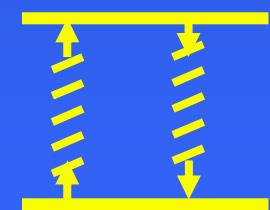
{e} EEK: Electron Excitation Kinetics



{f} REK: Radiation Excitation Kinetics

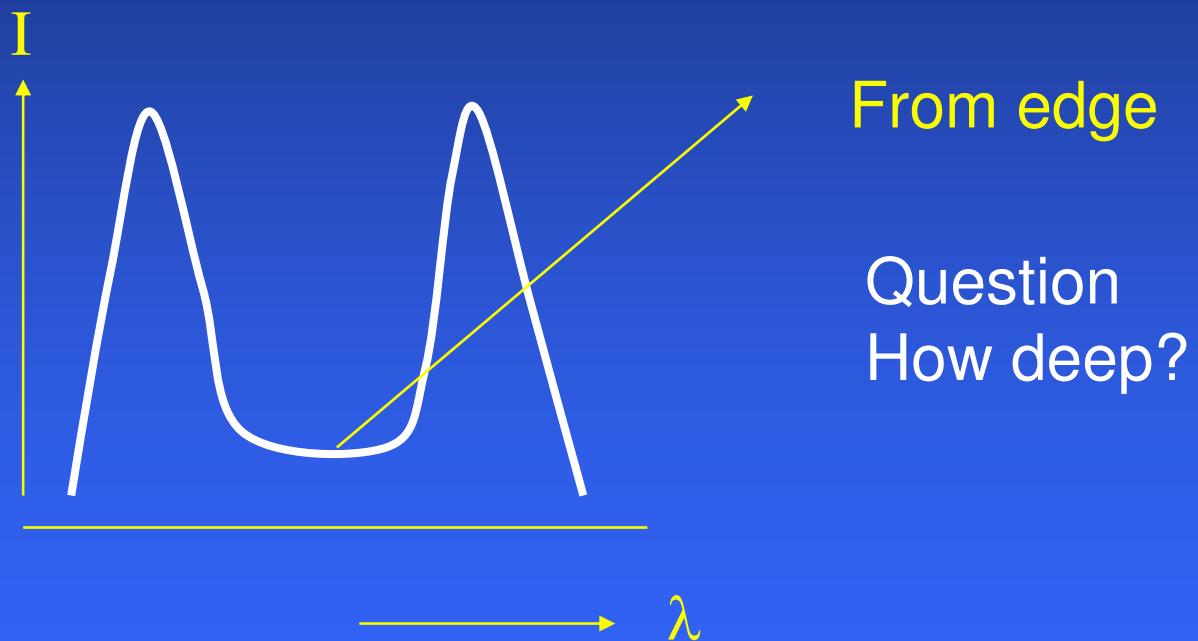


{h} HEK: Heavy particle Excitation Kinetics



Heating by radiation

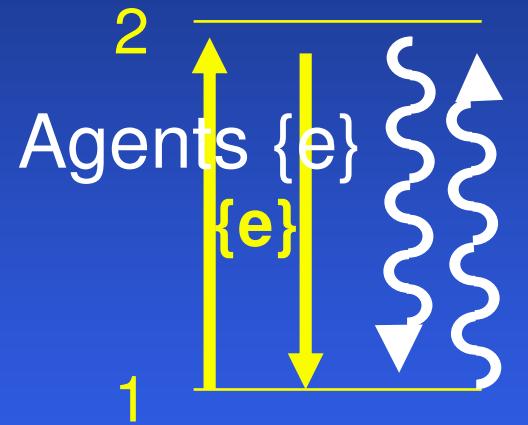
Take 254 nm line



Question
How deep?

Irradiation of the plasma edge

Plasma center



Local trapping
Planck Balance
Forced to equi

Plasma Edge

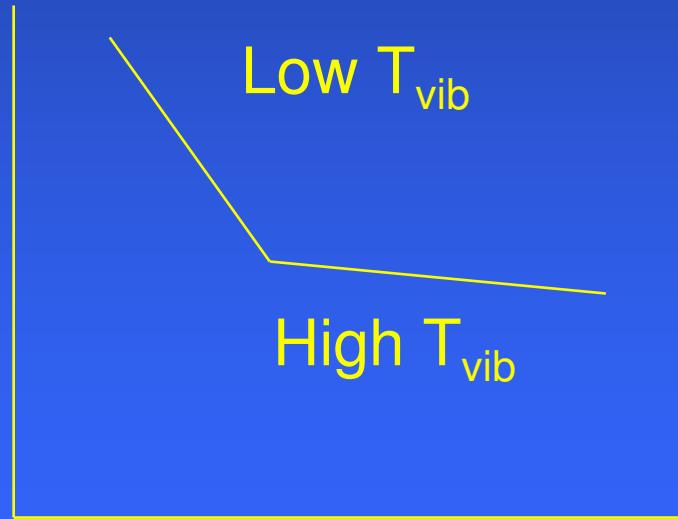


Radiation temperature
Impose
Boltzmann balance
Quasi equi

Molecular states

Known at the wall or afterglow
in time
in space

High population of highly vib states



Association via
Higher vibrational states

The Sulfur lamp

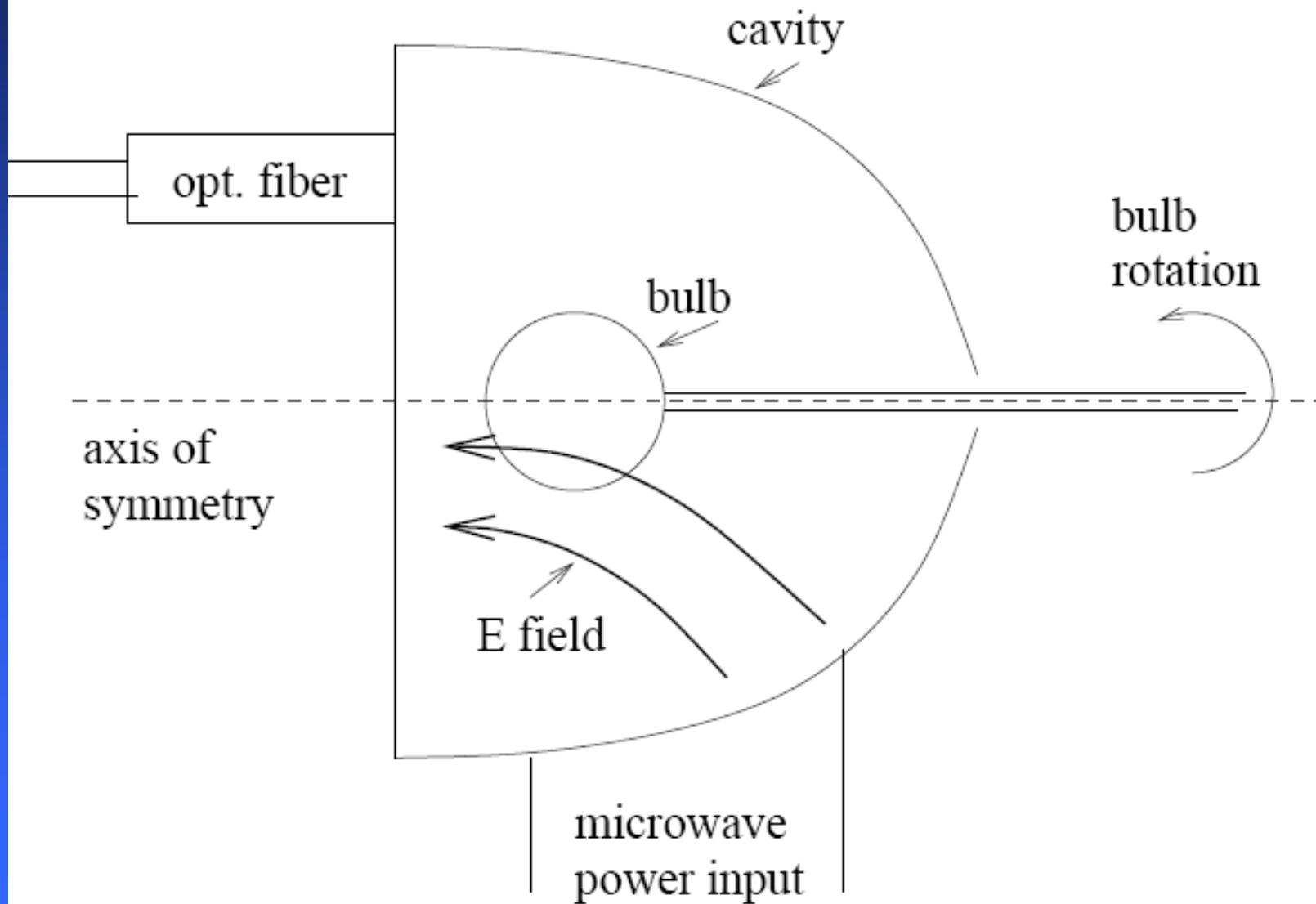
Harm van der Heijden 2003
Colin Johnston 2003

A microwave plasma $p = 10 \text{ bar}$
 $P = 1\text{k Watt}$

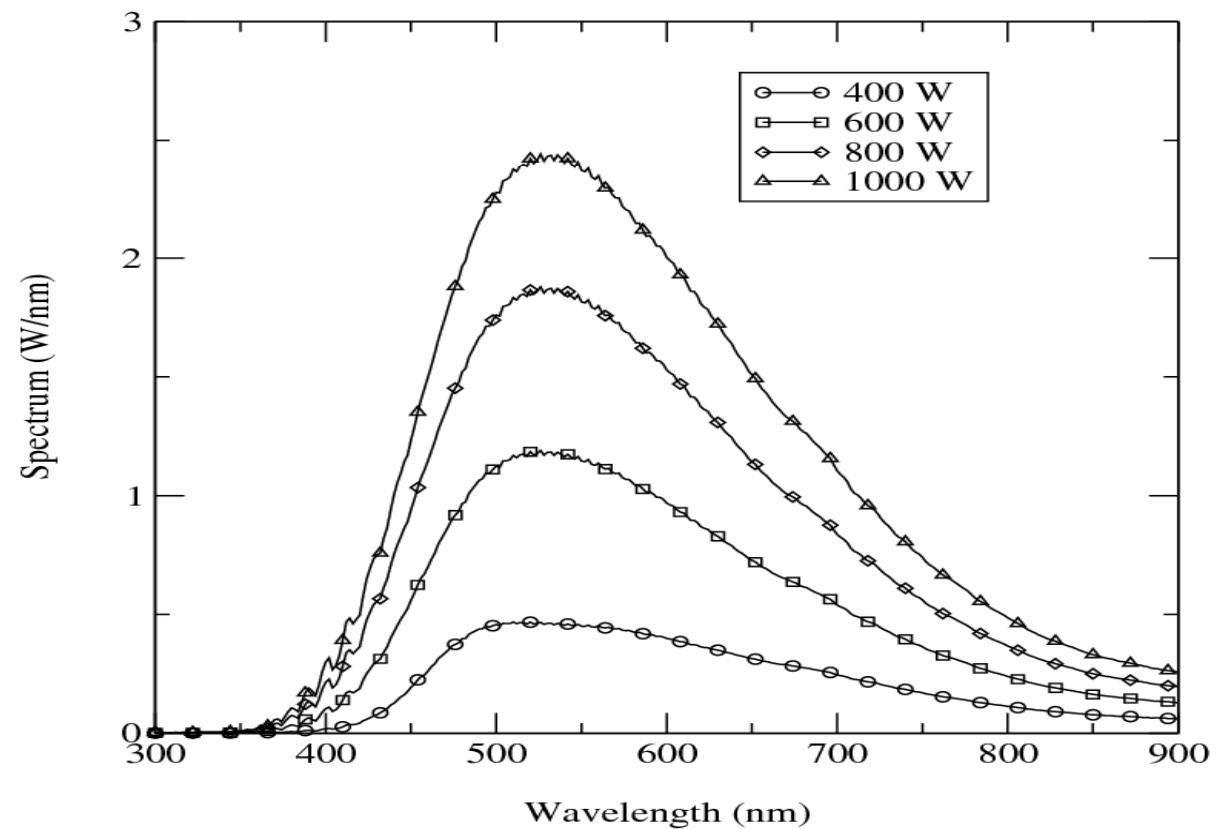
Experimental: not (laser) accessible
 only passive spectroscopy

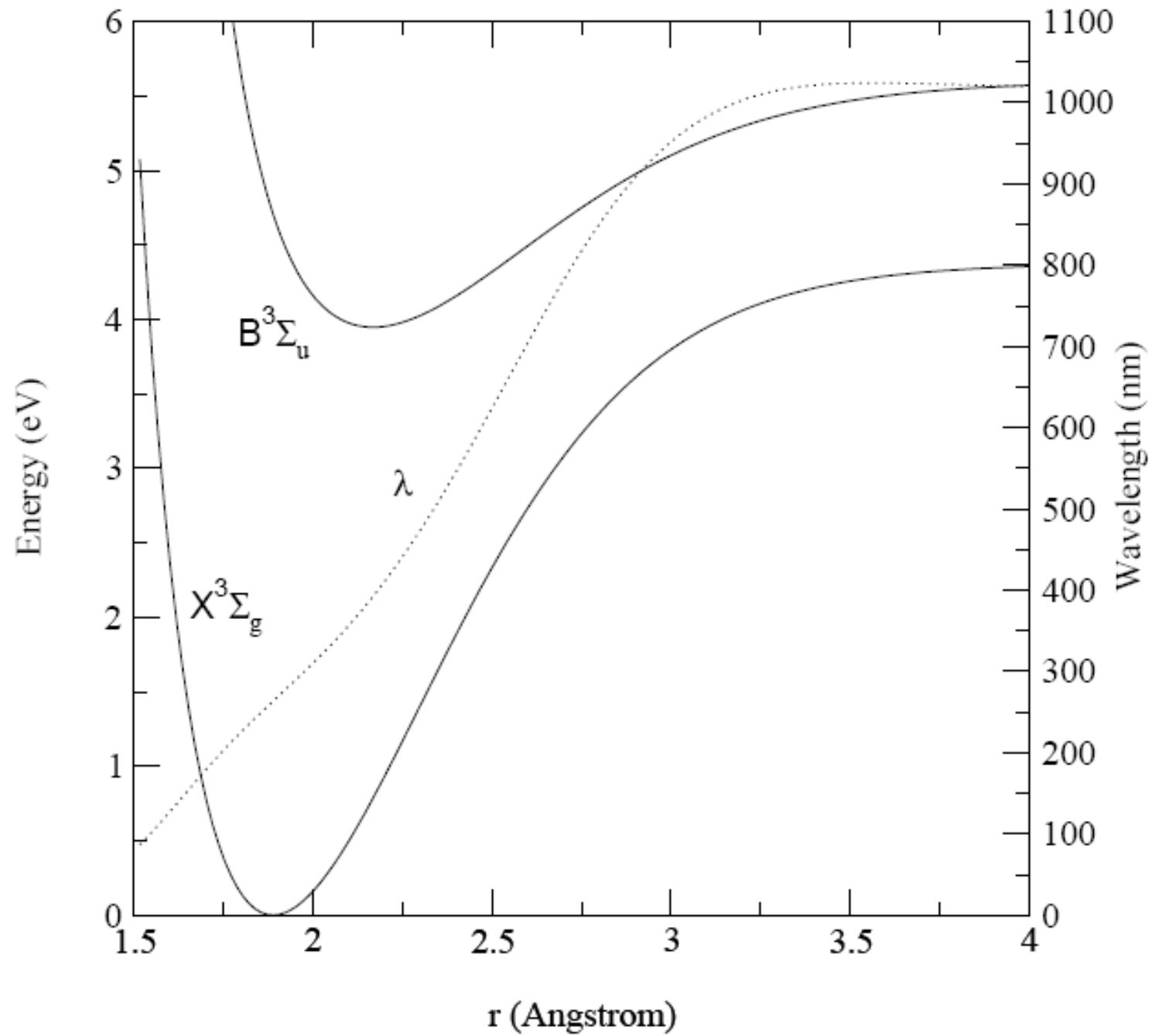
Modeling: LTE microwave power coupling
 heat conductivity
 radiation generation and transport

In 1992 a candidate for the illumination Sidney 2000

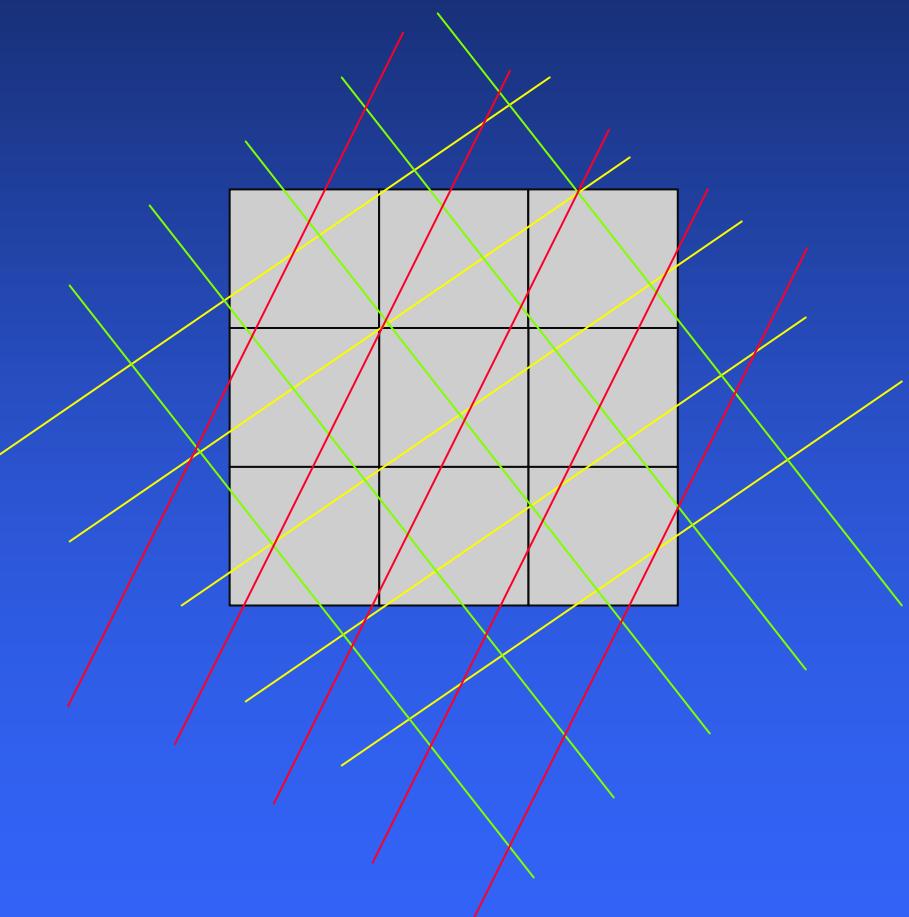


Matching the eye-sensitivity





Radiation: Ray Tracing

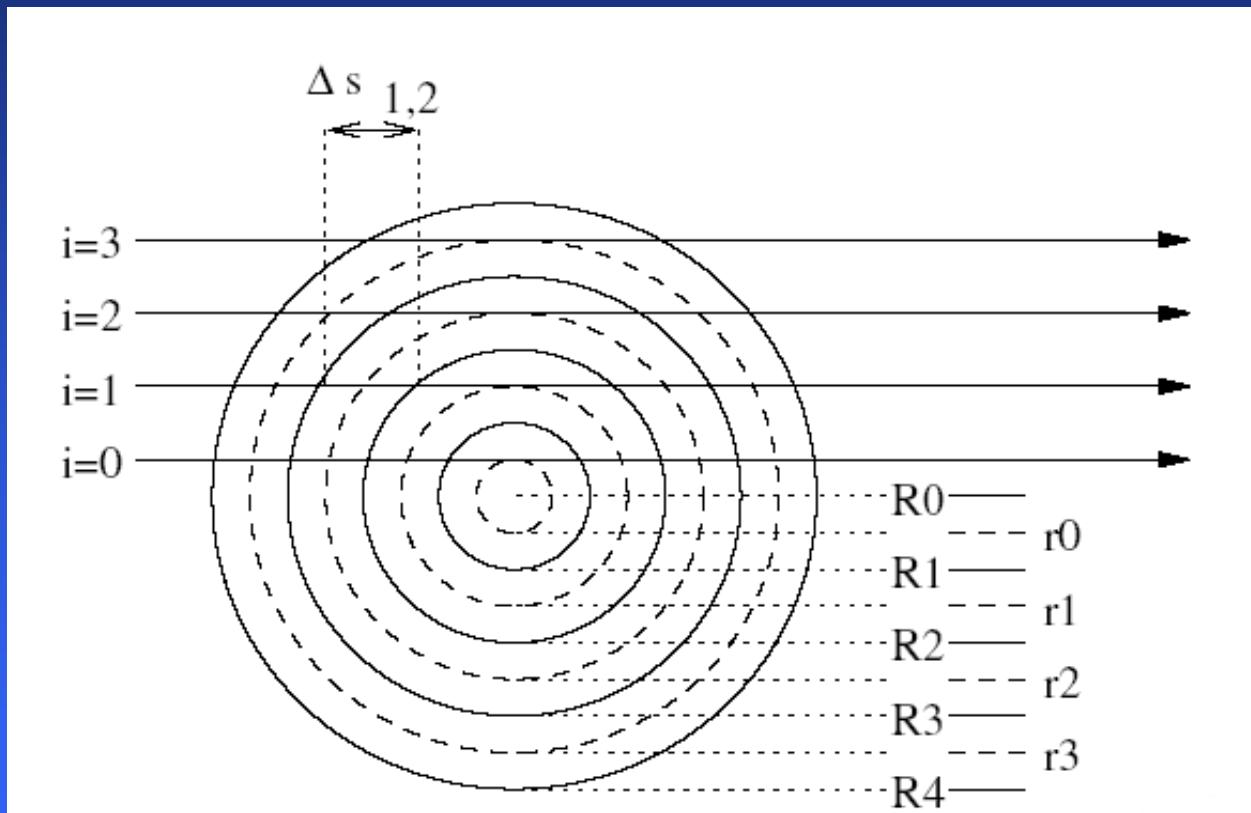


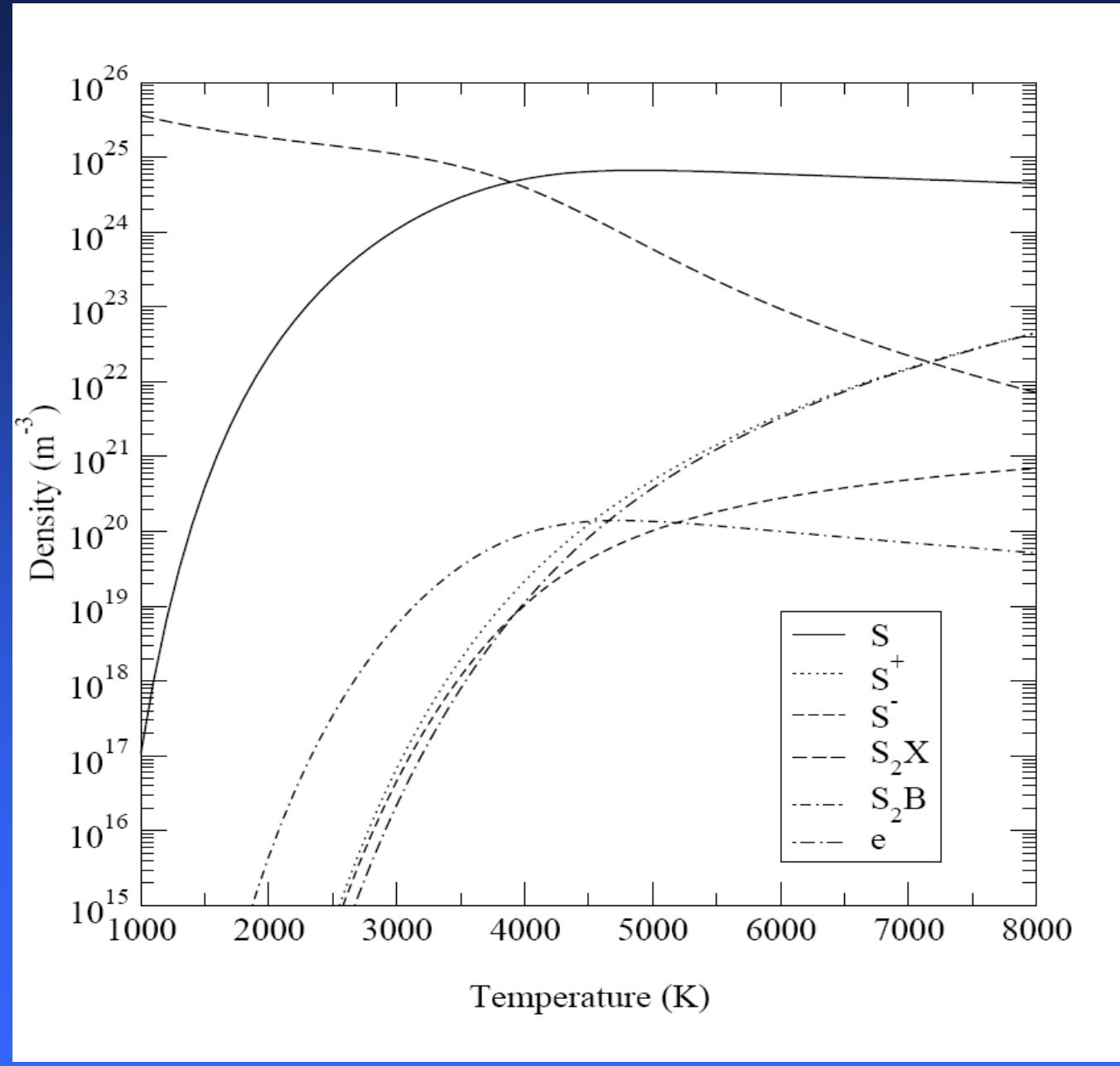
Evolve

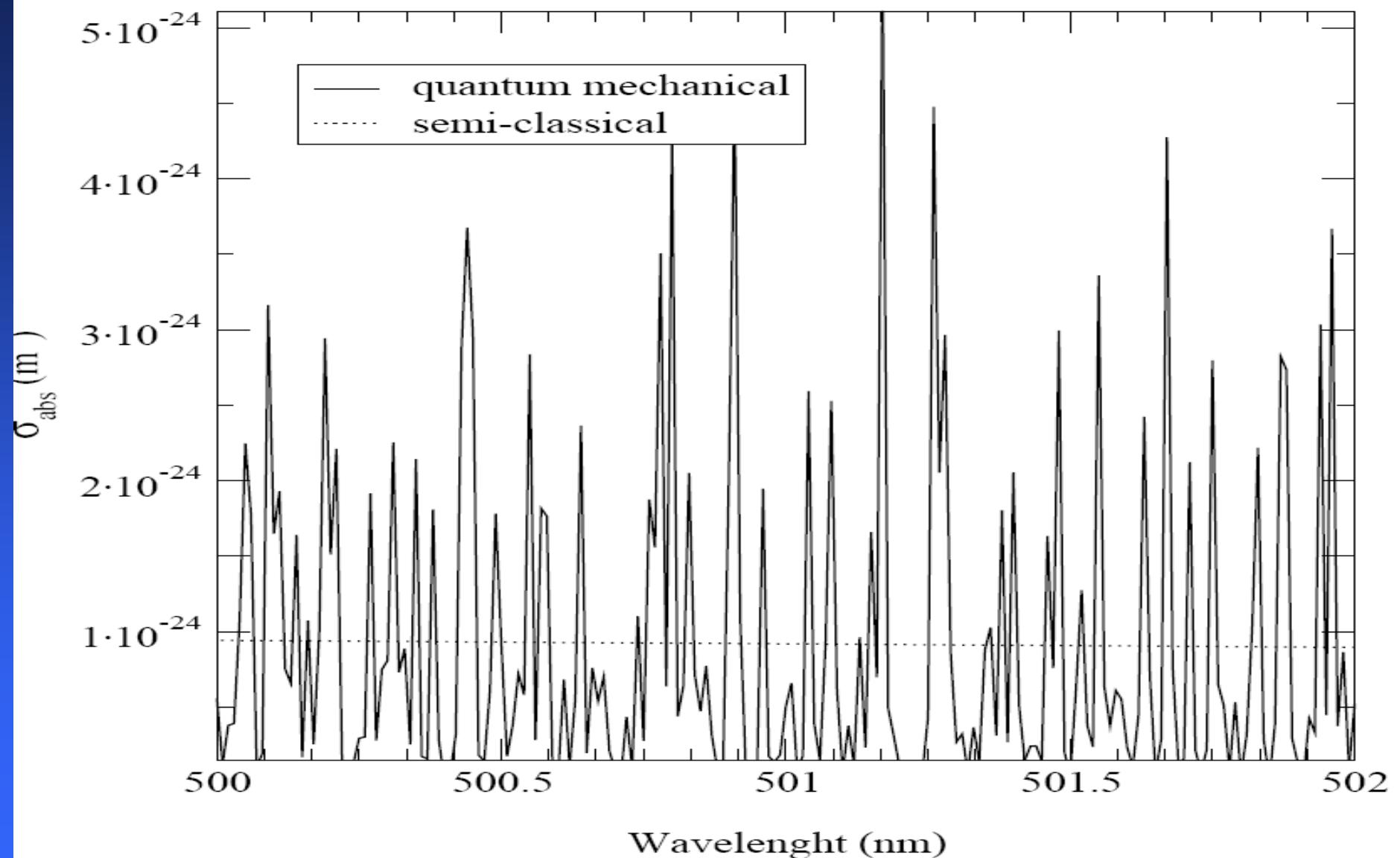
$$\frac{dI_v(v)/ds}{ds} = j_v - k(v)I_v(v)$$

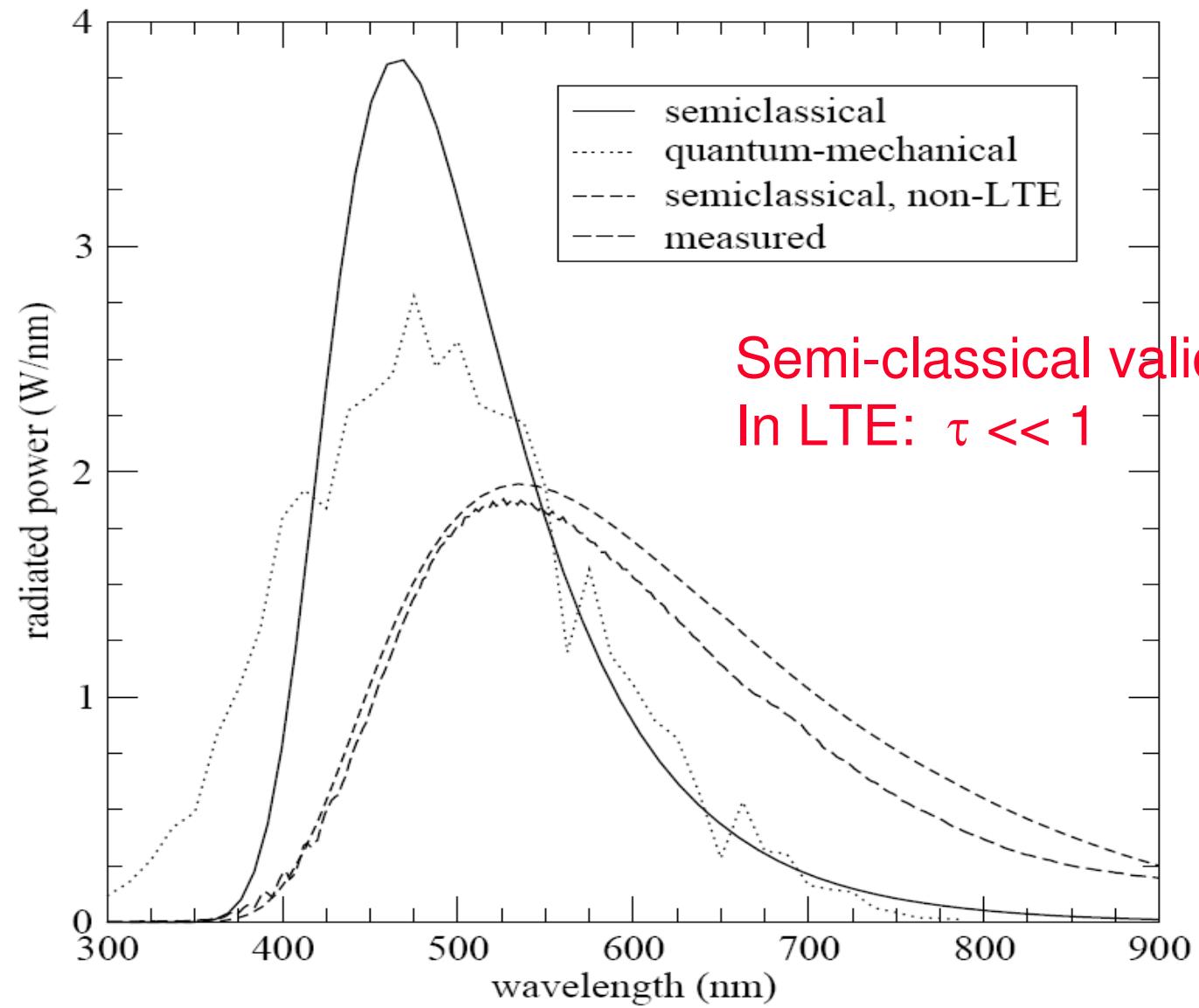
- For different lines
- Different v -values
- Compute plasma irradiation
- Solve Fluid equation
- Find new $k(v)$ and $j_v(v)$
- Evolve Radtrans eqn again
- Etc.

Using Symmetry









Main result

The IR excess can (only) be explained
By a non-equi distribution of the rot-vib population

B-Molecules are formed in higher states
During coll-decay to lower state radiation takes place

Inversion in the B molecule !!

Evidences for high-pressure non-LTE

CO_2 and CO lasers.

Sulfur lamp

Ball-lightening

Pink Afterglow N_2

Concluding

High pressure does not guarantee Equilibrium.

Electrons principal agents in center/initiation
where they create radiation
and heavy internal states:

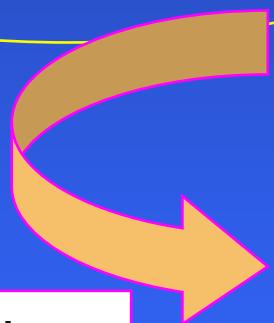
Side/after/behind glow ruled in many case by
heavies
photons

Thomson scattering difficult but of high value

Outlook

T P P
Thermal
Plasma
Processes

Equilibrium departure
a challenge



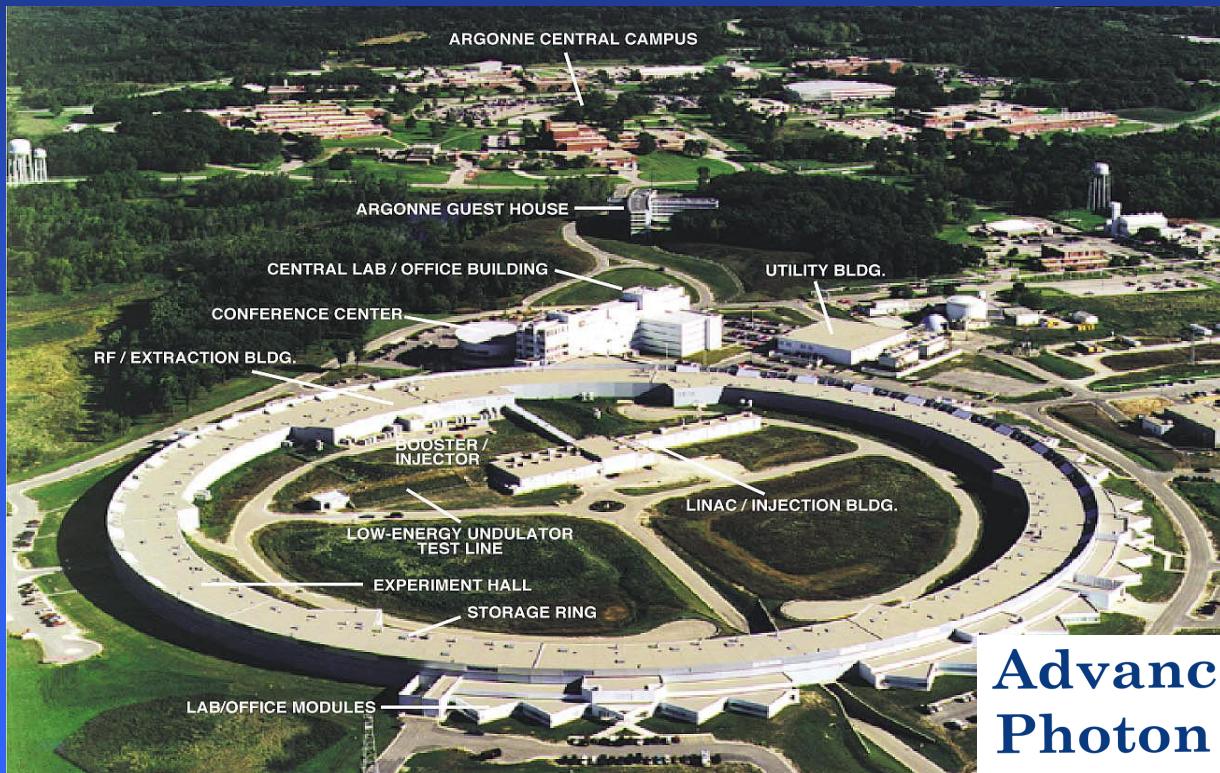
H T P P
High
Tech
Plasma
Processes

Thank you for your attention



XRF

Tanya Nimalasuriya (TU/e)
Evert Ridderhof (TU/e)
John J. Curry (NIST)
Craig J. Sansonetti (NIST)
Sharvjit Shastri (APS)

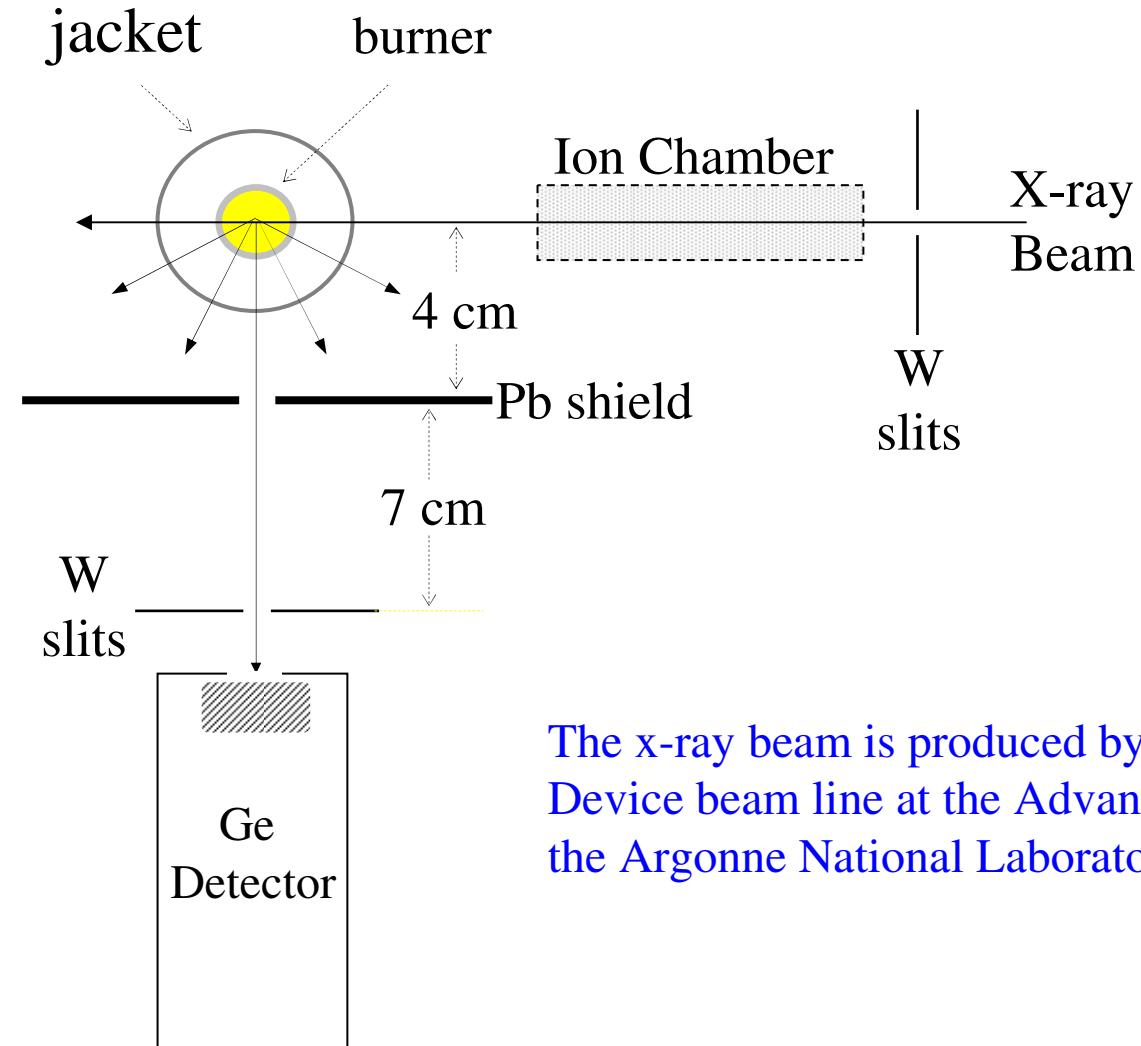


**Advanced
Photon
Source**

ARGONNE NATIONAL LABORATORY



XRF sketch

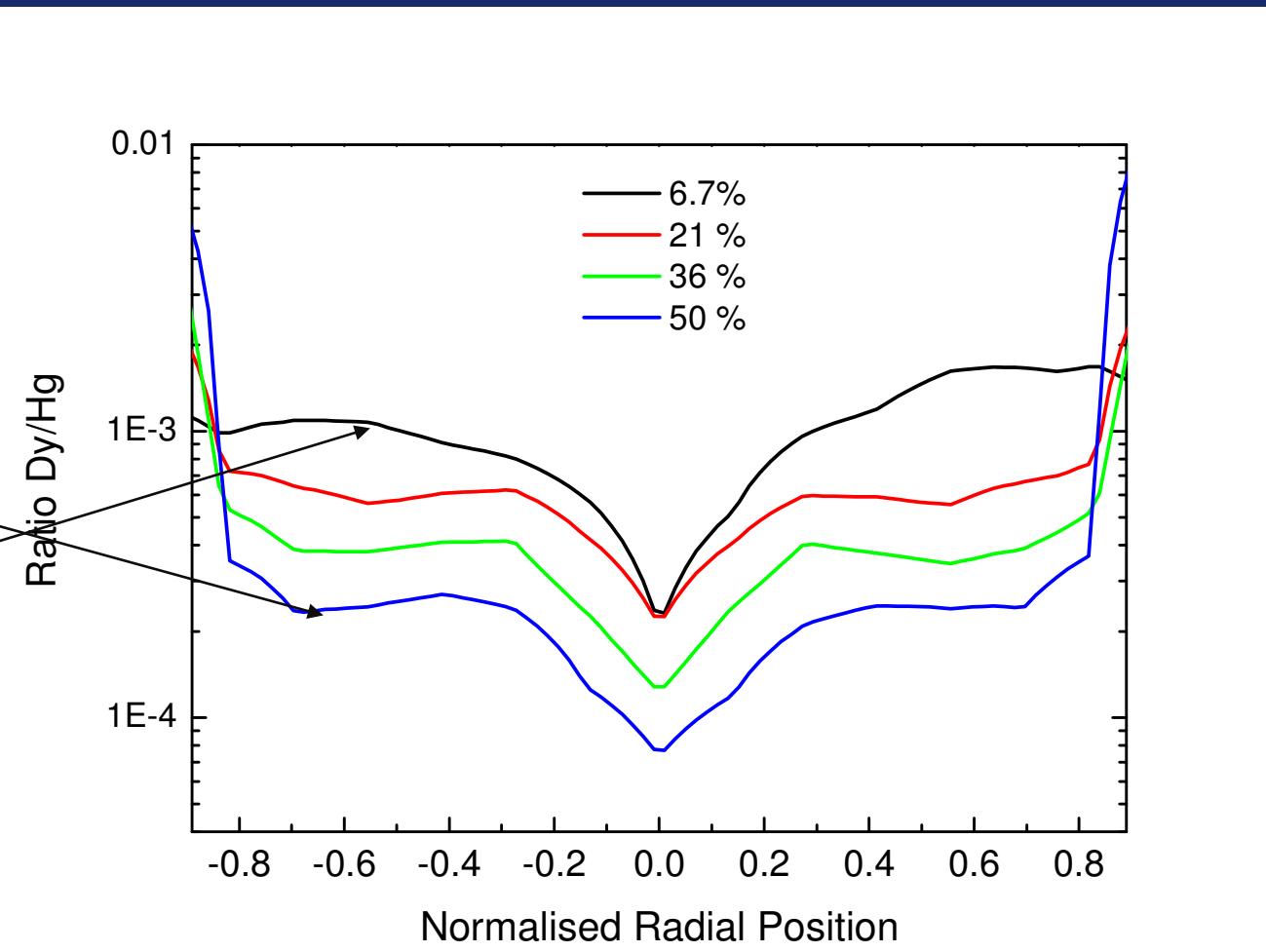
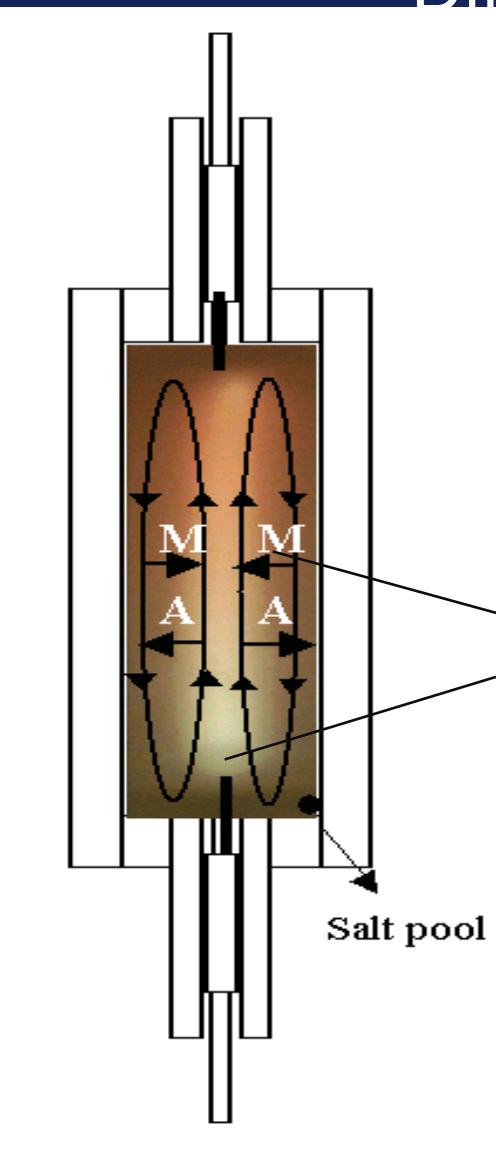


The x-ray beam is produced by the Sector 1 Insertion Device beam line at the Advanced Photon Source at the Argonne National Laboratory

XRF advantages

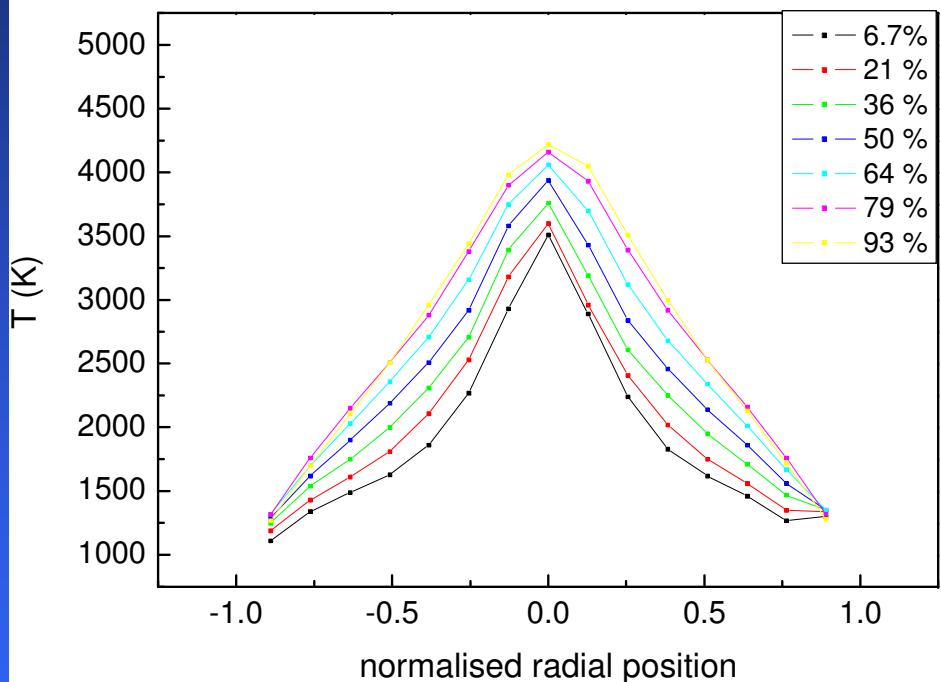
- X-ray induced fluorescence:
 - determines **elemental densities** of Dy,Hg
 - is effective anywhere in the burner
- No inversion technique is needed
- T profile with Hg densities

Diffusion versus (radial) convection

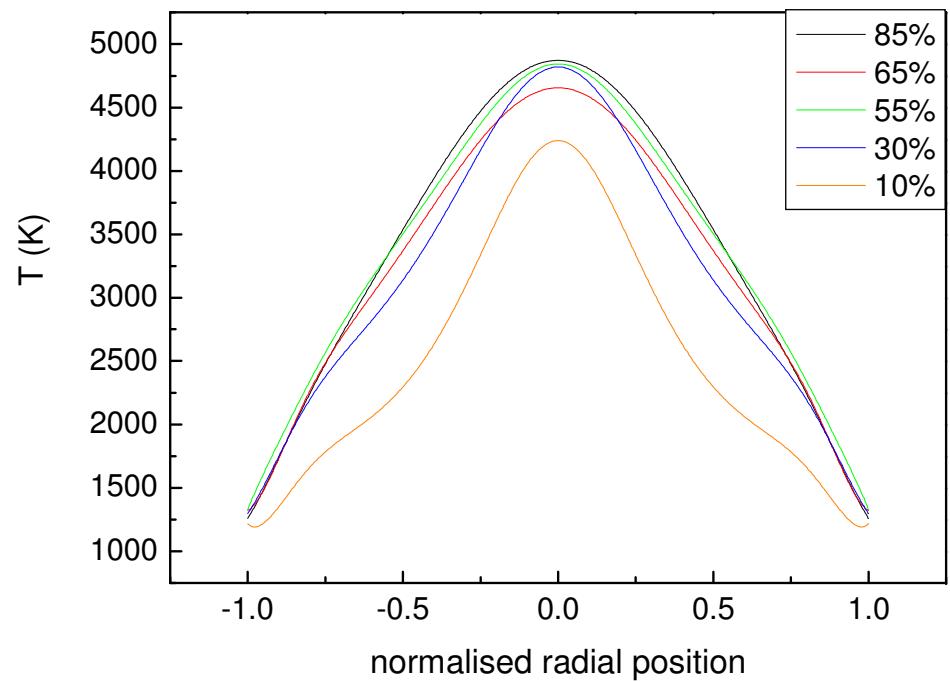


Temperature profile from Hg density

$T_{XRF} \approx T_{XRA}$ but low !!

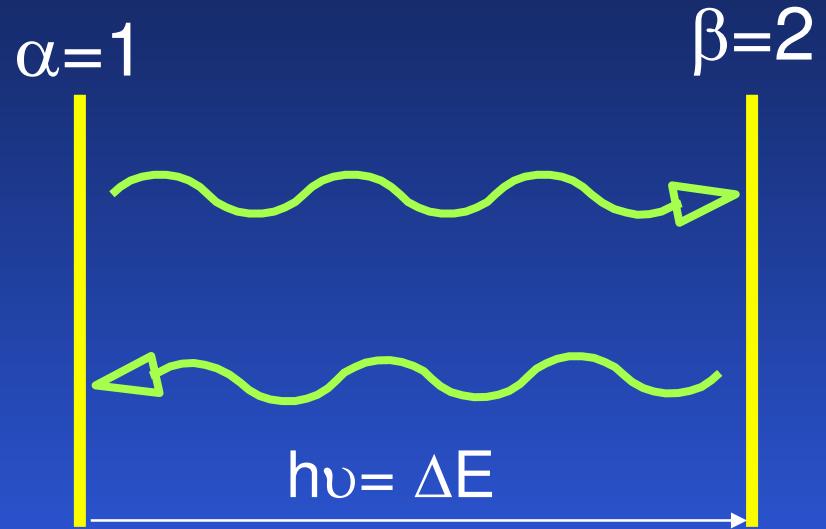


XRF 140 W



XRA 142 W, X.Y. Zhu

Example pLPE



Intense laser irradiates transition:

Proper balance Absorption St.Emission

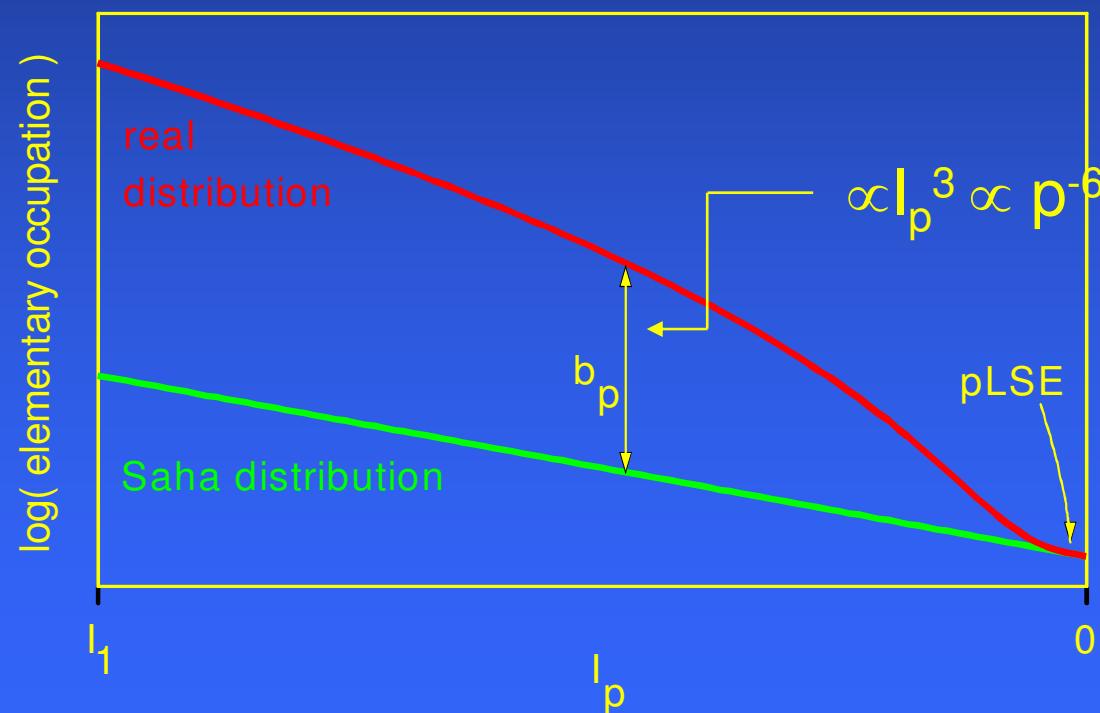
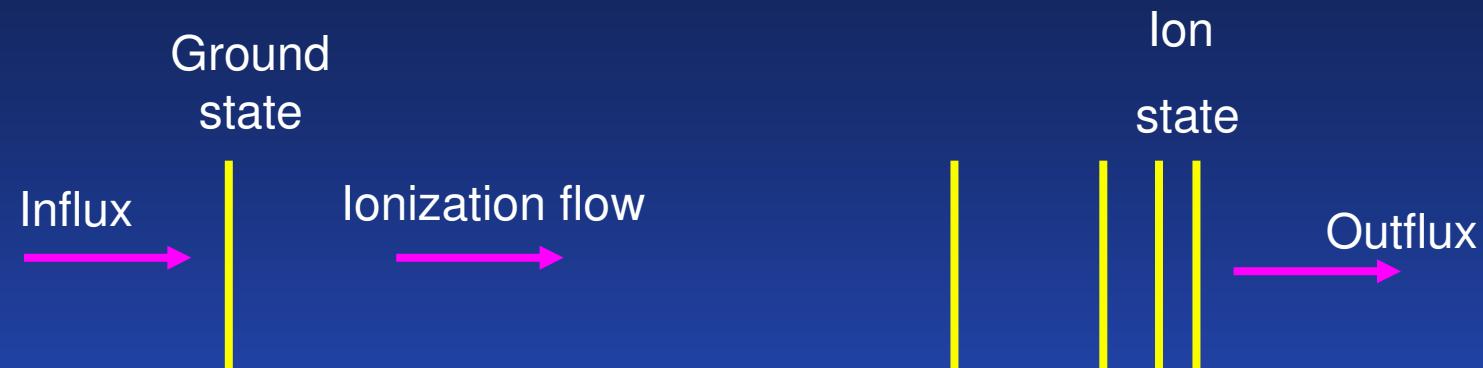
Look for comparable TE situation

$T \rightarrow \infty: \exp{-\Delta E/kT} = 1 \rightarrow$

$$\eta(1) = \eta(2)$$

TU/e

Ion Efflux Effecting the ASDF



Example pLSE



Approaching continuum:

Equi. restoration rates increase

Look for comparable TE situation

Saha equation ruled by electrons from continuum

$$\eta^s(p) = \eta_+ \eta_e \mathcal{V}_e \exp(I_p/kT_e)$$

The efficiency

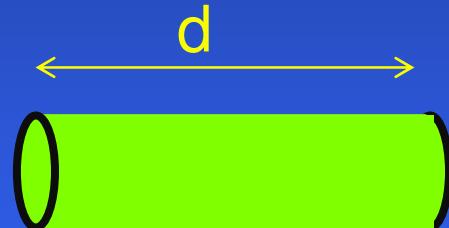
How many photons detected ?

$$\xi = F_{\text{scat}} \cdot F_{\text{coll}} \cdot F_{\text{det}}$$

Combined
Absorption
Emission

F_{scat} the scatter-fraction

$$\tau = kd = n_e \sigma d$$

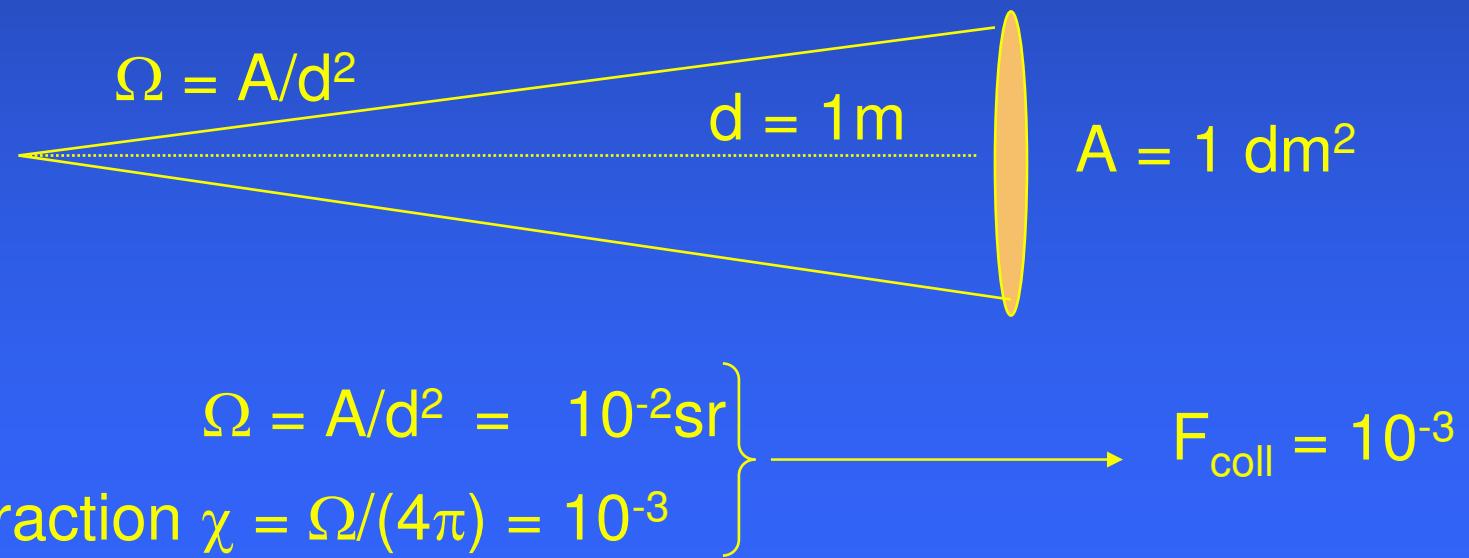


Take :

$$\left. \begin{array}{l} \sigma = 6.65 \cdot 10^{-29} \text{ m}^2 \\ n_e = 10^{20} \text{ m}^{-3} \\ L = 1 \text{ cm} = 10^{-2} \end{array} \right\} \rightarrow \tau = 10^{-10}$$

The collected fraction

The collected fraction: possible 1 dm² lens at 1 m



The detected fraction

The detection fraction $F_{\text{det}} = 10^{-2}$

$$\xi = F_{\text{scat}} \cdot F_{\text{coll}} \cdot F_{\text{det}} = 10^{-10} \cdot 10^{-3} \cdot 10^{-2} = 10^{-15} !!!$$

Laser needed

e.g. 1 mJ $\approx 5 \cdot 10^{15}$ photons.

Several TS competitors

Collectivity
Rayleigh scattering
False stray light, vessels
Plasma photons
Laser produced plasma

The competitors

Collectivity

Rayleigh scattering
False stray light, vessels

Plasma photons
Laser produced plasma

Change of Gaussian

At λ_0
At λ_0

At λ_0 and in $\Delta \lambda$ - TS
A different plasma

The competitors

Collectivity

Rayleigh scattering
False stray light, vessels

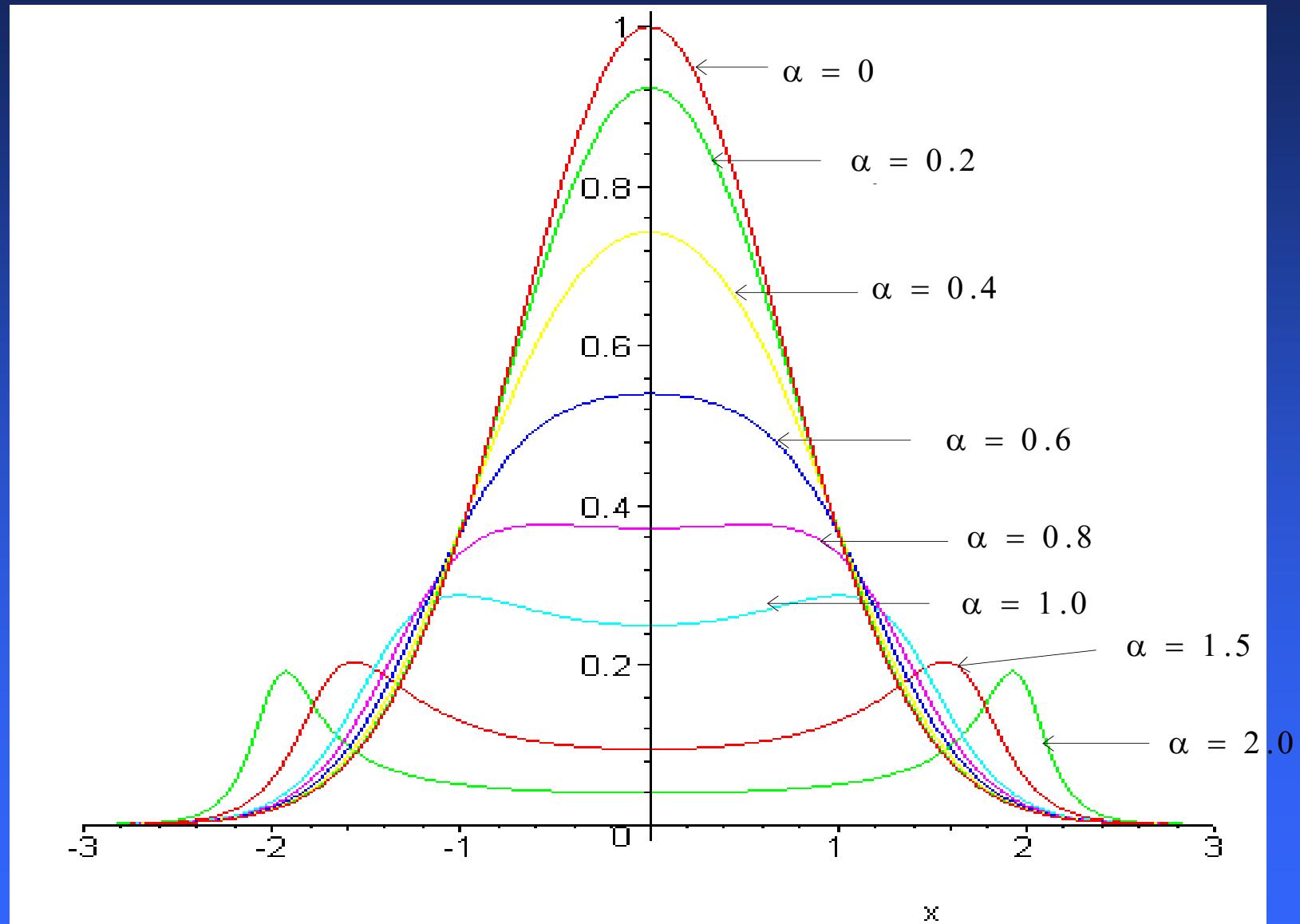
Plasma photons
Laser produced plasma

Change of Gaussian

At λ_0
At λ_0

At λ_0 and in $\Delta \lambda$ - TS
A different plasma

TS spectra shapes for collective behavior



The competitors

Collectivity

Rayleigh scattering
False stray light, vessels

Plasma photons
Laser produced plasma

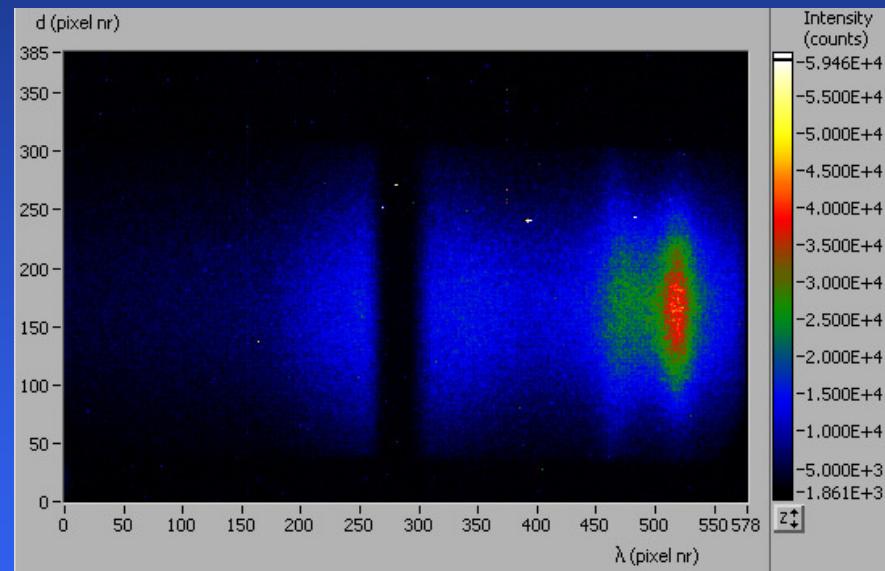
Change of Gaussian

At λ_0
At λ_0

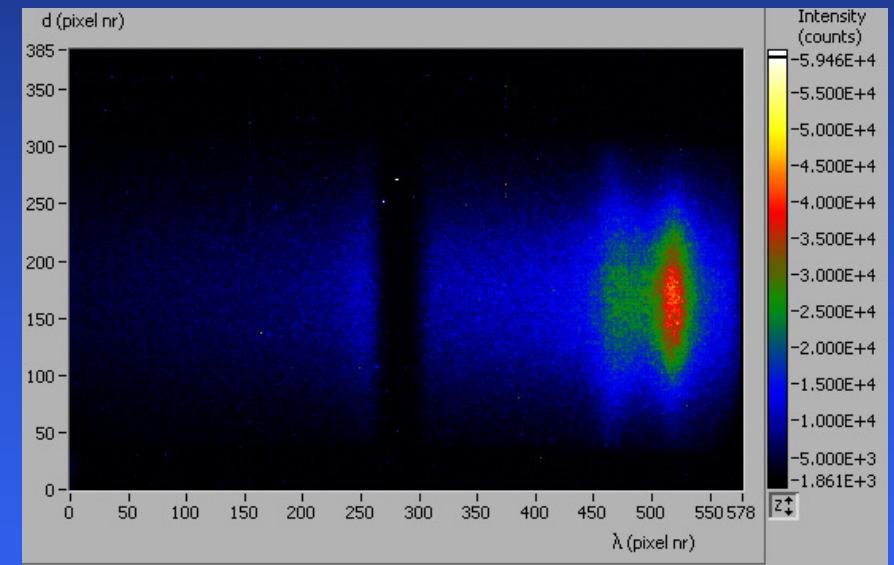
At λ_0 and in $\Delta \lambda$ - TS
A different plasma

Plasma Background Light 1

Thomson + Plasma



Plasma Background



The competitors

Collectivity

Rayleigh scattering
False stray light, vessels

Plasma photons
Laser produced plasma

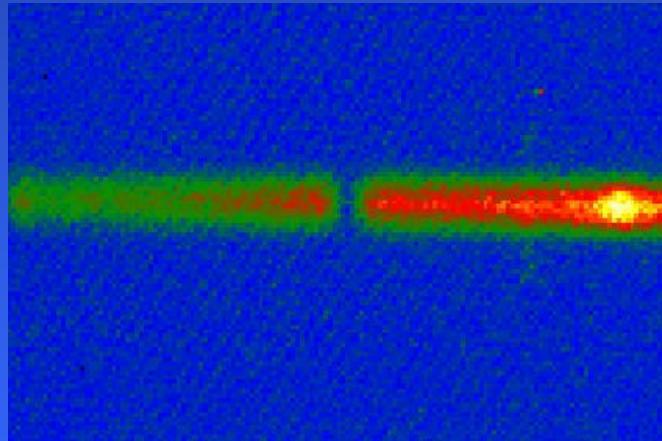
Change of Gaussian

At λ_0
At λ_0

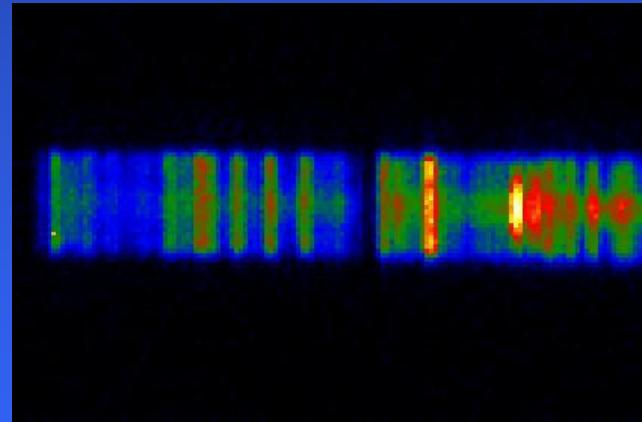
At λ_0 and in $\Delta \lambda - TS$
A different plasma

Plasma Background Light 2

150W Hg lamp



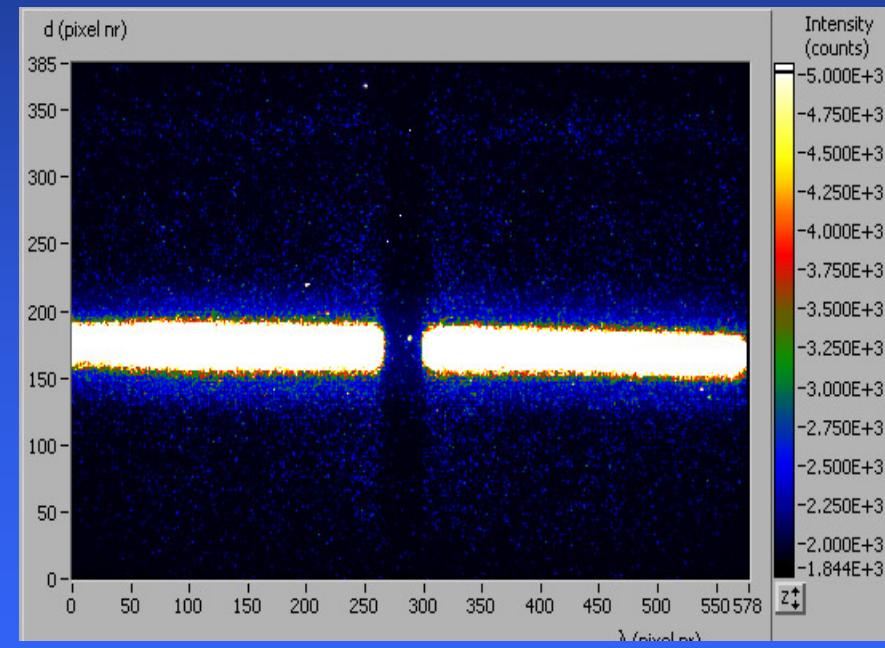
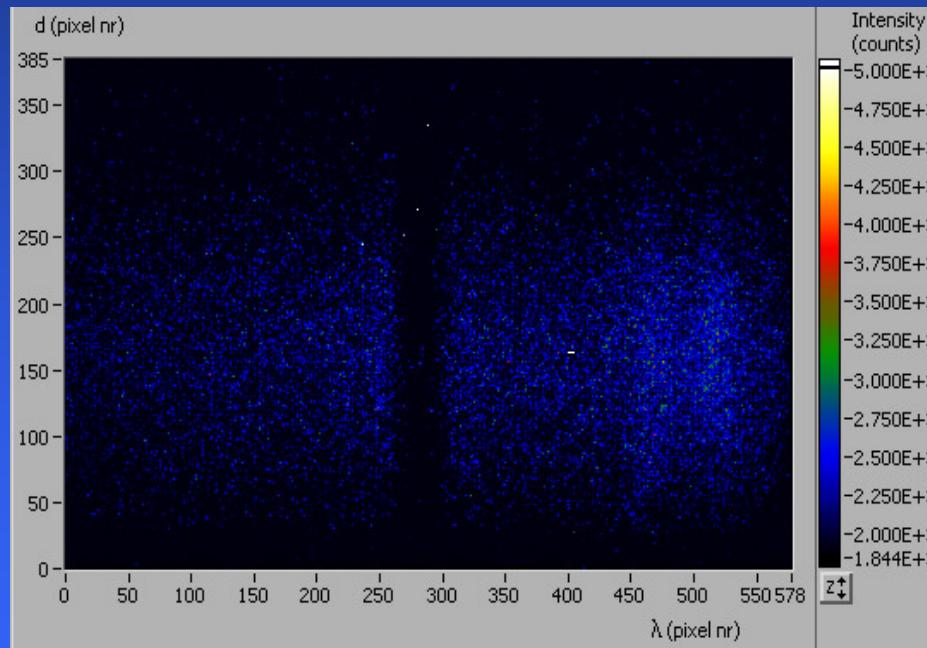
150W Hg+CeI₃ lamp



Laser Produced Plasma

Laser power: 4mJ.

Focal length:15cm.



Lamp_background_60s

TS_60s

TU/e

The competitors

Collectivity

Rayleigh scattering
False stray light, vessels

Plasma photons
Laser produced plasma

Change of Gaussian

At λ_0
At λ_0

At λ_0 and in $\Delta \lambda$ - TS
A different plasma

Solution Plasma light

Avoid LPP: Pulse energy < 1 mJ: **no extra plasma**

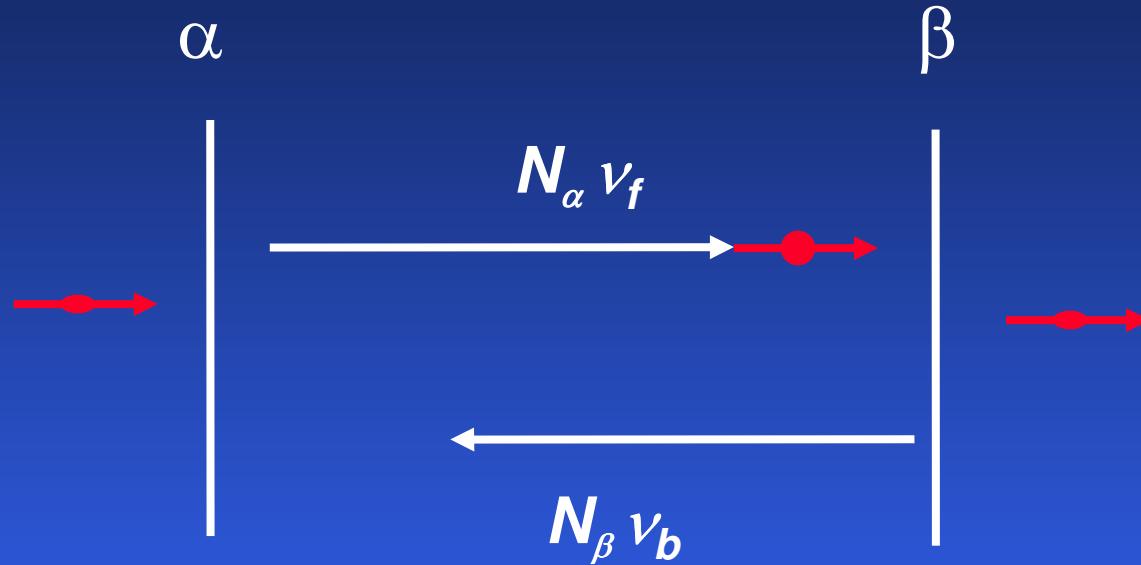
Compress **laser pulse** period

Method employed by Erik Kieft ASML

Using the Ekspla $\tau < 0.2 \text{ ns} !!$

Instead of “normal” $\tau < 8 \text{ ns}$

Equilibrium Departure



Non-Equilibrium
Equilibrium

$$N_\alpha v_f = N_\beta v_b + N_\beta v_t$$

$$N_\alpha^{eq} v_f = N_\beta^{eq} v_b$$

$$y = N/N^{eq}$$

$$y(\alpha) = y(\beta)[1 + \textcircled{v_t \tau_b}]$$

TU/e