# A new space of generalised functions with bounded variation motivated by fracture mechanics 

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#### Abstract

We introduce a new space of generalised functions with bounded variation to prove the existence of a solution to a minimum problem that arises in the variational approach to fracture mechanics in elastoplastic materials. We study the fine properties of the functions belonging to this space and prove a compactness result. In order to use the Direct Method of the Calculus of Variations we prove a lower semicontinuity result for the functional occurring in this minimum problem. Moreover, we adapt a nontrivial argument introduced by Friedrich to show that every minimizing sequence can be modified to obtain a new minimizing sequence that satisfies the hypotheses of our compactness result.


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## 1. Introduction

The variational approach to rate-independent evolution problems developed in [10] and [11] is based on a time discretization scheme, where the approximate solution at a given time is obtained by solving an incremental minimum problem which involves the solution at the previous time. The same approach was introduced independently in fracture mechanics in [8] (we refer also to [3] for further developments in this field).

In this framework, the study of crack growth in linearly elastic-perfectly plastic materials in the small strain regime leads to incremental minimization problems that involve the crack $\Gamma$ as well as the elastic part $e$ and the plastic part $p$ of the strain. In the (generalised) antiplane case, the reference configuration is a bounded Lipschitz domain $\Omega \subset \mathbb{R}^{d}$, the crack is a Borel set $\Gamma \subset \bar{\Omega}$,
with $\mathcal{H}^{d-1}(\Gamma)<+\infty$, and the displacement is a function $u: \Omega \backslash \Gamma \rightarrow \mathbb{R}$, whose gradient is additively decomposed as $D u=e+p$, where $e$, the elastic part, is an $L^{2}$-function defined in $\Omega \backslash \Gamma$ and $p$, the plastic part, is a bounded Radon measure defined on $\bar{\Omega} \backslash \Gamma$.

Given a Borel set $\Gamma_{0} \subset \bar{\Omega}$ (the crack at the previous time), with $\mathcal{H}^{d-1}\left(\Gamma_{0}\right)$ $<+\infty$, and a bounded Radon measure $p_{0}$ in $\bar{\Omega} \backslash \Gamma_{0}$ (the plastic strain at the previous time), the incremental minimum problem takes the form

$$
\begin{equation*}
\min \left\{\frac{1}{2} \int_{\Omega}|e|^{2} d x+\left|p-p_{0}\right|(\bar{\Omega} \backslash \Gamma)+\mathcal{H}^{d-1}\left(\Gamma \backslash \Gamma_{0}\right)\right\} \tag{1.1}
\end{equation*}
$$

where the minimum is taken over all competing cracks $\Gamma \supset \Gamma_{0}$ (irreversibility condition) and all pairs ( $e, p$ ) such that $e$ is an $L^{2}$ - function, $p$ is a Radon measure, and $e+p=D u$ in $\Omega \backslash \Gamma$ for some displacement $u$ satisfying prescribed boundary conditions (see Sect. 4 for a precise formulation).

The purpose of this paper is to prove the existence of a solution to problem (1.1). In [6] we considered the same problem only in the case $d=2$, with the additional constraint that $\Gamma$ and $\Gamma_{0}$ are compact and satisfy an a priori bound on the number of their connected components. In this case, given a minimizing sequence $\left(\Gamma_{k}, e_{k}, p_{k}\right)_{k}$, we can extract a subsequence (not relabelled) such that $\Gamma_{k} \rightarrow \Gamma$ in the Hausdorff metric, $e_{k} \rightharpoonup e$ weakly in $L^{2}$, and $p_{k} \stackrel{*}{\rightharpoonup} p$ locally weakly* as measures on $\bar{\Omega} \backslash \Gamma$. Therefore the existence of a solution to (1.1) follows from the Direct Method of the Calculus of Variations, since all terms in (1.1) are lower semicontinuous.

This simplified approach cannot be applied when $d>2$, nor when $d=2$ without bounds on the number of connected components of $\Gamma$. For this reason we prefer to rewrite the minimum problem (1.1) in terms of the displacement $u$, considered as a function defined $\mathcal{L}^{d}$-a.e. in $\Omega$.

Let $u_{0}$ be the displacement at the previous time, let $e_{0}$ be its elastic strain, so that $D u_{0}=e_{0}+p_{0}$, and let $f: \mathbb{R}^{d} \rightarrow \mathbb{R}$ be defined by $f(\xi)=\frac{1}{2}|\xi|^{2}$, if $|\xi| \leq 1$, and $f(\xi)=|\xi|-\frac{1}{2}$, if $|\xi| \geq 1$. Setting $v=u-u_{0}$, if $v \in B V(\Omega)$, the space of functions with bounded variation in $\Omega$, then $(\Gamma, u)$ is a solution to (1.1) if and only if $v$ is a solution of

$$
\begin{equation*}
\min \left\{\int_{\Omega} f\left(\nabla v+e_{0}\right) d x+\left|D^{c} v\right|(\Omega)+\int_{J_{v} \backslash \Gamma_{0}}|[v]| \wedge 1 d \mathcal{H}^{d-1}\right\} \tag{1.2}
\end{equation*}
$$

with suitable boundary conditions, and $\Gamma=\Gamma_{0} \cup\left\{x \in J_{v}:|[v](x)| \geq 1\right\}$, see Lemmas 4.1 and 4.2. Here and in the rest of the paper $\nabla v$ is the approximate gradient of $v$ (see (2.6)), $J_{v}$ and $[v]$ are the jump set and the jump of $v$ (see Sect. 2), $D^{c} v$ is the Cantor part of the distributional gradient $D v$ of $v$ (see (2.5)), and $a \wedge b:=\min \{a, b\}$ for every $a, b \in \mathbb{R}$.

Unfortunately, there are boundary conditions for which the minimum problem (1.2) has no solution in the space of functions of bounded variation, as shown by the example provided in Proposition 6.5. The reason is that, while the first term in (1.2) controls $\nabla v$ and the second one controls $D^{c} v$, the third term does not control the whole jump part of $D v$, which is given by the integral of $|[v]|$ on $J_{v}$.

Therefore, in order to prove the existence of a solution to the minimum problem (1.2), and hence (1.1), we consider a larger functional space for the admissible displacements, which we denote by $G B V_{\star}(\Omega)$. This is a subset of the space $G B V(\Omega)$ of generalised functions of bounded variation introduced in [1, Sect. 1] (see also [2, Definition 4.26]). In Sect. 3 we study the fine properties of functions in $G B V_{\star}(\Omega)$, as well as some structure properties of this space. In particular, we prove in Theorem 3.11 that, if $\left(v_{k}\right)_{k}$ is a minimizing sequence of (1.2) in $G B V_{\star}(\Omega)$ and

$$
\begin{equation*}
\sup _{k} \int_{\Omega} \psi\left(\left|v_{k}\right|\right) d x<+\infty \tag{1.3}
\end{equation*}
$$

for some continuous function $\psi$ with $\psi(t) \rightarrow+\infty$ as $t \rightarrow+\infty$, then a subsequence of $\left(v_{k}\right)_{k}$ converges pointwise $\mathcal{L}^{d}$-a.e. to a function $v \in G B V_{\star}(\Omega)$.

In Sect. 5 we prove that every minimizing sequence of problem (1.2) can be modified in order to obtain a new minimizing sequence which satisfies (1.3) for a suitable function $\psi$, depending on the sequence. The construction of $\psi$ is not trivial and is achieved by adapting to $G B V_{\star}(\Omega)$ the arguments introduced in [9] for $G S B V^{p}(\Omega)$.

To use the Direct Method of the Calculus of Variations we prove in Theorem 6.1 that the functional considered in (1.2) is lower semicontinuous. If $e_{0}$ is constant we can easily reduce the problem to the case $e_{0}=0$, which was studied in [4]. The result can be easily extended to the piecewise constant case by a localization argument. The general case is obtained by approximation.

The existence of a minimizing sequence satisfying (1.3) and the semicontinuity result imply that there exists a solution to (1.2) in $G B V_{\star}(\Omega)$, see Theorem 6.2. Since problem (1.1) is equivalent to problem (1.2) in $G B V_{\star}(\Omega)$, see Lemmas 4.1 and 4.2, we conclude that problem (1.1) has a solution, see Corollary 6.3.

## 2. Preliminaries on $B V$-spaces

In this section we fix the general notation used in the paper and we recall the fine properties of $B V$ and $G B V$ functions that will be used in the sequel.

For every topological space $X$, for every Borel set $Y \subset X$, and for every finite-dimensional Hilbert space $\Xi$, the space of $\Xi$-valued bounded Borel measures on $Y$ is denoted by $\mathcal{M}_{b}(Y ; \Xi)$.

Throughout this section $U$ is an open set in $\mathbb{R}^{d}$. For every set $E \subset U$ the characteristic function $\chi_{E}: U \rightarrow \mathbb{R}$ is defined by $\chi_{E}(x)=1$ if $x \in E$ and $\chi_{E}(x)=0$ if $x \notin E$. For every Borel measure $\mu$ on $U$ and every Borel set $E \subset U$ the Borel measure $\mu\llcorner E$ on $U$ is defined by $\mu\llcorner E(B)=\mu(B \cap E)$ for every Borel set $B \subset U$. If $f \in L^{1}(U ; \mu)$ the Borel measure $f \mu$ on $U$ is defined by $(f \mu)(B):=\int_{B} f d \mu$ for every Borel set $B \subset U$.

The Lebesgue measure is denoted by $\mathcal{L}^{d}$ and the $(d-1)$-dimensional Haussdorff measure by $\mathcal{H}^{d-1}$. For every $\mathcal{L}^{d}$-measurable set $E \subset U$ and every $\alpha \in[0,1]$ the set $E^{(\alpha)}$ of points in $U$ of $\mathcal{L}^{d}$-density $\alpha$ for $E$ is defined by

$$
E^{(\alpha)}:=\left\{x \in U: \lim _{\rho \rightarrow 0+} \frac{\mathcal{L}^{d}\left(E \cap B_{\rho}(x)\right)}{\mathcal{L}^{d}\left(B_{\rho}(x)\right)}=\alpha\right\}
$$

where $B_{\rho}(x)$ is the ball in $\mathbb{R}^{d}$ of radius $\rho$ and centre $x$.
Given an $\mathcal{L}^{d}$-measurable set $E \subset U$ and an $\mathcal{L}^{d}$-measurable function $u: E \rightarrow \overline{\mathbb{R}}$, we say that $a \in \overline{\mathbb{R}}$ is the approximate limit of $u(y)$ as $y$ tends to a point $x \in E^{(\alpha)}$ for some $\alpha>0$, in symbols

$$
\underset{y \rightarrow x}{\operatorname{ap}} \lim _{x} u(y)=a,
$$

if for every neighbourhood $A$ of $a$ in $\overline{\mathbb{R}}$ we have

$$
\begin{equation*}
\lim _{\rho \rightarrow 0+} \rho^{-d} \mathcal{L}^{d}\left(\left\{y \in E \cap B_{\rho}(x): u(y) \notin A\right\}\right)=0 \tag{2.1}
\end{equation*}
$$

It follows from the definition that if $f: \overline{\mathbb{R}} \rightarrow \overline{\mathbb{R}}$ is a continuous function and $\underset{y \rightarrow x}{\operatorname{ap}} \lim _{x} u(y)$ exists, then

$$
\begin{equation*}
\operatorname{ap}_{y \rightarrow x} \lim _{f} f(u(y))=f\left(\operatorname{ap}_{y \rightarrow x} \lim _{y} u(y)\right) . \tag{2.2}
\end{equation*}
$$

Moreover, if $u, v: E \rightarrow \overline{\mathbb{R}}$ are $\mathcal{L}^{d}$-measurable functions, $F$ is an $\mathcal{L}^{d}$-measurable subset of $E, u=v$ in $F$, and $x \in F^{(1)}$, then $\operatorname{app}_{y \rightarrow x} \lim _{x} u(y)$ exists if and only if ap $\lim v(y)$ exists, and in this case
$y \rightarrow x$

$$
\begin{equation*}
\underset{y \rightarrow x}{\operatorname{ap}} \lim _{x} v(y)=\underset{y \rightarrow x}{\operatorname{ap}} \lim _{x} u(y) . \tag{2.3}
\end{equation*}
$$

Let $u: U \rightarrow \overline{\mathbb{R}}$ be an $\mathcal{L}^{d}$-measurable function. We define the jump set $J_{u}$ as the set of all points $x \in U$ such that there exist $u^{+}(x), u^{-}(x) \in \overline{\mathbb{R}}$ with $u^{+}(x) \neq u^{-}(x)$ and a unit vector $\nu_{u}(x) \in \mathbb{R}^{d}$ such that, setting

$$
U^{+}:=\left\{y \in U:(y-x) \cdot \nu_{u}(x)>0\right\} \quad \text { and } \quad U^{-}:=\left\{y \in U:(y-x) \cdot \nu_{u}(x)<0\right\},
$$

we have

$$
u^{+}(x)=\left.\underset{y \rightarrow x}{\operatorname{ap} \lim } u\right|_{U^{+}}(y) \quad \text { and } \quad u^{-}(x)=\left.\underset{y \rightarrow x}{\operatorname{ap}} \lim _{y} u\right|_{U^{-}}(y) .
$$

It is easy to see that the triple $\left(u^{+}(x), u^{-}(x), \nu_{u}(x)\right)$ is uniquely defined up to a swap of the first two terms and a change of sign in the third one. For every $x \in J_{u}$ we set $[u](x):=u^{+}(x)-u^{-}(x)$. It can be proved that $J_{u}$ is a Borel set and $[u]: J_{u} \rightarrow \overline{\mathbb{R}}$ is a Borel function.

For the general properties of the space $B V(U)$ of functions of bounded variation on $U$ we refer to [7, Chapter 5] and [2, Chapter 3]. Let us fix $u \in$ $B V(U)$. As a consequence of [2, Theorem 3.78], we have that for $\mathcal{H}^{d-1}$-a.e. $x \in U \backslash J_{u}$ there exists

$$
\tilde{u}(x):=\underset{y \rightarrow x}{\operatorname{ap} \lim } u(y) .
$$

Moreover, the function $\tilde{u}$ defined by this formula is finite for $\mathcal{H}^{d-1}$-a.e. $x \in$ $U \backslash J_{u}$. For every $m \in \mathbb{R}_{+}$we set

$$
\begin{equation*}
\{|\tilde{u}| \leq m\}:=\left\{x \in U \backslash J_{u}: \tilde{u}(x) \text { exists and }|\tilde{u}(x)| \leq m\right\} . \tag{2.4}
\end{equation*}
$$

The set $\{|\tilde{u}|>m\}$ is defined in a similar way.
By the definition of $B V(U)$ the distributional gradient of $u$, denoted by $D u$, belongs to $\mathcal{M}_{b}\left(U ; \mathbb{R}^{d}\right)$. It can be decomposed as

$$
\begin{equation*}
D u=\nabla u \mathcal{L}^{d}+D^{c} u+D^{j} u \tag{2.5}
\end{equation*}
$$

where $\nabla u \in L^{1}\left(U ; \mathbb{R}^{d}\right), D^{j} u:=D u\left\llcorner J_{u}\right.$ is the jump part of $D u$, while $D^{c} u$, called the Cantor part of $D u$, is a bounded Radon measure with values in $\mathbb{R}^{d}$, singular with respect to $\mathcal{L}^{d}$ and such that $D^{c} u(B)=0$ on every Borel set $B \subset U$ with $\mathcal{H}^{d-1}(B)<+\infty$. Moreover, for $\mathcal{L}^{d}$ - a.e. $x \in U$ the vector $\nabla u(x)$ is the approximate gradient of $u$ at $x$, i.e.,

$$
\begin{equation*}
\operatorname{ap}_{y \rightarrow x} \lim _{x} \frac{u(y)-\tilde{u}(x)-\nabla u(x) \cdot(y-x)}{|y-x|}=0 . \tag{2.6}
\end{equation*}
$$

Finally, we have

$$
\begin{equation*}
D^{j} u=[u] \nu_{u} \mathcal{H}^{d-1}\left\llcorner J_{u}\right. \tag{2.7}
\end{equation*}
$$

For every $t \in \mathbb{R}$ and $m \in \mathbb{R}_{+}$we set $t^{(m)}:=(t \wedge m) \vee(-m)$ and note that the function $t \mapsto t^{(m)}$ is Lipschitz continuous with constant 1 . If $u$ is an $\overline{\mathbb{R}}$ valued function defined in $U, u^{(m)}$ is the function defined by $u^{(m)}(x)=u(x)^{(m)}$ for every $x \in U$.

The following theorem gives a formula for the distributional gradient of the truncation of a $B V$ function.
Theorem 2.1. Let $u \in B V(U)$ and let $m \in \mathbb{R}_{+}$. Then $u^{(m)} \in B V(U)$ and

$$
D u^{(m)}=\chi_{\{|u| \leq m\}} \nabla u \mathcal{L}^{d}+\chi_{\{|\tilde{u}| \leq m\}} D^{c} u+\left[u^{(m)}\right] \nu_{u} \mathcal{H}^{d-1}\left\llcorner J_{u}\right.
$$

Proof. It is enough to apply [2, Theorem 3.99] to $f(t)=t^{(m)}$.
Remark 2.2. Let $u \in B V(U)$ and let $m \in \mathbb{R}_{+}$. It follows from Theorem 2.1 that

$$
\begin{align*}
& \nabla u^{(m)}=0 \quad \mathcal{L}^{d} \text {-a.e. in }\{|u|>m\},  \tag{2.8}\\
& D^{c} u^{(m)}(B)=0 \quad \text { for every Borel set } B \subset\{|\tilde{u}|>m\} . \tag{2.9}
\end{align*}
$$

The following lemma provides a strong localization property for $D u$.
Lemma 2.3. Let $u, v \in B V(U)$ and let $E$ be a Borel set contained in $U \backslash\left(J_{u} \cup\right.$ $\left.J_{v}\right)$. Suppose that $\tilde{u}=\tilde{v} \mathcal{H}^{d-1}$-a.e. in $E$. Then $\nabla u=\nabla v \mathcal{L}^{d}$-a.e. in $E$ and $D^{c} u(B)=D^{c} v(B)$ for every Borel set $B \subset E$.

Proof. It is enough to apply [2, Proposition 3.92 and Remark 3.93].
We refer to [2, Chapter 3] for the definition and the main properties of the sets of finite perimeter. If $E \subset U$ is a $\mathcal{L}^{d}$-measurable set, its perimeter in $U$ is denoted by $P(E, U)$. When $P(E, U)<+\infty$, the reduced boundary of $E$ in $U$ is denoted by $\partial^{*} E$ and for every $x \in \partial^{*} E$ the approximate inner unit normal vector is denoted by $\nu_{E}(x)$.

The following lemma provides a precise formula for the gradient of the product between a bounded $B V$ function and the characteristic function of a set with finite perimeter.

Lemma 2.4. Let $u \in B V(U) \cap L^{\infty}(U)$ and let $E \subset U$ be a $\mathcal{L}^{d}$-measurable set with $P(E, U)<+\infty$. Then there exists $\gamma_{E} u \in L^{\infty}\left(\partial^{*} E, \mathcal{H}^{d-1}\right)$ such that for $\mathcal{H}^{d-1}$-a.e. $x \in \partial^{*} E$ we have

$$
\lim _{\rho \rightarrow 0+} \frac{1}{\rho^{d}} \int_{B_{\rho}^{+}(x)}\left|u(y)-\left(\gamma_{E} u\right)(x)\right| d y=0
$$

where $B_{\rho}^{+}(x)=\left\{y \in B_{\rho}(x):(y-x) \cdot \nu_{E}(x)>0\right\}$. Moreover, setting $v:=u \chi_{E}$ we have $v \in B V(U) \cap L^{\infty}(U)$ and

$$
\begin{align*}
\nabla v & =\chi_{E} \nabla u  \tag{2.10}\\
D^{c} v & =\chi_{E^{(1)}} D^{c} u  \tag{2.11}\\
D^{j} v=\chi_{E^{(1)}} D^{j} u & +\left(\gamma_{E} u\right) \nu_{E} \mathcal{H}^{d-1}\left\llcorner\partial^{*} E .\right. \tag{2.12}
\end{align*}
$$

Proof. The statement about $\gamma_{E} u$ is proved in [2, Theorem 3.77]. The properties concerning $v$ can be easily deduced from [2, Theorem 3.84].

We shall use the space $G B V(U)$ of generalised functions of bounded variation introduced in [1, Sect. 1] (see also [2, Definition 4.26]).

Remark 2.5. Since $u^{(m)}=\left(u^{(n)}\right)^{(m)}$ for $0<m<n$, Remark 2.2 implies that (2.8) and (2.9) hold also for every $u \in G B V(U)$ and every $m \in \mathbb{R}_{+}$.

In the following proposition we summarize the fine properties of functions in $G B V(U)$.

Proposition 2.6. Let $u \in G B V(U)$. Then the following properties hold:
(a) (precise values) for $\mathcal{H}^{d-1}$-a.e. $x \in U \backslash J_{u}$ there exists

$$
\begin{equation*}
\tilde{u}(x):=\underset{y \rightarrow x}{\operatorname{ap} \lim } u(y) \in \overline{\mathbb{R}} ; \tag{2.13}
\end{equation*}
$$

moreover, if $x \in\{|u| \leq m\}^{(1)} \backslash J_{u}$ for some $m \in \mathbb{R}_{+}$and $\tilde{u}(x)$ exists, then $\tilde{u}(x)=\widetilde{u^{(m)}}(x) \in \mathbb{R}$; in particular, $\tilde{u}=\widetilde{u^{(m)}} \mathcal{H}^{d-1}$-a.e. on $\{|u| \leq$ $m\}^{(1)} \backslash J_{u}$;
(b) (approximate differentiability) there exists a Borel function, denoted by $\nabla u: U \rightarrow \mathbb{R}^{d}$, such that for $\mathcal{L}^{d}$-a.e. $x \in U$ we have $\tilde{u}(x) \in \mathbb{R}$ and (2.6) holds; moreover, for every $m \in \mathbb{R}_{+}$we have

$$
\begin{equation*}
\nabla u(x)=\nabla u^{(m)}(x) \quad \text { for } \mathcal{L}^{d} \text {-a.e. } x \in\{|u| \leq m\} ; \tag{2.14}
\end{equation*}
$$

(c) (jumps) the set $J_{u}$ is countably $\mathcal{H}^{d-1}$-rectifiable and for $\mathcal{H}^{d-1}$-a.e. $x \in J_{u}$ the vector $\nu_{u}(x)$ is orthogonal to the approximate tangent space to $J_{u}$ at $x$ (according to [2, Definition 2.86]); moreover, for every $m \in \mathbb{R}_{+}$we have $J_{u^{(m)}} \subset J_{u}$ up to a set of $\mathcal{H}^{d-1}$-measure zero and $\left|\left[u^{(m)}\right]\right| \leq|[u]| \mathcal{H}^{d-1}$ a.e. on $J_{u^{(m)}} \cap J_{u}$; finally, for $\mathcal{H}^{d-1}$-a.e. $x \in J_{u}$, there exists $m_{x} \in \mathbb{N}$ such that $x \in J_{u^{(m)}}$ for every $m \in \mathbb{N}$ with $m \geq m_{x}$ and $\left[u^{(m)}\right](x) \rightarrow[u](x)$ as $m \rightarrow \infty$ with $m \in \mathbb{N}$;
(d) (Cantor part) for every $m, n \in \mathbb{R}_{+}$, with $m \leq n$, we have $D^{c} u^{(m)}(B)=$ $D^{c} u^{(n)}(B)$ for every Borel set $B \subset\{|\tilde{u}| \leq m\}$ and $\left|D^{c} u^{(m)}\right|(E) \leq$ $\left|D^{c} u^{(n)}\right|(E)$ for every Borel set $E \subset U$.

Proof. Properties (a), (b), and (c) can be deduced from [2, Theorem 4.34]. Equality $\tilde{u}(x)=\widetilde{u^{(m)}}(x)$ in (a) follows from (2.3).

Let us prove (d). For every $x \in\{|\tilde{u}| \leq m\}$ we have $\tilde{u}(x)^{(m)}=\tilde{u}(x)$, therefore the Lipschitz continuity of $t \mapsto t^{(m)}$ implies that $\left|u^{(m)}(y)-\tilde{u}(x)\right| \leq \mid u(y)-$ $\tilde{u}(x) \mid$ for every $y \in U$. The definition of ap $\lim$ then gives that $\operatorname{ap~}_{y \rightarrow x} \lim ^{(m)}(y)=$ $\tilde{u}(x)$. Hence $\widetilde{u^{(m)}}(x)=\tilde{u}(x)$, and the same result holds for $u^{(n)}$. The conclusion follows from Lemma 2.3 and Remark 2.5.

In the following theorem we show that, if $u \in G B V(U)$ satisfies condition (2.15) below, then we can define an $\mathbb{R}^{d}$-valued Radon measure that plays the role of the Cantor part of $D u$, even if the measure $D u$ cannot be defined.

Theorem 2.7. Let $u \in G B V(U)$ be such that

$$
\begin{equation*}
\sup _{m>0}\left|D^{c} u^{(m)}\right|(U)<+\infty . \tag{2.15}
\end{equation*}
$$

Then there exists a unique measure $\mu \in \mathcal{M}_{b}\left(U ; \mathbb{R}^{d}\right)$ such that for every $m \in \mathbb{R}_{+}$

$$
\begin{equation*}
\mu(B)=D^{c} u^{(m)}(B) \quad \text { for every Borel set } B \subset\{|\tilde{u}| \leq m\} \tag{2.16}
\end{equation*}
$$

and

$$
\begin{equation*}
\mu(B)=0 \quad \text { for every Borel set } B \subset\{|\tilde{u}|=+\infty\} \cup J_{u} \tag{2.17}
\end{equation*}
$$

Proof. Let $E_{0}:=\emptyset$ and for every $k \in \mathbb{N}$ let $E_{k}:=\{|\tilde{u}| \leq k\}$. By Proposition 2.6(d) for every $n \in \mathbb{N}$ we have

$$
\begin{equation*}
\sum_{k=1}^{n}\left|D^{c} u^{(k)}\right|\left(E_{k} \backslash E_{k-1}\right)=\sum_{k=1}^{n}\left|D^{c} u^{(n)}\right|\left(E_{k} \backslash E_{k-1}\right) \leq\left|D^{c} u^{(n)}\right|(U) \tag{2.18}
\end{equation*}
$$

hence, by (2.15),

$$
\begin{equation*}
\sum_{k=1}^{+\infty}\left|D^{c} u^{(k)}\right|\left(E_{k} \backslash E_{k-1}\right)<+\infty \tag{2.19}
\end{equation*}
$$

If $B \subset U$ is a Borel set we define

$$
\begin{equation*}
\mu(B):=\sum_{k=1}^{+\infty} D^{c} u^{(k)}\left(B \cap E_{k} \backslash E_{k-1}\right) \tag{2.20}
\end{equation*}
$$

Recalling (2.4) it follows immediately from the definition that (2.17) holds. By (2.19) the series converges absolutely and its sum is finitely additive with respect to $B$. To prove the countable additivity of $\mu$ it is enough to show that

$$
\begin{equation*}
\mu\left(B^{j}\right) \rightarrow 0 \tag{2.21}
\end{equation*}
$$

whenever $\left(B^{j}\right)_{j}$ is a decreasing sequence of Borel sets with empty intersection. In this case by definition we have

$$
\mu\left(B^{j}\right)=\sum_{k=1}^{+\infty} D^{c} u^{(k)}\left(B^{j} \cap E_{k} \backslash E_{k-1}\right)
$$

and for every $k \in \mathbb{N}$ we have $D^{c} u^{(k)}\left(B^{j} \cap E_{k} \backslash E_{k-1}\right) \rightarrow 0$ as $j \rightarrow \infty$. This implies (2.21) by the dominated convergence theorem for series, which can be applied thanks to (2.19).

To prove (2.16) we fix $m \in \mathbb{R}_{+}$and a Borel set $B \subset\{|\tilde{u}| \leq m\}$. Let $n$ be the smallest integer larger than or equal to $m$. Since $\{|\tilde{u}| \leq m\} \subset E_{n}$, we have $B \subset \bigcup_{k=1}^{n}\left(E_{k} \backslash E_{k-1}\right)$ and $B \cap E_{k} \backslash E_{k-1}=\emptyset$ for $k>n$. Therefore (2.20) and Proposition 2.6(d) give

$$
\begin{aligned}
\mu(B) & =\sum_{k=1}^{n} D^{c} u^{(k)}\left(B \cap E_{k} \backslash E_{k-1}\right)=\sum_{k=1}^{n} D^{c} u^{(n)}\left(B \cap E_{k} \backslash E_{k-1}\right) \\
& =D^{c} u^{(n)}(B)=D^{c} u^{(m)}(B) .
\end{aligned}
$$

This concludes the proof.
Definition 2.8. Assume that $u \in G B V(U)$ satisfies (2.15). The measure $\mu$ introduced in Theorem 2.7 is denoted by $D^{c} u$.

Lemma 2.3 ensures that this definition is consistent with (2.5) whenever $u \in B V(U)$. The following proposition shows that the total variation $\left|D^{c} u\right|$ of $D^{c} u$ coincides with the measure introduced in [2, Definition 4.33].

Proposition 2.9. Assume that $u \in G B V(U)$ satisfies (2.15). Then

$$
\begin{align*}
& D^{c} u^{(m)}(B) \rightarrow D^{c} u(B) \text { as } m \rightarrow+\infty,  \tag{2.22}\\
& \lim _{m \rightarrow+\infty}\left|D^{c} u^{(m)}\right|(B)=\sup _{m>0}\left|D^{c} u^{(m)}\right|(B)=\left|D^{c} u\right|(B), \tag{2.23}
\end{align*}
$$

for every Borel set $B \subset U$.
Proof. Using the notation in the proof of Theorem 2.7, by (2.16) and (2.17) we have

$$
\begin{equation*}
D^{c} u^{(m)}(B \cap\{|\tilde{u}| \leq m\})=D^{c} u(B \cap\{|\tilde{u}| \leq m\}) \rightarrow D^{c} u(B) . \tag{2.24}
\end{equation*}
$$

From Remark 2.5 it follows that $D^{c} u^{(m)}\left(B \backslash\left(\{|\tilde{u}| \leq m\} \cup J_{u}\right)\right)=0$. Since $D^{c} u^{(m)}\left(J_{u}\right)=0$, we conclude that $D^{c} u^{(m)}(B)=D^{c} u^{(m)}(B \cap\{|\tilde{u}| \leq m\})$. Together with (2.24) this implies (2.22).

The first equality in (2.23) follows from Proposition 2.6(d). To prove the other equality, for every Borel set $B \subset U$ we define

$$
\begin{equation*}
\nu(B):=\lim _{m \rightarrow+\infty}\left|D^{c} u^{(m)}\right|(B)=\sup _{m>0}\left|D^{c} u^{(m)}\right|(B) . \tag{2.25}
\end{equation*}
$$

Using the monotonicity with respect to $m$ stated in Proposition 2.6(d) we can prove that $\nu$ is a Borel measure. Therefore, it is enough to prove (2.23) for every Borel set $B \subset\{|\tilde{u}| \leq m\}$ and for every Borel set $B \subset\{|\tilde{u}|=+\infty\} \cup J_{u}$. The former follows from (2.16), while the latter follows from (2.17) and Remark 2.5, taking into account the fact that $\left|D^{c} u^{(m)}\right|\left(J_{u}\right)=0$ by the general properties of the Cantor part of the gradient of a $B V$ function.

## 3. The function space used in our problem

We now introduce the function space that will be used to formulate and solve problem (1.1). Throughout this section $U$ is a bounded open set in $\mathbb{R}^{d}$.

Definition 3.1. The space $G B V_{\star}(U)$ is defined as the space of functions $u: U \rightarrow$ $\mathbb{R}$ such that $u^{(m)} \in B V(U)$ for every $m \in \mathbb{R}_{+}$and

$$
\begin{equation*}
\sup _{m \in \mathbb{R}_{+}}\left(\int_{U}\left|\nabla u^{(m)}\right| d x+\left|D^{c} u^{(m)}\right|(U)+\int_{J_{u(m)}}\left|\left[u^{(m)}\right]\right| \wedge 1 d \mathcal{H}^{d-1}\right)<+\infty \tag{3.1}
\end{equation*}
$$

It follows from the definition of $G B V(U)$ that $G B V_{\star}(U) \subset G B V(U)$. Moreover, using Remark 2.5 and Proposition 2.6 it is easy to see that the supremum in (3.1) can be taken over $m \in \mathbb{N}$.

Remark 3.2. If $d=1$ and $u \in G B V_{\star}(U)$, then (3.1) implies that there exist at most a finite number of jump points $x$ of $u$ with $|[u](x)| \geq 1$. From this property and from (3.1) we can deduce that $u \in B V(U)$. Hence $G B V_{\star}(U)=B V(U)$ if $d=1$.

On the contrary, when $d \geq 2$ we have $G B V_{\star}(U) \neq B V(U)$. Indeed, let $x_{0} \in U$ and let $R>0$ be such that $B_{R}\left(x_{0}\right) \subset U$. For every $k \in \mathbb{N}$ let $R_{k}:=2^{-k} R$ and let $u: U \rightarrow \mathbb{R}$ be defined by $u(x):=1 / R_{k}^{d-1}$, if $x \in$ $B_{R_{k}}\left(x_{0}\right) \backslash B_{R_{k+1}}\left(x_{0}\right)$ for $k \in \mathbb{N}$, and $u(x):=0$, if $x \in U \backslash B_{R}\left(x_{0}\right)$. For every $m \in \mathbb{R}_{+}$we have $\nabla u^{(m)}=0, D^{c} u^{(m)}=0$, and

$$
\int_{J_{u(m)}}\left|\left[u^{(m)}\right]\right| \wedge 1 d \mathcal{H}^{d-1} \leq \sigma_{d-1} \sum_{k} R_{k}^{d-1}<+\infty
$$

where $\sigma_{d-1}:=\mathcal{H}^{d-1}\left(B_{1}(0)\right)$. This shows that $u \in G B V_{\star}(U)$. Since
$\int_{J_{u}}|[u]| d \mathcal{H}^{d-1}=\sigma_{d-1} \sum_{k=1}^{\infty}\left(\frac{1}{R_{k}^{d-1}}-\frac{1}{R_{k-1}^{d-1}}\right) R_{k}^{d-1}=\sigma_{d-1} \sum_{k=1}^{\infty}\left(1-\frac{1}{2^{d-1}}\right)=+\infty$,
we conclude that $u \notin B V(U)$.
For every $u \in G B V(U)$ we define

$$
\begin{equation*}
J_{u}^{1}:=\left\{x \in J_{u}:|[u](x)| \geq 1\right\} . \tag{3.2}
\end{equation*}
$$

Proposition 3.3. The space $G B V_{\star}(U)$ coincides with the set of functions $u \in$ $G B V(U)$ such that

$$
\begin{align*}
& \nabla u \in L^{1}\left(U ; \mathbb{R}^{d}\right),  \tag{3.3}\\
& \sup _{m \in \mathbb{R}_{+}}\left|D^{c} u^{(m)}\right|(U)<+\infty  \tag{3.4}\\
& \int_{J_{u}}|[u]| \wedge 1 d \mathcal{H}^{d-1}<+\infty \tag{3.5}
\end{align*}
$$

Remark 3.4. Inequality (3.5) is equivalent to

$$
\begin{equation*}
\int_{J_{u} \backslash J_{u}^{1}}|[u]| d \mathcal{H}^{d-1}<+\infty \quad \text { and } \quad \mathcal{H}^{d-1}\left(J_{u}^{1}\right)<+\infty \tag{3.6}
\end{equation*}
$$

Therefore, every $u \in G B V_{\star}(U)$ satisfies (3.6).
Proof of Proposition 3.3. Assume $u \in G B V_{\star}(U)$. Since $U$ is the union of the sets $\{|u| \leq m\}$ for $m \in \mathbb{N}$ (recall that $u$ is finite valued), from (2.14) and (3.1) we deduce (3.3). Inequality (3.4) follows from (3.1), while (3.5) can be deduced from (3.1) and Proposition 2.6(c).

Conversely, assume $u \in G B V(U)$ and that (3.3)-(3.5) are satisfied. From the definition of $G B V(U)$, see [2, Definition 4.26], for every $m \in \mathbb{R}_{+}$we have $u^{(m)} \in B V_{l o c}(U)$. To prove that $u^{(m)} \in B V(U)$ it is enough to show that $\left|D u^{(m)}\right|(U)<+\infty$. By the extension of (2.5) to $B V_{l o c}(U)$ we have

$$
\begin{equation*}
D u^{(m)}=\nabla u^{(m)} \mathcal{L}^{d}+D^{c} u^{(m)}+\left[u^{(m)}\right] \nu_{u^{(m)}} \mathcal{H}^{d-1} L J_{u^{(m)}} . \tag{3.7}
\end{equation*}
$$

By (2.14) we have that $\nabla u^{(m)}=\nabla u \mathcal{L}^{d}$-a.e. in $\{|u| \leq m\}$, while by Remark $2.5 \nabla u^{(m)}=0 \mathcal{L}^{d}$-a.e. in $\{|u|>m\}$. Therefore (3.3) ensures that $\nabla u^{(m)} \in L^{1}\left(U ; \mathbb{R}^{d}\right)$ and

$$
\begin{equation*}
\sup _{m \in \mathbb{R}_{+}} \int_{U}\left|\nabla u^{(m)}\right| d x<+\infty . \tag{3.8}
\end{equation*}
$$

By Proposition 2.6(c) we have $J_{u^{(m)}} \subset J_{u}$ up to a set of $\mathcal{H}^{d-1}$ measure zero and $\left|\left[u^{(m)}\right]\right| \leq|[u]| \mathcal{H}^{d-1}$-a.e. on $J_{u^{(m)}}$. Therefore

$$
\begin{align*}
\int_{J_{u}(m)}\left|\left[u^{(m)}\right]\right| d \mathcal{H}^{d-1} & \leq \int_{J_{u}(m) \backslash J_{u}^{1}}\left|\left[u^{(m)}\right]\right| d \mathcal{H}^{d-1}+2 m \mathcal{H}^{d-1}\left(J_{u}^{1}\right) \\
& \leq \int_{J_{u} \backslash J_{u}^{1}}|[u]| d \mathcal{H}^{d-1}+2 m \mathcal{H}^{d-1}\left(J_{u}^{1}\right)<+\infty \tag{3.9}
\end{align*}
$$

where the last inequality follows from (3.6). By (3.4), (3.7), (3.8), and (3.9) we conclude that $\left|D u^{(m)}\right|(U)<+\infty$, which implies $u^{(m)} \in B V(U)$.

Finally, recalling again Proposition 2.6(c), by (3.5) we have

$$
\begin{equation*}
\int_{J_{u}(m)}\left|\left[u^{(m)}\right]\right| \wedge 1 d \mathcal{H}^{d-1} \leq \int_{J_{u}}|[u]| \wedge 1 d \mathcal{H}^{d-1}<+\infty \tag{3.10}
\end{equation*}
$$

Together with (3.4) and (3.8), this implies (3.1), hence $u \in G B V_{\star}(U)$.
Remark 3.5. Let $\lambda>0$ and $u \in G B V(U)$. Then

$$
\int_{J_{u}}|[u]| \wedge 1 d \mathcal{H}^{d-1}<+\infty \quad \Longleftrightarrow \quad \int_{J_{u}}|[u]| \wedge \lambda d \mathcal{H}^{d-1}<+\infty
$$

Indeed, if $\int_{J_{u}}|[u]| \wedge 1 d \mathcal{H}^{d-1}<+\infty$ and $\lambda \leq 1$ we have $\int_{J_{u}}|[u]| \wedge \lambda d \mathcal{H}^{d-1} \leq$ $\int_{J_{u}}|[u]| \wedge 1 d \mathcal{H}^{d-1}<+\infty$. If $\int_{J_{u}}|[u]| \wedge 1 d \mathcal{H}^{d-1}<+\infty$ and $\lambda>1$ we have $\int_{J_{u}}|[u]| \wedge \lambda d \mathcal{H}^{d-1} \leq \int_{J_{u} \backslash J_{u}^{1}}|[u]| d \mathcal{H}^{d-1}+\lambda \mathcal{H}^{d-1}\left(J_{u}^{1}\right) \leq \lambda \int_{J_{u}}|[u]| \wedge 1 d \mathcal{H}^{d-1}<$ $+\infty$. The converse implication can be proved in the same way.

Given a function $u \in G B V_{\star}(U)$, by Theorem 2.7 and Proposition 3.3 the measures $\nabla u \mathcal{L}^{d}$ and $D^{c} u$ are well-defined and belong to $\mathcal{M}_{b}\left(U ; \mathbb{R}^{d}\right)$. Since in general $|[u]| \notin L^{1}\left(J_{u}, \mathcal{H}^{d-1}\right)$, we cannot use (2.5) and (2.7) to define a measure which plays the role of $D u$. However, this is possible on a suitable subset of $U$ and leads to a measure which will be crucial in the sequel.

Definition 3.6. Let $u \in G B V_{\star}(U)$ and let $\Gamma \subset U$ be a Borel set with

$$
\begin{equation*}
\int_{J_{u} \backslash \Gamma}|[u]| d \mathcal{H}^{d-1}<+\infty \tag{3.11}
\end{equation*}
$$

The measure $D^{\Gamma} u \in \mathcal{M}_{b}\left(U \backslash \Gamma ; \mathbb{R}^{d}\right)$ is defined by

$$
\begin{equation*}
D^{\Gamma} u:=\nabla u \mathcal{L}^{d}+D^{c} u+[u] \nu_{u} \mathcal{H}^{d-1}\left\llcorner\left(J_{u} \backslash \Gamma\right) .\right. \tag{3.12}
\end{equation*}
$$

Remark 3.7. If $u \in B V(U)$, using (2.5) and (2.7) we see that $D^{\Gamma} u$ coincides with the restriction of $D u$ to $U \backslash \Gamma$.

It is known that, if $u \in B V(U)$, then the approximate limit $\tilde{u}(x)$ is finite for $\mathcal{H}^{d-1}$-a.e. $x \in U \backslash J_{u}$, while $u^{+}(x)$ and $u^{-}(x)$ are finite for $\mathcal{H}^{d-1}$-a.e. $x \in J_{u}$ (see [7, Theorem 5.9.3]). These properties do not hold for an arbitrary function in $G B V(U)$. The following theorem shows that they hold for functions in $G B V_{\star}(U)$.

Theorem 3.8. Let $u \in G B V_{\star}(U)$. Then $\tilde{u}(x)$ is finite for $\mathcal{H}^{d-1}$-a.e. $x \in U \backslash J_{u}$, while $u^{+}(x)$ and $u^{-}(x)$ are finite for $\mathcal{H}^{d-1}$-a.e. $x \in J_{u}$.

Proof. It is enough to repeat the proof of [2, Theorem 4.40], replacing the hypothesis $\mathcal{H}^{d-1}\left(S_{u}^{*}\right)<+\infty$ with (3.5).

It is well-known that $G B V(U)$ is not a vector space (see [2, Remark 4.27]). The additional properties considered in the definition of $G B V_{\star}(U)$ lead to the following result.

Theorem 3.9. $G B V_{\star}(U)$ is a vector space.
Proof. It is obvious that, if $u \in G B V_{\star}(U)$ and $\lambda \in \mathbb{R}$, then $\lambda u \in G B V_{\star}(U)$. Given $u, v \in G B V_{\star}(U)$ we want to prove that $u+v \in G B V_{\star}(U)$. The first step is to show that, given $m \in \mathbb{R}_{+}$, we have $(u+v)^{(m)} \in B V(U)$. For every $r \in \mathbb{R}_{+}$, by the definition of $G B V_{\star}(U)$ the functions $u^{(r)}$ and $v^{(r)}$ belong to $B V(U)$ hence

$$
\begin{equation*}
w_{r}^{m}:=\left(u^{(r)}+v^{(r)}\right)^{(m)} \in B V(U) . \tag{3.13}
\end{equation*}
$$

By Theorem 2.1 we have

$$
\begin{align*}
\nabla w_{r}^{m} & =\chi_{\left\{\left|u^{(r)}+v^{(r)}\right| \leq m\right\}}\left(\nabla u^{(r)}+\nabla v^{(r)}\right)  \tag{3.14}\\
D^{c} w_{r}^{m} & =\chi_{\left\{\left|\tilde{u}^{(r)}+\tilde{v}^{(r)}\right| \leq m\right\}}\left(D^{c} u^{(r)}+D^{c} v^{(r)}\right) \tag{3.15}
\end{align*}
$$

To estimate $\left|D w_{r}^{m}\right|(U)$ we write

$$
\left|D w_{r}^{m}\right|(U)=\int_{U}\left|\nabla w_{r}^{m}\right| d x+\left|D^{c} w_{r}^{m}\right|(U)+\int_{J_{w_{r}^{m}}}\left|\left[w_{r}^{m}\right]\right| d \mathcal{H}^{d-1}
$$

Taking into account the definition of $G B V_{\star}(U)$, by (3.14), (3.15) there exists a constant $C>0$, independent of $m$, such that

$$
\begin{align*}
\int_{U}\left|\nabla w_{r}^{m}\right| d x & \leq \int_{U}\left|\nabla u^{(r)}\right| d x+\int_{U}\left|\nabla v^{(r)}\right| d x \leq C,  \tag{3.16}\\
\left|D^{c} w_{r}^{m}\right|(U) & \leq\left|D^{c} u^{(r)}\right|(U)+\left|D^{c} v^{(r)}\right|(U) \leq C, \tag{3.17}
\end{align*}
$$

for every $r \in \mathbb{R}_{+}$. Since by (3.13), $\left|\left[w_{r}^{m}\right]\right| \leq\left|\left[u^{(r)}\right]\right|+\left|\left[v^{(r)}\right]\right| \mathcal{H}^{d-1}$-a.e. on $J_{w_{r}^{m}}$ and $J_{w_{r}^{m}} \subset J_{u^{(r)}} \cup J_{v^{(r)}}$ up to a set of $\mathcal{H}^{d-1}$-measure zero, recalling also that $\left|w_{r}^{m}\right| \leq m$, we have

$$
\begin{aligned}
& \int_{J_{w_{r}^{m}}}\left|\left[w_{r}^{m}\right]\right| d \mathcal{H}^{d-1} \leq \int_{J_{u^{(r)}} \backslash J_{u}^{1}(r)}\left|\left[u^{(r)}\right]\right| d \mathcal{H}^{d-1}+\int_{J_{v^{(r)} \backslash J_{v}^{(r)}}}\left|\left[v^{(r)}\right]\right| d \mathcal{H}^{d-1} \\
& \quad+2 m \mathcal{H}^{d-1}\left(J_{u^{(r)}}^{1}\right)+2 m \mathcal{H}^{d-1}\left(J_{v^{(r)}}^{1}\right) \\
& \leq(2 m+1) \int_{J_{u^{(r)}}}\left|\left[u^{(r)}\right]\right| \wedge 1 d \mathcal{H}^{d-1}+(2 m+1) \int_{J_{v(r)}}\left|\left[v^{(r)}\right]\right| \wedge 1 d \mathcal{H}^{d-1} .
\end{aligned}
$$

Using the definition of $G B V_{\star}(U)$ we see that there exists a constant $C_{m}>0$ such that

$$
\begin{equation*}
\int_{J_{w_{r}^{m}}}\left|\left[w_{r}^{m}\right]\right| d \mathcal{H}^{d-1} \leq C_{m} \tag{3.18}
\end{equation*}
$$

for every $r \in \mathbb{R}_{+}$. Since $w_{r}^{m} \rightarrow(u+v)^{(m)} \in L^{1}(U)$ as $r \rightarrow \infty$, by (2.5) and (3.16)-(3.18) we deduce that $(u+v)^{(m)} \in B V(U)$.

To conclude the proof we have to show that

$$
\begin{align*}
& \sup _{m \in \mathbb{R}_{+}} \int_{U}\left|\nabla(u+v)^{(m)}\right| d x<+\infty,  \tag{3.19}\\
& \sup _{m \in \mathbb{R}_{+}}\left|D^{c}(u+v)^{(m)}\right|(U)<+\infty,  \tag{3.20}\\
& \sup _{m \in \mathbb{R}_{+}} \int_{J_{(u+v)}^{(m)}}\left|\left[(u+v)^{(m)}\right]\right| \wedge 1 d \mathcal{H}^{d-1}<+\infty . \tag{3.21}
\end{align*}
$$

To prove (3.19) for every $r \geq 0$ we set $A_{r}:=\{|u| \leq r\} \cap\{|v| \leq r\}$. Since $\left(u^{(r)}+v^{(r)}\right)^{(m)}=(u+v)^{(m)}$ in $A_{r}$, by Lemma 2.3 we have $\nabla\left(u^{(r)}+v^{(r)}\right)^{(m)}=$ $\nabla(u+v)^{(m)} \mathcal{L}^{d}$-a.e. in $A_{r}$. Recalling (3.16) we obtain

$$
\begin{aligned}
& \int_{A_{r}}\left|\nabla(u+v)^{(m)}\right| d x=\int_{A_{r}}\left|\nabla\left(u^{(r)}+v^{(r)}\right)^{(m)}\right| d x \\
& \quad \leq \int_{U}\left|\nabla\left(u^{(r)}+v^{(r)}\right)^{(m)}\right| d x \leq C
\end{aligned}
$$

and since $A_{r} \nearrow U$ as $r \rightarrow+\infty$ (recall that $u$ and $v$ have finite values) we get (3.19).

To prove (3.20) for every $r \geq 0$ we set, according to (2.4),

$$
\begin{equation*}
\left.\tilde{A}_{r}:=\{|\tilde{u}| \leq r)\right\} \cap\{|\tilde{v}| \leq r\} . \tag{3.22}
\end{equation*}
$$

By (2.2) applied to $f(t):=t^{(m)}$ we have $\tilde{u}(x)=\operatorname{ap}_{y \rightarrow x} \lim u^{(r)}(y)$ for every $x \in \tilde{A}_{r}$. Hence $\left.\widetilde{\left(u^{(r)}\right.}+\widetilde{v^{(r)}}\right)^{(m)}=(\tilde{u}+\tilde{v})^{(m)}$ in $\tilde{A}_{r}$. By Lemma 2.3 we have $D^{c}\left(u^{(r)}+\right.$
$\left.v^{(r)}\right)^{(m)}=D^{c}(u+v)^{(m)}$ as measures on $\tilde{A}_{r}$ and from (3.17) we get
$\left|D^{c}(u+v)^{(m)}\right|\left(\tilde{A}_{r}\right)=\left|D^{c}\left(u^{(r)}+v^{(r)}\right)^{(m)}\right|\left(\tilde{A}_{r}\right) \leq\left|D^{c}\left(u^{(r)}+v^{(r)}\right)^{(m)}\right|(U) \leq C$.
Recalling that $\tilde{A}_{r} \nearrow U \backslash(\{|\tilde{u}|=+\infty\} \cup\{|\tilde{v}|=+\infty\})$ and that $\mathcal{H}^{d-1}(\{|\tilde{u}|=$ $+\infty\} \cup\{|\tilde{v}|=+\infty\})=0$ by Theorem 3.8, we can pass to the limit as $r \rightarrow+\infty$ and we obtain that (3.20) holds.

It remains to prove (3.21). To this end we observe that by Proposition 3.3 there exists a constant $C>0$ such that

$$
\int_{J_{u}}|[u]| \wedge 1 d \mathcal{H}^{d-1} \leq C \quad \text { and } \quad \int_{J_{v}}|[v]| \wedge 1 d \mathcal{H}^{d-1} \leq C
$$

Since $\left|\left[(u+v)^{(m)}\right]\right| \wedge 1 \leq|[u+v]| \wedge 1 \leq|[u]| \wedge 1+|[v]| \wedge 1$, we have that

$$
\int_{J_{(u+v)(m)}}\left|\left[(u+v)^{(m)}\right]\right| \wedge 1 d \mathcal{H}^{d-1} \leq 2 C
$$

for every $m$. This concludes the proof.
Proposition 3.10. Let $u, v \in G B V_{\star}(U), \lambda \in \mathbb{R}$, and let $\Gamma \subset U$ be a Borel set. Then

$$
\begin{align*}
\nabla(u+v) & =\nabla u+\nabla v \quad \text { and } \quad \nabla(\lambda u)=\lambda \nabla u \quad \mathcal{L}^{d} \text {-a.e. in } U,  \tag{3.23}\\
D^{c}(u+v) & =D^{c} u+D^{c} v \quad \text { and } \quad D^{c}(\lambda u)=\lambda D^{c} u \quad \text { on } U,  \tag{3.24}\\
D^{\Gamma}(u+v) & =D^{\Gamma} u+D^{\Gamma} v \quad \text { and } \quad D^{\Gamma}(\lambda u)=\lambda D^{\Gamma} u \quad \text { on } U \backslash \Gamma . \tag{3.25}
\end{align*}
$$

Proof. Equalities (3.23) follow immediately from the definition of the approximate limit. The second equality in (3.24) follows from the definition of $D^{c} u$. To prove the first equality, we set $w:=u+v$ and we fix $m$ and $s$ with $0 \leq 2 s \leq m$. For every $r \geq s$ we have $\tilde{w}^{(m)}=\tilde{w}=\tilde{u}+\tilde{v}=\tilde{u}^{(r)}+\tilde{v}^{(r)} \mathcal{H}^{d-1}$-a.e. in $\tilde{A}_{s}$, where $\tilde{A}_{s}$, is defined in (3.22). Since $w^{(m)} \in B V(U)$ by Theorem 3.9 and $u^{(r)}$, $v^{(r)} \in B V(U)$ by the definition of $G B V_{\star}(U)$, using (2.2) and Lemma 2.3 for every Borel set $B$ in $U$ we obtain
$D^{c} w^{(m)}\left(B \cap \tilde{A}_{s}\right)=D^{c}\left(u^{(r)}+v^{(r)}\right)\left(B \cap \tilde{A}_{s}\right)=D^{c} u^{(r)}\left(B \cap \tilde{A}_{s}\right)+D^{c} v^{(r)}\left(B \cap \tilde{A}_{s}\right)$.
By Proposition 2.9 we can pass to the limit as $r \rightarrow+\infty$ and we get

$$
D^{c}(u+v)^{(m)}\left(B \cap \tilde{A}_{s}\right)=D^{c} u\left(B \cap \tilde{A}_{s}\right)+D^{c} v\left(B \cap \tilde{A}_{s}\right) .
$$

Taking the limit as $m \rightarrow+\infty$ and using Proposition 2.9 again we obtain

$$
D^{c}(u+v)\left(B \cap \tilde{A}_{s}\right)=D^{c} u\left(B \cap \tilde{A}_{s}\right)+D^{c} v\left(B \cap \tilde{A}_{s}\right) .
$$

Finally, arguing as in the proof of Theorem 3.9 we can pass to the limit as $s \rightarrow+\infty$ and obtain the first equality in (3.24).

By (3.23) and (3.24) to prove the first equality in (3.25) it is enough to show that

$$
\int_{J_{u+v} \cap B}[u+v] \nu_{u+v} d \mathcal{H}^{d-1}=\int_{J_{u} \cap B}[u] \nu_{u} d \mathcal{H}^{d-1}+\int_{J_{v} \cap B}[v] \nu_{v} d \mathcal{H}^{d-1}
$$

for every Borel set $B \subset U \backslash \Gamma$. This follows easily from the linearity of the jump and the locality property of approximate tangent spaces (see, e.g., [2, (2.65)]),
which gives (up to a sign) $\nu_{u+v}=\nu_{u} \mathcal{H}^{d-1}$-a.e. on $J_{u+v} \cap J_{u}$ and $\nu_{u+v}=\nu_{v}$ $\mathcal{H}^{d-1}$-a.e. on $J_{u+v} \cap J_{v}$. The second equality in (3.25) is trivial.

Theorem 3.11. (Compactness) Let $\left(u_{k}\right)_{k}$ be a sequence in $G B V_{\star}(U)$. Assume that there exist a constant $M>0$ and a continuous function $\psi: \mathbb{R}_{+} \rightarrow \mathbb{R}_{+}$ with $\psi(t) \rightarrow+\infty$ as $t \rightarrow+\infty$, such that

$$
\begin{align*}
& \int_{U}\left|\nabla u_{k}\right| d x+\left|D^{c} u_{k}\right|(U)+\int_{J_{u_{k}}}\left|\left[u_{k}\right]\right| \wedge 1 d \mathcal{H}^{d-1} \leq M  \tag{3.26}\\
& \sup _{k} \int_{U} \psi\left(\left|u_{k}\right|\right) d x<+\infty \tag{3.27}
\end{align*}
$$

Then there exist a subsequence, not relabelled, and a function $u \in G B V_{\star}(U)$ such that $u_{k} \rightarrow u \mathcal{L}^{d}$-a.e. in $U$.

Proof. We claim that for every $m \in \mathbb{N}$ the truncated functions $u_{k}^{(m)}$ are bounded in $B V(U)$. Indeed, by Proposition 2.6(b) we have that $\nabla u_{k}^{(m)}=\nabla u_{k} \mathcal{L}^{d}$-a.e. on $\left\{\left|u_{k}\right| \leq m\right\}$, while by Remark 2.5 we have $\nabla u_{k}^{(m)}=0 \mathcal{L}^{d}$-a.e. on $\left\{\left|u_{k}\right|>m\right\}$. By (3.26) this implies that

$$
\begin{equation*}
\int_{U}\left|\nabla u_{k}^{(m)}\right| d x \leq \int_{U}\left|\nabla u_{k}\right| d x \leq M . \tag{3.28}
\end{equation*}
$$

By (2.23) we have also

$$
\begin{equation*}
\left|D^{c} u_{k}^{(m)}\right|(U) \leq\left|D^{c} u_{k}\right|(U) \leq M . \tag{3.29}
\end{equation*}
$$

As for the estimate on the jump part, we observe that, by Proposition $2.6(\mathrm{c})$, we have that $J_{u_{k}^{(m)}} \subset J_{u_{k}}$ up to a set of $\mathcal{H}^{d-1}$-measure zero, and $\left|\left[u_{k}^{(m)}\right]\right| \leq\left|\left[u_{k}\right]\right| \mathcal{H}^{d-1}$-a.e. on $J_{u_{k}^{(m)}} \cap J_{u_{k}}$. Then

$$
\begin{align*}
& \int_{J_{u_{k}^{(m)}}}\left|\left[u_{k}^{(m)}\right]\right| \wedge 1 d \mathcal{H}^{d-1} \leq \int_{J_{u_{k}}}\left|\left[u_{k}\right]\right| \wedge 1 d \mathcal{H}^{d-1}  \tag{3.30}\\
& \int_{J_{u_{k}^{(m)}}}\left|\left[u_{k}^{(m)}\right]\right| d \mathcal{H}^{d-1} \leq \int_{J_{u_{k}^{(m)}} \backslash J_{u_{k}^{(m)}}^{1}}\left|\left[u_{k}^{(m)}\right]\right| d \mathcal{H}^{d-1}+2 m \mathcal{H}^{d-1}\left(J_{u_{k}^{(m)}}^{1}\right) \\
& \leq(1+2 m) \int_{J_{u_{k}}}\left|\left[u_{k}\right]\right| \wedge 1 d \mathcal{H}^{d-1} \leq(1+2 m) M \tag{3.31}
\end{align*}
$$

Therefore, by (2.5), (3.28), (3.29), and (3.31), the functions $u_{k}^{(m)}$ are bounded in $B V(U)$ uniformly with respect to $k$.

By the compactness of the embedding of $B V(U)$ into $L_{l o c}^{1}(U)$, using a diagonal argument we can extract a subsequence of $\left(u_{k}\right)_{k}$, not relabelled, such that for every $m \in \mathbb{N}$, the sequence $\left(u_{k}^{(m)}\right)_{k}$ converges $\mathcal{L}^{d}$-a.e. in $U$ to a function $v^{m} \in L^{\infty}(U)$. Since the $B V$-norm is lower semicontinuous with respect to $L^{1}$ convergence, we obtain that $v^{m} \in B V(U)$.

We observe that

$$
\begin{equation*}
m<n \quad \Longrightarrow \quad\left(v^{n}\right)^{(m)}=v^{m} \tag{3.32}
\end{equation*}
$$

This is an obvious consequence of the fact that $\left(u_{k}^{(n)}\right)^{(m)}=u_{k}^{(m)}$. From (3.32) we have that

$$
m<n \quad \Longrightarrow \quad\left\{\left|v^{n}\right|=n\right\} \subset\left\{\left|v^{m}\right|=m\right\}
$$

Let $E_{\infty}$ be the intersection of the sets $\left\{\left|v^{m}\right|=m\right\}$ for $m \in \mathbb{N}$. We claim that

$$
\begin{equation*}
\mathcal{L}^{d}\left(E_{\infty}\right)=0 . \tag{3.33}
\end{equation*}
$$

To prove this property we observe that it is not restrictive to assume that the function $\psi$ in (3.27) is increasing. For every $m \in \mathbb{N}$ by the Fatou Lemma we have
$\psi(m) \mathcal{L}^{d}\left(E_{\infty}\right)=\int_{E_{\infty}} \psi\left(\left|v^{m}\right|\right) d x \leq \liminf _{k} \int_{E_{\infty}} \psi\left(\left|u_{k}^{(m)}\right|\right) d x \leq \sup _{k} \int_{U} \psi\left(\left|u_{k}\right|\right) d x$,
where in the last inequality we used the monotonicity of $\psi$. Since, by assumption, $\psi(m) \rightarrow+\infty$ as $m \rightarrow+\infty$, from (3.27) we obtain (3.33).

If $x \in U \backslash E_{\infty}$ there exists $m \in \mathbb{N}$ such that $\left|v^{m}(x)\right|<m$. We set

$$
u(x):=v^{m}(x)
$$

and we observe that, by (3.32), the function $u$ is well-defined on $U \backslash E_{\infty}$ and $u^{(m)}=v^{m}$ in $U \backslash E_{\infty}$ for every $m \in \mathbb{N}$. We also set $u(x)=0$ for $x \in E_{\infty}$. Since $u^{(m)}=v^{m} \mathcal{L}^{d}$-a.e. on $U$ we conclude that $u^{(m)} \in B V(U)$ for every $m \in \mathbb{N}$ and $u_{k}^{(m)} \rightarrow u^{(m)}$ strongly in $L^{1}(U)$ as $k \rightarrow+\infty$. By (3.26), (3.28), (3.29), and (3.30) we obtain that

$$
\begin{equation*}
\int_{U}\left|\nabla u_{k}^{(m)}\right| d x+\left|D^{c} u_{k}^{(m)}\right|(U)+\int_{J_{u_{k}^{(m)}}}\left|\left[u_{k}^{(m)}\right]\right| \wedge 1 d \mathcal{H}^{d-1} \leq M . \tag{3.34}
\end{equation*}
$$

By [4, Theorem 2.1] we deduce that

$$
\begin{equation*}
\int_{U}\left|\nabla u^{(m)}\right| d x+\left|D^{c} u^{(m)}\right|(U)+\int_{J_{u^{(m)}}}\left|\left[u^{(m)}\right]\right| \wedge 1 d \mathcal{H}^{d-1} \leq M \tag{3.35}
\end{equation*}
$$

hence $u \in G B V_{\star}(U)$.

## 4. The incremental minimum problem

In this section we present a precise formulation of the incremental minimum problem (1.1), which appears in the variational approach to the quasistatic crack growth in elastic-perfectly plastic materials. The reference configuration is a bounded open set $\Omega \subset \mathbb{R}^{d}$ with Lipschitz boundary. The crack in the reference configuration is represented by a Borel set $\Gamma \subset \bar{\Omega}$, with $\mathcal{H}^{d-1}(\Gamma)<$ $+\infty$. The set $\bar{\Omega} \backslash \Gamma$ represents the elasto-plastic part of the body.

Since we are dealing with the antiplane case, the displacement of each material point is described by a function $u: \bar{\Omega} \backslash \Gamma \rightarrow \mathbb{R}$. Regarding $u$ as a function defined $\mathcal{L}^{d}$-a.e. in $\Omega$, we assume that

$$
\begin{equation*}
u \in G B V_{\star}(\Omega), \int_{J_{u} \backslash \Gamma}|[u]| d \mathcal{H}^{d-1}<+\infty, \text { and } \int_{\partial \Omega \backslash \Gamma}|u| d \mathcal{H}^{d-1}<+\infty . \tag{4.1}
\end{equation*}
$$

Here and in the rest of the paper the trace on $\partial \Omega$ of a function $v \in G B V(\Omega)$ is still denoted by $v$. The strain corresponding to the displacement $u$ is given by the measure $D^{\Gamma} u \in \mathcal{M}_{b}\left(\Omega \backslash \Gamma ; \mathbb{R}^{d}\right)$ introduced in Definition 3.6 with $U$ replaced by $\Omega$ and $\Gamma$ replaced by $\Gamma \cap \Omega$.

The Dirichlet boundary condition is assigned using the trace on $\partial \Omega$ of a function $w \in H^{1}(\Omega)$. The elastic part of the strain is denoted by $e$ and the plastic part by $p$. We assume that

$$
\begin{align*}
& e \in L^{2}\left(\Omega ; \mathbb{R}^{d}\right) \quad \text { and } \quad p \in \mathcal{M}_{b}\left(\bar{\Omega} \backslash \Gamma ; \mathbb{R}^{d}\right),  \tag{4.2}\\
& D^{\Gamma} u=e+p \text { as measures on } \Omega \backslash \Gamma  \tag{4.3}\\
& p=(w-u) \nu_{\Omega} \mathcal{H}^{d-1} \text { as measures on } \partial \Omega \backslash \Gamma, \tag{4.4}
\end{align*}
$$

where $\nu_{\Omega}$ is the outer unit normal to $\partial \Omega$. Here and in the rest of the paper we identify an $L^{1}$-function $\varphi$ and the measure $\varphi \mathcal{L}^{d}$. To simplify the exposition, for every Borel set $\Gamma \subset \bar{\Omega}$ and every $w \in H^{1}(\Omega)$, it is convenient to introduce the set $\mathcal{A}(\Gamma, w)$ of all triples $(u, e, p)$ which satisfy (4.1)-(4.4).

From the definition of $D^{\Gamma} u$ it follows that if $\mathcal{H}^{d-1}(\Gamma)<+\infty$, the absolutely continuous part $p^{a}$ of $p$ with respect to $\mathcal{L}^{d}$ satisfies

$$
\begin{equation*}
\nabla u=e+p^{a} \quad \mathcal{L}^{d} \text {-a.e. in } \Omega \tag{4.5}
\end{equation*}
$$

while the singular part $p^{s}$ of $p$ with respect to $\mathcal{L}^{d}$ satisfies

$$
\begin{align*}
& p^{s}(B)=D^{c} u(B)+\int_{J_{u} \cap B}[u] \nu_{u} d \mathcal{H}^{d-1} \text { for every Borel set } B \subset \Omega \backslash \Gamma,  \tag{4.6}\\
& p^{s}(B)=p(B)=\int_{B}(w-u) \nu_{\Omega} d \mathcal{H}^{d-1} \text { for every Borel set } B \subset \partial \Omega \backslash \Gamma . \tag{4.7}
\end{align*}
$$

In our incremental minimum problem the data at the previous time are

$$
\begin{align*}
& \text { a Borel set } \Gamma_{0} \subset \bar{\Omega} \quad \text { with } \mathcal{H}^{d-1}\left(\Gamma_{0}\right)<+\infty,  \tag{4.8}\\
& w_{0} \in H^{1}(\Omega) \quad \text { and } \quad\left(u_{0}, e_{0}, p_{0}\right) \in \mathcal{A}\left(\Gamma_{0}, w_{0}\right) . \tag{4.9}
\end{align*}
$$

Given

$$
\begin{equation*}
w \in H^{1}(\Omega) \tag{4.10}
\end{equation*}
$$

the precise formulation of the incremental minimum problem (1.1) is

$$
\begin{equation*}
\min _{\substack{\Gamma \operatorname{Borel}, \Gamma_{\mathcal{O}} \subset \Gamma \subset \bar{\Omega} \\(u, e, p) \in \mathcal{A}(\Gamma, w)}}\left\{\frac{1}{2} \int_{\Omega}|e|^{2} d x+\left|p-p_{0}\right|(\bar{\Omega} \backslash \Gamma)+\mathcal{H}^{d-1}\left(\Gamma \backslash \Gamma_{0}\right)\right\} . \tag{4.11}
\end{equation*}
$$

To solve this problem we introduce the function $f: \mathbb{R}^{d} \rightarrow \mathbb{R}$ defined by

$$
f(\xi):=\min _{\eta \in \mathbb{R}^{d}}\left\{\frac{1}{2}|\eta|^{2}+|\xi-\eta|\right\}= \begin{cases}\frac{1}{2}|\xi|^{2} & \text { if }|\xi| \leq 1  \tag{4.12}\\ |\xi|-\frac{1}{2} & \text { if }|\xi| \geq 1\end{cases}
$$

The minimum in the definition of $f(\xi)$ is attained for $\eta= \begin{cases}\xi & \text { if }|\xi| \leq 1, \\ \xi /|\xi| & \text { if }|\xi| \geq 1 .\end{cases}$ It is convenient to introduce the maps $\pi_{1}, \pi_{2}: \mathbb{R}^{d} \rightarrow \mathbb{R}^{d}$ defined by

$$
\pi_{1}(\xi)=\left\{\begin{array}{ll}
\xi & \text { if }|\xi| \leq 1 \\
\xi /|\xi| & \text { if }|\xi| \geq 1
\end{array} \quad \pi_{2}(\xi)= \begin{cases}0 & \text { if }|\xi| \leq 1 \\
\xi-\xi /|\xi| & \text { if }|\xi| \geq 1\end{cases}\right.
$$

We note that

$$
\begin{equation*}
\xi=\pi_{1}(\xi)+\pi_{2}(\xi), \quad\left|\pi_{1}(\xi)\right| \leq 1, \quad \text { and } \quad f(\xi)=\frac{1}{2}\left|\pi_{1}(\xi)\right|^{2}+\left|\pi_{2}(\xi)\right| \tag{4.13}
\end{equation*}
$$

For later use we observe that the definition (4.12) of $f$ implies

$$
\begin{equation*}
|\xi|-\frac{1}{2} \leq f(\xi) \leq|\xi| \quad \text { for every } \xi \in \mathbb{R}^{d} \tag{4.14}
\end{equation*}
$$

To deal with the boundary condition (4.4) in (4.11) it is convenient to introduce a bounded open set $\Omega^{\prime}$ with

$$
\begin{equation*}
\bar{\Omega} \subset \Omega^{\prime} \tag{4.15}
\end{equation*}
$$

and to extend $w, w_{0}, e_{0}$ in such a way that

$$
\begin{equation*}
w, w_{0} \in H^{1}\left(\Omega^{\prime}\right) \quad \text { and } \quad e_{0} \in L^{2}\left(\Omega^{\prime} ; \mathbb{R}^{d}\right) \tag{4.16}
\end{equation*}
$$

We now prove that problem (4.11) is equivalent to the following minimum problem

$$
\begin{align*}
& \min _{\substack{\Gamma \text { Borel, } \Gamma_{0} \subset \Gamma \subset \bar{\Omega} \\
v \in G B V_{\star}\left(\Omega^{\prime}\right)}}\left\{\int_{\Omega^{\prime}} f\left(\nabla v+e_{0}\right) d x+\left|D^{c} v\right|\left(\Omega^{\prime}\right)+\int_{J_{v} \backslash \Gamma}|[v]| d \mathcal{H}^{d-1}+\mathcal{H}^{d-1}(\Gamma)\right\} \\
& v=w-w_{0} \text { a.e. in } \Omega^{\prime} \backslash \Omega
\end{align*}
$$

Lemma 4.1. Assume (4.8), (4.9), (4.15), and (4.16). Let $\Gamma$ and ( $u, e, p$ ) be a solution of (4.11) and let $v:=u-u_{0}$ in $\Omega$ and $v:=w-w_{0}$ in $\Omega^{\prime} \backslash \Omega$. Then $\Gamma$ and $v$ solve (4.17).

Conversely, assume that $\Gamma$ and $v$ solve (4.17) and let $u:=\left.v\right|_{\Omega}+u_{0}$, $e:=\pi_{1}\left(\left.\nabla v\right|_{\Omega}+\left.e_{0}\right|_{\Omega}\right), p:=D^{\Gamma} u-e$ in $\Omega \backslash \Gamma$, and $p:=(w-u) \nu_{\Omega} \mathcal{H}^{d-1}$ on $\partial \Omega \backslash \Gamma$. Then $\Gamma$ and ( $u, e, p$ ) solve (4.11).

Proof. Let $\Gamma$ and $(u, e, p)$ be a solution of (4.11). It is clear that $\mathcal{H}^{d-1}\left(\Gamma \backslash \Gamma_{0}\right)<$ $+\infty$, hence (4.8) implies that $\mathcal{H}^{d-1}(\Gamma)<+\infty$. Let $v$ be as in the statement of the lemma. To prove that $\Gamma$ and $v$ solve (4.17) we fix a Borel set $\hat{\Gamma}$, with $\Gamma_{0} \subset \hat{\Gamma} \subset \bar{\Omega}$, and $\hat{v} \in G B V_{\star}\left(\Omega^{\prime}\right)$, with

$$
\begin{equation*}
\hat{v}=w-w_{0} \quad \mathcal{L}^{d} \text {-a.e. in } \Omega^{\prime} \backslash \Omega . \tag{4.18}
\end{equation*}
$$

We want to show that

$$
\begin{align*}
& \int_{\Omega^{\prime}} f\left(\nabla v+e_{0}\right) d x+\left|D^{c} v\right|\left(\Omega^{\prime}\right)+\int_{J_{v} \backslash \Gamma}|[v]| d \mathcal{H}^{d-1}+\mathcal{H}^{d-1}(\Gamma) \\
& \quad \leq \int_{\Omega^{\prime}} f\left(\nabla \hat{v}+e_{0}\right) d x+\left|D^{c} \hat{v}\right|\left(\Omega^{\prime}\right)+\int_{J_{\hat{v}} \backslash \hat{\Gamma}}|[\hat{v}]| d \mathcal{H}^{d-1}+\mathcal{H}^{d-1}(\hat{\Gamma}) . \tag{4.19}
\end{align*}
$$

It is not restrictive to assume that

$$
\begin{equation*}
\int_{J_{\hat{v}} \backslash \hat{\Gamma}}|[\hat{v}]| d \mathcal{H}^{d-1}<+\infty \quad \text { and } \quad \mathcal{H}^{d-1}(\hat{\Gamma})<+\infty \tag{4.20}
\end{equation*}
$$

We set

$$
\begin{equation*}
\hat{e}:=\pi_{1}\left(\left.\nabla \hat{v}\right|_{\Omega}+\left.e_{0}\right|_{\Omega}\right) \text { and } \hat{g}:=\left.\nabla \hat{v}\right|_{\Omega}+\nabla u_{0}-\hat{e}=\left.\nabla \hat{v}\right|_{\Omega}+\left.e_{0}\right|_{\Omega}+p_{0}^{a}-\hat{e}, \tag{4.21}
\end{equation*}
$$

where the last equality follows from (4.5) and (4.9). Then $\hat{e} \in L^{2}\left(\Omega ; \mathbb{R}^{d}\right)$ and $\hat{g} \in L^{1}\left(\Omega ; \mathbb{R}^{d}\right)$. We now define $\hat{u}:=\left.\hat{v}\right|_{\Omega}+u_{0}$ and note that $\hat{u} \in G B V_{\star}(\Omega)$ by Theorem 3.9. Moreover we define $\hat{p} \in \mathcal{M}_{b}\left(\bar{\Omega} \backslash \hat{\Gamma} ; \mathbb{R}^{d}\right)$ by

$$
\hat{p}:=\hat{g}+D^{c} \hat{u}+[\hat{u}] \nu_{\hat{u}} \mathcal{H}^{d-1}\left\llcorner\left(J_{\hat{u}} \backslash \hat{\Gamma}\right)+(w-\hat{u}) \nu_{\Omega} \mathcal{H}^{d-1}\llcorner(\partial \Omega \backslash \hat{\Gamma}) .\right.
$$

We remark that $\int_{J_{\hat{u}} \backslash \hat{\Gamma}}|[\hat{u}]| d \mathcal{H}^{d-1}<+\infty$ and $\int_{\partial \Omega \backslash \hat{\Gamma}}|\hat{u}| d \mathcal{H}^{d-1}<+\infty$ by (4.1), (4.9), (4.18), and (4.20). This shows that the definition of $\hat{p}$ makes sense. We note that $D^{\hat{\Gamma}} \hat{u}=\hat{e}+\hat{p}$ in $\Omega \backslash \hat{\Gamma}$ and $\hat{p}=(w-\hat{u}) \nu_{\Omega} \mathcal{H}^{d-1}$ on $\partial \Omega \backslash \hat{\Gamma}$, hence $(\hat{u}, \hat{e}, \hat{p}) \in \mathcal{A}(\hat{\Gamma}, w)$. Consequently, the minimality of $\Gamma$ and $(u, e, p)$ gives

$$
\begin{align*}
& \frac{1}{2} \int_{\Omega}|e|^{2} d x+\left|p-p_{0}\right|(\bar{\Omega} \backslash \Gamma)+\mathcal{H}^{d-1}\left(\Gamma \backslash \Gamma_{0}\right) \\
& \quad \leq \frac{1}{2} \int_{\Omega}|\hat{e}|^{2} d x+\left|\hat{p}-p_{0}\right|(\bar{\Omega} \backslash \hat{\Gamma})+\mathcal{H}^{d-1}\left(\hat{\Gamma} \backslash \Gamma_{0}\right) . \tag{4.22}
\end{align*}
$$

Since $\hat{v}=w-w_{0} \mathcal{L}^{d}$-a.e. in $\Omega^{\prime} \backslash \Omega$ and $\hat{u}=\left.\hat{v}\right|_{\Omega}+u_{0}$ we have that
$\int_{J_{\hat{v}} \backslash \hat{\Gamma}}|[\hat{v}]| d \mathcal{H}^{d-1}=\int_{J_{\hat{u}-u_{0}} \backslash \hat{\Gamma}}\left|\left[\hat{u}-u_{0}\right]\right| d \mathcal{H}^{d-1}+\int_{\partial \Omega \backslash \hat{\Gamma}}\left|w-\hat{u}-w_{0}+u_{0}\right| d \mathcal{H}^{d-1}$.
On the other hand by the definition of $\hat{p}$ we have

$$
\begin{aligned}
\mid \hat{p}- & p_{0}\left|(\bar{\Omega} \backslash \hat{\Gamma})=\int_{\Omega}\right| \hat{g}-p_{0}^{a}\left|d x+\left|D^{c}\left(\hat{u}-u_{0}\right)\right|(\Omega)\right. \\
& +\int_{J_{\hat{u}-u_{0}} \backslash \hat{\Gamma}}\left|\left[\hat{u}-u_{0}\right]\right| d \mathcal{H}^{d-1}+\int_{\partial \Omega \backslash \hat{\Gamma}}\left|w-\hat{u}-w_{0}+u_{0}\right| d \mathcal{H}^{d-1} \\
= & \int_{\Omega}\left|\hat{g}-p_{0}^{a}\right| d x+\left|D^{c} \hat{v}\right|(\Omega)+\int_{J_{\hat{v}} \backslash \hat{\Gamma}}|[\hat{v}]| d \mathcal{H}^{d-1}
\end{aligned}
$$

hence, by (4.13) and (4.21),

$$
\begin{align*}
& \int_{\Omega^{\prime}} f\left(\nabla \hat{v}+e_{0}\right) d x+\left|D^{c} \hat{v}\right|\left(\Omega^{\prime}\right)+\int_{J_{\hat{v}} \backslash \hat{\Gamma}}|[\hat{v}]| d \mathcal{H}^{d-1}+\mathcal{H}^{d-1}(\hat{\Gamma}) \\
& \quad=\gamma+\frac{1}{2} \int_{\Omega}|\hat{e}|^{2} d x+\left|\hat{p}-p_{0}\right|(\bar{\Omega} \backslash \hat{\Gamma})+\mathcal{H}^{d-1}(\hat{\Gamma}), \tag{4.23}
\end{align*}
$$

where

$$
\begin{equation*}
\gamma:=\int_{\Omega^{\prime} \backslash \Omega} f\left(\nabla w-\nabla w_{0}+e_{0}\right) d x . \tag{4.24}
\end{equation*}
$$

Similarly, using the definition (4.12) of $f$ instead of (4.13), we obtain

$$
\begin{align*}
& \int_{\Omega^{\prime}} f\left(\nabla v+e_{0}\right) d x+\left|D^{c} v\right|\left(\Omega^{\prime}\right)+\int_{J_{v} \backslash \Gamma}|[v]| d \mathcal{H}^{d-1}+\mathcal{H}^{d-1}(\Gamma) \\
& \quad \leq \gamma+\frac{1}{2} \int_{\Omega}|e|^{2} d x+\left|p-p_{0}\right|(\bar{\Omega} \backslash \Gamma)+\mathcal{H}^{d-1}(\Gamma) \tag{4.25}
\end{align*}
$$

Then from (4.22), (4.23), and (4.25) we obtain

$$
\begin{align*}
& \int_{\Omega^{\prime}} f\left(\nabla v+e_{0}\right) d x+\left|D^{c} v\right|\left(\Omega^{\prime}\right)+\int_{J_{v} \backslash \Gamma}|[v]| d \mathcal{H}^{d-1}+\mathcal{H}^{d-1}(\Gamma) \\
& \quad \leq \int_{\Omega^{\prime}} f\left(\nabla \hat{v}+e_{0}\right) d x+\left|D^{c} \hat{v}\right|\left(\Omega^{\prime}\right)+\int_{J_{\hat{v}} \backslash \hat{\Gamma}}|[\hat{v}]| d \mathcal{H}^{d-1}+\mathcal{H}^{d-1}(\hat{\Gamma}), \tag{4.26}
\end{align*}
$$

which shows that $\Gamma$ and $v$ solve (4.17).
Conversely, assume that $\Gamma$ and $v$ solve (4.17). By (4.8) it is clear that $\mathcal{H}^{d-1}(\Gamma)<+\infty$. We observe that the triple $(u, e, p)$ defined in the second part of the statement of the lemma belongs to $\mathcal{A}(\Gamma, w)$ and that $p=p^{a}+D^{c} u+$ $[u] \nu_{u} \mathcal{H}^{d-1}\left\llcorner\left(J_{u} \backslash \Gamma\right)\right.$ in $\Omega \backslash \Gamma$ while $p=(w-u) \nu_{\Omega} \mathcal{H}^{d-1}$ on $\partial \Omega \backslash \Gamma$. To prove that $\Gamma$ and $(u, e, p)$ solve (4.11) we fix a Borel set $\hat{\Gamma}$ with $\Gamma_{0} \subset \hat{\Gamma} \subset \bar{\Omega}$ and a triple $(\hat{u}, \hat{e}, \hat{p}) \in \mathcal{A}(\hat{\Gamma}, w)$. We want to show that

$$
\begin{align*}
& \frac{1}{2} \int_{\Omega}|e|^{2} d x+\left|p-p_{0}\right|(\bar{\Omega} \backslash \Gamma)+\mathcal{H}^{d-1}\left(\Gamma \backslash \Gamma_{0}\right) \\
& \quad \leq \frac{1}{2} \int_{\Omega}|\hat{e}|^{2} d x+\left|\hat{p}-p_{0}\right|(\bar{\Omega} \backslash \hat{\Gamma})+\mathcal{H}^{d-1}\left(\hat{\Gamma} \backslash \Gamma_{0}\right) . \tag{4.27}
\end{align*}
$$

Let $\hat{v}:=\hat{u}-u_{0}$ in $\Omega$ and $\hat{v}:=w-w_{0}$ in $\Omega^{\prime} \backslash \Omega$. Then, arguing as in the first part of the proof we obtain in this case (4.25) with an equality and (4.23) with the inequality $\leq$. Then (4.27) follows from (4.22) and (4.26), which is an obvious consequence of (4.17).

We now prove that (4.17) is equivalent to the minimum problem

$$
\begin{equation*}
\min _{\substack{v \in G B V_{\star}\left(\Omega^{\prime}\right) \\ v=w-w_{0} \text { a.e. in } \Omega^{\prime} \backslash \Omega}}\left\{\int_{\Omega^{\prime}} f\left(\nabla v+e_{0}\right) d x+\left|D^{c} v\right|\left(\Omega^{\prime}\right)+\int_{J_{v} \backslash \Gamma_{0}}|[v]| \wedge 1 d \mathcal{H}^{d-1}\right\} . \tag{4.28}
\end{equation*}
$$

We recall that in (3.2) we defined $J_{v}^{1}:=\left\{x \in J_{v}:|[v](x)| \geq 1\right\}$.
Lemma 4.2. If $\Gamma$ and $v$ solve (4.17), then $v$ is a solution of (4.28). Conversely, if $v$ is a solution of (4.28) and $\Gamma:=\left(J_{v}^{1} \cup \Gamma_{0}\right) \cap \bar{\Omega}$, then $\Gamma$ and $v$ solve (4.17).

Proof. Assume that $\Gamma$ and $v$ solve (4.17). Let $\hat{v} \in G B V_{\star}\left(\Omega^{\prime}\right)$ be such that $\hat{v}=w-w_{0} \mathcal{L}^{d}$-a.e. in $\Omega^{\prime} \backslash \Omega$ and let $\hat{\Gamma}=\left(J_{\hat{v}}^{1} \cup \Gamma_{0}\right) \cap \bar{\Omega}$. Then

$$
\int_{J_{v} \backslash \Gamma_{0}}|[v]| \wedge 1 d \mathcal{H}^{d-1} \leq \int_{J_{v} \backslash \Gamma}|[v]| d \mathcal{H}^{d-1}+\mathcal{H}^{d-1}\left(\Gamma \backslash \Gamma_{0}\right)
$$

and

$$
\int_{J_{\hat{v}} \backslash \hat{\Gamma}}|[\hat{v}]| d \mathcal{H}^{d-1}+\mathcal{H}^{d-1}\left(\hat{\Gamma} \backslash \Gamma_{0}\right)=\int_{J_{J_{\hat{v}} \backslash \Gamma_{0}}|[\hat{v}]| \wedge 1 d \mathcal{H}^{d-1}, ., ~ . ~}
$$

where we used the fact that $\mathcal{H}^{d-1}\left(J_{\hat{v}} \backslash \bar{\Omega}\right)=0$ since $\hat{v} \in H^{1}\left(\Omega^{\prime} \backslash \bar{\Omega}\right)$. Therefore, by the minimality of $\Gamma$ and $v$ we have

$$
\begin{aligned}
& \int_{\Omega^{\prime}} f\left(\nabla v+e_{0}\right) d x+\left|D^{c} v\right|\left(\Omega^{\prime}\right)+\int_{J_{v} \backslash \Gamma_{0}}|[v]| \wedge 1 d \mathcal{H}^{d-1} \\
& \quad \leq \int_{\Omega^{\prime}} f\left(\nabla v+e_{0}\right) d x+\left|D^{c} v\right|\left(\Omega^{\prime}\right)+\int_{J_{v} \backslash \Gamma}|[v]| d \mathcal{H}^{d-1}+\mathcal{H}^{d-1}\left(\Gamma \backslash \Gamma_{0}\right) \\
& \quad \leq \int_{\Omega^{\prime}} f\left(\nabla \hat{v}+e_{0}\right) d x+\left|D^{c} \hat{v}\right|\left(\Omega^{\prime}\right)+\int_{J_{\hat{v}} \backslash \hat{\Gamma}}|[\hat{v}]| d \mathcal{H}^{d-1}+\mathcal{H}^{d-1}\left(\hat{\Gamma} \backslash \Gamma_{0}\right) \\
& \quad=\int_{\Omega^{\prime}} f\left(\nabla \hat{v}+e_{0}\right) d x+\left|D^{c} \hat{v}\right|\left(\Omega^{\prime}\right)+\int_{J_{\hat{v}} \backslash \Gamma_{0}}|[\hat{v}]| \wedge 1 d \mathcal{H}^{d-1},
\end{aligned}
$$

which proves that $v$ solves (4.28).
Conversely, assume now that $v$ solves (4.28) and let $\Gamma:=\left(J_{v}^{1} \cup \Gamma_{0}\right) \cap \bar{\Omega}$. Let $\hat{\Gamma}$ be a Borel set with $\Gamma_{0} \subset \hat{\Gamma} \subset \bar{\Omega}$, let $\hat{v} \in G B V_{\star}\left(\Omega^{\prime}\right)$ with $\hat{v}=w-w_{0}$ $\mathcal{L}^{d}$-a.e. in $\Omega^{\prime} \backslash \Omega$, and let $\hat{\Gamma}_{1}:=J_{\hat{v}}^{1} \cup \Gamma_{0}$. By the minimality of $v$ we have

$$
\begin{aligned}
& \int_{\Omega^{\prime}} f\left(\nabla v+e_{0}\right) d x+\left|D^{c} v\right|\left(\Omega^{\prime}\right)+\int_{J_{v} \backslash \Gamma}|[v]| d \mathcal{H}^{d-1}+\mathcal{H}^{d-1}\left(\Gamma \backslash \Gamma_{0}\right) \\
& \quad=\int_{\Omega^{\prime}} f\left(\nabla v+e_{0}\right) d x+\left|D^{c} v\right|\left(\Omega^{\prime}\right)+\int_{J_{v} \backslash \Gamma_{0}}|[v]| \wedge 1 d \mathcal{H}^{d-1} \\
& \quad \leq \int_{\Omega^{\prime}} f\left(\nabla \hat{v}+e_{0}\right) d x+\left|D^{c} \hat{v}\right|\left(\Omega^{\prime}\right)+\int_{J_{\hat{v}} \backslash \Gamma_{0}}|[\hat{v}]| \wedge 1 d \mathcal{H}^{d-1}
\end{aligned}
$$

Since

$$
\begin{aligned}
& \int_{J_{\hat{v}} \backslash \Gamma_{0}}|[\hat{v}]| \wedge 1 d \mathcal{H}^{d-1}=\int_{J_{\hat{v}} \backslash \hat{\Gamma}_{1}}|[\hat{v}]| d \mathcal{H}^{d-1}+\mathcal{H}^{d-1}\left(\hat{\Gamma}_{1} \backslash \Gamma_{0}\right) \\
& \quad=\int_{\left(J_{\hat{v}} \backslash \hat{\Gamma}_{1}\right) \backslash \hat{\Gamma}}|[\hat{v}]| d \mathcal{H}^{d-1}+\int_{\left(J_{\hat{v}} \backslash \hat{\Gamma}_{1}\right) \cap \hat{\Gamma}}|[\hat{v}]| d \mathcal{H}^{d-1}+\mathcal{H}^{d-1}\left(\hat{\Gamma}_{1} \backslash \Gamma_{0}\right) \\
& \leq \int_{\left(J_{\hat{v}} \backslash \hat{\Gamma}_{1}\right) \backslash \hat{\Gamma}}|[\hat{v}]| d \mathcal{H}^{d-1}+\mathcal{H}^{d-1}\left(\hat{\Gamma} \backslash \hat{\Gamma}_{1}\right)+\mathcal{H}^{d-1}\left(\hat{\Gamma}_{1} \backslash \Gamma_{0}\right) \\
& \leq \int_{J_{\hat{v}} \backslash \hat{\Gamma}}|[\hat{v}]| d \mathcal{H}^{d-1}+\mathcal{H}^{d-1}\left(\hat{\Gamma} \backslash \Gamma_{0}\right)
\end{aligned}
$$

we obtain that

$$
\begin{aligned}
& \int_{\Omega^{\prime}} f\left(\nabla v+e_{0}\right) d x+\left|D^{c} v\right|\left(\Omega^{\prime}\right)+\int_{J_{v} \backslash \Gamma}|[v]| d \mathcal{H}^{d-1}+\mathcal{H}^{d-1}\left(\Gamma \backslash \Gamma_{0}\right) \\
& \quad \leq \int_{\Omega^{\prime}} f\left(\nabla \hat{v}+e_{0}\right) d x+\left|D^{c} \hat{v}\right|\left(\Omega^{\prime}\right)+\int_{J_{\hat{v}} \backslash \hat{\Gamma}}|[\hat{v}]| d \mathcal{H}^{d-1}+\mathcal{H}^{d-1}\left(\hat{\Gamma} \backslash \Gamma_{0}\right),
\end{aligned}
$$

which shows that $\Gamma$ and $v$ solve (4.17).
The results of this section show that the existence of a solution to the minimum problem (4.11) can be obtained by proving that the minimum problem (4.28) has a solution. To this aim we shall use the Direct Method of the Calculus of Variations. Unfortunately, not every energy-bounded sequence for (4.28) is relatively compact. For instance, if $w=w_{0}=0, \Gamma_{0}=\emptyset, e_{0}=0$, and $v_{k}=k \chi_{E}$, where $E$ is a set of finite perimeter with $\mathcal{L}^{d}(E)>0$, then $\left(v_{k}\right)_{k}$ is
energy-bounded for (4.28), but it has no subsequence which converges $\mathcal{L}^{d}$-a.e. to a finite-valued function.

The origin of this problem is the fact that, in general, an energy-bounded sequence does not satisfy (3.27). In the next section we shall construct a relatively compact minimizing sequence for problem (4.28), while in Sect. 6 we shall prove a lower semicontinuity result, which will allow us to obtain the existence of a minimizer.

## 5. Construction of a relatively compact minimizing sequence

In this section $\Omega$ is a bounded open subset of $\mathbb{R}^{d}$ with Lipschitz boundary, $c_{1}, c_{2}$ are constants with $0<c_{1} \leq c_{2}$, and $a_{1}, a_{2} \in L^{1}(\Omega)$. Given a Borel set $\Gamma_{0} \subset \bar{\Omega}$, with $\mathcal{H}^{d-1}\left(\Gamma_{0}\right)<+\infty$, and a Borel measurable function $g: \Omega \times \mathbb{R}^{d} \rightarrow \mathbb{R}$, with

$$
\begin{equation*}
c_{1}|\xi|-a_{1}(x) \leq g(x, \xi) \leq c_{2}|\xi|+a_{2}(x) \tag{5.1}
\end{equation*}
$$

for $\mathcal{L}^{d}$-a.e. $x \in \Omega$ and every $\xi \in \mathbb{R}^{d}$, we consider the functional $\mathcal{G}_{\Gamma_{0}}^{g}$ defined by

$$
\begin{equation*}
\mathcal{G}_{\Gamma_{0}}^{g}(u):=\int_{\Omega} g(x, \nabla u) d x+\left|D^{c} u\right|(\Omega)+\int_{J_{u} \backslash \Gamma_{0}}|[u]| \wedge 1 d \mathcal{H}^{d-1} \tag{5.2}
\end{equation*}
$$

for every $u \in G B V_{\star}(\Omega)$. The aim of this section is to show that, if $\left(u_{k}\right)_{k}$ is a minimizing sequence for $\mathcal{G}_{\Gamma_{0}}^{g}$, then we can modify it by means of piecewise constant translations obtaining a new minimizing sequence which satisfies the hypotheses of the compactness Theorem 3.11. The construction of the modified sequence follows the lines of [9] and requires several steps. We begin by constructing a suitable Caccioppoli partition (see [2, Definition 4.16]).

Lemma 5.1. ( $L^{\infty}$-approximation with piecewise constant functions) For every $M>0$ and for every $u \in G B V_{\star}(\Omega)$, with

$$
\begin{equation*}
\|\nabla u\|_{L^{1}\left(\Omega ; \mathbb{R}^{d}\right)}+\left|D^{c} u\right|(\Omega)+\int_{J_{u}}|[u]| \wedge 1 d \mathcal{H}^{d-1} \leq M \tag{5.3}
\end{equation*}
$$

there exist a Caccioppoli partition $\left(P_{j}\right)_{j}$ of $\Omega$ and corresponding translations $\left(b_{j}\right)_{j} \subset \mathbb{R}$ such that the function

$$
v:=u-\sum_{j=1}^{\infty} b_{j} \chi_{P_{j}}
$$

belongs to $B V(\Omega) \cap L^{\infty}(\Omega)$ and the following estimates hold:

$$
\begin{align*}
& \sum_{j=1}^{\infty} \mathcal{H}^{d-1}\left(\partial^{*} P_{j}\right) \leq 2+2 M+\mathcal{H}^{d-1}(\partial \Omega)  \tag{5.4}\\
& \|v\|_{L^{\infty}(\Omega)} \leq 2 M \tag{5.5}
\end{align*}
$$

Proof. We may assume that

$$
A:=\int_{\Omega}|\nabla u| d x+\left|D^{c} u\right|(\Omega)+\int_{J_{u} \backslash J_{u}^{1}}|[u]| d \mathcal{H}^{d-1}>0 .
$$

Indeed, if this is not the case, by (3.1) for every $m$ we can apply [2, Theorem 4.23] to the truncated function $u^{(m)}$, obtaining that $u^{(m)}$ is piecewise constant (see [2, Definition 4.21]). This implies that the function $u$ itself is piecewise constant and there is nothing to prove.

By the coarea formula in $\operatorname{GBV}(\Omega)$ (see [2, Theorem 4.34(d)] applied with $B=\Omega \backslash J_{u}^{1}$ ) for a.e. $t \in \mathbb{R}$ the set $\{u>t\}$ has finite perimeter in $\Omega$ and

$$
\int_{-\infty}^{+\infty} \mathcal{H}^{d-1}\left(\Omega \cap \partial^{*}\{u>t\} \backslash J_{u}^{1}\right) d t=A
$$

For every $i \in \mathbb{Z}$ there exists $t_{i} \in(i A,(i+1) A)$ such that the set $\left\{u>t_{i}\right\}$ has finite perimeter in $\Omega$ and

$$
\begin{equation*}
\mathcal{H}^{d-1}\left(\Omega \cap \partial^{*}\left\{u>t_{i}\right\} \backslash J_{u}^{1}\right) \leq \frac{1}{A} \int_{i A}^{(i+1) A} \mathcal{H}^{d-1}\left(\Omega \cap \partial^{*}\{u>t\} \backslash J_{u}^{1}\right) d t \tag{5.6}
\end{equation*}
$$

For every $i \in \mathbb{Z}$ let $E_{i}:=\left\{u>t_{i}\right\} \backslash\left\{u>t_{i+1}\right\}$. It is clear that $\left(E_{i}\right)_{i}$ is a partition of $\Omega$, that each $E_{i}$ has finite perimeter in $\mathbb{R}^{d}$ (recall that $\Omega$ has Lipschitz boundary, hence we can apply [2, Remark 4.20]), and that

$$
\begin{equation*}
\sum_{i \in \mathbb{Z}} \mathcal{H}^{d-1}\left(\Omega \cap \partial^{*} E_{i} \backslash J_{u}^{1}\right) \leq \frac{2}{A} \int_{-\infty}^{+\infty} \mathcal{H}^{d-1}\left(\Omega \cap \partial^{*}\{u>t\} \backslash J_{u}^{1}\right) d t=2 \tag{5.7}
\end{equation*}
$$

Let us prove now that

$$
\begin{equation*}
\sum_{i \in \mathbb{Z}} \mathcal{H}^{d-1}\left(\partial^{*} E_{i}\right) \leq 2+2 M+\mathcal{H}^{d-1}(\partial \Omega) . \tag{5.8}
\end{equation*}
$$

First of all, we claim that every $x \in J_{u}^{1}$ belongs at most to two sets $\partial^{*} E_{i}$. Indeed, it is known (see, e.g., [2, Theorem 3.61]) that for every $x \in \partial^{*} E_{i}$ we have $\lim _{\rho \rightarrow 0} \mathcal{L}^{d}\left(E_{i} \cap B_{\rho}(x)\right) / \mathcal{L}^{d}\left(B_{\rho}(x)\right)=\frac{1}{2}$. Therefore, if $x \in \partial^{*} E_{i} \cap$ $\partial^{*} E_{j} \cap \partial^{*} E_{k}$ for some $i<j<k$, we would have $\lim _{\rho \rightarrow 0} \mathcal{L}^{d}\left(\left(E_{i} \cup E_{j} \cup E_{k}\right) \cap\right.$ $\left.B_{\rho}(x)\right) / \mathcal{L}^{d}\left(B_{\rho}(x)\right)=\frac{3}{2}$, which is clearly impossible. This proves our claim, which implies that

$$
\begin{equation*}
\sum_{i \in \mathbb{Z}} \mathcal{H}^{d-1}\left(\partial^{*} E_{i} \cap J_{u}^{1}\right) \leq 2 \mathcal{H}^{d-1}\left(J_{u}^{1}\right) \leq 2 M \tag{5.9}
\end{equation*}
$$

A similar argument shows that every $x \in \partial \Omega$ belongs to at most one set $\partial^{*} E_{i}$, hence

$$
\begin{equation*}
\sum_{i \in \mathbb{Z}} \mathcal{H}^{d-1}\left(\partial^{*} E_{i} \cap \partial \Omega\right) \leq \mathcal{H}^{d-1}(\partial \Omega) \tag{5.10}
\end{equation*}
$$

Therefore (5.7), (5.9), and (5.10) give (5.8), which shows that $\left(E_{i}\right)_{i}$ is a Caccioppoli partition of $\Omega$.

Let us define $v:=u-\sum_{i \in \mathbb{Z}} t_{i} \chi_{E_{i}}$. For every $x \in E_{i}$ we have

$$
\begin{equation*}
0 \leq v(x)=u(x)-t_{i} \leq t_{i+1}-t_{i} \leq 2 A \leq 2 M \tag{5.11}
\end{equation*}
$$

which shows that

$$
\begin{equation*}
\|v\|_{L^{\infty}(\Omega)} \leq 2 M . \tag{5.12}
\end{equation*}
$$

We show that $v \in B V(\Omega)$. To this end let us consider $v_{k}:=\sum_{|i| \leq k} z_{i}$, with $z_{i}:=\left(u-t_{i}\right) \chi_{E_{i}}$. By (5.11) we have $z_{i}=\left(u^{\left(m_{i}\right)}-t_{i}\right) \chi_{E_{i}}$, with $m_{i}=2 M+$ $\left|t_{i}\right|$. Since both $u^{\left(m_{i}\right)}-t_{i}$ and $\chi_{E_{i}}$ belong to $B V(\Omega) \cap L^{\infty}(\Omega)$, by Lemma 2.4 we have $z_{i} \in B V(\Omega)$. Recalling (5.11) we have $\left[u^{\left(m_{i}\right)}\right]=[u]$ and $0 \leq\left[u^{\left(m_{i}\right)}-t_{i}\right] \leq$ $2 M$ on $E_{i}^{(1)} \cap J_{u}$ by the definition of $E^{(1)}$, while the trace operator $\gamma_{E_{i}}$ defined in Lemma 2.4 satisfies $\left|\gamma_{E_{i}}\left(u^{m_{i}}-t_{i}\right)\right| \leq 2 M \mathcal{H}^{d-1}$-a.e. on $\partial^{*} E_{i}$. Using Lemma 2.4 again, from these properties we obtain

$$
\begin{aligned}
\left|D z_{i}\right|(\Omega)= & \int_{\Omega}\left|\nabla z_{i}\right| d x+\left|D^{c} z_{i}\right|(\Omega)+\left|D^{j} z_{i}\right|(\Omega) \\
\leq & \int_{E_{i}}\left|\nabla u^{\left(m_{i}\right)}\right| d x+\left|D^{c} u^{\left(m_{i}\right)}\right|\left(E_{i}^{(1)}\right) \\
& +\int_{E_{i}^{(1)} \cap J_{u}^{\left(m_{i}\right)}}\left|\left[u^{\left(m_{i}\right)}\right]\right| d \mathcal{H}^{d-1}+2 M \mathcal{H}^{d-1}\left(\partial^{*} E_{i}\right) \\
\leq & \int_{E_{i}}|\nabla u| d x+\left|D^{c} u\right|\left(E_{i}^{(1)}\right)+\int_{E_{i}^{(1)} \cap\left(J_{u} \backslash J_{u}^{1}\right)}|[u]| d \mathcal{H}^{d-1} \\
& +2 M \mathcal{H}^{d-1}\left(E_{i}^{(1)} \cap J_{u}^{1}\right)+2 M \mathcal{H}^{d-1}\left(\partial^{*} E_{i}\right),
\end{aligned}
$$

where the last inequality follows from Proposition 2.6(c) and (2.23). Since the sets $E_{i}^{(1)}$ are pairwise disjoint, by (5.8)

$$
\left|D v_{k}\right|(\Omega) \leq \int_{\Omega}|\nabla u| d x+\left|D^{c} u\right|(\Omega)+\int_{J_{u} \backslash J_{u}^{1}}|[u]| d \mathcal{H}^{d-1}+\hat{M}
$$

where $\hat{M}:=2 M\left(2+3 M+\mathcal{H}^{d-1}(\partial \Omega)\right)$. Since the right-hand side is finite, we obtain that $\left|D v_{k}\right|(\Omega)$ is bounded uniformly with respect to $k$. On the other hand, since $\left(E_{i}\right)_{i}$ is a partition, inequality (5.12) implies that the sequence $\left(v_{k}\right)_{k}$ is bounded in $L^{\infty}(\Omega)$ and that $v_{k} \rightarrow v$ strongly in $L^{1}(\Omega)$. Therefore $v \in B V(\Omega)$.

To conclude the proof it is enough to take $P_{j}=E_{\sigma(j)}$ and $b_{j}=t_{\sigma(j)}$ where $\sigma: \mathbb{N} \rightarrow \mathbb{Z}$ is bijective.

In the following lemma the Caccioppoli partition is finite and we provide a precise estimate on the translations.

Lemma 5.2. (Piecewise Poincaré inequality) Let $\alpha \geq 1$ and let $0<\theta<1$. Then there exist positive constants $C_{\Omega}$ and $C_{\theta, \alpha, d}$ such that for every $u \in G B V_{\star}(\Omega)$ there exist a finite Caccioppoli partition $\Omega=\bigcup_{j=1}^{J} P_{j} \cup R_{1} \cup R_{2}$, a finite family of translations $\left(b_{j}\right)_{j=1}^{J} \subset \mathbb{R}$, and a constant $\lambda \in\left[1, C_{\theta, \alpha, d}\right]$, depending on $u$, satisfying the following estimates:

$$
\begin{align*}
& \mathcal{L}^{d}\left(R_{1} \cup R_{2}\right) \leq \theta C_{\Omega} \mathcal{H}^{d-1}\left(J_{u}^{1} \cup \partial \Omega\right),  \tag{5.13}\\
& \mathcal{H}^{d-1}\left(\partial^{*} R_{1}\right) \leq \theta C_{\Omega} \mathcal{H}^{d-1}\left(J_{u}^{1} \cup \partial \Omega\right),  \tag{5.14}\\
& \sum_{j=1}^{J} \mathcal{H}^{d-1}\left(\partial^{*} P_{j}\right)+\mathcal{H}^{d-1}\left(\partial^{*} R_{2}\right) \leq C_{\Omega} \mathcal{H}^{d-1}\left(J_{u}^{1} \cup \partial \Omega\right), \tag{5.15}
\end{align*}
$$

$$
\begin{align*}
& \max _{1 \leq j \leq J}\left\|u-b_{j}\right\|_{L^{\infty}\left(P_{j}\right)} \leq \lambda\left(\int_{\Omega}|\nabla u| d x+\left|D^{c} u\right|(\Omega)+\int_{J_{u} \backslash J_{u}^{1}}|[u]| d \mathcal{H}^{d-1}\right),  \tag{5.16}\\
& \min _{1 \leq j \leq J} \operatorname{essinf}\left|u-b_{R_{2}}\right| \geq \alpha \lambda\left(\int_{\Omega}|\nabla u| d x+\left|D^{c} u\right|(\Omega)+\int_{J_{u} \backslash J_{u}^{1}}|[u]| d \mathcal{H}^{d-1}\right),  \tag{5.17}\\
& \min _{1 \leq i<j \leq J}\left|b_{i}-b_{j}\right| \geq \alpha \lambda\left(\int_{\Omega}|\nabla u| d x+\left|D^{c} u\right|(\Omega)+\int_{J_{u} \backslash J_{u}^{1}}|[u]| d \mathcal{H}^{d-1}\right) . \tag{5.18}
\end{align*}
$$

Proof. It is enough to repeat the proof of [9, Lemma 3.5] replacing the space $G S B V^{p}\left(\Omega ; \mathbb{R}^{m}\right)$ by $G B V_{\star}(\Omega), \mathcal{H}^{d-1}\left(J_{u} \cup \partial \Omega\right)$ by $\mathcal{H}^{d-1}\left(J_{u}^{1} \cup \partial \Omega\right),\|\nabla u\|_{L^{1}(\Omega)}$ by $\|\nabla u\|_{L^{1}\left(\Omega ; \mathbb{R}^{d}\right)}+\left|D^{c} u\right|(\Omega)+\int_{J_{u} \backslash J_{u}^{1}}|[u]| d \mathcal{H}^{d-1}$, and $[9$, Theorem 2.5] by Corollary 5.1.

The following theorem shows that we can modify a function $u$ by means of piecewise constant translations, with a precise control on the value taken by the functional $\mathcal{G}_{\Gamma_{0}}^{g}$ defined in (5.2) on the modified function.

Theorem 5.3. (Piecewise translated functions) Let $M>0$ and $0<\theta<1$. Then there exist positive constants $C_{M, \Omega}$ and $C_{M, \theta, \Omega}$ with the following property: for every $u \in G B V_{\star}(\Omega)$ with

$$
\begin{equation*}
\|\nabla u\|_{L^{1}\left(\Omega ; \mathbb{R}^{d}\right)}+\left|D^{c} u\right|(\Omega)+\int_{J_{u}}|[u]| \wedge 1 d \mathcal{H}^{d-1} \leq M \tag{5.19}
\end{equation*}
$$

there exist a finite Caccioppoli partition $\Omega=\bigcup_{j=1}^{J} P_{j} \cup R$ and a finite family of translations $\left(t_{j}\right)_{j=1}^{J} \subset \mathbb{R}$ such that the function

$$
\begin{equation*}
v:=\sum_{j=1}^{J}\left(u-t_{j}\right) \chi_{P_{j}} \tag{5.20}
\end{equation*}
$$

belongs to $B V(\Omega) \cap L^{\infty}(\Omega)$ and the following estimates hold:

$$
\begin{align*}
& \mathcal{L}^{d}(R) \leq \theta C_{M, \Omega},  \tag{5.21}\\
& \sum_{j=1}^{J} \mathcal{H}^{d-1}\left(\partial^{*} P_{j}\right)+\mathcal{H}^{d-1}\left(\partial^{*} R\right) \leq C_{M, \Omega},  \tag{5.22}\\
& \|v\|_{L^{\infty}(\Omega)} \leq C_{M, \theta, \Omega}  \tag{5.23}\\
& \mathcal{G}_{\Gamma_{0}}^{g}(v) \leq \mathcal{G}_{\Gamma_{0}}^{g}(u)+\theta C_{M, \Omega}+\|a\|_{L^{1}(R)}, \tag{5.24}
\end{align*}
$$

for every $g$ satisfying (5.1) and for every Borel set $\Gamma_{0} \subset \Omega$ with $\mathcal{H}^{d-1}\left(\Gamma_{0}\right)<$ $+\infty$, where $a:=\left|a_{1}\right|+\left|a_{2}\right|$. Moreover, $\{v=0\} \supset\{u=0\}$ (up to a set of negligible $\mathcal{L}^{d}$ measure $)$. Finally, we can choose $\left(P_{j}\right)_{j=1}^{J}$ and $\left(t_{j}\right)_{j=1}^{J}$ so that the following additional property holds: for every collection $\left(t_{j}^{\prime}\right)_{j=1}^{J}$, with $\left|t_{j}-t_{j}^{\prime}\right| \leq$ $\theta^{-1}\|v\|_{L^{\infty}(\Omega)}$, the function $v^{\prime}:=\sum_{j=1}^{J}\left(u-t_{j}^{\prime}\right) \chi_{P_{j}}$ belongs to $B V(\Omega) \cap L^{\infty}(\Omega)$ and satisfies (5.24).

Proof. It is enough to repeat the proof of [9, Theorem 3.2] replacing $c_{3}$ by 1, $G S B V_{M}^{p}\left(\Omega ; \mathbb{R}^{m}\right)$ by $\left\{u \in G B V_{\star}(\Omega):(5.19)\right.$ holds $\}, \mathcal{H}^{d-1}\left(J_{u}\right)$ by $\mathcal{H}^{d-1}\left(J_{u}^{1}\right)$, $\mathcal{H}^{d-1}\left(J_{u} \cup \partial \Omega\right)$ by $\mathcal{H}^{d-1}\left(J_{u}^{1} \cup \partial \Omega\right),\|\nabla u\|_{L^{1}(\Omega)}$ by $\|\nabla u\|_{L^{1}\left(\Omega ; \mathbb{R}^{d}\right)}+\left|D^{c} u\right|(\Omega)+$ $\int_{J_{u} \backslash J_{u}^{1}}|[u]| d \mathcal{H}^{d-1}$, and [9, Lemma 3.5] by Lemma 5.2 above, obtaining that (5.21)-(5.23) hold.

To prove that $v \in B V(\Omega)$ it is enough to show that $\left(u-t_{j}\right) \chi_{P_{j}} \in B V(\Omega)$. We observe that for $m \geq C_{M, \theta, \Omega}+\left|t_{j}\right|$ we have $\left(u-t_{j}\right) \chi_{P_{j}}=\left(u^{(m)}-t_{j}\right) \chi_{P_{j}}$ by (5.23). Since $u^{(m)}-t_{j}$ and $\chi_{P_{j}}$ belong to $B V(\Omega) \cap L^{\infty}(\Omega)$, we conclude that $\left(u-t_{j}\right) \chi_{P_{j}} \in B V(\Omega)$.

It remains to prove (5.24). More precisely, we shall prove that

$$
\begin{align*}
& \int_{\Omega} g(x, \nabla v) d x \leq \int_{\Omega} g(x, \nabla u) d x+\int_{R} a d x  \tag{5.25}\\
& \left|D^{c} v\right|(\Omega) \leq\left|D^{c} u\right|(\Omega)  \tag{5.26}\\
& \int_{J_{v} \backslash \Gamma_{0}}|[v]| \wedge 1 d \mathcal{H}^{d-1} \leq \int_{J_{u} \backslash \Gamma_{0}}|[u]| \wedge 1 d \mathcal{H}^{d-1}+\theta C_{M, \Omega} \tag{5.27}
\end{align*}
$$

which gives (5.24). Inequality (5.25) can be proved as in the proof of $[9$, formula (14)], while (5.27) can be proved as in the last part of [9, Theorem 3.2]. As for (5.26) we begin by observing that by (5.23) there exists a constant $m>0$ such that $v=\sum_{j=1}^{J}\left(u^{(m)}-t_{j}\right) \chi_{P_{j}}$. By (2.11) and (2.23) we obtain

$$
\left|D^{c} v\right|(\Omega)=\sum_{j=1}^{J}\left|D^{c}\left(u^{(m)}-t_{j}\right)\right|\left(P_{j}^{(1)}\right) \leq \sum_{j=1}^{J}\left|D^{c} u\right|\left(P_{j}^{(1)}\right) \leq\left|D^{c} u\right|(\Omega),
$$

which gives (5.26).
The previous result can be extended to the case of functions satisfying prescribed boundary conditions in the usual $B V$ sense considered in (4.17) and (4.28). To this aim we introduce a bounded open set $\Omega^{\prime} \subset \mathbb{R}^{d}$ with Lipschitz boundary and containing $\bar{\Omega}$.

Corollary 5.4. (The case of boundary conditions) Let $M>0$ and $0<\theta<1$. Then there exist positive constants $C_{M, \Omega^{\prime}}$ and $C_{M, \theta, \Omega^{\prime}}$ with the following property: for each $h \in W^{1,1}\left(\Omega^{\prime}\right)$ with $\|\nabla h\|_{L^{1}\left(\Omega^{\prime} ; \mathbb{R}^{d}\right)} \leq M$ and each $u \in G B V_{\star}\left(\Omega^{\prime}\right)$ with $u=h \mathcal{L}^{d}$-a.e. on $\Omega^{\prime} \backslash \bar{\Omega}$ and

$$
\begin{equation*}
\int_{\Omega^{\prime}}|\nabla u| d x+\left|D^{c} u\right|\left(\Omega^{\prime}\right)+\int_{J_{u}}|[u]| \wedge 1 d \mathcal{H}^{d-1} \leq M \tag{5.28}
\end{equation*}
$$

there exist a finite Caccioppoli partition $\Omega^{\prime}=\bigcup_{j=1}^{J} P_{j} \cup R$ and a finite family of translations $\left(t_{j}\right)_{j=1}^{J}$ such that the function

$$
\begin{equation*}
v:=h \chi_{R}+\sum_{j=1}^{J}\left(u-t_{j}\right) \chi_{P_{j}} \tag{5.29}
\end{equation*}
$$

satisfies

$$
\begin{equation*}
v \in B V\left(\Omega^{\prime}\right), v=h \mathcal{L}^{d} \text {-a.e. on } \Omega^{\prime} \backslash \bar{\Omega}, v-h \in B V\left(\Omega^{\prime}\right) \cap L^{\infty}\left(\Omega^{\prime}\right), \tag{5.30}
\end{equation*}
$$

and the following estimates hold:

$$
\begin{align*}
& \|v-h\|_{L^{\infty}\left(\Omega^{\prime}\right)} \leq C_{M, \theta, \Omega^{\prime}}  \tag{5.31}\\
& \mathcal{L}^{d}(R) \leq \theta C_{M, \Omega^{\prime}}  \tag{5.32}\\
& \sum_{j=1}^{J} \mathcal{H}^{d-1}\left(\partial^{*} P_{j}\right)+\mathcal{H}^{d-1}\left(\partial^{*} R\right) \leq C_{M, \Omega^{\prime}}  \tag{5.33}\\
& \mathcal{L}^{d}\left(P_{j} \cap\left(\Omega^{\prime} \backslash \bar{\Omega}\right)\right)>0 \quad \text { for at most one index } j  \tag{5.34}\\
& \mathcal{G}_{\Gamma_{0}}^{g}\left(v, \Omega^{\prime}\right) \leq \mathcal{G}_{\Gamma_{0}}^{g}\left(u, \Omega^{\prime}\right)+\theta C_{M, \Omega^{\prime}}+\|a\|_{L^{1}(R)}+c_{2}\|\nabla h\|_{L^{1}\left(R ; \mathbb{R}^{d}\right)} \tag{5.35}
\end{align*}
$$

for every Borel measurable $g: \Omega^{\prime} \times \mathbb{R}^{d} \rightarrow \mathbb{R}$ satisfying (5.1) on $\Omega^{\prime} \times \mathbb{R}^{d}$, and for every Borel set $\Gamma_{0} \subset \Omega^{\prime}$, with $\mathcal{H}^{d-1}\left(\Gamma_{0}\right)<+\infty$, where $\mathcal{G}_{\Gamma_{0}}^{g}\left(u, \Omega^{\prime}\right)$ is defined as in (5.2), with $\Omega$ replaced by $\Omega^{\prime}$, and $a=\left|a_{1}\right|+\left|a_{2}\right|$. Moreover, we can choose $\left(P_{j}\right)_{j=1}^{J}$ and $\left(t_{j}\right)_{j=1}^{J}$ so that the following additional property holds: for every collection $\left(t_{j}^{\prime}\right)_{j=1}^{J}$, with $\left|t_{j}-t_{j}^{\prime}\right| \leq \theta^{-1}\|v-h\|_{L^{\infty}\left(\Omega^{\prime}\right)}$, the function $v^{\prime}:=h \chi_{R}+\sum_{j=1}^{J}\left(u-t_{j}^{\prime}\right) \chi_{P_{j}}$ belongs to $B V\left(\Omega^{\prime}\right)$ and satisfies (5.35).

Proof. We apply Theorem 5.3 to the function $u-h$ on $\Omega^{\prime}$, which belongs to $G B V_{\star}\left(\Omega^{\prime}\right)$ by Theorem 3.9. Arguing as in the proof of [9, Corollary 3.3] we obtain a finite Caccioppoli partition $\Omega^{\prime}=\bigcup_{j=1}^{J} P_{j} \cup R$ and a finite family of translations $\left(t_{j}\right)_{j=1}^{J} \subset \mathbb{R}$ such that (5.32)-(5.34) hold, the function

$$
\begin{equation*}
z:=\sum_{j=1}^{J}\left(u-h-t_{j}\right) \chi_{P_{j}}, \tag{5.36}
\end{equation*}
$$

belongs to $B V\left(\Omega^{\prime}\right) \cap L^{\infty}\left(\Omega^{\prime}\right)$, and

$$
\begin{equation*}
\|z\|_{L^{\infty}\left(\Omega^{\prime}\right)} \leq C_{M, \theta, \Omega^{\prime}} \tag{5.37}
\end{equation*}
$$

Moreover, by (5.26) and (5.27), $z$ satisfies

$$
\begin{align*}
& \left|D^{c} z\right|\left(\Omega^{\prime}\right) \leq\left|D^{c}(u-h)\right|\left(\Omega^{\prime}\right)=\left|D^{c} u\right|\left(\Omega^{\prime}\right)  \tag{5.38}\\
& \int_{J_{z} \backslash \Gamma_{0}}|[z]| \wedge 1 d \mathcal{H}^{d-1} \leq \int_{J_{u-h} \backslash \Gamma_{0}}|[u-h]| \wedge 1 d \mathcal{H}^{d-1}+\theta C_{M, \Omega^{\prime}} \\
& \quad=\int_{J_{u} \backslash \Gamma_{0}}|[u]| \wedge 1 d \mathcal{H}^{d-1}+\theta C_{M, \Omega^{\prime}} \tag{5.39}
\end{align*}
$$

Let $v:=z+h$. Then $v \in B V\left(\Omega^{\prime}\right)$ and satisfies (5.29), (5.30), and (5.31).
It remains to prove (5.35). More precisely, we shall prove that

$$
\begin{align*}
& \int_{\Omega^{\prime}} g(x, \nabla v) d x \leq \int_{\Omega^{\prime}} g(x, \nabla u) d x+\int_{R} a d x+c_{2} \int_{R}|\nabla h| d x,  \tag{5.40}\\
& \left|D^{c} v\right|\left(\Omega^{\prime}\right) \leq\left|D^{c} u\right|\left(\Omega^{\prime}\right),  \tag{5.41}\\
& \int_{J_{v} \backslash \Gamma_{0}}|[v]| \wedge 1 d \mathcal{H}^{d-1} \leq \int_{J_{u} \backslash \Gamma_{0}}|[u]| \wedge 1 d \mathcal{H}^{d-1}+\theta C_{M, \Omega^{\prime}}, \tag{5.42}
\end{align*}
$$

which gives (5.35).
Inequalities (5.41) and (5.42) follow from (5.38) and (5.39), respectively, since $D^{c} v=D^{c} z, J_{v}=J_{z}$, and $[z]=[v]$ (recall that $\left.h \in W^{1,1}\left(\Omega^{\prime}\right)\right)$.

To prove (5.40) we observe that $\nabla v=\nabla u \mathcal{L}^{d}$-a.e. on $\Omega^{\prime} \backslash R$, while $\nabla v=$ $\nabla h \mathcal{L}^{d}$-a.e. on $R$, so that

$$
\begin{aligned}
& \int_{\Omega^{\prime}} g(x, \nabla v) d x=\int_{\Omega^{\prime} \backslash R} g(x, \nabla u) d x+\int_{R} g(x, \nabla h) d x \\
& \quad \leq \int_{\Omega^{\prime}} g(x, \nabla u) d x+\int_{R} a d x+c_{2} \int_{R}|\nabla h| d x,
\end{aligned}
$$

which concludes the proof.
We are now in a position to prove the main result of this section.
Theorem 5.5. (Existence of modifications satisfying (3.27)) Let $h \in W^{1,1}\left(\Omega^{\prime}\right)$, let $g: \Omega^{\prime} \times \mathbb{R}^{d} \rightarrow \mathbb{R}$ be a Borel measurable function satisfying (5.1) on $\Omega^{\prime} \times \mathbb{R}^{d}$, and let $\Gamma_{0} \subset \Omega^{\prime}$ be a Borel set with $\mathcal{H}^{d-1}\left(\Gamma_{0}\right)<+\infty$.

Let $\left(u_{k}\right)_{k}$ be a sequence in $G B V_{\star}\left(\Omega^{\prime}\right)$ with $u_{k}=h$ on $\Omega^{\prime} \backslash \bar{\Omega}$. Assume that there exists $C>0$ such that

$$
\begin{equation*}
\mathcal{G}_{\Gamma_{0}}^{g}\left(u_{k}, \Omega^{\prime}\right) \leq C \quad \text { for every } k . \tag{5.43}
\end{equation*}
$$

Then there exist a subsequence of $\left(u_{k}\right)_{k}$, not relabelled, modifications $y_{k} \in$ $G B V_{\star}\left(\Omega^{\prime}\right)$ of $u_{k}$, with $y_{k}=h$ on $\Omega^{\prime} \backslash \bar{\Omega}$, and a continuous function $\psi: \mathbb{R}_{+} \rightarrow$ $\mathbb{R}_{+}$with $\psi(t) \rightarrow+\infty$ as $t \rightarrow \infty$ such that

$$
\begin{align*}
& \mathcal{G}_{\Gamma_{0}}^{g}\left(y_{k}, \Omega^{\prime}\right) \leq \mathcal{G}_{\Gamma_{0}}^{g}\left(u_{k}, \Omega^{\prime}\right)+\frac{1}{k},  \tag{5.44}\\
& \sup _{k} \int_{\Omega^{\prime}} \psi\left(\left|y_{k}\right|\right) d x<+\infty . \tag{5.45}
\end{align*}
$$

Remark 5.6. By (5.44), if $\left(u_{k}\right)_{k}$ is a minimizing sequence for the functional $\mathcal{G}_{\Gamma_{0}}^{g}$ with $u_{k}=h \mathcal{L}^{d}$-a.e. in $\Omega^{\prime} \backslash \bar{\Omega}$, then the same is true for $\left(y_{k}\right)_{k}$. Inequalities (5.43) and (5.44), together with (5.1), imply that $\left(y_{k}\right)_{k}$ satisfies (3.26), while (5.45) guarantees that (3.27) also holds. Hence, by Theorem 3.11 there exists a subsequence of $\left(y_{k}\right)_{k}$, not relabelled, and a function $u \in G B V_{\star}\left(\Omega^{\prime}\right)$ such that $y_{k} \rightarrow u \mathcal{L}^{d}$-a.e. in $\Omega^{\prime}$.
Proof of Theorem 5.5. We repeat the proof of [9, Theorem 3.8] with some modifications. By (5.43) we have

$$
\begin{equation*}
c_{1} \int_{\Omega^{\prime}}\left|\nabla u_{k}\right| d x+\left|D^{c} u_{k}\right|\left(\Omega^{\prime}\right)+\int_{J_{u_{k}}}\left|\left[u_{k}\right]\right| \wedge 1 d \mathcal{H}^{d-1} \leq M_{1} \tag{5.46}
\end{equation*}
$$

with $M_{1}:=C+\|a\|_{L^{1}\left(\Omega^{\prime}\right)}+\mathcal{H}^{d-1}\left(\Gamma_{0}\right)$, where $a=\left|a_{1}\right|+\left|a_{2}\right|$. We define $\theta_{\ell}:=2^{-\ell}$ and apply Corollary 5.4. Let us remark that, since we will pass to subsequences (not relabelled), we will eventually have only the inequality

$$
\begin{equation*}
\theta_{\ell} \leq 2^{-\ell} \tag{5.47}
\end{equation*}
$$

Step 1 (Application of Corollary 5.4) We apply Corollary 5.4 to the functions $u_{k}$ and the boundary data $h$ with parameters $\theta_{\ell}$ and $M:=\frac{M_{1}}{\min \left\{c_{1}, 1\right\}}$. We find finite Caccioppoli partitions $\Omega^{\prime}=\cup_{j \geq 1} P_{j}^{k, \ell} \cup R_{k}^{\ell}$ and piecewise translated functions $v_{k}^{\ell} \in B V\left(\Omega^{\prime}\right)$ defined by

$$
\begin{equation*}
v_{k}^{\ell}:=h+\sum_{j \geq 1}\left(u_{k}-t_{j}^{k, \ell}-h\right) \chi_{P_{j}^{k, \ell}}=h \chi_{R_{k}^{\ell}}+\sum_{j \geq 1}\left(u_{k}-t_{j}^{k, \ell}\right) \chi_{P_{j}^{k, \ell}} \tag{5.48}
\end{equation*}
$$

where $\left(t_{j}^{k, \ell}\right)_{j \geq 1}$ are suitable finite families of translations. For notational convenience we shall also use the notation $P_{0}^{k, \ell}=R_{k}^{\ell}$ so that $\left(P_{j}^{k, \ell}\right)_{j \geq 0}$ is a partition of $\Omega^{\prime}$. By Corollary 5.4 we have

$$
\begin{align*}
& v_{k}^{\ell} \in B V\left(\Omega^{\prime}\right), v_{k}^{\ell}=h \mathcal{L}^{d} \text {-a.e. on } \Omega^{\prime} \backslash \bar{\Omega}, v_{k}^{\ell}-h \in L^{\infty}\left(\Omega^{\prime}\right),  \tag{5.49}\\
& \left\|v_{k}^{\ell}-h\right\|_{L^{\infty}\left(\Omega^{\prime}\right)} \leq C_{M, \theta_{\ell}, \Omega^{\prime}},  \tag{5.50}\\
& \mathcal{L}^{d}\left(R_{k}^{\ell}\right) \leq \theta_{\ell} C_{M, \Omega^{\prime}}  \tag{5.51}\\
& \sum_{j \geq 0} \mathcal{H}^{d-1}\left(\partial^{*} P_{j}^{k, \ell}\right) \leq C_{M, \Omega^{\prime}}, \tag{5.52}
\end{align*}
$$

for every $k, \ell$ there is at most one $j$ with $\mathcal{L}^{d}\left(P_{j}^{k, \ell} \cap \Omega^{\prime} \backslash \bar{\Omega}\right)>0$,
$\mathcal{G}_{\Gamma_{0}}^{g}\left(v_{k}^{\ell}, \Omega^{\prime}\right) \leq \mathcal{G}_{\Gamma_{0}}^{g}\left(u_{k}, \Omega^{\prime}\right)+\theta_{\ell} C_{M, \Omega^{\prime}}+\|a\|_{L^{1}\left(R_{k}^{\ell}\right)}+c_{2}\|\nabla h\|_{L^{1}\left(R_{k}^{\ell} ; \mathbb{R}^{d}\right)}$.

By (5.51) there exists a decreasing sequence $\eta_{\ell}$ converging to zero such that

$$
\|a\|_{L^{1}\left(R_{k}^{\ell}\right)}+c_{2}\|\nabla h\|_{L^{1}\left(R_{k}^{\ell} ; \mathbb{R}^{d}\right)} \leq \eta_{\ell} C_{M, \Omega^{\prime}},
$$

which together with (5.54) gives

$$
\begin{equation*}
\mathcal{G}_{\Gamma_{0}}^{g}\left(v_{k}^{\ell}, \Omega^{\prime}\right) \leq \mathcal{G}_{\Gamma_{0}}^{g}\left(u_{k}, \Omega^{\prime}\right)+\left(\theta_{\ell}+\eta_{\ell}\right) C_{M, \Omega^{\prime}} \quad \text { for every } k \text { and } \ell . \tag{5.55}
\end{equation*}
$$

For later use we recall that for every family $\left(\hat{t}_{j}^{k, \ell}\right)_{j \geq 1}$, with $\left|t_{j}^{k, \ell}-\hat{t}_{j}^{k, \ell}\right| \leq$ $\theta_{\ell}^{-1}\left\|v_{k}^{\ell}-h\right\|_{L^{\infty}\left(\Omega^{\prime}\right)}$, we have that the functions $\hat{v}_{k}^{\ell}:=h \chi_{R_{k}^{\ell}}+\sum_{j \geq 1}\left(u_{k}-\right.$ $\left.\hat{t}_{j}^{k, \ell}\right) \chi_{P_{j}^{k, \ell}}$ belong to $B V\left(\Omega^{\prime}\right)$ and satisfy

$$
\begin{equation*}
\mathcal{G}_{\Gamma_{0}}^{g}\left(\hat{v}_{k}^{\ell}, \Omega^{\prime}\right) \leq \mathcal{G}_{\Gamma_{0}}^{g}\left(u_{k}, \Omega\right)+\left(\theta_{\ell}+\eta_{\ell}\right) C_{M, \Omega^{\prime}}, \tag{5.56}
\end{equation*}
$$

see (5.35).
Step 2 (Limiting objects for each $\ell$ ) By (5.43), (5.50), and (5.55) we obtain that for every $\ell$ the sequence $\left(v_{k}^{\ell}\right)_{k}$ is bounded in $B V\left(\Omega^{\prime}\right)$. Indeed, arguing as in the proof of (5.46), by (5.43) and (5.55) we have that

$$
c_{1} \int_{\Omega^{\prime}}\left|\nabla v_{k}^{\ell}\right| d x+\left|D^{c} v_{k}^{\ell}\right|\left(\Omega^{\prime}\right)+\int_{J_{v_{k}^{\ell}}}\left|\left[v_{k}^{\ell}\right]\right| \wedge 1 d \mathcal{H}^{d-1} \leq M_{1}+\left(\theta_{\ell}+\eta_{\ell}\right) C_{M, \Omega^{\prime}}
$$

This implies that $\left\|\nabla v_{k}^{\ell}\right\|_{L^{1}\left(\Omega^{\prime} ; \mathbb{R}^{d}\right)},\left|D^{c} v_{k}^{\ell}\right|\left(\Omega^{\prime}\right)$, and $\mathcal{H}^{d-1}\left(J_{v_{k}^{\ell}}^{1}\right)$ are bounded. Since $[h]=0$ we have $\int_{J_{v_{k}^{\ell}}}\left|\left[v_{k}^{\ell}\right]\right| d \mathcal{H}^{d-1}=\int_{J_{v_{k}^{\ell}} \backslash J_{v_{k}^{\ell}}^{1}}\left|\left[v_{k}^{\ell}\right]\right| \wedge 1 d \mathcal{H}^{d-1}+\int_{J_{v_{k}^{\ell}}^{1}} \mid\left[v_{k}^{\ell}-\right.$ $h] \mid d \mathcal{H}^{d-1} \leq M_{1}+\left(\theta_{\ell}+\eta_{\ell}\right) C_{M, \Omega^{\prime}}+2 \mathcal{H}^{d-1}\left(J_{v_{k}^{\ell}}^{1}\right)\left\|v_{k}^{\ell}-h\right\|_{L^{\infty}\left(\Omega^{\prime}\right)}$. By (5.50) we obtain that $\int_{v_{v_{k}^{\ell}}}\left|\left[v_{k}^{\ell}\right]\right| d \mathcal{H}^{d-1}$ is bounded with respect to $k$. Together with the previous bounds this implies that $\left|D v_{k}^{\ell}\right|\left(\Omega^{\prime}\right)$ is bounded uniformly with respect to $k$. Since $v_{k}^{\ell}=h$ on $\Omega^{\prime} \backslash \bar{\Omega}$, by the Poincaré inequality we deduce that $\left(v_{k}^{\ell}\right)_{k}$ is bounded in $B V\left(\Omega^{\prime}\right)$.

Using a diagonal argument we obtain a subsequence of $(k)_{k}$ (not relabelled) such that for every $\ell$ there exist a function $v^{\ell} \in B V\left(\Omega^{\prime}\right)$ and a constant
$L_{\ell} \in\left[0, C_{M, \theta_{\ell}, \Omega^{\prime}}\right]($ see (5.50)) such that

$$
\begin{equation*}
v_{k}^{\ell} \rightarrow v^{\ell} \quad \text { in } L^{1}\left(\Omega^{\prime}\right) \quad \text { and } \quad\left\|v_{k}^{\ell}-h\right\|_{L^{\infty}\left(\Omega^{\prime}\right)} \rightarrow L_{\ell} . \tag{5.57}
\end{equation*}
$$

By the semicontinuity of the $L^{\infty}$-norm we obtain

$$
\begin{equation*}
\left\|v^{\ell}-h\right\|_{L^{\infty}\left(\Omega^{\prime}\right)} \leq L_{\ell} \tag{5.58}
\end{equation*}
$$

Arguing as in Step 2 of the proof of [9, Theorem 3.8] we find Caccioppoli partitions $\left(P_{j}^{\ell}\right)_{j \geq 0}$ and $\left(P_{j}\right)_{j \geq 0}$ such that after extracting (not relabelled) subsequences in $\ell$ and $k$, we get

$$
\begin{equation*}
\sum_{j \geq 0} \mathcal{L}^{d}\left(P_{j}^{\ell} \triangle P_{j}\right) \leq 2^{-\ell} \text { and } \sum_{j \geq 0} \mathcal{L}^{d}\left(P_{j}^{k, \ell} \triangle P_{j}^{k}\right) \leq 2^{-\ell} \quad \text { for all } k \geq \ell \tag{5.59}
\end{equation*}
$$

Step 3 (Conclusion of the proof) If $\left(L_{\ell}\right)_{\ell}$ does not tend to $+\infty$ as $\ell \rightarrow+\infty$, by (5.58) there exists a subsequence, not relabelled, such that $\left(v^{\ell}-h\right)_{\ell}$ is bounded in $L^{\infty}\left(\Omega^{\prime}\right)$. Then $\left(v^{\ell}\right)_{\ell}$ is bounded in $L^{1}\left(\Omega^{\prime}\right)$ and we can take $\psi(t)=t$ to obtain

$$
\begin{equation*}
\sup _{\ell} \int_{\Omega^{\prime}} \psi\left(\left|v^{\ell}\right|\right) d x<+\infty \tag{5.60}
\end{equation*}
$$

The conclusion can now be obtained by repeating Step 5 of the proof of [9, Theorem 3.8] replacing $\hat{v}^{\ell}, \hat{v}_{k}^{\ell}$, and $E_{k}$ by $v^{\ell}, v_{k}^{\ell}$, and $\mathcal{G}_{\Gamma_{0}}^{g}\left(\cdot, \Omega^{\prime}\right)$, respectively.

If $L_{\ell} \rightarrow+\infty$, passing to a subsequence, not relabelled, we may assume that $L_{\ell}<L_{\ell+1}$. By the definition of $L_{\ell}$, for every $\ell$ we can find an increasing sequence $\left(k_{\ell}\right)_{\ell}$ such that $\left\|v_{k}^{\ell}-h\right\|_{L^{\infty}\left(\Omega^{\prime}\right)}<\left\|v_{k}^{\ell+1}-h\right\|_{L^{\infty}\left(\Omega^{\prime}\right)}$ for every $\ell$ and for every $k \geq k_{\ell}$. This allows us to follow the lines of Step 3 of the proof of [9, Theorem 3.8]. Namely we replace the translations $t_{j}^{k, \ell}$ by the translations $\hat{t}_{j}^{k, \ell}$ introduced in that paper and we consider the corresponding functions $\hat{v}_{k}^{\ell}$ defined as in (5.48) with $t_{j}^{k, \ell}$ replaced by $\hat{t}_{j}^{k, \ell}$. This construction leads to the fact that $\hat{v}_{k}^{\ell}$ satisfies (5.56) and

$$
\begin{equation*}
\left\|\hat{v}_{k}^{\ell}-h\right\|_{L^{\infty}\left(\Omega^{\prime}\right)} \leq 2 \sum_{m=1}^{\ell} C_{M, \theta_{m}, \Omega^{\prime}} \quad \text { for every } k \geq k_{\ell} . \tag{5.61}
\end{equation*}
$$

Hence we can repeat the argument leading to (5.57) and we obtain a subsequence of $(k)_{k}$ (not relabelled) and, for every $\ell$, a function $\hat{v}^{\ell} \in B V\left(\Omega^{\prime}\right)$ such that

$$
\begin{equation*}
\hat{v}_{k}^{\ell} \rightarrow \hat{v}^{\ell} \quad \text { in } L^{1}\left(\Omega^{\prime}\right) \tag{5.62}
\end{equation*}
$$

The conclusion can now be obtained by repeating Steps 4 and 5 in the proof of $\left[9\right.$, Theorem 3.8] with $E_{k}$ replaced by $\mathcal{G}_{\Gamma_{0}}^{g}\left(\cdot, \Omega^{\prime}\right)$.

## 6. Existence result

In this section we shall prove that the minimum problem (4.28) has a solution. As observed at the end of Sect. 4, this will lead to the proof of the existence of a solution to problem (4.11).

Let $\Omega$ and $\Omega^{\prime}$ be bounded open sets in $\mathbb{R}^{d}$ with Lipschitz boundary and with $\bar{\Omega} \subset \Omega^{\prime}$, and let $f: \mathbb{R}^{d} \rightarrow \mathbb{R}$ be the function defined by (4.12). We begin by proving a lower semicontinuity result.
Theorem 6.1. Let $\Phi \in L^{1}\left(\Omega^{\prime} ; \mathbb{R}^{d}\right)$ and let $\Gamma_{0} \subset \Omega^{\prime}$ be a Borel set with $\mathcal{H}^{d-1}\left(\Gamma_{0}\right)$ $<+\infty$. Then the functional $\mathcal{F}_{\Gamma_{0}}^{\Phi}: G B V_{\star}\left(\Omega^{\prime}\right) \rightarrow[0,+\infty]$ defined by

$$
\mathcal{F}_{\Gamma_{0}}^{\Phi}(v):=\int_{\Omega^{\prime}} f(\nabla v+\Phi) d x+\left|D^{c} v\right|\left(\Omega^{\prime}\right)+\int_{J_{v} \backslash \Gamma_{0}}|[v]| \wedge 1 d \mathcal{H}^{d-1}
$$

is lower semicontinuous with respect to the convergence in measure on $\Omega^{\prime}$.
Proof. Let us fix a bounded open set $U \subset \mathbb{R}^{d}$. The first step in the proof is to show that the functional

$$
\mathcal{F}_{0, U}(v):=\int_{U} f(\nabla v) d x+\left|D^{c} v\right|(U)+\int_{J_{v} \cap U}|[v]| \wedge 1 d \mathcal{H}^{d-1}
$$

is lower semicontinuous on $G B V_{\star}(U)$ with respect to the convergence in measure on $U$. Let $\left(u_{k}\right)_{k} \subset G B V_{\star}(U)$ be a sequence converging in measure to some $u \in G B V_{\star}(U)$ such that $\left(\mathcal{F}_{0, U}\left(u_{k}\right)\right)_{k}$ is bounded. For every $m>0$ the sequence of truncations $\left(u_{k}^{(m)}\right)_{k}$ converges to $u^{(m)}$ in measure on $U$ and $\mathcal{F}_{0, U}\left(u_{k}^{(m)}\right) \leq \mathcal{F}_{0, U}\left(u_{k}\right)$ by Proposition 2.6.

Let us show that $\left(u_{k}^{(m)}\right)_{k}$ is bounded in $B V(U)$ by a constant depending on $m$. The first inequality in (4.14) implies that

$$
\int_{U}\left|\nabla u_{k}^{(m)}\right| d x+\left|D^{c} u_{k}^{(m)}\right|(U)+\int_{J_{u_{k}^{(m)}}}\left|\left[u_{k}^{(m)}\right]\right| d \mathcal{H}^{d-1} \leq c_{m} \mathcal{F}_{0, U}\left(u_{k}\right)+\frac{1}{2} \mathcal{L}^{d}(U)
$$

where $c_{m}:=1+2 m$, hence the boundedness of $\left(u_{k}^{(m)}\right)_{k}$ in $B V(U)$ follows from (2.5).

By [4, Theorem 2.1] we have

$$
\begin{equation*}
\mathcal{F}_{0, U}\left(u^{(m)}\right) \leq \liminf _{k \rightarrow \infty} \mathcal{F}_{0, U}\left(u_{k}^{(m)}\right) \leq \liminf _{k \rightarrow \infty} \mathcal{F}_{0, U}\left(u_{k}\right) \tag{6.1}
\end{equation*}
$$

Passing to the limit as $m \rightarrow+\infty$ and using Propositions 2.6 and 2.9 we obtain

$$
\mathcal{F}_{0, U}(u) \leq \liminf _{k \rightarrow \infty} \mathcal{F}_{0, U}\left(u_{k}\right)
$$

which concludes the proof of the lower semicontinuity of $\mathcal{F}_{0, U}$ on $G B V_{\star}(U)$ with respect to the convergence in measure on $U$.

Given $\xi \in \mathbb{R}^{d}$, for every $v \in G B V_{\star}(U)$ let

$$
\mathcal{F}_{\xi, U}(v):=\int_{U} f(\nabla v+\xi) d x+\left|D^{c} v\right|(U)+\int_{J_{v}}|[v]| \wedge 1 d \mathcal{H}^{d-1}
$$

Since $\mathcal{F}_{\xi, U}(v)=\mathcal{F}_{0, U}\left(v+\ell_{\xi}\right)$, where $\ell_{\xi}(x):=\xi \cdot x$, we deduce that $\mathcal{F}_{\xi, U}$ is lower semicontinuous on $G B V_{\star}(U)$ with respect to the convergence in measure on $U$.

To prove a similar result for $\mathcal{F}_{\Gamma_{0}}^{\Phi}$ we fix a sequence $\left(v_{k}\right)_{k} \subset G B V_{\star}\left(\Omega^{\prime}\right)$ which converges in measure on $\Omega^{\prime}$ to a function $v \in G B V_{\star}\left(\Omega^{\prime}\right)$, and an increasing sequence $\left(K_{j}\right)_{j}$ of compact subsets of $\Gamma_{0}$ such that $\mathcal{H}^{d-1}\left(\Gamma_{0} \backslash K_{j}\right) \rightarrow 0$. It is easy to construct a sequence $\left(\Phi_{j}\right)_{j}$ of piecewise constant functions converging
to $\Phi$ in $L^{1}\left(\Omega^{\prime} ; \mathbb{R}^{d}\right)$ such that for every $j$ there exists a partition $U_{j}^{1}, \ldots, U_{j}^{i_{j}}, N_{j}$ of $\Omega^{\prime} \backslash K_{j}$, with $U_{j}^{i}$ open and $\mathcal{H}^{d-1}\left(N_{j}\right)<+\infty$, such that $\Phi_{j}=\xi_{j}^{i}$ in $U_{j}^{i}$ for suitable constant vectors $\xi_{j}^{i} \in \mathbb{R}^{d}$. It is not restrictive to assume also that $\mathcal{H}^{d-1}\left(J_{v} \cap N_{j}\right)=0$. By the previous step of the proof, for every $j$ we have

$$
\begin{equation*}
\mathcal{F}_{K_{j}}^{\Phi_{j}}(v)=\sum_{i=1}^{i_{j}} \mathcal{F}_{\xi_{j}^{i}, U_{j}^{i}}(v) \leq \sum_{i=1}^{i_{j}} \liminf _{k \rightarrow \infty} \mathcal{F}_{\xi_{j}^{i}, U_{j}^{i}}\left(v_{k}\right) \leq \liminf _{k \rightarrow \infty} \mathcal{F}_{K_{j}}^{\Phi_{j}}\left(v_{k}\right) . \tag{6.2}
\end{equation*}
$$

Since $f$ is Lipschitz continuous with constant 1 , for every $u \in G B V_{\star}\left(\Omega^{\prime}\right)$ we have

$$
\left|\mathcal{F}_{K_{j}}^{\Phi_{j}}(u)-\mathcal{F}_{\Gamma_{0}}^{\Phi}(u)\right| \leq\left\|\Phi_{j}-\Phi\right\|_{L^{1}\left(\Omega^{\prime} ; \mathbb{R}^{d}\right)}+\mathcal{H}^{d-1}\left(\Gamma_{0} \backslash K_{j}\right),
$$

hence

$$
\mathcal{F}_{\Gamma_{0}}^{\Phi}(v) \leq \liminf _{k \rightarrow \infty} \mathcal{F}_{\Gamma_{0}}^{\Phi}\left(v_{k}\right)+2\left(\left\|\Phi_{j}-\Phi\right\|_{L^{1}\left(\Omega^{\prime} ; \mathbb{R}^{d}\right)}+\mathcal{H}^{d-1}\left(\Gamma_{0} \backslash K_{j}\right)\right) .
$$

Passing to the limit as $j \rightarrow \infty$ we obtain the lower semicontinuity inequality along the sequence $\left(v_{k}\right)_{k}$.

We are now ready to prove the existence of a solution to the minimum problem (4.28).

Theorem 6.2. Let $w \in H^{1}\left(\Omega^{\prime}\right)$, let $\Phi \in L^{1}\left(\Omega^{\prime} ; \mathbb{R}^{d}\right)$, and let $\Gamma_{0} \subset \Omega^{\prime}$ be a Borel set with $\mathcal{H}^{d-1}\left(\Gamma_{0}\right)<+\infty$. Then the minimum problem

$$
\begin{equation*}
\min _{\substack{v \in G B V_{\star}\left(\Omega^{\prime}\right) \\ v=w \text { a.e. in } \Omega^{\prime} \backslash \Omega}}\left\{\int_{\Omega^{\prime}} f(\nabla v+\Phi) d x+\left|D^{c} v\right|\left(\Omega^{\prime}\right)+\int_{J_{v} \backslash \Gamma_{0}}|[v]| \wedge 1 d \mathcal{H}^{d-1}\right\} \tag{6.3}
\end{equation*}
$$

has a solution.
Proof. Since $\mathcal{F}_{\Gamma_{0}}^{\Phi}$ coincides with the functional $\mathcal{G}_{\Gamma_{0}}^{g}$ introduced in (5.2), with $g(x, \xi):=f(\xi+\Phi(x))$, and by (4.14) $g$ satisfies (5.1), we can apply Theorem 5.5 and obtain that there exist a minimizing sequence $\left(u_{k}\right)_{k} \subset G B V_{\star}\left(\Omega^{\prime}\right)$, with $u_{k}=w \mathcal{L}^{d}$-a.e. in $\Omega^{\prime} \backslash \Omega$, and a continuous function $\psi: \mathbb{R}_{+} \rightarrow \mathbb{R}_{+}$, with $\psi(t) \rightarrow+\infty$ as $t \rightarrow+\infty$, such that (3.26) and (3.27) hold. Then by the Compactness Theorem 3.11 there exist a subsequence, not relabelled, and a function $u \in G B V_{\star}\left(\Omega^{\prime}\right)$ such that $u_{k} \rightarrow u \mathcal{L}^{d}$-a.e. in $\Omega^{\prime}$. By the Semicontinuity Theorem 6.1 we obtain that

$$
\begin{aligned}
& \int_{\Omega^{\prime}} f(\nabla u+\Phi) d x+\left|D^{c} u\right|\left(\Omega^{\prime}\right)+\int_{J_{u} \backslash \Gamma_{0}}|[u]| \wedge 1 d \mathcal{H}^{d-1} \\
& \quad \leq \underset{k}{\liminf }\left(\int_{\Omega^{\prime}} f\left(\nabla u_{k}+\Phi\right) d x+\left|D^{c} u_{k}\right|\left(\Omega^{\prime}\right)+\int_{J_{u_{k}} \backslash \Gamma_{0}}\left|\left[u_{k}\right]\right| \wedge 1 d \mathcal{H}^{d-1}\right) .
\end{aligned}
$$

Since $\left(u_{k}\right)_{k}$ is a minimizing sequence and $u=w \mathcal{L}^{d}$ - a.e. in $\Omega^{\prime} \backslash \Omega$ we conclude that $u$ is a solution of the minimum problem (6.3).

We now show that the minimum problem (4.11) has a solution.

Corollary 6.3. Let $\Gamma_{0} \subset \bar{\Omega}$ be a Borel set with $\mathcal{H}^{d-1}\left(\Gamma_{0}\right)<+\infty$, let $w_{0}, w \in$ $H^{1}(\Omega)$, and let $\left(u_{0}, e_{0}, p_{0}\right) \in \mathcal{A}\left(\Gamma_{0}, w_{0}\right)$. Then the minimum problem

$$
\begin{equation*}
\min _{\substack{\Gamma \text { Borel, } \Gamma_{0} \subset \Gamma \subset \bar{\Omega} \\(u, e, p) \in \mathcal{A}(\Gamma, w)}}\left\{\frac{1}{2} \int_{\Omega}|e|^{2} d x+\left|p-p_{0}\right|(\bar{\Omega} \backslash \Gamma)+\mathcal{H}^{d-1}\left(\Gamma \backslash \Gamma_{0}\right)\right\} \tag{6.4}
\end{equation*}
$$

has a solution.
Proof. By the equivalence results proved in Lemmas 4.1 and 4.2, the conclusion follows from Theorem 6.2.

We conclude the paper with two results which show that in general we cannot find a solution $v$ to the minimum problem (6.3) with $v \in B V\left(\Omega^{\prime}\right)$. In the following proposition we show that this may happen even if $w=0$.

Proposition 6.4. Assume that $d \geq 2$. Then there exist a function $\Phi \in L^{1}\left(\Omega^{\prime} ; \mathbb{R}^{d}\right)$ and a Borel set $\Gamma_{0} \subset \bar{\Omega}$, with $\mathcal{H}^{d-1}\left(\Gamma_{0}\right)<+\infty$, such that the solution of the minimum problem (6.3) with $w=0$ is unique, belongs to $L^{1}\left(\Omega^{\prime}\right)$, but does not belong to $B V\left(\Omega^{\prime}\right)$.

Proof. It is not restrictive to assume that $0 \in \Omega$ and that $\Omega^{\prime}$ is connected. Let $0<R<1$ be such that $[-R, R]^{d} \subset \Omega$, let $\psi, h:(0, R] \rightarrow[0,+\infty)$ be defined by

$$
\begin{equation*}
\psi(r):=\frac{1}{r}-\frac{1}{R} \quad \text { and } \quad h(r):=r^{2} \quad \text { for every } r \in(0, R], \tag{6.5}
\end{equation*}
$$

and let $\Omega_{R, h}:=\left\{\left(x_{1}, \ldots, x_{d}\right) \in(-R, R)^{d}: 0<x_{1}<R,\left|x_{2}\right|<h\left(x_{1}\right)\right\}$.
For every $x=\left(x_{1}, \ldots, x_{d}\right) \in \Omega^{\prime}$ we define

$$
v_{0}(x):= \begin{cases}\psi\left(x_{1}\right) & \text { if } x \in \Omega_{R, h} \\ 0 & \text { otherwise }\end{cases}
$$

By (6.5) it is easy to see that $v_{0} \in G B V_{\star}\left(\Omega^{\prime}\right) \cap L^{1}\left(\Omega^{\prime}\right), v_{0}=0$ in $\Omega^{\prime} \backslash \Omega, D^{c} v_{0}=$ $0, \mathcal{H}^{d-1}\left(J_{v_{0}}\right)<+\infty$, and $J_{v_{0}} \supset\left\{x \in(-R, R)^{d}: 0<x_{1}<R, x_{2}= \pm h\left(x_{1}\right)\right\}$. Since on this set $\left|\left[v_{0}\right](x)\right|=\psi\left(x_{1}\right)$, from (6.5) we deduce

$$
\int_{J_{v_{0}}}\left|\left[v_{0}\right]\right| d \mathcal{H}^{d-1}=+\infty
$$

which shows that $v_{0} \notin B V\left(\Omega^{\prime}\right)$.
Let us define $\Phi:=-\nabla v_{0}$ and $\Gamma_{0}:=J_{v_{0}}$. Then it is clear that $v_{0}$ is a solution of (6.3) since the value of the functional in $v_{0}$ is equal to zero. If $v$ is another solution we must have

$$
\int_{\Omega^{\prime}} f\left(\nabla v-\nabla v_{0}\right) d x+\left|D^{c} v\right|\left(\Omega^{\prime}\right)+\int_{J_{v} \backslash J_{v_{0}}}|[v]| \wedge 1 d \mathcal{H}^{d-1}=0 .
$$

This implies $\nabla v=\nabla v_{0} \mathcal{L}^{d}$-a.e. in $\Omega^{\prime}, D^{c} v=0$ in $\Omega^{\prime}$, and $\mathcal{H}^{d-1}\left(J_{v} \backslash J_{v_{0}}\right)=0$. Therefore $v-v_{0} \in W^{1,1}\left(\Omega^{\prime} \backslash \bar{J}_{v_{0}}\right), \nabla\left(v-v_{0}\right)=0 \mathcal{L}^{d}$-a.e. in $\Omega^{\prime}$, and $v-v_{0}=0$ $\mathcal{L}^{d}$-a.e. in $\Omega^{\prime} \backslash \Omega$. Since $\Omega^{\prime} \backslash \bar{J}_{v_{0}}$ is connected, we conclude that $v=v_{0} \mathcal{L}^{d}$-a.e. in $\Omega^{\prime}$.

In the following proposition we consider the case $\Phi=0$.

Proposition 6.5. Assume $d \geq 2$. Then there exist a function $w \in H^{1}\left(\Omega^{\prime}\right)$ and a Borel set $\Gamma_{0} \subset \bar{\Omega}$, with $\mathcal{H}^{d-1}\left(\Gamma_{0}\right)<+\infty$, such that the solution of the minimum problem (6.3) with $\Phi=0$ is unique, belongs to $L^{1}\left(\Omega^{\prime}\right)$, but does not belong to $B V\left(\Omega^{\prime}\right)$.

Proof. It is not restrictive to assume that $\Omega^{\prime}$ is connected. Since $\Omega$ has Lipschitz boundary, up to a change in the coordinate system, we may assume that there exist an open set $A \subset \mathbb{R}^{d-1}$, an interval $I \subset \mathbb{R}$, and a Lipschitz function $g: A \rightarrow I$ such that $\Omega \cap(A \times I)=\{(y, z) \in A \times I: z<g(y)\}$. It is not restrictive to assume that $\bar{A} \times \bar{I} \subset \Omega^{\prime}$. We fix a nonempty open set $A^{\prime} \subset \subset A$ and a sequence $\left(y_{k}\right)_{k \geq k_{0}} \subset A^{\prime}$ such that the balls in $\mathbb{R}^{d-1}$ of centre $y_{k}$ and radius $1 / k^{2}$ are pairwise disjoint and contained in $A^{\prime}$.

For every $y_{0} \in \mathbb{R}^{d-1}, h>0$, and $r>0$ let

$$
C_{r}^{h}\left(y_{0}\right):=\left\{(y, z) \in \mathbb{R}^{d-1} \times \mathbb{R}:\left|y-y_{0}\right|<r,\left|z-g\left(y_{0}\right)\right|<h\right\}
$$

be the cylinder of centre $\left(y_{0}, g\left(y_{0}\right)\right)$, height $2 h$, and radius $r$. For every pair of open sets $U$ and $V$, with $U \subset V \subset \mathbb{R}^{d}$, let

$$
\begin{equation*}
\operatorname{cap}(U, V)=\min \left\{\int_{V}|D u|^{2} d x: u \in H_{0}^{1}(V), u=1 \mathcal{L}^{d} \text {-a.e. in } U\right\} \tag{6.6}
\end{equation*}
$$

be the relative capacity of $U$ in $V$.
Let $L>0$ be the Lipschitz constant of $g$. We may assume that $C_{1 / k^{2}}^{L / k^{2}}\left(y_{k}\right)$ $\subset A^{\prime} \times I$ for every $k \geq k_{0}$. Since we have $\operatorname{cap}\left(C_{r}^{L r}\left(y_{k}\right), C_{1 / k^{2}}^{L / k^{2}}\left(y_{k}\right)\right) \rightarrow 0+$ as $r \rightarrow 0+$, there exists $r_{k}$ such that

$$
\begin{equation*}
0<r_{k}<1 / k^{2} \quad \text { and } \quad \operatorname{cap}\left(C_{r_{k}}^{L r_{k}}\left(y_{k}\right), C_{1 / k^{2}}^{L / k^{2}}\left(y_{k}\right)\right)<\frac{1}{k^{4 d-2}} \tag{6.7}
\end{equation*}
$$

Let $w_{k}$ be the solution of the minimum problem (6.6) with $U=C_{r_{k}}^{L r_{k}}\left(y_{k}\right)$ and $V=C_{1 / k^{2}}^{L / k^{2}}\left(y_{k}\right)$, extended to zero out of $C_{1 / k^{2}}^{L / k^{2}}\left(y_{k}\right)$, and let

$$
w:=\sum_{k=k_{0}}^{\infty} k^{2 d-3} w_{k} .
$$

By (6.7) the series converges in $H_{0}^{1}\left(\Omega^{\prime}\right)$, hence $w \in H_{0}^{1}\left(\Omega^{\prime}\right)$. Moreover

$$
\begin{align*}
& w=k^{2 d-3} \quad \mathcal{L}^{d} \text {-a.e. in } C_{r_{k}}^{L r_{k}}\left(y_{k}\right),  \tag{6.8}\\
& w=0 \quad \mathcal{L}^{d} \text {-a.e. in } \Omega^{\prime} \backslash\left(A^{\prime} \times I\right) . \tag{6.9}
\end{align*}
$$

Let $\Gamma_{1}:=\partial \Omega \cap\left(\bar{A}^{\prime} \times I\right) \backslash \bigcup_{k} C_{r_{k}}^{L r_{k}}\left(y_{k}\right), \Gamma_{2}:=\bigcup_{k} \partial C_{1 / k^{2}}^{L / k^{2}}\left(y_{k}\right) \cap \Omega$, and $\Gamma_{0}:=$ $\Gamma_{1} \cup \Gamma_{2}$. Then

$$
\begin{aligned}
& \mathcal{H}^{d-1}\left(\Gamma_{0}\right) \leq \mathcal{H}^{d-1}(\partial \Omega)+\sum_{k} \mathcal{H}^{d-1}\left(\partial C_{1 / k^{2}}^{L / k^{2}}\left(y_{k}\right)\right) \\
& \quad \leq \mathcal{H}^{d-1}(\partial \Omega)+\sum_{k} \omega_{d-1} 1 / k^{2(d-1)}+2 \sum_{k} \sigma_{d-2} L / k^{2(d-2)+2}<+\infty
\end{aligned}
$$

where $\omega_{d-1}$ is the $(d-1)$-measure of the unit ball in $\mathbb{R}^{d-1}$ and $\sigma_{d-2}$ is the (d -2 )-measure of its boundary.

Let $v_{0}: \Omega^{\prime} \rightarrow \mathbb{R}$ be defined by $v_{0}:=w$ in $\Omega^{\prime} \backslash \Omega, v_{0}:=k^{2 d-3}$ in $C_{1 / k^{2}}^{L / k^{2}}\left(y_{k}\right) \cap$ $\Omega$ and $v_{0}:=0$ in $\Omega \backslash \bigcup_{k} C_{1 / k^{2}}^{L / k^{2}}\left(y_{k}\right)$. Since $\nabla v_{0}=0 \mathcal{L}^{d}$-a.e. in $\Omega, \nabla v_{0}=\nabla w$ $\mathcal{L}^{d}$-a.e. in $\Omega^{\prime} \backslash \Omega, D^{c} v_{0}=0$ in $\Omega^{\prime}, J_{v_{0}} \subset \Gamma_{0}$, and $\mathcal{H}^{d-1}\left(\Gamma_{0}\right)<+\infty$, we have that $v_{0} \in G B V_{\star}\left(\Omega^{\prime}\right)$. We observe that the functional in the minimum problem (6.3) with $\Phi=0$ attains the value $\int_{\Omega^{\prime} \backslash \Omega} f(\nabla w) d x$ at $v_{0}$, hence $v_{0}$ is a solution to this problem.

If $v$ is another solution of (6.3) we must have $v=w$ in $\Omega^{\prime} \backslash \Omega$ and

$$
\int_{\Omega} f(\nabla v) d x+\left|D^{c} v\right|\left(\Omega^{\prime}\right)+\int_{J_{v} \backslash \Gamma_{0}}|[v]| \wedge 1 d \mathcal{H}^{d-1}=0 .
$$

This implies $\nabla v=0 \mathcal{L}^{d}$-a.e. in $\Omega, D^{c} v=0$ in $\Omega^{\prime}$, and $\mathcal{H}^{d-1}\left(J_{v} \backslash \Gamma_{0}\right)=0$. Hence $v \in W^{1,1}\left(\Omega^{\prime} \backslash \Gamma_{0}\right)$. Since $\nabla v=\nabla v_{0} \mathcal{L}^{d}$-a.e. in $\Omega^{\prime} \backslash \Gamma_{0}$ and $\Omega^{\prime} \backslash \Gamma_{0}$ is connected, from the equality $v=v_{0} \mathcal{L}^{d}$-a.e. in $\Omega^{\prime} \backslash \Omega$ we conclude that $v=v_{0}$ $\mathcal{L}^{d}$-a.e. in $\Omega^{\prime}$.

Since

$$
\begin{aligned}
& \int_{\Omega^{\prime}}\left|v_{0}\right| d x \leq \int_{\Omega^{\prime} \backslash \Omega}|w| d x+\sum_{k=k_{0}}^{\infty} k^{2 d-3} \mathcal{L}^{d}\left(C_{1 / k^{2}}^{L / k^{2}}\left(y_{k}\right)\right) \\
& \quad \leq \int_{\Omega^{\prime} \backslash \Omega}|w| d x+2 \omega_{d-1} L \sum_{k=k_{0}}^{\infty}\left(\frac{1}{k^{2}}\right)^{d-1} \frac{1}{k^{2}} k^{2 d-3}<+\infty
\end{aligned}
$$

we have that $v_{0} \in L^{1}\left(\Omega^{\prime}\right)$.
To prove that $v_{0} \notin B V(\Omega)$ we estimate the integral of the jump of $v_{0}$. Since $\left|\left[v_{0}\right]\right|=k^{2 d-3} \mathcal{H}^{d-1}$-a.e. on $\partial C_{1 / k^{2}}^{L / k^{2}}\left(y_{k}\right) \cap \Omega$, and the base of this cylinder is contained in $\Omega$, we have

$$
\int_{J_{v_{0}}}\left|\left[v_{0}\right]\right| d \mathcal{H}^{d-1} \geq \omega_{d-1} \sum_{k=k_{0}}^{\infty}\left(\frac{1}{k^{2}}\right)^{d-1} k^{2 d-3}=+\infty
$$

This shows that $v_{0} \notin B V\left(\Omega^{\prime}\right)$, since for every $v \in B V\left(\Omega^{\prime}\right)$ we have $[v] \in$ $L^{1}\left(J_{v}, \mathcal{H}^{d-1}\right)$.

Remark 6.6. By the equivalence results proved in Lemmas 4.1 and 4.2, if $\Gamma_{0}$ and $w$ are as in Proposition 6.5, then the minimum problem (6.4) corresponding to $w_{0}=0$ and $\left(u_{0}, e_{0}, p_{0}\right)=(0,0,0)$ has a unique solution $(u, e, p)$ with $u \notin B V(\Omega)$.

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