

A large fraction of hydrogen-rich supernova progenitors experience elevated mass loss shortly prior to explosion

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Submitted to ApJ

ABSTRACT

Spectroscopic detection of narrow emission lines traces the presence of circumstellar mass distributions around massive stars exploding as core-collapse supernovae. Transient emission lines disappearing shortly after the supernova explosion suggest that the material spatial extent is compact and implies an increased mass loss shortly prior to explosion. Here, we present a systematic survey for such transient emission lines (Flash Spectroscopy) among Type II supernovae detected in the first year of the Zwicky Transient Facility (ZTF) survey. We find that at least six out of ten events for which a spectrum was obtained within two days of the estimated explosion time show evidence for such transient flash lines. Our measured flash event fraction ($> 30\%$ at 95% confidence level) indicates that elevated mass loss is a common process occurring in massive stars that are about to explode as supernovae.

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Keywords: supernova:general - methods: observational - stars: mass-loss - stars: massive

1. INTRODUCTION

Massive stars ($M > 8 M_{\odot}$) explode as core-collapse supernovae (CC SNe; Smartt 2015; Gal-Yam 2017), and often experience mass loss from their outer layers due to stellar winds, binary interaction, or eruptive mass-loss events (see, e.g., Smith 2014 and references within). The mass lost by these stars forms distributions of circumstellar medium (CSM). The CSM properties depend on the mass-loss rate, the velocity of the flow, and the duration of the process.

When a massive star surrounded by CSM explodes as a CC SN, signatures of the CSM may manifest as spectroscopic features with a narrow width reflecting the mass-loss velocity, which is typically low compared to the expansion velocity of the supernova ejecta (a few hundreds of km s^{-1} vs. $\approx 10000 \text{ km s}^{-1}$). Such features often persist for more than two days from the explosion, which sets the extent of the material to $> 10^{14} \text{ cm}$, a scale far above that of the atmospheres of the largest supergiants. In Type IIn SNe (e.g., Schlegel 1990, Filippenko 1997, Gal-Yam 2017, Kiewe et al. 2012, Taddia et al. 2013, Nyholm et al. 2019), narrow hydrogen lines persist for weeks to years after explosion, indicating an extensive CSM distribution. Type Ibn events (e.g., Pastorello et al. 2016, Gal-Yam 2017, Hosseinzadeh et al. 2015, Karamehmetoglu et al. 2019) show strong emission lines of helium, suggesting recent mass loss from stripped progenitors. In both Types IIn and Ibn, there is evidence that in at least some cases, the mass-loss is generated by precursor events, prior to the SN explosion (e.g. Pastorello et al. 2007, Foley et al. 2007, Ofek et al. 2013, Ofek et al. 2014, Strotjohann et al. 2020).

If the CSM extension is confined to a relatively compact location around an exploding star, the explosion shock-breakout flash may ionize the CSM. The resulting recombination emission lines will be transient, persisting only until the SN ejecta overtakes and engulfs the denser parts of the CSM (supernovae with “flash ionized” emission lines; Gal-Yam et al. 2014). Such events later evolve spectroscopically in a regular manner, e.g., presenting photospheric spectra with broad P-Cygni line profiles.

Several serendipitous observation of such “flash features” in early supernova spectra were made over the years (e.g., Niemela et al. 1985, Garnavich & Ann 1994, Quimby et al. 2007). We define flash features here as transient narrow emission lines (of the order of $\approx 10^2 \text{ km s}^{-1}$) of highly ionised species (e.g.: He II, C III, N III, N IV) in the early phases of the supernova event (in

general, less than a week from the estimated explosion). Gal-Yam et al. (2014) presented very early observations of the Type IIb SN 2013cu, and noted that such flash features could be routinely observed by modern high-cadence SN surveys. These features reveal the composition of the pre-explosion mass loss, and hence probe the surface composition of the progenitor star, which is hard to measure by other means. This work motivated additional studies on such flash objects. For example, Yaron et al. (2017) presented a time-series of early spectra which they used to constrain the CSM distribution around the spectroscopically normal SN 2013fs. They show that the CSM was lost from the progenitor during the year prior to its explosion. Hosseinzadeh et al. (2018) studied the low-luminosity Type II event, SN 2016bkv, which showed early flash ionisation features. They suggest that its early light-curve bump implies a contribution from CSM interaction to the early light curve. Such interpretations motivate the systematic study of early light curves of Type II SNe with flash features to distinguish between possible contributions of CSM interaction versus shock cooling emission, for example, by testing the correlation of peak luminosity and/or rise time with the existence of flash features. Several theoretical investigations also focused on such events (e.g., Groh 2014, Dessart et al. 2017, Kochanek 2019, Moriya et al. 2017 and Boian & Groh 2020).

A systematic study of such transient signatures of CSM around SN II progenitor stars has been limited by the challenge of routinely observing CC SNe early enough (typically within less than a few days from explosion), before these features disappear. Khazov et al. (2016) conducted the first sample study of flash ionisation features in Type II SNe using data from the PTF and iPTF surveys. They gathered twelve objects showing flash ionisation features and estimated that more than $\sim 20\%$ of SNe II show flash ionisation features, although their analysis was limited by the heterogeneity of their data.

Routine observations of young (“infant”) SNe was one of the main goals of the ZTF survey (Gal-Yam 2019; Graham et al. 2019). Here, we present our systematic search and follow-up observations of infant Type II SNe from ZTF. We use a sample of 28 events collected during the first year of ZTF operation to place a lower limit on the fraction of SN progenitor stars embedded in CSM. Ten of these objects were spectroscopically observed within two days of the estimated explosion.

In section § 2, we describe the properties of our infant SN survey and the construction of our sample of SNe II. In § 3 we present our analysis, in § 4 we discuss our findings, and we conclude in § 5.

2. OBSERVATIONS AND SAMPLE CONSTRUCTION

2.1. Selecting infant SNe from the ZTF partnership stream

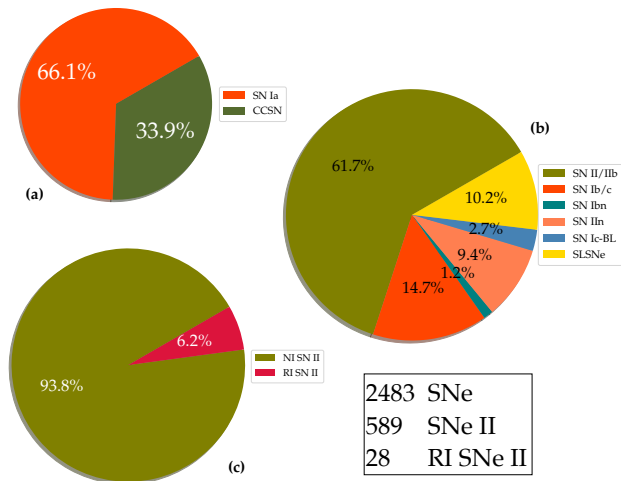


Figure 1. ZTF Spectroscopically-confirmed SN discovery statistics during 2018. (a) Most events (66%) are SNe Ia; CC SNe comprise about 34%. (b) The division among CC SN sub-classes (c) The fraction of real infant (RI) SNe II is 6.2% of the total Type II population. NI stands for the Non Infant SN II population (see text).

The Zwicky Transient Facility (ZTF) is a wide-field, high cadence, multiband survey that started operating in March 2018 (Bellm et al. 2019; Graham et al. 2019). ZTF imaging is obtained using the Samuel Oschin 48” Schmidt telescope at Palomar observatory (P48). ZTF observing time is divided into three programs: the public (MSIP) 3-day all-sky survey, partnership surveys, and Caltech programs. This paper is based on data obtained by the high-cadence partnership survey. As part of this program, during 2018, extra-galactic survey fields were observed in both the ZTF g - and r -bands 2-3 times per night, per band. New images were processed through the ZTF pipeline (Masci et al. 2019), and reference images, built by combining stacks of previous ZTF imaging in each band, were then subtracted using the Zackay et al. (2016) image subtraction algorithm (ZOGY). A 30s integration time was used in both g - and r -band exposures. A 5σ detection limit is adopted for estimating the limiting magnitude, typically reaching ~ 20.5 mag in r -band in a single observation.

We conducted our year-1 ZTF survey for infant SNe following the methodology of Gal-Yam et al. (2011). We selected potential targets via a custom filter running on the ZTF alert stream using the GROWTH Marshal platform (Kasliwal et al. 2019). The filter scheme was based on the criteria listed in Table 1.

Alerts that passed our filter (typically 50 – 100 alerts per day) were then visually scanned by a duty astronomer, in order to reject various artefacts (such as unmasked bad pixels or ghosts) and false positive signals (such as flaring M stars, CVs and AGN). Most spurious sources could be identified by cross-matching with additional catalogues, e.g., *WISE* IR photometry (Wright et al. 2010) to detect red M stars, the Gaia DR2 catalog (Gaia Collaboration et al. 2018), and catalogs from time-domain surveys such as the Palomar Transient Factory (PTF; Law et al. 2009) and the Catalina Real-Time Survey (CRTS; Drake et al. 2014) for previous variability of CVs and AGN.

Due to time-zone differences, our scanning team (located mostly at the Weizmann Institute in Israel and the Oskar Klein Center (OKC) in Sweden) could routinely monitor the incoming alert stream during the California night time. We aimed at triggering spectroscopic follow-up of promising infant SN candidates within hours of discovery (and thus typically within < 2 days from explosion), as well as *Swift* (Gehrels et al. 2004) Target-of-Opportunity (ToO) UV photometry.

2.2. Sample Construction

Figure 1 shows the SN Type distribution amongst the ~ 2500 spectroscopically-confirmed SNe gathered by ZTF between March and December 2018. About 34% are core-collapse events, and $\sim 62\%$ of those are of Type II. We can only place statistically meaningful constraints on the frequency of flash features among Type II SNe, since these mostly occur in this population. Hence, we choose to study only the SN II population from ZTF year 1.

Our infant SN program allowed us to obtain early photometric and spectroscopic follow-up of young SNe. However, we may have missed some relevant candidates. To ensure the completeness of our sample, we, therefore, inspected all spectroscopically classified SNe II (including subtypes IIn and IIb) from ZTF¹ using the ZTFquery package (Rigault 2018). We removed from this sample all events (the large majority) lacking a ZTF non-detection limit within 2.5 days prior to the first detection recorded on the ZTF Marshal. To include events in our

¹ between March 2018 and December 2018

Table 1. Filter criteria selecting infant SN candidates

<i>Stationary</i>	Reject solar-system objects using apparent motion
<i>Recent limit</i>	Require a non-detection limit within < 2.5 days from the first detection
<i>Extragalactic</i>	Reject alerts within 14 degrees from the Galactic plane
<i>Significant</i>	Require a ZOGY score of > 5
<i>Stellar</i>	Require a SG (star-galaxy) score ¹ of < 0.49

¹This parameter indicates whether the closest source in the PS1 catalogue is stellar.
See <https://zwockytransientfacility.github.io/ztf-avro-alert/schema.html>

final sample, we required that they show a significant and rapid increase in flux with respect to the last non-detection, as previously observed for very young SNe (e.g., Gal-Yam et al. 2014, Yaron et al. 2017). This excludes older events that are just slightly below our detection limit and are picked up by the filter when they slowly rise, or when weather conditions improve. We implemented a cut on the observed rise of Δr or $\Delta g > 0.5$ mag with respect to the recent limit in the same band, and labeled all events that satisfy this cut as “real infant” (RI; Fig. 1, panel C).

All in all, we gathered 43 candidates which fulfilled the RI criteria. Additional inspection led us to determine that 15 candidates were spurious (see Appendix A for details). Our final sample (Table 2) thus includes a total of 28 RI Type II SNe, or about 6.2% of all the SNe II found during 2018 by the ZTF survey. During its first year of operation (starting March 2018), ZTF obtained useful observations for our program during approximately 32 weeks, excluding periods of reference image building (initially), periods dedicated to Galactic observations, and periods of technical/weather closure. We find that the survey provided about one real infant SN II per week.

2.3. Spectroscopic Observations

Our goal was to obtain rapid spectroscopy of RI SN candidates following the methods of Gal-Yam et al. (2011). This was made possible using rapid ToO follow-up programs as well as on-request access to scheduled nights on various telescopes. During the scanning campaign, we applied the following criteria for rapid spectroscopic triggers. The robotic SEDm (see below) was triggered for all candidates brighter than a threshold magnitude of 19 mag in 2018. Higher-resolution spectra (using WHT, Gemini, or other available instruments) were triggered for events showing recent non-detection limits (within 2.5 d prior to first detection) as well as a significant rise in magnitude compared to a recent limit or within the observing night.

P60/SEDm—The Spectral Energy Distribution Machine (SEDm; Ben-Ami et al. 2012; Blagorodnova et al.

2018; Neill 2019) is a high-throughput, low-resolution spectrograph mounted on the 60” robotic telescope (P60; Cenko et al. 2006) at Palomar observatory. 65% of the time on the SEDm was dedicated to ZTF partnership follow up. SEDm data are reduced using an automated pipeline (Rigault et al. 2019). The co-location of the P60 and ZTF/P48 on the same mountain, as well as the P60 robotic response capability, enable very short (often same-night) response to ZTF events, sometimes very close to the time of first detection (e.g., see ZTF18abwls0i, below). However, the low resolution ($R \sim 100$) of the instrument limits our capability to characterise narrow emission lines. This, along with the overall sensitivity of the system, motivated us to try to obtain higher-resolution follow-up spectroscopy with other, larger, telescopes, particularly for all infant SNe detected below a magnitude cut of $r \sim 19$ mag.

P200/DBSP—We used the Double Beam SPectrograph (DBSP; Oke & Gunn 1982) mounted on the 5m Hale telescope at Palomar Observatory (P200) to obtain follow-up spectroscopy in either ToO mode or during classically scheduled nights. The default configuration used the 600/4000 grism on the blue side, the 316/7150 grating on the red side, along with the D55 dichroic, achieving a spectral resolution $R \sim 1000$. Spectra obtained with DBSP were reduced using the pyraf-dbsp pipeline (Bellm & Sesar 2016).

WHT-ISIS/ACAM—We obtained access to the 4.2m William Herschel Telescope (WHT) at the Observatorio del Roque de los Muchachos in La Palma, Spain, via the Optical Infrared Coordination Network for Astronomy (OPTICON²) program³. We used both single-slit spectrographs ISIS and ACAM (Benn et al. 2008) in ToO service observing mode. The delivered resolutions

² <https://www.astro-opticon.org/index.html>

³ Program IDs OPT/2017B/053, OPT/2018B/011, OPT/2019A/024, PI Gal-Yam

294 were $R \sim 1000$ and $R \sim 400$, respectively. Spectral data
295 were reduced using standard routines within IRAF⁴.

296 *Keck/LRIS*—We used the Low-Resolution Imaging
297 Spectrometer (LRIS; Oke et al. 1995) mounted on the
298 Keck-I 10m telescope at the W. M. Keck Observatory in
299 Hawaii in either ToO mode or during scheduled nights.
300 The data were reduced using the LRIS automated re-
301 duction pipeline Lpipe (Perley 2019).

302 *GMOS/Gemini*—We used the Gemini Multi-Object
303 Spectrograph (GMOS; Hook et al. 2004) mounted on
304 the Gemini North 8m telescope at the Gemini Obser-
305 vatory on Mauna Kea, Hawaii. All observations were
306 conducted at a small airmass ($\lesssim 1.2$). For each SN, we
307 obtained 2×900 s exposures using the B600 grating with
308 central wavelengths of 520 nm and 525 nm. The 5 nm
309 shift in the effective central wavelength was applied to
310 cover the chip gap, yielding a total integration time of
311 3600 s. A $1.0''$ -wide slit was placed on each target at
312 the parallactic angle. The GMOS data were reduced
313 following standard procedures using the Gemini IRAF
314 package.

315 *APO/DIS*—We used the Dual Imaging Spectrograph
316 (DIS) on the Astrophysical Research Consortium (ARC)
317 3.5 m telescope at Apache Point Observatory (APO)
318 during scheduled nights. The data were reduced using
319 standard procedures and calibrated to a standard star
320 obtained on the same night using the PyDIS package
321 (Davenport et al. 2016).

322 Spectra used for classification are presented in figures
323 14 and 15, and summarised in Table 8. All the data
324 presented in this paper will be made public on WIS-
325 eREP (Yaron & Gal-Yam 2012).

2.4. Photometry

328 The ZTF alert system (Patterson et al. 2018) provides
329 on the fly photometry (Masci et al. 2019) and astrom-
330 etry based on a single image for each alert. In order
331 to improve our photometric measurements (and in par-
332 ticular, to test the validity of non-detections just prior
333 to discovery) we performed forced PSF photometry at
334 the location of each event. As shown by Yaron et al.
335 (2019), the 95% astrometric scatter among ZTF alerts is
336 $\sim 0.44''$; for our events we had multiple detections, with
337 typically higher signal-to-noise ratio data around the SN

⁴ IRAF is distributed by the National Optical Astronomy Observa-
tories, which are operated by the Association of Universities for
Research in Astronomy, Inc., under cooperative agreement with
the National Science Foundation.

338 peak compared to the initial first detections. We there-
339 fore computed the median coordinates of all the alert
340 packages and performed forced photometry using this
341 improved astrometric location.

342 We used the pipeline developed by F. Masci and R.
343 Laher⁵ to perform forced PSF photometry at the median
344 SN centroid on the ZTF difference images available from
345 the IRSA database. For each light curve, we filtered out
346 measurements returned by the pipeline with non-valid
347 flux values.

348 We performed an additional quality cut on each light
349 curve by rejecting observations with a data quality pa-
350 rameter *scisigpix*⁶ that is more than five times the me-
351 dian absolute deviation (MAD) away from the median of
352 this parameter. We also removed faulty measurements
353 where the *infobitssci*⁷ parameter is not zero. According
354 to the Masci & Laher prescription, we rescaled the flux
355 errors by the square root of the χ^2 of the PSF fit esti-
356 mate in each image. We then corrected each measured
357 forced photometry flux value by the photometric zero
358 point of each image, as provided by the pipeline:

$$f_{zp,corrected} = f_{forced-phot} \times 10^{-0.4 \times z_p} \quad (1)$$

359 We determined our zero-flux baseline using forced
360 photometry observations obtained prior to the SN ex-
361 plosion. We calculated the median of these observations,
362 rejected outliers that are > 3 MAD away from the me-
363 dian, re-calculated the median and subtracted it from
364 our measured post-explosion flux values; these correc-
365 tions were small, of the order of $< 0.1\%$ of the super-
366 nova flux values.

367 If the ratio between the measured flux and the uncer-
368 tainty σ is below 3, we considered this measurement to
369 be a non-detection, and reported a 5σ upper limit. Oth-
370 erwise (if the flux to error ratio is above 3σ), we reported
371 the flux, magnitude and respective errors. We recov-
372 ered detections prior to the first detection by the real-
373 time pipeline using the forced photometry pipeline in 11
374 cases⁸. We redefined the first detection and last non-
375 detection according to the forced photometry pipeline
376 measurements in these cases.

⁵ <http://web.ipac.caltech.edu/staff/fmasci/ztf/forcedphot.pdf>

⁶ A parameter calculated by the pipeline that measures the pixel
noise in each science image

⁷ *infobitssci* is a quality assessment parameter on the processing
summary.

⁸ ZTF18aarqxbw, ZTF18aavpady, ZTF18aawyjjq,
ZTF18abcezmh, ZTF18abckutn, ZTF18abcptmt,
ZTF18abdbyisy, ZTF18abddjpt, ZTF18abokyfk, ZTF18abrljlc,
ZTF18abvvmfd

Table 2. Sample of Real Infant 2018 (28 objects)

IAU name (SN)	Internal ZTF name	Type ^a	Redshift z	Explosion JD Date [d]	Error [d]	First detection [d] ^b	Last non detection [d]	First spectrum [d]	Telescope/ instrument	Flash
2018grf	18abwlsoi	SN II ¹	0.054	2458377.6103	0.0139	0.0227	-0.8725	0.1407	P60/SEDm	✓
2018fzn	18abojpnr	SN I Ib ²	0.037	2458351.7068	0.0103	0.0102	-0.0103	0.1902	P60/SEDm	✗
2018dfi	18abffypq	SN I Ib ³	0.031	2458307.2540	0.4320	0.4320	-0.4320	0.6180	P200/DBSP	✓
2018cxn	18abckutn	SN II ⁴	0.041	2458289.8074	0.4189	0.0576	-0.0494	0.9406	P200/DBSP	✗
2018dfc	18abeajml	SN II ⁵	0.037	2458303.7777	0.0118	0.0213	-0.9806	1.0153	P60/SEDm	✓
2018fif	18abokyfk	SN II ⁶	0.017	2458350.9535	0.3743	-0.0635	-1.0525	1.0525	P200/DBSP	✓
2018gts	18abvmdf	SN II ⁷	0.030	2458375.1028	0.5551	-0.4688	-1.3648	1.5162	P60/SEDm	✓
2018cyg	18abdbyys	SN II ⁸	0.011	2458294.7273	0.2034	0.0297	0.0147	1.6727	WHT/ACAM	?
2018cug	18abcptmt	SN II ⁹	0.050	2458290.9160	0.0250	-0.0066	-0.0670	1.7960	P60/SEDm	✓
2018egh	18abgqvww	SN II ¹⁰	0.038	2458312.7454	0.4351	0.9846	0.0931	1.8236	WHT/ISIS	?
2018bqs	18aarpttw	SN II ¹¹	0.047	2458246.8133	0.0071	0.0087	-0.9926	2.0867	APO/DIS	✗
2018fsm	18absldfl	SN II ¹²	0.035	2458363.4226	0.4565	0.4564	-0.4564	2.3674	P60/SEDm	✗
2018bge	18aaqkoyr	SN II ¹³	0.024	2458243.1671	0.5180	0.5179	-0.5180	2.5169	P200/DBSP	✗
2018leh	18adbmrug	SN I In ¹⁴	0.024	2458481.7505	0.9485	0.9485	-0.9485	3.6985	WHT/ISIS	✓
2018iua	18acploez	SN II ¹⁵	0.042	2458439.9877	0.9784	0.9783	-0.9783	3.7933	P60/SEDm	✗
2018gvn	18abyvenk	SN II ¹⁶	0.043	2458385.6198	0.0011	0.0012	-0.8565	6.1122	P60/SEDm	✗
2018clq	18aatlfus	SN II ¹⁷	0.045	2458248.8967	0.9564	0.9564	-0.9564	6.9274	P60/SEDm	✗
2018ccp	18aawyjjq	SN II ¹⁸	0.040	2458263.7743	0.1241	0.0106	-0.8684	8.1087	P60/SEDm	✗
2018lth	18aayxxew	SN II ¹⁹	0.061	2458278.6531	0.9154	0.0509	-1.9102	8.1589	Keck/LRIS	✗
2018inm	18achtnvk	SN II ²⁰	0.040	2458432.9113	0.6895	1.9927	1.9497	9.0137	P60/SEDm	✗
2018iwe	18abufaej	SN II ²¹	0.062	2458368.8561	0.0179	0.0179	-0.0179	12.0159	P60/SEDm	✗
2018fso	18abrlljc	SN II ²²	0.050	2458357.6987	0.8255	-0.0177	-0.9157	14.0113	P60/SEDm	✗
2018efd	18abgrbjb	SN I Ib ²³	0.030	2458312.8922	0.3938	0.8568	0.8244	14.9388	P60/SEDm	✗
2018cyh	18abcezmh	SN II ²⁴	0.057	2458286.3752	0.6050	0.4348	0.3898	16.5678	P60/SEDm	✗
2018ltg	18aarqxbw	SN II ²⁵	0.048	2458241.4360	3.4950	3.4950	-3.4950	37.5310	P200/DBSP	✗
2018lti	18abddjpt	SN II ²⁶	0.070	2458294.6217	0.1224	0.1693	-0.7917	40.2333	P60/SEDm	✗
2018efj	18abimhfu	SN II ²⁷	0.050	2458320.6574	0.0210	0.0096	-0.9028	42.0096	P60/SEDm	✗
2018cfj	18aavpady	SN II ²⁸	0.047	2458256.4531	0.4771	0.4771	-0.4771	55.0469	Keck/LRIS	✗

¹ Fremling et al. (2018a)	² Fremling et al. (2018g)	³ Hiramatsu et al. (2018)	⁴ Fremling & Sharma (2018a)	⁵ Fremling & Sharma (2018b)
⁶ Gal-Yam et al. (2018)	⁷ Fremling et al. (2018b)	⁸ Fremling & Sharma (2018c)	⁹ Fremling & Sharma (2018d)	¹⁰ Bruch (2020)
¹¹ Bruch (2020)	¹² Fremling et al. (2018c)	¹³ Prentice (2018)	¹⁴ Dugas et al. (2019)	¹⁵ Bruch (2020)
¹⁶ Fremling et al. (2018d)	¹⁷ Fremling et al. (2018h)	¹⁸ Fremling et al. (2019)	¹⁹ Bruch (2020)	²⁰ Fremling et al. (2018e)
²¹ Bruch (2020)	²² Fremling et al. (2018i)	²³ Fremling et al. (2018j)	²⁴ Bruch (2020)	²⁵ Bruch (2020)
²⁶ Bruch (2020)	²⁷ Fremling et al. (2018f)	²⁸ Bruch (2020)		

^a TNS Classification reports are referenced at the end of the table

^b All times reported relative to the estimated explosion date in fractional days

377 We present our photometry for all RI objects in Ta-
378 ble 3⁹.

379 3. ANALYSIS AND RESULTS

⁹ The light curves presented in this table are not corrected for MW extinction nor for redshift. The absolute magnitude is calculated using the package `Distance` from `astropy` (Price-Whelan et al. 2018)

380 In this section, we study the 28 RI SNe that passed
381 our selection criteria, excluding spurious candidates (see
382 Appendix A for details). In order to measure the frac-
383 tion of objects showing flash features and thus evidence
384 for CSM, we estimated the explosion time based on ZTF
385 forced photometry light curves. We then defined sub-
386 samples based on the SN age (relative to the estimated
387 explosion) when the first spectrum was obtained.

Table 3. Forced photometry of the RI sample

IAU name	ZTF name	Filter	JD [day]	Flux [10^{-8} mJy]	Flux error [10^{-8} mJy]	Apparent magnitude [AB mag]	Absolute magnitude [AB mag]	Magnitude error [AB mag]
...
18bge	ZTF18aaqkoyr	r	2458260.6754	7.1542	0.1350	17.86	-17.21	0.02
18bge	ZTF18aaqkoyr	r	2458260.6830	7.0166	0.1336	17.89	-17.19	0.02
18ccp	ZTF18aawyjjq	r	2458261.8319	-1.0936	0.1595	99.00	nan	nan
18ccp	ZTF18aawyjjq	r	2458261.8377	-0.5241	0.1723	99.00	nan	nan
18ccp	ZTF18aawyjjq	r	2458261.8387	0.1034	0.1696	99.00	nan	nan
...

NOTE—This table includes the flux measurements returned by the forced photometry pipeline. In this table, we report the last non detections within 2.5 days from the first Marshal detection and all the measurements which follow. The full version of this table is electronic. Light curves are plotted in the annex, Figures 12 and 13.

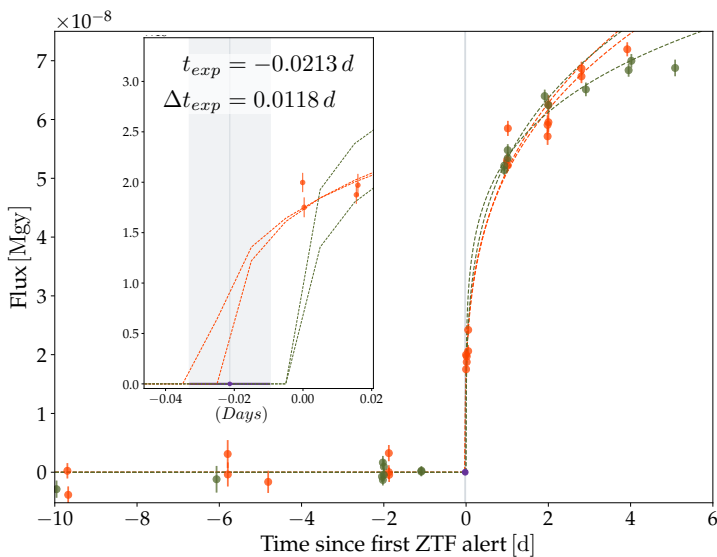


Figure 2. Early light curve fits used to determine the explosion date for SN 2018dfc. Power-law fits to the observations during the first 2 or 5 days are shown in both the g (green points) and r (red points) bands. The mean and standard deviation of the fits (inset) are adopted as the explosion time and the error. The time origin is defined as the time of the first alert from ZTF.

3.1. Explosion time estimation

In order to estimate the explosion time, which we define here as the time of zero flux, we fitted the following general power law to our flux measurements:

$$f(t) = a \times (t - t_{exp})^n \quad (2)$$

using the routine `curvefit` within the `astropy` python package (Astropy Collaboration et al. 2013). We fitted the first two days of data following the first detection as well as the first 5 days (see Fig. 2, for example) in both the g and r -bands. The estimated explosion time

is taken as the weighted mean¹⁰ of the four fits, and we adopted the standard deviation as the error on this value. In ten cases, however, there were not enough data in either band to perform the fit. In those cases, we set the explosion date as the mean between the last non detection and the first detection (Fig. 12). In all but four cases the estimated explosion date (EED) is within less than a day from the first detection (Fig. 3; Table 2).

3.2. Peak magnitude

Following Khazov et al. (2016), we also tested whether events showing flash features are, on average, more luminous. As shown in Table 2, the relevant events to consider are only those with relatively early spectra. We therefore compute the peak magnitude of all seventeen events with a first spectrum obtained within seven days from explosion. In the literature, we rarely found flash ionisation features which last more than a week from the EED. A first spectrum obtained a week after the EED could miss potential flash features. We hence chose the seven-day sub-sample in order to increase our pool of objects for this analysis while maintaining a realistic estimation of the percentage of flash ionisation events. We used the forced photometry lightcurves to evaluate the peak magnitude. We fitted a polynomial of order 3 to the flux measurements over several intervals of time. The lower bound of these fits is within the first few days from explosion time and the upper bound between 10 to 40 days after the estimated explosion time (Fig. 4). We varied randomly the lower and upper boundaries and repeated the fit a hundred times. We adopted the mean

¹⁰ Each fit is weighted according to the value of the fit on the estimated explosion time.

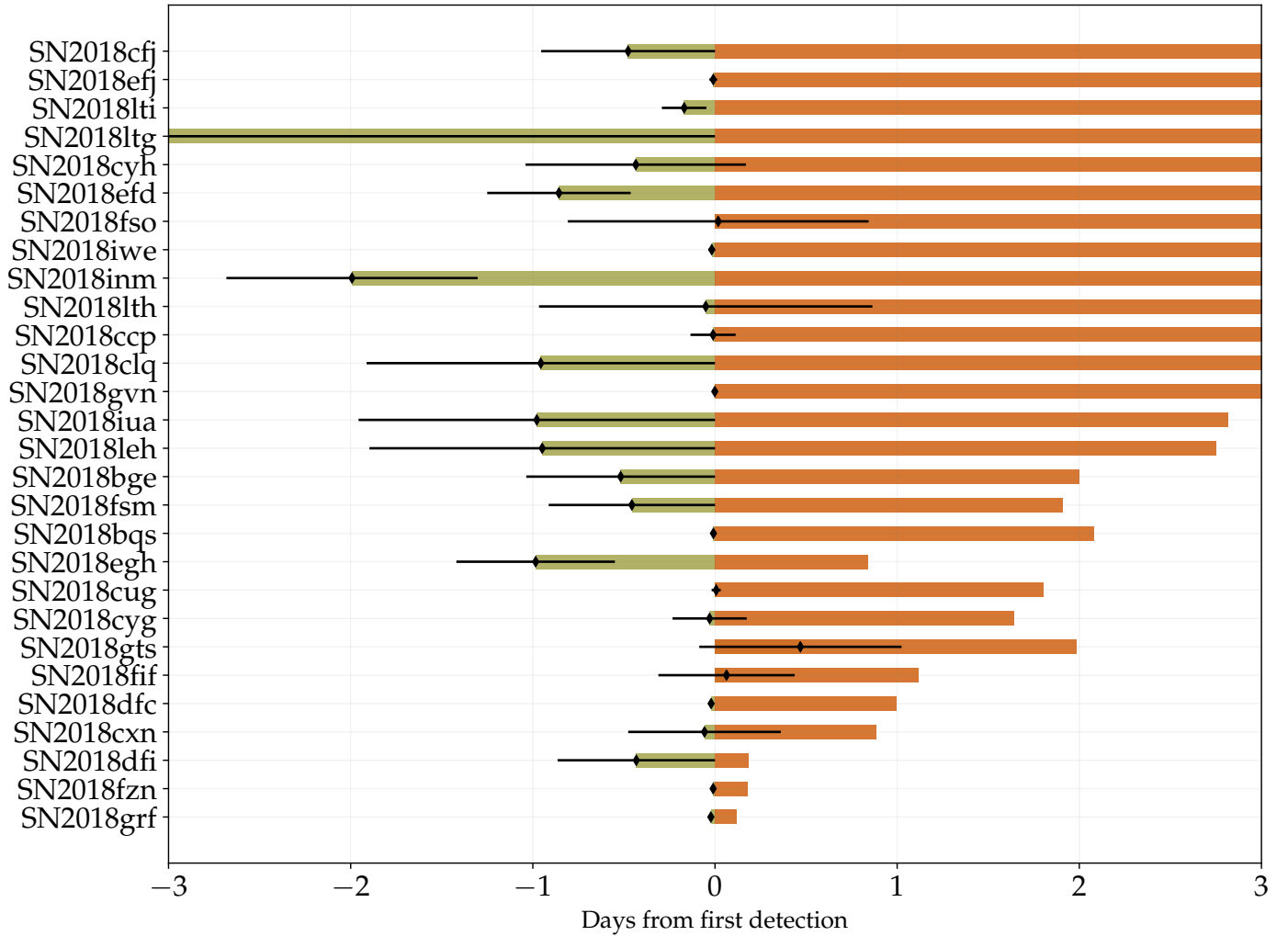


Figure 3. A graphic summary of the sample timeline, from the estimated explosion date (green) to the time of the first spectrum (red). The x-axis origin (“0” time) corresponds to the first photometric detection of each candidate. The black diamonds correspond to the estimated explosion time. SN 2018ltg was included in the sample of RI SNe II since its non-detection limit from the Marshal alert system was $< 2.5 d$ even though the explosion time estimation from the forced photometry lightcurve puts the last limit more than three days earlier.

428 and standard deviation of peak times obtained as the
 429 peak date and its error (vertical grey band in Fig. 4) and
 430 took the mean and standard deviation of the flux value
 431 as the the peak flux and error (horizontal grey band in
 432 Fig. 4). The absolute peak magnitude is computed as:

$$M_{peak} = m_{peak} - dm - A_{\lambda} \quad (3)$$

433 with dm , the distance modulus and A_{λ} the milky way
 434 extinction. We report these values for each event in
 435 each available band in Table 4. We obtained the dm
 436 using the python package `astropy.cosmology` (Price-
 437 Whelan et al. 2018) with the Planck18 cosmology. The

438 extinction was calculated using the packages `sdfmap`¹¹
 439 to estimate $E(B - V)$ and `extinction` (Barbary 2016)
 440 for A_{λ} . We assumed $R_V = 3.1$ and the Cardelli et al.
 441 (1989) extinction law. The errors on the absolute mag-
 442 nitude were calculated with:

$$\delta_{peak} = \sqrt{(\delta m_{peak})^2 + (\delta dm)^2} \quad (4)$$

443 with δm_{peak} , the error from the fit. We assume here that
 444 the error on the distance modulus is linear with the red-
 445 shift, hence : $\delta dm = \frac{\delta z \times dm}{z}$. Redshift errors were gath-
 446 ered from NED, when available. We estimated the red-

¹¹ <https://github.com/kbarbary/sdfmap>

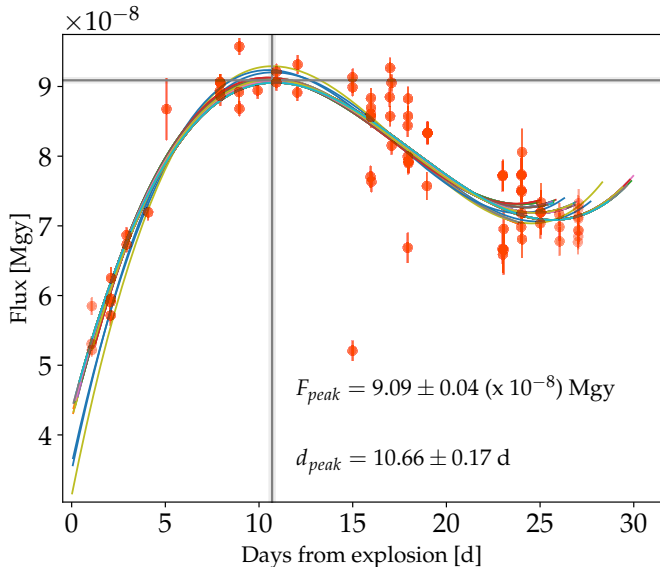


Figure 4. Example of the peak estimation in the red band for SN2018dfc. The different curves correspond to a polynomial of order 3 fitted over the time intervals noted in the legend. The cross corresponds to the peak date and flux estimated from the mean of all the values obtained, and the grey bands note the estimated errors, see text for details.

447 shift of the remaining supernovae by fitting a gaussian
 448 shaped line to narrow H_{α} emission line. We favoured
 449 spectra contaminated by host galaxy lines. We used
 450 the package `minuit` (Dembinski et al. 2020) to fit the
 451 H_{α} line. We remark that the redshift errors are bigger
 452 whenever we were using low-resolution SEDm spectra.

3.3. Early spectroscopy

454 We sorted the 28 RI SNe in our sample according to
 455 the difference between the estimated explosion time and
 456 the time of the first spectrum (Table 2, “First spectrum”
 457 column; Fig. 3). From previous work (Gal-Yam et al.
 458 2014, Yaron et al. 2017, Khazov et al. 2016), we know
 459 that flash features are typically present from the time
 460 of explosion up to several days later. We, therefore,
 461 defined a sub-sample including events with spectra ob-
 462 tained within 2 d from explosion (top of Table 2), Which
 463 includes about one third of the total sample (ten ob-
 464 jects).

465 Throughout the 2018 campaign, we found that seven
 466 infant supernovae of Type II show flash features (Ta-
 467 ble 2; Fig. 5). Two additional infant objects were
 468 marked as potential flash events (Fig. 8; see below).
 469 Four of the seven confirmed flashers had their first spec-
 470 trum obtained with SEDm.

471 The two-day sub-sample includes six events showing
 472 flash features, two potential flashers (SN 2018cyg and

473 SN2018egh, Fig. 8), and two events which have high
 474 signal to noise early spectra that show no flash features
 475 (Fig. 7). One object, SN 2018leh, shows flash features
 476 but its first spectrum was obtained > 3 days after explo-
 477 sion, see Table 2, Fig. 6.

3.3.1. The Flash events

478
 479 The identification of flash features in this work is
 480 focused on the spectral range surrounding the strong
 481 He II emission line at 4686 Å. This follows previous work
 482 (Khazov et al. 2016) and is also supported by theoret-
 483 ical model grids (Boian & Groh 2020) which show that
 484 this feature is ubiquitous in early spectra (< 2 d). We
 485 chose not to use hydrogen lines as a marker for flash fea-
 486 tures since contribution from host galaxy lines is likely
 487 to complicate the analysis.

488 In previous well-studied cases of events with high-
 489 quality early spectra, such as SN 2013fs (Yaron et al.
 490 2017) and SN 2013cu (Gal-Yam et al. 2014), the He II
 491 $\lambda 4686$ line is very prominent with a profile that is often
 492 well described by a narrow core with broad Lorentzian
 493 wings, which could be attributed to electron scattering
 494 within the CSM (Huang & Chevalier 2018).

495 As discussed in detail by Soumagnac et al. (2019),
 496 while the spectra of such events evolve with time, the
 497 strong He II emission line is replaced by a ledge-shaped
 498 feature that is probably composed of blended high-
 499 ionization lines of C, N and O. The He II line and some
 500 other lines (e.g. C III or N III) are sometimes detected
 501 as a narrow emission line on top of the ledge-shaped
 502 feature (see Fig 5 and Fig. 7 of Soumagnac et al. 2019).

503 As several of our early spectra were obtained with the
 504 low-resolution SEDm instrument (in particular those of
 505 SN 2018grf, SN 2018gts, and SN 2018cug), we could
 506 not easily differentiate between the various manifesta-
 507 tions of the excess emission around 4686 Å. We there-
 508 fore adopted the detection of excess emission around
 509 this wavelength as our criterion for defining an object
 510 as having flash features. Analysis of the cases where we
 511 have both early SEDm spectra and high spectral res-
 512 olution data from larger telescopes (e.g., SN 2018dfc),
 513 confirm the nature of the emission we see in the SEDm
 514 spectra and support our approach (Fig. 5).

515 SN 2018leh is the seventh object which displayed flash
 516 features. It does not belong to the sub-sample we consid-
 517 ered for this study since its first spectrum was obtained
 518 ≈ 3.7 days after the estimated explosion time. This ob-
 519 ject shows the Balmer emission lines H_{α} , H_{β} and H_{γ}
 520 that persist for an extended period of time, ≈ 10 days.
 521 This led us to classify this event as a SN IIn. The first
 522 spectrum also shows a strong He II line, which disap-
 523 peared about ten days later, see Fig. 6. The transient

Table 4. Peak absolute magnitudes of the 17 objects within the 7-day spectroscopic sub-sample

IAU name	filter	$z \pm \delta z$	dm [AB mag]	$m_{peak} \pm \delta m_{peak}$ [AB mag]	$d_{peak} \pm \delta d_{peak}$ [days]*	extinction [AB mag]	$M_{peak} \pm \delta M_{peak}$ [AB mag]
SN2018bge	r	0.02389 ± 0.00011	35.159	17.823 ± 0.005	19.039 ± 0.837	0.044	-17.380 ± 0.162
	g			17.900 ± 0.003	10.310 ± 0.859	0.062	-17.321 ± 0.162
SN2018bqs	r	0.04730 ± 0.00060 †	36.675	18.776 ± 0.015	8.099 ± 0.295	0.017	-17.916 ± 0.499
	g			18.759 ± 0.021	6.146 ± 0.308	0.024	-17.941 ± 0.499
SN2018clq	g	0.04509 ± 0.00001	36.572	18.081 ± 0.004	5.158 ± 0.961	0.250	-18.742 ± 0.009
	r			18.118 ± 0.022	3.541 ± 1.005	0.176	-18.631 ± 0.023
SN2018cxn	r	0.04070 ± 0.00012	36.343	18.860 ± 0.006	15.844 ± 0.795	0.040	-17.523 ± 0.107
	g			18.864 ± 0.012	10.414 ± 0.686	0.057	-17.536 ± 0.108
SN2018cug	g	0.05000 ± 0.00373 ‡	36.804	18.580 ± 0.006	8.560 ± 0.467	0.129	-18.352 ± 2.746
	r			18.592 ± 0.009	12.140 ± 0.488	0.091	-18.303 ± 2.746
SN2018cyg	r	0.01127 ± 0.00001	33.507	18.176 ± 0.009	16.008 ± 0.814	0.041	-15.372 ± 0.031
	g			19.171 ± 0.007	11.240 ± 0.355	0.059	-14.394 ± 0.031
SN2018dfc	r	0.03653 ± 0.00009	36.102	17.603 ± 0.005	10.663 ± 0.168	0.193	-18.692 ± 0.089
	g			17.555 ± 0.007	7.553 ± 0.368	0.274	-18.821 ± 0.089
SN2018dfi	g	0.03130 ± 0.00016	35.758	17.987 ± 0.002	1.803 ± 0.432	0.070	-17.841 ± 0.183
	r			18.161 ± 0.037	3.482 ± 2.804	0.049	-17.646 ± 0.187
SN2018egh	r	0.03773 ± 0.00010	36.174	19.384 ± 0.001	17.785 ± 0.911	0.070	-16.860 ± 0.096
	g			19.548 ± 0.007	14.102 ± 1.349	0.100	-16.725 ± 0.096
SN2018fzn	g	0.03740 ± 0.00036 †	36.154	18.846 ± 0.019	19.266 ± 0.984	0.269	-17.577 ± 0.350
	r			18.505 ± 0.009	22.641 ± 0.549	0.189	-17.838 ± 0.350
SN2018fif	g	0.01719 ± 0.00003	34.434	17.471 ± 0.011	12.400 ± 0.644	0.351	-17.314 ± 0.061
	r			17.227 ± 0.006	17.048 ± 0.682	0.248	-17.454 ± 0.060
SN2018fsm	g	0.03500 ± 0.00366 ‡	36.006	17.939 ± 0.009	6.334 ± 0.631	0.325	-18.393 ± 3.765
	r			18.051 ± 0.011	9.460 ± 3.238	0.229	-18.184 ± 3.765
SN2018gts	g	0.029600 ± 0.00018	35.634	18.906 ± 0.009	6.248 ± 0.623	0.059	-16.787 ± 0.217
	r			18.315 ± 0.004	8.525 ± 0.579	0.041	-17.360 ± 0.217
SN2018grf	r	0.05380 ± 0.00307 ‡	36.969	18.463 ± 0.006	7.438 ± 0.279	0.081	-18.587 ± 2.110
	g			18.406 ± 0.005	5.624 ± 0.509	0.115	-18.678 ± 2.110
SN2018gvn	g	0.04330 ± 0.00333 ‡	36.481	18.359 ± 0.028	7.604 ± 1.791	0.137	-18.259 ± 2.806
SN2018iua	r	0.04150 ± 0.00284 ‡	36.386	18.943 ± 0.005	15.001 ± 1.016	0.083	-17.527 ± 2.490
	g			19.114 ± 0.015	4.043 ± 2.824	0.118	-17.391 ± 2.490
SN2018leh	g	0.02390 ± 0.00003	35.160	17.092 ± 0.005	13.735 ± 0.983	0.838	-18.905 ± 0.044
	r			17.156 ± 0.045	16.186 ± 1.282	0.590	-18.594 ± 0.063

NOTE— * : measured with respect to the EED; ‡ : the redshift was measured based on a spectrum from SEDm; † : the redshift was measured based on a higher resolution spectrum (e.g. DBSP and APO, here).

524 He II line would technically qualify this event as a mem-
525 ber of the flash class. The flash features of this object
526 seem to last longer than the rest of our flasher sample. A
527 discussion of the group of objects displaying long-lived
528 flash features, and their relation to some SNe IIn (e.g.,
529 SN 1998S; [Lentz et al. 2001](#), and SN 2018zd; [Zhang et al.](#)
530 [2020](#)), is outside the scope of this paper.

531 3.3.2. The Non-flashers

532 We defined an event as lacking flash features when we
533 had early, high-quality spectra (i.e. high S/N or higher

534 resolution than SEDm) that did not show any excess
535 emission around He II 4686 Å. Often, this meant that
536 the spectrum was blue and featureless. Among the ten
537 events included in our 2-day sub-sample, SN 2018fzn
538 was observed shortly after explosion (0.19 d, Table 2)
539 with SEDm. While the resolution was low, the signal to
540 noise was sufficient to determine that we could not find
541 any hint of possible excess emission (Fig. 7). Based on
542 the few previous events with spectra that were obtained
543 less than two days from EED (in particular SN 2013fs;

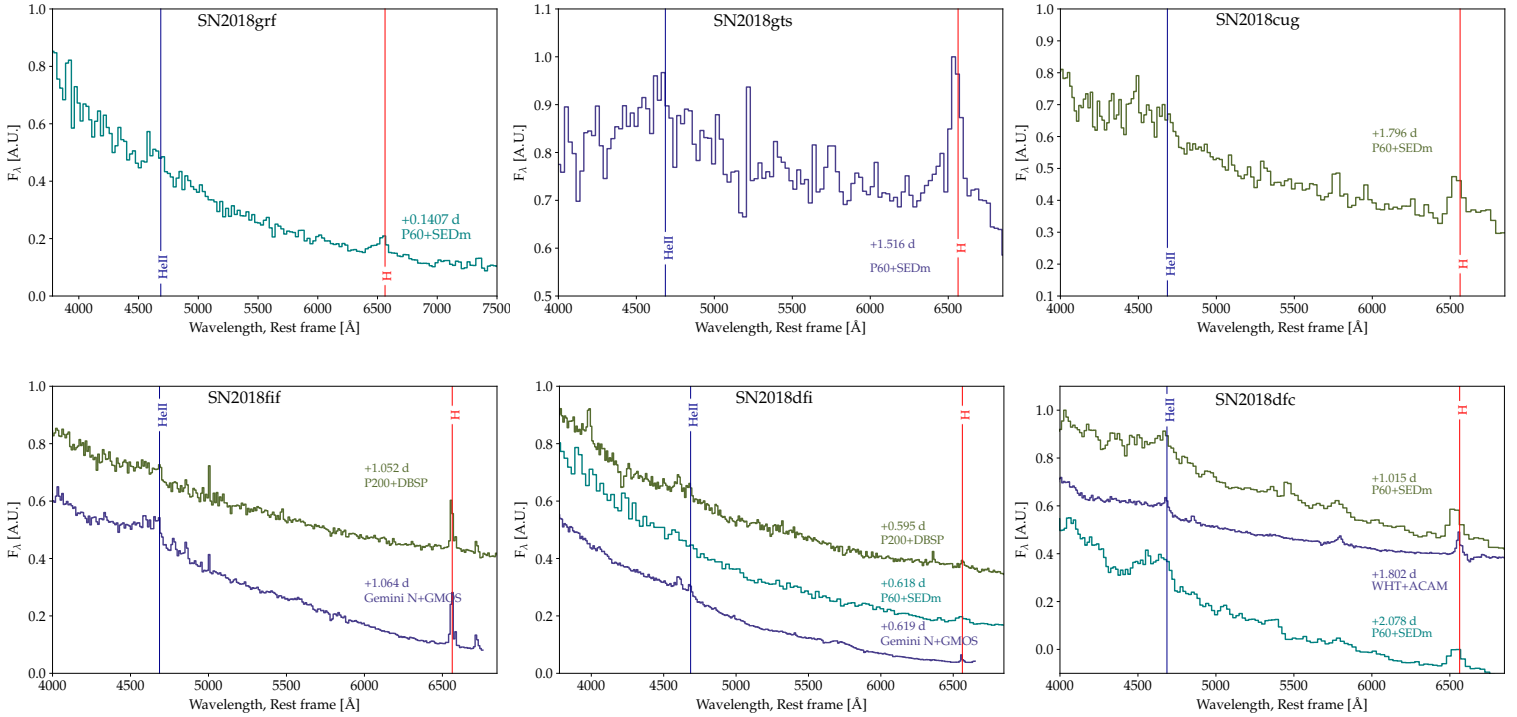


Figure 5. A collection of spectra of six confirmed Flashers. The acquisition time of the spectra are given with regard to the estimated explosion date.

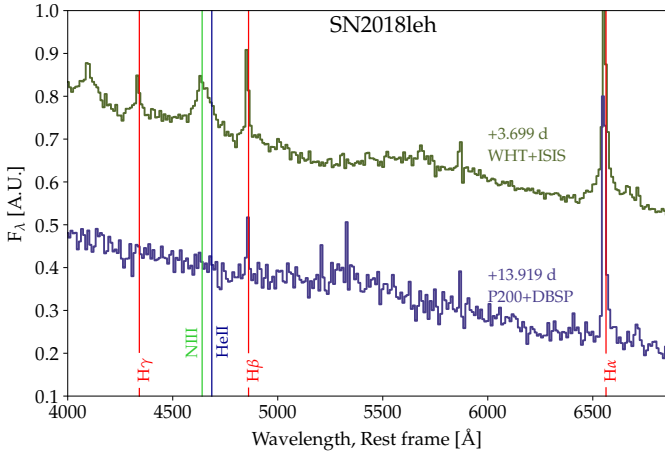


Figure 6. Spectroscopic evolution of SN 2018leh, a Type II_n SN that shows transient He II emission 4 days after its estimated explosion time.

544 Yaron et al. 2017), we expected strong emission lines
 545 that would be observable with SEDm (see the simulation
 546 in Extended Data Figure 2 of Gal-Yam et al. 2014).
 547 The first spectrum of SN 2018cxn was obtained with
 548 P200/DBSP less than a day past explosion. The higher
 549 resolution and the complete absence of He II emission
 550 (Fig. 7) imply no flash feature. For both cases, we con-

clude that there were no indications for a circumstellar
 shell.

3.3.3. The dubious flashers

551 SN 2018cyg and SN 2018egh both show excess flux
 552 around 4686 Å (Fig. 8). However, this excess does not
 553 resemble the ledge-shaped feature seen, for example, in
 554 the spectra of SN 2018fif (Soumagnac et al. 2019), and
 555 discussed above. An additional complication is that the
 556 spectra of SNe II at the early phase (prior to the appear-
 557 ance of strong and broad hydrogen Balmer lines)
 558 sometime show an absorption complex extending be-
 559 tween $\approx 4000 - 4500$ Å. Such a complex appears in the
 560 spectra of both SN 2018cyg and SN 2018egh. It was dif-
 561 ficult to determine whether the apparent bump around
 562 4600 Å represents an actual excess, or if it rather was
 563 the continuum edge redward of an absorption feature. In
 564 addition, even though we secured early, high resolution
 565 spectra for these objects (Table 2), they both lacked a
 566 narrow emission component from He II. The broad fea-
 567 tures were, however, transient and did not appear at
 568 later times. These issues made it difficult to determine
 569 if these events displayed flash features.

570 As an additional test of whether these two objects
 571 show a flux excess around 4600 Å, we conducted the fol-
 572 lowing test: we constructed model spectra composed of
 573 black body continua, over which we superposed model
 574 Gaussian emission lines whose width was a free param-
 575 eter.

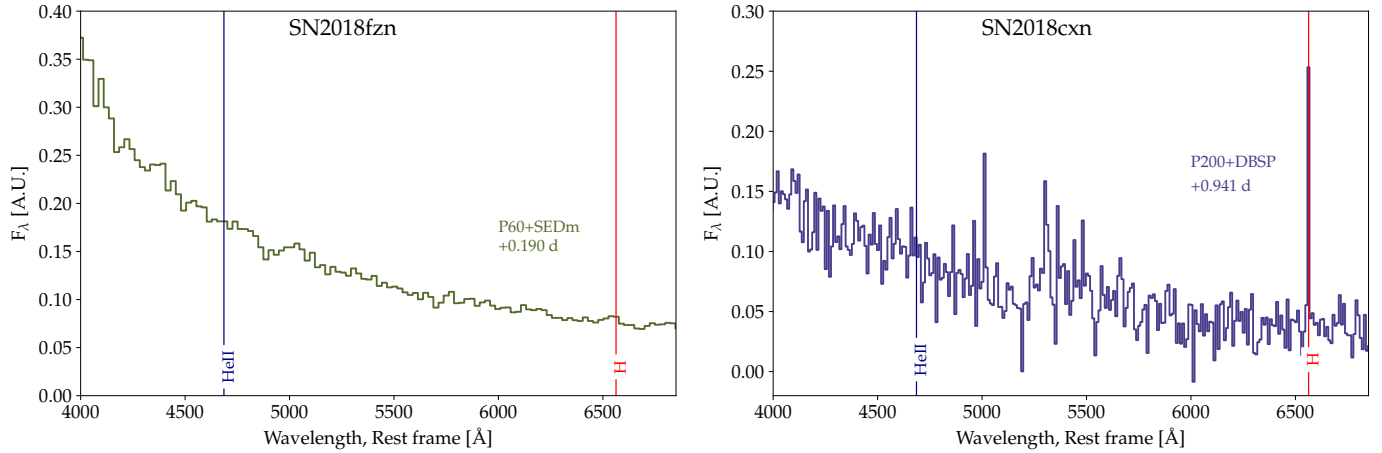


Figure 7. Early spectra of non-flashers SN 2018fzn and SN 2018cxn. These spectra were both obtained within less than a day from the estimated time of explosion. Only a smooth continuum is observed.

eter (with typical best fits of $\approx 100 \text{ km s}^{-1}$), in those cases (in particular, SN 2018dfc) where such lines were apparent. In addition, we added a broad feature extending between $4200 - 4750 \text{ \AA}$, which we defined by fitting a third-order polynomial to the ledge-shaped feature appearing in the SN 2018fif WHT spectrum (Fig. 8). The data were fitted using the python package *iminuit* (Dembinski et al. 2020). We then performed a χ^2 test to determine whether the bump feature is significantly detected (in the sense that $\Delta\chi^2 > 1$ between models) by comparing the goodness of fit over the intervals given in Table 5.

The results of these model comparisons are reported in Table 5 and Figure 9. As can be seen, the bump was strongly detected in the spectra of SN 2018dfc (and was also recovered for SN 2018fif), but neither for SN 2018cyg nor SN 2018egh. The results did not change if we fitted narrow lines, although no obvious additional lines (e.g. H_γ) were identified in the spectra. For SN 2018dfc, the bump feature was detected both in the earlier low-resolution SEDm spectrum (at low significance) and clearly in the later high-resolution WHT spectrum. In conclusion, we can not ascertain that SN 2018cyg and SN 2018egh showed flash features. We report below our results on flash statistics, considering all possible options (i.e., both, one, or neither of these show evidence for CSM).

4. DISCUSSION AND CONCLUSIONS

4.1. How common are flash features

Based on our systematic survey of infant SNe II with spectra obtained within two days of discovery, we found that at least 60%, and perhaps as many as 80% of the

sample of ten events showed evidence of flash-ionized emission. Taking into account our limited sample size and assuming binomial statistics $\mathcal{B}(k, n, p)$, we infer the true fraction of SNe with CSM which manifests as flash features, using a Bayesian model. The true probability p to observe an event with flash features given the observed fraction D is :

$$P(p|D) = \frac{P(D|p) \times \pi(p)}{P(D)} \quad (5)$$

Where p is the probability of observing a flash ionised event (here $p \in [0; 1]$), D is the observation presented in this paper (i.e.: 6 out of 10 candidates are showing flash features). The probability of our observation, $P(D)$ can be calculated with the formula of total probability, i.e. $P(D) = \int_0^1 \mathcal{B}(6, 10, p) \times \pi(p) dp$. We assumed a uniform distribution for the prior $\pi(p)$, which allowed us to write the posterior function as:

$$P(p|D) = \frac{\binom{10}{6} p^6 (1-p)^4}{\int_0^1 \binom{10}{6} p^6 (1-p)^4 dp} \quad (6)$$

which results in a Beta distribution (see Figure 10). We can put a strict lower limit on the fraction of infant SNe II showing flash features of $> 30.8\%$ ($> 23.5\%$) at the 95% (99%) confidence level (CL). The lower limit rises to 39.1% if either 18cyg or 18egh was a flasher, and to 48.3% if both were, at the 95% CL. This fraction rapidly drops when events with spectra obtained within 7 d from explosion are considered (the lower bound drops to 21.5% at the 95% confidence level); presumably the fraction could be even higher for events with even earlier spectra.

These results are broadly consistent with previous work by Khazov et al. (2016), which estimated that

Table 5. Results of test fits for models with and without the broad bump feature.

Name	Spectrum	Lines fit	χ_2/dof with bump	χ_2/dof without bump	Fit Interval []
SN 2018dfc	P60+SEDm +1.015 d	[HeII, H β]	0.76	1.43	4000-5300
SN 2018dfc	WHT+ACAM +1.082 d	[HeII, H β]	1.66	4.09	4000-5300
SN 2018fif	Gemini+GMOS +1.064 d	[HeII, H β]	2.12	3.34	4000-5000
SN 2018egh	WHT+ISIS +1.824 d	[HeII, H β]	0.87	0.91	4000-5300
SN 2018egh	WHT+ISIS +1.824 d	No Lines	0.87	0.93	4000-5300
SN 2018cyg	WHT+ACAM +1.673 d	No Lines	0.90	0.90	4000-5300

639 7–36% of SNe II show flash features in spectra obtained
640 within < 2 d from explosion (68% confidence level). It is
641 also consistent with the low observed frequency of flash
642 features among the general population of Type II SNe
643 reported in the literature, as these events very rarely
644 have a spectrum obtained < 2 d after explosion. Ta-
645 ble 2 shows that the fraction of flash events falls rapidly
646 at ages > 2 d. The unique nightly cadence of the ZTF
647 partnership survey enabled us to discover infant SNe
648 routinely, rapidly obtain spectra, and robustly measure
649 the frequency of this phenomenon.

4.2. Possible biases

651 Khazov et al. (2016) (see their Fig. 8) show that Type
652 II SNe showing flash-ionized features tend to be brighter
653 at peak than other events. We cannot confirm that this
654 is also true for our sample. We considered here the sub-
655 sample of infant supernovae whose first spectrum was
656 obtained within less than 7 days from the estimated
657 explosion time. The peak magnitudes were obtained
658 following the method described in 3.3.2. Figure 11,
659 top panel, shows the peak magnitudes in both g and
660 r bands for flashers and non-flashers. Flashers appear
661 to be brighter in both bands. However, when one con-
662 siders SN 2018cyg as a flasher, the average peak mag-
663 nitude of both groups is inverted, and non-flashers ap-
664 pear brighter than flashers (see Table 6, top section).
665 Since SN 2018cyg is strongly reddened, we repeated this
666 same analysis but with SN 2018cyg being host extinc-
667 tion corrected. To apply the extinction correction, we
668 used the spectrum from 2018 August, 4¹² and applied
669 the method described in Poznanski et al. 2012, using the
670 line doublet of sodium. We considered the doublet not
671 to be resolved and apply the following formula:

$$\log_{10}(E_{B-V}) = 1.17 \times EW(D_1 + D_2) - 1.85 \pm 0.08 \quad (7)$$

672 We estimated the EW of the D₁+D₂ lines using the
673 built-in tool from WISeREP by measuring it several
674 times. The mean EW is 1.64 with an error of 0.17 .

¹² see on WISeREP : <https://wiserep.weizmann.ac.il/object/698>

675 Following Eq. (5), the final peak magnitudes for SN
676 2018cyg are : $M_{peak,r} = -18.45 \pm 0.50$ and $M_{peak,g} =$
677 -18.77 ± 0.80 . Table 6 summarises the different cases:
678 whether SN 2018cyg is a flasher and whether SN 2018cyg
679 was corrected for estimated host extinction. We find
680 that flash events are not inherently brighter than non-
681 flash events.

682 We also inspected in Fig. 11 (lower panel) the distri-
683 bution of apparent magnitudes at discovery for our < 7 d
684 sample. As can be seen there, we found that the flash
685 events were not significantly brighter at discovery than
686 other events. Thus neither were more likely to be dis-
687 covered, nor to be followed up, as both aspects depend
688 on the apparent magnitude of the object at discovery.

4.3. Implications

690 We showed here that a significant fraction, and possi-
691 bly most, Type II SN progenitors, show transient emis-
692 sion lines in their early spectra, which provides evidence
693 that these stars are embedded in a compact distribution
694 of CSM (Yaron et al. 2017). The narrow width of these
695 emission lines indicates a slow expansion speed for the
696 CSM ($100 - 800 \text{ km s}^{-1}$, Boian & Groh 2020), and com-
697 bined with its compact radial dimension (< 10^{15} cm)
698 we have evidence that the CSM was deposited by the
699 stars within months to a few years prior to its termi-
700 nal explosion. Assuming these progenitors are mostly
701 red supergiants (RSGs; Smartt 2015), this would sug-
702 gest that most exploding RSGs experience an enhanced
703 mass loss shortly prior to explosion.

704 While RSGs certainly lose mass during their final
705 stages of evolution (Smith 2014), such a period of en-
706 hanced mass loss shortly (months to a year) prior to
707 explosion is not explained by standard stellar evolution
708 models. Thus, our work indicates that additional phys-
709 ical processes leading to such pre-explosion instabilities
710 (e.g., Arnett & Meakin 2011, Shiode & Quataert 2014)

711 not only exist, but are ubiquitous among massive stars.
712 As we have shown that most SN II progenitors likely
713 undergo a remarkable evolution shortly prior to explo-
714 sion, it may be needed to re-examine the stellar mod-
715 els used as initial conditions to explosion simulations.

Table 6. Peak magnitude comparison between the flash events and the non flash events.

	$M_{peak, flasher}$	$M_{peak, non flasher}$	
<i>18cyg not corrected for extinction</i>			
r band	18cyg \subset flasher	-17.58 ± 0.96	-17.76 ± 0.42
	18cyg $\not\subset$ flasher	-17.91 ± 0.48	-17.46 ± 0.90
<i>18cyg corrected for extinction</i>			
	18cyg \subset flasher	-17.97 ± 0.48	-17.76 ± 0.42
	18cyg $\not\subset$ flasher	-17.91 ± 0.48	-17.85 ± 0.46
	$M_{peak, flasher}$	$M_{peak, non flasher}$	
<i>18cyg not corrected for extinction</i>			
g band	18cyg \subset flasher	-17.30 ± 1.31	-17.64 ± 0.57
	18cyg $\not\subset$ flasher	-17.73 ± 0.71	-17.31 ± 1.13
<i>18cyg corrected for extinction</i>			
	18cyg \subset flasher	-17.86 ± 0.75	-17.64 ± 0.57
	18cyg $\not\subset$ flasher	-17.76 ± 0.75	-17.75 ± 0.64

NOTE— This analysis is performed with the subsample which has a first spectrum within less than seven days from the estimated explosion time.

At least some of the effects proposed to explain such pre-explosion mass loss may render the spherical pre-explosion stellar models used in explosion simulations less realistic (Arnett & Meakin 2016). Perhaps our work provides a clue about how to tackle some of the problems encountered in reproducing the observed properties of SN explosions using numerical explosion models.

5. CONCLUSIONS

We report the results from the first year (2018) of our systematic survey for infant Type II SNe in the ZTF partnership survey. We collected 28 such objects (at a rate of about one per week) and obtained rapid follow-up spectroscopy within 2 d from explosion for 10 events. Between 6 – 8 of these show evidence for transient emission from a surrounding distribution of CSM. Thus we can place a strict lower limit of $> 30\%$ (at 95% C.L.) on the fraction of SN II progenitors that explode within compact CSM distributions. This finding is inconsistent with predictions from standard stellar evolution models. It suggests that additional physics is required to explain the final stages (~ 1 year prior to explosion) of massive star evolution. The structural changes that may accompany such final episodes of intense mass loss can modify the stellar structure prior to explosion and may require adjusting the initial conditions assumed for core-collapse SN explosion simulations, and may thus shed light on the yet unsolved question of how massive stars end their lives in supernova explosions.

6. ACKNOWLEDGEMENTS

AGY’s research is supported by the EU via ERC grant No. 725161, the ISF GW excellence center, an IMOS

space infrastructure grant and BSF/Transformative, Minerva and GIF grants, as well as The Benoziyo Endowment Fund for the Advancement of Science, the Deloro Institute for Advanced Research in Space and Optics, The Kimmel Center for planetary science, The Veronika A. Rabl Physics Discretionary Fund, Paul and Tina Gardner, Yeda-Sela and the WIS-CIT joint research grant; AGY is the recipient of the Helen and Martin Kimmel Award for Innovative Investigation. The ztfquery code was funded by the European Research Council (ERC) under the European Union’s Horizon 2020 research and innovation programme (grant agreement No. 759194 - USNAC, PI: Rigault). The ZTF forced-photometry service was funded under the Heising-Simons Foundation grant #12540303 (PI: Graham). Based on observations obtained with the Samuel Oschin 48-inch Telescope at the Palomar Observatory as part of the Zwicky Transient Facility project. ZTF is supported by the National Science Foundation under Grant No. AST-1440341 and a collaboration including Caltech, IPAC, the Weizmann Institute for Science, the Oskar Klein Center at Stockholm University, the University of Maryland, the University of Washington, Deutsches Elektronen-Synchrotron and Humboldt University, Los Alamos National Laboratories, the TANGO Consortium of Taiwan, the University of Wisconsin at Milwaukee, and Lawrence Berkeley National Laboratories. Operations are conducted by COO, IPAC, and UW. The data presented here were obtained [in part] with ALFOSC, which is provided by the Instituto de Astrofísica de Andalucía (IAA) under a joint agreement with the University of Copenhagen

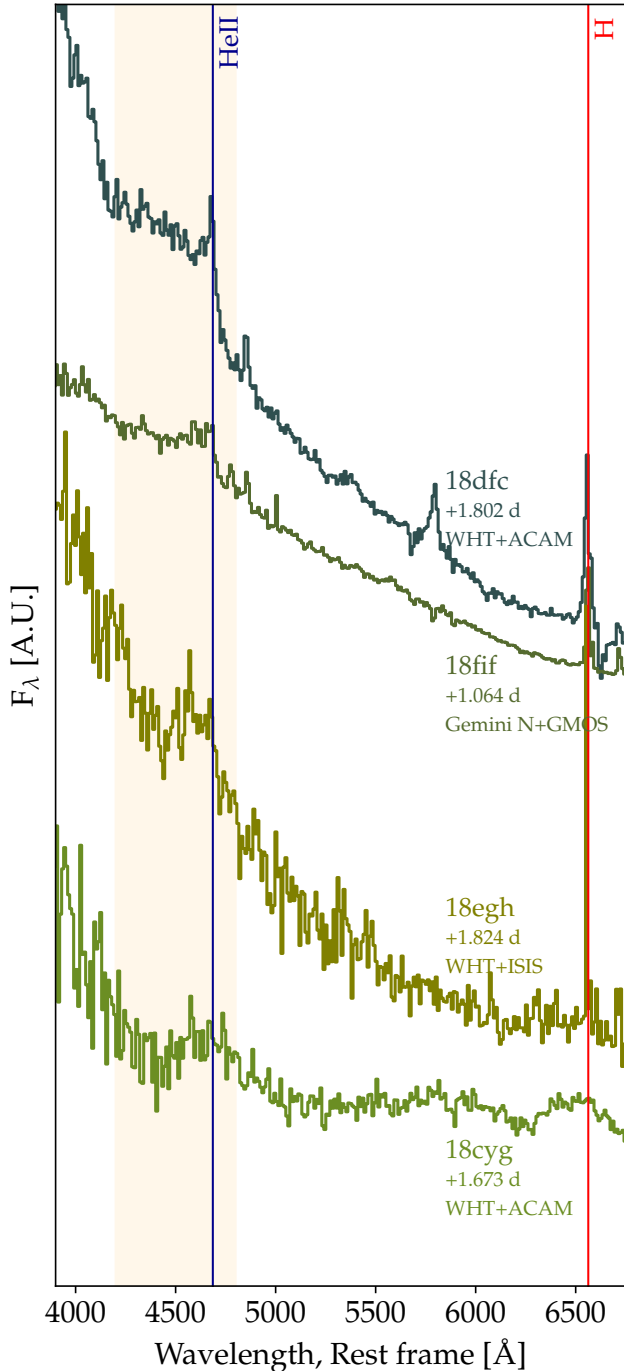


Figure 8. Candidates showing a wide bump-like structure close to the He II emission line. We highlight in orange the region we searched for excess emission. The spectra of 18cyg and 18egh were both binned to 10

779 and NOTSA. A.A.M. is funded by the LSST Corpora-
 780 tion, the Brinson Foundation, and the Moore Founda-
 781 tion in support of the LSSTC Data Science Fellowship
 782 Program; he also receives support as a CIERA Fellow
 783 by the CIERA Postdoctoral Fellowship Program (Cen-
 784 ter for Interdisciplinary Exploration and Research in
 785 Astrophysics, Northwestern University). Based on ob-
 786 servations obtained at the international Gemini Obser-
 787 vatory, a program of NSF's NOIRLab, which is man-
 788 aged by the Association of Universities for Research
 789 in Astronomy (AURA) under a cooperative agreement
 790 with the National Science Foundation. on behalf of the
 791 Gemini Observatory partnership: the National Science
 792 Foundation (United States), National Research Council
 793 (Canada), Agencia Nacional de Investigación y De-
 794 sarrollo (Chile), Ministerio de Ciencia, Tecnología e In-
 795 novación (Argentina), Ministério da Ciência, Tecnolo-
 796 gia, Inovações e Comunicações (Brazil), and Korea As-
 797 tronomy and Space Science Institute (Republic of Ko-
 798 rea). This research has made use of the NASA/IPAC Ex-
 799 tragalactic Database (NED), which is funded by the Na-
 800 tional Aeronautics and Space Administration and oper-
 801 ated by the California Institute of Technology. The re-
 802 search leading to these results has received funding from
 803 the European Union Seventh Framework Programme
 804 (FP7/2013-2016) under grant agreement No. 312430
 805 (OPTICON).

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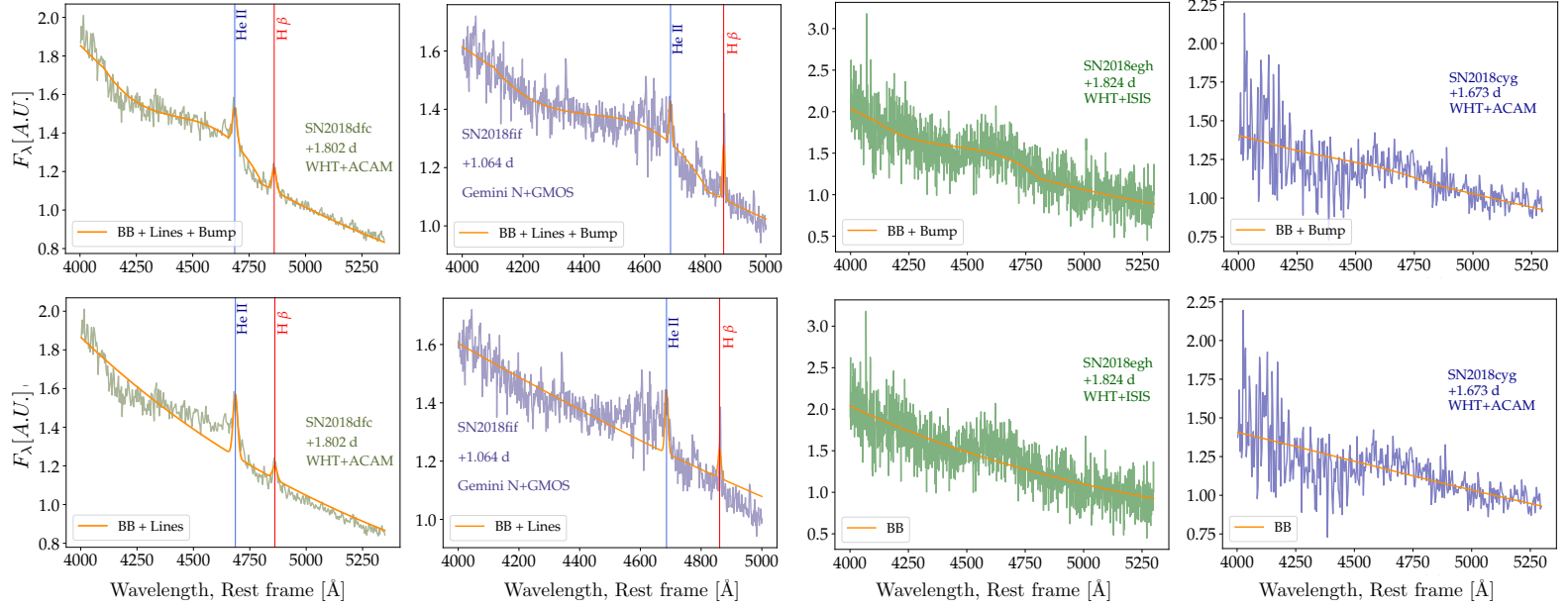


Figure 9. Fit results with (top panels) and without (bottom panels) the broad feature component for SNe 2018dfc, 2018fff, 2018egh and 2018cyg (from left to right). No narrow emission lines are seen in the spectra of 2018egh and 2018cyg, and neither provides a significant detection of a bump component.

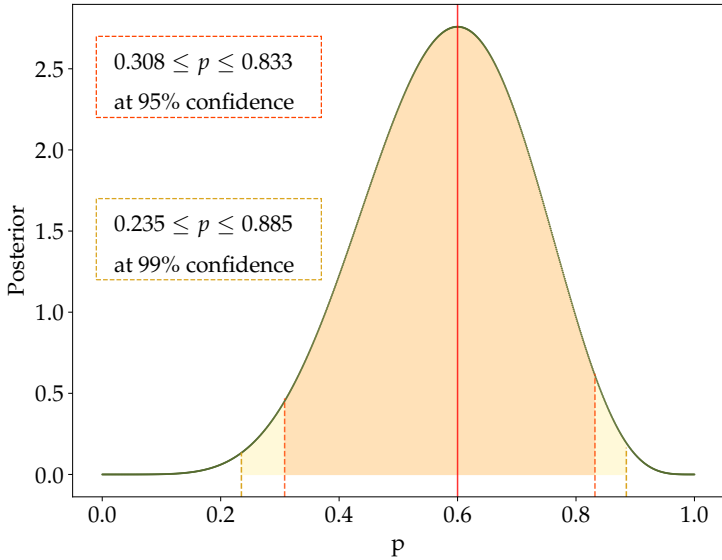


Figure 10. Posterior probability distribution vs. the probability to observe a flash ionised event. This analysis is based on the subsample of infant candidates which had a first spectrum within < 2 days from the EED, and considering that 18cyg and 18egh are not flashers. The lower limit is 30.8% (23.5%) at 95%(99%) confidence interval.

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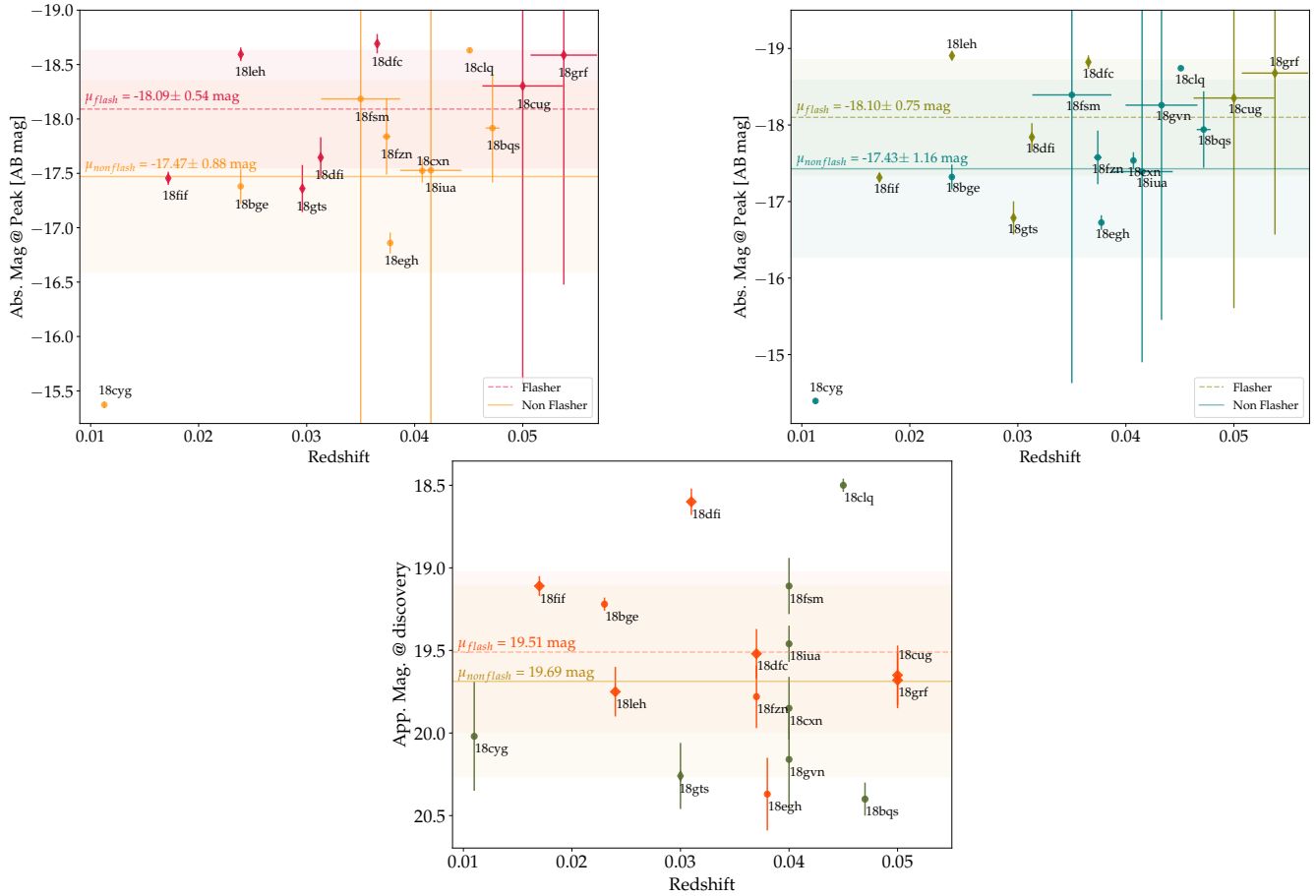


Figure 11. Top: absolute magnitude in r band (left) and g band (right) vs. redshift. Bottom: apparent magnitude at discovery vs. redshift. Color bands represent the error on the mean peak magnitude for both flash and non flash groups. SN 18cyg is host reddened and hence appears very faint, see text.

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990 84

7. APPENDIX

7.1. *Justification of candidate rejection*

991 The full list of candidate infant SNe II returned by `ztfquery` (see § 2.2) is given in Table 7. Of the 43 candidates,
 992 inspection shows that 15 are spurious, and these have been removed from our sample. We provide some comments on
 993 removed objects.
 994
 995

996 *Early false positives*—A group of objects detected right at the start of the survey (during March 2018 till early April)
 997 suffered from unreliable photometry, manifest as a mix of detections and non-detections during the same period,
 998 and often during the same night. This is likely due to problematic early references. The mix of detections and
 999 non-detections created artificial triggers due to a spurious non-detection just prior to the first detection. This group
 1000 includes ZTF18aaayemw, ZTF18aaccmnh, ZTF18aagrdd (which was also detected by ATLAS 3 days prior to the
 1001 ZTF false non-detection, and reported to the TNS as AT2018ahi), ZTF18aahrzrb, ZTF18aainvic, and ZTF18aaogibq.

1002 *ZTF18aaqkdwu*—This trigger resulted from a spurious photometry point generated by the pipeline at the location of
 1003 SN 2019eoe a year prior to the explosion of the actual SN.

1004 *ZTF18aasxvsg*—Additional analysis recovered several clear detections prior to the spurious non-detection that triggered
 1005 this event.

1006 *ZTF18abcqhgr*—This event is likely a real infant SN II, but we could not recover it using the forced photometry pipeline
 1007 and it was therefore removed from the sample. This object does not have an early spectrum.

1008 *ZTF18acbvwsp*—This event was detected by SNHunt and reported to the TNS as AT 2018hqm a few days prior to the
 1009 only ZTF non-detection, indicating it is likely not a RI SN.

1010 *ZTF18acecuwq*—The early photometry of this event shows a mix of detections and non-detections during the same
 1011 nights, and was deemed unreliable. A spectrum obtained within a day of the false non-detection (A. Tzanidakis, in
 1012 preparation) is that of an old SN II, supporting this conclusion.

1013 *ZTF18acgvqiq*—This event was detected by ATLAS and reported to the TNS as SN 2018fru more than 2 months prior
 1014 to the ZTF non-detection, indicating our non-detections preceding the ZTF first detection were spurious.

1015 *ZTF18acefuhk*—Updated photometry does not recover a non-detection prior to first detection that satisfies our criteria.
 1016 This object does not have early spectra.

1017 *ZTF18acqxyiq*—The forced photometry pipeline did not recover the non-detection by the real-time pipeline, leaving
 1018 the explosion time poorly constrained.

1019 *ZTF18adbikdz*—This object was detected by Gaia and reported to the TNS as AT2017isr over a month prior to the
 1020 first detection by ZTF (when it was already declining). Our single non-detection is spurious.

Table 7. Results of the search for infant SN II using ZTFquery

Name	RA	Dec	Redshift	First Detection	First spectrum	Real?
	[deg]	[deg]		[days]	[days]	
ZTF18aaayemw	134.8982936	45.6116267	0.052 (NED; Helou et al. 1991)	2458156.7621	0.024	✗
ZTF18aaccmnh	194.9769678	37.8589965	0.035 (NED; Helou et al. 1991)	2458184.8604	0.018	✗
ZTF18aagrdded	209.8414748	46.0317554	0.047 ¹	2458198.8809	0.011	✗
ZTF18aahrzrb	181.397224	34.3888035	0.040 ¹	2458217.7371	1.001	✗
ZTF18aainvic	256.5204624	29.6683607	0.032 ¹	2458218.9088	0.019	✗
ZTF18aaogibq	253.5409858	24.721127	0.037 (NED; Helou et al. 1991)	2458231.8783	0.020	✗
ZTF18aaqkdwu	199.7588529	45.0263019	0.060 (NED; Helou et al. 1991)	2458243.677	0.001	✗
ZTF18aaqkoyr	166.0666639	50.0306275	0.023 (NED; Helou et al. 1991)	2458243.6854	1.036	✓
ZTF18aarpttw	247.2599041	43.6268239	0.047 ¹	2458246.822	1.001	✓
ZTF18aarqxbw	276.4265403	34.6584885	0.048 ¹	2458246.8404	1.878	✓
ZTF18aasxvsg	217.1290246	37.0678367	0.025 (NED; Helou et al. 1991)	2458244.8361	0.018	✗
ZTF18aatlfus	257.1764284	28.5206128	0.045 (NED; Helou et al. 1991)	2458249.8534	1.913	✓
ZTF18aavpady	273.0031098	44.3602114	0.047 ¹	2458257.8452	0.870	✓
ZTF18aawyjjq	263.0587448	36.0740074	0.040 ¹	2458263.796	0.011	✓
ZTF18aayxxew	197.1395703	45.9861525	0.061 ¹	2458278.7043	1.961	✓
ZTF18abcezmh	269.4519011	40.0764001	0.057 ¹	2458288.7881	0.874	✓
ZTF18abckutn	237.0269066	55.7148077	0.040 (NED; Helou et al. 1991)	2458290.6992	0.834	✓
ZTF18abcptmt	267.3298968	49.4124315	0.050 ¹	2458291.7869	0.878	✓
ZTF18abcqhgr	254.818188	60.4317998	0.070 (NED; Helou et al. 1991)	2458291.8048	0.021	✗
ZTF18abdbysy	233.5352962	56.6968517	0.011 (NED; Helou et al. 1991)	2458295.7208	0.016	✓
ZTF18abdjdjpt	278.7048393	38.2987246	0.070 ¹	2458295.7913	0.021	✓
ZTF18abeajml	252.0323502	24.3041089	0.037 (NED; Helou et al. 1991)	2458303.7989	1.002	✓
ZTF18abffyqp	252.7086818	45.397907	0.031 (NED; Helou et al. 1991)	2458307.6862	0.864	✓
ZTF18abgqvww	254.3164613	31.9632993	0.038 (NED; Helou et al. 1991)	2458313.7295	0.891	✓
ZTF18abgrbjb	274.9986631	51.7965471	0.030 ¹	2458313.7492	0.032	✓
ZTF18abimhfu	240.1422651	31.6429838	0.050 ¹	2458320.6667	0.912	✓
ZTF18abojpnr	297.4871203	59.5928266	0.037 ¹	2458351.7166	0.021	✓
ZTF18abokyfk	2.3606444	47.3540929	0.017 (NED; Helou et al. 1991)	2458351.8659	0.887	✓
ZTF18abrlljc	253.1840255	70.0882366	0.050 ¹	2458359.7	0.054	✓
ZTF18absldff	33.5997507	30.811929	0.035 ¹	2458363.8793	0.913	✓
ZTF18abufaej	4.4825733	12.0916007	0.062 ¹	2458368.8738	0.036	✓
ZTF18abvmdf	249.1975409	55.7358424	0.030 (NED; Helou et al. 1991)	2458375.7154	0.016	✓
ZTF18abwlsoi	261.8976711	71.5302584	0.054 ¹	2458377.6334	0.895	✓
ZTF18abyvenk	273.9764532	44.6964862	0.043 ¹	2458385.6212	0.858	✓
ZTF18acbwvsp	341.9067649	39.8806077	0.017 (NED; Helou et al. 1991)	2458423.6368	0.907	✗
ZTF18acecuxq	68.8323442	17.1948085	0.026 (NED; Helou et al. 1991)	2458431.8168	1.011	✗
ZTF18acefuhk	136.7936282	43.9207446	0.058 (NED; Helou et al. 1991)	2458426.9469	0.951	✗
ZTF18acvggiq	204.0157722	66.3012068	0.010 (NED; Helou et al. 1991)	2458432.0181	1.966	✗
ZTF18achtnvk	96.1687142	46.5039037	0.040 ¹	2458434.9036	0.043	✓
ZTF18acploez	130.03737	68.9031912	0.042 ¹	2458440.9658	1.957	✓
ZTF18acqxyiq	149.8258285	34.895493	0.038 (NED; Helou et al. 1991)	2458443.9437	0.001	✗
ZTF18adbikdz	252.014493	26.2118328	0.034 (NED; Helou et al. 1991)	2458482.0504	0.004	✗
ZTF18adbmrug	61.2637352	25.2619198	0.024 (NED; Helou et al. 1991)	2458482.6991	1.897	✓

NOTE—43 candidates were found, of which 15 ($\sim 35\%$) were spurious, leaving 28 infant SNe II in our sample¹ This work

7.2. *Forced photometry light curves*

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1022 Figure 12

1023 Figure 13

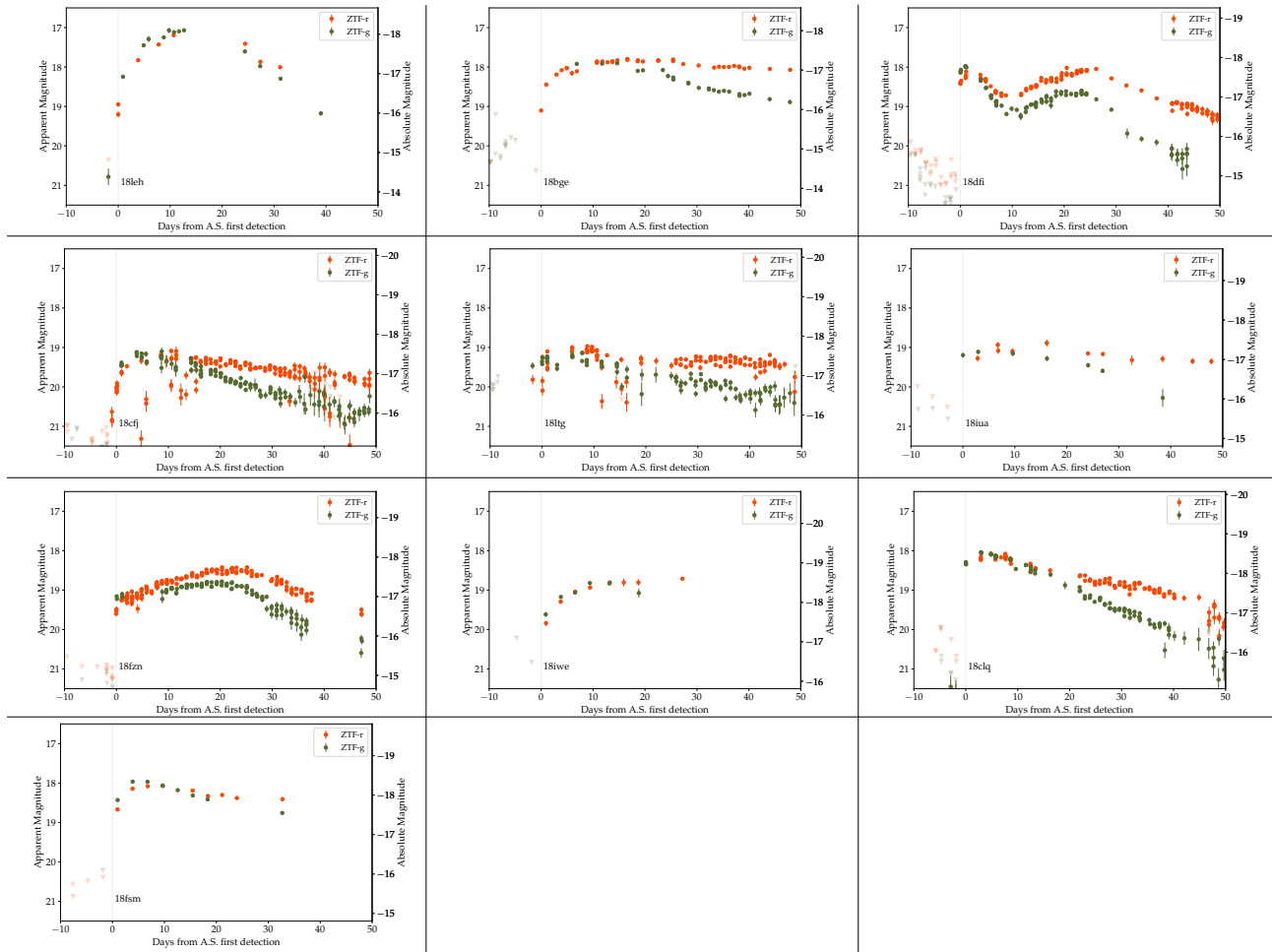


Figure 12. Forced photometry light curves of our Real Infant SN II sample. The grey line represents the first detection from the alert system (i.e. time "0"). Any detection prior to this line was recovered by the forced photometry pipeline. The left y-axis corresponds to the apparent magnitude; the right y-axis to the absolute magnitude. The explosion date of these objects was estimated as the middle date between the last non-detection and the first detection

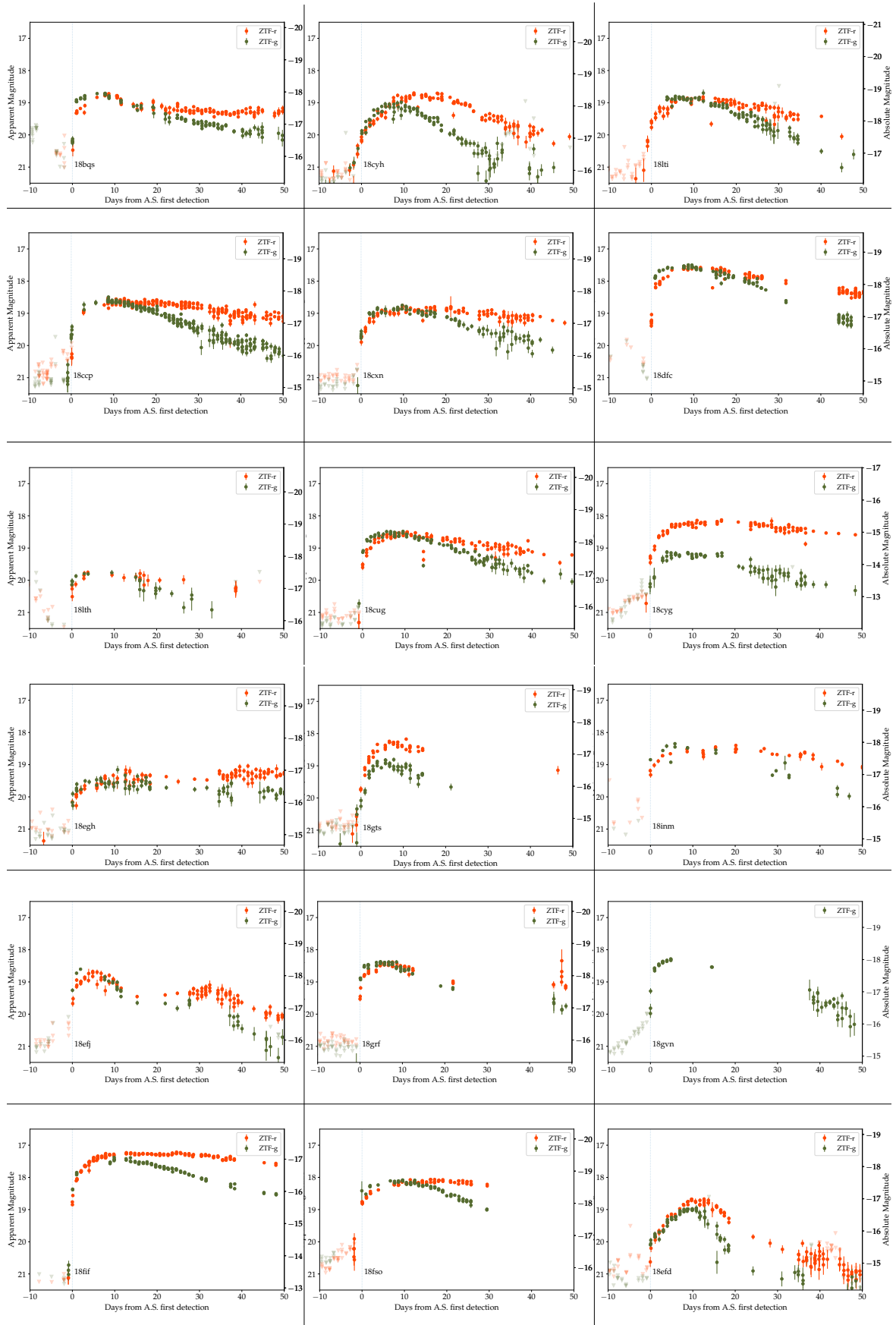


Figure 13. Figure 13 Forced-photometry light curves in both r and g band [continued]. The explosion date of these objects was estimated using the method described in 3.3.1.

1024

7.3. *Classification spectra of the Real Infant sample*

1025 Figure 14

1026 Figure 15

1027 Table 8

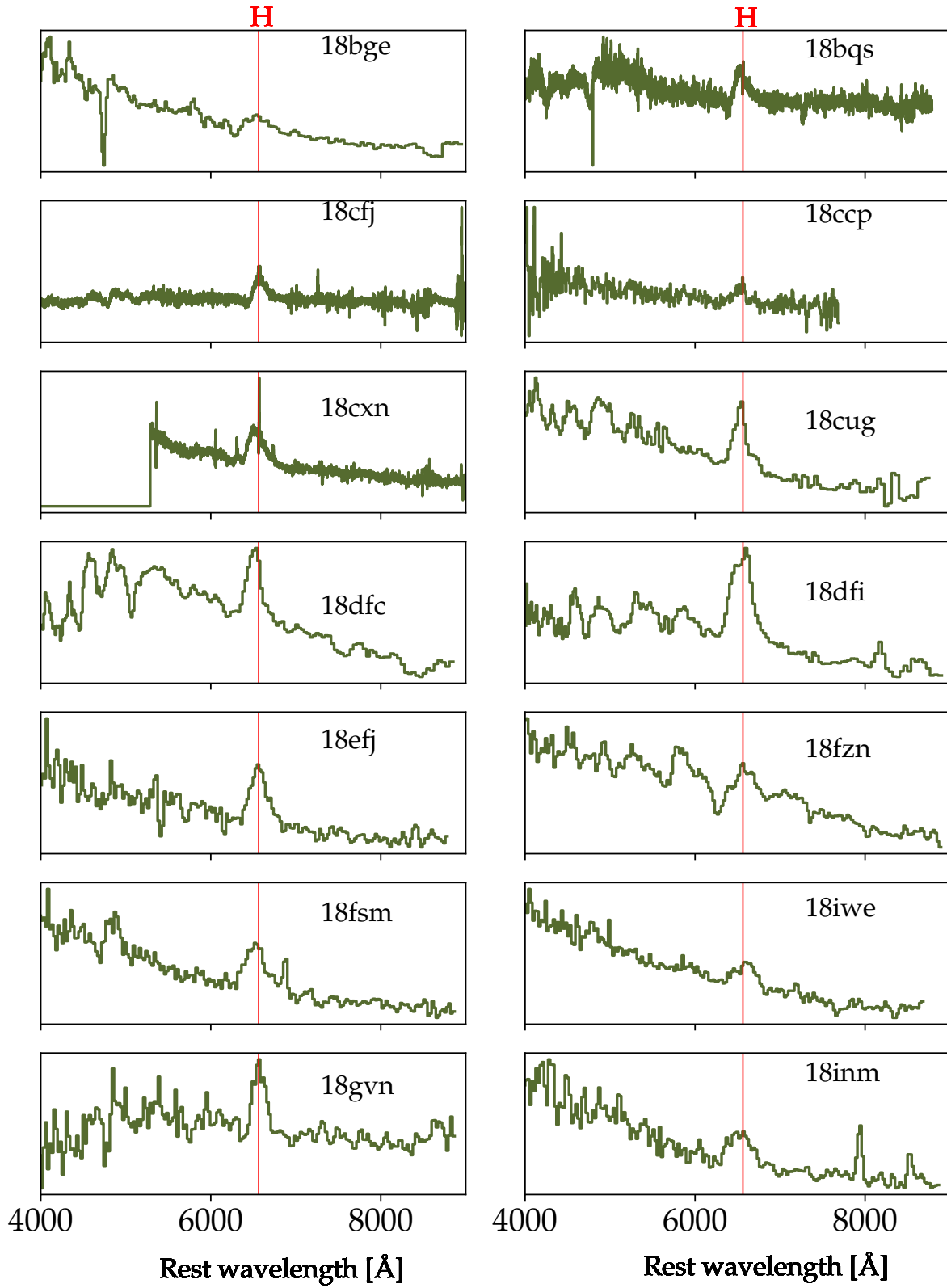


Figure 14. Classification spectra of the real infant sample. The red vertical line marks the H α line. See detailed in Table 8.

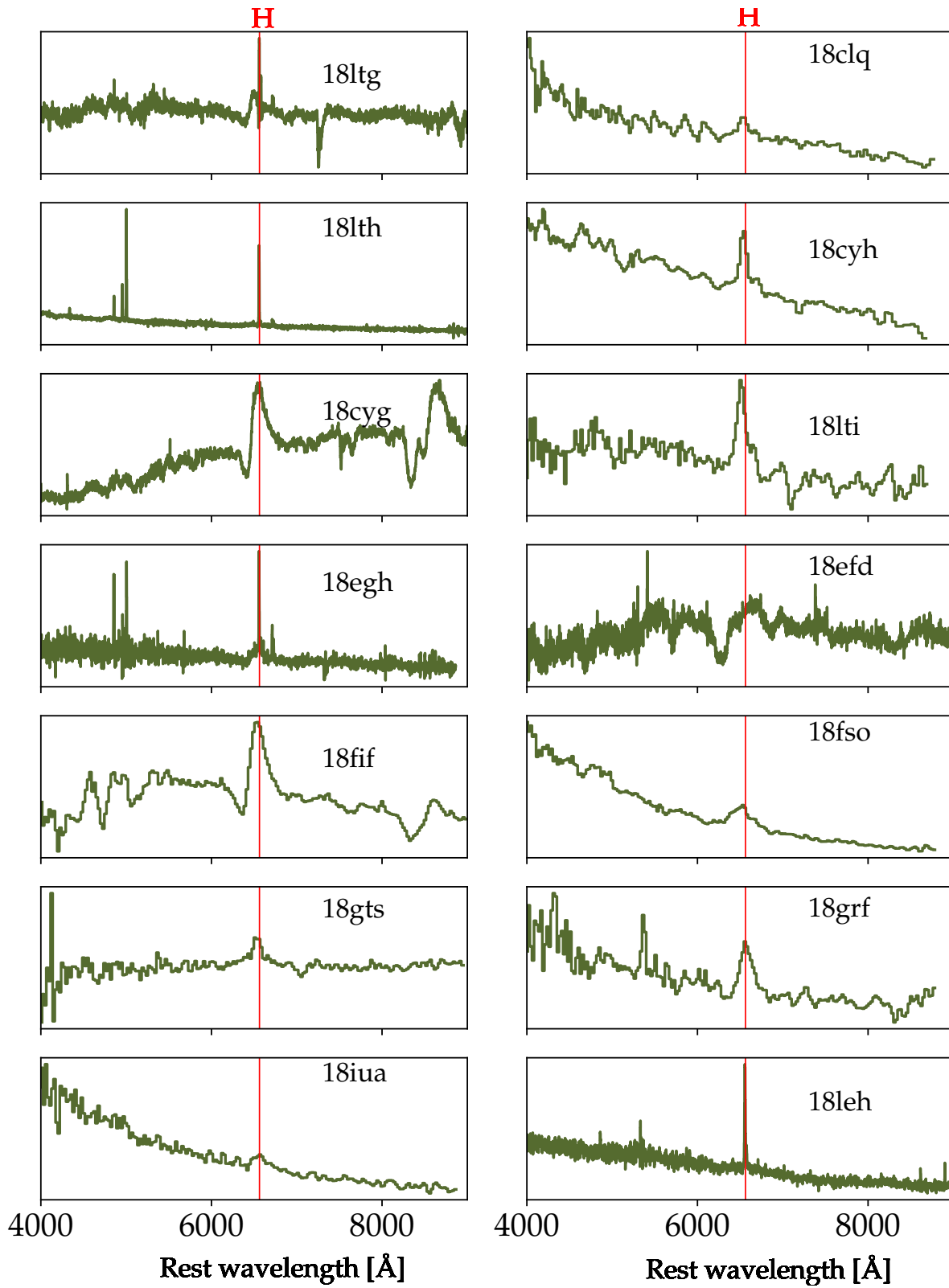


Figure 15. [continued] Classification spectra of the real infant sample.

Table 8. List of photospheric spectra corresponding to Figures 14 and 15

IAU name	Estimated explosion time [JD]	redshift	Instrument	Time to spectrum [d]
18bge	2458243.1671	0.024	SEDm+P60	3.33
18bqs	2458246.8133	0.047	DBSP+P200	38.69
18ltg	2458241.436	0.048	DBSP+P200	37.06
18clq	2458248.8967	0.045	SEDm+P60	7.60
18cfj	2458256.4531	0.047	LRIS+Keck1	55.05
18ccp	2458263.7743	0.040	SPRAT+LT	15.73
18lth	2458278.6531	0.061	LRIS+Keck1	7.85
18cyh	2458286.3752	0.057	SEDm+P60	16.12
18cxn	2458289.8074	0.041	DBSP+P200	17.69
18cug	2458290.916	0.050	SEDm+P60	24.58
18cyg	2458294.7273	0.011	DBSP+P200	39.77
18lti	2458294.6217	0.070	SEDm+P60	39.88
18dfc	2458303.7777	0.037	SEDm+P60	27.72
18dfi	2458307.254	0.031	SEDm+P60	27.25
18egh	2458312.7454	0.038	DBSP+P200	38.75
18efd	2458312.8922	0.030	DBSP+P200	21.61
18efj	2458320.6574	0.050	SEDm+P60	41.84
18fzn	2458351.7068	0.037	SEDm+P60	18.79
18fif	2458350.9535	0.017	SEDm+P60	35.55
18fso	2458357.6987	0.050	SEDm+P60	13.80
18fsm	2458363.4226	0.035	SEDm+P60	21.08
18iwe	2458368.8561	0.062	SEDm+P60	11.64
18gts	2458375.1028	0.030	SEDm+P60	41.40
18grf	2458377.6103	0.054	SEDm+P60	65.89
18gvn	2458385.6198	0.043	SEDm+P60	63.88
18inm	2458432.9113	0.040	SEDm+P60	15.59
18iua	2458439.9877	0.042	SEDm+P60	3.51
18leh	2458481.7505	0.024	DBSP+P200	13.75