

LETTER TO THE EDITOR

TESS's first planet: a super-Earth transiting the naked-eye star π Mensae

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ABSTRACT

We report on the confirmation and mass determination of π Men c, the first transiting planet discovered by NASA's *TESS* space mission. π Men is a naked-eye ($V=5.65$ mag), quiet G0 V star that was previously known to host a sub-stellar companion (π Men b) on a long-period ($P_{\text{orb}}=2091$ days), eccentric ($e=0.64$) orbit. Using *TESS* time-series photometry, combined with *Gaia* data, published UCLES@AAT Doppler measurements, and archival HARPS@ESO-3.6m radial velocities, we found that π Men c is a close-in planet with an orbital period of $P_{\text{orb}}=6.27$ days, a mass of $M_c=4.52 \pm 0.81 M_{\oplus}$, and a radius of $R_c=2.06 \pm 0.03 R_{\oplus}$. Based on the planet's orbital period and size, π Men c is a super-Earth located at, or close to, the radius gap, while its mass and bulk density suggest it may have held on to a significant atmosphere. Because of the brightness of the host star, this system is highly suitable for a wide range of further studies to characterize the planetary atmosphere and dynamical properties. We also performed an asteroseismic analysis of the *TESS* data and detected a hint of power excess consistent with the seismic values expected for this star, although this result depends on the photometric aperture used to extract the light curve. This marginal detection is expected from pre-launch simulations hinting at the asteroseismic potential of the *TESS* mission for longer, multi-sector observations and/or for more evolved bright stars.

Key words. Planetary systems – Planets and satellites: individual: π Mensae b, π Mensae c – Stars: fundamental parameters – Stars: individual: π Mensae – Techniques: photometric – Techniques: radial velocities

1. Introduction

Successfully launched on 18 April 2018, NASA's Transiting Exoplanet Survey Satellite (*TESS*; Ricker et al. 2015) will provide us with a leap forward in understanding the diversity of small planets ($R_p < 4 R_{\oplus}$). Unlike previous space missions, *TESS* is performing an all-sky transit survey focusing on bright stars ($5 < V < 11$ mag) so that detailed characterizations of the planets and their atmospheres can be performed. In its two-year prime mission, *TESS* observes first the southern and then the northern ecliptic hemisphere. The survey is broken up into 26 anti-solar sky sectors. *TESS* uses four cameras to observe each sector, resulting in a combined field of view of $24^{\circ} \times 96^{\circ}$. The overlap between sectors towards the ecliptic poles provides greater sensitivity to smaller and longer-period planets in those regions of the celestial sphere. *TESS* records full-frame images of its entire field of view every 30 minutes and observes approximately 200 000 pre-selected main-sequence stars with a cadence of ~ 2 minutes. The mission will certainly open a new era in the study of close-in small planets, providing us with cornerstone objects amenable to both mass determination – via Doppler spectroscopy – and atmospheric characterization – via transmission spectroscopy with NASA's James Webb space telescope (*JWST*) and the next generation of extremely large ground-based telescopes (*ELT*, *TMT*, and *GMT*).

Following a successful commissioning of 3 months, *TESS* started the science operation on 25 July 2018 by photometrically monitoring its first sector (Sector 1), which is centered at coordinates $\alpha = 352.68^{\circ}$, $\delta = -64.85^{\circ}$ (J2000). Shortly after ~ 30 days of (almost) continuous observations in Sector 1, 73 transiting planet candidates were detected in the two-minute cadence light curves by the *TESS* team and made available to the scientific community upon registration to a dedicated web portal hosted at the Massachusetts Institute of Technology (MIT) web page¹.

In this letter, we present the spectroscopic confirmation of π Men c, the first transiting planet discovered by the *TESS* space mission. The host star is π Mensae (HD 39091; Table 1), a naked-eye ($V=5.65$ mag), relatively inactive ($\log R'_{\text{HK}} = -4.941$; Gray et al. 2006) G0 V star already known to host a sub-stellar companion (π Men b) on a ~ 2100 -day eccentric ($e \approx 0.6$) orbit (Jones et al. 2002). π Men c is a $2.06 R_{\oplus}$ planet with an orbital period of 6.27 days. Using *Gaia* photometry, archival HARPS Doppler data, and published UCLES high-precision radial velocities (RVs) we confirmed the planetary nature of the transiting signal detected by *TESS* and derived the planet's mass. We note that in the final stage of preparing this manuscript, an independent investigation of this system was publicly announced by Huang et al. (2018).

¹ Available at <https://tess.mit.edu/alerts/>.

Table 1. Main identifiers, coordinates, and parallax, optical, and infrared magnitudes of π Men.

Parameter	Value	Source
HD	39091	
TIC ID	261136679	TIC
TOI ID	144	TESS Alerts
Gaia DR2 ID	4623036865373793408	Gaia DR2
RA (J2000)	05 37 09.885	Gaia DR2
RA (J2000)	-80 28 08.831	Gaia DR2
π	54.705 ± 0.0671 mas	Gaia DR2
V	5.65 ± 0.01	Mermilliod (1987)
B	6.25 ± 0.01	Mermilliod (1987)
J	4.869 ± 0.272	2MASS
H	4.424 ± 0.226	2MASS
K_s	4.241 ± 0.027	2MASS
G	5.4907 ± 0.0014	Gaia DR2
G_{BP}	5.8385 ± 0.0041	Gaia DR2
G_{RP}	5.0643 ± 0.0034	Gaia DR2

Note – Values marked with TIC, Gaia DR2, and 2MASS are from Stassun et al. (2018), Gaia Collaboration et al. (2018), and Cutri et al. (2003), respectively.

2. TESS photometry

We downloaded the TESS Sector 1 light curves from the MIT web site. For the TESS object of interest TOI-144 (π Men, HD 39091, TIC 261136679), the light curve is provided by the NASA Ames Science Processing Operations Center (SPOC). The time-series includes 18 036 short-cadence ($T_{\text{exp}} = 2$ min) photometric measurements. TESS observations started on 25 July 2018 and ended on 22 August 2018. We removed any measurements that have a non-zero “quality” flag, that is, those suffering from cosmic rays or instrumental issues. We removed any long-term stellar variability by fitting a cubic spline with a width of 0.75 days. We searched the light curve for transit signals using the Box-least-squares algorithm (BLS; Kovács et al. 2002). We detected the signal of π Men c with a signal-to-noise ratio (S/N) of 9.1 and our ephemeris is consistent with that reported by the TESS team. We did not find any additional transit signal with $S/N > 6$. We also performed a periodogram and auto-cross-correlation analysis in an attempt to extract the rotation period of the star from the out-of-transit TESS light curve, but we found no significant rotation signal in the light curve.

3. Limits on photometric contamination

As a result of the $\sim 21''$ pixel scale of the TESS detectors, photometric contamination due to chance alignment with a background source is more likely than in previous transit surveys, such as *Kepler*. We investigated this possibility using archival images of π Men from the SERC-J and AAO-SES surveys² and Gaia DR2 (Gaia Collaboration et al. 2018). The TESS photometric aperture used to create the SPOC light curve is approximately 6×6 TESS pixels in extent. We thus executed a query of Gaia DR2 centered on the coordinates of π Men from the TESS Input Catalog³ (TIC; Stassun et al. 2018) using a search radius of $2'$. The archival images were taken in 1978 and 1989, and therefore π Men appears significantly offset from its current position

² Available at http://archive.stsci.edu/cgi-bin/dss_form.

³ Available at <https://mast.stsci.edu/portal/Mashup/Clients/Mast/Portal.html>.

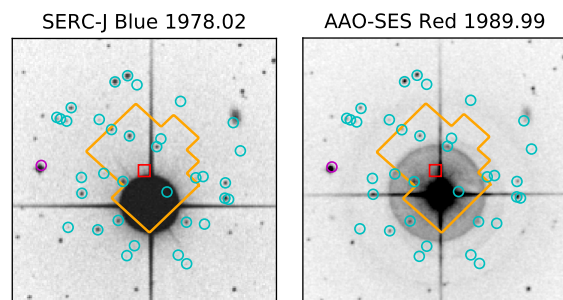


Fig. 1. $5' \times 5'$ archival images with the SPOC photometric aperture overlaid in orange, the Gaia DR2 position (J2015.5) of π Men indicated by a red square, and other Gaia DR2 sources within $2'$ of π Men indicated by circles. The magenta circle indicates the position of Gaia DR2 4623036143819289344, a nearby source bright enough to be the host of the observed transit signals (see Section 3), and cyan circles indicate sources that are too faint.

due to proper motion; no background source is visible near the current position of π Men. Figure 1 shows Gaia DR2 source positions overlaid on the archival images, along with the SPOC photometric aperture.

Assuming a maximum eclipse depth of 100%, the measured transit depth (see Section 9) puts an upper limit on the magnitude of a putative eclipsing binary within the photometric aperture, since a fainter source would be overly diluted by the flux from π Men. As the Gaia G_{RP} band-pass is a good approximation to the TESS band-pass, we find a limiting magnitude of $G_{RP,\text{max}} = 14.1$ mag. Assuming an aperture radius of $60''$ ($120''$), a simulated stellar population along the line of sight to π Men from TRILEGAL⁴ (Girardi et al. 2005) implies a frequency of 0.3578 (1.4312) stars brighter than $G_{RP,\text{max}}$. Indeed, only one other Gaia DR2 source within $2'$ of π Men is brighter than $G_{RP,\text{max}}$, consistent with the expectation from TRILEGAL: Gaia DR2 4623036143819289344 ($G_{RP} = 12.1644 \pm 0.0011$ mag, separation $\approx 118''$). As this source is clearly outside of the TESS photometric aperture, we conclude that π Men is the true host of the transit signal as seen by TESS, and that photometric dilution from sources other than π Men is negligible.

4. Custom light curve preparation

Having established that π Men is the host of the transit signal based on the above analysis of the SPOC aperture and the corresponding light curve, we then performed an analysis of the TESS pixel data by downloading the target pixel file (TPF) for π Men from the MAST web page⁵. We extracted a series of light curves from the pixel data by producing aperture masks containing different contiguous sets of pixels centered on π Men, which appeared elongated due to CCD blooming preferentially along the columns of the detector. The apertures were produced by computing the 50th to 95th percentiles of the median image and selecting pixels with median counts above each percentile value. We then selected an optimal aperture by minimizing the 6.5 hour combined differential photometric precision (CDPP) noise metric (Christiansen et al. 2012). We found that the light curves exhibited elevated noise levels between BJD–2450000 \approx 8347–8350, which corresponded to a deviation in spacecraft pointing.

⁴ Available at <http://stev.oapd.inaf.it/cgi-bin/trilegal>.

⁵ Available at <https://archive.stsci.edu/prepds/tess-data-alerts/>.

As no transits occur during this time we decided to mask it from each light curve before computing the value of the CDPD noise metric. An aperture somewhat larger than the SPOC aperture yielded the light curve with the best CDPD value, and we use this for our subsequent transit analysis. We note that, although our custom aperture is larger than the SPOC aperture, π Men is so much brighter than other nearby stars that photometric dilution remains negligible compared to the uncertainties of the data.

5. UCLES and HARPS archival spectra

Jones et al. (2002) reported on the detection of a long-period ($P_{\text{orb}} \approx 2100$ days), eccentric ($e \approx 0.6$), sub-stellar companion to π Men with a minimum mass of $M_b \approx 10 M_{\text{Jup}}$. Their discovery is based on 28 RV measurements obtained between November 1998 and April 2002 using the UCLES spectrograph mounted at the 3.92-m Anglo-Australian Telescope at Siding Spring Observatory. Fourteen additional UCLES RVs taken between August 2002 and October 2005 were published by Butler et al. (2006). For the sake of clarity, we list the 42 UCLES RVs in Table 2.

We also retrieved from the ESO public archive 145 high-resolution ($R \approx 115\,000$) spectra of π Men, taken with the HARPS spectrograph (Mayor et al. 2003) mounted at the ESO-3.6m telescope of La Silla observatory (Chile). The observations were carried out between 28 December 2012 and 17 March 2017 UTC, as part of the observing programs 072.C-0488, 183.C-0972, and 192.C-0852. The retrieved data-set includes Echelle and order-merged spectra in flexible image transport system (FITS) format, along with additional FITS files containing the cross-correlation function (CCF) and its bisector, computed from the HARPS pipeline using a G2 numerical mask. From the FITS headers, we extracted the barycentric Julian dates, the RVs, and their uncertainties, along with the full width at half maximum (FWHM) and bisector span (BIS) of the CCF, and the S/N per pixel at 5500 Å. On June 2015, the HARPS fiber bundle was upgraded (Lo Curto et al. 2015). To account for the RV offset caused by the instrument refurbishment, we treated the HARPS RVs taken before/after June 2015 as two different data sets (Table 3 and 4). Following Eastman et al. (2010), we converted the heliocentric Julian dates (HJD_{UTC}) of the UCLES time stamps and the barycentric Julian (BJD_{UTC}) of the HARPS time stamps into barycentric Julian dates in barycentric dynamical time (BJD_{TDB}).

6. Stellar fundamental parameters

We determined the spectroscopic parameters of π Men from the co-added HARPS spectrum, which has a S/N per pixel of ~ 1880 at 5500 Å. We used Spectroscopy Made Easy (SME; Valenti & Piskunov 1996; Valenti & Fischer 2005; Piskunov & Valenti 2017), a spectral analysis tool that calculates synthetic spectra and fits them to high-resolution observed spectra using a χ^2 minimizing procedure. The analysis was performed with the non-LTE SME version 5.2.2, along with ATLAS 12 one-dimensional model atmospheres (Kurucz 2013).

In order to estimate the micro-turbulent (v_{mic}) and macro-turbulent (v_{mac}) velocities, we used the empirical calibration equations for Sun-like stars from Bruntt et al. (2010) and Doyle et al. (2014), respectively. The effective temperature T_{eff} was measured fitting the wings of the H_α line (Fuhrmann et al. 1993; Axer et al. 1994; Fuhrmann et al. 1994, 1997b,a). We excluded the core of H_α because of its origin in higher layers of stellar photospheres. The surface gravity $\log g_\star$ was determined from

the wings of the Ca I $\lambda 6102$, $\lambda 6122$, $\lambda 6162$ Å triplet, and the Ca I $\lambda 6439$ Å line. We measured the iron abundance $[\text{Fe}/\text{H}]$ and projected rotational velocity $v \sin i_\star$ by simultaneously fitting the unblended iron lines in the spectral region 5880–6600 Å.

We found an effective temperature of $T_{\text{eff}} = 5870 \pm 50$ K, surface gravity $\log g_\star = 4.33 \pm 0.09$ (cgs), and an iron abundance relative to solar of $[\text{Fe}/\text{H}] = 0.05 \pm 0.09$ dex. We measured a $[\text{Ca}/\text{H}]$ abundance of 0.07 ± 0.10 dex. The projected rotational velocity was found to be $v \sin i_\star = 3.3 \pm 0.5$ km s⁻¹, with $v_{\text{mic}} = 1.06 \pm 0.10$ km s⁻¹ and $v_{\text{mac}} = 3.35 \pm 0.4$ km s⁻¹. These values were confirmed by modeling the Na I doublet at 5889.95 and 5895.924 Å. We detected no interstellar sodium, as expected given the vicinity of the star ($d=18.3$ pc).

We used the BAYesian STellar Algorithm (BASTA, Silva Aguirre et al. 2015) with a large grid of GARSTEC stellar models (Weiss & Schlattl 2008) to derive the fundamental parameters of π Men. We built the spectral energy distribution (SED) of the star using the magnitudes listed in Table 1, and then fitted the SED along with our spectroscopic parameters (T_{eff} , $\log g_\star$, $[\text{Fe}/\text{H}]$) and *Gaia* parallax to a grid of GARSTEC models. Following Luri et al. (2018), we quadratically added 0.1 mas to the nominal uncertainty of *Gaia* parallax to account for systematic uncertainties of *Gaia* astrometry. We adopted a minimum uncertainty of 0.01 mags for the *Gaia* magnitudes to account for systematic uncertainties in the *Gaia* photometry. Given the proximity of the star ($d=18.3$ pc), we assumed no interstellar reddening.

We found that π Men has a mass of $M_\star = 1.02 \pm 0.03 M_\odot$ and a radius of $R_\star = 1.10 \pm 0.01 R_\odot$, implying a surface gravity of $\log g_\star = 4.36 \pm 0.02$ (cgs), in agreement with the spectroscopically derived value of 4.33 ± 0.09 . The stellar models constrain the age of the star to be 5.2 ± 1.1 Gyr. The fundamental parameters of π Men are given in Table 5. We stress that the uncertainties on the derived parameters are internal to the stellar models used and do not include systematic uncertainties related to input physics or to the bolometric correction.

7. Seismic analysis

In order to better characterize the star, we performed an asteroseismic analysis of the SPOC light curve as well as the custom light curve optimized for characterization of the exoplanet (Sect. 4). For the former, the corrections we made consist of three steps. First, we corrected the SPOC flux⁶ performing a robust locally weighted regression as described in Cleveland (1979) in order to smooth long-period variation from the light curve without removing any transit signal. We also calibrated the data following the methods described in García et al. (2011). The results of both analyses provided similar seismic results, although the corrections applied were very different. As a second step we removed the transits by folding the light curve at the period of the planet transit and filtering it with a wavelet transform using an “à trous” algorithm (Starck & Murtagh 2002, 2006). Finally as the last step, the gaps of the resultant light curve were interpolated using inpainting techniques following García et al. (2014) and Pires et al. (2015). For the custom aperture (Sect. 4), the first step consisted of correcting the light curve following only García et al. (2011) (as the two corrections methods led to similar results) and we applied the same steps two and three.

We used the FliPer metric (Bugnet et al. 2018) to estimate $\log g_\star$ directly from the global power of the power spectrum density of both light curves. Unfortunately, due to the high level of

⁶ We used the pre-search data conditioning simple aperture photometry (Smith et al. 2012).

noise and the filters applied to the light curve to flatten it and properly remove the transits, part of the power below $\sim 100 \mu\text{Hz}$ is removed providing only a loose limit of the value of surface gravity or the frequency of maximum power of the modes. We applied the standard seismic A2Z pipeline (Mathur et al. 2010) to look for the power excess due to acoustic modes. While the blind search did not yield a significant detection in either light curve, we estimated where we would expect the acoustic modes given the spectroscopic parameters derived in this paper. The modes are expected around $2500 \mu\text{Hz}$. The power spectrum of SPOC light curve shows a slight excess of power around $2600 \mu\text{Hz}$ (frequency of maximum power or ν_{max}) and the A2Z pipeline that computes the power spectrum of the power spectrum detects a signal at $119.98 \pm 9.25 \mu\text{Hz}$, which could be the large frequency separation ($\Delta\nu$ is the frequency difference between 2 modes of same degree and consecutive orders) with a 95% confidence level. This value corresponds to the $\Delta\nu$ expected from the global seismic scaling relations for modes at $2607 \pm 16 \mu\text{Hz}$. Although quite unlikely, this signal could be due to noise. If we consider that this is a real seismic signature, using $\Delta\nu$, ν_{max} , and T_{eff} , along with the solar scaling relations, we found a stellar mass of $M_{\star} = 1.02 \pm 0.15 M_{\odot}$ and a stellar radius of $R_{\star} = 1.09 \pm 0.10 R_{\odot}$, in agreement with the spectroscopic values (Sect. 6). However, the asteroseismic analysis of the custom light curve does not reproduce the previous results, as the seismic signal is not statistically significant. A better knowledge of the optimal aperture mask for asteroseismology is required at this point, which is beyond the scope of this paper. We note that this marginal detection is expected for a star of this magnitude and evolutionary stage observed over the course of one month. Asteroseismic analysis with *TESS* will require multiple-sector data and/or more evolved bright stars, as predicted by pre-launch simulations (Campante 2017).

8. Frequency analysis of the Doppler data

We performed a frequency analysis of the UCLES and HARPS RVs in order to search for the presence of the transiting planet in the Doppler data, and look for possible additional periodic signals. The generalized Lomb-Scargle (GLS; Zechmeister & Kürster 2009) periodogram of the combined UCLES and HARPS RVs⁷ shows a very significant peak at the orbital frequency of the outer sub-stellar companion ($f_b = 0.0005 \text{ c/d}$), with a false-alarm probability (FAP) lower than 10^{-10} .

The upper panel of Fig. 2 displays the GLS periodogram of the RV residuals following the subtraction of the Doppler signal induced by the outer sub-stellar object. We found that the most significant peak is seen at the frequency of the transit signal detected by *TESS* ($f_c = 0.16 \text{ c/d}$), with a FAP $< 10^{-5}$ and an RV semi-amplitude of $\sim 1.5 \text{ m s}^{-1}$. The peak has no counterparts in the periodograms of the HARPS activity indicators (second and third panels of Fig. 2), suggesting that the signal is induced by the presence of an orbiting planet with a period of 6.27 days.

We also subtracted the Doppler reflex motion induced by the transiting planet from our RV data and searched the residuals for additional periodic signal but found no peak with FAP $< 10^{-4}$.

9. Joint analysis of the transit and Doppler data

We performed a joint analysis of our custom light curve (Sect. 4) and RV measurements (Sect. 5) using the software suite

⁷ We accounted for the instrumental offsets using the values derived in Sect. 9.

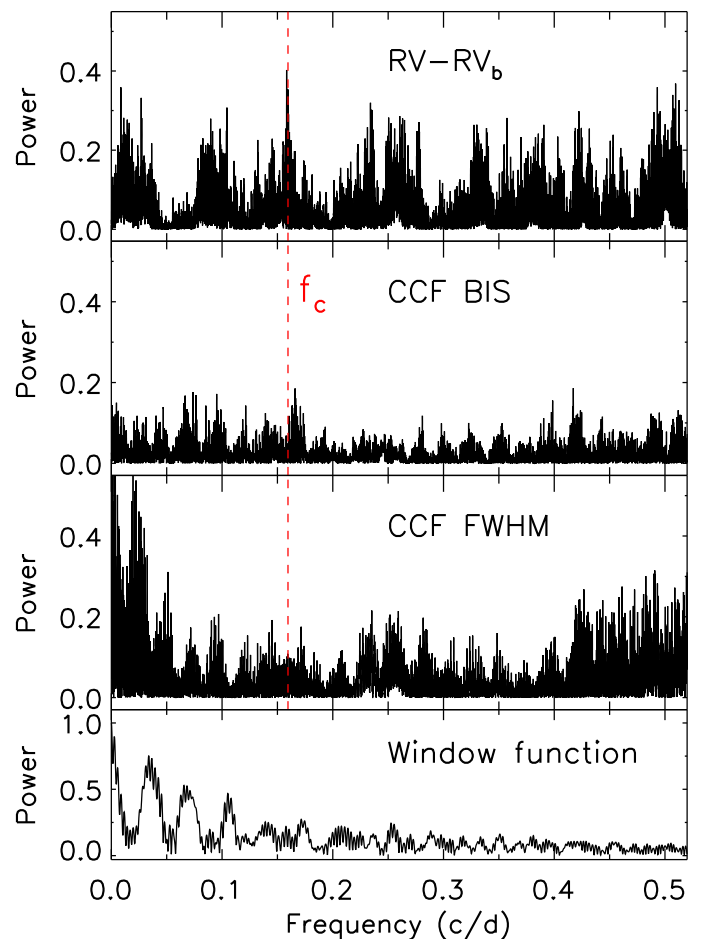


Fig. 2. *First panel:* GLS periodogram of the UCLES and HARPS RV residuals following the subtraction of the Doppler reflex motion induced by the outer sub-stellar companion. *Second and third panels:* GLS periodogram of the BIS and FWHM of the HARPS CCF (data acquired with the old fiber bundle). *Fourth panel:* Periodogram of the window function of the combined RV measurements. The dashed vertical red line marks the orbital frequency of the transiting planet ($f_c = 0.16 \text{ c/d}$).

pyaneti (Barragán et al. 2019), which allows for parameter estimation from posterior distributions calculated using Markov chain Monte Carlo methods.

We extracted 10 hours of *TESS* data points centered around each of the five transits observed by *TESS*. The five segments were de-trended using the code *exotrending* (Barragán & Gandolfi 2017), fitting a second-order polynomial to the out-of-transit data. We used all 187 Doppler measurements presented in Sect. 5 and accounted for the RV offsets between the different instruments and the two HARPS setups.

The RV model consists of two Keplerians to account for the Doppler signal induced by planets *b* and *c*. We fitted for a RV jitter term for each instrument/setup to account for instrumental noise not included in the nominal uncertainties, and/or to account for any stellar-activity-induced RV variation. We modeled the *TESS* transit light curves using the limb-darkened quadratic model of Mandel & Agol (2002). For the limb darkening coefficients, we set Gaussian priors using the values derived by Claret (2017) for the *TESS* pass-band. We imposed conservative error bars of 0.1 on both the linear and the quadratic limb-darkening term. A preliminary analysis showed that the transit light curve poorly constrains the scaled semi-major axis (a/R_{\star}). We therefore set a Gaussian prior on a/R_{\star} using Kepler’s third’s law, the

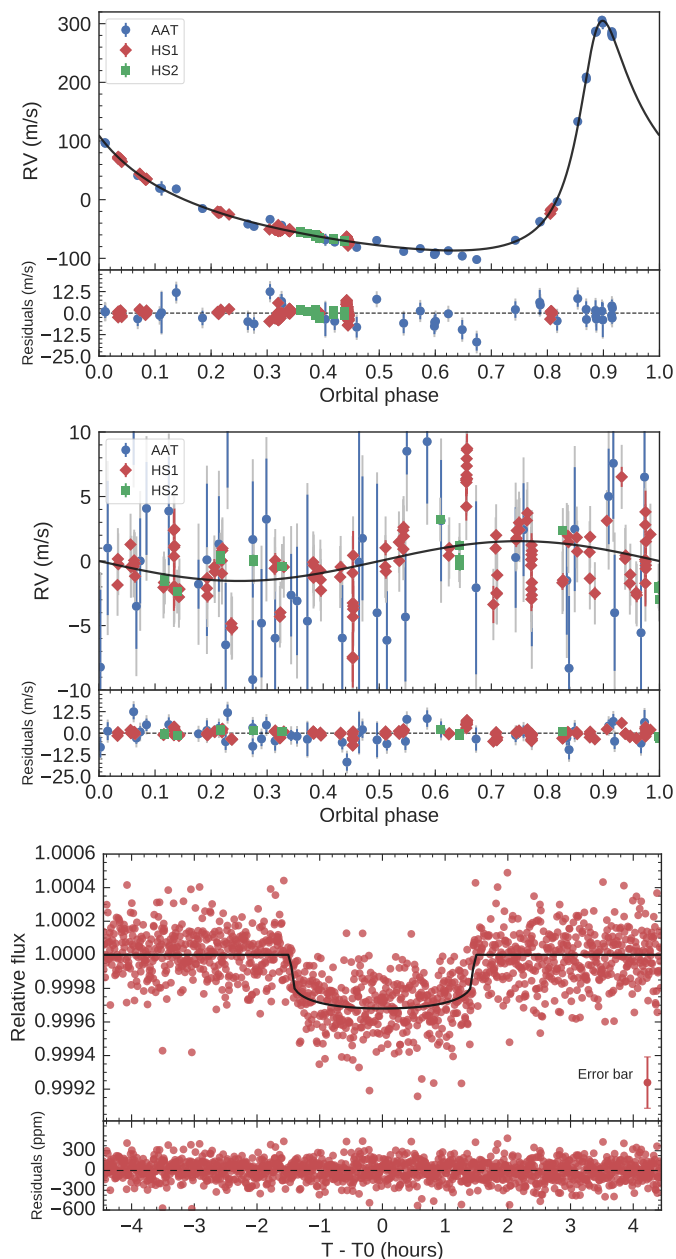


Fig. 3. Phase-folded RV curves of π Men b (upper panel) and c (middle panel), and transit light curve of π Men c (lower panel). The best-fitting transit and Keplerian models are overplotted with thick black lines. The zero phases of the RV curves of π Men b (upper panel) and c (lower panel) correspond to the time of inferior conjunction and transit epoch, respectively (Table 5). The *TESS* data points are shown with red circles (lower panel). The AAT data and the two sets of HARPS RVs (HS1 and HS2) are shown with circles, diamonds, and squares, respectively. The vertical gray lines mark the error bars including jitter.

orbital period, and the derived stellar parameters (Sect. 6). Because the eccentricity of planet *c* is poorly constrained by the observations, we fixed it to zero for our analysis (see also Section 10). We imposed uniform priors for the remaining fitted parameters. Details of the fitted parameters and prior ranges are given in Table 5. Before performing the final analysis, we ran a numerical experiment to check if the *TESS* two-minute integration time needed to be taken into account following Kipping (2010). We found no differences in the posterior distributions

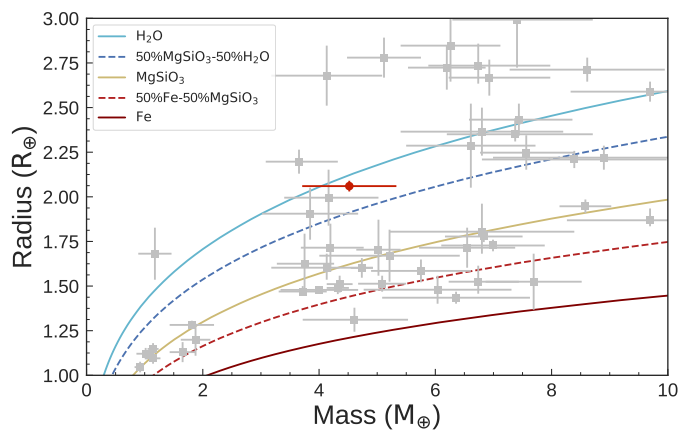


Fig. 4. Mass-radius for low-mass ($M_p < 10 M_\oplus$), small ($R_p < 3 M_\oplus$) planets with mass-radius measurements better than 25% (from <http://www.astro.keele.ac.uk/jkt/tepcat/>; Southworth et al. 2007). Composition models from Zeng et al. (2016) are displayed with different lines and colors. The solid red circle marks the position of π Men c.

for fits with and without re-sampling; we therefore proceeded with our analysis without re-sampling. We used 500 independent Markov chains initialized randomly inside the prior ranges. Once all chains converged, we used the last 5 000 iterations and saved the chain states every ten iterations. This approach generates a posterior distribution of 250 000 points for each fitted parameter. Table 5 lists the inferred planetary parameters. They are defined as the median and 68% region of the credible interval of the posterior distributions for each fitted parameter. The transit and RV curves are shown in Fig. 3.

10. Discussion and conclusion

π Men is a bright ($V=5.65$ mag) Sun-like star (SpT=G0 V) known to host a sub-stellar companion (π Men b) on a long-period eccentric orbit (Jones et al. 2002). Combining *Gaia* photometry with archival RV measurements we confirmed that the $P=6.27$ day transit signal detected in the *TESS* light curve of π Men is caused by a *bona fide* transiting super-Earth and derived its mass. π Men c becomes *TESS*'s first confirmed planet.

π Men joins the growing number of stars known to host both long-period Jupiter analogs and close-in small planets ($R_p < 4 R_\oplus$). Zhu & Wu (2018) pointed out that long-period, gas-giant planets are common around stars hosting super-Earths. Bryan et al. (2018) recently found that the occurrence rate of companions in the range $0.5\text{--}20 M_{\text{Jup}}$ at $1\text{--}20$ AU in systems known to host inner small planets is $39 \pm 7\%$, suggesting that the presence of outer gas giant planets does not prevent the formation of inner Earth- and Neptune-size planets. We performed a dynamical stability analysis of π Men using the software *mercury6* (Chambers 1999). Assuming co-planar orbits, we let the system evolve for 100 000 yr. For π Men b we found negligible changes of the semi-major axis and eccentricity of $< 2.6 \times 10^{-3}$ AU and 3×10^{-4} , respectively. For π Men c we found no variation larger than 1×10^{-5} of its semi-major axis, with changes of its eccentricity $\lesssim 0.05$.

The actual orientation of the outer planet's orbit is unknown. While we know the inner planet's inclination, because it transits, its eccentricity is poorly constrained by the data. Compact multi-planet systems have been observed to have near-zero eccentricities (e.g., Hadden & Lithwick 2014; Van Eylen & Albrecht 2015; Xie et al. 2016). However, planets with only a single transiting

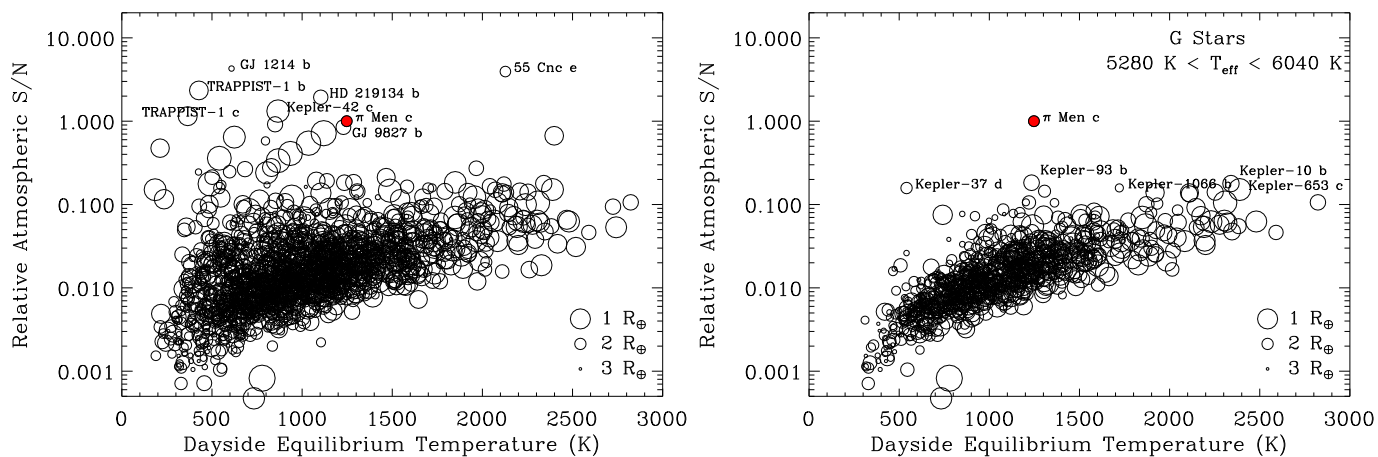


Fig. 5. *Left Panel:* Relative S/N of an atmospheric signal for all exoplanets with $R_p < 3 R_\oplus$ as a function of planetary equilibrium temperature. The symbol size is inversely proportional to the planetary radius, emphasizing those closest to an Earth radius. The π Men c planet is used as the atmospheric signal reference and it is indicated by the filled circle. It is among the top ten most favorable planets for atmospheric characterization. *Right panel:* As in the left panel, simply limited to solar type host stars (i.e., G-type stars; $5280 < T_{\text{eff}} < 6040$ K). The π Men c planet is by far the most favorable planet around such a star for atmospheric characterization. The other optimal atmospheric targets all transit K and M stars. For this reason, the coronal and stellar wind properties that interact with the π Men c atmosphere may be distinctly different to those experienced by the rest of the sample.

planet often appear to be “dynamically hotter”, and many have a non-zero eccentricity, which can, for example, be described by the positive half of a zero-mean Gaussian distribution, with a dispersion $\sigma_e = 0.32 \pm 0.06$ (Van Eylen et al. 2018b). The outer planet, π Men c, has an orbital eccentricity of ~ 0.64 . A far-out giant planet, such as planet c, may in fact increase the orbital eccentricity of a close-in super-Earth, such as planet b (see, e.g., Mustill et al. 2017; Hansen 2017; Huang et al. 2017). Following Van Eylen et al. (2018b), we found an orbital eccentricity based on the transit data alone of $[0, 0.45]$ at 68% confidence. Because the current RV observations cannot constrain the eccentricity either, we fixed it to zero in the above analysis (see Section 9). The brightness of the host star makes this planetary system an exciting target for further RV follow-up to measure the inner planet’s eccentricity.

The transiting planet π Men c has a mass of $M_c = 4.52 \pm 0.81 M_\oplus$ and a radius of $R_c = 2.06 \pm 0.03 R_\oplus$, yielding a mean density of $\rho_c = 2.82 \pm 0.53 \text{ g cm}^{-3}$. Figure 4 shows the mass-radius diagram for small planets whose masses have been determined with a precision better than 25%. Theoretical models from Zeng et al. (2016) are overplotted using different lines and colors. The position of π Men c suggests a composition of Mg silicates and water. Alternatively, the planet might have a solid core surrounded by a gas envelope. At short orbital periods, super-Earths and sub-Neptunes are separated by a radius gap at $\approx 1.8 R_\oplus$ (Fulton et al. 2017; Van Eylen et al. 2018a). The exact location of the radius gap is observed to be a function of the orbital period (Van Eylen et al. 2018a), as predicted by models of photoevaporation (e.g., Owen & Wu 2013; Lopez & Fortney 2013). Van Eylen et al. (2018a) found that the radius gap is located at $\log R = m \times \log P + a$, where $m = -0.09^{+0.02}_{-0.04}$ and $a = 0.37^{+0.04}_{-0.02}$. At the orbital period of π Men c, that is, $P_{\text{orb}} = 6.27$ days, the radius gap is then located at $R_p = 1.99 \pm 0.20 R_\oplus$. This implies that π Men c, with a radius of $R_c = 2.06 \pm 0.03 R_\oplus$, is located just around the radius gap, although the measured density suggests that the planet may have held on to (part of) its atmosphere.

The naked-eye brightness of π Men immediately argues that any transiting planet will be attractive for atmospheric char-

acterization. Observations of a planetary atmosphere through transmission spectroscopy during transit provide opportunities to measure the extent, kinematics, abundances, and structure of the atmosphere (Seager & Deming 2010). Such measurements can be utilized to address fundamental questions such as planetary atmospheric escape and interactions with the host star (Cauley et al. 2017), formation and structure of planetary interiors (Owen et al. 1999), planetary and atmospheric evolution (Öberg et al. 2011), and biological processes (Meadows & Seager 2010).

The left panel of Fig. 5 displays a relative atmospheric detection S/N metric normalized to π Men c for all known small exoplanets with $R_p < 3 R_\oplus$. The sample is taken from the Exoplanet Orbit Database⁸ as of September 2018. The atmospheric signal is calculated in a similar way to Gillon et al. (2016) and Niraula et al. (2017). In particular, the relative atmospheric S/N plotted in Fig. 5 is calculated from Eq. 1a and 1b from Niraula et al. (2017). This calculation assumes similar atmospheric properties (e.g., Bond albedo, mean molecular weight) for all planets. Large atmospheric signals result from hot, extended atmospheres of planets that transit small, cool stars. For this reason, planets transiting such stars as GJ 1214 and TRAPPIST-1 are excellent targets for this kind of study. Nevertheless, the brightness of the host star along with the period and duration of the transit also significantly contribute to the ability to build up a sufficiently high S/N to detect atmospheric signals. We used the J-band flux (e.g., H_2O measurements with *JWST*; Beichman et al. 2014) and weighted the metric to optimize the S/N ratio over a period of time rather than per transit.

In the context of all small planets ($R_p < 3 R_\oplus$), π Men c has the seventh strongest atmospheric signal, behind GJ 1214 b, 55 Cnc e, the TRAPPIST-1 b and c planets, HD219134 b, and Kepler-42 c. However, π Men is unique among this notable group of stars in that it is the only G-type star (Fig. 5, right panel). All of the other exoplanets transit K- or M-type stars. The brightness of π Men is able to overcome the disadvantage of a small

⁸ Available at exoplanets.org.

planet transiting a slightly larger star, to provide the best opportunity of probing the atmospheric properties of a super-Earth orbiting a solar type star. Given the significant changes in the structure of stellar coronae and stellar winds between G- and M-type stars, the atmospheric properties and evolution of π Men c may be distinctly different from the atmospheres detected around the sample of very-low-mass M-type stars (e.g., GJ 1214 and TRAPPIST-1). For example, the TRAPPIST-1 e, f, and g planets essentially orbit within the stellar corona of the host star and may be subject to a substantial stellar wind, which will result in a strong injection of energy in the atmosphere and may prevent the formation of a significant atmosphere (Cohen et al. 2018). When inferring the properties of coronae and winds of stars other than the Sun, we often have to use poorly constrained models and empirical correlations, the validity of which are best for stars that are quite similar to the Sun. In this respect, π Men is a unique laboratory because of its greater similarities to the Sun with respect to all the other stars known to host mini-Neptunes and Super-Earths amenable to multi-wavelength atmospheric characterization.

We further studied the long-term stability of a possible hydrogen-dominated atmosphere by estimating the mass-loss rates. To this end, we employed the interpolation routine described by Kubyshkina et al. (2018), which interpolates the mass-loss rate among those obtained with a large grid of one-dimensional upper-atmosphere hydrodynamic models for super-Earths and sub-Neptunes. Employing the values listed in Table 5 and a Sun-like high-energy emission, which is a reasonable assumption given that π Men has a temperature and age similar to those of the Sun, we obtained a mass-loss rate of $1.2 \times 10^{10} \text{ g s}^{-1}$, which corresponds to $\sim 1.4\%$ of the estimated planetary mass per gigayear. This indicates that the question of whether the planet still holds a hydrogen-dominated atmosphere or not greatly depends on the initial conditions, namely how much hydrogen the planet accreted during its formation. If the planet originally accreted a small hydrogen-dominated atmosphere with a mass of only a few percent of the total planetary mass, we can expect it to be, for the vast majority, lost, particularly taking into account that the star was more active in the past. In contrast, a significant hydrogen mass fraction would still be present if the planet originally accreted a large amount of hydrogen. The inferred bulk density hints at the possible presence of a hydrogen-dominated atmosphere, but it does not give a clear indication. Ultraviolet observations aiming at detecting hydrogen Ly α absorption and/or carbon and oxygen in the upper planetary atmosphere would be decisive in identifying its true nature.

While this manuscript was in preparation, Huang et al. (2018) publicly reported an independent analysis of the π Men system, using a custom *TESS* photometry pipeline. They derived a mass and radius for π Men c consistent within less than 1- σ with our measurements.

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(ESA) mission *Gaia* (<https://www.cosmos.esa.int/gaia>), processed by the *Gaia* Data Processing and Analysis Consortium (DPAC, <https://www.cosmos.esa.int/web/gaia/dpac/consortium>). Funding for the DPAC has been provided by national institutions, in particular the institutions participating in the *Gaia* Multilateral Agreement. We acknowledge the traditional owners of the land on which the AAT stands, the Gamilaraay people, and pay our respects to elders past and present. We are also very grateful to the referee, Steve Howell, for his helpful comments and suggestions, which resulted in an improved paper. We also thank him for his rapid review of the paper.

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Table 2. UCLES RV measurements of π Men.

BJD _{TDB} ^a -2450000	RV (km s ⁻¹)	$\pm\sigma$ (km s ⁻¹)
829.993723	-0.0410	0.0048
1119.251098	-0.0674	0.0098
1236.033635	-0.0792	0.0060
1411.325662	-0.0858	0.0058
1473.267712	-0.0800	0.0048
1526.081162	-0.0930	0.0046
1527.082805	-0.0898	0.0041
1530.128708	-0.0879	0.0045
1629.912366	-0.0927	0.0056
1683.842991	-0.1005	0.0050
1828.188260	-0.0674	0.0048
1919.099660	-0.0350	0.0072
1921.139081	-0.0373	0.0047
1983.919846	-0.0028	0.0056
2060.840355	0.1361	0.0048
2092.337359	0.2120	0.0047
2093.352231	0.2094	0.0044
2127.328562	0.2878	0.0059
2128.336410	0.2861	0.0042
2130.339049	0.2899	0.0067
2151.292440	0.3079	0.0052
2154.305009	0.3030	0.0100
2187.196618	0.2857	0.0039
2188.236606	0.2893	0.0037
2189.223031	0.2797	0.0033
2190.145881	0.2835	0.0037
2387.871387	0.1009	0.0036
2389.852023	0.0974	0.0033
2510.307394	0.0417	0.0042
2592.126975	0.0202	0.0032
2599.155380	0.0210	0.0120
2654.099326	0.0188	0.0047
2751.918480	-0.0117	0.0042
2944.224628	-0.0434	0.0038
3004.075458	-0.0321	0.0044
3042.078745	-0.0440	0.0042
3043.018085	-0.0463	0.0045
3047.050110	-0.0408	0.0043
3048.097508	-0.0444	0.0036
3245.311649	-0.0697	0.0050
3402.035747	-0.0669	0.0018
3669.244092	-0.0863	0.0019

Notes:

^a Barycentric Julian dates are given in barycentric dynamical time.

Table 3. HARPS RV measurements of π Men acquired with the old fiber bundle. The entire RV data set is available in a machine-readable table in the online journal.

BJD _{TDB} ^a -2450000	RV (km s ⁻¹)	$\pm\sigma$ (km s ⁻¹)	BIS (km s ⁻¹)	FWHM (km s ⁻¹)	T _{exp} (s)	S/N ^b
3001.830364	10.6600	0.0014	-0.0019	7.6406	109	69.0
3034.607261	10.6665	0.0008	0.0040	7.6368	200	120.7
3289.869718	10.6448	0.0012	-0.0013	7.6378	60	79.6
3289.870782	10.6428	0.0011	-0.0012	7.6406	60	89.4
3289.871836	10.6446	0.0011	-0.0007	7.6394	60	90.7
3289.872866	10.6449	0.0011	-0.0026	7.6439	60	86.5
...

Notes:

^a Barycentric Julian dates are given in barycentric dynamical time.^b S/N per pixel at 550 nm.**Table 4.** HARPS RV measurements of π Men acquired with the new fiber bundle.

BJD _{TDB} ^a -2450000	RV (km s ⁻¹)	$\pm\sigma$ (km s ⁻¹)	BIS (km s ⁻¹)	FWHM (km s ⁻¹)	T _{exp} (s)	S/N ^b
7298.853243	10.6750	0.0005	0.0081	7.6856	450	187.3
7298.858243	10.6747	0.0004	0.0083	7.6842	450	242.8
7327.755817	10.6744	0.0003	0.0089	7.6870	900	305.7
7354.783687	10.6674	0.0002	0.0104	7.6867	900	538.0
7357.725912	10.6727	0.0002	0.0105	7.6858	900	542.9
7372.705131	10.6664	0.0004	0.0094	7.6825	300	273.0
7372.708997	10.6662	0.0004	0.0118	7.6822	300	247.3
7372.712758	10.6654	0.0003	0.0104	7.6831	300	320.9
7423.591772	10.6630	0.0005	0.0113	7.6796	450	217.1
7423.597918	10.6628	0.0005	0.0115	7.6782	450	214.1
7424.586637	10.6645	0.0004	0.0104	7.6814	450	299.7
7424.592367	10.6643	0.0004	0.0113	7.6816	450	288.5
7462.517924	10.6612	0.0003	0.0116	7.6822	450	326.8
7462.523491	10.6612	0.0003	0.0106	7.6816	450	337.8
7464.499915	10.6616	0.0005	0.0083	7.6812	300	217.7
7464.503781	10.6627	0.0004	0.0112	7.6818	300	276.2
7464.507474	10.6611	0.0004	0.0108	7.6820	300	286.4

Notes:

^a Barycentric Julian dates are given in barycentric dynamical time.^b S/N per pixel at 550 nm.

Table 5. π Men system parameters.

Parameter	Prior ^(a)	Derived value
Stellar parameters		
Star mass M_\star (M_\odot)	...	1.02 ± 0.03
Star radius R_\star (R_\odot)	...	1.10 ± 0.01
Effective Temperature T_{eff} (K)	...	5870 ± 50
Surface gravity ^(b) $\log g_\star$ (cgs)	...	4.36 ± 0.02
Surface gravity ^(c) $\log g_\star$ (cgs)	...	4.33 ± 0.09
Iron abundance [Fe/H] (dex)	...	0.05 ± 0.09
Projected rotational velocity $v \sin i_\star$ (km s^{-1})	...	3.3 ± 0.5
Age (Gyr)	...	5.2 ± 1.1
Model parameters of π Men b		
Orbital period $P_{\text{orb,b}}$ (days)	$\mathcal{U}[2076.0, 2106.0]$	2091.2 ± 2.0
Time of inferior conjunction $T_{\text{inf,b}}$ (BJD _{TDB} - 2 450 000)	$\mathcal{U}[6528.0, 6568.0]$	6548.2 ± 2.7
$\sqrt{e_b} \sin \omega_{\star,b}$	$\mathcal{U}[-1, 1]$	-0.3919 ± 0.0076
$\sqrt{e_b} \cos \omega_{\star,b}$	$\mathcal{U}[-1, 1]$	0.6970 ± 0.0052
Radial velocity semi-amplitude variation K_b (m s^{-1})	$\mathcal{U}[0, 500]$	195.8 ± 1.4
Model parameters of π Men c		
Orbital period $P_{\text{orb,c}}$ (days)	$\mathcal{U}[6.24, 6.29]$	6.26834 ± 0.00024
Transit epoch $T_{0,c}$ (BJD _{TDB} - 2 450 000)	$\mathcal{U}[8325.47, 8325.52]$	8325.503055 ± 0.00077
Scaled semi-major axis a_c/R_\star	$\mathcal{N}[13.11, 0.17]$	13.10 ± 0.18
Planet-to-star radius ratio R_c/R_\star	$\mathcal{U}[0, 0.1]$	0.01721 ± 0.00024
Impact parameter b_c	$\mathcal{U}[0, 1]$	0.614 ± 0.0180
$\sqrt{e_c} \sin \omega_{\star,c}$	$\mathcal{F}[0]$	0
$\sqrt{e_c} \cos \omega_{\star,c}$	$\mathcal{F}[0]$	0
Radial velocity semi-amplitude variation K_c (m s^{-1})	$\mathcal{U}[0, 10]$	1.55 ± 0.27
Additional model parameters		
Parameterized limb-darkening coefficient q_1	$\mathcal{N}[0.36, 0.1]$	0.33 ± 0.10
Parameterized limb-darkening coefficient q_2	$\mathcal{N}[0.25, 0.1]$	0.23 ± 0.10
Systemic velocity γ_{AAT} (km s^{-1})	$\mathcal{U}[-0.20, 0.20]$	0.0021 ± 0.0011
Systemic velocity γ_{HS1} (km s^{-1})	$\mathcal{U}[10.53, 10.88]$	10.70916 ± 0.00039
Systemic velocity γ_{HS2} (km s^{-1})	$\mathcal{U}[10.55, 10.90]$	10.73156 ± 0.00071
RV jitter term σ_{AAT} (m s^{-1})	$\mathcal{U}[0, 100]$	$4.26^{+1.10}_{-0.96}$
RV jitter term σ_{HS1} (m s^{-1})	$\mathcal{U}[0, 100]$	$2.35^{+0.19}_{-0.17}$
RV jitter term σ_{HS2} (m s^{-1})	$\mathcal{U}[0, 100]$	$1.69^{+0.39}_{-0.29}$
Derived parameters of π Men b		
Planet minimum mass $M_b \sin i_b$ (M_{Jup})	...	9.66 ± 0.20
Semi-major axis of the planetary orbit a_b (AU)	...	3.22 ± 0.03
Orbit eccentricity e_b	...	0.6394 ± 0.0025
Argument of periastron of stellar orbit $\omega_{\star,b}$ (deg)	...	330.65 ± 0.65
Time of periastron passage $T_{\text{per,b}}$ (BJD _{TDB} - 2 450 000)	...	6306.1 ± 4.6
Derived parameters of π Men c		
Planet mass M_c (M_\oplus)	...	4.52 ± 0.81
Planet radius R_c (R_\oplus)	...	2.06 ± 0.03
Planet mean density ρ_c (g cm^{-3})	...	2.82 ± 0.53
Semi-major axis of the planetary orbit a_c (AU)	...	0.06702 ± 0.00109
Orbit eccentricity e_c	...	0 (fixed)
Orbit inclination i_c (deg)	...	87.31 ± 0.11
Equilibrium temperature ^(d) $T_{\text{eq,c}}$ (K)	...	1147 ± 12
Transit duration $\tau_{14,c}$ (hours)	...	$2.969^{+0.030}_{-0.032}$

Note – ^(a) $\mathcal{U}[a, b]$ refers to uniform priors between a and b , \mathcal{U} and $\mathcal{F}[a]$ to a fixed a value. ^(b) From spectroscopy and isochrones. ^(c) From spectroscopy. ^(d) Assuming zero albedo and uniform redistribution of heat.