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Continuous-time orbit problems are decidable in polvnomial-time

^b Department of Computer Science and Information Systems, Birkbeck, University of London, United Kingdom

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1. Introduction

In this paper, we study the linear dynamical system whose dynamics is described by a linear differential equa*tion*. Formally, given a matrix $A \in \mathbb{K}^{n \times n}$ and a vector $\vec{\zeta} \in \mathbb{K}^n$, the trajectory of the system, $\vec{x}(t)$ for $t \in \mathbb{R}_{>0}$, is defined as the solution of the following Cauchy problem:

$$\begin{cases} \frac{d\bar{x}}{dt} = A\bar{x} \\ \bar{x}(0) = \bar{\zeta}. \end{cases}$$
(1)

Here \mathbb{K} is an arbitrary field and \mathbb{R} is the real field.

Linear dynamical systems have found applications in a wide range of scientific areas, for instance, theoretical biology, economics, and quantum computing. One of the basic algorithmic questions regarding a linear dynamical system is the orbit problem, which can be formulated as follows: Given the trajectory $\vec{x}(t)$ determined by $A \in \mathbb{K}^{n \times n}$

* Corresponding author.

http://dx.doi.org/10.1016/j.ipl.2014.08.004 0020-0190/© 2014 Elsevier B.V. All rights reserved. and $\vec{\zeta} \in \mathbb{K}^n$, and a point $\vec{\xi} \in \mathbb{K}^n$, decide whether there exists some time $t \in \mathbb{R}_{>0}$ such that $\vec{x}(t) = \vec{\xi}$. Namely, whether $\vec{\xi}$ can be reached from $\vec{\zeta}$.

The decidability of the orbit problem has been shown by Hainry [7], when \mathbb{K} is the rational field. In this note we improve this result by showing that it is in P. Our algorithm follows Hainry [7] in general, i.e., by Jordan norm forms and results from transcendental number theory such as the Gelfond-Schneider theorem and the Lindemann-Weierstrass theorem. However, our arguments are considerably simpler. In particular, it turns out that the distinction of two Jordan norm forms based on eigenvalues of A in [7] is unnecessary, neither is the use trigonometric functions. These simplifications enable us to perform a complexity analysis which appeared to be hard and was lacking by Hainry's arguments.

Related work. Ref. [8] studied the discrete-time orbit problem and showed that the problem is in P. The upper-bound was improved to the logspace counting hierarchy (together with a $C_{=L}$ lower-bound) [1]. The techniques employed

Taolue Chen^{a,*}, Nengkun Yu^{c,d}, Tingting Han^b

^a Department of Computer Science, Middlesex University London, United Kingdom

^c Institute for Quantum Computing, University of Waterloo, Canada

^d Department of Mathematics & Statistics, University of Guelph, Canada

ABSTRACT

We place the continuous-time orbit problem in P, sharpening the decidability result shown by Hainry [7].

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there are considerably different from the current paper. Ref. [5] considered a generalisation of the orbit problem, i.e. the orbit problem in higher dimensions, and related the problem to the celebrated Skolem problem. The authors showed that this problem is in P when the dimension is one, and is in NP^{RP} for dimension two or three. Ref. [3] studied the continuous-time Skolem problem. The authors identified decidability for this problem in some special cases, and showed that the related nonnegativity problem is NP-hard in general (whereas the decidability is left open).

2. Preliminaries

Throughout the paper, we write \mathbb{C} , \mathbb{Q} , \mathbb{A} , and \mathbb{R} for the set of complex, rational, algebraic, and real numbers, respectively. For any complex number z = a + bi where $a, b \in \mathbb{R}$ and i is the imaginary unit, we denote the real part and the imaginary part of z by $\Re(z) = a$ and $\Im(z) = b$ respectively.

Definition 1. An *algebraic number* is a number that is a root of a non-zero polynomial in one variable with rational coefficients. An algebraic number α is represented by $(P, (a, b), \rho)$ where *P* is the *minimal polynomial* of α , a + bi is an approximation of α such that $|\alpha - (a + bi)| < \rho$ and α is the only root of *P* in the open ball $\mathcal{B}(a + bi, \rho)$.

It is well known that a root of a non-zero polynomial in one variable with coefficients of algebraic numbers is also algebraic. Moreover, given the representations of two algebraic numbers α and β , the representations of $\alpha \pm \beta$, $\alpha \cdot \beta$, $\frac{\alpha}{\beta}$ can be computed in polynomial time, so is the equality checking [6].

In the sequel, we list some basic facts from transcendental number theory [2].

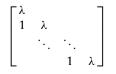
Theorem 1 (*Gelfond–Schneider*). Assume $a, b \in \mathbb{A}$ with $a \neq 0, 1$ and $b \notin \mathbb{Q}$, then any value of a^b is a transcendental number.

Corollary 1. Assume $a, b \in \mathbb{A}$ with $\ln(a), \ln(b)$ being linearly independent over \mathbb{Q} , then they are linearly independent over \mathbb{A} .

Theorem 2 (Lindemann–Weierstrass). If $\alpha_1, ..., \alpha_n$ are algebraic numbers which are linearly independent over the rational numbers \mathbb{Q} , then $e^{\alpha_1}, ..., e^{\alpha_n}$ are algebraically independent over \mathbb{Q} .

Corollary 2. For any $\alpha \neq 0$, one of α and e^{α} must be transcendental.

Definition 2. A *Jordan block* is a square matrix of the following form



A square matrix J is in Jordan norm form if

$$J = \begin{bmatrix} J_1 & & \\ & \ddots & \\ & & J_k \end{bmatrix}$$

where each J_i for $1 \le i \le k$ is a Jordan block.

The following proposition is a basic fact of linear algebra.

Proposition 1. Any matrix $A \in \mathbb{Q}^{n \times n}$ is similar to a matrix in Jordan form. Namely, there exist some $P \in \mathbb{A}^{n \times n}$ and $J \in \mathbb{A}^{n \times n}$ in Jordan form such that $A = P^{-1}JP$.

For any matrix $A \in \mathbb{C}^{n \times n}$, the *exponential* of A, denoted by e^A , is the $n \times n$ matrix given by

$$e^A = \sum_{k=0}^{\infty} \frac{1}{k!} A^k.$$

For the differential equation (1), the solution can be written as

$$\vec{x}(t) = e^{tA}\vec{\zeta},$$

and evidently the orbit problem is to determine whether there exists $t \in \mathbb{R}_{\geq 0}$ such that $e^{tA}\vec{\zeta} = \vec{\xi}$.

3. Main results

In this section we fix an instance of the orbit problem, i.e., $A \in \mathbb{Q}^{n \times n}$ and $\vec{\zeta}, \vec{\xi} \in \mathbb{Q}^n$. We consider the Jordan norm form of A such that $A = P^{-1}JP$, where $P \in \mathbb{A}^{n \times n}$ and $J \in \mathbb{A}^{n \times n}$, i.e.,

$$J = \begin{bmatrix} J_1 & 0 & \cdots & 0 \\ 0 & J_2 & \ddots & \vdots \\ \vdots & \ddots & \ddots & 0 \\ 0 & \cdots & 0 & J_k \end{bmatrix}$$

Moreover, we denote the eigenvalues for the Jordan blocks by $\lambda_1, \dots, \lambda_k$, and we write

$$\vec{x} = P\vec{\zeta} = \begin{bmatrix} \vec{x}_1 \\ \vdots \\ \vec{x}_k \end{bmatrix}$$
 and $\vec{y} = P\vec{\xi} = \begin{bmatrix} \vec{y}_1 \\ \vdots \\ \vec{y}_k \end{bmatrix}$

such that for each $1 \le i \le k$, \vec{x}_i or \vec{y}_i is of the size of J_i . For simplicity, we group the eigenvalue λ_i and the corresponding vectors \vec{x}_i and \vec{y}_i together and refer to block B_i . We say $B_i = (\lambda_i, \vec{x}_i, \vec{y}_i)$ is oblivious if $\vec{x}_i = \mathbf{0}$; otherwise, it is non-oblivious.

Theorem 3. To determine whether there exists $t \in \mathbb{R}_{\geq 0}$ such that $e^{tA}\vec{\zeta} = \vec{\xi}$ for $A \in \mathbb{Q}^{n \times n}$ and $\vec{\zeta}, \vec{\xi} \in \mathbb{Q}^n$ is in *P*.

Proof. Observe that

$$e^{tA} = e^{tP^{-1}JP} = P^{-1}e^{tJ}P,$$

and thus

$$e^{tA}\vec{\zeta} = \vec{\xi}$$
 iff $e^{tJ}(P\vec{\zeta}) = P\vec{\xi}$

Namely, $e^{tJ}\vec{x} = \vec{y}$, and thus for each $1 \le i \le k$ we have

$$e^{tJ_i}\vec{x}_i = \vec{y}_i$$

In the case that B_i is oblivious (i.e., $\vec{x}_i = \mathbf{0}$), it must be the case that $\vec{y}_i = \mathbf{0}$. In the sequel, we shall focus on the non-oblivious blocks.

Observe that

$$e^{tJ_i} = e^{t\lambda_i} \begin{bmatrix} 1 & & & \\ t & 1 & & \\ \frac{t^2}{2} & t & 1 & \\ \vdots & \ddots & \ddots & \ddots & \\ \frac{t^s}{s!} & \cdots & \frac{t^2}{2} & t & 1 \end{bmatrix},$$

where s is the size of J_i . We consider the following two cases.

(i) $\lambda_i = 0$. Then it must be the case that

$$\begin{bmatrix} 1 & & & \\ t & 1 & & \\ \frac{t^2}{2} & t & 1 & \\ \vdots & \ddots & \ddots & \ddots & \\ \frac{t^s}{s!} & \cdots & \frac{t^2}{2} & t & 1 \end{bmatrix} \vec{x}_i = \vec{y}_i.$$

Recall that entries of \vec{x}_i and \vec{y}_i are all algebraic numbers. Hence, as we assume that $\vec{x}_i \neq \mathbf{0}$, we have that $t \in \mathbb{A}$.

(ii) $\lambda_i \neq 0$. Then

$$e^{t\lambda_{i}}\begin{bmatrix}1\\t&1\\\frac{t^{2}}{2}&t&1\\\vdots&\ddots&\ddots&\ddots\\\frac{t^{s}}{t^{s}}&\cdots&\frac{t^{2}}{2}&t&1\end{bmatrix}\vec{x}_{i}=\vec{y}_{i}.$$

Recall that $\vec{x}_i \neq \mathbf{0}$. Clearly $e^{\lambda_i t} \in \mathbb{A}$. Note that Corollary 2 asserts that either $e^{\lambda_i t} \notin \mathbb{A}$ or $\lambda_i t \notin \mathbb{A}$. Hence $\lambda_i t \notin \mathbb{A}$ and thus $t \notin \mathbb{A}$. Furthermore, we claim that the size of the Jordan block (i.e., *s*) must be 1, because otherwise clearly $t \in \mathbb{A}$ which is a contradiction.

We distinguish the following two cases:

(a) All non-oblivious blocks are of eigenvalue 0. By case (i), $t \in \mathbb{A}$. Choose one of such blocks, we have an equation of the form

$$\begin{bmatrix} 1 & & & \\ t & 1 & & \\ \frac{t^2}{2} & t & 1 & \\ \vdots & \ddots & \ddots & \ddots & \\ \frac{t^s}{s!} & \cdots & \frac{t^2}{2} & t & 1 \end{bmatrix} \vec{u} = \vec{v}$$

and $\vec{u} \neq \mathbf{0}$. Let $i^* = \min\{i \mid \vec{u}_i \neq 0\}$ (such i^* must exist). Hence it must be the case that $t = \frac{\vec{v}_{i*}}{\vec{u}_{i*}}$.

(b) There exists at least one non-oblivious block whose eigenvalue is nonzero. Then by case (ii), $t \notin \mathbb{A}$. It follows that all non-oblivious blocks must have nonzero eigenvalues and all such Jordan blocks are of size 1.

That is, without loss of generality we have an equation of the form

$$\begin{bmatrix} e^{t\lambda_1} & & \\ & e^{t\lambda_2} & \\ & & \ddots & \\ & & & e^{t\lambda_\ell} \end{bmatrix} \begin{bmatrix} u_1 \\ \vdots \\ u_\ell \end{bmatrix} = \begin{bmatrix} v_1 \\ \vdots \\ v_\ell \end{bmatrix}$$
(2)

such that for each $1 \le i \le \ell$, $u_i \ne 0$ and $\lambda_i \ne 0$. Here ℓ is the number of non-oblivious blocks. Writing $z_i = \frac{v_i}{u_i}$, we have that, for $1 \le i \le \ell$,

 $e^{\lambda_i t} = z_i.$

We then claim that Eq. (2) has a solution $t \in \mathbb{R}_{\geq 0}$ iff 1. for any $1 \leq i, j \leq \ell, \frac{\lambda_i}{\lambda_j} \in \mathbb{Q}$ and $z_i^{\lambda_i} = z_j^{\lambda_j}$; and

2. there exist λ_i and z_i such that (2a) Either $\Re(z_i) > 0$, $\Re(\lambda_i) = 0$, and $\Im(z_i) = 0$; (2b) or $\Re(\lambda_i) = 0$ and $|z_i| = 1$.

The "if" part is obvious. To see the "only if" part, firstly it is easy to see that for $1 \le i, j \le \ell, z_i^{\lambda_j} = z_j^{\lambda_i}$. Namely, $\lambda_j \ln(z_i) - \lambda_i \ln(z_j) = 0$. By Corollary 1, $\ln(z_i)$ and $\ln(z_j)$ are linear independent over \mathbb{Q} . Hence $\frac{\lambda_i}{\lambda_j} = \frac{\ln(z_i)}{\ln(z_j)} \in \mathbb{Q}$.

Now let's focus on any $e^{\lambda_i t} = z_i$. Assume that $\lambda = a + bi$ and z = c + di, where $a, b, c, d \in \mathbb{R} \cap \mathbb{A}$. Recall that $\lambda_i \neq 0$. We consider the following cases:

- $a \neq 0$ and b = 0. Then *t* exists iff c > 0 and d = 0. This is equivalent to the case (2a).
- *a* = 0 and *b* ≠ 0. Then *t* exists iff $c^2 + d^2 = 1$. This is equivalent to the case (2b).
- $-a \neq 0$ and $b \neq 0$. It follows that

$$\begin{cases} e^{at} = \sqrt{c^2 + d^2} \in \mathbb{A} \\ e^{bti} = \frac{c + di}{\sqrt{c^2 + d^2}} \in \mathbb{A} \end{cases}$$

It follows that $(\sqrt{c^2 + d^2})^{i\frac{b}{a}} = \frac{c+di}{\sqrt{c^2+d^2}}$. By Theorem 1 we must have that $i\frac{b}{a} \in \mathbb{Q}$ which is a contradiction. Hence this case is actually vacuous.

Based on the above arguments, the algorithm is rather straightforward and we can analyse its complexity. By the result of [4], there is a polynomial-time algorithm to perform the Jordan decomposition for *A*, namely, one can compute the λ_i 's, \vec{x} and \vec{y} in polynomial time. Hence we can check for each oblivious block (λ_i , x_i , y_i) whether $y_i = \mathbf{0}$. If this is not the case, the algorithm is terminated and returns "No". Otherwise, we can determine either case (a) or case (b).

- In case (a), we can check whether $t = \frac{\bar{v}_{i^*}}{u_{i^*}}$ is the solution for all non-oblivious blocks. This can be done easily in polynomial time.
- In case (b), we can check whether conditions 1 and 2 are satisfied. To check $\frac{\lambda_i}{\lambda_j} \in \mathbb{Q}$, it suffices to check whether the degree of the minimal polynomial of $\frac{\lambda_i}{\lambda_j}$ is at most 1, which can be done in polynomial time. On top of this, checking $z_i^{\lambda_i} = z_j^{\lambda_j}$ amounts to checking

 $z_i^{r_{ij}} = z_j$ where $r_{ij} = \frac{\lambda_i}{\lambda_j}$, which can done in polynomial time as well. Furthermore it is trivial to check, for some λ_i and z_i whether (2a) or (2b) holds.

This completes the proof. \Box

4. Conclusion

In this paper, we have shown that the continuoustime orbit problem is decidable in polynomial-time. A very natural question is to consider the continuous-time orbit problem in higher dimensions. Combining the arguments of [5] and this paper, one can settle the case of dimension two or three; one can also link this problem to the continuous-time Skolem problem. However, solving this problem thoroughly seems to be difficult without a breakthrough (cf. [3]), notwithstanding some recent development for the discrete-time case [9]. It is also interesting to see whether the P upper-bound established here can be improved further, along the line of [1]. The main difficulty seems to lie in factoring polynomials which is needed for [ordan decomposition in [4]. To the best of our knowledge, the best upper-bound is P (by, e.g., the LLL algorithm) which obstructs further improvement inside P. We leave it an interesting open problem how to circumvent this difficulty.

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