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Two-dimensional effects in Fowler-Nordheim field emission from transition metal dichalcogenides

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Abstract. We report field emission from bilayer MoS_2 and monolayer WSe_2 synthesized by CVD on SiO_2/Si substrate. We show that the emitted current follows a Fowler-Nordheim model modified to account for the two-dimensional confinement of charge carriers. We derive the figures of merit of field emission and demonstrate that few-layer transition-metal dichalcogenides are suitable for field emission applications.

1. Introduction

Transition - metal dichalcogenides (TMDs), chemical formula MX_2 , have been largely investigated during the last few years. They present a layered structure, characterized by a honeycomb lattice where each transition metal M ($Mo, W, We, etc.$) is covalently bonded to two chalcogenide atoms X (S, Se, Te). TMDs have been intensively studied for the ease of fabrication and the tuneable and sizeable bandgap suitable for optoelectronic applications [1,2]. Like graphene, TMDs can be mechanically exfoliated from a bulk material and transferred on a substrate [3]; however, single-crystal and large-scale flakes, with controlled number of layers, are more easily produced by chemical vapor deposition (CVD) [4]. Those materials are considered good candidates for technological applications because they offer good control on of the energy bandgap through the number of layers [5], which increases passing from bulk to monolayer form, high on/off ratio ($> 10^7$) [6] and low subthreshold swing (below 70 mV/decade) [7,8], even though their mobility is typically only few tens $cm^2V^{-1}s^{-1}$ [9,10]. Their strong light absorption and emission open the route for high-responsivity photodetectors, high-intensity light emitters and efficient photovoltaic cells [11,12]. Furthermore, the intrinsic sharp edges due to the 2D nature and the low work function enable field emission applications. Anyway, unlike graphene and carbon nanotubes (CNTs) [13,14], TMDs have not been much explored as cold cathodes for applications in vacuum electronics, flat panel displays, electron microscopy, x-ray tubes [15–17], etc. Only few studies focussed on field emission



from MoS_2 and WSe_2 have been reported to date [18–20]. In this paper, we investigate the field emission (FE) behaviour of MoS_2 and WSe_2 bilayer and monolayer flakes. Specifically, we report FE measurements from their flat part. We show that the experimental data are well fitted by a Fowler-Nordheim (FN) model [21] modified to take into account the two-dimensional (2D) nature of the material. We derive the characteristic FE parameters such as the field enhancement factor β and the turn on field to show that TMDs are promising materials for field emission applications.

2. Materials and Methods

We grew the two analysed samples, MoS_2 and WSe_2 flakes, by chemical vapor deposition on a highly p-doped Si substrates covered by 300 nm of SiO_2 (silicon dioxide). The CVD growths, at $850\text{ }^\circ\text{C}$, were performed in a two-zone quartz tube furnace, with the substrate placed face down on the transition-metal precursor and the chalcogen carried by an Ar gas flow [22]. Flakes were randomly formed on the substrates. Figure 1(a) and Figure 1(b) show SEM images of two selected flakes, MoS_2 and WSe_2 , respectively. The field emission measurements were performed in the SEM chamber, at a pressure $< 10^{-6}\text{ mbar}$, exploiting the piezo-controlled tungsten tip (W -tip) movement for the accurate positioning of the anode in front of the emitting flake (cathode) at a close distance d , as displayed in insets of Figure 1(a) and (b).

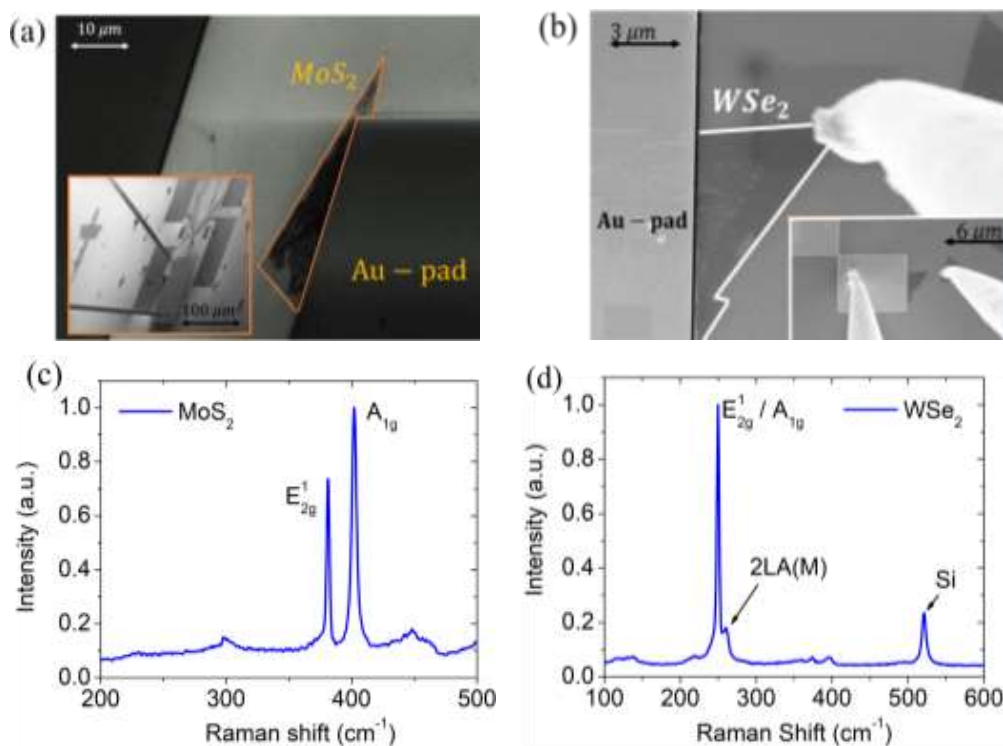


Figure 1. SEM images of the flakes used for field emission measurements and in insets the position of the tungsten tips: the lower one on the metal pad contacting the flake, the other placed in front of the flake at a close distance d , respectively for MoS_2 (a) and WSe_2 (b). Raman spectra of the MoS_2 (c) and WSe_2 (d).

The flakes were characterized by Raman spectroscopy with an unpolarized incident laser to identify the number of layers. The Raman spectrum reported in Figure 1(c), referring to MoS_2 , shows two peaks at $\sim 381\text{ cm}^{-1}$ and $\sim 401\text{ cm}^{-1}$ which are associated with the in-plane vibrations of sulfur atoms (E_{2g}^1) and with the out-of-plane vibrations of Mo and S atoms (A_{1g}), respectively [23]. The separation of $\sim 20\text{ cm}^{-1}$ indicates bilayer flakes. Similarly, Figure 1(d) shows the Raman spectrum of the WSe_2 flake. It exhibits two peaks around $\sim 250\text{ cm}^{-1}$ and $\sim 260\text{ cm}^{-1}$, corresponding the

first to a combination of the in-plane vibrations of W and Se atoms (E_{2g}) and out-of-plane vibrations of Se atoms (A_{1g}), and the other one to a second order resonant Raman mode ($2LA(M)$) due to LA phonons at the M point in the Brillouin zone [24]. The peak frequency positions are compatible with WSe_2 monolayer, which was further confirmed by photoluminescence measurements (not shown here) [18].

3. Results

Considering that the MoS_2 and WSe_2 flakes are naturally n-doped and that the sharp edges or nano-protrusions of the flat-part can originate a high electric field amplification [14,25], we expect easy extraction of electrons upon application of a voltage.

According to the standard FN theory [26] the field emission current usually results in a linear $\ln(I_{ds}/(V_{ds}^2))$ vs. $1/V_{ds}$ FN plot. However, for 2D materials, it has been proposed [21] that a modified FN model has to be applied to the vertical electron emission to take into account the reduced dimensionality, the energy spectrum, the non-conserving in-plane electron momentum, the finite-temperature and the space-charge-limited effects.

The modified FN model results in the following field emission current

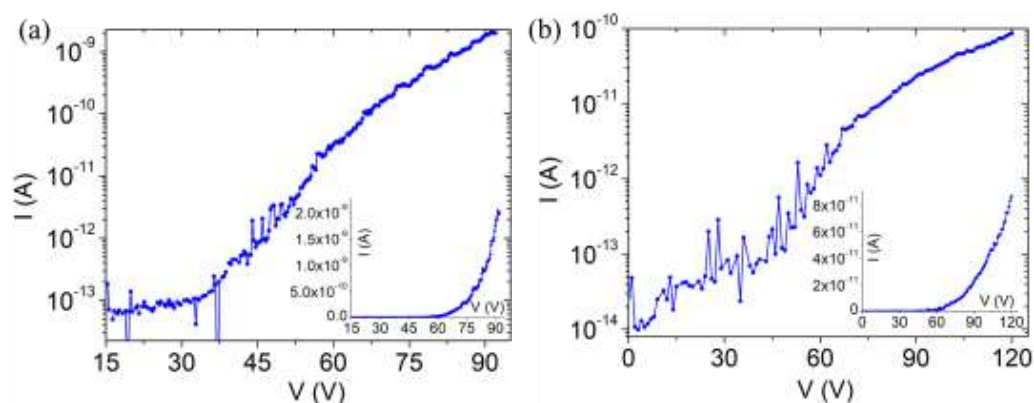
$$I_{FN}^{2D} = A_{FN}^{2D} \exp \left[-\frac{b\phi^3}{\beta E} \right] \quad (1)$$

where A_{FN}^{2D} and b are constants, S and ϕ represent the emitting surface and the electron affinity (in our case $\phi = 4.2$ eV for MoS_2 [27] and 3.9 eV for WSe_2 [28]), β is the so-called field enhancement factor, a typical figure of merit for the qualification of field emitter material. Finally, the electric field is $E = V/d$, where V represents the anode-cathode voltage and d their distance.

For both samples the FE current was measured with the (anode) W-tip positioned at a fixed distance ($d \sim 200$ nm for MoS_2 and $d \sim 400$ nm for WSe_2) from the surface (inner flat part of the flake). We disregarded the edges after having verified that the oxidation after air exposure and the higher concentration of unreacted precursor (MO_x and WO_x) had hampered the extraction of electrons from them. The current was monitored during the voltage sweep up to 120 V (a constraint to avoid the oxide breakdown).

Figure 2(a) and 2(b) show that, when V is below a critical value, the current fluctuates around the noise floor $< 10^{-13}$ A, and there is no measurable signal. For $V_{ds} > V_{th}$ an exponential increase (well visible in insets of Figure 2(a) and (b)) in the current is observed. The current-voltage characteristics show the typical fluctuations of emission current, related to desorption of adsorbates caused by Joule heating.

The fits of the experimental data at high-field regime of eq.1 are shown in Figure 2(c) and 2(d) compared with the 3D Fowler-Nordheim model [26], reported in inset. The accuracy of the linear fit is evaluated with the $\chi^2 \sim 0.99$, defining a good agreement in a wide range of voltage values, slightly better for WSe_2 respect to the classical 3D Fowler-Nordheim model.



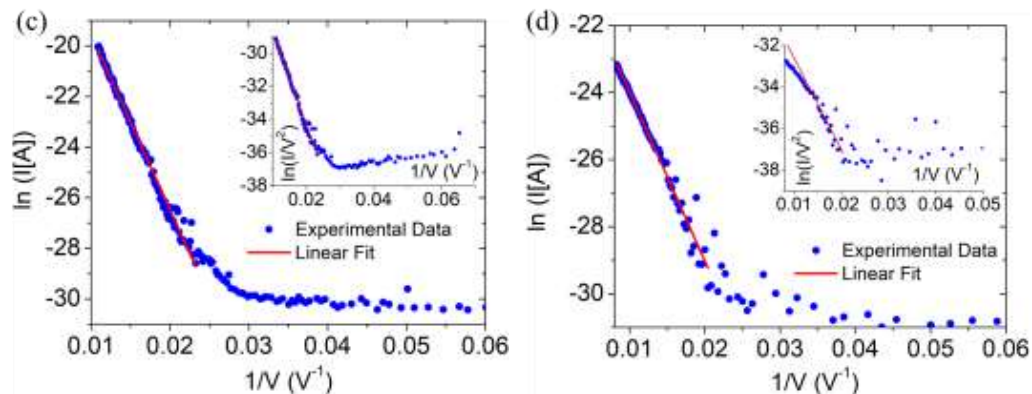


Figure 2. Field emission curves plotted in logarithmic scale as a function of anode-cathode voltage and linear scale (inset) for MoS_2 (a) and WSe_2 (b). Experimental data (blue circles) and model fit (red lines) of field emission measures for MoS_2 (c) and WSe_2 (d). The insets show the comparison of the experimental data fitted with standard Fowler-Nordheim model.

Indeed, this provides an indication in favour of the 2D modified FN model proposed in ref. [21], that should be further addressed and confirmed in future experiments.

From the slopes of the linear fits (eq.1) we derive the field enhancement factor $\beta = -b d \phi^{3/2}/m \sim 17$ and ~ 40 for MoS_2 and WSe_2 respectively.

The seemingly low values of the amplification factor (good emitters reach $\beta \sim 1000$ or more) are remarkable if we consider that emission happens from the inner part of the flake, without taking advantage from the 2D shape of the flake. Despite that, the turn-on field, defined as the field applied to extract a current of 1 pA has a value of $\sim 230 V/\mu m$ and $150 V/\mu m$ for MoS_2 and WSe_2 , respectively, which are rather low compared to the typical turn on field of several $kV/\mu m$ needed to extract electrons from flat surfaces. These results demonstrate the MoS_2 and WSe_2 suitability as field emitters.

4. Conclusion

In conclusion, we have fabricated MoS_2 bilayers and WSe_2 monolayer flakes and analysed their field emission properties. We have demonstrated that a modified field emission model, which considers the 2D nature of the flakes, provides a slightly better agreement with the experimental data compared to the original Fowler-Nordheim model, opening the route for further investigations. We have highlighted that TMDs are promising materials not only for electronics and optoelectronics but also for field emission applications.

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