

Back-gated Nb-doped MoS₂ junctionless field-effect-transistors

Gioele Mirabelli, Michael Schmidt, Brendan Sheehan, Karim Cherkaoui, Scott Monaghan, Ian Povey, Melissa McCarthy, Alan P. Bell, Roger Nagle, Felice Crupi, Paul K. Hurley, and Ray Duffy

Citation: *AIP Advances* **6**, 025323 (2016); doi: 10.1063/1.4943080

View online: <http://dx.doi.org/10.1063/1.4943080>

View Table of Contents: <http://aip.scitation.org/toc/adv/6/2>

Published by the [American Institute of Physics](#)

Back-gated Nb-doped MoS₂ junctionless field-effect-transistors

Gioele Mirabelli,^{1,2} Michael Schmidt,¹ Brendan Sheehan,¹ Karim Cherkaoui,¹ Scott Monaghan,¹ Ian Povey,¹ Melissa McCarthy,¹ Alan P. Bell,³ Roger Nagle,¹ Felice Crupi,² Paul K. Hurley,¹ and Ray Duffy¹

¹Tyndall National Institute, University College Cork, Cork, Ireland

²University of Calabria, Arcavacata di Rende, Cosenza, Italy

³CRANN and School of Chemistry, Trinity College Dublin, Dublin 2, Ireland

(Received 7 January 2016; accepted 12 February 2016; published online 26 February 2016)

Electrical measurements were carried out to measure the performance and evaluate the characteristics of MoS₂ flakes doped with Niobium (Nb). The flakes were obtained by mechanical exfoliation and transferred onto 85 nm thick SiO₂ oxide and a highly doped Si handle wafer. Ti/Au (5/45 nm) deposited on top of the flake allowed the realization of a back-gate structure, which was analyzed structurally through Scanning Electron Microscopy (SEM) and Transmission Electron Microscopy (TEM). To best of our knowledge this is the first cross-sectional TEM study of exfoliated Nb-doped MoS₂ flakes. In fact to date TEM of transition-metal-dichalcogenide flakes is extremely rare in the literature, considering the recent body of work. The devices were then electrically characterized by temperature dependent I_{ds} versus V_{ds} and I_{ds} versus V_{bg} curves. The temperature dependency of the device shows a semiconductor behavior and, the doping effect by Nb atoms introduces acceptors in the structure, with a p-type concentration $4.3 \times 10^{19} \text{ cm}^{-3}$ measured by Hall effect. The p-type doping is confirmed by all the electrical measurements, making the structure a junctionless transistor. In addition, other parameters regarding the contact resistance between the top metal and MoS₂ are extracted thanks to a simple Transfer Length Method (TLM) structure, showing a promising contact resistivity of $1.05 \times 10^{-7} \Omega/\text{cm}^2$ and a sheet resistance of $2.36 \times 10^2 \Omega/\text{sq}$. © 2016 Author(s). All article content, except where otherwise noted, is licensed under a Creative Commons Attribution (CC BY) license (<http://creativecommons.org/licenses/by/4.0/>). [<http://dx.doi.org/10.1063/1.4943080>]

INTRODUCTION

One of the most promising candidates as a substitution for silicon in modern electronics devices are Transition-Metal-Dichalcogenides (TMDs). These are semiconductors in the form of MX₂, where M is a transition metal (Mo, W, Hf), and X is a chalcogen (S, Se, Te). The weak Van-Der-Waals forces between each 2D monolayer of which they are formed, allow thin multilayers to be easily exfoliated by a scotch tape technique from their bulk form, as first was shown for graphene.¹ However one of the main limitations of graphene is its zero bandgap and resulting low on/off current ratio, which limits its usefulness for field-effect-transistor (FET) applications. On the contrary, a selection of TMDs show a finite direct bandgap between 1-2 eV, predicted by density-functional-theory (DFT) calculations,² which makes them suitable for logic devices. Their ultra-thin nature would allow a better electrostatic control for device applications with respect to the state-of-the-art devices, and better immunity to short-channel effects, earning the attention of the electronics and optoelectronics community.

Several groups have studied processes more compatible with large area coverage than mechanical exfoliation to obtain thin film or few monolayers of TMD materials. These include Chemical Vapor Deposition (CVD) used to grow MoS₂ by sulphurization of the metal,³⁻⁵ Molecular Beam Epitaxy (MBE) methods to grow HfSe₂,⁶ or MoSe₂,⁷ which can be potentially extended to other 2D

materials, and Atomic Layer Deposition (ALD) of MoS₂.⁸ Even though these kinds of processes will be needed for future applications in larger scale production, scotch tape based exfoliation is widely used because of its easiness, and fresh flakes can be obtained quickly. In addition, this method doesn't compromise device performance in terms of gate control and current on/off ratios. Often, back-gate devices are prepared with mechanically exfoliated flakes, in which the substrate is used as a control gate. Top-gated TMD films have produced an on/off ratio exceeding 10⁸ and mobility of almost 200 cm²/V.s for MoS₂ based FETs.⁹

In many cases, after flake preparation, an important problem that can affect and limit device performance is the contact resistance between the metal and the TMD. A high resistance can be due to residues at the TMD-metal interface, such as moisture or oxidation states on the surface. Moreover, X-ray photoelectron spectroscopy (XPS) results have shown impurities of carbon and oxygen on geological MoS₂.⁸ A cleaner TMD surface can be obtained by 15 minutes ultra-high-vacuum (UHV) annealing at 400 °C, which can burn off the unwanted contaminants. This is important because a low contact resistance can be obtained if a clean TMD-metal interface is achieved.¹⁰

An important role is also played by the types of metals used because Fermi level pinning may occur at the interface. This would result in a Schottky-barrier modulated current¹¹ instead of a more linear ohmic characteristic. It has been shown how low work function metals¹² such as Sc, may be used to improve both carrier injection and the contact resistance of the device, but other metals, such as Ti¹¹⁻¹³ and Ni^{12,14,15} have been used as well.

In this article, we report the electrical measurements of a back-gated MoS₂ flake doped with Nb, mechanical exfoliated from a Nb-doped MoS₂ crystal. Overall, device performance shows p-type characteristics, as expected for substitutional Nb doping,¹⁶⁻¹⁸ and promising results in terms of MoS₂-metal contact resistance, demonstrating that high Nb doping levels can improve MoS₂ device characteristics, much like in other doped semiconductors.¹⁴

EXPERIMENTAL

Using classic scotch tape techniques thin multilayer MoS₂ flakes were mechanically exfoliated from a bulk MoS₂ crystal. During formation of the synthetic MoS₂ crystal by a chemical vapor method, Nb was incorporated by the supplier. The presence of the Nb in the crystal was confirmed by time of flight secondary ion mass spectrometry (SIMS) analysis (not shown). The exfoliated flakes were gently pressed onto a surface composed of 85 nm of SiO₂ and a highly-doped Si handle wafer. The height of the flake measured using a Tencor Profiler was approximately 9.8 nm, thus the number of layers of the flake, considering each layer as 0.65 nm thick, is approximately 15. Ti/Au metal contact pads and electrodes were defined by electron-beam lithography, followed by metal evaporation and a lift off process. 15 kV beam exposures were performed with a Zeiss SUPRA SEM with a Raith Elephy Plus blanker. An adhesion layer of 5 nm of Ti and 45 nm of Au were deposited with e-beam evaporation. Before and after metal deposition the sample wasn't subjected to any process to remove surface contaminants.

For structural analysis, cross-section samples were obtained by using FEI's Dual Beam Helios Nanolab 600i system using Ga ion beam. Three layers of protective material were used, namely electron beam C, electron beam Pt, and ion beam C. Lamellas were thinned and polished at 30 kV 100 pA and 5 kV 47 pA, respectively. Cross-sectional Transmission Electron Microscopy (XTEM) imaging was carried out using a JEOL 2100 HRTEM operated at 200 kV in Bright Field mode using a Gatan Double Tilt holder. For electrical characterization, the Agilent B-1500 parameter analyser was used. Hall effect measurements were obtained using a LakeShore HMS 8404 Hall Measurement System with both DC and AC magnetic field capability.¹⁹ We carried out Hall effect measurements on an approximately 190 μm thick flake exfoliated from the same bulk Nb-doped MoS₂ crystal. Both DC and AC magnetic field methods were used with field reversal strength of ±1.7 T for the DC technique and +1.2 T_{rms} for the AS method using a lock-in amplifier. Geometry and current reversal methods were employed to remove error factors.

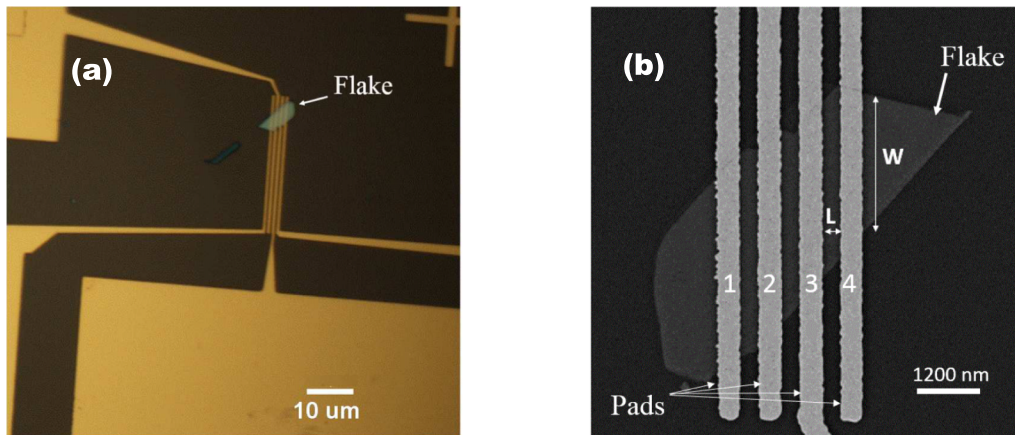


FIG. 1. (a) Optical image of the MoS₂ flake. (b) SEM image of the contacted device, with the contact tracks numbered 1 to 4.

RESULTS AND DISCUSSION

The Nb-doped MoS₂ device is shown in Figure 1(a) and 1(b). The image in Figure 1(a) was taken with an optical microscope. The large areas of metal are the contact pads, and the lines are the metal electrodes leading to the MoS₂ flake, which is labelled. In addition, a SEM image of the device is shown in Fig. 1(b), to be sure that metal deposition was successful, without any cracks or discontinuities of the metal tracks along the flake. For the sake of clarity, the pads are numbered 1 to 4 from left to right.

Representative TEM images are reported in order to have a deeper understanding of the material characteristics of the flake. Figure 2 shows a wide view of the device, where it is possible to notice the Ti/Au pads on top of the thin flake, and then the SiO₂ and the Si substrate. Interestingly, the flake is characterized by undulations along the surface of the substrate, which are most prominent around the edges of the Au/Ti metal. It is possible the weak Van der Waals forces between layers of MoS₂ allow for some bending of the thin MoS₂, which causes a bending in the MoS₂ film primarily around the edges of the Au/Ti metal electrodes, most probably occurring due to the lift off process.

Figure 3 shows four different representative TEM images with a higher magnification, taken at different locations along the flake. In Fig. 3(a) it is possible to clearly see the layered structure of the flake, composed by 16 layers, confirming the height of almost 10 nm found with the Tencor Profiler. Figure 3(b) shows another particular area of the flake, where the flake bends, even if it is not near the metal pads.

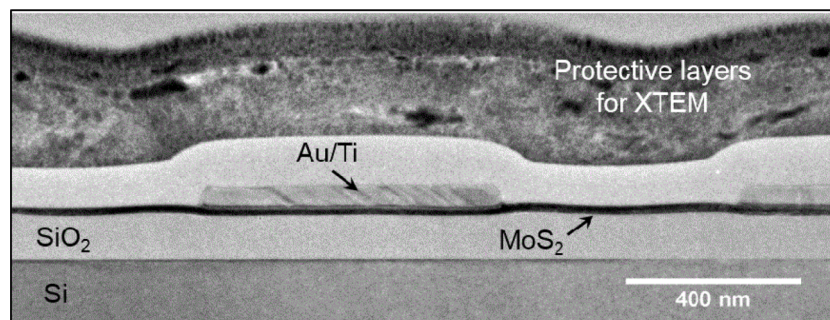


FIG. 2. Representative TEM image of the overall structure of the flake. It is possible to note the undulations of the 10 nm MoS₂ flake along the surface of the substrate.

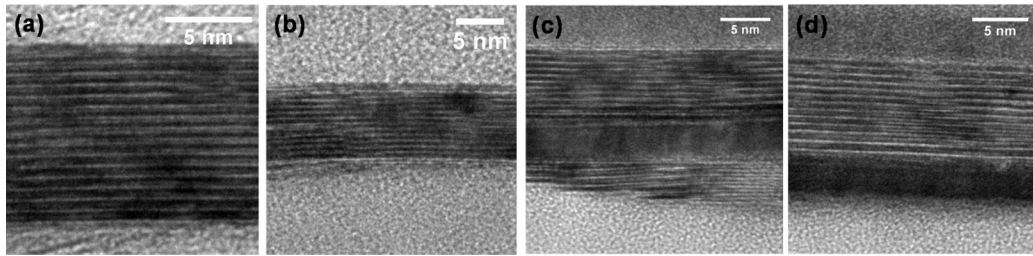


FIG. 3. Representative higher-resolution TEM images showing (a) the layered structure of the MoS₂, (b) bending of the MoS₂. Indicated by the white arrows in (c) and (d) uniformly grey regions which are crystalline but do not have the classic layered MoS₂ structure. This may be a side-effect of the mechanical exfoliation.

Figure 3(c) displays that the layered MoS₂ structure and the height are not constant along all the flake. There are regions within the MoS₂ that are crystalline, but not perfectly layered like that in Fig. 3(a). This is the uniformly grey region. This may be a side-effect of the mechanical exfoliation process, as cross sections of the bulk crystal (not shown) do not show this. This may be due to a different orientation of the MoS₂ layers in that area, so the flake may still be crystalline in that region. Moreover the height is clearly not perfectly constant, but some missing layers can be due to the stress induced on the flake through the mechanical exfoliation method. Figure 3(d) reports another region of thin MoS₂ which is crystalline but does not have the classic layered MoS₂ structure. This uniformly grey region might have a different orientation to the rest of the layered MoS₂.

Electrical measurements were carried out to characterise the electrical conductivity of the device. First of all, I_{ds} versus V_{ds} curves were measured at 25, 40 and 55 °C, considering all pad combinations to ensure that there were no open-circuits of any pad or short-circuits between pads. Figure 4(a) shows the IV characteristics between the first and the second pad, while Figure 4(b) considers the first and the last one. This effectively changes the channel length of the back-gated device. The current scales with gate length as expected. An increase in temperature corresponds to a higher current, consistent with a semiconductor behavior. The current is not linear with respect to source-drain voltage, so it is not an ohmic behavior, but it seems to be modulated by Schottky barrier between the flake and the metal, as was reported in other publications.^{11,12}

In order to further demonstrate the p-type behavior of the device, I_{ds} versus V_{ds} curves were extracted for different values of V_{bg} , considering a constant drain-source voltage at 1 V. As it is clear from the plot in Fig. 5, a more negative back-gate voltage results in a higher drain current, as it would be expected from a p-type doped semiconductor. Calculating the maximum depletion region

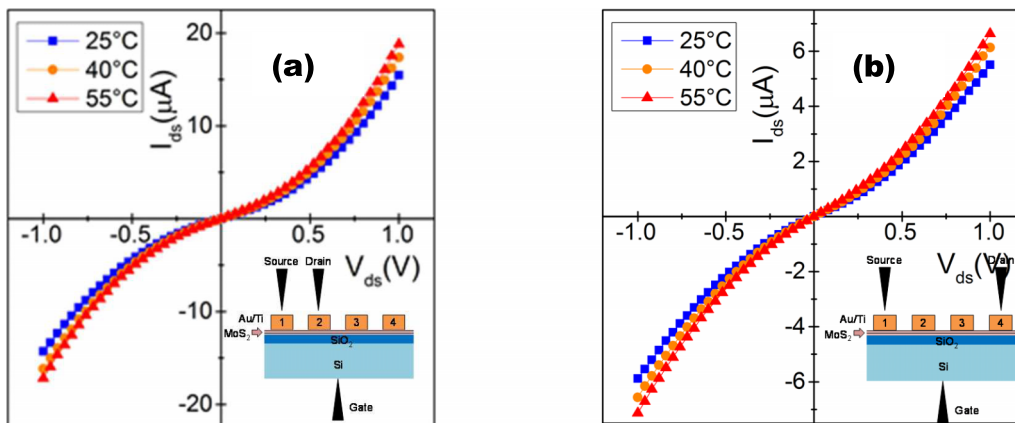


FIG. 4. IV characteristics between (a) first and second pad and between (b) first and fourth pad, as labelled in Fig. 1, for no voltage applied to the back gate. The insets show the contact scheme.

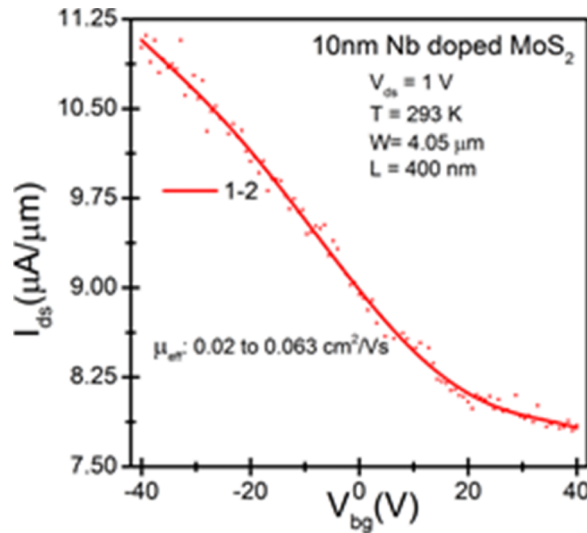


FIG. 5. I_{ds} - V_{bg} curve between the first and second electrodes as labelled in Fig. 1.

in the MoS₂ film based on the doping concentration of $4.3 \times 10^{19} \text{ cm}^{-3}$ and an out of plane dielectric constant of 10.5 for MoS₂, yields a value of 4.7 nm^{18} . This value is lower than the thickness of the MoS₂ flake ($\sim 10 \text{ nm}$) and as a consequence the device cannot be fully turned off compromising its I_{on}/I_{off} ratio.

Further proof of the p-type doping of the flake are the findings of our DC and AC Hall effect measurements. From other publications, the carrier concentration reported for this type of flake was 3×10^{19} – $3 \times 10^{20} \text{ cm}^{-3}$.^{17,18} Our Hall effect data confirm this, as the p-type doping concentration found is $4.3 \times 10^{19} \text{ cm}^{-3}$ with a Hall mobility of approximately $7.5 \text{ cm}^2/\text{V}\cdot\text{s}$.

Considering the overall structure of the device and such a high doping, it can be categorized as a “junctionless transistor”, whose architecture can achieve low leakage currents and very good response against short channel effects.^{20,21} The importance of this is not only related to the potential of the electrical characteristics, but also to the easier process and the higher device density that could be advantageous with respect to the modern silicon-based devices in production.

Another parameter extracted from the electrical measurements is the field effect mobility. The field effect mobility was determined from the derivative of the transfer characteristic, and the results for two separate pad combinations are shown in Figure 6, with the peak mobility values for all pad combinations included in Table I.

The field effect mobility is lower than the Hall effect mobility probably due to a number of factors. The Hall effect measurement was performed on a large crystal, while the device characteristics were extracted from a 10 nm flake, with all its imperfections (see XTEM in Fig. 3) and many surfaces. These imperfections are likely to degrade the carrier mobilities. The field effect mobility values were found to be in the range 0.02 and $0.2 \text{ cm}^2/\text{V}\cdot\text{s}$.

Previous studies have shown the thickness dependence of the mobility^{12,22,23} and the improvement after high-k dielectric passivation^{24,25} or vacuum annealing for removal of atmospheric adsorbates.²² In terms of optimum TMD thickness, Das *et al.*¹² reported a systematic study on MoS₂ flakes, back-gated n-type FET devices with Sc contact metallization. Based on drive current and mobility, MoS₂ thicknesses of ~ 6 – 12 nm , which comprises the height of our flake, appears to be optimum. Also, a thickness study showed a peak of mobility of approximately $54 \text{ cm}^2/\text{V}\cdot\text{s}$ for 8 layers. In that work electron mobility was $184 \text{ cm}^2/\text{V}\cdot\text{s}$ without high-k overlayers, but with Al₂O₃ high-k dielectric mobility increased to $700 \text{ cm}^2/\text{V}\cdot\text{s}$. Very recently Lee *et al.* reported field-effect mobility of $4.5 \text{ cm}^2/\text{V}\cdot\text{s}$ measured in MoS₂/SiO₂/Si back-gated transistors, based on a clean transfer of large-area MoS₂ films using deionized water.²⁶ Considering our process flow, there were no high-k dielectric passivation or vacuum annealing, so these could be some of the reasons for the low values of field effect mobility.

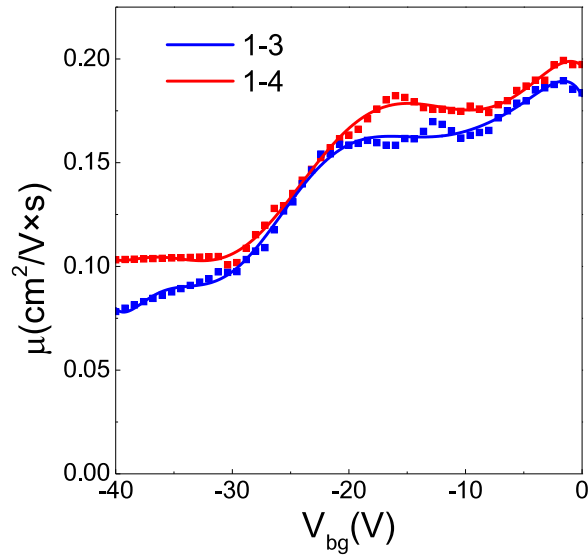


FIG. 6. Derivation of field effect mobility values versus V_{bg} for two different pad combinations: 1-3 and 1-4.

TABLE I. Voltages where peak $g_{m,max}$ values are recorded and the corresponding field-effect mobility values at those voltages for all pad combinations.

Pad combinations	Voltage (V)	μ ($\text{cm}^2/\text{V}\cdot\text{s}$)
1-2	-0.8	0.06
1-4	-1.6	0.20
1-3	-1.6	0.19
2-3	-3.2	0.04
2-4	-1.6	0.08
3-4	-2.4	0.02

Even if the current is not perfectly ohmic as said before, still the promising electrical response of the flake is worthy to be further studied. Then, from the SEM image reported in Fig. 1, one can notice how the pad configuration looks like a simple Transfer Length Method (TLM) structure, where the length and the width are reported as “L” and “W”, respectively. If the TLM method could be applied more parameters can be extracted from the device, allowing a deeper understanding of its characteristics.

Each track is 400 nm wide and the spacing between each track is 400 nm. In order to apply the TLM method three lengths can be considered: between first and second pad (spacing=400 nm), between first and third pad (spacing=1200 nm), and first and fourth pad (spacing=1600 nm). In addition, the width “W” is approximately 4.05 μm .

We are aware that TLM layouts are usually more optimized than here. Indeed, in order to create a TLM there should be more tracks, with varied spacing. In Fig. 7 it is possible to see that the total resistance plotted against the contact separations considered before is linear at different temperatures. With the TLM, from the data in Fig. 7, it is possible to extract the contact resistance and the transfer length of the AuTi/MoS₂ contacts. Considering these values it is possible to evaluate sheet resistance (R_{SH}), specific contact resistivity (ρ_C), and contact resistance (R_C) expressed in Ω and $\Omega \cdot \text{mm}$. All the results are shown in Table II. The transfer length values are between 207 and 211 nm, which is almost half the width of the contact tracks. The contact resistance values are ranged between 12.34 and 12.45 $\text{k}\Omega$, which is in agreement with other publications,^{27,28} despite the fact that we did not perform optimized process recipes to minimize contact resistance. With regards to R_C expressed as $\Omega \cdot \text{mm}$ our value reported at 25 °C is almost 6.3 times higher than previously

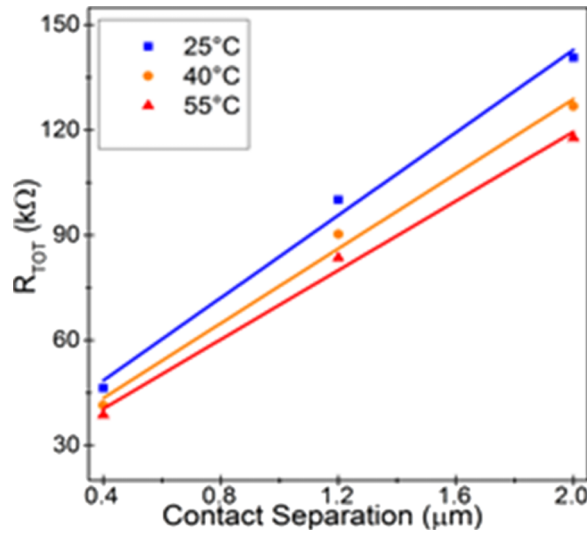


FIG. 7. Total resistance versus contact separation at different temperatures.

TABLE II. Results obtained from the MoS₂ device evaluated as a TLM structure.

Temp. (°C) :	25	40	55
1/kT	38.9	37.0	35.3
R _{SH} (Ω/sq)	2.36×10^2	2.13×10^2	1.97×10^2
L _T (nm)	211.1	207.4	209.2
ρ _C (Ω*cm ²)	1.05×10^{-7}	9.18×10^{-8}	8.65×10^{-8}
R _C (kΩ)	12.4	11.0	10.3
R _C (Ω*mm)	4.98	4.42	4.13

reported,²² but it can be related to the different type of flake studied and to different process applied to the flake before testing it.

A low specific contact resistivity of $5 \times 10^{-7} \Omega \cdot \text{cm}^2$ for MoS₂ was reported by English *et al.*¹⁰ In that work, the low contact resistivity was reached thanks to the low pressure adopted during the metal deposition by electron-beam lithography, in order to guarantee a cleaner interface between the MoS₂ and the metal. In terms of ρ_C, another relevant result of $2.2 \times 10^{-7} \Omega \cdot \text{cm}^2$ was reached by Kang *et al.*²⁹ through the use of Mo as metal contact, in order to form lower Schottky barrier and a high quality interface. Note that in these cases, MoS₂ flakes were mechanically exfoliated and no intentional or unintentional doping was reported. Our low value of $1.05 \times 10^{-7} \Omega \cdot \text{cm}^2$ indicates that high doping of MoS₂ is a very effective way to drive down ρ_C. Future work should combine high-doping levels with other processes such as pre-metal deposition cleaning and surface functionalisation techniques, along with dedicated TLM structures for accurate parameter extraction of ρ_C below $10^{-7} \Omega \cdot \text{cm}^2$.

CONCLUSIONS

In summary, highly doped MoS₂ flakes were mechanically exfoliated and contacted. The p-type behavior of the flake, expected since it is doped with Nb atoms, was confirmed by electrical device measurements and by Hall-effect measurements that indicated a high carrier concentration of $4.3 \times 10^{19} \text{ cm}^{-3}$ and Hall mobility of $7.5 \text{ cm}^2/\text{V.s}$. The low contact resistance (12.4 kΩ and 4.98 Ω. mm) and specific contact resistivity ($1.05 \times 10^{-7} \Omega \cdot \text{cm}^2$) data reported, even if extracted with a simple TLM structure, can be related to the high p-type doping of the flake, which, together with other processes widely used to improve metal contacts and obtain clean TMDs surfaces, can

be helpful in minimizing the specific contact resistance of the MoS₂/metal interface. Moreover, the high doping can be the key to the introduction of TMD-based junctionless transistors, in which the favorable characteristics of TMDs and the uncomplicated process required by this kind of device architecture may be combined.

ACKNOWLEDGEMENTS

We acknowledge the support of Science Foundation Ireland through the US-Ireland R&D Partnership Programme, “*Understanding the Nature of Interfaces in Two Dimensional Electronics (UNITE)*”, Grant Number SFI/13/US/I2862. The research was supported in part by the Higher Education Authority Programme for Research in Third Level Institutions in Ireland under Grant Agreement no. HEA PRTL15.

- ¹ K. S. Novoselov, A. K. Geim, S. V. Morozov, D. Jiang, Y. Zhang, S. V. Dubonos, I. V. Grigorieva, and A. A. Firsov, “Electric Field Effect in Atomically Thin Carbon Films,” *Science* **306**, 666–669 (2004).
- ² O. V. Yazyev and A. Kis, “MoS₂ and semiconductors in the flatland,” *Materials Today* **18**(Number 1 January/February), (2015).
- ³ J. Park, N. Choudhary, J. Smith, G. Lee, M. Kim, and W. Choi, “Thickness modulated MoS₂ grown by chemical vapor deposition for transparent and flexible electronic devices,” *Appl. Phys. Lett.* **106**, 012104 (2015).
- ⁴ H. Wang, L. Yu, Y. Lee, W. Fang, A. Hsu, P. Herring, M. Chin, M. Dubey, L. Li, J. Kong, and T. Palacios, “Large-scale 2D electronics based on single-layer MoS₂ grown by chemical vapor deposition,” in *Proceedings of the IEEE Electron Devices Meeting (IEDM), 2012* (2012), pp. 4.6.1–4.6.4.
- ⁵ K. Kang, S. Xie, L. Huang, Y. Han, P. Y. Huang, K. F. Mak, C.-J. Kim, D. Muller, and J. Park, “High-mobility three-atom-thick semiconducting films with wafer-scale homogeneity,” *Nature* **520**, 656 (2015).
- ⁶ R. Yue, A. T. Barton, H. Zhu, A. Azcatl, L. F. Pena, J. Wang, X Peng, N. Lu, L. Cheng, R. Addou, S. McDonnell, L. Colombo, J. W. P. Hsu, J. Kim, M. J. Kim, R. M. Wallace, and C. L. Hinkle, “HfSe₂ Thin Films: 2D Transition Metal Dichalcogenides Grown by Molecular Beam Epitaxy,” *ACS Nano* **9**(1), 474 (2015).
- ⁷ S. Vishwanath, S. Rouvimov, T. Orlova, X. Liu, J. K. Furdyna, D. Jena, and X. G. H. Xing, “Atomic Structure of Thin MoSe₂ Films Grown by Molecular Beam Epitaxy,” in *Proceedings of Microscopy and Microanalysis 2014* (2014).
- ⁸ L. K. Tan, B. Liu, J. H. Teng, S. Guo, H. Y. Low, and K. P. Loh, “Atomic Layer deposition of a MoS₂ film,” *Nanoscale* **6**, 10584 (2014).
- ⁹ B. Radisavljevic, A. Radenovic, J. Brivio, V. Giacometti, and A. Kis, “Single-layer MoS₂ transistors,” *Nature Nanotechnology* **6**, 147 (2011).
- ¹⁰ C. D. English, G. Shine, V. E. Dorgan, K. C. Saraswat, and E. Pop, “Improving contact resistance in MoS₂ field effect transistors,” in *Proceedings of Device Research Conference (DRC) 2014* (2014), pp. 193–194.
- ¹¹ H. Liu, M. Si, Y. Deng, A. T. Neal, Y. Du, S. Najmaei, P. M. Ajayan, J. Lou, and P. D. Ye, “Switching Mechanism in Single-Layer Molybdenum Disulfide Transistors: An Insight into Current Flow across Schottky Barriers,” *ACS Nano* **8**(1), 1031–1038 (2014).
- ¹² S. Das, H. -Y. Chen, A. V. Penumatcha, and J. Appenzeller, “High Performance Multilayer MoS₂ Transistors with Scandium Contacts,” *Nano Lett.* **13**(1), 100 (2013).
- ¹³ W. Liu, J. Kang, D. Sarkar, Y. Khatami, D. Jena, and K. Banerjee, “Role of metal contacts in designing high performance monolayer n-type WSe₂ field effect transistor,” *Nano Lett.* **13**(5), 1983 (2013).
- ¹⁴ L. Yang, K. Majumdar, H. Liu, Y. Du, H. Wu, M. Hatzistergos, P. Y. Hung, R. Tieckelmann, W. Tsai, C. Hobbs, and P. D. Ye, “Chloride molecular doping technique on 2D materials: WS₂ and MoS₂,” *Nano Lett.* **14**(11), 6275 (2014).
- ¹⁵ H. Liu, A. T. Neal, Y. Du, and P. D. Ye, “Fundamentals in MoS₂ transistors: dielectric, scaling and metal contacts,” *ACS Transactions* **58**(7), 203 (2013).
- ¹⁶ K. Dolui, I. Rungger, C. D. Pemmaraju, and S. Sanvito, “Possible doping strategies for MoS₂ monolayers: An ab initio study,” *Physical Review B* **88**, 075420 (2013).
- ¹⁷ M. R. Laskar, D. N. Nath, L. Ma, E. W. Lee, C. H. Lee, T. Kent, Z. Yang, R. Mishra, M. A. Roldan, J. -C. Idrobo, S. T. Pantelides, S. J. Pennycook, R. C. Myers, Y. Wu, and S. Rajan, “P-type doping of thin films using Nb,” *Appl. Phys. Lett.* **104**, 092104 (2014).
- ¹⁸ S. Das, M. Demarteau, and A. Roelofs, “Nb-doped single crystalline MoS₂ field effect transistor,” *Appl. Phys. Lett.* **106**, 173506 (2015).
- ¹⁹ J. Lindemuth and S. I. Mizuta, “Hall measurements on low mobility materials and high resistivity materials,” *Proc. Of SPIE*. **8110**, 811001 (2011).
- ²⁰ J. P. Colinge, A. Kranti, R. Yan, C. W. Lee, I. Ferain, R. Yu, N. D. Akhavan, and P. Razavi, “Junctionless Nanowire Transistor (JNT): Properties and design guidelines,” *Solid-State Electronics* **65–66**, 33–37 (2011).
- ²¹ J. P. Colinge, C. W. Lee, A. Afzalian, N. D. Akhavan, R. Yan, I. Ferain *et al.*, “Nanowire transistors without junctions,” *Nat Nanotechnol.* **5**, 225 (2010).
- ²² W. Liu, J. Kang, W. Cao, S. D. Khatami, Y. D. Jena, and K. Banerjee, “High performance few layer MoS₂ field-effect-transistor with record low contact resistance,” in *Proceedings of the IEEE Electron Devices Meeting (IEDM) 2013* (2013), pp. 19.4.1–19.4.4.
- ²³ D. Lembke, A. Allain, and A. Kis, “Thickness-dependent mobility in two-dimensional MoS₂ transistors,” *Nanoscale*. **14**, 6255 (2015).
- ²⁴ H. Liu, J. Gu, and P. D. Ye, “MoS₂ nanoribbon transistors: transition from depletion mode to enhancement mode by channel-width trimming,” *IEEE El. Dev. Lett.* **33**, 1273 (2012).

- ²⁵ H. Liu and P. D. Ye, "MoS₂ dual-gate MOSFET with atomic layer deposition Al₂O₃ as Top-Gate Dielectric," *IEEE El. Dev. Lett* **33**, 546 (2012).
- ²⁶ C. H. Lee, W. McCulloch, E. W. Lee II, Lu Ma, S. Krishnamoorthy, J. Hwang, Y. Wu, and S. Rajan, "Transferred large area single crystal MoS₂ field effect transistors," *Appl. Phys. Lett.* **107**, 193503 (2014).
- ²⁷ H.-J. Kwon, S. Kim, J. Jang, and C. P. Grigoropoulos, "Evaluation of pulsed laser annealing for flexible multilayer MoS₂ transistor," *Appl. Phys. Lett.* **106**, 113111 (2015).
- ²⁸ H. -J. Kwon, J. Jang, S. Kim, V. Subramanian, and C. P. Grigoropoulos, "Electrical Characteristics of MoS₂ transistors at real operating temperatures with different ambient conditions," *Appl. Phys. Lett.* **105**, 152105 (2014).
- ²⁹ J. Kang, W. Liu, and K. Banerjee, "High-performance MoS₂ transistors with low-resistance molybdenum contacts," *Appl. Phys. Lett.* **104**, 093106 (2014).