

# *Observational evidence of S-web source of the slow solar wind*

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# Observational Evidence of S-web Source of the Slow Solar Wind

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## Abstract

From 2022 March 18 to 21, NOAA Active Region (AR) 12967 was tracked simultaneously by Solar Orbiter at 0.35 au and Hinode/EIS at Earth. During this period, strong blueshifted plasma upflows were observed along a thin, dark corridor of open magnetic field originating at the AR's leading polarity and continuing toward the southern extension of the northern polar coronal hole. A potential field source surface model shows large lateral expansion of the open magnetic field along the corridor. Squashing factor  $Q$ -maps of the large-scale topology further confirm super-radial expansion in support of the S-web theory for the slow wind. The thin corridor of upflows is identified as the source region of a slow solar wind stream characterized by  $\sim 300 \text{ km s}^{-1}$  velocities, low proton temperatures of  $\sim 5 \text{ eV}$ , extremely high density  $> 100 \text{ cm}^{-3}$ , and a short interval of moderate Alfvénicity accompanied by switchback events. When the connectivity changes from the corridor to the eastern side of the AR, the in situ plasma parameters of the slow solar wind indicate a distinctly different source region. These observations provide strong evidence that the narrow open-field corridors, forming part of the S-web, produce some extreme properties in their associated solar wind streams.

*Unified Astronomy Thesaurus concepts:* [Slow solar wind \(1873\)](#)

*Supporting material:* animations

## 1. Introduction

The solar corona continuously expands outward as hot, magnetized streams of plasma that fill the heliosphere. These streams are generally classified as the fast and slow solar wind (SW) based on a velocity threshold of  $500 \text{ km s}^{-1}$  (e.g., Temmer 2021). The fast SW is quasi-uniform, exhibiting smooth, slow changes in plasma parameters, whereas the slow SW tends to be highly variable, especially in plasma composition and density. The contrasting characteristics of the two SW streams reflect their differing source regions. Fast SW originates from coronal holes (CHs); the plasma travels

unadulterated along open fields from the chromosphere/corona into the heliosphere (e.g., Zirker 1977; Cranmer 2009). However, the sources of the slow SW are found in the closed-field regions on the Sun, e.g., active regions (ARs), CH boundaries (CHBs), and streamers (e.g., Geiss et al. 1995; McComas et al. 1998; Brooks et al. 2015; Abbo et al. 2016). The highly variable nature of the plasma suggests that the slow SW is released onto open-field lines by stochastic processes such as magnetic reconnection. How plasma confined to the closed field of the Sun contributes to the SW remains an open question. Results from recent missions, such as ESA's Solar Orbiter (SO; Müller et al. 2020; García Marirrodriga et al. 2021) and NASA's Parker Solar Probe (PSP; Fox et al. 2016), which are parts of the Heliophysics System Observatory, are providing new insights into the physical processes and environment where the slow SW is formed and accelerated.



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The S-web model proposed by Antiochos et al. (2011) provides a theoretical framework for the origin of the slow SW. According to this model, the corona is filled with a network of narrow open-field corridors bounded by a web of separatrix surfaces or quasi-separatrix layers (QSLs; Démoulin et al. 1996). These boundaries are defined by discontinuous/dramatic changes of the field-line mapping from the photosphere to the source surface. The narrow, sometimes infinitesimal, corridors are topologically robust features that link CHs on the Sun (Antiochos et al. 2007, 2011). The topology of the S-web consists of giant arcs of open–closed magnetic flux boundary layers extending tens of degrees in longitude (Higginson et al. 2017b; Scott et al. 2018).

Interchange reconnection is proposed as the mechanism responsible for the release of the slow SW at the S-web arcs. Three-dimensional magnetohydrodynamic (MHD) simulations of Higginson et al. (2017a) demonstrated that photospheric convective motions on supergranular scales drive the magnetic field differently within open/closed-field regions. This results in ongoing dynamics along the S-web. Multiple interchange reconnection events are induced there, and plasma is exchanged between the open and closed flux domains. Interchange reconnection is also induced in pseudo-streamer configurations at their null points and separators (Aslanyan et al. 2021). Recently, Chitta et al. (2023) presented extreme-ultraviolet (EUV) and white-light observations of the highly structured mid-corona that show indications of the interchange reconnection dynamics consistent with the MHD simulations of the S-web.

In this work, we provide observational evidence of an S-web source region and its associated slow SW stream. Plasma parameters are determined within the narrow corridor, which is an open-field source region, from spectroscopic observations acquired by Hinode’s EUV Imaging Spectrometer (EIS; Culhane et al. 2007). The associated slow SW stream is characterized from in situ measurements made by SO’s Magnetometer (MAG; Horbury et al. 2020) and Solar Wind Analyser Proton and Alpha particle Sensor (SWA-PAS; Owen et al. 2020). SO’s Extreme Ultraviolet Imager (EUI)/Full Sun Imager (FSI; Rochus et al. 2020), coupled with magnetic field observations and modeling, are employed to connect the SW to its probable source region.

## 2. Solar Orbiter Target Region

In 2022 March, SO was at perihelion when the science observation phase of the mission commenced. During this time, a number of the Solar Orbiter Observing Plans (SOOPs) selected by the high-level Science Working Team (SWT) were run. SO’s complex payload, comprising six remote sensing and four in situ instruments, requires coordinated planning in order to ensure that the mission’s scientific objectives are met. SOOPs are the primary mechanism for planning and coordinating SO observations among its own instruments, as well as with other ground-based and space-borne assets. SOOPs are especially important for observations acquired during the three 10-day remote sensing windows scheduled for each orbit (for a complete account of SO’s planning, see Zouganelis et al. 2020). The L\_SMALL\_HRES\_HCAD\_Slow-Wind-Connection SOOP was selected to be among the first to be run. Its stated aim is to identify the sources of the SW, and it was designed to address in part one of the four top-level science questions of the mission: “What drives the solar wind, and

where does the coronal magnetic field originate?” (Müller et al. 2020).

The ideal target source region of the SOOP is at the boundary/interface of open and closed magnetic field either at the edges of ARs close to the low-latitude open field or at CHBs. The open field can provide pathways for plasma to travel from the source region to SO, where it is detected in situ. Remote sensing instruments are pointed at the target with the expectation that the spacecraft crosses the boundary/interface a number of days later, depending on a number of factors, including SO’s distance from the Sun and the SW velocity. Connectivity of the slow wind stream is then confirmed by matching plasma parameters measured at the Sun and at the SO spacecraft in combination with magnetic field modeling. NOAA AR 12967 provided the perfect slow wind source region target for SO when the spacecraft was  $\sim 0.35$  au from the Sun. Utilizing spectroscopic measurements, magnetic topology, and in situ observations, NOAA AR 12967 is identified as being the likely source region magnetically connected to the slow wind stream detected in situ (see Section 5). A complete account of the target selection, connectivity, and science observations taken when SO was at/near perihelion is provided in Yardley et al. (2023).

## 3. Remote Sensing Observations

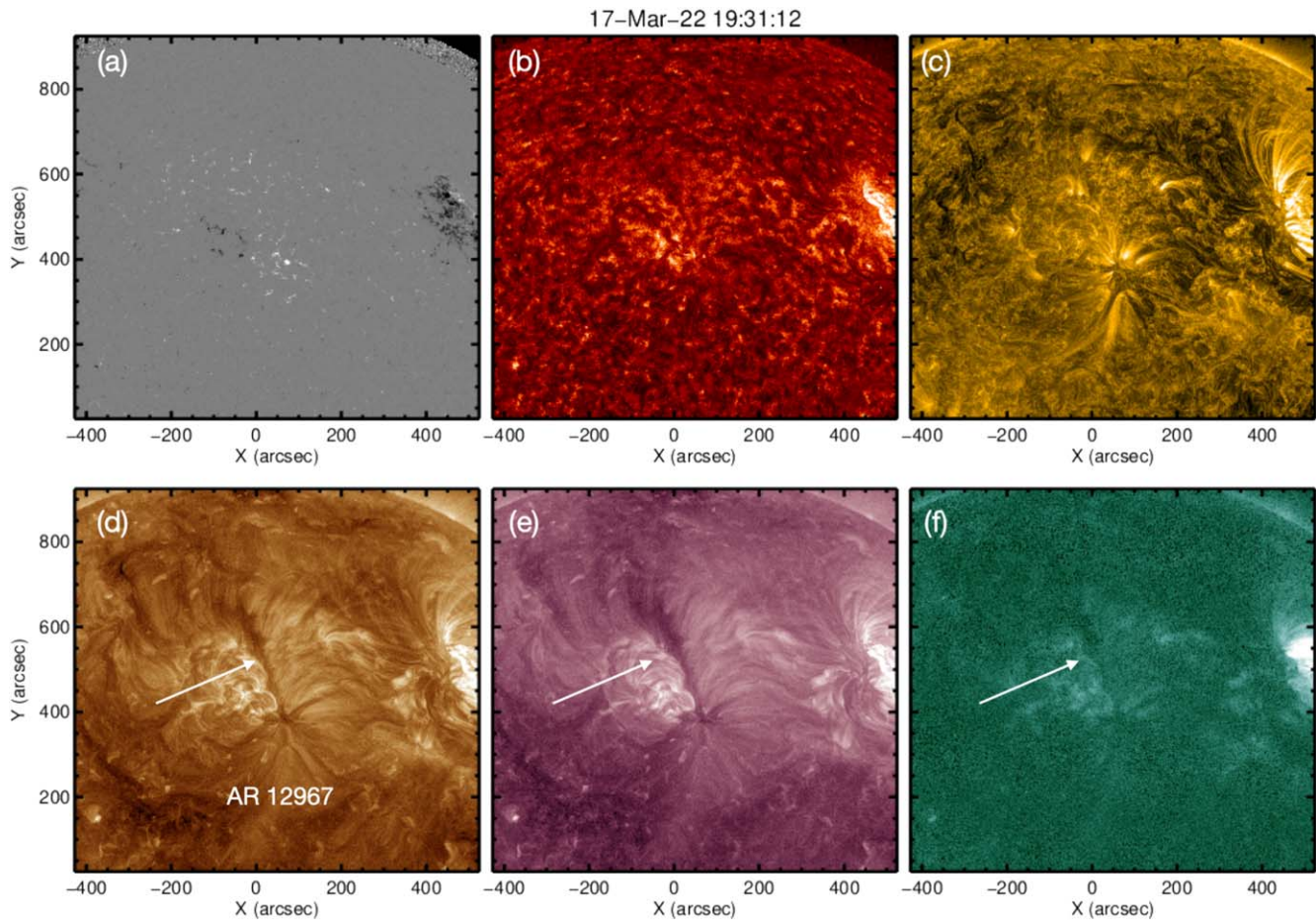
### 3.1. SDO/HMI Observations

NOAA AR 12967 rotated onto the solar disk on 2022 March 12 (from the viewpoint of Earth). Its leading polarity contained a small, coherent positive sunspot bordered to the north and south by dispersed positive field. This configuration is shown in the magnetogram from the Helioseismic and Magnetic Imager (HMI; Scherrer et al. 2012) on board the Solar Dynamics Observatory (SDO; Pesnell et al. 2012) in Figure 1(a). By 2022 March 18, the sunspot was no longer visible in the HMI magnetograms. The following negative polarity was fully dispersed from the time it was first observed at the east limb, suggesting that the AR was already in its decay phase. At central meridian passage, the AR’s total unsigned magnetic flux was  $\sim 3.5 \times 10^{21}$  Mx, typical for a medium-sized AR with a lifetime measured in weeks rather than months (van Driel-Gesztelyi & Green 2015).

The distribution of the photospheric magnetic field remained broadly unchanged throughout the AR’s transit across the solar disk (see the online animation associated with Figure 1). Its following negative polarity field was embedded in the surrounding positive field in the southern section of the AR. In the AR’s northern section, the dispersed positive field extended toward the northern polar CH.

### 3.2. SDO/AIA Observations

Images from the 304 Å ( $\log T = 4.7$ ), 171 Å ( $\log T = 5.8$ ), 193 Å ( $\log T = 6.2$  and 7.3), 211 Å ( $\log T = 6.3$ ), and 94 Å ( $\log T = 6.8$ ) channels of SDO’s Atmospheric Imaging Assembly (AIA; Lemen et al. 2012) are displayed in Figures 1(b)–(f). The multipanel image is taken from the online animation associated with this figure. The movie shows the chromospheric/coronal evolution of NOAA AR 12967 from 2022 March 14 12:00 UT to 2022 March 20 23:56 UT. Each image has a field of view (FOV) of  $900'' \times 850''$  encompassing NOAA AR 12967 and the surrounding dispersed magnetic field. Bright core loops connecting the AR’s main polarities are visible in the hotter



**Figure 1.** (a) SDO/HMI line-of-sight magnetogram and SDO/AIA (b) 304 Å, (c) 171 Å, (d) 193 Å, (e) 211 Å, and (f) 94 Å images of NOAA AR 12967 at 2022 March 17 19:31 UT. The arrow indicates the dark open-field corridor described in the text (see Section 3.2). An animation of these data is available in the online Journal. It shows the chromospheric/coronal evolution of NOAA AR 12967 from 2022 March 14 12:00 UT to 2022 March 20 23:56 UT (duration is 57 s).

(An animation of this figure is available.)

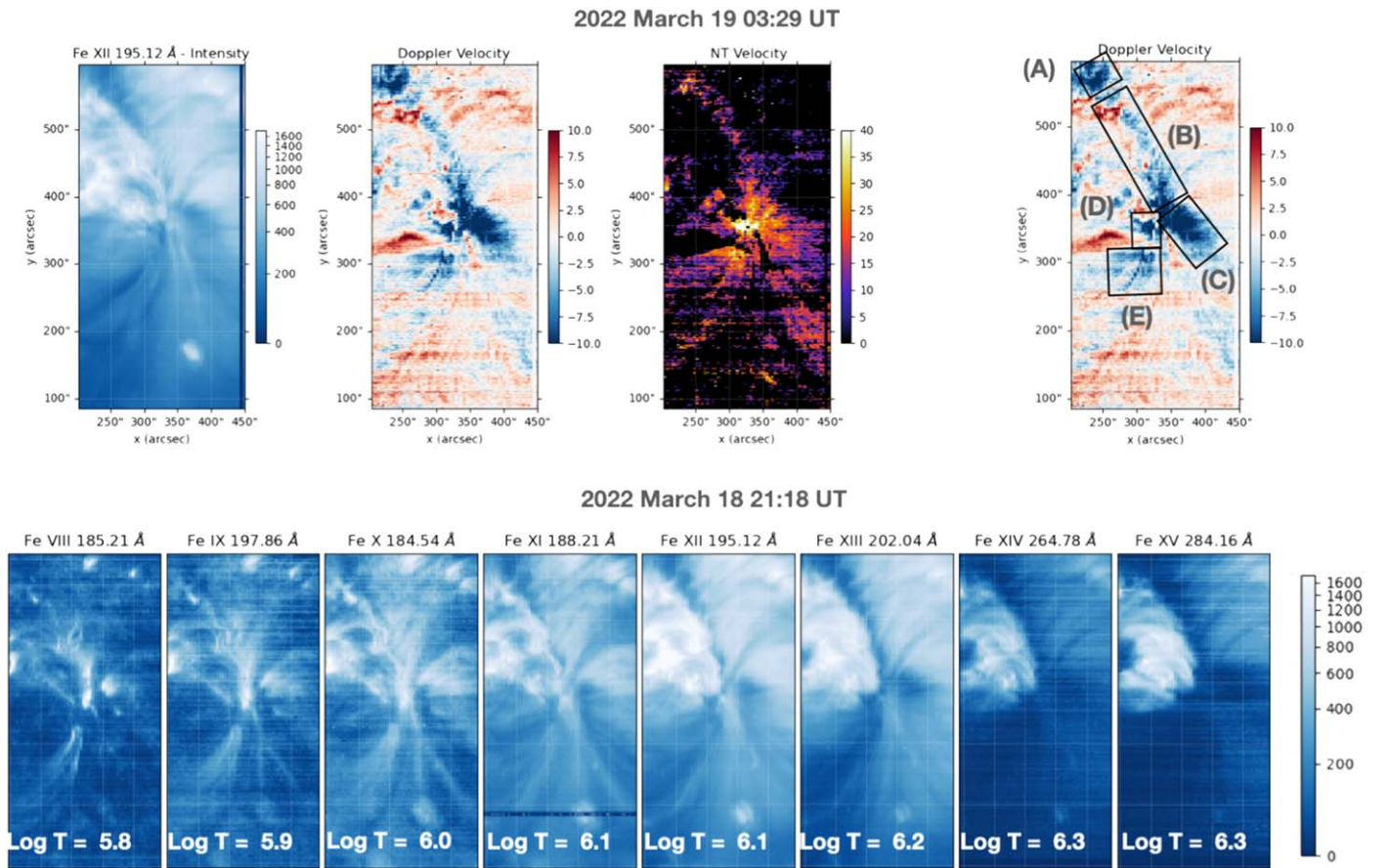
AIA channels. Less bright, longer loops extend from around the positive sunspot to the negative fields of the nearby NOAA AR 12965 in the west and quiet Sun in the south. Loops rooted along the channel of positive field in the north appear to connect with the negative polarity field of NOAA AR 12967 on the eastern edge and negative field of NOAA AR 12965 on the western edge, forming a dark, narrow corridor that extends from the positive sunspot toward the north polar CH.

The dark corridor is a striking feature in the hotter 193 Å, 211 Å, and 94 Å channels (see arrows in Figure 1) but less evident in the chromospheric 304 Å and upper transition region/lower corona 174 Å channels. As the AR rotates across the solar disk, projection effects mean that some of the corridor is obscured. This is especially the case in the 171 Å channel, where there is a mixture of short, low-lying loops and long fan loops in the vicinity of the corridor. However, when the AR is close to central meridian on 2022 March 17, the full length of the corridor is clearly observed in the 193 Å, 211 Å, and 94 Å channels. In these channels, the corridor significantly widens from where it starts at the positive sunspot until it blends into the northern polar CH extension. The corridor persists and remains remarkably stable during the period of 5.5 days covered by the included movie, even after a filament eruption along the northern external polarity inversion line at 2022 March 16 ~13:00 UT.

### 3.3. Hinode/EIS Observations

Hinode/EIS provided coordinated observations during the slow wind connection SOOP. EIS tracked the positive polarity of NOAA AR 12967 from 2022 March 18 to 21. Study ID 600/acronym DHB\_007\_v2 was employed for a total of 12 observations. The 2'' slit was rastered in 4'' steps taking 60 s exposures at each of the 62 pointing positions. A single observation took 62 minutes to build an FOV of 248'' × 512''. The north–south extent of the FOV covered the upflow region over the positive polarity and the dark corridor until the spacecraft pointing was shifted further south on 2022 March 20. Cotemporal observations from SO's Spectral Imaging of the Coronal Environment (SPICE; SPICE Consortium et al. 2020) instrument do not cover the corridor; therefore, they are not included in this analysis.

All spectroscopic data were processed with the EIS Python Analysis Code (EISPAC; <https://github.com/USNavalResearchLaboratory/eispac/>) software package using Python compatible level1 HDF5 files available at <https://eis.nrl.navy.mil>. The Fe ion emission lines were fitted with either single- or double-Gaussian functions applying the template files included in EISPAC. Single-Gaussian functions were fitted to all lines with the exception of the blended Fe XI 188.21 Å, Fe XII 195.12 Å, and Fe XV 284.16 Å lines.



**Figure 2.** Hinode/EIS Fe XII 195.12 Å intensity, Doppler velocity, and nonthermal velocity maps (top left to right) at 2022 March 19 03:29 UT and Doppler velocity map overlotted with boxes indicating regions of interest in the corridor (top right). The date is close to the beginning of connectivity between the corridor and SO (see B1 in Figure 8). Bottom: intensity maps of Fe ions at 2022 March 18 21:18 UT showing the expansion of the corridor width with temperature/height (see Section 3.3). At this time, the corridor is close to solar central meridian, where the corridor is best viewed.

The top panel of Figure 2 shows a sample of the Fe XII 195.12 Å emission-line intensity, Doppler velocity, and nonthermal velocity maps at 2022 March 19 03:29 UT, when the full length of the dark corridor is evident in the EIS observations. This time at the Sun is roughly equivalent to when the SO spacecraft first encountered the slow SW stream taking into account the magnetic connectivity and the propagation time of the SW (see Section 6). Consistent with the observed stability of the dark corridor in images from the hotter AIA channels, the Doppler and nonthermal velocities for each of the 12 observations are very similar to those shown in the top panel of Figure 2. Blueshifted upflows in the range =  $[-20, 0]$  km s $^{-1}$  are observed in the Doppler velocity map over the positive polarity of NOAA AR 12967, along the dark corridor of the positive field. Nonthermal velocities within the upflow regions are in the range =  $[0, 44]$  km s $^{-1}$ .

Table 1 contains measurements of the plasma and magnetic field parameters for subregions defined in the Fe XII 195.12 Å Doppler velocity map in the top right panel of Figure 2. Only those pixels containing Doppler upflows were considered for each parameter; pixels with downflows were masked. Values were determined using the summed spectra within a region for each emission line to improve the signal-to-noise ratio. The parameters are First Ionization Potential (FIP) bias, Doppler and nonthermal velocities, plasma density using Fe XIII ions, temperature of the peak differential emission measure (DEM), and magnetic flux density (from the corresponding HMI

**Table 1**

Plasma and Magnetic Parameters in Box Regions Defined in the Upper Right of Figure 2 with the Fe XII 195.12 Å Doppler Velocity Map at 2022 March 19 03:29 UT

Boxes	A	B	C	D	E
FIP bias	1.2	2.0	2.1	1.8	2.1
Doppler velocity (km s $^{-1}$ )	-4.1	-11.3	-12.2	-13.3	-10.1
Nonthermal velocity (km s $^{-1}$ )	19.7	25.6	28.1	34.3	25.3
Log $_{10}$ density (cm $^{-3}$ )	8.3	8.6	8.2	8.5	8.3
Log $_{10}$ temperature (K)	6.2	6.2	6.2	6.0	6.2
Peak DEM (10 $^{21}$ cm $^{-5}$ K $^{-1}$ )	2.4	5.7	3.1	6.3	1.8
Magnetic flux density $B$ (G)	14.4	23.2	65.9	76.4	17.7

**Note.** (See Section 3.3).

magnetogram). The temperatures corresponding to the peak DEMs and Si X 258.37 Å – S X 264.22 Å FIP bias were calculated using the method of Brooks & Warren (2011) and Brooks et al. (2015).

Region A is located in the relatively broad northern section of the corridor. The broader width of the region and its low magnetic flux density are more consistent with CH or polar CH extensions. As expected, it has photospheric FIP bias with low upflow and nonthermal velocities (e.g., Harra et al. 2015; Fazakerley et al. 2016). The density is slightly lower in A than in regions B and D and similar to the other two regions; however, there is little contrast in density among the five

regions. The temperatures corresponding to the peak DEMs are similar across the regions with  $\log T$  in the range = [6.0, 6.2]. Where the corridor is narrow in region B, the FIP bias and the upflow and nonthermal velocities are comparable to regions C–E within the AR. However, the magnetic flux there is in between CH and AR values. Overall, the plasma parameters measured in the corridor are closer to those found at the edges of ARs (e.g., Doschek et al. 2008; Harra et al. 2008; Brooks & Warren 2011, 2016) than in CHs.

The images from different AIA channels show that the dark corridor is expanding with height/temperature in the solar atmosphere. In order to verify whether this is the case, EIS intensity maps of emission lines sensitive to narrower temperature ranges are plotted for each ion from Fe VIII 185.21 Å to Fe XV 284.16 Å at 2022 March 18 21:18 UT in the bottom panel of Figure 2. At this time, the corridor is close to the solar central meridian so that projection effects are minimized. The Fe emission lines represent  $\log T$  in the range = [5.8, 6.3] (see the values written at the bottom of the maps in the figure). The corridor is not clearly observed in the Fe VIII intensity map. In the northern section of the Fe IX map, the dark corridor becomes evident, though it is short and narrow; it continues to lengthen and widen in the intensity maps as the  $\log T$  increases from 5.9 to 6.3.

## 4. Magnetic Field Modeling

### 4.1. PFSS Extrapolation

A potential field source surface (PFSS) extrapolation was used to investigate the large-scale coronal magnetic configuration of NOAA AR 12967 and the dark corridor. The PFSS extrapolation of the photospheric magnetic field at 2022 March 18 18:03 UT was created using the IDL SolarSoft PFSS package (Schrijver & De Rosa 2003). With this package, PFSS extrapolations are generated using the evolving surface flux transport (SFT) model based on assimilated HMI data and known methods for evolving the observed magnetic flux. The SFT model tracks unipolar flux elements that are subject to differential rotation and meridional flows. Flux elements will coalesce and cancel. For each time step, the positions and flux amounts of the elements are output. The HMI magnetogram image resolution is then binned down by a factor of four, and all pixels within 60° of Sun center are converted to flux elements. A full-Sun flux map based on the HMI magnetogram inserted into the SFT model is used as the input into the PFSS code to create a 384 × 192 full-Sun magnetogram. Spherical decomposition is performed on the magnetogram to reconstruct the magnetic field as a function of height. The top of the PFSS model is the source surface, set at a radial distance  $R_{ss}$ , defined here to be at 2.5  $R_{\odot}$ .

Various viewpoints of the extrapolation are shown in Figure 3. Open-field lines are overplotted on the Hinode/EIS Fe XII 195.12 Å Doppler velocity map embedded in an SDO/HMI LOS synoptic map of CR 2255 closest in time to that of the EIS raster. Open-field lines are located along the corridor of blueshifted upflows in the Doppler velocity map. The top view in panel (a) includes field lines computed along the central section of the open-field channel, and the side view in panel (b) shows the coherence along the channel with only moderate meridional expansion in the northern and southern sections. A strong lateral expansion is present as shown in panel (c), where

two rows of field lines are drawn on the sides of the open-field corridor.

Two cross sections of the PFSS extrapolation orthogonal to the corridor, one in the north and the other in the south, show that the open-field lines are bordered on either side by closed loops (Figures 4(a), (b)). The field extends radially in the center of the corridor and then diverges away from the radial direction closer to the edges/closed loops.

The expansion properties of the open field are quantified by computing the expansion factor:

$$f(R) = \frac{R_{\odot}^2 B(R_{\odot})}{R^2 B(R)}, \quad (1)$$

where  $R$  is the radial distance from the Sun center,  $R_{\odot}$  is the solar radius, and  $B(R_{\odot})$  and  $B(R)$  are the magnetic field strength at  $R_{\odot}$  and  $R$ . With magnetic flux conservation,  $f(R)$  provides the flux tube cross-section expansion corrected for the spherical expansion. With  $R = R_{ss}$ , the radius at the source surface (set at 2.5  $R_{\odot}$ ),  $f(R_{ss})$  is the expansion factor used in previous studies of SW sources (e.g., Wang & Sheeley 1990; Wang et al. 2009).

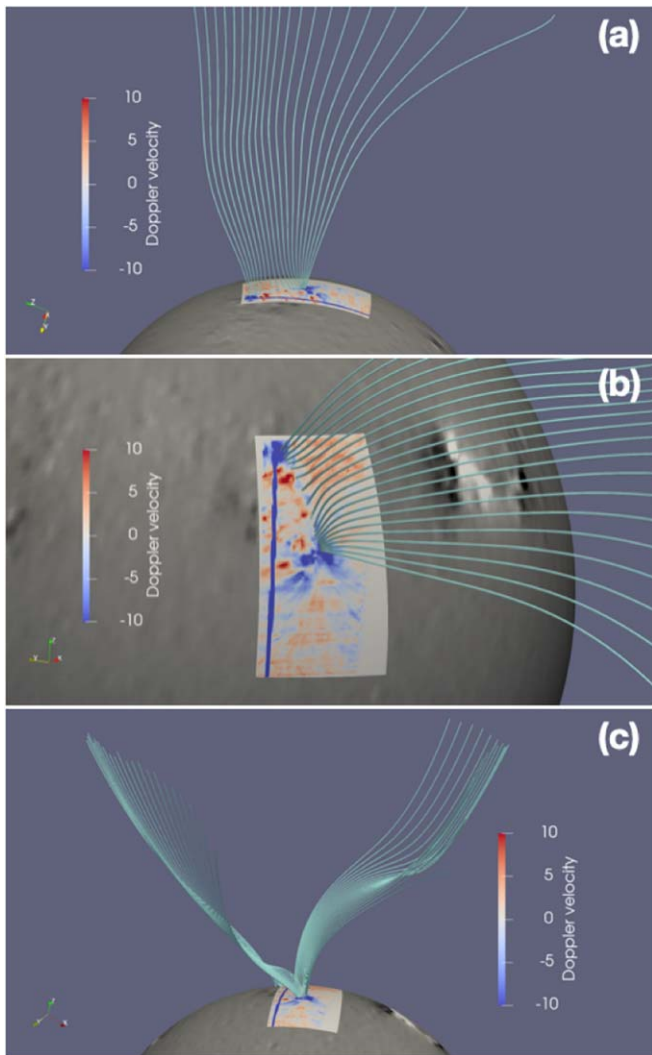
The expansion factor  $f(R)$  is plotted at three representative heights (1.2, 1.5, and 2.5  $R_{\odot}$ ) in Figures 4(c) and (d) for the northern and southern cross sections of the corridor. Cyan squares along each field line correspond to these heights. The X-axis in each plot is the distance in Mm in the photosphere of each field line from the leftmost (easternmost) one. In general, the width of the open-field corridor decreases, and the asymmetry of the closed field increases from north to south. The northern slice has an expansion that is smaller in the central part with respect to the lateral sides. At the edges, the open flux can expand rapidly when the top of the lateral closed field is reached. This magnetic field behavior is quantified in Figure 4(c). The plot of  $f(R)$  shows a smaller expansion at low heights while reaching larger values at greater heights on both sides.

Expansion properties change as the surrounding closed field and the width of the corridor vary (see abscissae at the bottom of panels (c) and (d)). An extreme in field-line asymmetry is reached along the southern slice (Figure 4(b)), where the small size of the eastern arcade field allows an expansion of the corridor field at lower heights ( $\lesssim 1.2 R_{\odot}$  in Figure 4(d)). In contrast, the much larger western arcade field limits the corridor expansion at low heights, but when the maximum height of the closed field is exceeded (approx 1.5  $R_{\odot}$ ), the open field has a bigger volume into which it can expand until it reaches pressure balance with the nearby open field. This leads to larger  $f(R_{ss})$  on the western side than on the eastern side.

Previous studies have shown that different expansion profiles, in particular  $f(R_{ss})$ , have direct implications for the acceleration profiles of the SW, especially its terminal speed (Wang et al. 2009; Pinto et al. 2016; Pinto & Rouillard 2017). Depending on which flux tube is crossed by SO, different SW properties, such as plasma speed, are to be expected with in situ measurements.

### 4.2. Q-Maps

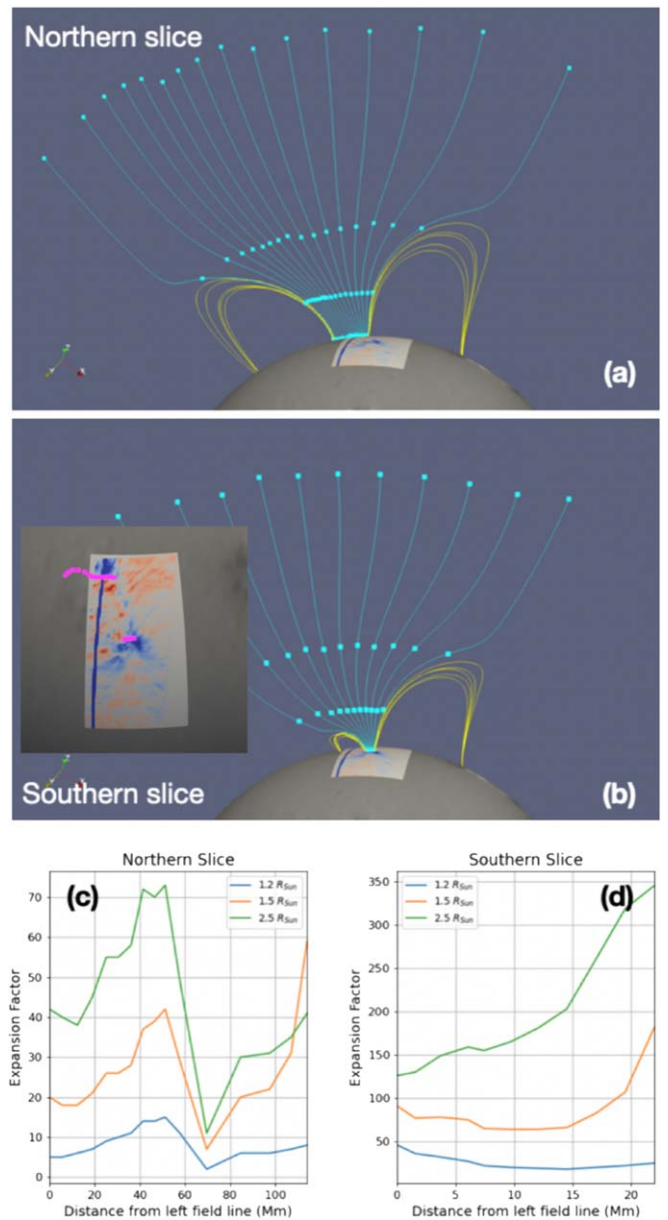
The open flux corridor is bordered by a drastic change in magnetic connectivity (Figure 4), with field lines changing from reaching the source surface to reaching back to the photosphere as closed field. Such a magnetic configuration is typical of the presence of a separatrix and more generally of



**Figure 3.** PFSS extrapolation with Hinode/EIS Fe XII 195.12 Å Doppler velocity map at 2022 March 18 19:42 UT overlaid on SDO/HMI LOS magnetogram. Panels (a)–(c) show open-field lines from varying viewpoints of the dark corridor (see Section 4.1). The viewpoints are (a) top view, (b) side view, and (c) along the corridor from south to north. In panels (a) and (b) field lines are set on the central part of the corridor, and they are set on both sides of it in panel (c).

QSLs. Initially, QSLs were defined in closed-field regions to understand the location of flare ribbons in magnetic configurations where there are no magnetic bald patches or null points present, i.e., where the magnetic topology is a priori as simple as a magnetic arcade. In fact, for ARs, typically thin coronal volumes are present where the magnetic connectivity changes are extreme. These thin regions define QSLs (Démoulin et al. 1996). In the limit where the connectivity change is discontinuous, a QSL reduces to a separatrix. A convenient way to identify strong QSLs is by using the squashing factor,  $Q$ , which quantifies the change in the magnetic connectivities at a boundary (Titov et al. 2002). More generally,  $Q$  is a geometric measure of the degree to which neighboring field lines diverge in a 3D volume.

The concept of QSLs was extended to open/closed magnetic configurations by considering the magnetic connectivity between two surfaces, e.g., the photosphere and the source surface (Titov et al. 2011), as was previously done for straightened coronal configurations (e.g., Milano et al. 1999).

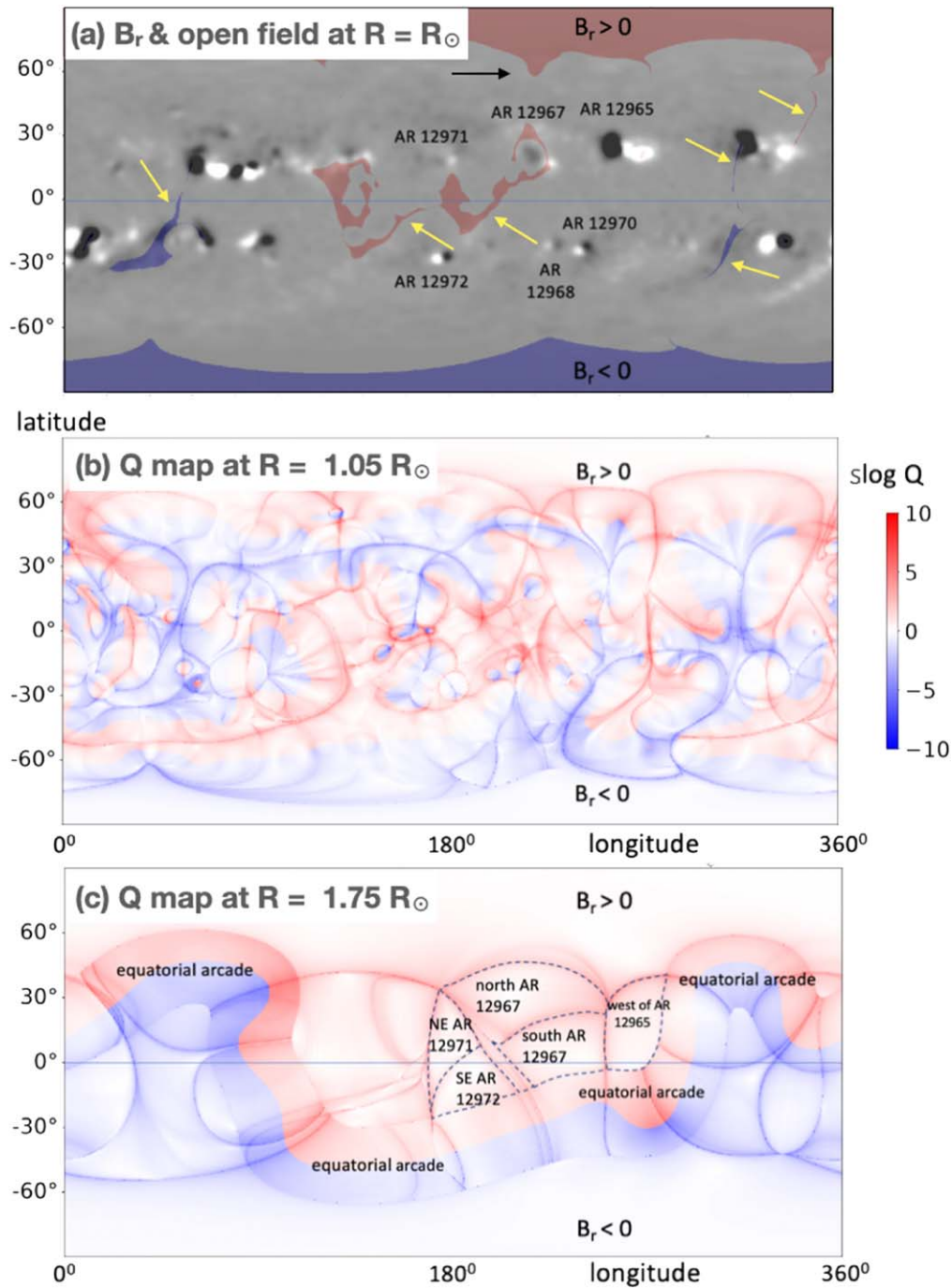


**Figure 4.** Field lines along slices orthogonal to the corridor of upflows in the (a) north and (b) south. Blue/yellow field lines represent the open/closed field. Pink circles in the inset of panel (b) identify the locations of the slices. Plots of expansion factors for (c) northern and (d) southern slices vs. distance in Mm between footpoints of field lines in the photosphere starting from the left. The three squares along each field line, in panels (a) and (b), correspond to radii of 1.2, 1.5, and 2.5  $R_{\odot}$ , and the expansion factor at each of the heights is plotted in blue, orange, and green, respectively, in the bottom panels (see Section 4.1).

Since field lines extending outward from the source surface are supposed to follow a Parker spiral, there is no change in the  $Q$ -maps for a surface more distant than the source surface.

Synoptic maps of  $Q$  for each Carrington rotation (CR) are available from the SDO’s Joint Science Operations Center (JSOC; see <http://hmi.stanford.edu/QMap/>). Figure 5(a) shows the SDO/HMI synoptic magnetogram for CR 2255, with the photospheric footpoints ( $r = R_{\odot}$ ) of the open field in red/blue designating positive/negative radial  $B$  field components. The main ARs on the Sun near NOAA AR 12967 are labeled. Open positive field encompasses the negative polarity of NOAA AR 12967 in the configuration of a classic pseudo-streamer (Wang et al. 2007). In the northern direction, there is



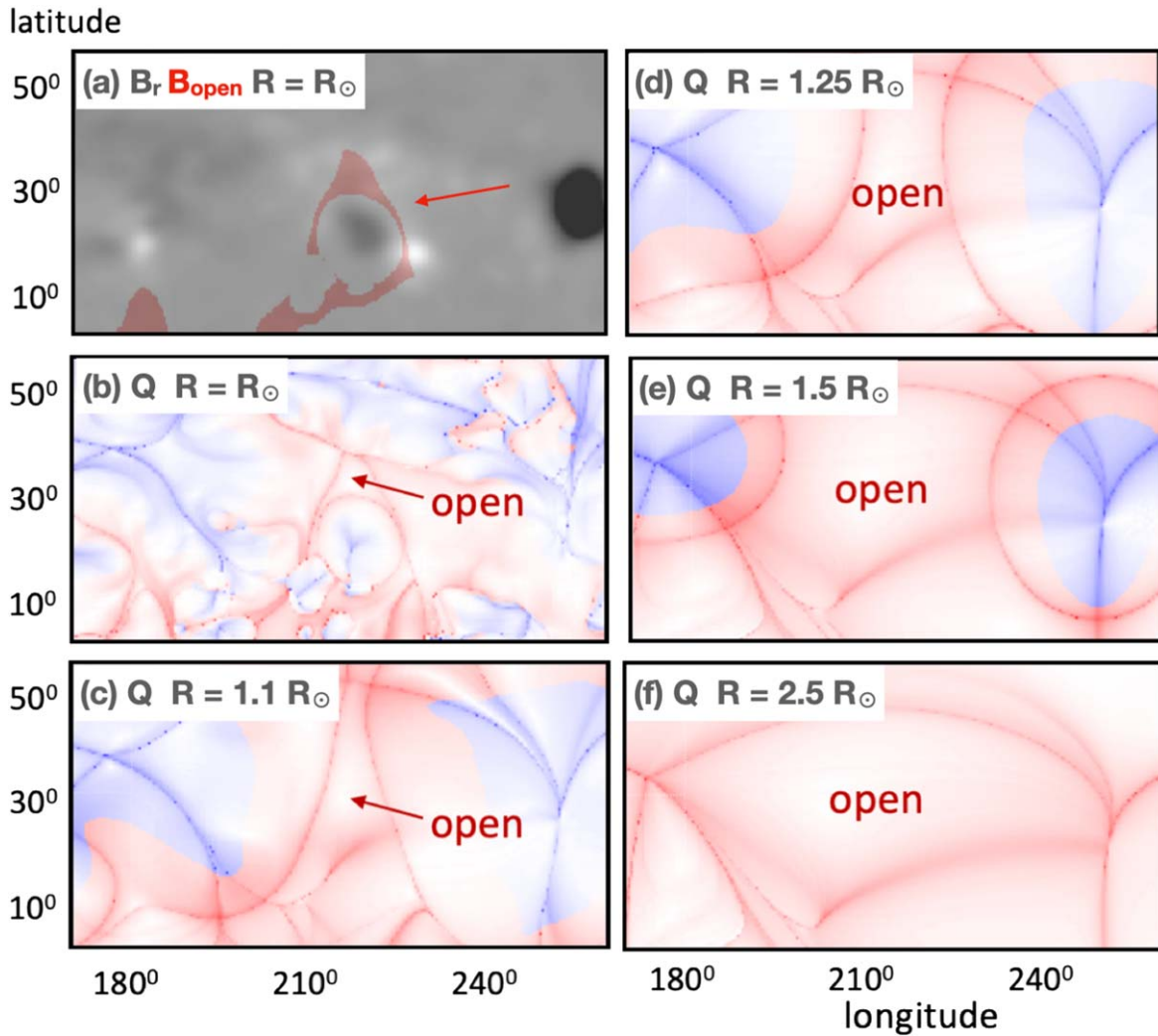


**Figure 5.** Full synoptic maps from <http://hmi.stanford.edu/QMap/> (see Section 4.2). (a) Full synoptic SDO/HMI magnetogram showing photospheric footpoints of the open field (red/blue is positive/negative smoothed radial  $B$  field component); (b) full synoptic  $Q$ -map at a height of  $1.05 R_{\odot}$  for CR 2255, where  $Q$  is the squashing factor; and (c) full synoptic  $Q$ -map at a height of  $1.75 R_{\odot}$ , with dashed lines superposed on high  $Q$  values showing the limits of open-field regions in the vicinity of NOAA AR 12967. The “equatorial” arcade is visible as a red/blue serpentine-like ribbon surrounding the Sun (beneath the heliospheric current sheet). The black arrow in panel (a) indicates the northern polar CH extending toward NOAA AR 12967. Yellow arrows in panel (a) identify many of the S-web corridors in the synoptic map.  $\text{slog } Q$  is defined by  $\text{sign}(B_r) \times \log [Q/2 + (Q^2/4-1)^{0.5}]$ . Large values of  $\text{slog } Q$  define the boundaries between regions of similar connectivities.

an extension of the northern polar CH indicated by the black arrow in the figure. It is plausible that the extension is connected to the open-field region by a thin channel (Antiochos et al. 2011) that is not detectable owing to the low spatial resolution of the connectivity computations.

Even though magnetic synoptic maps have highly degraded spatial resolution compared to the original HMI magnetograms and with the spatial variations further smoothed with height by a potential field extrapolation, the synoptic  $Q$ -map at a height

of  $1.05 R_{\odot}$  in Figure 5(b) gives a sense of the complexity of the magnetic connectivities. This is due to the fact that the photospheric magnetic fields of numerous sources are linked together. Local magnetic flux imbalances, such as with the negative polarity of NOAA AR 12967, are typically associated with closed magnetic connections to the surrounding field of the dominant opposite polarity. It results in a round flower-like  $Q$  pattern, with the “petal” boundaries defined by high  $Q$  values (see, e.g., Figure 5(b) and Figures 6(d) and (e)). The boundaries



**Figure 6.** (a) Zoomed-in synoptic SDO/HMI magnetogram showing photospheric footpoints of the open field (red/blue is positive/negative smoothed radial  $B$  field component) around NOAA AR 12967. (b–f) Zoomed-in  $Q$ -maps at heights ranging from the photosphere at  $R_{\odot}$  to the source surface at  $2.5 R_{\odot}$ . Dark colors in  $Q$ -maps represent high  $Q$  values (see the scale to the right of Figure 5(b)). The red arrow in panel (a) indicates the open-field corridor.

separate the central polarity connecting to surrounding opposite polarities of larger magnetic flux. This is the case around the negative polarity of NOAA AR 12967, and large-scale examples are also present in Figure 5(b) (e.g., NOAA AR 12965 and NOAA AR 12970). Boundaries of high  $Q$  are the same as the giant arcs of the S-web mentioned in the Introduction. The  $Q$ -maps in Figures 5(b) and (c) can be compared to Figure 7(a) of Antiochos et al. (2011), Figure 3 of Scott et al. (2018), and Extended Data Figure 4(c) of Chitta et al. (2023), revealing comparable patterns of  $Q$ .

When  $Q$ -maps are computed at greater heights, the flower boundaries retract as shown in Figure 5(c) ( $R = 1.75 R_{\odot}$ ). At such heights, all of the closed-field regions have disappeared except for a large-scale arcade that encircles the Sun (below the heliospheric current sheet). Without the presence of the ARs, this would correspond to an equatorial arcade with a streamer above it. In this active period, the arcade is warped to significant latitudes by the AR flux distribution. The rest of the sphere surface at  $1.75 R_{\odot}$  in Figure 5(c) is covered by open fields coming from different photospheric sources, so they are separated by high  $Q$  values. The photospheric linkage is indicated by labels around NOAA AR 12967. At the source surface and farther away, only these domains of different

connectivities remain. In summary,  $Q$ -maps show the extension of the various sources of the SW (which come into contact at the high  $Q$  locations).

Figure 6(a) is a zoomed-in version of the synoptic magnetogram centered on the AR target and its surroundings. The open-field corridor introduced in Section 3.2 is the thin strip of positive field on the western side of the AR that broadens in its northern extension. The closed-field region surrounding the negative polarity of NOAA AR 12967, present in Figure 6(b), disappears in the map of Figure 6(c), which corresponds to a height difference of  $0.1 R_{\odot}$ . For other AR polarities, the dome of closed field can extend to much greater heights such as the eastern and western ARs (NOAA AR 12971 and NOAA AR 12965), where the round QSLs are still present at  $r = 1.5 R_{\odot}$  (right and left sides of Figure 6(e)). It is the retraction of the closed field with height that allows the broad lateral expansion of the open field present in between the ARs. Figure 6 provides a complementary view to the lateral expansion of the open field shown in Figure 4.

In the  $Q$ -map at  $r = 1.1 R_{\odot}$  (Figure 6(c)), the negative polarity of NOAA AR 12967 is fully covered by open field forming a triangular-like shape. The open-field region continues to expand up to  $r = 2.5 R_{\odot}$  (Figures 6(d)–(f)). After

starting as a narrow region at  $r \approx R_\odot$ , the corridor broadens greatly in the longitudinal direction, in contrast to only marginally increasing in the latitudinal direction. This is in agreement with the observed broadening of the corridor with increasing temperature and height in the Hinode/EIS ion intensity maps (Figure 2). The significant broadening with height is related to the large expansion factors found in Figure 4. As a consequence, a spacecraft orbiting close to the solar equatorial plane is expected to cross a wide range of longitudes with plasma originating from the corridor even if remote observations only show a thin latitudinal extent of the source region (Figures 1, 2). Indeed, the extent of the magnetic field at  $r = 2.5 R_\odot$  is up to about  $60^\circ$  of longitude (Figure 6(f)), which is comparable to the typical distance between neighboring ARs in Figure 5(a).

### 5. Magnetic Connectivity

The Magnetic Connectivity Tool (MCT; Rouillard et al. 2020) was used to connect the in situ measurements at SO to the possible on-disk source regions. The MCT models the corona with a potential field computed from a synoptic map of the photospheric field, and then it performs a PFSS reconstruction similar to what is described in Section 4.1. The photospheric field could be either an NSO/GONG synoptic map derived directly from observed line-of-sight magnetograms or an ADAPT synoptic map using flux transport schemes to evolve the observed magnetic field near the central meridian to the required time (Hickmann et al. 2015). Examples of NSO and ADAPT synoptic magnetograms are provided in the Appendix. For the PFSS reconstruction, the magnetic field is computed up to the source surface ( $R = R_{ss}$ ) where the magnetic field is supposed to be radial. At larger distances, the interplanetary magnetic field is assumed to be a Parker spiral with a constant SW speed.

Based on ADAPT synoptic maps, the MCT establishes partial or full connectivity between the positive polarity region of NOAA AR 12967 and SO during the period 2022 March 20–26 (spacecraft time corresponding to 2022 March 18–24 at Sun time, i.e., an SW parcel departs from the Sun at the Sun time and arrives at SO at the spacecraft time). From 2022 March 20–22, connectivity is found to be to the positive fields of both NOAA AR 12967 and NOAA AR 12965 (to the west). A full and stable connection to NOAA AR 12967 is achieved only after the beginning of 2022 March 22 (spacecraft time). Partial disconnection begins early on 2022 March 25, when SO is connected to the positive field on either side of the negative polarity of NOAA AR 12967. SO is fully disconnected from NOAA AR 12967 from 2022 March 27.

Figure 7 illustrates the connectivities and their probabilities from MCT ADAPT maps during each interval (2022 March 17, 20, 23 at Sun times) plotted on SDO/AIA 193 Å, SDO/AIA 171 Å, and SO EUV/FSI 174 Å maps. Distributions of connectivity points are calculated using multiple field lines generated by sampling an uncertainty ellipse region around the target spacecraft position and then remapping the respective probability densities across the interplanetary and coronal magnetic field (Rui Pinto, 2022, private communication). The percentages in the SO/EUV column of images are the probabilities for each grouping of circles. It should be noted that connectivity depends on a number of factors, including the data input (e.g., types of magnetograms, cadence of magnetograms) at the base of the PFSS reconstruction, the height of  $R_{ss}$

(Panasenco et al. 2020; Koukras et al. 2022), the wind speed and its acceleration profile (e.g., as deduced from a statistical analysis; Dakeyo et al. 2022), and the inclusion of more realistic physics. Differences in the timing and connectivities produced by varying the input parameters demonstrate the limitations of setting the connectivity with such basic models (PFSS and Parker spiral). Nonetheless, this method provides clear hints of the magnetic connectivity between NOAA AR 12967 and SO, while the timing is refined below with the in situ data.

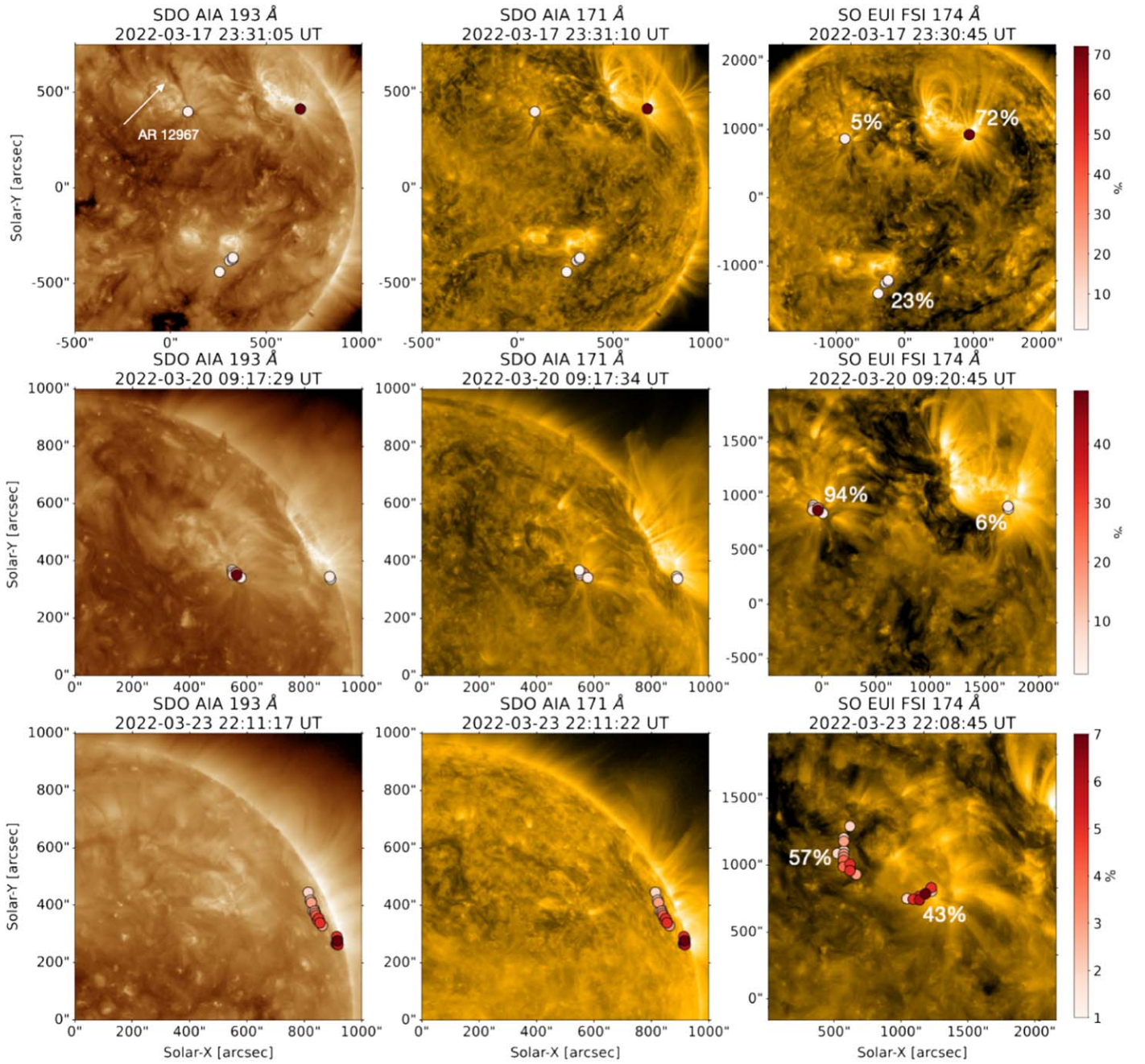
### 6. In Situ Observations

SW data were obtained from SWA-PAS and MAG in situ instruments on board SO. Figure 8 shows SWA-PAS 4 s time resolution radial velocity  $V_R$ , proton number density  $N_p$  (black), and proton temperature  $T_p$  (blue) data in panels (a)–(b) and total  $|B|$  and radial  $B_R$  MAG magnetic field data in panels (d)–(e) for the period 2022 March 19–25. As is also the case for the synoptic maps,  $B_R$  in the outward/inward direction from the Sun is defined as positive/negative field. Panel (c) is the absolute value of the correlation coefficient  $C_{vb}$  between  $V_N$  and  $b_N$  (with  $b_N$  being the normal component of the magnetic field in Alfvén units,  $b_N = B_N/(4\pi\rho)^{1/2}$ , with  $\rho$  being the mass density and  $B_N$  the normal component of the magnetic field) computed over a running window of 30 minutes, which is a typical Alfvénic scale in the SW (Tu & Marsch 1995; Bavassano et al. 1998). All in situ data are plotted against measurement time at the SO spacecraft (in the range = [0.32, 0.38] au).

It should be noted that for SW velocities below  $\sim 275$ – $300$  km s $^{-1}$  a part of the velocity distribution function may fall into the low-energy measurement range of the SWA-PAS sensor. It is known that there is an instrumental issue at these low energies at which the counting efficiency is reduced. This means that under these unusual circumstances the low-energy ions can be undersampled. If a significant part of the distribution lies in this low-energy range, then there is likely to be an effect on the subsequent calculation of the moments (velocity, density, and temperature) of the distribution. A quality factor parameter is provided by the sensor teams as an assessment on the potential level of impacted data points: a quality factor of  $<0.2$  means that there are no significant issues with the derived moments;  $<2.0$  means that there can be some undersampling of the low-energy data;  $>2.0$  means that undersampling may be significant and caution must be used in quantitative interpretation of the data (A. Fedorov, 2023, private communication).<sup>20</sup> The data in interval B1–B2 (Figure 8) have a median quality factor = 1.6 and  $\sigma = 1.6$ . The high densities of this very slow SW period may be underestimated.

Magnetic mapping techniques produce a relatively wide window of partial/full connectivity between the open-field corridor of NOAA AR 12967 and SO (Section 5). The sharp transitions observed in the in situ data provide for better-defined boundaries of this interval (B1–B2 in Figure 8).  $B_R$  (panel (e)) changes from negative to positive at 2022 March 21  $\sim 04:00$  UT. It remains clearly positive on the 22nd and generally positive on the 23rd, though showing far more variability. However, early on 2022 March 24,  $B_R$  has low

<sup>20</sup> See also “SWA/PAS Tutorial” at [https://drive.google.com/drive/folders/1d2Y-G0BiAyAyQqTXL6x9zI39PosoWay\\_](https://drive.google.com/drive/folders/1d2Y-G0BiAyAyQqTXL6x9zI39PosoWay_).



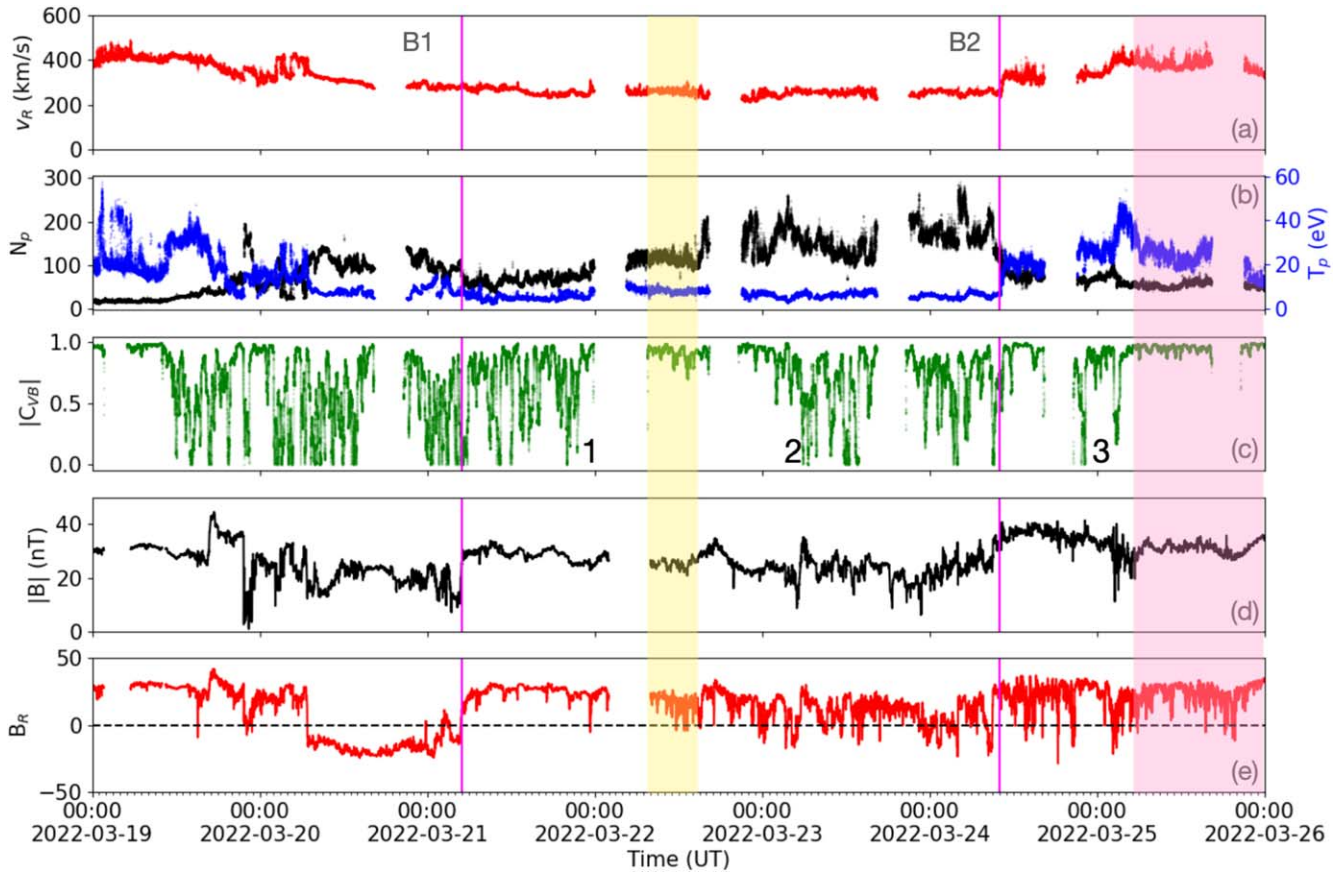
**Figure 7.** SDO/AIA 193 Å, SDO/AIA 171 Å, and SO/EUI FSI 174 Å images (left to right) on 2022 March 17, 20, and 23 (top to bottom, Sun times) showing connectivities from the MCT at Sun times. This is between 1.5 and 2.25 days earlier than spacecraft times for this data set. Separation angle between SDO and SO is  $40^\circ$ . Colors of circles show the probability of the connectivity (see color-coding on the right side and Section 5). The percentages are the probabilities for each region. The white arrow indicates the open-field corridor of S-web.

mean positive values with several excursions into the negative domain before becoming suddenly mostly positive after B2.  $|B|$  slightly increases at this transition and thereafter has a nearly constant value. Abrupt changes in both  $B_R$  and  $|B|$  are the result of a change in magnetic connectivities. Boundary B2 is further confirmed by the fact that the plasma parameters show discontinuities at the time associated with this boundary in Figures 8(a) and (b).

The time period covered between B1 and B2 is from 2022 March 21  $\sim$ 04:00 UT to 2022 March 24  $\sim$ 09:00 UT, so (with a duration of 3 days and 5 hours) slightly less than the duration of about 3.5 days of stable and unique connection to the

western positive polarity of NOAA AR 12967 given by the mapping from both ADAPT and NSO synoptic maps. During the B1–B2 interval, upflows along the corridor observed on the Sun (Figure 2) are expected to be outflows as determined from the magnetic field extrapolations (Section 4.1) and in situ data.

The highly variable density of the slow SW stream in the B1–B2 interval increases from approximately  $60\text{--}80\text{ cm}^{-3}$  at B1 to  $150\text{--}200\text{ cm}^{-3}$  at B2, with spikes approaching  $300\text{ cm}^{-3}$ . The level of  $N_p$  is high since  $100\text{ cm}^{-3}$  at a distance of 0.35 au translates to  $12\text{ cm}^{-3}$  at 1 au, assuming a classical  $r^{-2}$  dependence; therefore, this stream would have a density between 7 and  $25\text{ cm}^{-3}$  at 1 au.  $N_p$  is comparable to the



**Figure 8.** SO SWA-PAS (panels (a)–(b)) and MAG (panels (d)–(e)) in situ data: (a) proton velocity  $V_R$  ( $\text{km s}^{-1}$ ) in red, (b) proton number density  $N_p$  ( $\text{cm}^{-3}$ ) in black and proton temperature  $T_p$  (eV) in blue, (c) absolute value of correlation coefficient  $|C_{vb}|$  in green, (d) total magnetic field  $|B|$  (nT), and (e) radial magnetic field component  $B_R$  (nT) in red.  $|C_{vb}|$  is computed over a running window of 30 minutes. The pink vertical lines B1 and B2 denote the start and end of the interval of connectivity with the narrow corridor. Shaded areas indicate periods of moderate Alfvénicity. Times are at the spacecraft, which is between 1.5 and 2.25 days later than at Sun times (see Figure 7). Numbers correspond to locations in the extrapolation in Figure 9. These locations indicate the inverted-U-shaped Alfvénicity pattern discussed in the text.

denser slow SW measured at 0.35 au by Helios (see Figure 5 of Maksimovic et al. 2020). Furthermore, apart from local small spikes,  $V_R$  is remarkably uniform. It is close to the slowest wind velocities measured at this solar distance with  $V_R$  around  $300 \text{ km s}^{-1}$ . The proton temperature  $T_p$  of a few eV ( $\sim 10^4 \text{ K}$ ) is similar to the Helios slow SW measurements at about 0.3 au (Perrone et al. 2020). This is consistent with the known correlation between  $T$  and  $V_R$  at 1 au (Lopez & Freeman 1986; Elliott et al. 2005, 2012; Perrone et al. 2019).

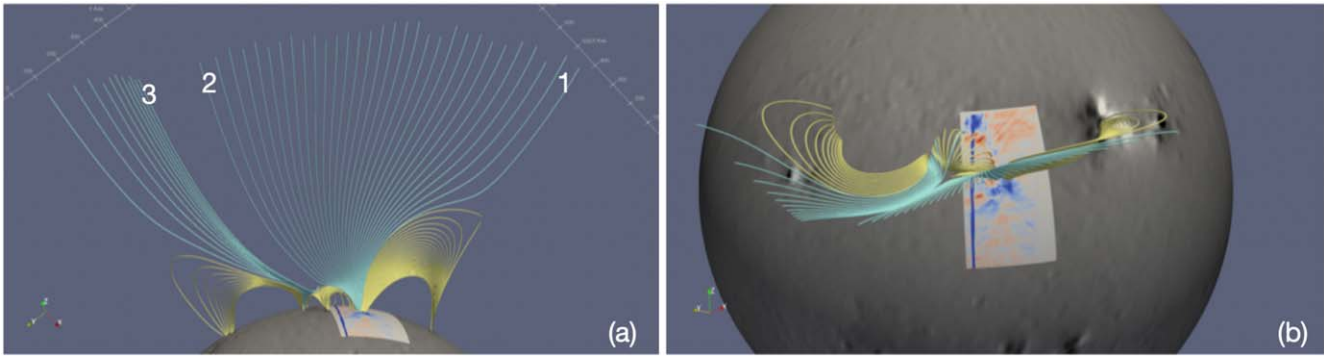
The magnetic field strength in the interval B1–B2 is homogeneous apart from fast spikes of lower values (Figure 8(d)). The radial component  $B_R$  is mostly antisunward directed as expected from the positive polarity of the corridor. The large variability and occasional reversal events are likely to be switchbacks (e.g., Neugebauer & Goldstein 2013; Horbury et al. 2018, 2020). More globally,  $B_R$  is progressively decreasing to smaller values before increasing at the end of the interval close to B2.

An indication of Alfvénicity is a strong correlation between  $V$  and  $B$  components; however, for the sake of simplicity, only the normal component is considered here since it is more Alfvénic than the radial and tangential components (Tu et al. 1989). In the B1–B2 interval,  $|C_{vb}|$  oscillates in the range of  $[0, 1]$ , with spikes below  $\approx 0.5$  indicating poor Alfvénic correlation. However, on 2022 March 22 the  $V$ – $B$  (component) correlation is more stable and in the range  $= [0.7, 1.0]$  for

approximately 1/3 of a day (see the yellow shaded area in Figure 8(c)). In addition,  $N_p$  and  $B$  are roughly constant during the Alfvénic period, suggesting low compressibility.  $N_p$  then exceeds  $200 \text{ cm}^{-3}$  accompanied by rising  $|B|$  on either side of the second data gap, indicating more mixed compressibility. In combination, the in situ observations indicate the presence of Alfvénic fluctuations in the shaded region during the B1–B2 interval. Next, from early on 2022 March 23,  $|C_{vb}|$  oscillates between 0 and 1 again until approximately 2022 March 25 05:00 UT, after which the slow wind stream is more Alfvénic (red shaded area).

After B2, SO is mainly connected to the diffuse positive polarity on the east side of NOAA AR 12967’s negative polarity according to ADAPT synoptic maps. The open magnetic field there does not expand to the same extent as on the western side (see Figure 9, which shows the magnetic field along an east–west slice of the extrapolation through the negative polarity of NOAA AR 12967). This is indicative of a much lower expansion factor compared to the western side, in agreement with higher in situ velocities and proton temperatures and lower plasma density. A detailed analysis of the Alfvénic SW intervals during the 30-day perihelion passage of SO is available in R. D’Amicis et al. (2023, in preparation).

Linking SW streams to their source regions is not straightforward. In this case, the SO MCT connectivity predictions after the B1–B2 interval are widely different,



**Figure 9.** Same PFSS extrapolation and Hinode/EIS map as in Figure 3 from (a) side and (b) top viewpoints. Blue/yellow field lines represent open/closed field. A topological dome encompasses the negative polarity of NOAA AR 12967 in panel (b). Numbers correspond to those in Figure 8.

depending on whether the input magnetograms are NSO or ADAPT (see images and movies in the [Appendix](#)). The NSO LOS magnetogram shows that the connectivity shifts from the positive polarity of NOAA AR 12967 along the channel of positive field leading to dispersed field in the southern hemisphere near NOAA AR 12972 (see Figure 5(a)), whereas the ADAPT connectivity prediction is to the positive field immediately to the east of embedded negative polarity of NOAA AR 12967.

The Alfvénicity level provided by  $|C_{vb}|$  has several inverted-U-shape patterns. The first one is observed from March 22 00:00 UT until March 23 05:00 UT (corresponding to labels 1 and 2 in Figure 8), and a second one (half of an inverted U shape) is present from March 24 03:00 UT until after March 26 00:00 UT (corresponding to label 3 in Figure 8). These inverted-U-shape patterns are consistent with SO changing connectivity from the CH-like field on the western side of the NOAA AR 12967 negative polarity to the CH-like field on the eastern side. More precisely, the locations of these numbers are shown in the coronal extrapolation of Figure 9. It matches the ADAPT connectivity prediction rather than that of the NSO (where the connection is extending progressively to the channel of positive field to the southeast of NOAA AR 12967). The central part of the corridors has a lower expansion factor (e.g., Figure 4) and is associated with  $|C_{vb}|$  close to 1. The sharp fall in Alfvénicity at 1–2 and 3 in Figure 8 is comparable to what is observed in equatorial CHs (Wang & Ko 2019; D’Amicis et al. 2021). These inverted-U-shape patterns are an important marker of the magnetic field differential expansion across open-field regions, i.e., of individual open fields in the low corona where they are spatially separated.

In contrast, the NSO connectivity prediction is in between NOAA AR 12967 and NOAA AR 12972 (see the [Appendix](#) and associated movies). The photospheric magnetic field along this path is highly dispersed compared to that of the narrow corridor; therefore, a lower flux tube expansion with height is to be expected. This is confirmed in the PFSS extrapolation by computing an ensemble of field lines (not shown). Furthermore, there is flux tube contraction with height moving away from NOAA AR 12967 toward NOAA AR 12972. This behavior is different from that within the corridor and the associated yellow shaded region, while the yellow and red shaded regions have broadly similar Alfvénicity. Finally, the distance between the two ARs relative to the distance across the negative polarity of NOAA AR 12967 would require a much larger time period between the yellow and red shaded regions in Figure 8. We

conclude that the in situ observations do not agree with the PFSS computed from NSO synoptic magnetograms.

Plasma composition and charge state are well-known in situ parameters that are useful for refining/confirming connectivity; however, these data were not available from SO’s perihelion. Instead, a combination of the expansion factors in different areas of the PFSS extrapolation and Alfvénicity of the slow SW provides a high confidence level that the connectivity shifts from the west to east side of the negative polarity of NOAA AR 12967 as predicted by the MCT using ADAPT magnetograms.

## 7. Discussion and Conclusions

### 7.1. Characterization of Narrow Corridor and Its Associated Slow Wind Stream

SO tracked the decaying NOAA AR 12967 during its solar disk crossing and obtained both remote sensing and in situ observations. SO was also magnetically connected to NOAA AR 12967 and its vicinity for about 5 days. A distinguishing feature of the AR was a dark and narrow corridor observed in the EUV that expanded with height and temperature in the corona (Figures 1, 2). Plasma upflows of approximately 10 km s<sup>-1</sup> ran along the full extent of the corridor as observed by Hinode/EIS. NOAA AR 12967’s magnetic field structure has a classical pseudo-streamer configuration with the following negative polarity surrounded by positive field (its leading polarity counterpart in the southwest and remnant fields elsewhere, Figures 5–6). The open positive field in the west is spatially coincident with the plasma upflows of the narrow corridor, implying that the upflows are outflows (Figures 3–4, 9).

In situ measurements at SO show that the SW stream emanating from the narrow corridor had exceptionally slow radial velocities of  $\sim 300$  km s<sup>-1</sup> with correspondingly low proton temperatures of about 5 eV and increasingly high proton density from  $\sim 60$  to 300 cm<sup>-3</sup> in the time interval defined by boundaries B1–B2 (Figure 8). For approximately 1/3 days during this interval, the slow SW exhibited moderate Alfvénicity with the correlation coefficient between  $V$  and  $B$ ,  $|C_{vb}|$ , in the range [0.7, 1]. This period was preceded and followed by more variable and non-Alfvénic periods, especially closer to the B1 and B2 boundaries.

Alfvénic slow SW was first reported by Marsch et al. (1981) using Helios 2 measurements at 0.3 au. Subsequent in situ observations confirmed the presence of Alfvénic slow SW streams at  $\sim 0.3$  au (e.g., D’Amicis & Bruno 2015; Bale et al. 2019; Stansby et al. 2019; Perrone et al. 2020; Stansby et al.

2020) and at 1 au (e.g., D’Amicis et al. 2011, 2019; Wang & Ko 2019; D’Amicis et al. 2021). Alfvénicity is a key discriminator for classifying the slow SW and identifying its source regions, especially at close distances like that of SO at its perihelion (Stansby et al. 2019; D’Amicis et al. 2021). In this case, it was especially important in the absence of in situ composition data.

While proton temperatures and magnetic field flux densities are consistent with the average values measured by Helios within non-Alfvénic slow SW at a distance of 0.3–0.35 au, the velocity and density parameters detected at SO are extreme (see, e.g., Table 1 of Perrone et al. 2020). In fact, they are comparable to the values of the so-called very slow SW described by Sanchez-Diaz et al. (2016), where velocities  $<300 \text{ km s}^{-1}$  are associated with much higher and more variable densities in contrast with those reported by Perrone et al. (2020). The very slow SW was observed by Helios 1 and 2 at distances mainly  $<0.5$  au and not above 0.7 au. Sanchez-Diaz et al. (2016) proposed that the origin of the very slow SW may be small and isolated CHs, i.e., equatorial CHs, that become more widespread during active periods of the solar cycle.

Exceedingly high densities combined with very low velocities provide crucial information for further refining the source region–spacecraft connectivity during the B1–B2 interval. SW velocities of  $\sim 300 \text{ km s}^{-1}$  are more likely to originate at the southern end of the narrow corridor, where  $f(R_{ss})$  is an order of magnitude higher compared to that of the north (see Figures 4(c) and (d)). Furthermore, the southern end of the corridor is adjacent to the core of NOAA AR 12967. Typically, the core loops connecting opposite polarities have the highest density (e.g., Tripathi et al. 2008) and coronal FIP bias (e.g., Del Zanna & Mason 2014; Baker et al. 2015, 2021) values within the various loop populations of ARs. The in situ plasma during B1–B2 is more characteristic of the core loops in the south compared to the CH and quiet Sun found at the northern end of the corridor. Finally, the positive polarity of NOAA AR 12967 is a natural location for reconnection to take place between core loops and nearby open field (Baker et al. 2009), so that the plasma can be transferred from the closed core loops to nearby open field.

A comparison of the plasma parameters and magnetic field within the narrow corridor and its associated SW stream with that of equatorial CHs provides insight as to where they are located on the scale of CHs. Wang & Ko (2019) compiled the basic properties of 14 equatorial CHs and their associated slow wind streams ( $\sim 350\text{--}450 \text{ km s}^{-1}$ ) at 1 au from 2014 to 2017 (during solar maximum/beginning of decay phase of cycle 24). They found the following: (1) longitudinal widths  $\sim [3^\circ, 10^\circ]$ , (2) relatively high Alfvénicity  $|C_{vb}| \sim [0.6, 0.7]$  that tended to fall sharply approaching the edges or boundaries of the slow SW streams, (3) footpoint magnetic field strengths  $\sim [15, 30] \text{ G}$ , (4) proton temperatures  $\sim [1, 3] \text{ eV}$ , and (5) large expansion factors  $>9$ . These results, where applicable, correspond to those of D’Amicis et al. (2019), who identified equatorial CHs as the source regions of highly Alfvénic slow SW at 1 au.

The comparable values for key properties of the narrow corridor and its associated very slow SW are more extreme than those of equatorial CHs: the longitudinal width of the corridor in the 1.5 MK corona is up to 2 orders of magnitude smaller, the expansion factors are as much as an order of magnitude larger in the south, and the SW velocities are lower by  $\sim 100 \text{ km s}^{-1}$ . Alfvénicity in the B1–B2 interval possibly reaches higher levels

compared to the equatorial CHs in the study of Wang & Ko (2019); however, the extent of Alfvénicity appears to be shorter. The extremely high expansion factor related to this short period of Alfvénicity is expected, as there is low expansion only in the center of the flux tube compared to the edges. Closer to the edges, the field becomes more curved so that more of the Alfvén waves are reflected back rather than traveling outward along the open field.

## 7.2. Narrow Corridor and the S-web

The narrow corridor observed by Hinode/EIS has key features predicted by the S-web model. First, at typical coronal temperatures, it is a few orders of magnitude thinner in the longitudinal direction compared to typical equatorial CHs that were confirmed source regions of the slow SW. Second, the corridor is conjectured to be topologically robust (Antiochos et al. 2007, 2011). In this case, the dark corridor was observed to be highly stable for at least 7 days. Third, the corridor expands to a latitudinal width of  $60^\circ$  at the source surface ( $=2.5 R_\odot$ ). Fourth, the overall magnetic configuration is a pseudo-streamer formed of the embedded negative polarity of NOAA AR 12967 with the corridor formed within the surrounding positive field on the west. Open-field corridors are intrinsically associated with pseudo-streamers (Antiochos et al. 2011; Crooker et al. 2012; Scott et al. 2019).

Some of the extreme properties of the slow SW stream observed by SO are a direct consequence of the very nature of the S-web comprising open-field corridor source regions. The very slow SW recounted by Sanchez-Diaz et al. (2016) is plausibly produced by such regions. Extremely narrow corridors are characterized by super-radial expansion and therefore very slow wind velocities. Following on from this, Alfvénic slow SW is associated with overexpanded flux tubes (see D’Amicis et al. 2021, and references therein).  $|C_{vb}|$  was in the range  $[0.7, 1]$  for 1/3 days during the central part of the B1–B2 interval. Large variability in  $B_R$  and occasional reversal events, i.e., switchbacks, occurred while the very slow SW was moderately Alfvénic. Recent PSP and SO observations have established that switchbacks are linked to the Alfvénic slow wind from equatorial CHs (e.g., Bale et al. 2019; D’Amicis et al. 2021; Fedorov et al. 2021). It is reasonable that the induced interchange reconnection events along the S-web arcs (Higginson et al. 2017b) are a possible source of the switchbacks ever present in the inner heliosphere (D’Amicis et al. 2021, and references therein) and observed in the Alfvénic very slow SW emanating from the narrow open-field corridor.

The corridor associated with NOAA AR 12967 is not rare. On the contrary, the full synoptic HMI magnetogram shown in Figure 5(a) contains a number of possible S-web corridors identified by the yellow arrows. There is a spectrum of widths, mainly as a consequence of the  $B_R$  distribution in the photosphere. The narrowest of corridors will be beyond the spatial resolution of the magnetic models used to determine the open-field regions. Furthermore, limitations of EUV observations, e.g., instruments with limited FOV and projection effects (corridors are masked by nearby brighter regions), also contribute to the mistaken impression that they are not common features.

In conclusion, the combination of observations from SO and Hinode/EIS has made it possible to characterize the S-web narrow open-field corridor source region and its slow SW

stream. The intrinsic topology of the corridors means that super-radial expansion of the open field is likely to yield extremely slow SW velocities, with moderate Alfvénic content (along the central section of the corridor), and switchbacks. Other SW properties such as composition, charge state, and density are more likely to be governed by the surroundings of the corridors as interchange reconnection at the corridor boundaries opens up pathways for closed-field plasma of the nearby quiet Sun and ARs to reach the heliosphere.

## 8. Future Work

The present study is in line with previous studies, thereby improving our understanding of the link between solar and in situ observations as summarized in Section 1. Connectivity computations need to be pursued in a variety of other cases to test the statistical validity of the above conclusions. This can be achieved with the data currently accumulated by SO, PSP, SDO, Hinode, and ground-based observatories. For example, a natural extension of this work can be achieved during the 30-day perihelion passage of SO based on the study of R. D’Amicis et al. (2023, in preparation). PFSS extrapolation and magnetic field topology computations, as well as the outward tracing with a Parker spiral, adapted to the in situ measured velocity, are well-developed tools and sufficiently light in computational time to analyze a large amount of data. This method has known limitations that will need to be carefully tested in a large set of cases. Such tests that are worth conducting include how to further constrain where to set the source surface radius (Panasenco et al. 2020; Koukras et al. 2022) and how to better understand the open flux deficit (Linker et al. 2021). These can best be realized when the HCS is significantly warped so that the spacecraft successively scans different and identifiable source regions (e.g., of opposite magnetic polarity and/or of different spatial extents, thereby containing different SW types). Furthermore, the transit time from the corona to the spacecraft needs to be improved. One possibility is to use the statistical results of the wind profiles obtained by combining Helios and PSP observations (Dakeyo et al. 2022).

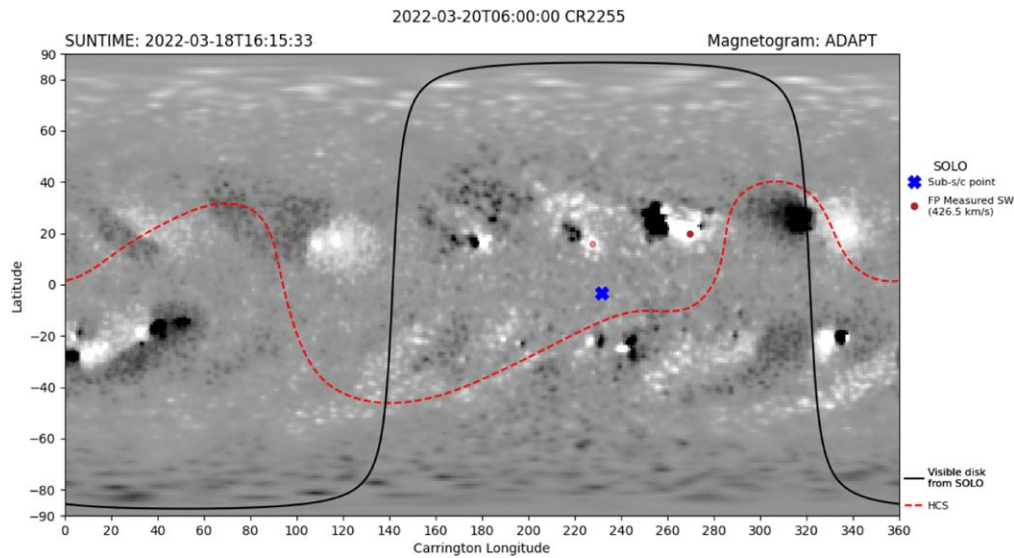
The precision of the deduced connectivities with various models and photospheric observations needs to be quantitatively evaluated in order to derive what is the best mapping function of the studied coronal configuration (Badman et al. 2022). However, the approach used in this study, as well as most other ones, involves a static approach to connectivity while the corona and the SW are known to be dynamic over a broad range of times (seconds to years); therefore, the above analysis is only providing the steady part of the connectivities, on a timescale of a few weeks. Further improvements call for the development of MHD numerical codes (coronal and interplanetary) to deal with the temporal evolution, e.g., the formation of stream interaction regions. This is a long-term effort that requires an improvement both in the boundary conditions (constrained by observations) and in the input physics (such as the required extra heating and SW acceleration; see, e.g., Perri et al. 2023, for more information). The already-impressive harvest of SO and PSP observations, together with the development of numerical tools, will allow us to deeply improve our knowledge of coronal and interplanetary physics, which are currently summarized in Rouillard et al. (2021) and Raouafi et al. (2023).

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## Appendix

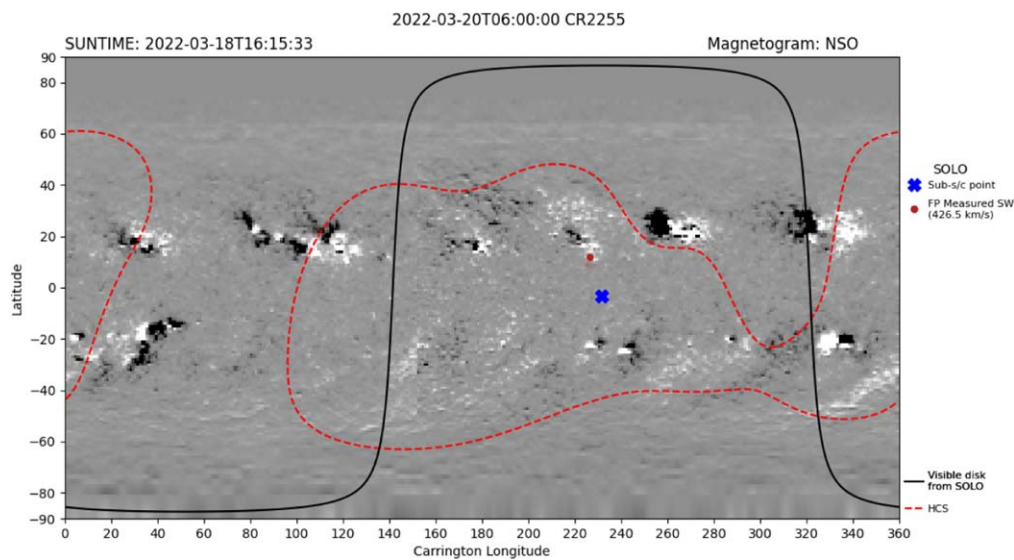
The SO MCT described in Section 6 identifies the possible source regions of the SW streams. Sample MCT ADAPT (Figure 10) and NSO (Figure 11) input magnetograms are shown below. The images are from animations included with the online versions of these figures. A comparison of the two movies shows differences between the potential magnetic field and the connectivities generated from different input magnetograms. The most striking feature is the dissimilarity in the





**Figure 10.** ADAPT connectivity maps generated using the SO MCT found at <http://connect-tool.irap.omp.eu/>. Spacecraft/Sun times are 2022 March 20 06:00 UT/2022 March 18 16:15 UT. Slow SW connectivities are indicated by the red circles. An animation of these data is available in the online Journal. The animation covers the period beginning on Spacecraft/Sun times of March 20 00:00 UT/2022 March 18 03:16 UT and ending on March 27 00:00 UT/2022 March 25 04:58 UT (duration  $\sim 2$  s).

(An animation of this figure is available.)



**Figure 11.** NSO connectivity maps generated using the SO MCT found at <http://connect-tool.irap.omp.eu/>. Spacecraft/Sun times are 2022 March 20 06:00 UT/2022 March 18 16:15 UT (bottom). Slow SW connectivities are indicated by the red circles. An animation of these data is available in the online Journal. The animation covers the period beginning on Spacecraft/Sun times of March 19 00:00 UT/2022 March 17 09:26 UT and ending on March 27 00:00 UT/2022 March 25 04:58 UT (duration  $\sim 2$  s).




















(An animation of this figure is available.)

heliospheric current sheets (HCS). The HCS in the ADAPT map runs from NW to SE, crossing the solar equator, whereas the one in the NSO map “wraps” around the northern (NOAA AR 12967 and NOAA AR 12971) and southern (NOAA AR 12668, NOAA AR 12670, and NOAA AR 12672) activity belts. Until early on 2022 March 24, SO connects to both NOAA AR 12965 and NOAA AR 12967, and then only to NOAA AR 12967. The connectivities from the two different MCT models are reasonably consistent, though there are slight timing differences. There is significant divergence in connectivities using the different input magnetograms after early

on 2022 March 24. In the ADAPT map SO is partially/fully connected to the positive polarity on the eastern side of the negative polarity of NOAA AR 12967, whereas in the NSO map SO is connected to magnetic field present in the southeast direction from the positive polarity of NOAA AR 12967 and across the solar equator.

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## References

- Abbo, L., Ofman, L., Antiochos, S. K., et al. 2016, *SSRv*, 201, 55  
 Antiochos, S. K., DeVore, C. R., Karpen, J. T., & Mikić, Z. 2007, *ApJ*, 671, 936  
 Antiochos, S. K., Mikić, Z., Titov, V. S., Lionello, R., & Linker, J. A. 2011, *ApJ*, 731, 112  
 Aslanyan, V., Pontin, D. I., Wyper, P. F., et al. 2021, *ApJ*, 909, 10  
 Badman, S. T., Brooks, D. H., Poirier, N., et al. 2022, *ApJ*, 932, 135  
 Baker, D., Brooks, D. H., Démoulin, P., et al. 2015, *ApJ*, 802, 104  
 Baker, D., Stangalini, M., Valori, G., et al. 2021, *ApJ*, 907, 16  
 Baker, D., van Driel-Gesztelyi, L., Mandrini, C. H., Démoulin, P., & Murray, M. J. 2009, *ApJ*, 705, 926  
 Bale, S. D., Badman, S. T., Bonnell, J. W., et al. 2019, *Natur*, 576, 237  
 Bavassano, B., Pietropaolo, E., & Bruno, R. 1998, *JGR*, 103, 6521  
 Brooks, D. H., Ugarte-Urra, I., & Warren, H. P. 2015, *NatCo*, 6, 5947  
 Brooks, D. H., & Warren, H. P. 2011, *ApJL*, 727, L13  
 Brooks, D. H., & Warren, H. P. 2016, *ApJ*, 820, 63  
 Chitta, L. P., Seaton, D. B., Downs, C., DeForest, C. E., & Higginson, A. K. 2023, *NatAs*, 7, 133  
 Cranmer, S. R. 2009, *LRSP*, 6, 3  
 Crooker, N. U., Antiochos, S. K., Zhao, X., & Neugebauer, M. 2012, *JGRA*, 117, A04104  
 Culhane, J. L., Harra, L. K., James, A. M., et al. 2007, *SoPh*, 243, 19  
 Dakeyo, J.-B., Maksimovic, M., Démoulin, P., Halekas, J., & Stevens, M. L. 2022, *ApJ*, 940, 130  
 D'Amicis, R., & Bruno, R. 2015, *ApJ*, 805, 84  
 D'Amicis, R., Bruno, R., & Bavassano, B. 2011, *JASTP*, 73, 653  
 D'Amicis, R., Matteini, L., & Bruno, R. 2019, *MNRAS*, 483, 4665  
 D'Amicis, R., Perrone, D., Bruno, R., & Velli, M. 2021, *JGRA*, 126, e28996  
 Del Zanna, G., & Mason, H. E. 2014, *A&A*, 565, A14  
 Démoulin, P., Hénoux, J. C., Priest, E. R., & Mandrini, C. H. 1996, *A&A*, 308, 643  
 Doschek, G. A., Warren, H. P., Mariska, J. T., et al. 2008, *ApJ*, 686, 1362  
 Elliott, H. A., Henney, C. J., McComas, D. J., Smith, C. W., & Vasquez, B. J. 2012, *JGRA*, 117, A09102  
 Elliott, H. A., McComas, D. J., Schwadron, N. A., et al. 2005, *JGRA*, 110, A04103  
 Fazakerley, A. N., Harra, L. K., & van Driel-Gesztelyi, L. 2016, *ApJ*, 823, 145  
 Fedorov, A., Louarn, P., Owen, C. J., et al. 2021, *A&A*, 656, A40  
 Fox, N. J., Velli, M. C., Bale, S. D., et al. 2016, *SSRv*, 204, 7  
 García Marirrodriga, C., Pacros, A., Strandmo, S., et al. 2021, *A&A*, 646, A121  
 Geiss, J., Gloeckler, G., & von Steiger, R. 1995, *SSRv*, 72, 49  
 Harra, L., Baker, D., Edwards, S. J., et al. 2015, *SoPh*, 290, 3203  
 Harra, L. K., Sakao, T., Mandrini, C. H., et al. 2008, *ApJL*, 676, L147  
 Hickmann, K. S., Godinez, H. C., Henney, C. J., & Arge, C. N. 2015, *SoPh*, 290, 1105  
 Higginson, A. K., Antiochos, S. K., DeVore, C. R., Wyper, P. F., & Zurbuchen, T. H. 2017a, *ApJ*, 837, 113  
 Higginson, A. K., Antiochos, S. K., DeVore, C. R., Wyper, P. F., & Zurbuchen, T. H. 2017b, *ApJL*, 840, L10  
 Horbury, T. S., Matteini, L., & Stansby, D. 2018, *MNRAS*, 478, 1980  
 Horbury, T. S., O'Brien, H., Carrasco Blazquez, I., et al. 2020, *A&A*, 642, A9  
 Koukras, A., Dolla, L., & Keppens, R. 2022, arXiv:2212.11553  
 Lemen, J. R., Title, A. M., Akin, D. J., et al. 2012, *SoPh*, 275, 17  
 Linker, J. A., Heinemann, S. G., Temmer, M., et al. 2021, *ApJ*, 918, 21  
 Lopez, R. E., & Freeman, J. W. 1986, *JGR*, 91, 1701  
 Maksimovic, M., Bale, S. D., Berčić, L., et al. 2020, *ApJS*, 246, 62  
 Marsch, E., Rosenbauer, H., Schwenn, R., Muehlhaeuser, K. H., & Denskat, K. U. 1981, *JGR*, 86, 9199  
 McComas, D. J., Bame, S. J., Barraclough, B. L., et al. 1998, *GeoRL*, 25, 1  
 Milano, L. J., Dmitruk, P., Mandrini, C. H., Gómez, D. O., & Démoulin, P. 1999, *ApJ*, 521, 889  
 Müller, D., St., Cyr, O. C., Zouganelis, I., et al. 2020, *A&A*, 642, A1  
 Neugebauer, M., & Goldstein, B. E. 2013, in AIP Conf. Ser. 1539, Solar Wind 13, ed. G. P. Zank et al. (Melville, NY: AIP), 46  
 Owen, C. J., Bruno, R., Livi, S., et al. 2020, *A&A*, 642, A16  
 Panasenco, O., Velli, M., D'Amicis, R., et al. 2020, *ApJS*, 246, 54  
 Perri, B., Kuźma, B., Brchneľova, M., et al. 2023, *ApJ*, 943, 124  
 Perrone, D., D'Amicis, R., De Marco, R., et al. 2020, *A&A*, 633, A166  
 Perrone, D., Stansby, D., Horbury, T. S., & Matteini, L. 2019, *MNRAS*, 488, 2380  
 Pesnell, W. D., Thompson, B. J., & Chamberlin, P. C. 2012, *SoPh*, 275, 3  
 Pinto, R. F., Brun, A. S., & Rouillard, A. P. 2016, *A&A*, 592, A65  
 Pinto, R. F., & Rouillard, A. P. 2017, *ApJ*, 838, 89  
 Raouafi, N. E., Matteini, L., Squire, J., et al. 2023, *SSRv*, 219, 8  
 Rochus, P., Auchère, F., Berghmans, D., et al. 2020, *A&A*, 642, A8  
 Rouillard, A. P., Pinto, R. F., Vourlidis, A., et al. 2020, *A&A*, 642, A2  
 Rouillard, A. P., Viall, N., Pierrard, V., et al. 2021, in Solar Physics and Solar Wind, ed. N. E. Raouafi & A. Vourlidis, Vol. 1 (Washington, DC: American Geophysical Union), 1  
 Sanchez-Diaz, E., Rouillard, A. P., Lavraud, B., et al. 2016, *JGRA*, 121, 2830  
 Scherrer, P. H., Schou, J., Bush, R. I., et al. 2012, *SoPh*, 275, 207  
 Schrijver, C. J., & De Rosa, M. L. 2003, *SoPh*, 212, 165  
 Scott, R. B., Pontin, D. I., & Wyper, P. F. 2019, *ApJ*, 882, 125  
 Scott, R. B., Pontin, D. I., Yeates, A. R., Wyper, P. F., & Higginson, A. K. 2018, *ApJ*, 869, 60  
 SPICE Consortium, Anderson, M., Appourchaux, T., et al. 2020, *A&A*, 642, A14  
 Stansby, D., Horbury, T. S., & Matteini, L. 2019, *MNRAS*, 482, 1706  
 Stansby, D., Matteini, L., Horbury, T. S., et al. 2020, *MNRAS*, 492, 39  
 Temmer, M. 2021, *LRSP*, 18, 4  
 Titov, V. S., Hornig, G., & Démoulin, P. 2002, *JGRA*, 107, 1164  
 Titov, V. S., Mikić, Z., Linker, J. A., Lionello, R., & Antiochos, S. K. 2011, *ApJ*, 731, 111  
 Tripathi, D., Mason, H. E., Young, P. R., & Del Zanna, G. 2008, *A&A*, 481, L53  
 Tu, C. Y., & Marsch, E. 1995, *SSRv*, 73, 1  
 Tu, C. Y., Marsch, E., & Thieme, K. M. 1989, *JGR*, 94, 11739  
 van Driel-Gesztelyi, L., & Green, L. M. 2015, *LRSP*, 12, 1  
 Wang, Y. M., & Ko, Y. K. 2019, *ApJ*, 880, 146  
 Wang, Y. M., Ko, Y. K., & Grappin, R. 2009, *ApJ*, 691, 760  
 Wang, Y. M., & Sheeley, N. R., Jr. 1990, *ApJ*, 365, 372  
 Wang, Y. M., Sheeley, N. R., Jr., & Rich, N. B. 2007, *ApJ*, 658, 1340  
 Yardley, S. L., Owen, C. J., Long, D. M., et al. 2023, arXiv:2304.09570  
 Zirker, J. B. 1977, *RvGSP*, 15, 257  
 Zouganelis, I., De Groof, A., Walsh, A. P., et al. 2020, *A&A*, 642, A3