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Asymptotic analysis of a class of functional equations and applications

by P. J. GRABNER†, H. PRODINGER AND R. F. TICHY†

ABSTRACT. Flajolet and Richmond have invented a method to solve a large class of divide-and-conquer recursions. The essential part of it is the asymptotic analysis of a certain generating function for $z \rightarrow \infty$ by means of the Mellin transform. In this paper this type of analysis is performed for a reasonably large class of generating functions fulfilling a functional equation with polynomial coefficients. As an application, the average life time of a party of N people is computed, where each person advances one step or dies with equal probabilities, and an additional “killer” can kill at any level up to d survivors, according to his probability distribution.

1. Introduction

In [3] Flajolet and Richmond presented an ingenious method to deal with a class of recursions where a typical example looks like

$$(1.1) \quad f_{n+d} = 1 + \sum_{k=0}^n 2^{-n} \binom{n}{k} (f_k + f_{n-k}).$$

Here, $d \geq 0$ is a fixed integer. For $d = 0$ and $d = 1$ such recursions (and their solutions!) occur frequently as *divide-and-conquer* recursions in the *Analysis of Algorithms*, (see [6] as a general reference on the subject. However, the new approach allows it for the first time to deal with the cases $d = 2, 3, \dots$; it consists of several stages. First, the recursion is translated into the language of exponential generating functions. Then, a related function (sometimes called the *Poisson transform*) is considered. Then, from its coefficients, an *ordinary generating function* $G(z)$ is built up. The analytic behaviour of the latter for $z \rightarrow \infty$ is then to be analyzed by *Mellin transform* methods, (this is the heart of the method). Then, by a substitution, the asymptotic behaviour as $z \rightarrow 1$ of the ordinary generating

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function of the desired numbers f_n is found. As the last step, this local information is translated to the asymptotic behaviour of the coefficients by *transfer theorems*.

The relevant function $G(z)$ fulfills a functional equation

$$(1.2) \quad G(z)(1+z)^d = 2z^d G\left(\frac{z}{2}\right) + P_0(z),$$

where the polynomial $P_0(z)$ depends on the starting values.

We devote this paper to the general study of recursions as in (1.2), with general polynomials as factors and “ $\frac{1}{2}$ ” replaced by an arbitrary parameter $0 < \lambda < 1$.

This is not only interesting by itself, but can be applied in the last section to a stochastic process which is a generalization of one presented in [10]. Although formulated in terms of “recreational mathematics” the recursions describe either average “trie” parameters or the behaviour of an “approximate counter”. In order to keep this paper short, we refrain from describing those algorithms and data structures and refer for all computer science algorithms to [6] and [7]. For the Mellin transform, which has proven to be very useful especially in number theory and in theoretical computer science, we refer to [1, chapter 13], [6], [8] and to the survey [4]. We note here that the classical applications of the Mellin transform occur in prime number theory and, more recently, in the analysis of digital problems (cf. [9] and for a detailed survey see [2]).

Let us recall the fundamental properties of the Mellin transform

$$(1.3) \quad f^*(s) = \int_0^\infty f(x)x^{s-1} dx.$$

If $f(x)$ is a piece-wise continuous function with

$$f(x) = O(x^\alpha) \quad \text{for } x \rightarrow 0 \quad \text{and} \quad f(x) = O(x^\beta) \quad \text{for } x \rightarrow \infty,$$

and $\alpha > \beta$, then the Mellin transform exists as a holomorphic function in the vertical strip $-\alpha < \Re s < -\beta$ (“fundamental strip”). Our basic tool is Mellin’s inversion formula

$$(1.4) \quad f(x) = \frac{1}{2\pi i} \int_{c-i\infty}^{c+i\infty} f^*(s)x^{-s} ds$$

for any $-\alpha < c < -\beta$. This formula yields a correspondence between the singularities of $f^*(s)$ and the asymptotics of $f(x)$ (cf. [8, chapter 1.5]). This technique is especially useful in the asymptotic analysis of *harmonic sums* $h(x) := \sum_k \lambda_k f(a_k x)$ since

$$(1.5) \quad h^*(s) = f^*(s) \sum_k \frac{\lambda_k}{a_k^s}.$$

Thus the Mellin transform of a harmonic sum is the product of a generalized Dirichlet series and the transform of the base function.

In Section 2 we introduce a class of functional equations with polynomial coefficients generalizing (1.2) and perform the asymptotic analysis by the Mellin transform approach. Two cases have to be distinguished: in the first one the main term does not contain an oscillating part (Section 2) and in the second one there are oscillations in the main term (Section 3). Section 4 is devoted to a direct extension of [3]. In the final section this method is applied to analyze a special evolution process where N persons are climbing up an infinite staircase as follows: at each step either a person goes to the next step or the person dies. We compute the average of the maximum lifetime.

2. A class of functional equations

We consider the functional equation

$$(2.1) \quad G(z)P_1(z) = G(\lambda z)P_2(z) + P_0(z),$$

where $P_0(z) = a_0 + a_1z + \dots + a_{d_0}z^{d_0}$, $P_1(z) = b_0 + b_1z + \dots + b_{d_1}z^{d_1}$ and $P_2(z) = c_0 + c_1z + \dots + c_{d_2}z^{d_2}$ are polynomials of degrees d_0, d_1, d_2 (with $d_1 \leq d_2, d_0 \geq 0$), respectively, assume that P_1 and P_2 have no non-negative real roots, $P_1(0) \neq P_2(0)$ and $0 < \lambda < 1$. Our main result gives an asymptotic expansion for the unique analytic solution of (2.1). Formally iterating the functional equation we get

$$G(z) = \sum_{n=0}^{\infty} P_0(\lambda^n z) \frac{P_2(z) \cdots P_2(\lambda^{n-1} z)}{P_1(z) \cdots P_1(\lambda^n z)}.$$

Let now

$$Q_1(z) = \prod_{n=0}^{\infty} \frac{1}{b_0} P_1(\lambda^n z) \quad \text{and} \quad Q_2(z) = \prod_{n=0}^{\infty} \frac{1}{c_0} P_2(\lambda^n z).$$

Thus we can write

$$(2.2) \quad G(z) = \frac{Q_2(z)}{b_0 Q_1(z)} \sum_{n=0}^{\infty} P_0(\lambda^n z) \left(\frac{c_0}{b_0}\right)^n \frac{Q_1(\lambda^{n+1} z)}{Q_2(\lambda^n z)}.$$

Our next step is to establish an asymptotic expansion of $G(z)$ for $z \rightarrow \infty$. This will be done by a Mellin-transform technique. Let for the following

$$(2.3) \quad \frac{1}{b_0} P_1(z) = \prod_{k=1}^{r_1} (1 + \alpha_k z)^{\mu_k} \quad \text{and} \quad \frac{1}{c_0} P_2(z) = \prod_{k=1}^{r_2} (1 + \beta_k z)^{\nu_k}$$

(μ_k and ν_k denoting the multiplicities of the roots) and take the Mellin-transforms of the logarithms of these equations

$$\begin{aligned} \left(\log \frac{1}{b_0} P_1\right)^*(s) &= \frac{\pi}{s \sin \pi s} \sum_{k=1}^{r_1} \mu_k \alpha_k^{-s} \quad \text{and} \\ \left(\log \frac{1}{c_0} P_2\right)^*(s) &= \frac{\pi}{s \sin \pi s} \sum_{k=1}^{r_2} \nu_k \beta_k^{-s} \end{aligned}$$

for $-1 < \Re s < 0$. Observing that $\log Q_1$ and $\log Q_2$ are harmonic sums yields

$$(2.4) \quad \begin{aligned} (\log Q_1)^*(s) &= \frac{1}{1 - \lambda^{-s}} \frac{\pi}{s \sin \pi s} \sum_{k=1}^{r_1} \mu_k \alpha_k^{-s} \quad \text{and} \\ (\log Q_2)^*(s) &= \frac{1}{1 - \lambda^{-s}} \frac{\pi}{s \sin \pi s} \sum_{k=1}^{r_2} \nu_k \beta_k^{-s}. \end{aligned}$$

In order to establish an asymptotic expansion for Q_1 and Q_2 we apply Mellin’s inversion formula

$$(2.5) \quad \log Q_1(z) = \frac{1}{2\pi i} \int_{-\frac{1}{2}-i\infty}^{-\frac{1}{2}+i\infty} (\log Q_1)^*(s) z^{-s} ds$$

(and similarly for Q_2). Inserting (2.4) and shifting the line of integration to $\Re s = \frac{1}{2}$ (cf. [8, chapter 1.5]) yields

(2.6) $\log Q_1(z) =$

$$\sum_{m \in \mathbb{Z}, s = \chi_m} \operatorname{Res} \frac{1}{\lambda^{-s} - 1} \frac{\pi}{s \sin \pi s} \sum_{k=1}^{r_1} \mu_k \alpha_k^{-s} z^{-s} + \frac{1}{2\pi i} \int_{\frac{1}{2} - i\infty}^{\frac{1}{2} + i\infty} (\log Q_1)^*(s) z^{-s} ds,$$

where $\chi_m = \frac{2m\pi i}{\log \lambda}$ are the roots of $1 - \lambda^{-s} = 0$. Note that for $s = 0$ we have a triple pole and in the other cases simple poles. Computing residues (and applying Vieta's rule) we get

$$\begin{aligned} \operatorname{Res}_{s=0} \frac{1}{1 - \lambda^{-s}} \frac{\pi}{s \sin \pi s} \sum_{k=1}^{r_1} \mu_k \alpha_k^{-s} z^{-s} = \\ - \frac{d_1}{2 \log \frac{1}{\lambda}} \log^2 z - \left(\frac{\log \frac{b_0}{b_{d_1}}}{\log \frac{1}{\lambda}} - \frac{d_1}{2} \right) \log z - B, \end{aligned}$$

where

$$B = \frac{1}{2 \log \frac{1}{\lambda}} \sum_{k=1}^{r_1} \mu_k \log^2 \alpha_k - \frac{1}{2} \log \frac{b_0}{b_{d_1}} - \frac{\log \frac{1}{\lambda}}{12} - \frac{\pi^2}{6 \log \frac{1}{\lambda}}.$$

For the residues at $\chi_m \neq 0$ we obtain

$$\begin{aligned} \operatorname{Res}_{s=\chi_m} \frac{1}{1 - \lambda^{-s}} \frac{\pi}{s \sin \pi s} \sum_{k=1}^{r_1} \mu_k \alpha_k^{-s} z^{-s} = \\ \frac{\pi}{\chi_m (\log \lambda) \sin \pi \chi_m} \sum_{k=1}^{r_1} \mu_k \alpha_k^{-\chi_m} z^{-\chi_m}. \end{aligned}$$

Thus we have

$$\log Q_1(z) = \frac{d_1}{2 \log \frac{1}{\lambda}} \log^2 z + \left(\frac{\log \frac{b_0}{b_{d_1}}}{\log \frac{1}{\lambda}} - \frac{d_1}{2} \right) \log z + B + \varphi_1 \left(\frac{\log z}{\log \frac{1}{\lambda}} \right) + O \left(z^{-\frac{1}{2}} \right),$$

where $\varphi_1(u)$ is a continuous periodic function of period 1, whose Fourier expansion is given by

(2.7) $\varphi_1(u) = - \sum_{m \in \mathbb{Z} \setminus \{0\}} \frac{\pi}{\chi_m (\log \lambda) \sin \pi \chi_m} \sum_{k=1}^{r_1} \mu_k \alpha_k^{\chi_m} e^{2\pi i m u}.$

Similar calculations yield the asymptotic expansion for Q_2 :

$$\log Q_2(z) = \frac{d_2}{2 \log \frac{1}{\lambda}} \log^2 z + \left(\frac{\log \frac{c_0}{c_{d_2}}}{\log \frac{1}{\lambda}} + \frac{d_2}{2} \right) \log z + C + \varphi_2 \left(\frac{\log z}{\log \frac{1}{\lambda}} \right) + O \left(z^{-\frac{1}{2}} \right),$$

where

$$C = \frac{1}{2 \log \frac{1}{\lambda}} \sum_{k=1}^{r_2} \nu_k \log^2 \beta_k - \frac{1}{2} \log \frac{c_0}{c_{d_2}} - \frac{\log \frac{1}{\lambda}}{12} - \frac{\pi^2}{6 \log \frac{1}{\lambda}}$$

and $\varphi_2(u)$ is a continuous periodic function of period 1, whose Fourier expansion is given by

$$(2.8) \quad \varphi_2(u) = - \sum_{m \in \mathbb{Z} \setminus \{0\}} \frac{\pi}{\chi_m(\log \lambda) \sin \pi \chi_m} \sum_{k=1}^{r_2} \nu_k \beta_k^{\chi_m} e^{2\pi i m u}.$$

Setting $\Phi(u) = \exp(C - B + \varphi_2(u) - \varphi_1(u))$, $A = \frac{d_2 - d_1}{2 \log \frac{1}{\lambda}}$, $L = \frac{\log \frac{c_0 b_{d_1}}{c_{d_2} b_0}}{\log \frac{1}{\lambda}}$ and observing the rate of growth of the transforms in (2.4), we derive

$$(2.9) \quad \frac{Q_2(z)}{Q_1(z)} = \begin{cases} z^L \Phi \left(\frac{\log z}{\log \frac{1}{\lambda}} \right) \left(1 + O(z^{-\frac{1}{2}}) \right) & \text{if } d_1 = d_2 \\ e^{A \log^2 z} z^{L + \frac{d_2 - d_1}{2}} \Phi \left(\frac{\log z}{\log \frac{1}{\lambda}} \right) \left(1 + O(z^{-\frac{1}{2}}) \right) & \text{if } d_1 < d_2 \end{cases}$$

for $z \rightarrow \infty$ inside an angular region

$$|\arg z| \leq \delta < \delta_0 = \pi - \min_k (|\arg \alpha_k|, |\arg \beta_k|).$$

As above we obtain

$$(2.10) \quad \frac{Q_1(\lambda z)}{Q_2(z)} = \begin{cases} \frac{b_{d_1}}{b_0 \lambda^{d_1}} z^{-L - d_1} \frac{1}{\Phi \left(\frac{\log z}{\log \frac{1}{\lambda}} \right)} \left(1 + O(z^{-\frac{1}{2}}) \right) & \text{if } d_1 = d_2 \\ e^{-A \log^2 z + O(\log z)} & \text{if } d_1 < d_2 \end{cases}$$

for $|\arg z| \leq \delta < \delta_0$.

Let us first consider the case $d_1 = d_2 = d$ and $a_0 \neq 0$. Then we need $|c_0| < |b_0|$ to ensure convergence in (2.2). Setting

$$(2.11) \quad H(z) = \sum_{n=0}^{\infty} P_0(\lambda^n z) \left(\frac{c_0}{b_0}\right)^n \frac{Q_1(\lambda^{n+1} z)}{Q_2(\lambda^n z)},$$

and observing that this is a harmonic sum yields for its Mellin transform

$$(2.12) \quad H^*(s) = \frac{b_0}{b_0 - c_0 \lambda^{-s}} \left(P_0(z) \frac{Q_1(\lambda z)}{Q_2(z)} \right)^* (s).$$

Notice that this Mellin transform only exists if $-L + d_0 - d < 0$, since “the order at ∞ ” has to be smaller than “the order at 0”. From our asymptotic expansion (2.10) we want to extract some information on the singularities of the transformed harmonic sum $H^*(s)$. The fundamental strip of this transform is

$$0 < \Re s < -\max \left(-L + d_0 - d, \frac{\log \frac{c_0}{b_0}}{\log \frac{1}{\lambda}} \right).$$

Thus we have to distinguish two cases. First we consider the case occurring in our applications in Section 3, i.e.

$$(2.13) \quad b_d \lambda^{d_0} \leq c_d \lambda^d.$$

This is the case where the singularities originate directly from the individual summands (in the other case the singularities are caused by the harmonic summation and a special example of this type was extensively studied by Flajolet and Richmond [3]). For $b_d \lambda^{d_0} < c_d \lambda^d$, by (2.10) we have

$$(2.14) \quad H(z) = \frac{b_d^2}{(b_d - c_d \lambda^{d-d_0}) b_0 \lambda^d} z^{-L-d+d_0} \frac{1}{\Phi \left(\frac{\log z}{\log \frac{1}{\lambda}} \right)} (1 + O(z^{-\varepsilon})),$$

where ε is a suitable positive number. Combining (2.14) and (2.9) the fluctuation Φ cancels out in the asymptotics of $G(z)$ and we obtain

$$(2.15) \quad G(z) = \frac{b_d^2 a_{d_0}}{b_0^2 (b_d - c_d \lambda^{d-d_0}) \lambda^d} z^{d_0-d} (1 + O(z^{-\varepsilon})).$$

For $b_d \lambda^{d_0} = c_d \lambda^d$ the function $H^*(s)$ has double poles at $-L-d+d_0 + \chi_m$. Thus the asymptotic formula (2.14) has to be replaced by

$$(2.16) \quad H(z) = \left(\frac{\log z}{\log \frac{1}{\lambda}} \frac{1}{\Phi \left(\frac{\log z}{\log \frac{1}{\lambda}} \right)} + \Psi \left(\frac{\log z}{\log \frac{1}{\lambda}} \right) \right) z^{-L-d+d_0} (1 + O(z^{-\varepsilon})).$$

Combining (2.16) and (2.9) we derive

$$(2.17) \quad G(z) = \left(\frac{1}{b_0} \frac{\log z}{\log \frac{1}{\lambda}} + \frac{1}{b_0} \Psi \left(\frac{\log z}{\log \frac{1}{\lambda}} \right) \right) z^{d_0-d} (1 + O(z^{-\varepsilon})),$$

where Ψ is an unspecified periodic fluctuation.

Until now we have restricted ourselves to the case $a_0 \neq 0$. In general let ℓ be defined by $P_0(z) = z^\ell \tilde{P}_0(z)$ with a polynomial $\tilde{P}_0(z)$ such that $\tilde{P}_0(0) \neq 0$. Then the equation corresponding to (2.11) is

$$(2.18) \quad H(z) = z^\ell \sum_{n=0}^{\infty} \tilde{P}_0(\lambda^n z) \left(\frac{\lambda^\ell c_0}{b_0} \right)^n \frac{Q_1(\lambda^{n+1} z)}{Q_2(\lambda^n z)}.$$

The sum in this expression is exactly of the type considered above and therefore we need the condition $\lambda^\ell |c_0| < |b_0|$ to ensure convergence. Thus we obtain the same asymptotic expansions also in the general case $\ell \geq 0$.

Summing up what we have proved until now yields

THEOREM 1. *Let $G(z)$ be the unique analytic solution of a functional equation of type (2.1) and let ℓ be the smallest number such that $a_\ell \neq 0$. Furthermore assume the additional properties $d_1 = d_2 = d$, $\lambda^\ell |c_0| < |b_0|$ and $b_d \lambda^{d_0} \leq c_d \lambda^d$. Then $G(z)$ has an asymptotic expansion for $z \rightarrow \infty$, ($|\arg z| \leq \delta < \pi - \min_k(|\arg \alpha_k|, |\arg \beta_k|)$), of the form*

$$G(z) = \begin{cases} \frac{b_d^2 a_{d_0}}{b_0^2 (b_d - c_d \lambda^{d-d_0}) \lambda^d} z^{d_0-d} (1 + O(z^{-\varepsilon})) & \text{if } b_d \lambda^{d_0} < c_d \lambda^d \\ \left(\frac{1}{b_0} \frac{\log z}{\log \frac{1}{\lambda}} + \frac{1}{b_0} \Psi \left(\frac{\log z}{\log \frac{1}{\lambda}} \right) \right) z^{d_0-d} (1 + O(z^{-\varepsilon})) & \text{if } b_d \lambda^{d_0} = c_d \lambda^d, \end{cases}$$

where ε is a suitable positive number and Ψ denotes a continuous periodic function of period 1.

Remark 1: We note that in the case covered by Theorem 1 the periodic fluctuations do not occur in the main term.

3. Fluctuations in the main term

First we consider functional equations of type (2.1) with $d_1 = d_2 = d$ satisfying the condition $b_d \lambda^{d_0} > c_d \lambda^d$ converse to (2.13). In the previous section the convergence of the Fourier series for φ_1 and φ_2 was immediate. In the case treated now we establish a preparatory lemma generalizing Lemma 5 in [3], from which the convergence of the corresponding Fourier series can be deduced. In the following let v denote the maximal order of a zero of the entire function Q_2 . Note that v is finite, but may be larger than $\max_k \nu_k$.

LEMMA 1. Let $F(z) = P_0(z) \frac{Q_1(\lambda z)}{Q_2(z)}$. Then the Mellin transform $F^*(s)$ admits the representation

$$\frac{\pi}{\sin \pi s} (E_0(\lambda^{-s}) + (s - 1)E_1(\lambda^{-s}) + \dots + (s - 1) \cdots (s - v + 1)E_{v-1}(\lambda^{-s}))$$

where the E_i 's are entire functions. Thus for $\sigma > 0$

$$(3.1) \quad |F^*(\sigma + it)| = O(|t|^{v-1} e^{-\pi|t|}) \quad \text{as } |t| \rightarrow \infty.$$

Proof. Proceeding as in [3] let

$$J(s) = \int_{\mathcal{H}} F(z)(-z)^{s-1} dz,$$

where \mathcal{H} denotes a Hankel contour that goes from $+\infty - i0$, circles around 0 clockwise and returns to $+\infty + i0$. Then a standard argument (cf. [11]) yields $J(s) = 2i \sin \pi s F^*(s)$. Now we want to evaluate the loop integral by residues. For this purpose we have to know polynomial upper bounds for $F(z)$, which follow from

$$(3.2) \quad \left| \prod_{n=0}^{\infty} \left(\frac{1 + \alpha \lambda^n z}{1 + \beta \lambda^n z} \right) \right| \leq K_1 |z|^{K_2} \quad \text{for } |1 + \beta \lambda^n z| \geq \varepsilon,$$

where α, β and $\frac{\lambda}{4} > \varepsilon > 0$ are given numbers and K_1 and K_2 are positive constants depending on those numbers and on λ .

For proving this assertion let $\gamma = \max\{1, 2|\alpha|, 2|\beta|\}$ and $n_0 = \left\lfloor \frac{\log 2\gamma z}{\log \frac{\lambda}{4}} \right\rfloor + 1$. Thus we have $|\lambda^{n_0} z \gamma| \leq \frac{1}{2}$. We split the product (3.2) into two parts $n < n_0$ and $n \geq n_0$. The second product can be estimated by taking logarithms

$$\begin{aligned} & \left| \sum_{n=n_0}^{\infty} \log \left(\frac{1 + \alpha \lambda^n z}{1 + \beta \lambda^n z} \right) \right| \\ &= \left| \sum_{n=n_0}^{\infty} \left((\alpha - \beta) z \lambda^n - (\alpha^2 - \beta^2) \frac{z^2}{2} \lambda^{2n} + (\alpha^3 - \beta^3) \frac{z^3}{3} \lambda^{3n} \pm \dots \right) \right| \\ &\leq \gamma |z| \lambda^{n_0} \frac{1}{1 - \lambda} + \gamma^2 |z|^2 \lambda^{2n_0} \frac{1}{1 - \lambda^2} + \dots \leq \frac{1}{1 - \lambda}. \end{aligned}$$

For estimating the first part of the product in the region $|1 + \beta\lambda^n z| \geq \varepsilon$ for $n = 0, \dots, n_0 - 1$ we note that the function $f(w) = \frac{1+\alpha w}{1+\beta w}$ is continuous on the compact subregion $|1 + \beta w| \geq \varepsilon$ of the Riemannian sphere, and hence attains its maximum $M = M(\varepsilon, \alpha, \beta, \lambda) \geq 1$. Thus we obtain

$$\left| \prod_{n=0}^{\infty} \left(\frac{1 + \alpha\lambda^n z}{1 + \beta\lambda^n z} \right) \right| \leq M^{n_0} e^{\frac{1}{1-\lambda}} \leq K_1 |z|^{K_2},$$

and (3.2) is proved.

Observing that $\frac{Q_1}{Q_2}$ can be decomposed into finitely many factors of the type (3.2) we get a polynomial upper bound for $F(z)$ in the region

$$|1 + \beta_k \lambda^n z| \geq \varepsilon, \quad (k = 1, \dots, r_2 \text{ and } n = 0, 1, \dots).$$

Notice that ε is chosen small enough that arbitrarily large circles centered at the origin are contained in this region. Again by a standard argument we derive

$$J(s) = 2\pi i \sum_{k=1}^{r_2} \sum_{n=0}^{\infty} \operatorname{Res}_{z=-\frac{\lambda^{-n}}{\beta_k}} F(z)(-z)^{s-1}$$

for $\Re s < \sigma_0 < 0$. Next we calculate the residues and observe that the maximum order of the poles is v . Let us consider a pole $-\frac{\lambda^{-n}}{\beta_k}$ of multiplicity b and set

$$X_{k,n}(z) = F(z)(1 + \beta_k \lambda^n z)^b.$$

First we have for $w = z + \frac{\lambda^{-n}}{\beta_k}$

$$(3.3) \quad (-z)^{s-1} = \lambda^{-n(s-1)} \beta_k^{-s+1} \left(1 + \frac{(1-s)}{1!} (\beta_k \lambda^n w) + \frac{(1-s)(2-s)}{2!} (\beta_k \lambda^n w)^2 + \dots \right).$$

Expanding $X_{k,n}(-\frac{\lambda^{-n}}{\beta_k} + w)$ around $w = 0$ we get

$$\begin{aligned} X_{k,n} \left(-\frac{\lambda^{-n}}{\beta_k} + w \right) &= P_0(z) Q_1(\lambda z) \prod \left(1 - \frac{\beta_l}{\beta_k} \lambda^{m-n} \right)^{-1} \prod \left(1 + \frac{w}{\frac{\lambda^{-m}}{\beta_l} - \frac{\lambda^{-n}}{\beta_k}} \right)^{-1} \\ &= \sum_{j=0}^{\infty} \Lambda_j(k, n) w^j, \end{aligned}$$

where the products range over all (l, m) such that $\beta_l \lambda^m \neq \beta_k \lambda^n$. We observe

$$(3.4) \quad \Lambda_j(k, n) = O\left(\lambda^{\frac{n(n+1)}{2} + j^n} \eta_k^{-j}\right)$$

with an absolute O -constant, where η_k is the minimal distance of $\frac{1}{\beta_k}$ to the other roots of Q_2 .

Multiplying (3.3) and the expansion of $X_{k,n}$ yields for the coefficient of w^{b-1}

$$\lambda^{-n(s-1)} \beta_k^{-s+1} \sum_{j=0}^{b-1} \Lambda_{b-j-1}(k, n) \frac{(1-s)(2-s)\cdots(j-s)}{j!}.$$

Each residue provides one term in the power series expansion of the $E_i(\lambda^{-s})$. Because of the upper bound (3.4) the power series expansion of E_i is hyperexponentially convergent and E_i is an entire function. \square

Now we are able to treat functional equations of type (2.1) with equal degrees and fluctuations in the main term.

THEOREM 2. *Let $G(z)$ be the unique analytic solution of a functional equation of type (2.1) and let ℓ be the smallest number such that $a_\ell \neq 0$. Furthermore assume the additional properties $d_1 = d_2 = d$, $\lambda^\ell |c_0| < |b_0|$ and $b_d \lambda^{d_0} > c_d \lambda^d$. Then $G(z)$ has an asymptotic expansion for $z \rightarrow \infty$, ($|\arg z| \leq \delta < \pi - \min_k(|\arg \alpha_k|, |\arg \beta_k|)$), of the form*

$$G(z) = z^M \Psi_1\left(\frac{\log z}{\log \frac{1}{\lambda}}\right) (1 + O(z^{-\varepsilon})),$$

where ε is a suitable positive number, $M = \frac{\log \frac{b_d}{c_d}}{\log \frac{1}{\lambda}}$ and Ψ_1 some continuous periodic fluctuation (the Fourier expansion of which is discussed in the proof).

Proof. We apply Mellin's inversion formula to (2.12) and observe that the first singularities encountered when shifting the line of integration to the right originate from the factor $\frac{b_0}{b_0 - c_0 \lambda^{-s}}$. Notice that evaluations of the E_i 's of Lemma 1 at the poles are constants and that the Fourier expansion of the arising fluctuation is exponentially convergent because of the estimate (3.1). The fluctuation Ψ_1 is obtained by multiplying with the periodic

function $\frac{1}{b_0}\Phi$ as in (2.9). The proof is completed by multiplying with the first case in the asymptotic formula (2.9). \square

Up to now we have only considered polynomials P_1 and P_2 of equal degree. What remains for a complete asymptotic analysis of functional equations of type (2.1) is the case $d_1 < d_2$, (we note here that the converse case $d_1 > d_2$ cannot be treated by the Mellin transform approach).

Remark 2: Lemma 1 remains true in the case $d_1 < d_2$. The only problem is to find a polynomial upper bound for $\frac{Q_1(\lambda z)}{Q_2(z)}$. This product can be split into products of type (3.2) and $d_2 - d_1$ additional linear factors in the denominator. For these additional factors we easily find the polynomial lower bound $K\varepsilon^{(d_2-d_1)\log z}$.

By the same reasoning as above we obtain

THEOREM 3. *Let $G(z)$ be the unique analytic solution of a functional equation of type (2.1) with $d_1 < d_2$. Then $G(z)$ has an asymptotic expansion for $z \rightarrow \infty$, ($|\arg z| \leq \delta < \pi - \min_k(|\arg \alpha_k|, |\arg \beta_k|)$), of the form*

$$G(z) = e^{A \log^2 z} z^{M + \frac{d_2 - d_1}{2}} \Psi_2 \left(\frac{\log z}{\log \frac{1}{\lambda}} \right) (1 + O(z^{-\varepsilon})),$$

where ε is a suitable positive number, $A = \frac{d_2 - d_1}{2 \log \frac{1}{\lambda}}$, $M = \frac{\log \frac{b_d}{c_d}}{\log \frac{1}{\lambda}}$ and Ψ_2 some continuous periodic fluctuation.

4. The Flajolet–Richmond case

In [3] Flajolet and Richmond analyze a tree-partitioning process in which n elements split into d at the root of a tree (d is a design parameter), the rest going recursively into two subtrees with a binomial probability distribution. They prove an asymptotic formula for the expected number f_n of non-empty nodes in such a random tree, by studying the generating function $G(z) = \sum_{n \geq 0} g_n z^n$ where $f_n = \sum_{k=0}^n \binom{n}{k} g_k$. $G(z)$ satisfies the functional equation

$$(4.1) \quad G(z)(1+z)^d = 2z^d G\left(\frac{z}{2}\right) + P_0(z),$$

where P_0 is a polynomial of degree d . This is not a functional equation of type (2.1) since $c_0 = 0$. Thus the above theorems cannot be applied. In the following we extend our general approach to the case $c_0 = 0$, $d_1 = d_2 = d$ which covers the special case (4.1) (in the case $b_0 = 0$ there is no analytic solution, for $d_1 \neq d_2$ the Mellin transform approach cannot be applied).

THEOREM 4. *Let $G(z)$ be the unique analytic solution of the functional equation*

$$G(z)P_1(z) = G(\lambda z)P_2(z) + P_0(z) \quad (0 < \lambda < 1),$$

where $P_1(z) = b_0 + b_1z + \dots + b_dz^d$, $P_2(z) = c_\kappa z^\kappa + \dots + c_dz^d$ (with $b_0 \neq 0$, $c_\kappa \neq 0$, $\kappa > 0$) are polynomials of equal degree d with no positive real roots and let $P_0(z) = a_0 + a_1z + \dots + a_{d_0}z^{d_0} \neq 0$. Then $G(z)$ has an asymptotic expansion for $z \rightarrow \infty$, ($|\arg z| \leq \delta < \pi - \min_k(|\arg \alpha_k|, |\arg \beta_k|)$), of the form

$$G(z) = \begin{cases} z^{M+d_0-d} \Psi_3 \left(\frac{\log z}{\log \frac{1}{\lambda}} \right) (1 + O(z^{-\varepsilon})) & \text{if } M < d - d_0 \\ \left(\frac{a_d \log z}{c_d \log \frac{1}{\lambda}} + \Psi_4 \left(\frac{\log z}{\log \frac{1}{\lambda}} \right) \right) z^{d_0-d} (1 + O(z^{-\varepsilon})) & \text{if } M = d - d_0 \\ \frac{a_{d_0}}{c_d} z^{d_0-d} (1 + O(z^{-\varepsilon})) & \text{if } M > d - d_0, \end{cases}$$

where ε is a suitable positive number, $M = \frac{\log \frac{b_d}{c_d}}{\log \frac{1}{\lambda}}$ and Ψ_3, Ψ_4 some continuous periodic fluctuations.

Proof. Using the substitution $t = \frac{1}{z}$ we can rewrite the explicit formula for $G(z)$

$$G\left(\frac{1}{t}\right) = t^{d-d_0} \sum_{n=0}^{\infty} \bar{P}_0(\lambda^{-n}t) \lambda^{n(d_0-d)} \frac{\bar{P}_2(t) \cdots \bar{P}_2(\lambda^{-(n-1)}t)}{\bar{P}_1(t) \cdots \bar{P}_1(\lambda^{-n}t)},$$

where $\bar{P}_0(z) = z^{d_0}P_0\left(\frac{1}{z}\right)$, $\bar{P}_1(z) = z^dP_1\left(\frac{1}{z}\right)$ and $\bar{P}_2(z) = z^dP_2\left(\frac{1}{z}\right)$ are polynomials with $\deg \bar{P}_1 = d > d - \kappa = \deg \bar{P}_2$. Proceeding as in Section 2 we set

$$\bar{Q}_1(t) = \prod_{n=0}^{\infty} \frac{1}{b_d} \bar{P}_1(\lambda^n t)$$

$$\bar{Q}_2(t) = \prod_{n=0}^{\infty} \frac{1}{c_d} \bar{P}_2(\lambda^n t)$$

and obtain

$$(4.2) \quad G\left(\frac{1}{t}\right) = t^{d-d_0} \frac{\bar{Q}_1(\lambda t)}{c_d \bar{Q}_2(\lambda t)} \sum_{n=0}^{\infty} \bar{P}_0(\lambda^{-n}t) \lambda^{n(d_0-d)} \frac{\bar{Q}_2(\lambda^{-(n-1)}t)}{\bar{Q}_1(\lambda^{-n}t)}.$$

Let $F(t) = \bar{P}_0(t) \frac{\bar{Q}_2(\lambda t)}{Q_1(t)}$. Notice that Lemma 1 can be applied to this function. Furthermore we have, for $t \rightarrow \infty$, $F(t) \leq \exp(-\eta \log^2 t + O(\log z))$ with some constant $\eta > 0$ and $F(0) = a_{d_0}$. Thus the fundamental strip of $F^*(s)$ is $0 < \Re s < \infty$. Denoting the sum in (4.2) by $H(t)$ we obtain

$$(4.3) \quad H^*(s) = F^*(s) \frac{b_d}{b_d - c_d \lambda^{d_0-d+s}}$$

with the fundamental strip $\max(0, -M - d + d_0) < \Re s < \infty$. This yields the desired asymptotic expansion for $H(t)$, $t \rightarrow 0$ which implies the expansion for $G(z)$, $z \rightarrow \infty$. \square

Remark 3: The computer science problem originally considered by Flajolet and Richmond is covered by the first alternative of Theorem 4.

5. An evolution process with killer

We consider the following stochastic process. N persons (starting at level 1) are climbing up an infinite staircase. At every step for each person there are two possibilities of equal probability: (i) go to the next step (ii) the person dies. After this at any level an outside killer kills j persons with given probability p_j , ($0 \leq j \leq d$, d is fixed; $p_0 \neq 0$), (if there are only $< j$ persons available, all of them are killed). This situation is a generalization of [10] where the case $d = 1$ was discussed (emerging from some computer science problems).

We ask for the expectation C_N of the maximum lifetime. To this end, we introduce the probability generating functions $F_N(z)$ where the coefficient of z^k is the probability that the maximum level of N persons is k . We have the recursion ($N \geq 1$)

$$(5.1) \quad F_N(z) = z \sum_{k=1}^N 2^{-N} \binom{N}{k} [p_0 F_k(z) + \dots + p_{k-d} F_{k-d}] + z 2^{-N},$$

(observe that $2^{-N} \binom{N}{k}$ is the probability that k out of N people have made it to the next level. The variable z marks the level, and the effect of the killer is described by $[p_0 F_k(z) + \dots + p_{k-d} F_{k-d}]$). The cumbersome special cases when there are less persons than the killer wants to kill can be avoided by setting $F_i(z) = z$ for $i \leq 0$.

Now the expectations C_N are obtained by differentiating the $F_N(z)$ and then evaluating at $z = 1$. Doing this we obtain the recursion ($N \geq 1$)

$$(5.2) \quad C_N = 1 + \sum_{j=0}^d p_j \sum_{k=1}^N 2^{-N} \binom{N}{k} C_{k-j},$$

and the starting values $C_i = 1$ for $i \leq 0$.

Setting up the exponential generating function $C(z) = \sum_{N \geq 0} C_N z^N / N!$ we find

$$(5.3) \quad C(z) = e^z - e^{z/2} + \sum_{j=0}^d p_j e^{z/2} \sum_{k \geq 0} C_{k-j} \frac{(z/2)^k}{k!}.$$

Now it is customary (in order to get an easier equation) to set $D(z) = e^{-z} C(z)$ (the ‘‘Poisson transform’’). For simplicity, we replace z by $2z$, multiply the equation by e^{-z} and differentiate it d times, yielding

$$(5.4) \quad e^z \sum_{k=0}^d \binom{d}{k} 2^k D^{(k)}(2z) = e^z + \sum_{j=0}^d p_{d-j} C^{(j)}(z).$$

Furthermore, we can express (by Leibniz’ rule)

$$C^{(j)}(z) = e^z \sum_{i=0}^j \binom{j}{i} D^{(i)}(z),$$

so that we obtain after another multiplication by e^{-z}

$$(5.5) \quad \sum_{j=0}^d \binom{d}{j} 2^j D^{(j)}(2z) = 1 + \sum_{j=0}^d p_{d-j} \sum_{i=0}^j \binom{j}{i} D^{(i)}(z).$$

Now set $D(z) = \sum_{N \geq 0} D_N z^N / N!$ and $G(z) = \sum_{N \geq 0} D_N z^N$. Comparing coefficients we find

$$(5.6) \quad \sum_{j=0}^d \binom{d}{j} 2^{N+j} D_{N+j} = \delta_{N,0} + \sum_{j=0}^d p_{d-j} \sum_{i=0}^j \binom{j}{i} D_{N+i}.$$

Finally, we multiply by $(\frac{z}{2})^{N+d}$ and sum up for all $N \geq 0$ to obtain

$$(5.7) \quad G(z) \left(1 + \frac{z}{2}\right)^d = G\left(\frac{z}{2}\right) P_2(z) + P_0(z)$$

with

$$(5.8) \quad P_2(z) = \sum_{j=0}^d p_{d-j} \sum_{i=0}^j \binom{j}{i} \left(\frac{z}{2}\right)^{d-i} = \sum_{j=0}^d p_j \left(\frac{z}{2}\right)^j \left(1 + \frac{z}{2}\right)^{d-j}$$

and

$$(5.9) \quad P_0(z) = \left(\frac{z}{2}\right)^d + \sum_{j=0}^d \binom{d}{j} \left(\frac{z}{2}\right)^{d-j} \sum_{i=0}^{j-1} D_i z^i - \sum_{j=0}^d p_{d-j} \sum_{i=0}^j \binom{j}{i} \left(\frac{z}{2}\right)^{d-i} \sum_{h=0}^{i-1} D_h \left(\frac{z}{2}\right)^h.$$

We want to apply Theorem 4 and find $d_0 = d$, $a_d = b_d = c_d = 2^{-d}$ and thus $M = 0$. Hence the second alternative of Theorem 4 gives

$$(5.10) \quad G(z) = \log_2 z + \Psi_4(\log_2 z) + O(z^{-\epsilon}).$$

This leads to

$$(5.11) \quad F(z) = \sum_{N \geq 0} C_N z^N = \frac{1}{1-z} G\left(\frac{z}{1-z}\right) \\ = \frac{1}{1-z} \log_2 \frac{1}{1-z} + \frac{1}{1-z} \Psi_4\left(\log_2 \frac{1}{1-z}\right) + O((1-z)^{-1+\epsilon}),$$

for $z \rightarrow 1$. Now, since

$$[z^N] \frac{1}{1-z} \log \frac{1}{1-z} = 1 + \frac{1}{2} + \dots + \frac{1}{N} = \log N + \gamma + O\left(\frac{1}{N}\right),$$

singularity analysis (cf. [5]) leads to Theorem 5.

THEOREM 5. *The average level C_N that a party of N people with an additional killer will reach is asymptotically given by*

$$C_N = \log_2 N + \varphi(\log_2 N) + O(N^{-\epsilon})$$

where $\varphi(x)$ is a periodic function with period 1. Its Fourier coefficients could be given in principle by evaluating an appropriate Mellin integral at the points $2k\pi i / \log 2$, $k \in \mathbb{Z}$. $\varphi(x)$ depends on the probability distribution (p_0, \dots, p_d) , (with $p_d \neq 0$) of the killer.

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