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Zhenlu Wang, Bingchen Liang, Guoxiang Wu, B.W. Borsje



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1 Modeling the formation and migration of sand waves: The role of tidal forcing, sediment  
2 size and bed slope effects

3 Zhenlu Wang<sup>1,2</sup>, Bingchen Liang<sup>1,2,\*</sup>, Guoxiang Wu<sup>1,2</sup>, Borsje B. W.<sup>3</sup>

4 <sup>1</sup> College of Engineering, Ocean University of China, Qingdao, China

5 <sup>2</sup> Shandong Provincial Key Laboratory of Ocean Engineering, Ocean University of China, Qingdao, China

6 <sup>3</sup> Water Engineering and Management, University of Twente, Enschede, The Netherlands

7 Abstract:

8 Tidal sand waves are rhythmic bedforms existing widely in shallow shelf seas and are formed by the  
9 interaction of tidal currents and topography. Using a process-based numerical model, Delft3D, the wave  
10 lengths and migration rates of sand waves were simulated and verified with field measurements. The  
11 physical mechanism that controls the evolution of sand waves was mainly the balance between bedload  
12 transport, suspended load transport and the slope effect. It was found that the bedload transport multiplier,  
13 which reflects the bed slope effect, was a key parameter to reproduce the observed sand wave dynamics  
14 accurately. If the bedload transport multiplier is tuned with the actual grain size, it fits the observations on  
15 wavelength of sand wave much better. Both the migration rates and wavelengths were better predicted by  
16 the process-based numerical Delft3D model compared to a linear stability analysis sand wave model, since  
17 the former adopted sophisticated process formulations necessary for accurate field predictions. Next, sand  
18 wave formation and evolution under different environment settings, including tidal forcing and sediment  
19 sizes, were examined systematically. It was found that the preferred wavelength ( $L_{FGM}$ , fastest growing  
20 mode) of the sand wave increased with increasing tidal current magnitudes and decreasing sand diameters.  
21 Sand waves were only formed within a certain range and combination of tidal current magnitude and sand

22 diameters. Downstream- and upstream- migration of sand waves were predicted by considering residual  
23 currents or tidal constituent of higher harmonics.

24 Keywords: Sand waves; Delft3D; bed slope effect; migration rate; Morphodynamics; tidal forcing.

## 25 **1. Introduction**

26 The bed of shallow shelf seas inhabits regular bed forms of various sizes. According to their spatial scale, the  
27 bed forms can be categorized as sand ripples, sand waves or sand ridges (Németh et al., 2003). The wave  
28 lengths and heights of sand ripples are often too small (order of tens of centimeters) to influence the ocean  
29 engineering facilities (such as offshore underwater pipelines and cables) significantly. Tidal sand ridges  
30 typically have wavelengths of the order of 5~10 km and can be up to 30 m high, which is more than half of  
31 local water depth. The crests of sand ridges are usually rotated 10~30° anticlockwise with respect to the  
32 principal tidal current direction in the northern hemisphere (Hulscher, 1996). Since sand ridges hardly move,  
33 they do not interfere with ocean engineering facilities. From an engineering perspective, more attentions are  
34 drawn to the sand waves with wave lengths of up to hundreds of meters, the heights of several meters, and the  
35 migration rates of up to several meters per year (McCave, 1971; Terwindt, 1971; Huntly et al., 1993; Borsje et  
36 al., 2014). The crests of the sand waves are almost perpendicular to the direction of the dominating currents  
37 (Besio et al., 2003). Sand wave migrations have a significant effect on many offshore activities, for example,  
38 the navigation, pipelines and cables, and sand mining (Németh et al., 2003; de Jong et al., 2016; Roetert et al.,  
39 2017). Understanding the characteristics of sand wave dynamics is therefore critical for smart and sustainable  
40 design in offshore operations (Damveld et al., 2016).

41 Theoretical model analysis and field observation were widely used in sand wave studies. The most prevalent  
42 model approach is the stability analysis, considering the sandy bottom and the hydrodynamic as a coupled

43 system to study whether regular features are associated with free instabilities of the system. Huthnance (1982)  
44 was the first to use this kind of model to predict the evolution of sand banks. However, no sand waves were  
45 found with the model which is probably because the vertical structure of the flow was neglected (Hulscher et  
46 al.,1993). The model was further improved and modified to include the vertical dimension by Hulscher et al.  
47 (1993,1996), enabling it to predict the typical characteristics of sand waves. Hulscher (1996) described the  
48 formation of sand waves as a free instability of sandy seabed subject to tidal motions. The formation  
49 mechanism of sand waves was explained by that the interactions of the oscillatory tidal flows with sinusoidal  
50 bed perturbations give rise to a tidally-averaged residual current, in the form of vertical recirculating cells  
51 directed from the trough to the crest of sand waves (Borsje et al., 2013). In the model proposed by Hulscher  
52 (1996) only bed load was considered, and a constant vertical eddy viscosity with a partial slip condition at the  
53 bed instead of no-slip condition was used. By introducing a height- and flow-dependent model for turbulent  
54 viscosity, Komarova et al. (2000) found long waves can be damped and a finite separation can be found which  
55 is agreement with the observed data. Besio et al. (2003) studied the effect of sediment characteristics on  
56 formation of sand waves. Next, many more physical processes were included in the model, such as graded  
57 sediment (Roos et al., 2007; van Oyen et al., 2009), suspended sediment transport (Blondeaux et al., 2005;  
58 Blondeaux et al., 2016), biology (Borsje et al., 2009a; 2009b; Damveld et al., 2019), and wind wave effects  
59 (Campmans et al., 2017). Moreover, models were developed to predict the migration of sand waves by  
60 including residual current (Németh et al., 2002) or different tidal components (Besio et al., 2004). In order to  
61 investigate the long-term behavior of sand waves, non-linear stability models were developed (Németh et al.,  
62 2007; van den Berg, 2007; van den Berg et al., 2012; Campmans et al., 2018).

63 Recently, process-based numerical models, such as Delft3D, ROMS (Zang et al., 2011), have been widely used

64 in the study of coastal morphodynamics. Tonnon et al. (2007) studied the morphodynamic behavior of an  
65 artificial sand wave using Delft3D. Borsje et al. (2013, 2014) used Delft3D to study the formation mechanism  
66 of sand waves, and investigated the effect of turbulence formulation and suspended sediment transport. The  
67 role of suspended sediment transport and tidal asymmetry on the equilibrium sand wave height were further  
68 examined by van Gerwen et al. (2018).

69 In summary, the models showed that the characteristics of sand waves (wavelength, growth rate, and migration  
70 rate) depend on various environment factors, such as sediment grain sizes, tidal characteristics and the  
71 associated residual currents. However, the sensitivity of sand wave dynamics to the environment factors is  
72 hardly explored. The characteristics of sand waves and the environment parameters in two areas of North Sea  
73 were summarized by Menninga (2012). With these field data, a Delft3D model was set up and verified. Next  
74 the model was used to systemically investigate the role of variations in various environmental factors on sand  
75 wave characteristics.

76 The aim of the paper is threefold: (1) to explain the main dynamic characteristics of sand waves in the  
77 study-sites of IJmuiden area and Rotterdam area (Menninga, 2012) by applying a process-based numerical  
78 model Delft3D, and compare it with a linear stability sand wave model of Blondeaux et al. (2016), (2) to  
79 examine the sensitivity of sand wave characteristics to environment factors, (3) to investigate the migration  
80 rates under different combinations of tidal components and related phase differences systematically and the  
81 impact on wavelength of the sand wave. The remainder of the paper is organized as follows:

82 Section 2 introduces the model equations and model set-up. In section 3, the comparison between field data  
83 and model results are given. Subsequently sensitivity of environment factors on characteristics of sand waves  
84 are analyzed. Finally, section 4 and 5 contain the discussion and conclusions, respectively.

## 85 2. Methods and model description

### 86 2.1 Model formulation

87 The Delft3D modeling system consists of a hydrodynamic model based on the nonlinear shallow water  
88 equations, a suspended sediment and bedload transport model and a bed evolution model. Because most sand  
89 wave crests are oriented almost perpendicular to the tidal currents, the main physical process controlling the  
90 sand wave morphodynamics can be captured with the 2DV mode of Delft3D. Coriolis effects are assumed to  
91 have a negligible effect on the evolution of sand waves (Hulscher, 1996) and therefore not taken into account  
92 in the sand wave model. The basic hydrodynamics equations are given in Appendix A.

93 Sediment can be transported as bedload and suspended load. Sediment transport and morphological evolution  
94 equations are given in Appendix B. A detailed model formulations and solutions can be referred to Lesser et al.  
95 (2004) and Delft3D Flow manual.

### 96 2.2 Model set-up

97 The length of the model domain is about 50 km. 6 km of sand waves are designed to lay in the center of the  
98 model domain. The horizontal grid size in the sand wave area is 10 m and increases gradually up to 1500 m at  
99 the lateral boundaries (Figure 1(a)). In the vertical dimension, 20  $\sigma$ -layers is set with high resolution near the  
100 bed and decreasing towards to the water surface (Figure 1(b)). The initial amplitude of the sand waves is 0.5 m,  
101 which is small compared with the water depth. The sand wave field is multiplied by a tapering function to  
102 create a transition region from the sand wave domain to flat-bed. Riemann boundary conditions are imposed at  
103 lateral boundaries so that the tidal waves can travel out of the model domain without reflecting back. We used  
104 S2 tide instead of M2 for the convenient with the FFT analysis, because the tidal period of this constituent is  
105 exactly 12 hours. Also, we considered M0 tide (residual current) and the S4 tide (instead of M4) as the input

106 parameters. The Chézy coefficient is set to a constant value of  $65 \text{ m}^{1/2}/\text{s}$  (Borsje et al., 2013). The time step is  
 107 12 s, and the model runs for two tidal circles. In the first tidal circle, only hydrodynamics is calculated to  
 108 achieve a steady state (spin-up time), then in the second tidal circle bed updating is activated. The bed  
 109 evolution can be described as  $z_b = A(t)\exp(2\pi i x / L) + c.c.$  (here  $L$  is the sand wave length). Assuming  
 110 exponential growth (Besio et al., 2008; Borsje et al., 2013), the growth rate  $\gamma_R$  and migration rate  $\gamma_I$  is  
 111 calculated by:

$$112 \quad \gamma_R = \frac{1}{T} \text{Re}(\log(A1/A0)) \quad (1)$$

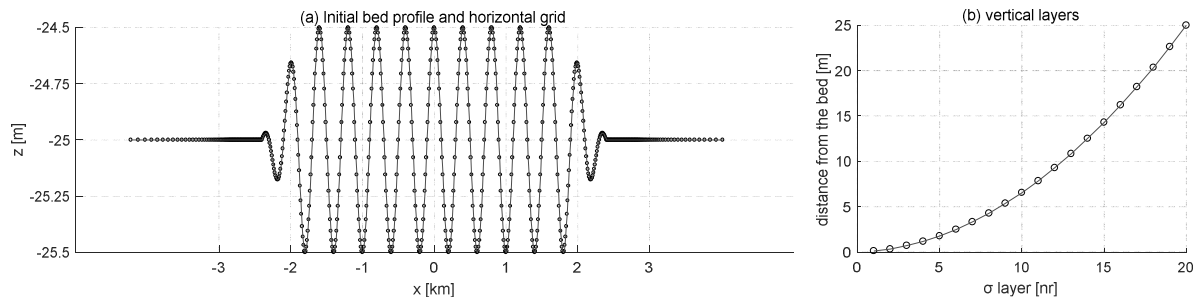
$$113 \quad \gamma_I = \frac{-1}{kT} \text{Im}(\log(A1/A0)) \quad (2)$$

114 The model parameters used was presented in Table 1.

115 Table 1. Overview of values of model parameters.

Description	symbol	value	unit
gravity acceleration	$g$	9.81	$\text{m/s}^2$
Chézy roughness	$C_D$	65	$\text{m}^{1/2}\text{s}^{-1}$
water density	$\rho_w$	1025	$\text{kg/m}^3$
sediment density	$\rho_s$	2650	$\text{kg/m}^3$
Kármán's constant	$\kappa$	0.41	-
bed porosity	$p_{or}$	0.4	-
initial sand wave amplitude	$A_0$	0.5	m

116



117

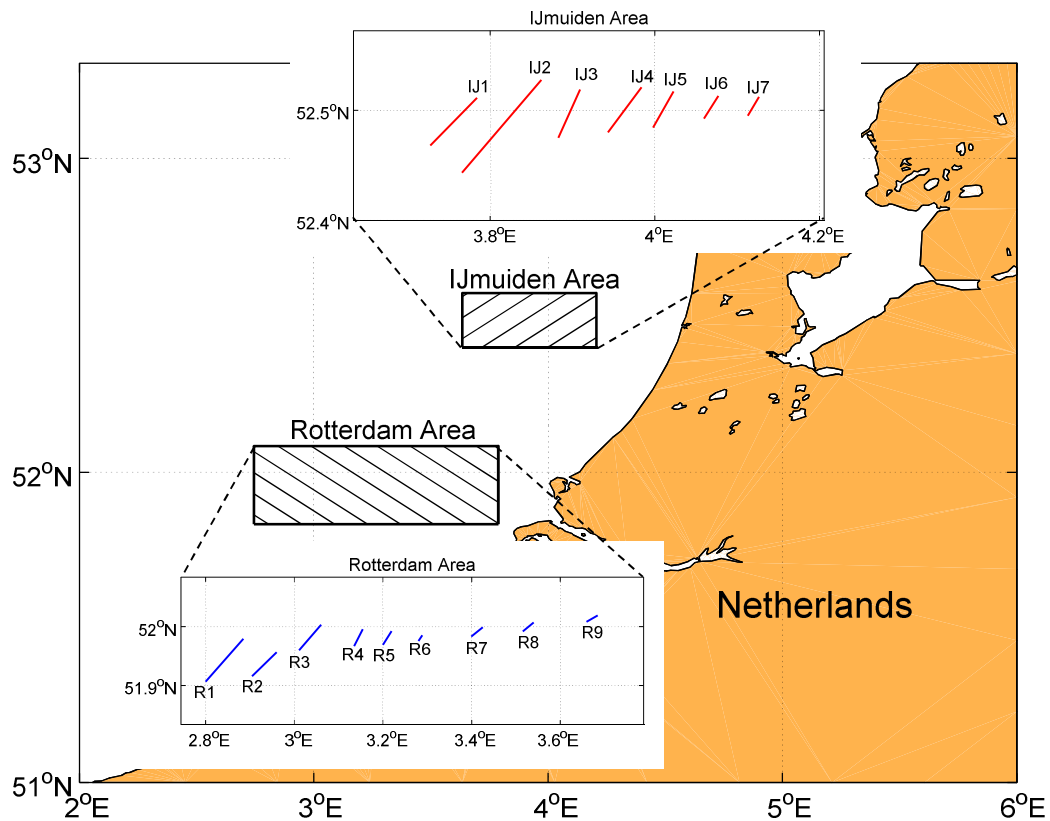
118 Figure 1. the grid set-up with (a) initial bed profile for wavelength of 400 m sand wave and the horizontal grid using

119 the fine grid in the center 6 km where the sand waves calculated and a coarser grid towards the boundaries and (b)  
120 the distribution of vertical  $\sigma$ -layers over the water depth. Note that the wavelength of the initial bed profiles varies  
121 in a range to determine the fastest growing mode and the water depth changed for the different model simulations.

### 122 2.3 Environment parameters and characteristics of observed sand waves

123 The model presented above was used to investigate the effects of suspended sediment transport (Borsje et al.,  
124 2014) and different turbulence formulations (Borsje et al., 2013) on the formation of sand waves. The model  
125 showed fairly good agreement with field data qualitatively. As Figure 2 shows, data sets of two areas in the  
126 North Sea was collected and analyzed by Menninga (2012). IJmuiden area comprises 170 km<sup>2</sup>, with latitudinal  
127 coordinates (in UTM-WGS84) ranging from 5810100 to 5820200 and longitudinal coordinates from 546500 to  
128 604200, and Rotterdam area comprises 310 km<sup>2</sup>, with latitudinal coordinates ranging from 5742600 to  
129 5766300 and longitudinal coordinates ranging from 476900 to 564300.





130

131

Figure 2. Study areas: IJmuiden area and Rotterdam area (adapted from Menninga (2012)).

132

Table 2 shows the environmental parameters at two sites in the North Sea collected by Menninga (2012),

133

namely, IJmuiden area (denoted by IJ-, including 7 sites), and Rotterdam area (denoted by R-, including 9

134

sites), respectively. The two sites have quite different environmental parameters, and the water depth in

135

Rotterdam area is between 26.6~36 m, which is larger than that in IJmuiden (23.6~27 m). The tidal-velocity

136

amplitude in IJmuiden area is between 0.63 m/s and 0.66 m/s which is smaller than that in Rotterdam area

137

(0.69~0.79 m/s). The mean sediment diameter in Rotterdam area is about 0.358 mm, which is also larger than

138

that in IJmuiden (0.284 mm). The characteristics of sand waves are different in two areas. The main difference

139

in characteristics of sand waves is that the wave length and the migration rate in IJmuiden area are both larger

140

than those in Rotterdam area. Refer to Menninga (2012) and Blondeaux et al. (2016) for more details on

141 environmental parameters and characteristics of sand waves.

142 Table 2. Environmental parameters and Characteristics of sand waves of two sites (Menninga, 2012)

Site	$H$ (m)	$d_{50}$ (mm)	$U_{M2}$ (m/s)	$ecc_{M2}$	$U_{M4}$ (m/s)	$ecc_{M4}$	$\Phi_{M2-M4}$ (degrees)	$U_{M0}$ (m/s)	$L_{sw}$ (m)	$H_{sw}$ (m)	$c_{crest}$ (m/year)	$c_{trough}$ (m/year)	$c_{mean}$ (m/year)
IJ1	27.0	0.309	0.66	0.018	0.047	0.51	339	0.019	380	3.4	2.4	0.2	1.30
IJ2	26.5	0.308	0.66	0.009	0.047	0.54	342	0.018	395	3.3	2.3	-0.1	1.10
IJ3	26.0	0.288	0.65	0.001	0.046	0.58	346	0.019	450	3.4	2.6	0.2	1.40
IJ4	24.5	0.278	0.64	0.006	0.046	0.62	352	0.018	428	2.6	2.8	0.4	1.60
IJ5	24.3	0.268	0.64	0.011	0.045	0.64	355	0.017	395	2.3	3.3	0.4	1.85
IJ6	23.8	0.271	0.63	0.012	0.043	0.67	354	0.022	310	1.5	3.1	1.2	2.15
IJ7	23.6	0.266	0.63	0.005	0.041	0.76	0.84	0.029	615	2.8	2.7	0.2	1.45
R1	36.0	0.445	0.79	0.11	0.060	0.07	115	0.019	315	4.0	-0.1	0.0	-0.05
R2	35.4	0.421	0.73	0.11	0.053	0.09	130	0.04	200	3.2	0.1	0.0	0.05
R3	33.3	0.398	0.71	0.10	0.054	0.11	124	0.044	270	4.1	0.4	0.3	0.35
R4	32.6	0.372	0.73	0.11	0.057	0.16	122	0.041	220	3.5	0.7	0.5	0.60
R5	31.8	0.317	0.72	0.11	0.056	0.19	121	0.039	205	3.1	1.0	0.9	0.95
R6	32.0	0.332	0.72	0.12	0.057	0.21	120	0.039	225	3.3	0.7	0.6	0.65
R7	29.0	0.306	0.71	0.12	0.057	0.25	116	0.039	245	3.6	1.0	1.5	1.25
R8	28.4	0.336	0.69	0.12	0.058	0.27	110	0.034	250	3.3	1.0	0.9	0.95
R9	26.6	0.301	0.67	0.11	0.059	0.27	102	0.034	315	2.1	2.1	2.1	2.10

143 where  $H$  is water depth,  $d_{50}$  is the median sediment diameter,  $U_{M2}$  is the amplitude of M2 tidal current,  
 144  $ecc_{M2}$  and  $ecc_{M4}$  are M2 ellipticity and M4 ellipticity respectively,  $U_{M4}$  is the amplitude of M4 tidal current,  
 145  $\phi_{M2-M4}$  is the phase between M2 and M4 tide,  $U_{M0}$  denotes the residual current,  $L_{sw}$  and  $H_{sw}$  are the  
 146 wavelength and height of sand waves respectively,  $c_{crest}$ ,  $c_{trough}$ ,  $c_{mean}$  are the crest/trough/mean migration rate  
 147 of sand waves.

148 Predictions using a stability analysis matched the field data qualitatively in terms of various characteristics of  
 149 the sand waves (Blondeaux, 2016). In their model the wind waves and different tide constituents were taken  
 150 into account, a time-independent eddy viscosity function was used to close the hydrodynamic problem, and  
 151 suspended sediment transport was included in addition to bedload transport. In their stability analysis model,  
 152 default values were used for all parameters, although the values of some key parameters (such as sediment

153 sizes) fall into a large range. Considering that modeled sand wave dynamics may be very sensitive to these  
154 environment parameters, in this study we use the Delft3D model to analyze the characteristics of the sand  
155 waves systematically with different settings of environmental parameters and sediment characteristics.

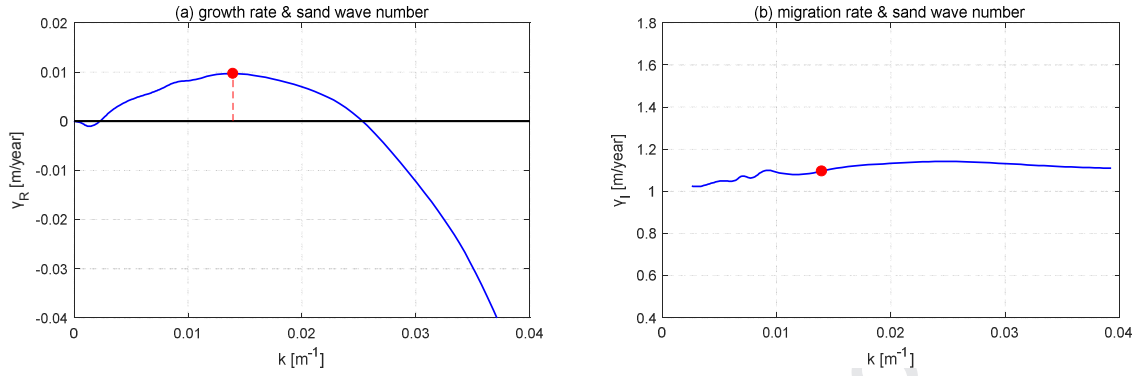
### 156 **3. Results**

#### 157 3.1 Model validation: Comparison with field data

158 The model was established using the environmental parameters of case IJ3 and case R3 in the Table 2,  
159 respectively. Bed evolution was calculated with different sand wave numbers as the initial bed level. For the  
160 IJ3 case, the slope factor  $\alpha_{bs}$  was set to 1.4. With Fast Fourier Transform (FFT) analysis, the relationship  
161 between sand wave growth rates and sand wave numbers was obtained, as summarized in Figure 3 (a). The red  
162 dot represents the fastest growing mode, which corresponds the preferred wavelength of the sand wave under  
163 the specified environmental parameters. From Figure 3 (a), the  $L_{FGM}$  reads 450 m which matches perfectly  
164 with the observed data. The migration rates versus sand wave numbers are shown in Figure 3 (b). The  
165 migration rate corresponding to the fastest growing mode (red dot) is 1.1 m/year, which is close to the  
166 observed migration rate of 1.4 m/year. The same method was used on R3 case, with the same slope factor of  
167 1.4, the  $L_{FGM}$  reads 400 m which is much larger than the observed data (270 m). After calibrating the model  
168 with the slope factor of 1.0, it turns out a better prediction of  $L_{FGM}$  (290 m). The migration rate is 1.49 m/year  
169 with slope factor of 1.0, which is almost the same (1.51 m/year) as that with a slope parameter of 1.4.

170 The verified model was used on other sites of the two areas to calculate the preferred wavelength and  
171 migration rate of the sand wave. As the first 7 sites located the same area, i.e. IJmuiden, we used the same  
172 slope factor of 1.4. For the rest 9 sites located in the Rotterdam area, the slope factor of 1.0 was used.  
173 Blondeaux et al. (2016) used stability analysis model to study the characteristics of sand waves and the results

174 can be found in Table 3. The results of numerical simulation we used are also listed in Table 3.



175

176 Figure 3. Growth rates (a) and migration rates (b) under different sand wave numbers. The curves are corresponding

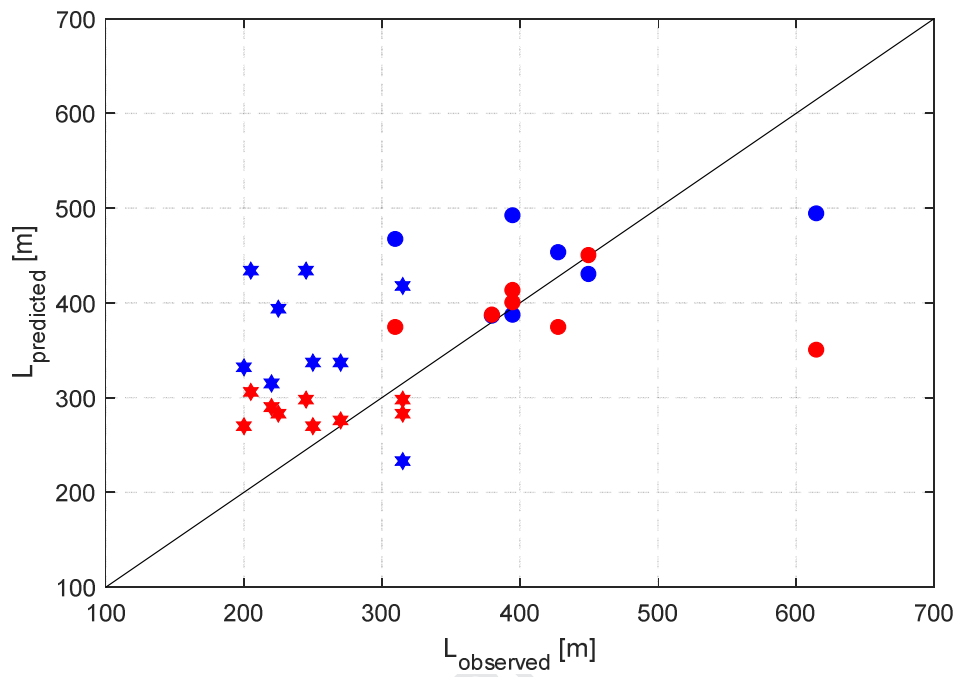
177

to the case of J3 and the red dot indicates the fastest growing mode.

178

Table 3 Sand wave characteristics predicted in Blondeaux et al. (2016) and in this study

Site	results in Blondeaux et al. (2016)		results in this study	
	$L_{sw\_S}$ (m)	$c_S$ (m/year)	$L_{sw\_N}$ (m)	$c_N$ (m/year)
IJ1	386	3.77	387	2.73
IJ2	387	3.63	400	2.72
IJ3	430	3.44	450	1.10
IJ4	453	3.18	374	2.77
IJ5	492	2.85	413	2.70
IJ6	467	3.45	374	2.91
IJ7	494	4.39	350	3.34
R1	233	-0.35	270	-0.09
R2	332	3.91	276	1.17
R3	337	4.32	290	1.63
R4	315	4.29	306	1.49
R5	434	4.00	283	1.35
R6	394	3.76	298	1.39
R7	434	4.08	270	1.59
R8	337	2.67	283	1.39
R9	418	2.63	270	1.68

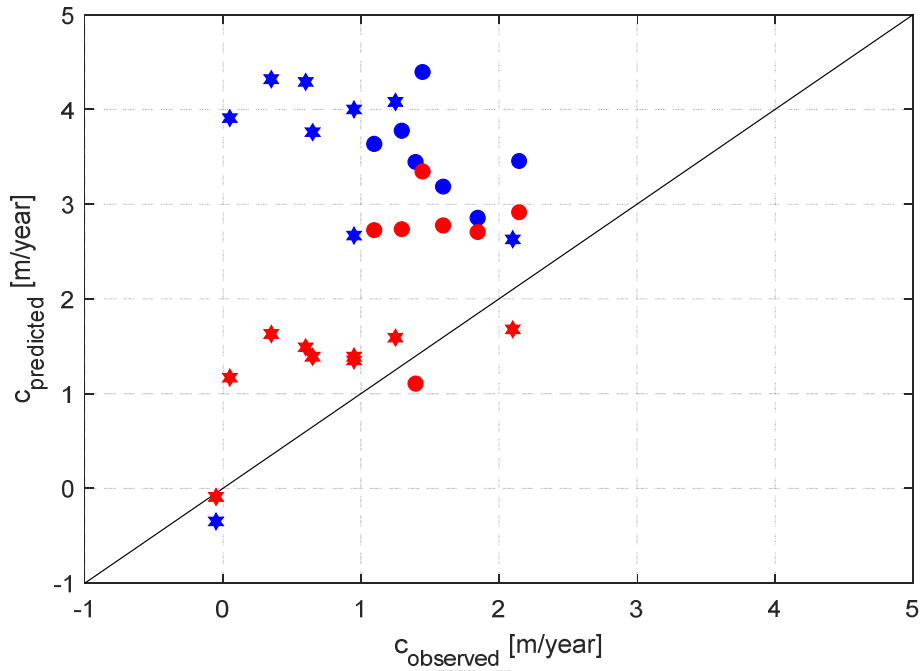


179

180 Figure 4. Comparison of predictions on wavelengths of sand waves by the stability model (blue marks, dots for

181 IJmuiden site and hexagrams for Rotterdam site, Blondeaux et al. 2016), Delft3D (red marks, dots for IJmuiden site

182 with  $\alpha_{bs} = 1.4$ , hexagrams for Rotterdam site with  $\alpha_{bs} = 1.0$ ).



183

184 Figure 5. Comparison of predictions of sand wave migration rates by the stability model (blue marks, dots for

185 IJmuiden site and hexagrams for Rotterdam site, Blondeaux et al. 2016), Delft3D (red marks, dots for IJmuiden site

186 with  $\alpha_{bs} = 1.4$ , hexagrams for Rotterdam site with  $\alpha_{bs} = 1.0$ ).

187 Figure 4 shows the predicted wavelengths of sand waves of stability model and Delft3D versus observed field

188 values. All the predicted values seemed well agreement with the observed values qualitatively. Three indexes,

189 namely, the Bias, the Root-mean-square-errors (RMSE), and the Correlation relationship, were used to make

190 quantitate comparisons between the two models.

191

$$Bias = \frac{1}{N_d} \sum_{i=1}^{N_d} (S_i - M_i) \quad (3)$$

192

$$RMSE = \sqrt{\frac{1}{N_d} \sum_{i=1}^{N_d} (S_i - M_i)^2} \quad (4)$$

$$r = \frac{\sum_{i=1}^{N_d} [(S_i - \bar{S})(M_i - \bar{M})]}{\sqrt{\sum_{i=1}^{N_d} (S_i - \bar{S})^2 \sum_{i=1}^{N_d} (M_i - \bar{M})^2}} \quad (5)$$

193 where  $N_d$  is the total number of the values of observations and predictions,  $S_i$  is the predicted value and  
 194  $M_i$  is the observed value. Bias is also referred as the mean error. A smaller Bias value indicates a better  
 195 agreement between observed and predicted values. The  $r$  denotes the correlation coefficients.

197 Table 4. statistical indicator values for characteristics of sand waves

Statistical indicators	Test <sup>a</sup>	Values for wave length of sand waves	Values for migration of sand waves
RMSE	T1	117.9	2.56
	T2	81.2	0.99
Bias	T1	70.3	2.27
	T2	6.5	0.76
$R$	T1	0.51	0.16
	T2	0.69	0.68

198 <sup>a</sup> Two tests are performed, and test T1 represents for the stability results of Blondeaux et al. (2016). T2 represent for numerical results in this  
 199 paper,  $\alpha_{bs}=1.4$  for IJmuiden area and 1.0 for Rotterdam area.

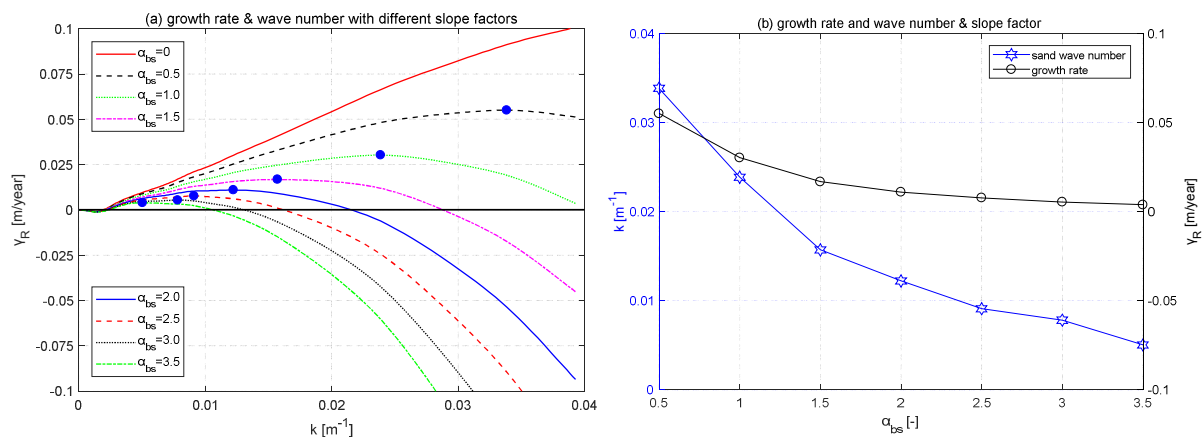
200 Table 4 shows the statistical indicator values for characteristics of sand waves. For the wave length of the  
 201 stability model, RMSE is 117.9, Bias is 70.3 while the correlation coefficient is 0.51. For the numerical model,  
 202 the indicators are 81.2, 6.5 and 0.69 respectively when the slope factor  $\alpha_{bs}$  is 1.4 for IJmuiden area and 1.0  
 203 for Rotterdam area. The prediction of wave lengths of sand waves was much better using the numerical model.  
 204 The same method was used to estimate the difference between predicted values and observed values of  
 205 migration rate (Figure 5). For the stability model, the values of three indicators are 2.56, 2.27 and 0.16 which  
 206 means predicted values are not in agreement well with observed values. Then for numerical model the three  
 207 indicators are 0.99, 0.76 and 0.68 respectively. These values are much better than that predicted by stability

208 models. With different slope factors in two areas, a better prediction on wavelengths and migration rates of  
 209 sand waves can be obtained.

210 The bed slope effect can change the transport rate once the sediment is in motion (van Rjin, 1993). This effect  
 211 can be formulated by Equation. 28 in Appendix B. In general, the repose angle of coarser sand is larger than  
 212 finer sand which will induce the slope coefficient  $\alpha_s$  smaller. From Equation. 28 of Appendix B we can use a  
 213 smaller slope factor  $\alpha_{bs}$  to tune the slope coefficient  $\alpha_s$  smaller. As can be found in Table 2, the sediment  
 214 diameter in the Rotterdam site is larger than that in the IJmuiden site, so it is reasonable to use a smaller slope  
 215 coefficient for the Rotterdam area.

### 216 3.2 The effects of bed slope

217 The effects of different slope factors on sand wave growth and  $L_{FGM}$  were examined. A series of idealized  
 218 model settings were designed, in which only an S2 tide was considered. Other parameters were kept fixed as  
 219 follows:  $H = 25$  m,  $U_{S2} = 0.65$  m/s,  $d_{50} = 0.35$  mm. These parameters resemble a typical North Sea situation for  
 220 sand wave occurrence (Borsje, et al., 2009a). We varied the slope factor  $\alpha_{bs}$  in the range of  $0 \sim 3.5$  with an  
 221 interval of 0.5.



222  
 223 Figure 6. (a) The growth rates versus sand wave numbers with different slope factors, the blue dots indicating the



224 fastest growing mode. (b) The growth rates and sand wave numbers versus slope factors under the fastest growing  
 225 mode conditions.

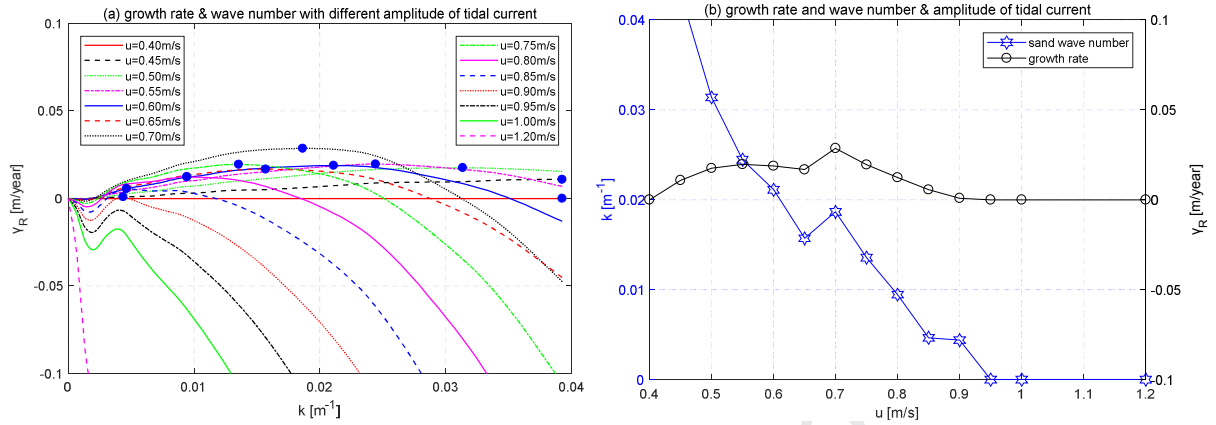
226 As Figure 6(a) shows, without considering slope effect (i.e.  $\alpha_{bs} = 0$ ), the sand waves will always be growing  
 227 in the range of the specified sand wave numbers of 0~ 0.04. For large wave numbers, the tide-averaged vertical  
 228 circulation flow is strongest (Borsje et al., 2013) leading to the largest growth of the sand waves. However,  
 229 when the slope factor  $\alpha_{bs}$  is larger than 1.0, the curves between the growth rate and sand wave number start  
 230 to display the parabolic shape as a result of slope-induced transport. A critical point exists at the peak of the  
 231 parabolic (blue dots in Figure 6(a)) which represents the fastest sand wave growth rates. The bed slope effect  
 232 acts as a decaying mechanism which makes the upslope sand transport more difficult than downslope transport.  
 233 The decaying mechanism will be enhanced under a larger slope factor, resulting in a slower growth rate.  
 234 Besides, the range of possible wavelengths of sand waves will be narrower under a larger slope factor.

235 Figure 6(b) shows the relationship of growth rates and the preferred wave numbers of sand waves versus slope  
 236 factors. It is clear that the wave length decreased rapidly while the growth rate of sand waves only decreases  
 237 slightly with the increase of the slope factor. When the slope factor is large enough, the sand waves hardly  
 238 grow.

### 239 3.3 The effects of tidal currents and sediment sizes

240 The effects of different tidal currents and sediment sizes on sand wave growth and  $L_{FGM}$  were examined. Two  
 241 series of numerical experiments were carried out. For the first series, some parameters were kept fixed as  
 242 follows:  $H = 25$  m,  $d_{50} = 0.35$  mm,  $\alpha_{bs} = 1.5$ . We varied the amplitude of tidal currents in the range of  
 243 0.4~1.2 m/s with interval of 0.05 m/s. For the second series, the parameters were kept fixed as follows:  $H = 25$   
 244 m,  $U_{S2} = 0.65$  m/s,  $\alpha_{bs} = 1.5$ . We varied the sediment sizes in the range of 0.15~0.55 mm with an interval of

245 0.05 mm.



246

247 Figure 7. (a) the growth rate versus sand wave number with different amplitude of tidal current and the blue dots

248 indicate the fastest growing mode. (b) the growth rate and sand wave number versus amplitude of tidal current under

249

the fastest growing mode conditions.

250 The red solid line in Figure 7(a) shows that the sediment stayed still when the tidal current strength was less

251 than 0.4 m/s. For increasing tidal amplitudes of the tidal current, sediment will start to move as soon as the bed

252 shear stress became larger than the critical bed shear stress. When the amplitude of tidal current was larger than

253 0.45 m/s, the sand waves keep growing in the range of specified wave numbers. The fastest growing mode

254 (blue dots) appeared in a specific range of the tidal currents. On the other side, if the velocity was too large

255 (larger than 0.9 m/s in this study), the growth rate of sand waves stays negative, indicating that no sand waves

256 were formed. This test also proved that sand waves can only be formed within certain range of tidal current

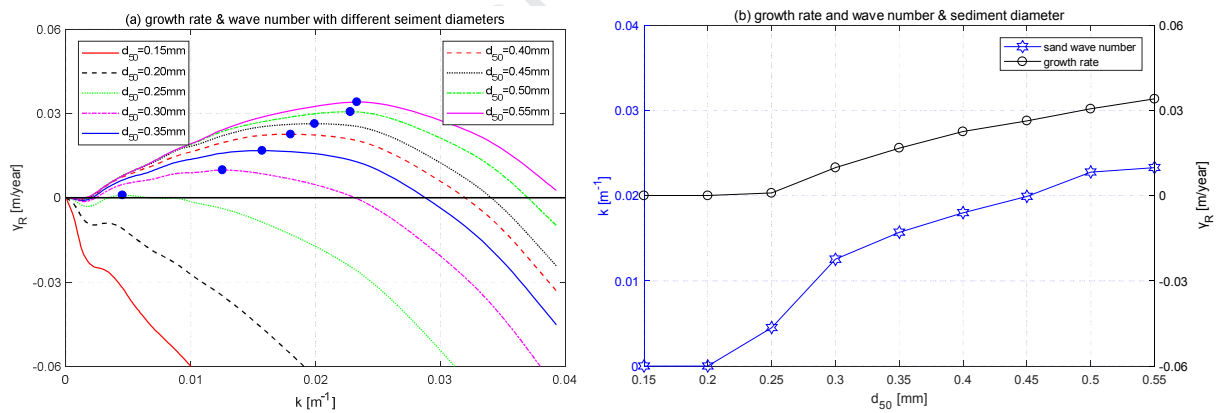
257 magnitudes in accordance with previous studies (Borsje et al., 2014). Generally, the range of the wavelength of

258 the sand wave which may be appeared will become narrow with the increase of the tidal current strength until

259 the sand waves disappear. Figure 7(b) shows the sand wave number and the growth rate of fastest growing

260 mode under different tidal current magnitudes. The growth rates first increased, and then decreased, with

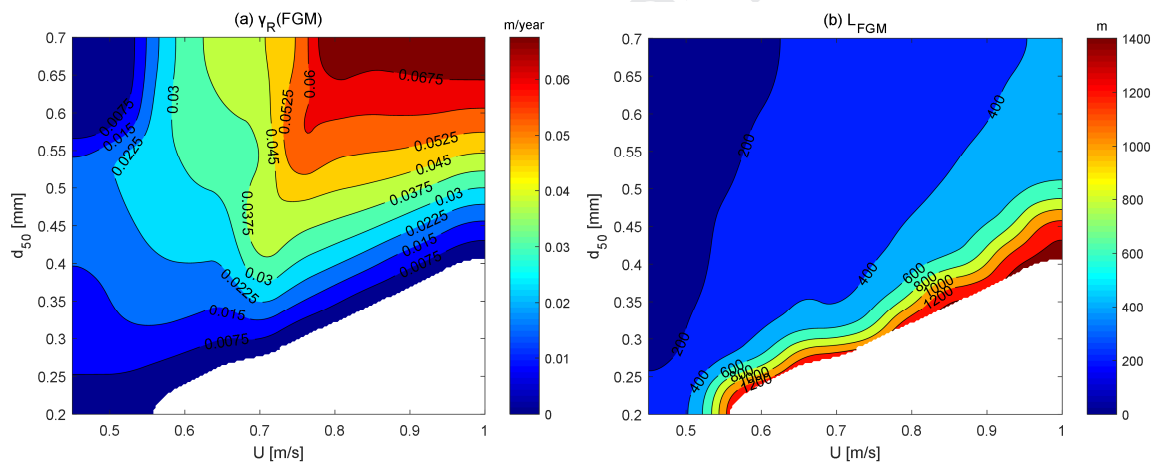
261 increasing tidal currents. For this parameter setting, the maximum value occurred at the tidal velocity of 0.7  
 262 m/s. When the velocity was smaller, less sediment was moved thus the growth rates of sand waves were small  
 263 as well. When the velocity was larger, most sediment moves as suspended load, which may decay the growth  
 264 of the sand waves (Borsje et al., 2014). The wavelength of the sand wave increased rapidly in general with  
 265 increasing tidal currents. Note there appears a slight discontinuous behavior when the amplitude of tidal  
 266 current is 0.7 m/s. This is due to a combination of bedload transport and suspended sediment transport. With  
 267 increasing tidal currents, the growth rate of sand waves due to bedload increases while the growth rate due to  
 268 suspended sediment decreases. There is a small difference in the rate of increase and decrease which induced  
 269 the discontinuous behavior. As shown in Figure 7, the sand wave most likely occurred in the range of 0.45~0.9  
 270 m/s for this parameter setting.



271 Figure 8. (a) the growth rates versus wave numbers with different sediment diameters and the blue dots indicate the  
 272 fastest growing mode. (b) the growth rates and the wave numbers of sand waves versus sediment diameters under  
 273 the fastest growing mode conditions.  
 274

275 It is found from Figure 8(a) that when the sediment diameter was smaller than 0.2 mm, no sand waves were  
 276 formed and all sand waves decayed, resulting in a flat bed. It is because that suspended load dominates when

277 the sediment diameter is small, preventing the formation of sand waves. The growth rates become positive and  
 278 the sand waves appear when the sediment diameter is larger than 0.25 mm. As the sediment diameter continued  
 279 to increase, bedload transport began to dominate, which promotes sand waves growth. As shown in Figure 8(a),  
 280 the range of possible wavelengths of sand waves was generally wider for larger sediment diameters. Figure 8(b)  
 281 shows the sand wave numbers and growth rates of the fastest growing mode versus sediment diameters. The  
 282 growth rate of the sand wave increased and the wavelength decreased with the increase of the sediment  
 283 diameter. To conclude, our simulations showed that finer sediment tend to produce longer wave lengths, while  
 284 coarser sediment tend to produce shorter wave lengths.



285  
 286 Figure 9. (a) the relationship between growth rate of fastest growing mode and sediment diameter and tidal current.

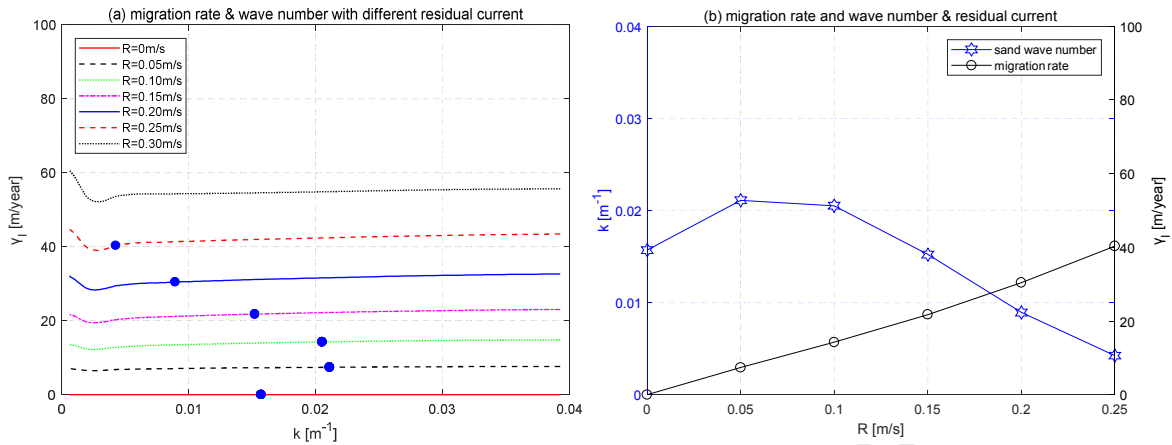
287 (b) the relationship between preferred sand wavelength and sediment diameter and tidal current.

288 Simulations with different combinations of sediment diameters and tidal currents were further conducted. The  
 289 water depth was 25 m, and the slope factor was set to 1.5. Figure 9 shows that at the upper left and lower right  
 290 the sand waves are absent. When the velocity is weak and the sediment is coarse (the upper left), the sediment  
 291 hardly can move by the tidal current, so the bed form will maintain stable. When the velocity is too strong and  
 292 the sediment is fine (the lower right), the dominated sediment transport is suspended load transport which

293 damps the growth of the sand waves, the bed form also stays stable. The parameters of the most likely  
294 appearance of the sand waves are coarser sediment and relative larger tidal current which can be found at the  
295 upper right of the Figure 9(a). Figure 9(b) shows the wavelength of fastest growing mode of sand waves  
296 mainly in the range of 200 m to 600 m. The longer wavelength of sand waves appears on the condition of the  
297 high velocity and fine sediment where the suspended sediment transport dominate the bed form evolution.  
298 Previous research and observed field data indicate that the wavelength is up to hundreds of meters which  
299 support the results here we found.

### 300 3.4 Sand wave migration: the role of tidal asymmetry

301 Sand waves can be formed and grow under a symmetry tide (such as the S2 used in above simulations).  
302 However, sand wave migrations may occur in case of the distortion and loss of tidal symmetry. Here we  
303 analyzed the relationship between sand wave migrations and tidal asymmetry. Two series of numerical  
304 experiments were conducted. In the first series only an S2 tide was considered, the parameters were kept fixed  
305 as follows:  $H = 25$  m,  $U_{S2} = 0.65$  m/s,  $d_{50} = 0.35$  mm,  $\alpha_{bs} = 1.5$ . We varied the residual currents in the range  
306 of 0~0.3 m/s with an interval of 0.05 m/s. In the second series, the parameters were kept fixed as follows:  $H =$   
307  $25$  m,  $U_{S2} = 0.65$  m/s,  $d_{50} = 0.35$  mm,  $U_{S4} = 0.05$  m/s,  $\alpha_{bs} = 1.5$ , and the residual current were set to 0 m/s.  
308 We varied the phases of S4 tide in the range of 0~360° with an interval of 30°.



309

310 Figure 10. (a) the migration rate versus wave number with different residual current and the blue dots indicate the

311 fastest growing mode. (b) the migration rate and the wave number of sand waves versus residual current under the

312

fastest growing mode conditions.

313 The red line in Figure 10(a) shows that no migration occurs under a single S2 tide. When the residual currents

314 are added, sand waves start to migrate in the direction of the residual currents. The symbols in Figure 10(b)

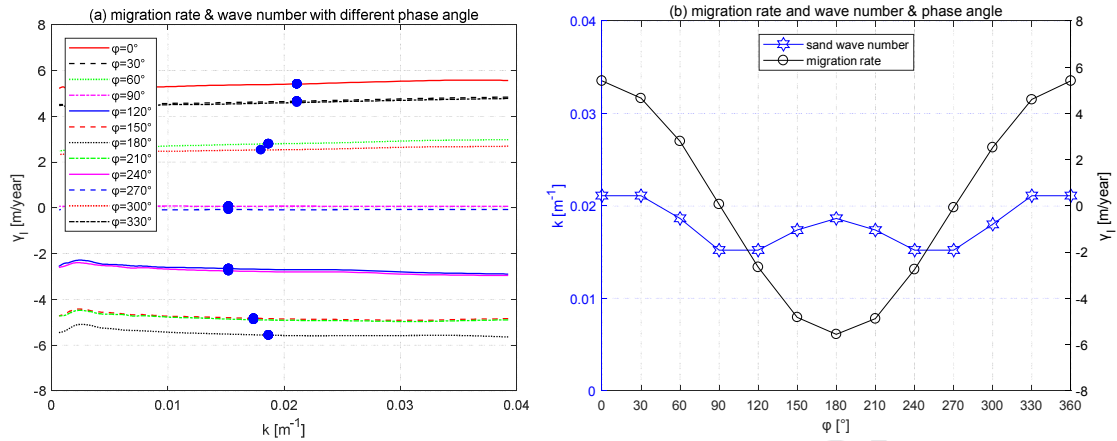
315 corresponding to the fastest growing mode which is the most preferred bed forms that may appear. It is found

316 that migration direction always followed the direction of residual currents and the migration rates increased

317 with increasing residual currents. However, when the residual current was too strong, the fastest growing mode

318 disappeared, as shown by the black dash line in Figure 10(a). The wavelength of the sand wave first increased

319 and then decreased with the increase of the residual currents.



320

321 Figure 11. (a) the migration rate versus wave number with different phases between S2 and S4 and the blue dots

322 indicate the fastest growing mode. (b) the migration rate and the wave number of sand waves versus phases under

323 the fastest growing mode conditions.

324 Previous studies suggested that the migration direction of the sand waves may be opposite to the residual

325 current (Besio, 2004). Figure 11 shows the migration rate of sand waves with different phases between the tide

326 component S2 and S4. The curve of sand wave number and migration rate corresponding to fastest growing

327 mode is symmetrical on both side taking 180° as the centerline. When the phase is 90° or 270°, the sand wave

328 cannot migrate. While the phase is falling in the scope of 90°~270°, the migration direction will be negative.

329 Although the current induced by S2 and S4 tide constituents will vanish during the time averaged, the

330 nonlinear relationship between fluid velocity and the sediment transport will induce the sand waves move

331 towards a specific direction. As a result, the migration direction of the sand waves will depend on the relative

332 intendance of phase of different tidal components and residual current. If the residual current is small and the

333 higher harmonic tidal component is large or the phase between tidal components is appropriate, then migration

334 direction of sand waves may be opposite to the residual current. The  $L_{FGM}$  varied according to the phases335 between the tidal components. The  $L_{FGM}$  decreased with the increasing of the migration rate no matter the sand

336 wave migrates forward or backward. When the phase is  $90^\circ$ , the wave length is the longest and the sand wave  
337 cannot migrate.

#### 338 **4. Discussion**

339 In this paper, we studied the effect of some environment parameters on formation and migration of sand waves.  
340 The model was calibrated with slope factor for two specific sites, the wavelength of the sand wave is  
341 reproduced good compared to observed data. The migration rate is considerably better predicted with the  
342 numerical model compared to the Blondeaux et al. (2016) model. In Blondeaux et al. (2016), the eddy  
343 viscosity was assumed to be time-independent following Besio et al. (2006). The eddy viscosity we used in the  
344 Delft3D model was derived from a more accurate turbulence formulation ( $k - \varepsilon$ ), and it may gave a better  
345 prediction on the wavelengths and the migration rates of sand waves.

346 In general, the migration rate of sand waves in IJmuiden is larger than that in Rotterdam area for both the field  
347 data and model results. As discussed above, the migration of sand waves is mainly caused by residual currents  
348 and tidal asymmetries which typically result from different tidal components and phases. In the IJmuiden area,  
349 the phase between M2 and M4 is  $339^\circ \sim 0.84^\circ$ , under which the migration of sand waves is positive according  
350 to Figure 9(b). The migration direction controlled by the residual current is also positive, so the sand waves  
351 move at a fast rate in the direction of the residual current. On the other side, in the Rotterdam area, the phase of  
352 M2 and M4 is  $102^\circ \sim 130^\circ$ , resulting in a negative migration of sand waves based on Figure 11.

353 The negative effect exerted by tidal phases counteracts the positive effect induced by residual currents, thus the  
354 migration rates are smaller in the Rotterdam area compared with that in the IJmuiden area. Take the parameters  
355 of R1 site for example in the Table 2, the negative effect induced by the tidal components' interaction  
356 overwhelms the positive effect caused by residual currents, making the sand waves migrate in the opposite



357 direction of the residual current (0.019 m/s). Then focus on R2 site, the residual current (0.04 m/s) is much  
358 larger than that in R1 site. The positive effect overwhelms the negative effect which may makes the sand  
359 waves migrate in the direction of the residual current. The difference of quantitative results between the  
360 migration rate of observed data and model data in IJmuiden area may be caused by the approach of evaluation  
361 of the migration. Table 2 also shows that in Rotterdam area crest and trough migration is almost the same while  
362 in the IJmuiden area there is quite difference in crest and trough migration. The mean migration rate of the  
363 observed data was obtained by the average of the crest and trough migration which may be underestimated  
364 compared with the model result which was evaluated by FFT analysis.

365 According to van Santen et al. (2011), the wavelength of sand waves decreased with increasing grain size  
366 which concurs with the results of our model. However, in their study, the larger tidal current amplitudes will  
367 induce the decrease of tidal sand wavelength which is opposed to ours. They only used the bedload transport to  
368 evaluate the sediment transport while the suspended sediment transport was also included in our model. The  
369 suspended transport may induce the longer wavelength of sand waves as discussed by Borsje et al. (2014). So,  
370 when the tidal current amplitude was large, suspended transport would be the dominant transport which may  
371 cause the longer wavelength.

372 The role of wind waves was not included in our model which may influence the prediction of characteristics of  
373 sand waves (Campmans et al., 2017). The interaction between surface wave and the tidal current may results  
374 the change of the flow fields. The nonlinear of wave action in shallow water should not be neglected on the  
375 sediment transport. So, a next step is to study the role of wind waves in the formation and evolution of sand  
376 waves with the process-based model Delft3D.

## 377 **5. Conclusions**

378 The sand wave dynamics were investigated with a process-based numerical model, Delft3D. Two sites in the  
379 North Sea were used for model validations. The sensitivities of modeled sand wave growth and migration to  
380 environmental settings and model parameters, namely, sediment sizes, tidal currents and tidal asymmetries, and  
381 bed slope effects, were examined systematically.

382 The conclusions can be summarized as follows:

383 1. The numerical model is effective to research the characteristics of the sand waves, including wave length  
384 and migration rate of sand waves. The bed slope factor is critical to predict the wavelength and migration  
385 rate of sand waves and need to be calibrated and is related to the local sediment grain size. In general,  
386 when the grain size is finer, a higher slope factor can be adopted to perform a better result in predicting the  
387 characteristics of sand waves.

388 2. Sensitive analysis of environment parameters shows that the preferred wavelength increased and the  
389 corresponding growth rate decreased with the increase of bed slope factors. In general, the sand waves  
390 were absent when the tidal current was too large and the sediment was too fine. The preferred wavelength  
391 of sand waves decreased with the increase of the sediment diameter and the decrease of tidal current. For  
392 coarser sediment, the growth rate of sand waves will increase first and then decrease with the increase of  
393 the tidal current. For finer sediment the growth rate will always decrease with the increase of the tidal  
394 current.

395 3. The interaction between the residual current and the tidal components can explain the migration of sand  
396 waves. Considering the phases between different tidal components the sand waves may migrate in the  
397 opposite of the residual current which is agreement with the observed field data. Residual currents can also  
398 influence the fastest growing modes and a larger residual current can even up to total disappearance of the

399 sand waves. A larger residual current may induce a longer wave length of sand waves. The  $L_{FGM}$  is varied  
400 with different phases between the different tidal components.

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483 **Appendix A. Hydrodynamics model**

484 In the  $\sigma$  coordinate system, the 2DV shallow water equations read,

485 
$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + \frac{\omega}{(H + \zeta)} \frac{\partial u}{\partial \sigma} = -\frac{1}{\rho_w} P_u + F_u + \frac{1}{(H + \zeta)^2} \frac{\partial}{\partial \sigma} \left( \nu_T \frac{\partial u}{\partial \sigma} \right) \quad (6)$$

486 
$$\frac{\partial \omega}{\partial \sigma} = -\frac{\partial \zeta}{\partial t} - \frac{\partial [(H + \zeta)u]}{\partial x} \quad (7)$$

487 where  $x$  denotes the horizontal direction,  $\sigma$  denotes the vertical direction related to  $\sigma$  coordinate system,

488  $u$  is the horizontal velocity and  $\omega$  is vertical velocity related to  $\sigma$  coordinate system,  $\rho_w$  is the water

489 density,  $H$  is the water depth below the reference plane,  $\zeta$  is the free water surface elevation,  $\nu_T$  is the

490 vertical eddy viscosity which can be evaluated by a specified turbulence closure model,  $P_u$  is the hydrostatic

491 pressure gradient,  $F_u$  in the momentum equation represents the horizontal Reynold's stress, which is the

492 gradient of shear stress and has the form:

493 
$$F_u = \frac{\partial \tau_u}{\partial x} \quad (8)$$

494 The shear stress  $\tau_u$  is,

495 
$$\tau_u = 2\nu_H \left( \frac{\partial u}{\partial x} + \frac{\partial u}{\partial \sigma} \frac{\partial \sigma}{\partial x} \right) \quad (9)$$

496 where  $\nu_H$  is the horizontal eddy viscosity.

497 Under the so-called "shallow water assumption", the vertical momentum equation reduces to hydrostatic

498 pressure equation. The hydrostatic pressure gradient can be expressed as:

499 
$$P_u = -\rho_w g \frac{\partial (H + \zeta)}{\partial x} \quad (10)$$

500 where  $g$  is the gravity acceleration.

501 About the determination of eddy viscosity, most studies on sand waves using stability models adopted a  
 502 constant or highly simplified relationship, while Borsje et al. (2013) showed that the  $k - \varepsilon$  turbulence model  
 503 could improve the prediction of wavelength and migration rate of sand waves considerably. Turbulent kinetic  
 504 energy  $k$  and energy dissipation  $\varepsilon$  can be solved by transport equations. For details on turbulence model  
 505 and  $k - \varepsilon$  equations see Rodi (1984) and Burchard et al. (2008).

506 At the surface, a no-stress condition is applied and vertical velocity  $\omega$  is set to zero:

$$507 \quad \rho_w \frac{v_T}{(H + \zeta)} \frac{\partial u}{\partial \sigma} \Big|_{\sigma=0} = 0, \quad \omega \Big|_{\sigma=0} = 0 \quad (11)$$

508 At the sea bed which is the closed boundary, the tangential shear stress is calculated based on a logarithmic law  
 509 of the wall:

$$510 \quad \tau_b \equiv \rho_w \frac{v_T}{(H + \zeta)} \frac{\partial u}{\partial \sigma} \Big|_{\sigma=-1} = \rho_w u_* |u_*|, \quad \omega \Big|_{\sigma=-1} = 0 \quad (12)$$

511 where  $\tau_b$  is the bed shear stress,  $u_*$  is the friction velocity which is determined by the logarithmic-law  
 512 related to the velocity gradient at the lowest grid point.

513 Under the assumption of a logarithmic velocity profile, the relationship between depth-averaged velocity and  
 514 friction velocity is as follows:

$$515 \quad u_r = \frac{u_*}{\kappa} \ln \left( 1 + \frac{H + \zeta}{ez_0} \right) \quad (13)$$

516 where  $\kappa$  is the Kármán's constant,  $z_0$  is the bed roughness height, which reads:

$$517 \quad z_0 = \frac{H + \zeta}{e \frac{1 + \kappa C_D}{\sqrt{g}} - e} \quad (14)$$

518 where  $C_D$  is Chézy coefficient.

## 519 **Appendix B. Sediment transport and morphological evolution model**



520 Sediment can be transported as bedload and suspended load. The coordinate system used to determine the  
 521 sediment transport is based on Cartesian  $z$  coordinate system.

522 The vertical velocity  $\omega$  solved in hydrodynamics model is related to the moving  $\sigma$  plane (Lesser et al.,  
 523 2004). The vertical velocity  $w$  in the  $z$  coordinate system can be expressed in the horizontal velocities,  
 524 water depths, water levels, and vertical coordinate velocity according to:

$$525 \quad w = \omega + u \left( \sigma \frac{\partial(H + \zeta)}{\partial x} + \frac{\partial \zeta}{\partial x} \right) + \left( \sigma \frac{\partial(H + \zeta)}{\partial t} + \frac{\partial \zeta}{\partial t} \right) \quad (15)$$

526 Suspended sediment transport is modeled by the 2DV advection-diffusion equation, which reads:

$$527 \quad \frac{\partial c}{\partial t} + \frac{\partial(cu)}{\partial x} + \frac{\partial(w - w_s)c}{\partial z} = \frac{\partial}{\partial x} \left( \epsilon_{s,x} \frac{\partial c}{\partial x} \right) + \frac{\partial}{\partial z} \left( \epsilon_{s,z} \frac{\partial c}{\partial z} \right) \quad (16)$$

528 where  $c$  is the suspended sediment concentration,  $w$  is the vertical velocity related to  $z$  coordinate  
 529 system,  $w_s$  is the sediment settling velocity,  $\epsilon_{s,x}$  and  $\epsilon_{s,y}$  are horizontal and vertical eddy diffusivities,  
 530 respectively, which are assumed to be equal to the horizontal and vertical eddy viscosity.

$$531 \quad w_s = \frac{10\nu}{d_{50}} \left( \sqrt{1 + \frac{0.01(\rho_s / \rho_w - 1)gd_{50}}{\nu^2}} - 1 \right) \quad (17)$$

532 where  $\nu$  is the water viscosity coefficient,  $\rho_s$  is sediment density.

533 According to Van Rijn (2007), sediment transport above the reference height ( $a = 0.01H$ ) is regarded as  
 534 suspended load, which can be calculated as:

$$535 \quad S_s = \int_a^{(H+\zeta)} uc \, dz \quad (18)$$

536 where  $S_s$  is suspended load transport.

537 The bedload can be calculated using the approximation method proposed by Van Rijn (1993), which reads:

$$538 \quad S'_b = 0.5 \rho_s d_{50} u_*' d_*^{-0.3} T \quad (19)$$

539 where  $S'_b$  is the bedload transport without considering slope effect,  $u'_*$  is the effective bed-shear velocity,

540  $d_*$  is non-dimensional particle diameter and  $T$  is non-dimensional bed-shear stress.

$$541 \quad u'_* = u_* \mu_c^{0.5} \quad (20)$$

542 where  $\mu_c$  is efficiency factor current,

$$543 \quad \mu_c = \frac{f'_c}{f_c} \quad (21)$$

544 where  $f'_c$  and  $f_c$  are grain-related and total current-related friction factors, respectively.

$$545 \quad f'_c = 0.24 \left[ \log_{10} \left( \frac{12(H + \zeta)}{3d_{90}} \right) \right]^{-2} \quad (22)$$

$$546 \quad f_c = 0.24 \left[ \log_{10} \left( \frac{12(H + \zeta)}{k_s} \right) \right]^{-2} \quad (23)$$

547 where  $d_{90}$  is 90% sediment passing size,  $d_{90} = 1.5d_{50}$ ,  $k_s$  is the Nikuradse roughness length,  $k_s = 30z_0$ .

$$548 \quad d_* = d_{50} \left[ \frac{(\rho_s / \rho_w - 1)g}{v^2} \right]^{1/3} \quad (24)$$

$$549 \quad T = \frac{\mu_c \tau_b - \tau_{cr}}{\tau_{cr}} \quad (25)$$

550 where  $\tau_{cr}$  is critical bed shear stress, which reads,

$$551 \quad \tau_{cr} = (\rho_s - \rho_w) g d_{50} \theta_{cr} \quad (26)$$

552 where  $\theta_{cr}$  is the threshold parameter calculated according to van Rijn (1993) as a function of the

553 non-dimensional grain size.

554 Bedload transport should be adjusted by bed slope effect as it is affected by bed level gradients. The adjusted

555 bedload transport is defined as:

$$556 \quad S_b = \alpha_s S'_b \quad (27)$$

557 where  $S_b$  is the adjusted bedload transport and  $\alpha_s$  is a correction parameter for slope effect, which can be  
 558 calculated with Bagnold (1966):

$$559 \quad \alpha_s = 1 + \alpha_{bs} \left( \frac{\tan(\phi)}{\cos\left(\tan^{-1}\left(\frac{\partial z}{\partial s}\right)\right) \left(\tan(\phi) - \frac{\partial z}{\partial s}\right)} - 1 \right) \quad (28)$$

560 where  $\phi$  is the repose angle of sediment,  $\partial z / \partial s$  is the bed slope in the direction along to the unadjusted  
 561 bedload transport vector,  $\alpha_{bs}$  is a user-defined tuning parameter.

562 The bed level change can be evaluated by sediment continuity equation (Exner equation):

$$563 \quad (1 - p_{or}) \frac{\partial z_b}{\partial t} + \frac{\partial}{\partial x} (S_b + S_s) = 0 \quad (29)$$

564 in which  $p_{or} = 0.4$  is the bed porosity.

# Modeling the formation and migration of sand waves: The role of tidal forcing, sediment size and bed slope effects

Zhenlu Wang<sup>1</sup>, Bingchen Liang<sup>1,2,\*</sup>, Guoxiang Wu<sup>1,2</sup>, Borsje B. W.<sup>3</sup>

## Highlights:

- (1). After calibration of the slope parameter for a specific location, the sand wave length is reproduced good compared to field observations
- (2). The migration rate is considerably better predicted with the Delft3D model compared to the Blondeaux model. This is because of the more accurate turbulence formulation ( $k-\epsilon$ ) in the Delft3D model, while the constant eddy viscosity overestimates the transport rates and thereby the migration rates.
- (3). Residual currents and higher harmonics not only influence the migration rates but also the fastest growing modes (even up to total disappearance of the sand waves)

Dear Editor,

I am submitting our manuscript entitled “**Modeling the formation and migration of sand waves: The role of tidal forcing, sediment size and bed slope effects**” by Zhenlu Wang, Bingchen Liang, Guoxiang Wu and Borsje B. W. to *Continental Shelf Research* as a full-length article. We would be grateful if the manuscript could be reviewed and considered for publication in *Continental Shelf Research*.

This manuscript has not been published or presented elsewhere in part or in entirety and is not under consideration by another journal. All the authors have approved the manuscript and agree with submission to your esteemed journal. There are no conflicts of interest to declare.

**Thank you and best regards !**

Sincerely,

Bingchen Liang

College of Engineering

Ocean University of China

Qingdao 266100

China

Email: bingchen@ouc.edu.cn

Tel: 86-532-66783719

Fax: 86-532-66781550