Capillary origami of micro-machined micro-objects: Bi-layer conductive hinges

CrossMark

A. Legrain [a,*], J.W. Berenschot [a], N.R. Tas [a], L. Abelmann [a,b]

[a] MESA + Institute for Nanotechnology, University of Twente, Enschede, The Netherlands
[b] KIST Europe, Saarbrücken, Germany

ARTICLE INFO

ABSTRACT

Recently, we demonstrated controllable 3D self-folding by means of capillary forces of silicon-nitride micro-objects made of rigid plates connected to each other by flexible hinges (Legrain et al., 2014). In this paper, we introduce platinum electrodes running from the substrate to the plates over these bendable hinges. The fabrication yield is as high as $(77 \pm 2)\%$ for hinges with a length less than 75 μm. The yield reduces to $(18 \pm 2)\%$ when the length increases above 100 μm. Most of the failures in conductivity are due to degradation of the platinum/chromium layer stack during the final plasma cleaning step. The bi-layer hinges survive the capillary folding process, even for extremely small bending radii of 5 μm, nor does the bending have any impact on the conductivity. Stress in the different layers deforms the hinges, which does not affect the conductivity. Once assembled, the conductive hinges can withstand a current density of $(1.6 \pm 0.4) \times 10^6$ A/cm^2. This introduction of conductive electrodes to elastocapillary self-folded silicon-based micro-objects extends the range of their possible applications by allowing an electronic functionality of the folded parts.

1. Introduction

Pushed by the semiconductor industry, lithography-based fabrication techniques now include new methods other than UV illumination: mask-less projection of laser light, electrons and ion beams, etc. Extremely precise, these methods are nevertheless inherently two-dimensional. They essentially consist on patterning one single layer after the other. Precise three-dimensional (3D) shapes can be obtained using anisotropic etching of silicon, but they are constrained in depth or height because of a limited substrate thickness. Consequently, high aspect-ratio structures are difficult to fabricate using techniques based on lithography and etching [2,3].

One solution is to fold structures out of the plane of fabrication, preferably by self-folding. Self-folding broadly refers to the self-assembly of interconnected parts that fold themselves into predefined shapes without the need of active human control. The assembly is either triggered by external stimuli (e.g., pH or temperature variation), enabled by external forces such as magnetic forces, or by internal forces such as pre-stressed layers [3–5]. Three-dimensional assembly of micro-/nano-objects is possible using self-folding, where inherently two-dimensional micro-fabrication techniques have been shown to be inadequate [2].

Capillary origami designates the self-folding of flexible elastic material using surface tension as the enabling force [6,7]. Capillary origami is a particularly interesting tool for micro-fabrication, since at

small scales, surface forces dominate over bulk forces such as gravity [8,9]. This technique was first employed to assemble silicon-based micro-objects by Syms et al., who used melting solder to assemble hingeless silicon objects with integrated metal pads in flaps [10,11]. The scope of such solder assembly was then extended by Gracias et al., who demonstrated the folding of complex structures (cubes, pyramids...) of micrometer and nanometer sizes fabricated using standard lithography and deposition techniques [12–15]. Applications of such structures range from three-dimensional micro-opto-electro-mechanical systems (MOEMS) [16–18] to RF nano-antennae [19].

We have used capillary origami to fold silicon nitride micro-objects. Unlike in solder assembly, the structures are necessarily hinged and their folding relies on the deformation of thin flexible silicon nitride plates, therefore the method can also be called elastocapillary folding. The folding is driven by the surface tension of water, which is simply manually deposited [20,21] or brought to the origami pattern through a tube at its center [1]. The final shape of the object is predefined by the patterning of rigid silicon nitride plates that form the body of the 3D object, linked to each other by hinges. Objects remain assembled due to strong stiction between the flaps [20,1] or by designing complex stop-programmable hinges [21].

The applications of this self-folding technique are limited to situations where a passive mechanical structure is required that extends far above the wafer surface. To extend the range of applications, it would be extremely useful if electrical connection can be made to the moving parts of the folding structure. Electrical conductance over hinges has been demonstrated in experiments where folding is

* Corresponding author.
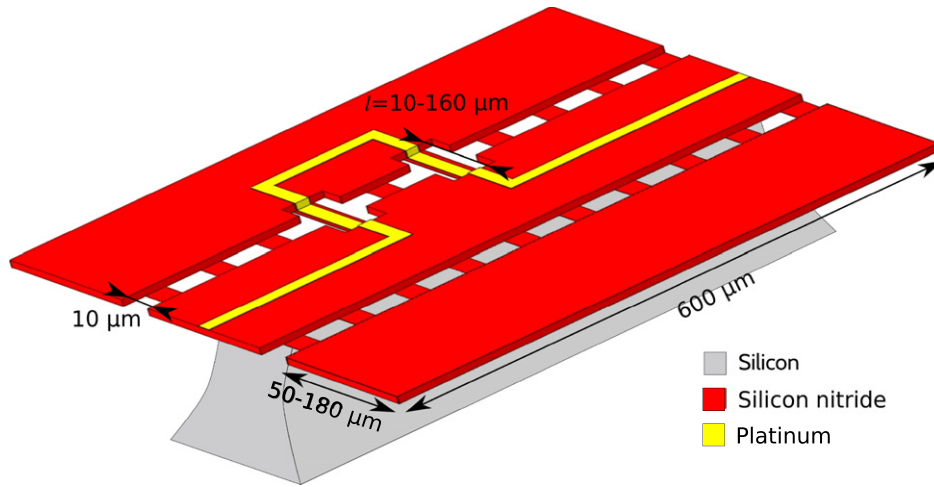E-mail address: a.b.h.legrain@utwente.nl (A. Legrain).

Fig. 1. Concept image of a typical test structure. Two free hanging silicon nitride flaps are connected to the central part by means of thin flexible silicon nitride hinges. The central part is fixed, resting on a silicon pillar. In this example, two decoupled bi-layer hinges run along the center of the structure. The length l of the metalized junctures can be tuned, while the plain SiRN hinges are always 10 µm long. The metal parts run to contact pads on the outside of the structures.

achieved by magnetic lifting [22,23] or stress gradients [24,25]. In this research we have combined electrical connectivity with folding by capillary forces. Fig. 1 shows the type of structures used to prove the feasibility of conductive hinges. Structures based on our previous publications [20,1] are extended with platinum wires that run from the substrate towards the flap via bendable hinges. After elastocapillary folding, a three-dimensional triangular prism structure ('Toblerone' [26]) is realized, which has electrical wiring on its movable parts.

The requirements for the electrode material are a high electrical and thermal conductivity to reduce power loss and unwanted heating, in combination with a low Young's modulus to keep the hinges as flexible

as possible. Gold would therefore be an excellent candidate. However, as possible application of elastocapillary folding we consider out-of-plane flow sensors. Since these sensors need to operate at high temperature, additional requirements such as high melting point ($T_m = 1768$ °C) [27], high temperature coefficient of resistance (TCR, 3.1/°C to 4.3×10^{-3}/°C) over a wide temperature range (0 °C to 800 °C) [28] and high oxidation resistance are crucial. In this work, we therefore selected platinum as the material of choice.

In the following we will discuss shortly the stress induced by folding and the effect on electrical resistance and explain the fabrication process. Subsequently we demonstrate that conductivity can be preserved
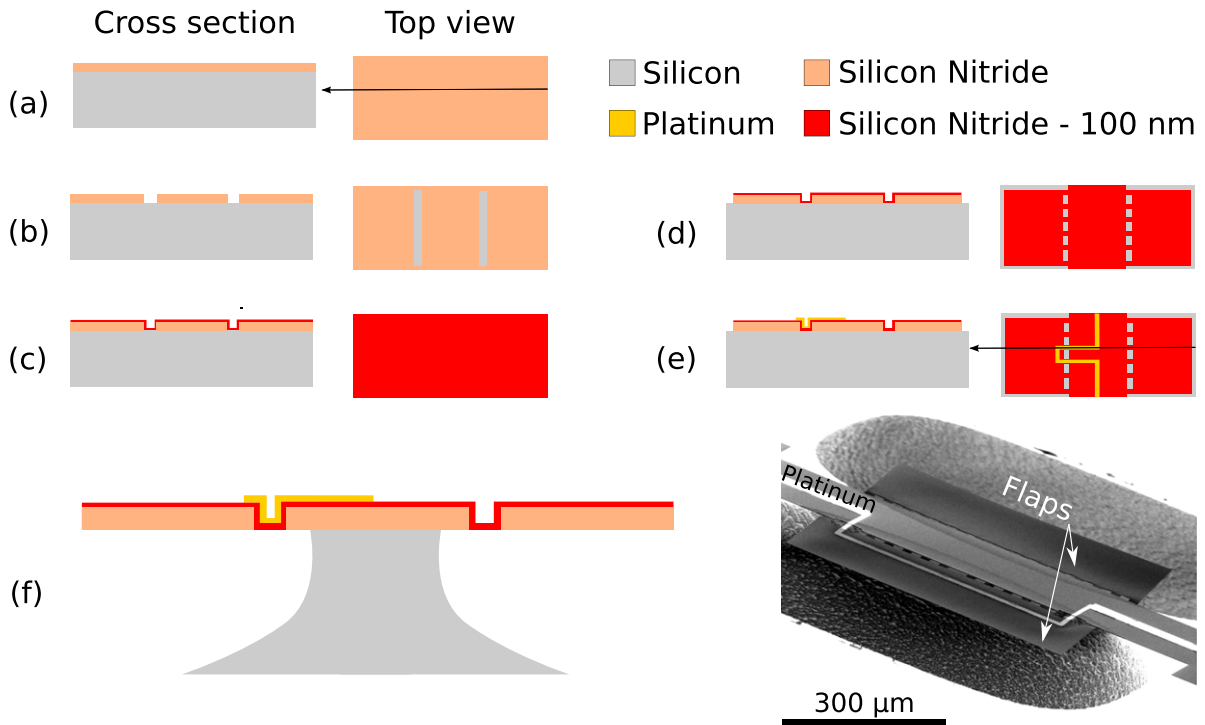

Fig. 2. Fabrication steps for silicon nitride (SiRN) structures with integrated conductive hinges capable of being folded out of the plane by capillary forces. (a): Deposition of SiRN by LPCVD. (b): First lithographic step; definition of the hinges. SiRN is either dry or wet etched. (c): Second deposition of SiRN. (d): Overall definition of the structures by a lithography step and subsequent dry etching. (e): Sputtering of metal followed by a standard lift-off procedure. (f): Release of the flexible objects by semi-isotropic etching of silicon. The SEM picture on the right hand side shows an example of a final structure.

under folding, and focus on the yield of the process and discuss various causes of failure during fabrication.

2. Theory

The theory of elastocapillary folding has been explained in [20,1]. Here we add the stress induced in the hinges after folding, and the consequences for the electrical resistivity of the metal layer on the hinge.

After folding, the flexible hinges are strained with a value $\epsilon = e / R$, where R is the radius of curvature of the neutral plane inside the plate and e is the distance from the neutral plane. The position of the neutral plane in a multilayered structure with respect to the bottom of the SiRN layer is [29]:

$$z_n = \frac{t_{SiRN} + t_{Pt}}{2} - \frac{t_{SiRN} t_{Pt} (E_{SiRN} - E_{Pt})}{2(E_{SiRN} t_{SiRN} + E_{Pt} t_{Pt})} \qquad (1)$$

where t and E are the thickness and Young's modulus of the SiRN and platinum layers ($t_{SiRN} = 100$ nm, $E_{SiRN} = 290$ GPa, $t_{Pt} = 75$ nm or 150 nm, $E_{Pt} = 168$ GPa). In the experiments in this paper the neutral plane is located inside the SiRN layer for $t_{Pt} = 75$ nm ($z_n = 77$ nm) and just inside the Pt layer for $t_{Pt} = 150$ nm ($z_n = 108$ nm). The platinum layer is therefore mostly under a negative (compressive) strain.

The strain in metal results in a change in resistivity through geometric change and the piezoresistance effect. This effect is characterized by the strain gauge factor:

$$GF = \frac{\Delta R / R}{\epsilon}. \qquad (2)$$

For platinum thin films, values of GF between 2.7 and 3.8 have been reported [30,31].

3. Experimental section

3.1. Fabrication

Fig. 2 shows the necessary steps to micro-machine test structures such as the one depicted in Fig. 1. Fabrication starts with a conformal deposition by low pressure chemical vapor deposition (LPCVD) of a low stress thick silicon nitride (SiRN, Fig. 2(a)). Three layers, 499, 797 and 1083 nm, were deposited on different samples to investigate the influence of the step height on the fabrication yield.

A first lithographic step follows to etch away the thick SiRN at the location of the future hinges. For this step, both dry etching and wet etching were used in order to study the impact on the folding of the shape of the transition between the hinges and the flaps. Dry etching is performed after patterning only the photoresist, and yields a straight profile in the thick SiRN layer. Wet etching requires first the (≈50 nm). The polysilicon layer is patterned using dry etching, and subsequently used as a mask to etch SiRN in isotropic wet etchant, either hydrogen fluoride (HF) or phosphoric acid (H_3PO_4). The resulting shape is a smooth circular transition, inherent to isotropic etching. Once silicon is reached, the thin polysilicon layer can be stripped in tetramethylammonium hydroxide (TMAH).

The next step is the deposition of a thin flexible layer of SiRN by LPCVD, Fig. 2(c). A thickness of 100 nm has proven to provide both flexibility and solidity to the hinges [20,1].

The second lithography step follows to define the geometry of the patterns, Fig. 2(d). In this design, combinations of different flap widths (50–180 µm), leading to different final folding angles, with different bi-layer hinge lengths (10–160 µm) were tried out. The location of the conductive hinges was also varied, and reference structures (structures without SiRN below metal or with no metal) were included.

Next, platinum is deposited and patterned on top of the SiRN by a standard lift-off procedure, Fig. 2(e). Chromium (≈10 nm) is used as an adhesive layer. Two different layers of platinum, 75 nm and 150 nm,
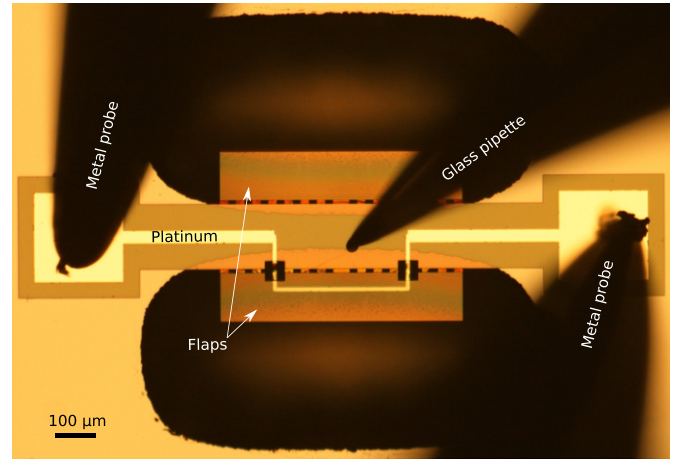

Fig. 3. Photograph of experimental setup, top view. The foldable object is placed at the center. A glass syringe is positioned on top of the structure by means of a translation table, and is used to apply water for the folding process. Metal probes are in contact with metal pads on both sides.

were sputtered on different samples at a pressure of 6.6×10^{-6} bar. The thicknesses of the final metal layers were checked using a mechanical surface profiler. Lift-off was achieved by first coating and patterning a 3.5 µm thick photoresist with no post-development baking, to keep the profile of the resist straight. After metal deposition, the wafers were placed in an ultrasonic bath and immersed in sequence in acetone and in isopropanol.

Finally, a last lithography step is performed to protect the SiRN objects during their release, Fig. 2(f). Under-etching of silicon is performed in semi-isotropic sulfur hexafluoride (SF_6) gas etchant. Since all the structures should be released simultaneously, the mask openings are designed to be of equal size and at the same distance from the final etch location.

Moreover, it should be noted that the etching selectivity between silicon and SiRN during this step is around 1000. The initial thickness of the SiRN layer from which the hinges are patterned (Fig. 2(c)) was therefore chosen accordingly thicker (≈20 nm), to compensate for thinning during the silicon etch step.

Finally, the photoresist is stripped in oxygen plasma for 2 h once the objects are released from the substrate.

3.2. Residual stress measurements

The stress in the deposited thin layers was measured using the wafer curvature method. The curvature of a dummy wafer is measured before and after deposition of the thin film by means of a surface profiler. Stress is then calculated using the Stoney equation [32]. In the case of SiRN, for which material is deposited on both sides of wafers by LPCVD, one side of the wafer is stripped before measuring the radius of curvature. The thickness of the substrates is measured using a dedicated tool (Heidenhain measuring station) while the thickness of the SiRN and sputtered metals are checked by ellipsometry and mechanical surface

Table 1
Fabrication yield of conductive hinges as a function of hinge length, before folding. A correlation between length and conductivity is obvious. Statistics made with N structures from 7 wafers, which were fabricated with different combinations of thicknesses for the thick SiRN flaps (499, 797 or 1083 nm), thicknesses of platinum (75 or 150 nm) and etching method for hinges (dry or wet, Fig. 2(b)). The latter parameters had no observable significant impact on the yield.

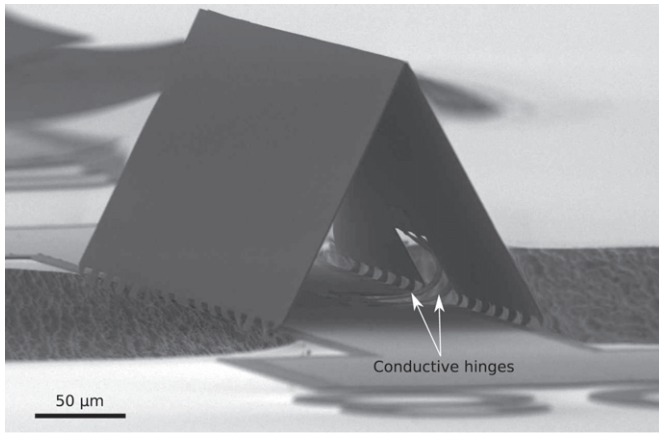| Hinge length	Conductivity yield	N
$10 \leq l \leq 75$ µm	$(77 \pm 2)\%$	541
$75 < l \leq 100$ µm	$(26 \pm 4)\%$	129
$l > 100$ µm	$(18 \pm 2)\%$	391

Fig. 4. Folded micro-object with conductive bi-layer hinges that allow for electrical connection to the folded flaps. The structure is 600 μm long and the flaps are 80 μm wide. The plain SiRN hinges are 10 μm long and 30 μm wide, while the two bi-layer hinges are 75 μm long.

profiler, respectively. The values found in the literature are used for Young's moduli and Poisson's ratios.

3.3. Experimental setup

Self-folding experiments and the electrical characterization are carried out in the same setup, shown in Fig. 3. Resistivity of the structures is measured by a multimeter in situ during assembly by means of two metal probes placed on metal contact pads on both sides of the structures. For breakdown measurements, a voltage source was used to force a current through the bi-layer hinges.

A glass micro-pipette of 10 μm diameter is positioned on top of the micro-origami pattern by an accurate x–y translation table. A low volume syringe (~10 μL), filled with ultra pure water, is used to propel a drop out of the micro-pipette. A hydrophobic polytetrafluoroethylene (teflon) coating is applied on the micro-pipette to avoid wetting of the outside of the glass pipette. Two cameras on top and on the side are used to monitor the folding process. Depending on the volume of liquid deposited, folding takes 10 to 120 s.

4. Results and discussion

4.1. Fabrication yield

As described in Section 3.1, various combinations of metal thicknesses and step heights have been fabricated. Moreover, the transition between the hinges and the flaps was machined by dry etching, resulting in a sharp step, as well as wet etching, resulting in a smooth circular transition.

None of the aforementioned parameters have any significant influence on the fabrication yield. The only feature that proved to have a statistically significant impact on the yield is the length of the composite hinges. Table 1 summarizes our results. Short hinges, with a length l below 75 μm, are the most robust, with a yield of $(77 \pm 2)\%$. The variance of this value is calculated from the number of structures tested assuming a binomial distribution. The yield drastically drops for longer junctures, and less than one out of five structures are conductive when $l \geq 100$ μm. The reasons will be made clear in Section 4.3. Short hinges are therefore the best option for the design of conductive bi-layer junctures.

4.2. Folding and electrical characterization of conductive hinges

Fifty structures with conductive bi-layer hinges were successfully folded and monitored using the setup presented in Fig. 3. Fig. 4 shows an example of a folded micro-object, with the two conductive hinges inside the right-hand side flap. Neither the folding angle nor the length of the bi-layer hinges has any impact on the folding.

As shown in Fig. 5, the conductive hinges do not break even when they are rotated by 180°. Picture (b) shows 10 μm bi-layer hinges that are fully rotated and yet do not break, therefore resulting in a radius of curvature $R \sim 5$ μm. These are the shortest that can be designed since they are as long as the plain SiRN hinges. Using the theory of Section 2, we estimate that at the smallest radius of curvature, the maximum tensile stress in the SiRN layer is 5.6 GPa, which is about the yield strength of (5.5 ± 8) GPa reported by Kaushik et al. [33]. It seems therefore likely that much smaller bending radii will lead to broken hinges.

The resistivity between the two metal probes was measured to be on the order of 110 Ω and 80 Ω for a platinum thickness of 75 nm and 150 nm, respectively. The value varies by 25% from measurement to measurement, depending on the total length of the metal and the exact landing position of the probes. These values are higher than the theoretical values of respectively (64 ± 18) Ω and (32 ± 5) Ω, calculated from the bulk value of the resistivity of platinum, $\rho_b = 1.06 \times 10^{-7}$ Ωm. The difference can be explained by additional contact resistances as well as the fact that the resistivity of platinum thin films is known to be higher than the bulk value. A factor of five is reported for 20 nm platinum thick films [34,35].

Even at the smallest radius of curvature, no difference in conductance could be observed upon folding. We estimate that the compressive strain in the platinum layer is smaller than $\epsilon = -2.0\%$, leading to a maximum change in resistance of the hinge of 8% (Eq. (2)). Since the hinge resistance is only 1% of the total resistance of the circuit, the effect is indeed too low to be easily observed.

Our results are in agreement with recent publications on flexible electronics that also achieve very small radii of curvature. Using ultra-thin substrates (<2 μm), fully printed organic TFTs could be folded to very small bending radii of about $R \sim 5$ μm [36]. Bendable organic solar cells with $R \sim 10$ μm [37], as well as polymer-based LEDs with $R < 10$ μm [38], have also been reported.
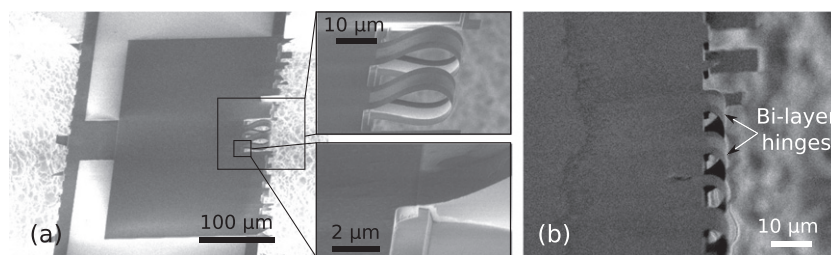

Fig. 5. (a): Extreme folding of 50 μm long hinges over an angle of 180. Top inset: a zoom on the formed loop, no deformation of the bi-layer junctures can be observed. Bottom inset: a zoom on the transition between the flaps and the bi-layer junctures. The standard sputtering method yields a nice conformal step coverage. (b): Even the shortest hinges of 10 μm length sustain a 180 rotation.
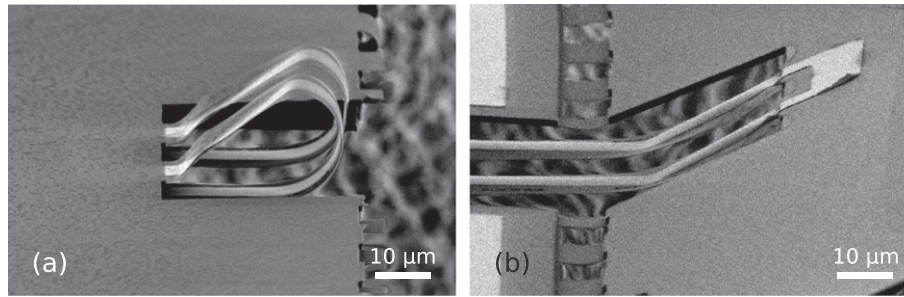
Fig. 6. Bi-layer hinges from two distinct structures. (a): l = 100 μm, was folded by 180 using the setup shown in Fig. 3 (b): l = 105 μm, was not folded. The flap is tilted out of plane (around 30). Both structures are still conductive but show deformations due to stress mismatch between the different layers. In both cases only the parts located at the level of the flaps are deformed.

4.3. Hinge deformation

While short hinges show no deformation and are nicely curved after folding, see Fig. 5, longer junctures are deformed, see Fig. 6. The deformation occurs only at the level of the flaps, with the hinge section connected to the fixed central part remaining intact. In photograph (b) we can see that the flap is tilted out of the plane, which suggests that the deformation in the bi-layer junctures is due to this pre-folding. Despite the strong deformation, the hinges are still conductive.

Fig. 7 shows similar hinges without a metal top layer. These plain decoupled SiRN hinges are deformed in a similar way as the bi-layer version of Fig. 6. As before, the section of the hinges connected to the central part remains intact. The other side, connected to the flap, remains parallel to the thick SiRN plane for a short distance before the curvature starts.

Why this deformation occurs is unclear. We can imagine several possible explanations for this effect, none of which seem adequate. The flaps are up to ten times thicker than the hinges, and the stress mismatch between the two SiRN layers could explain the plastic deformation. However, there is no evidence of a thickness dependency of the residual stress in the SiRN layers in the literature. SiO_2 layers between the two SiRN layers or under the SiRN formed as the wafers are loaded in the LPCVD equipment are also candidates. But such layers would be a few thick, and it is unlikely that they could be the cause of such large bending. We also observed that during the stripping of the photoresist in the oxygen plasma, the last step of fabrication, the temperature rises and the photoresist melts, which results in folding of the flaps. As will be discussed in the next section, we observed the degradation of metal during this same step. Such degradation could be attributed to the diffusion of chromium, which starts at 300 °C. The stresses due to melting photoresist at rather high temperatures may cause permanent plastic deformation of the hinges. The correct explanation however is

probably a combination of the previously described factors, coupled with the complex three-dimensional shape of the flap–hinge system.

Next to curvature along the hinges, buckling in the transverse direction is present in the case of bi-layer hinges, see Fig. 6. This buckling is not visible in the case of plain SiRN hinges, Fig. 7. Since we suspect that this is caused by stress gradients, we measured the residual stress using a wafer curvature method, see Section 3.2. The SiRN shows tensile stress, $\sigma_{r\text{-SiN}} = (-169 \pm 25)$ MPa, whereas the stress is compressive for platinum, $\sigma_{r\text{-Pt}} = (452 \pm 70)$ MPa. This significant mismatch causes a stress gradient along the bi-layer stack, which could explain the observed deformation.

4.4. Failure mechanism

Deformations caused by stress can cause the bi-layer hinges to break, as shown in Fig. 8(a). However, this is the case only for very long hinges, typically $l > 150$ μm. Most of the long hinges lose conductivity because of a poor metal coverage following the final oxygen plasma cleaning step, as shown in Fig. 8(b). Why this deterioration of the platinum happens is unclear. Electrical discharge and heating during the two hour final oxygen plasma cleaning step seems to be the only plausible explanation.

Even if the hinges survive the fabrication process, they might fail if the current applied is too high. Therefore we measured the current passing through the hinges while applying an increasing potential difference at their terminals. The resistance increases at higher voltages, as is illustrated by a linear fit through the points below 1 V. It is likely that this increase is caused by an increase in temperature. The measured resistance, however, includes the series resistances and therefore cannot be used to estimate the temperature of the metal layers.

The bi-layer hinges lost conductivity during the experiments when a current density of $j = (1.6 \pm 0.4) \times 10^6$ A/cm^2 was passing through
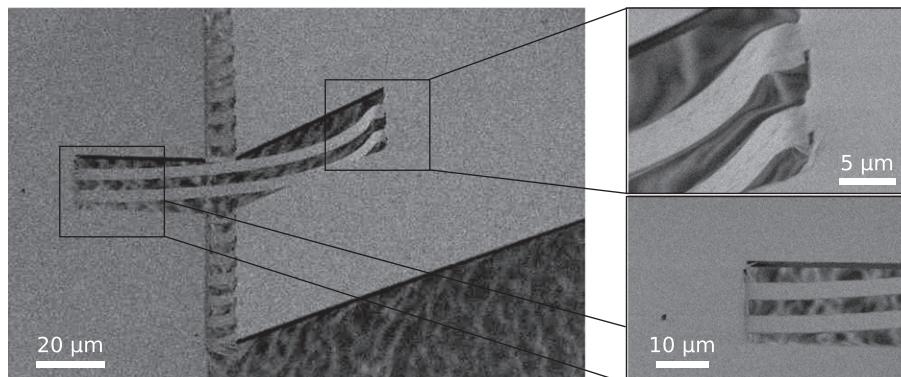

Fig. 7. Plain decoupled SiRN hinges. Although no metal is present on top of the junctures, the deformation already seen for bi-layer objects can be observed here. Most of the constraint seems to happen at the transition from the flap to the thin junctures.
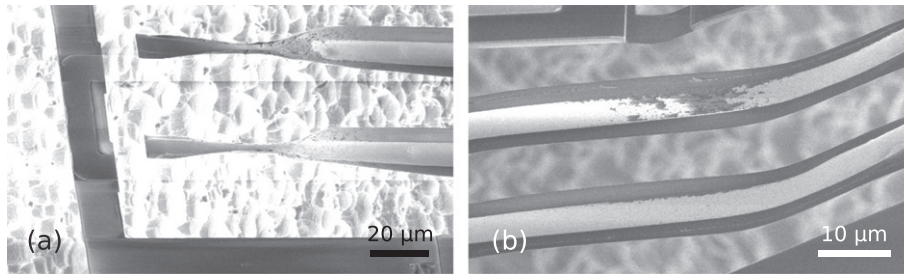
Fig. 8. Typical failures of long hinges that show no conductivity. (a): $l = 160\,\mu m$. For very long junctures, buckling and bending can cause the hinges to break at the transition with the flaps. (b): $l = 125\,\mu m$. Platinum burns at the middle of the hinges.

them, as shown in Fig. 9. This value of current density is close to the reported values in the literature at which electromigration has been observed [39,40]. However, the failure could also be attributed to the high temperature induced by the high current density, which is linked to the degradation of platinum films [41]. In this case the dominant mechanism is agglomeration, a surface-diffusion driven capillary process, whose effects are worsened by a mismatch of the temperature coefficients between the different layers [42]. It is possible to limit the impact of high temperatures on the hinges, through annealing or depositing platinum without an adhesion layer [28]. It has been proven that annealing of platinum thin film up to 300 °C yields reorganization of the atoms and a decrease of resistivity. At higher temperature (450 °C and up) the adhesion layer starts to diffuse into the platinum which creates hillocks and the conductivity quickly worsens [43–45]. On the other hand, such treatment modifies the stress in the layer, especially when there is a mismatch of thermal coefficient of expansion between the substrate and the thin film. Release of tensile stress [45] and large increase of compressive stress [44,46] have been reported. The stress is highly dependent on the conditions of deposition and the adhesion layer that has been used [47]. Further work is necessary to investigate the effect of annealing on the mechanical and electrical properties of bi-layer hinges.

5. Conclusion

We have successfully fabricated free-standing silicon-nitride plates with integrated platinum electrodes, which can be folded by elastocapillary forces into three-dimensional objects of several hundred μm in size. The electrodes are guided to the free-standing folding elements of the objects over two bi-layer silicon nitride/platinum hinges.
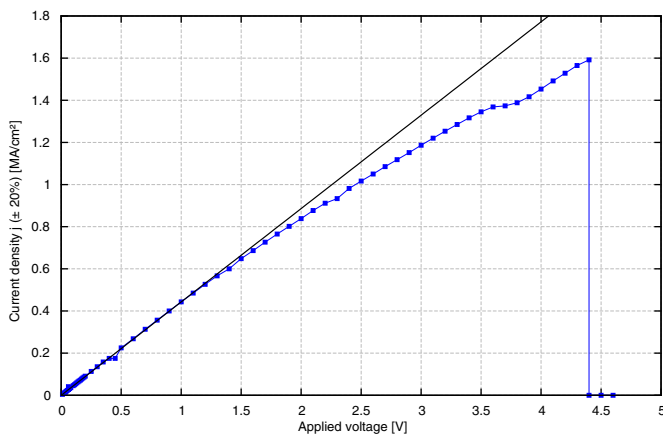
The resistance of the hinges does not change during the folding. A current density as high as $j = (1.6 \pm 0.4) \times 10^6$ A/cm² was passed through the hinges before conductivity was lost.

None of the fifty bi-layer hinges, fabricated with different lengths, broke during folding. Bending radii as small as 5 μm were achieved without mechanical failure or lost of conductivity.

If hinges break, it is during the fabrication process. Hinges with a length below 75 μm were mostly conductive with a high yield of (72 ± 2)%. The yield of the fabrication process drops to (18 ± 2)% for bi-layer hinges with a length above 100 μm. The yield is not influenced by the thickness of the connected flaps, the platinum layer, or the method used to etch the hinges' molds (wet or dry).

Short hinges with a length below 50 μm are nicely curved and show no deformation or burning after folding. Most of the mid-sized hinges, from 50 to 75 μm, do not break, but do show visible deformations such as buckling and bending. Similar deformations are observed even when there is no metal on top of the hinges. Although made of the same material as the hinges, the connected flap is folded out of plane by around 30°.

The failure mechanism during fabrication is either breaking or burning of the top platinum layer. In the case of long hinges with a length over 100 μm, the stress mismatch in the layers causes the breaking of the hinges at their connection with the flaps in rare cases. The majority of failures is, however, due to burning of the platinum top layer which is observed after the final oxygen plasma cleaning step, leading to a loss of conductivity.

Elastocapillary folding is a powerful technique that allows the 3D assembly of silicon based objects of mm size, something which is known to be difficult to accomplish using inherently two-dimensional fabrication techniques. However, the self-folding method has lacked applications so far. We believe that these conductive hinges will extend the application scope of elastocapillary folded structures, with possible applications to out-of-plane sensing, high aspect ratio coils, or even 3D electronics.

Fig. 9. Current–voltage characteristic of the bi-layer conductive hinges. The conductivity breaks down at $j = (1.6 \pm 0.4) \times 10^6$ A/cm². The solid black line is a linear fit through the points below 1 V.

References

[1] A. Legrain, T.G. Janson, J.W. Berenschot, L. Abelmann, N.R. Tas, J. Appl. Phys. 115 (2014) 214905.
[2] M.J. Madou, Fundamentals of Microfabrication, CRC Press, New York, 1997.
[3] T.G. Leong, A.M. Zarafshar, D.H. Gracias, Small 6 (2010) 792–806 (ISSN 1613–6829).
[4] V.B. Shenoy, D.H. Gracias, MRS Bull. 37 (09) (2012) 847–854 (ISSN 1938–1425).
[5] L. Ionov, Macromol. Chem. Phys. 214 (2013) 1178–1183 (ISSN 1521–3935).
[6] R. Syms, E.M. Yeatman, V.M. Bright, G.M. Whitesides, J. Microelectromech. Syst. 12 (2003) 387–417 (ISSN 10577157 (ISSN)).

[7] C. Py, P. Reverdy, L. Doppler, J. Bico, B. Roman, C.N. Baroud, Phys. Rev. Lett. 98 (2007) 156103 (ISSN 00319007 (ISSN)).
[8] M. Boncheva, D. Bruzewicz, G. Whitesides, Pure Appl. Chem. 75 (2003) 621–630.
[9] B. Roman, J. Bico, J. Phys. Condens. Matter 22 (2010) 493101.
[10] R. Syms, E. Yeatman, Electron. Lett. 29 (1993) 662–664 (ISSN 0013–5194).
[11] R. Syms, J. Microelectromech. Syst. 4 (1995) 177–184 (ISSN 1057–7157).
[12] D. Gracias, V. Kavthekar, J. Love, K. Paul, G. Whitesides, Adv. Mater. 14 (2002) 235–238 (ISSN 1521–4095).
[13] T.G. Leong, P.A. Lester, T.L. Koh, E.K. Call, D.H. Gracias, Langmuir 23 (2007) 8747–8751 (Preprint).
[14] J.H. Cho, D.H. Gracias, Nano Lett. 9 (2009) 4049–4052 (ISSN 15306984 (ISSN)).
[15] J.H. Cho, M.D. Keung, N. Verellen, L. Lagae, V.V. Moshchalkov, P. Van Dorpe, D.H. Gracias, Small 7 (2011) 1943–1948 (ISSN 1613–6829).
[16] R. Syms, IEEE Photon. Technol. Lett. 12 (2000) 1519–1521 (ISSN 1041–1135).
[17] R.J. Linderman, P.E. Kladitis, V.M. Bright, Sensors Actuators A Phys. 95 (2002) 135–142 (ISSN 0924–4247).
[18] Y. Hong, R. Syms, Sensors Actuators A Phys. 127 (2006) 381–391 (ISSN 0924–4247).
[19] J.R. Park, D. Choi, D. Gracias, T. Leong, N. Presser, G. Stupian, M. Leung, Y. Kim, Nanotechnology (2011) 22.
[20] J.W. van Honschoten, J.W. Berenschot, T. Ondarçuhu, R.G.P. Sanders, J. Sundaram, M. Elwenspoek, N.R. Tas, Appl. Phys. Lett. 97 (2010) 0141031–0141033.
[21] J. van Honschoten, A. Legrain, J. Berenschot, L. Abelmann, N. Tas, Micro-assembly of three dimensional tetrahedra by capillary forces Micro Electro Mechanical Systems (MEMS), IEEE 24th International Conference on 2011, pp. 288–291 (ISSN 1084–6999).
[22] J. Chen, Z. Fan, J. Zou, J. Engel, C. Liu, J. Aerosp. Eng. 16 (2003) 85–97.
[23] C.W. Chang, W. Hsu, J. Micromech. Microeng. (2009) 19.
[24] F. Iker, N. Andr/'e, T. Pardoen, J.P. Raskin, J. Microelectromech. Syst. 15 (2006) 1687–1697.
[25] N. Andr, B. Rue, G. Scheen, D. Flandre, L.A. Francis, J.P. Raskin, Sensors Actuators A Phys. 206 (2014) 67–74 (ISSN 0924–4247).
[26] C. Van Gulijk, Comput. Chem. Eng. 22 (1998) S767–S770 (ISSN 00981354).
[27] D.R. Lide, CRC Handbook of Chemistry and Physics, Internet Version, 86th ed. Taylor and Francis Group LLC, 2006.
[28] C. Alépée, Technologies for High-temperature Silicon Microreactors Series in Microsystems (Hartung-Gorre), 2001, ISBN 9783896496355.
[29] C.A. Klein, R.P. Miller, J. Appl. Phys. 87 (2000) 2265–2272 (ISSN 0021–8979).
[30] S. Fricke, A. Friedberger, G. Mueller, H. Seidel, U. Schmid, Strain Gauge Factor and TCR of Sputter Deposited Pt Thin Films up to 850 °C, Sensors, 2008 IEEE, 2008. 1532–1535 (ISSN 1930–0395).
[31] P. Schmid, C. Zarfl, G. Balogh, U. Schmid, Procedia Eng. 87 (2014) 172–175 (ISSN 1877–7058).
[32] G. Stoney, Proc. R. Soc. Lond. A 82 (1909) 172–175 (ISSN 0950–1207).
[33] A.B. Kaushik, H. Kahn, A. Heuer, J. Microelectromech. Syst. 14 (2005) 359–367 (ISSN 10577157).
[34] M. Avrekh, O. Monteiro, I. Brown, Appl. Surf. Sci. 158 (2000) 217–222.
[35] M. Salvadori, A. Vaz, R. Farias, M. Cattani, J. Metastable Nanocrystalline Mater. 20–21 (2004) 775–780.
[36] K. Fukuda, Y. Takeda, Y. Yoshimura, R. Shiwaku, L. Tran, T. Sekine, M. Mizukami, D. Kumaki, S. Tokito, Nat. Commun. (2014) 5.
[37] M. Kaltenbrunner, M.S. White, E.D. Głowacki, T. Sekitani, T. Someya, N.S. Sariciftci, S. Bauer, Nat. Commun. 3 (2012) 770.
[38] M.S. White, M. Kaltenbrunner, E.D. Głowacki, K. Gutnichenko, G. Kettlgruber, I. Graz, S. Aazou, C. Ulbricht, D.A. Egbe, M.C. Miron, et al., Nat. Photonics 7 (2013) 811–816.
[39] S. Kumar, M. Alam, Z. Connell, M. Haque, Scr. Mater. 65 (2011) 277–280.
[40] M. Rudneva, B. Gao, F. Prins, Q. Xu, H. Van Der Zant, H. Zandbergen, Microsc. Microanal. 19 (2013) 43–48.
[41] R. Tiggelaar, R. Sanders, A. Groenland, J. Gardeniers, Sensors Actuators A Phys. 152 (2009) 39–47.
[42] S. Firebaugh, K. Jensen, M. Schmidt, J. Microelectromech. Syst. 7 (1998) 128–135.
[43] G. Fox, S. Trolier-McKinstry, S. Krupanidhi, L. Casas, J. Mater. Res. 10 (1995) 1508–1515.
[44] S. Kweon, S. Choi, S. Yeom, J. Roh, Jpn. J. Appl. Phys. 40 (2001) 5850.
[45] U. Schmid, H. Seidel, Thin Solid Films 516 (2008) 898–906 (ISSN 0040–6090).
[46] P.M. Wang, M. Cornwell, M.R. Prausnitz, Diabetes Technol. Ther. 7 (2005) 131–141 (ISSN 15209156 (ISSN)).
[47] V. Branger, V. Pelosin, K. Badawi, P. Goudeau, Thin Solid Films 275 (1996) 22–24 (ISSN 0040–6090).