

Excitation quenching in Er³⁺-doped Al₂O₃ amplifiers

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Al₂O₃:Er³⁺ is an established material for integrated amplifier applications, offering several advantages such as low background losses, a broad emission spectrum, and a relatively high refractive index of 1.65, which enables the fabrication of small optical waveguide structures [1]. Our growth method allows for straightforward deposition on silicon wafers [2] and direct integration with silicon photonic circuits [3]. Net gain over a bandwidth of 80 nm (1500-1580 nm) and with a peak value of 2.0 dB/cm at 1533 nm was measured in our waveguide amplifiers [4]. By fitting the gain measurements with a simple Er³⁺ rate-equation model described in [4] one can extract values for the coefficient W_{ETU} of energy-transfer upconversion (ETU) between neighboring Er³⁺ ions, which has a detrimental influence on the amplifier performance. In the model, unlike other authors [5], we have discarded the excited-state absorption (ESA) process ${}^4\text{I}_{13/2} \rightarrow {}^4\text{I}_{9/2}$, because our direct ESA measurements show that this process lies outside the gain bandwidth. In further experiments we prove the existence of an additional ETU process (${}^4\text{I}_{13/2}, {}^4\text{I}_{11/2} \rightarrow {}^4\text{I}_{15/2}, {}^4\text{F}_{9/2}$), whose influence on gain, however, turned out to be small.

The resulting W_{ETU} values show a clear discrepancy of about one order of magnitude with the results obtained from luminescence decay measurements. The discrepancy is explained by the fact that a certain fraction of Er³⁺ ions undergoes fast quenching which is not resolved in our decay measurements owing to the time response of the set-up. This quenching can be induced by, e.g., active ion pairs and clusters, undesired impurities, or host material defects. Other investigations indicate that the characteristic time constant of such quenching may be on the order of 1-10 μs [6]. By modifying the model to account for these fast quenching effects, we derived that the fraction of quenched ions f_q increases from 10% to 33% when the Er³⁺ concentration increases from 1.17 to $3.66 \times 10^{20} \text{ cm}^{-3}$, as shown in Fig. 1 (a), suggesting a higher probability of fast quenching with increasing concentration. These results were afterwards confirmed by nonsaturable absorption experiments, see Fig. 1 (b), in which 1480-nm pump light was launched into the waveguide, and the transmission as a function of launched power was recorded. The transmission is strongly dependent on the amount of quenched ions, and our improved model could fit the results with good accuracy.

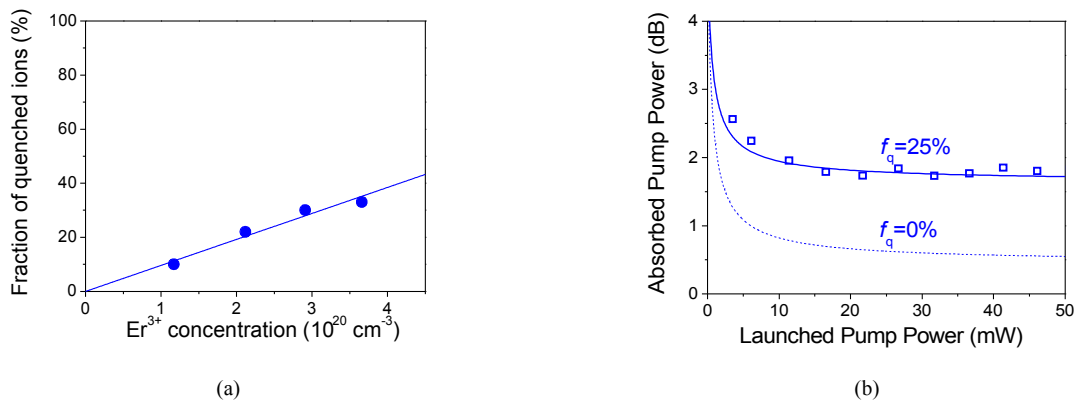


Fig. 1 (a) Fraction of quenched ions as a function of Er concentration; (b) absorbed versus launched pump power for a waveguide with an Er³⁺ concentration of $2.12 \times 10^{20} \text{ cm}^{-3}$ and 1 cm length. The squares are the experimental data, while the solid and dashed lines represent the calculations for 25% and 0% quenched ions, respectively.

The quantitative understanding of the influence of the different quenching processes leads to an improved simulation that predicts a gain of ~ 3 dB/cm for an optimized Er³⁺ concentration when pumping at 1480 nm.

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