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Nonequilibrium Statistical Physics | Physics Course Materials

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01. Introduction: Maps

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Abstract
Part one of course materials for Nonequilibrium Statistical Physics (Physics 626), taught by Gerhard Müller at the University of Rhode Island. Entries listed in the table of contents, but not shown in the document, exist only in handwritten form. Documents will be updated periodically as more entries become presentable.
Updated with version 2 on 5/3/2016.

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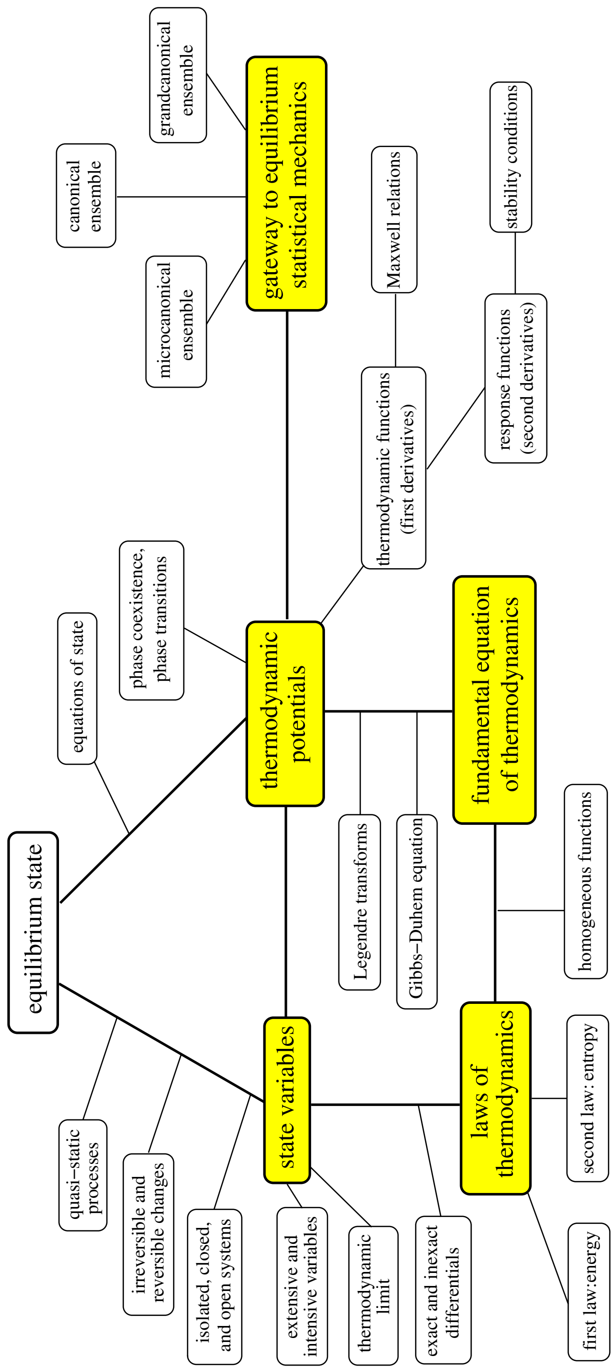
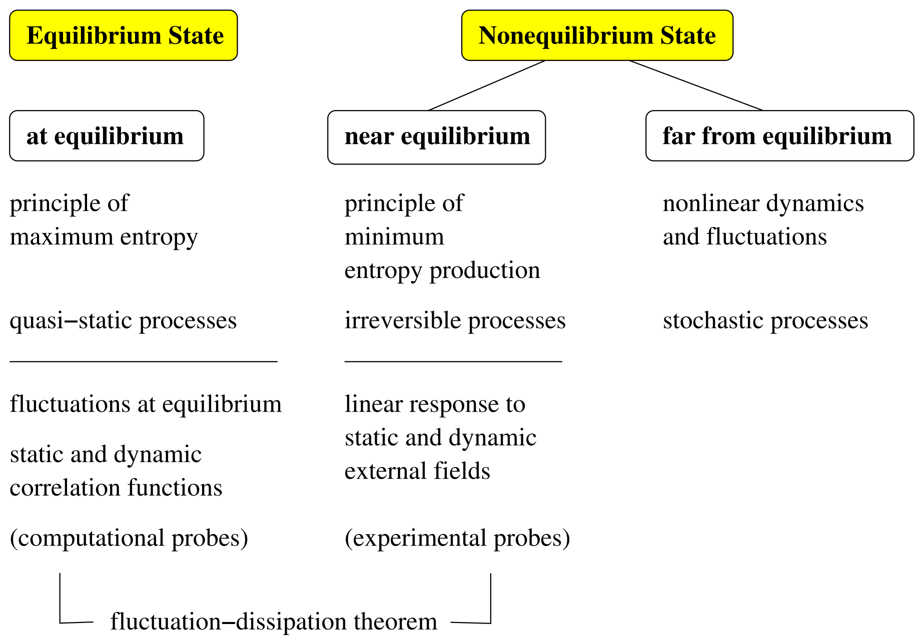


Thermal Equilibrium and Nonequilibrium [nln1]



Distinguish independently between

- *equilibrium* and *nonequilibrium* situations,
- *time-independent* and *time-dependent* phenomena.

	equilibrium situation	**nonequilibrium situation**
time–independent phenomena	equal–time correlations	equal–time correlations in steady states
time–dependent phenomena	delayed–time correlations	delayed–time correlations in steady states any correlations in non–steady states

Levels of Description in Statistical Physics [nln2]

microscopic level	**kinetic level**	**thermodynamic level**
N–particle phase space	1–particle phase space	configuration space
Liouville equation generalized Langevin equation	Boltzmann equation Fokker–Planck equation Langevin equation	hydrodynamic equations master equation Fokker–Planck equation Langevin equation
no contraction	some contraction	more contraction
deterministic time evolution	probabilistic time evolution	

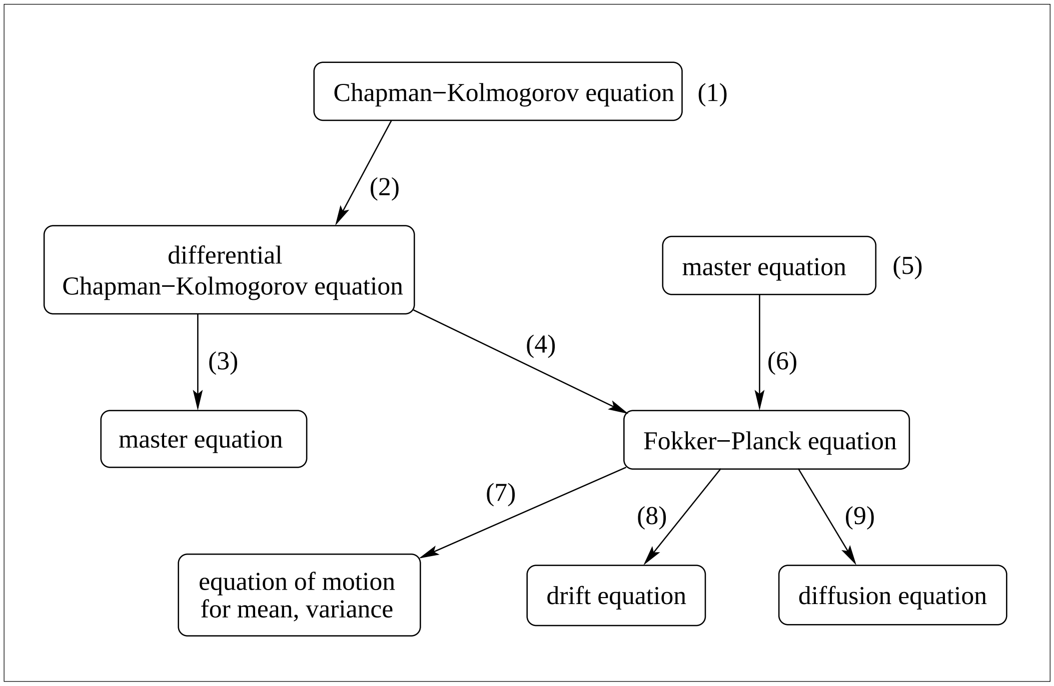
Contraction – memory – time scales [nln15]

microscopic dynamics	**⇒ contraction ⇒**	**stochastic dynamics**
future state determined by present state alone	focus on subset of dynamical variables	future state determined by present and past states
deterministic time evolution of dynamic variables	⇓	ignoring memory of past makes dynamics of selected variables probabilistic
	judicious choice: slow variables and long time scales	deterministic time evolution of probability distributions and mean values
	⇓ ⇒	short memory of fast variables has little impact on dynamics of slow variables at long times

Comments:

- In a classical Hamiltonian system the deterministic time evolution pertains to canonical coordinates and functions thereof.
- The time rate of change of any such variable depends on the instantaneous values of all canonical coordinates.
- On the contracted level of description we seek a way of describing an autonomous time evolution of a subset of variables.
- For that purpose the information contained in the instantaneous values of the variables that do not belong to the subset is transcribed into previous values of the variables that do belong to the subset.
- The autononmous time evolution of the variables belonging to the subset thus includes memory of its previous values.
- Slow variables contribute long memory and fast variables contribute short memory.
- If the subset contains all slow variables then any effects on its autonomous time evolution contributed by the remaining variables involve only short memory.
- Effects of short memory are more easily accounted for than effects of long memory.

Markov processes: map of specifications [nln16]



(1) Chapman-Kolmogorov equation imposes restrictions on permissible functions $P(x,t|x_0)$ but does not suggest a classification of processes.

(2) Particular solutions that are specified by

- $A(x,t)$ describing drift,
- $B(x,t)$ describing diffusion,
- $W(x|x';t)$ describing jumps.

(3) Jump processes exclusively.

(4) Processes with continuous sample paths, satisfying Lindeberg criterion (drift and diffusion, no jumps).

(5) Master equation with any $W(x|x';t)$ specifies a Markov process. Natural starting point for processes with discrete stochastic variables.

(6) Transition rates $W(x|x';t)$ of master equation approximated by two jump moments provided they exist. Approximation captures drift and diffusion parts (on some scale).

(7) Drift and diffusion determine mean $\langle\langle x\rangle\rangle$ and variance $\langle\langle x^2\rangle\rangle$ via equations of motion for jump moments.

(8) Deterministic process have no diffusive part: $B(x,t)=0$.

(9) Purely diffusive processes have no drift: $A(x,t)=0$.

Brownian motion: panoramic view [nln23]

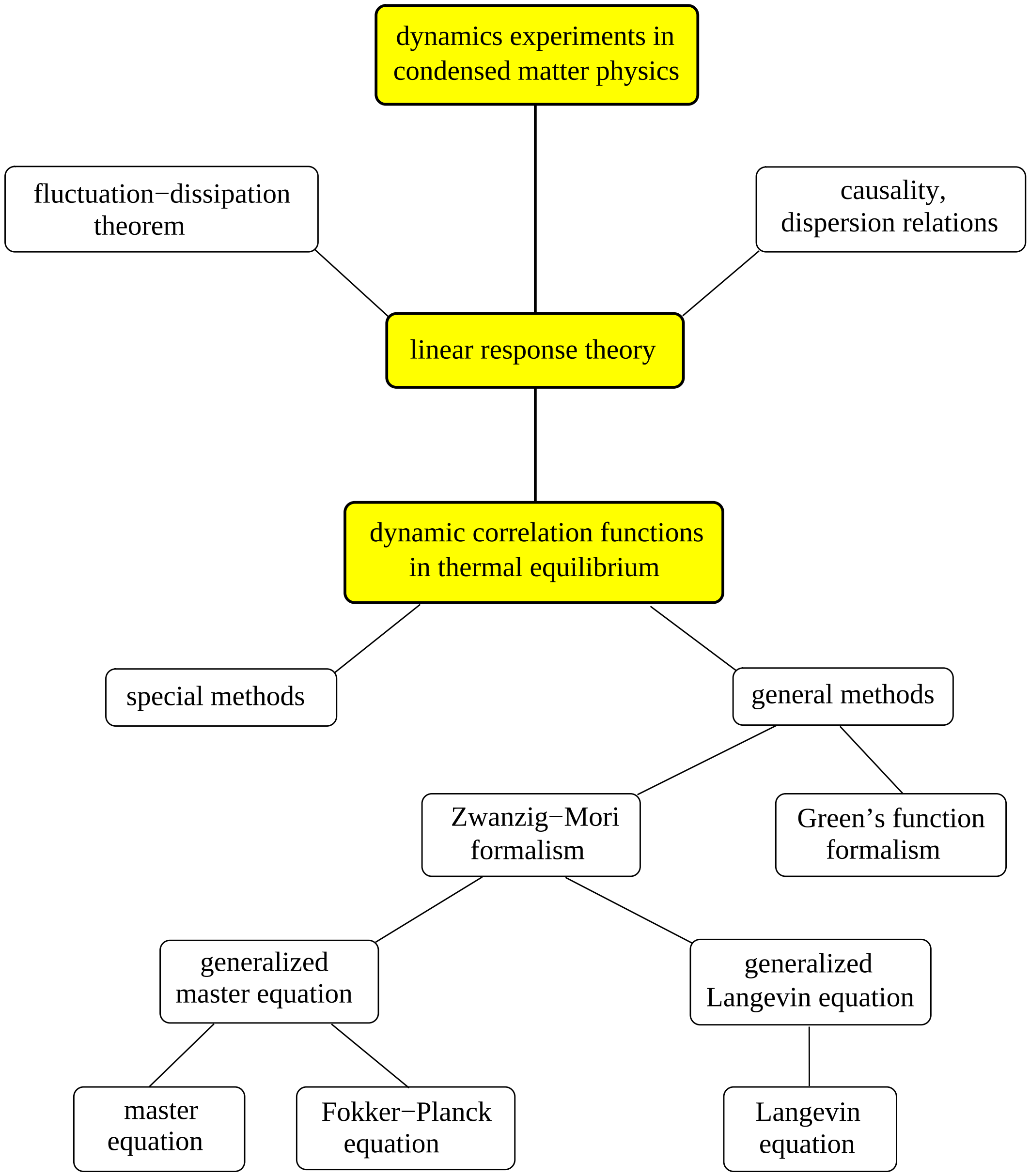
- Levels of contraction (horizontal)
- Modes of description (vertical)

$\longrightarrow$ contraction $\longrightarrow$

relevant space	N-particle phase space	1-particle phase space	configuration space
dynamical variables	$\{\mathbf{x}_i,\, \mathbf{p}_i\}$	$\mathbf{x}$, $\mathbf{p}$	$\mathbf{x}$
theoretical framework	Hamiltonian mechanics	Langevin theory	Einstein theory
... for dynamical variables	generalized Langevin equation	Langevin equation (for $dt \ll \tau_R$)	Langevin equation (for $dt \gg \tau_R$)
... for probability distribution	quant./class. Liouville equation	Fokker-Planck equation (Ornstein-Uhlenbeck process)	Fokker-Planck equation (diffusion process)

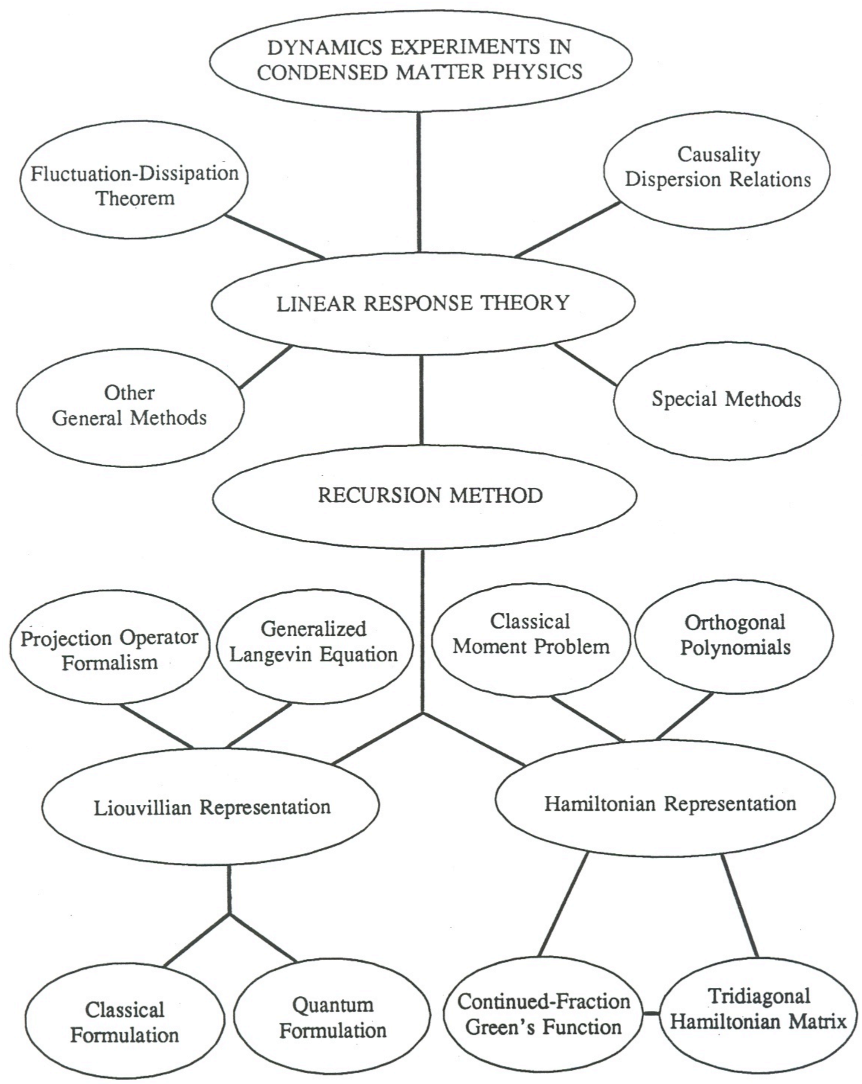
- Here dt is the time step used in the theory and τ_R is the relaxation time associated with the drag force the Brownian particle experiences.
- The generalized Langevin equation is equivalent to the Hamiltonian equation of motion for a generic classical many-body system and equivalent to the Heisenberg equation of motion for a generic quantum many-body system.

Linear response and equilibrium dynamics [nln24]



Stage for Recursion Method [nln79]

Recursion method as applied to many-body dynamics:
backdrop, props, protagonists.



[from Viswanath and Müller 1994]

Modules of Recursion Method [nln80]

Recursion method as applied to many-body dynamics:
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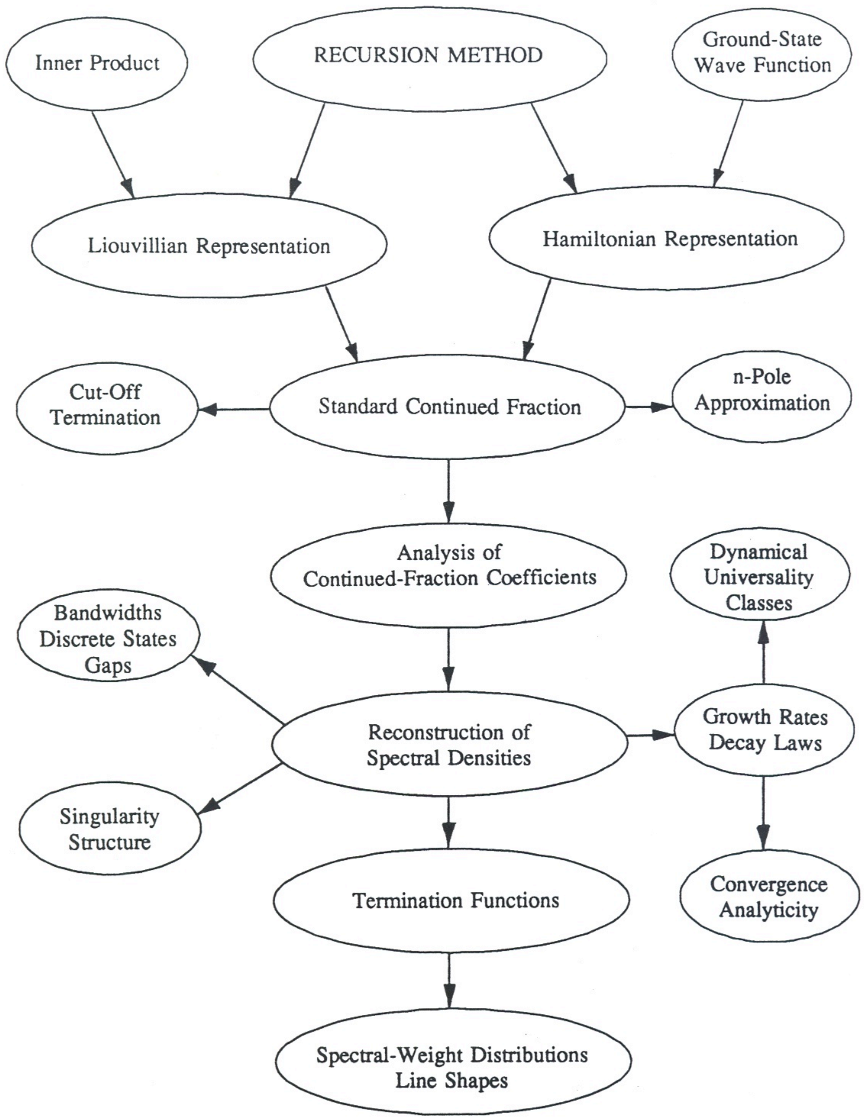

[from Viswanath and Müller 1994]