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# A Distributed Parabolic Control with Mixed Boundary Conditions

Jose-Luis Menaldi \* Domingo Alberto Tarzia<sup>†</sup>

#### Abstract

We study the asymptotic behavior of an optimal distributed control problem where the state is given by the heat equation with mixed boundary conditions. The parameter  $\alpha$  intervenes in the Robin boundary condition and it represents the heat transfer coefficient on a portion  $\Gamma_1$  of the boundary of a given regular *n*-dimensional domain. For each  $\alpha$ , the distributed parabolic control problem optimizes the internal energy *g*. It is proven that the optimal control  $\hat{g}_{\alpha}$  with optimal state  $u_{\hat{g}_{\alpha}\alpha}$  and optimal adjoint state  $p_{\hat{g}_{\alpha}\alpha}$  are convergent as  $\alpha \to \infty$  (in norm of a suitable Sobolev parabolic space) to  $\hat{g}$ ,  $u_{\hat{g}}$  and  $p_{\hat{g}}$ , respectively, where the limit problem has Dirichlet (instead of Robin) boundary conditions on  $\Gamma_1$ . The main techniques used are derived from the parabolic variational inequality theory.

**Keywords and phrases:** Parabolic variational inequalities, Distributed evolution optimal control, Mixed boundary conditions, Adjoint state, Optimality condition, Asymptotic.

**AMS (MOS) Subject Classification.** Primary: 49J20, 49J40, Secondary: 35R35, 35K20, 35B40.

# 1 Introduction

Let  $\Omega$  be a bounded domain in  $\mathbb{R}^n$  with a regular boundary  $\partial \Omega = \Gamma_1 \cup \Gamma_2$ , which is the union of two essentially disjoint (and regular) portions  $\Gamma_1$  and  $\Gamma_2$ , where  $\Gamma_1$  has a positive (n-1)-Hausdorff measure. Also suppose given a time interval [0, T], for some T > 0. Consider the following two-state evolution heat conduction problems with mixed boundary conditions,

$$\partial_t u - \Delta u = g \text{ in } \Omega, \qquad u \big|_{\Gamma_1} = b, \qquad -\partial_n u \big|_{\Gamma_2} = q,$$
(1.1)

and, for a parameter  $\alpha > 0$ ,

$$\partial_t u - \Delta u = g \text{ in } \Omega, \qquad -\partial_n u \big|_{\Gamma_1} = \alpha(u - b), \qquad -\partial_n u \big|_{\Gamma_2} = q, \qquad (1.2)$$

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both with an initial condition

$$u(0) = v_b, \tag{1.3}$$

where g is the internal energy in  $\Omega$ , b is the temperature (of the external neighborhood) on  $\Gamma_1$  for (1.1) (for (1.2)), q is the heat flux on  $\Gamma_2$  and  $\alpha$  is the heat transfer coefficient of  $\Gamma_1$  (Newton's law on  $\Gamma_1$ ). All data, g, q, b,  $v_b$  and the domain  $\Omega$  with the boundary  $\partial \Omega = \Gamma_1 \cup \Gamma_2$  are assumed to be sufficiently smooth so that the problems (1.1) and (1.2) admit variational solutions in Sobolev spaces.

The data b,  $v_b$  and q are fixed, sufficiently smooth and satisfy the compatibility condition  $v_b = b$  on  $\Gamma_1$ , while g is taken as a control variable in  $L^2(0,T;L^2(\Omega))$ , and  $\alpha$  as a (singular) parameter destined to approaches infinite. Thus, denote by  $u_g$  and  $u_{g\alpha}$  the solution of (1.1) and (1.2), respectively, with the initial condition (1.3) in the following standard variational form

$$\begin{cases} u_g - v_b \in L^2(0, T; V_0), & u_g(0) = v_b \text{ and } \dot{u}_g \in L^2(0, T; V_0') \\ \text{such that } \langle \dot{u}_g(t), v \rangle + a(u_g(t), v) = L_g(t, v), & \forall v \in V_0, \end{cases}$$
(1.4)

and

$$\begin{cases} u_{g\alpha} \in L^2(0,T;V), & u_{g\alpha}(0) = v_b \text{ and } \dot{u}_{g\alpha} \in L^2(0,T;V') \\ \text{such that } \langle \dot{u}_{g\alpha}(t), v \rangle + a_\alpha(u_{g\alpha}(t),v) = L_{g\alpha}(t,v), \quad \forall v \in V, \end{cases}$$
(1.5)

where

$$V_{0} := \{ v \in H^{1}(\Omega) : v |_{\Gamma_{1}} = 0 \},$$

$$H := L^{2}(\Omega), \qquad (g, h)_{H} := \int_{\Omega} gh \, \mathrm{d}x,$$

$$L_{g}(t, v) := (g(t), v)_{H} - \int_{\Gamma_{2}} q(t)v \, \mathrm{d}\gamma,$$

$$a(u, v) := \int_{\Omega} \nabla u \cdot \nabla v \, \mathrm{d}x,$$

$$a_{\alpha}(u, v) := a(u, v) + \alpha \int_{\Gamma_{1}} uv \, \mathrm{d}\gamma,$$

$$L_{g\alpha}(t, v) := L_{g}(t, v) + \alpha \int_{\Gamma_{1}} bv \, \mathrm{d}\gamma,$$
(1.6)

and  $\langle \cdot, \cdot \rangle$  denotes the duality bracket. Note that the dual space  $V'_0$  (and V') of  $V_0$ (and V) is not an space of distributions, since  $\mathcal{D}(\Omega)$  is not dense in  $V_0 \subset V$ , due to the non-zero boundary conditions on  $\Gamma_2$ . The norm in  $V_0$  is given by  $v \mapsto \|\nabla v\|_H$ , while the norm in V is  $(\|v\|_H^2 + \|\nabla v\|_H^2)^{1/2}$ . Nevertheless,  $v \mapsto L_g(t, v)$  and  $v \mapsto L_{g\alpha}(t, v)$  are linear continuous functional satisfying

$$\begin{split} \|L_g(t,\cdot)\|_{V'_0} &\leq \|g(t)\|_{V'_0} + \|q(t)\|_{H^{-1/2}(\Gamma_2)}, \quad \forall v \in V_0, \\ \|L_{g\alpha}(t,\cdot)\|_V &\leq \|g(t)\|_{V'} + \|q(t)\|_{H^{-1/2}(\Gamma_2)} + \alpha \|b\|_{H^{1/2}(\Gamma_1)}, \quad \forall v \in V, \end{split}$$

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and  $a(\cdot, \cdot)$  and  $a_{\alpha}(\cdot, \cdot)$  are bilinear symmetric continuous forms on  $V_0$  and V, respectively. Also, it is clear the compatibility assumption  $v_b = b$  on  $\Gamma_1$  and that if b = 0 then  $L_g(t, \cdot) = L_{g,\alpha}(t, \cdot)$ .

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One should remark that an element u of  $L^2(0,T;V)$  such that  $\dot{u}$  belongs to  $L^2(0,T;V')$  then u can be regarded as a continuous function from [0,T] into H. This makes clear the meaning of the initial condition at t = 0 (and idem with  $V_0$  replacing V).

On the space  $\mathcal{H} := L^2(\Omega \times ]0, T[)$  with norm  $\|\cdot\|_{\mathcal{H}}$  and inner product  $(\cdot, \cdot)_{\mathcal{H}}$ , i.e.,

$$(u,v)_{\mathcal{H}} = \int_0^T \left( u(t), v(t) \right)_H \mathrm{d}t, \quad \forall u, v \in \mathcal{H},$$

consider the nonnegative functional costs J and  $J_{\alpha}$ , defined by the expressions

$$J(g) := \frac{1}{2} \|u_g - z_d\|_{\mathcal{H}}^2 + \frac{m}{2} \|g\|_{\mathcal{H}}^2, \tag{1.7}$$

and

$$J_{\alpha}(g) := \frac{1}{2} \|u_{g\alpha} - z_d\|_{\mathcal{H}}^2 + \frac{m}{2} \|g\|_{\mathcal{H}}^2,$$
(1.8)

where  $z_d$  is a given element in  $\mathcal{H} = L^2(\Omega \times ]0, T[)$  and m is a strictly positive constant.

Our interest is on the distributed parabolic (or evolution) optimal control problems

Find 
$$\hat{g}$$
 such that  $J(\hat{g}) \le J(g), \quad \forall g \in \mathcal{H}$  (1.9)

and

 $\mathcal{B}$ 

Find 
$$\hat{g}_{\alpha}$$
 such that  $J_{\alpha}(\hat{g}_{\alpha}) \leq J_{\alpha}(g), \quad \forall g \in \mathcal{H},$  (1.10)

as well as the asymptotic behavior as the parameter  $\alpha$  approaches infinite.

This type of optimal distributed control problems have been extensively studied, e.g., see the book Lions [10] among others. As point out early, our interest is the convergence as  $\alpha \to \infty$ , a parabolic version of Gariboldi and Tarzia [8], which is related to Ben Belgacem et al. [4] and Tabacman and Tarzia [11].

# 2 Parabolic Equations with Mixed Conditions

Note that if via Riesz' representation H = H' then one has  $V \subset H \subset V'$  and  $V_0 \subset H \subset V'_0$  with a continuous and dense inclusion.

As mentioned early the control parameter g belongs to  $\mathcal{H}$ , and the data for the optimal control problems are  $z_d$  and m satisfying

$$z_d \in \mathcal{H} = L^2(0, T; L^2(\Omega)), \quad \text{and} \quad m > 0.$$

$$(2.1)$$

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The regularity of the domain  $\Omega$ , the boundary  $\Gamma_1 \cup \Gamma_2$  and the regularity of the boundary data  $v_b$ , b and q are summarized on the assumption

there exists 
$$\psi \in L^2(0,T;H^2(\Omega))$$
 with  $\psi \in L^2(0,T;L^2(\Omega))$   
such that  $\psi(0) = v_b$ ,  $\psi|_{\Gamma_1} = b$ ,  $\partial_n \psi|_{\Gamma_1} = 0$ ,  $-\partial_n \psi|_{\Gamma_2} = q$ , (2.2)

with the standard notation of Sobolev and Lebesgue spaces and the compatibility assumption  $v_b = b$  on  $\Gamma_1$ . Note the over conditioning for  $\psi$  on  $\Gamma_1$ , which is not necessary but convenient in some way (e.g., the adjoint state has a very similar equation with homogeneous boundary conditions).

Thus, the change of unknown function u into  $u - \psi$  reduces to analysis the case where the boundary data  $v_b$ , b and q are all zero, and g is replaced by  $g - (\partial_t - \Delta)\psi$ . However, for  $\alpha > 0$  a new term appears, namely,

$$\langle g_{\psi}(t), v \rangle = (g(t), v)_H + \int_{\Gamma_1} v \partial_n \psi(t) \, \mathrm{d}\gamma, \quad \forall v \in V,$$
(2.3)

i.e., the new Robin boundary condition is non-homogeneous and

$$\|g_{\psi}(t)\|_{V'} = \sup_{\|v\|_{V} \le 1} \left| \langle g_{\psi}(t), v \rangle \right| \le \|g(t)\|_{L^{2}(\Omega)} + \|\partial_{n}\psi(t)\|_{H^{-1/2}(\Gamma_{1})}.$$

Thus, because of the over conditioning on  $\Gamma_1$  one has  $g_{\psi} = g$ . Anyway, both problems, (1.4) and (1.5) become

$$\begin{cases} u_g \in L^2(0,T;V_0), & \text{with } u_g(0) = 0 \quad \text{and } \dot{u}_g \in L^2(0,T;V_0') \\ \text{such that } \langle \dot{u}_g(t), v \rangle + a(u_g(t),v) = (g(t),v)_H, \quad \forall v \in V_0 \end{cases}$$
(2.4)

and

$$\begin{cases} u_{g\alpha} \in L^2(0,T;V), & \text{with } u_{g\alpha}(0) = 0 \text{ and } \dot{u}_{g\alpha} \in L^2(0,T;V') \\ \text{such that } \langle \dot{u}_{g\alpha}(t), v \rangle + a_{\alpha} \big( u_{g\alpha}(t), v \big) = (g(t), v)_H, & \forall v \in V, \end{cases}$$
(2.5)

where  $(\cdot, \cdot)_H$ ,  $a(\cdot, \cdot)$  and  $a_{\alpha}(\cdot, \cdot)$  are as in (1.6). Again  $V_0 \subset V$  with inclusion continuous but not dense, so that V' is not identifiable with a subset of  $V'_0$ . However, by Hahn-Banach Theorem, any element in  $V'_0$  can be extended to an element in V' preserving its norm.

Recall that for any element u in  $L^2(0,T;V)$  with  $\dot{u}$  in  $L^2(0,T;V')$  such that the distribution  $(\partial_t - \Delta)u$  belongs to  $L^2(\Omega \times ]0,T[)$  one can integrate by parts to interpret  $\partial_n u$  as an element in  $L^2(0,T;H^{-1/2}(\partial\Omega))$ , where  $H^{-1/2}(\partial\Omega)$ is the dual space of  $H^{1/2}(\partial\Omega) = \gamma(H^1(\Omega))$  and  $\gamma$  is the trace operator from  $H^1(\Omega)$  onto  $H^{1/2}(\partial\Omega)$ . Again, to simplify the arguments, one may assume that  $\partial\Omega = \Gamma_1 \cup \Gamma_2$  such that for any  $v_i$  in  $H^{1/2}(\Gamma_i)$  there exists v in  $H^1(\Omega)$  satisfying  $v = v_i$  on  $\Gamma_i$ , for i = 1, 2, e.g., the two pieces of the boundary are strictly disjoint,  $\Gamma_1 \cap \Gamma_2 = \emptyset$  (i.e.,  $\Gamma_i = \partial\Omega_i$  and  $\overline{\Omega}_1 \subset \Omega_2$ ). Therefore, the parabolic equations (2.4) and (2.5) mean the following:

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- space of the solution:  $u_g$  in  $L^2(0,T;V_0)$  with  $\dot{u}_g$  in  $L^2(0,T;V'_0)$ , and  $u_{g\alpha}$  in  $L^2(0,T;V)$  with  $\dot{u}_{g\alpha}$  in  $L^2(0,T;V')$ ,
- *initial condition*: for either  $u = u_g$  or  $u = u_{g\alpha}$  the solution u belongs to  $C^0(0,T; L^2(\Omega))$  and so u(0) = 0 in  $L^2(\Omega)$ ,
- equation in  $\Omega \times ]0, T[:$  for either  $u = u_g$  or  $u = u_{g\alpha}$  the solution u is considered as a distribution so that  $(\partial_t \Delta)u = g$  in  $\mathcal{D}'(\Omega \times ]0, T[)$ ,
- boundary condition on  $\Gamma_2$ : for either  $u = u_g$  or  $u = u_{g\alpha}$  the trace of the solution u is defined and  $\partial_n u = 0$  in  $L^2(0,T; H^{-1/2}(\Gamma_2))$ ,
- boundary condition on  $\Gamma_1$ :  $u_g = 0$  in  $L^2(0,T; H^{1/2}(\Gamma_1))$  and  $\partial_n u_{g\alpha} + \alpha u_{g\alpha} = 0$  in  $L^2(0,T; H^{-1/2}(\Gamma_1))$ .

Firstly, note that  $u_{g\alpha}|_{\Gamma_1}$  belongs to  $L^2(0,T; H^{1/2}(\Gamma_1))$  and

$$L^{2}(0,T; H^{1/2}(\Gamma_{1})) \subset L^{2}(0,T; L^{2}(\Gamma_{1})) \subset L^{2}(0,T; H^{-1/2}(\Gamma_{1})),$$

with continuous and dense inclusion. Secondly, when comparing the solutions  $u_g$  and  $u_{g\alpha}$  one has both in the larger space  $L^2(0,T;V)$ . However, the continuous inclusion  $V_0 \subset V$  is not dense, and so the inclusion  $V' \subset V'_0$  is not injective, one has  $\dot{u}_g$  and  $\dot{u}_{g\alpha}$  elements in  $L^2(0,T;V'_0)$ , which are not identifiable as distributions.

## **3** State and Adjoint State Equations

To study the optimal control problem (1.9), denote by  $u_0$  the solution  $u_g$  of the parabolic variational equality either (1.4) or equivalently (2.4) corresponding to g = 0, and define the (linear) operator  $C: \mathcal{H} \to L^2(0,T;V_0)$ , given by  $C(g) := u_g - u_0$ . We have

**Proposition 3.1.** With the previous notation and assumptions, the functional (1.7) can be expressed as

$$J(g) = \frac{1}{2}\pi(g,g) - \ell(g) + \frac{1}{2} ||z_d - u_0||_{\mathcal{H}}^2, \quad \forall g \in \mathcal{H},$$

where  $\pi(g,h) := (C(g), C(h))_{\mathcal{H}} + m(g,h)_{\mathcal{H}}$  is a symmetric, continuous and coercive bilinear form on  $\mathcal{H}$  and  $\ell(g) := (C(g), z_d - u_0)_{\mathcal{H}}$  is a linear continuous functional on  $\mathcal{H}$ . Moreover, J is strictly convex and its Gateaux derivative is given by  $\langle J'(g), h \rangle = (u_g - z_d, C(g))_{\mathcal{H}} + m(g,h)_{\mathcal{H}}$ . Furthermore, as a consequence, the optimal control problem (1.9) has a unique minimizer  $\hat{g}$  in  $\mathcal{H}$ , i.e.,  $J(\hat{g}) \leq J(g)$ , for every g in  $\mathcal{H}$ , any solution  $\bar{g}$  of the equation  $J'(\bar{g}) = 0$  is indeed a minimizer. Also, if  $p_g$  is the adjoint state defined by the parabolic variational equality with a terminal condition

$$\begin{cases} p_g \in L^2(0,T;V_0), & with \quad p_g(T) = 0 \quad and \quad \dot{p}_g \in L^2(0,T;V_0') \\ such \ that \quad -\langle \dot{p}_g(t), v \rangle + a(u_g(t),v) = (u_g - z_d, v)_H, \quad \forall v \in V_0, \end{cases}$$
(3.1)

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then  $J'(g) = mg + p_g$  for every g in  $\mathcal{H}$  and  $J'(\hat{g}) = m\hat{g} + p_{\hat{g}} = 0$ .

*Proof.* Note the boundary conditions for the adjoint state  $p_g$  are

$$p_g(t) = 0$$
 on  $\Gamma_1$  and  $\partial_n p_g(t) = 0$  on  $\Gamma_2$ 

for almost every t in ]0, T[.

First, we check the expression of J, if  $z'_d := z_d - u_0$  then

$$J(g) = \frac{1}{2} \|C(g) - z'_d\|_{\mathcal{H}}^2 + \frac{m}{2} \|g\|_{\mathcal{H}}^2 =$$
  
=  $\frac{1}{2} \Big[ \|C(g)\|_{\mathcal{H}}^2 + \|z'_d\|_{\mathcal{H}}^2 - 2(C(g), z'_d)_{\mathcal{H}} \Big] + \frac{m}{2} \|g\|_{\mathcal{H}}^2 =$   
=  $\frac{1}{2} \pi(g, g) - L(g) + \frac{1}{2} \|z_d - u_0\|_{\mathcal{H}}^2.$ 

To verify that  $g \mapsto C(g)$  is a linear application, one checks that the function  $r_1u_{g_1} + r_2u_{g_2} + (1 - r_1 - r_2)u_0$  is a solution of the parabolic variational equality (1.4) with  $g = r_1g_1 + r_2g_2$ , for every real numbers  $r_1, r_2$ ; and by uniqueness one has

$$u_{r_1g_1+r_2g_2} = r_1 u_{g_1} + r_2 u_{g_2} + (1 - r_1 - r_2) u_0, aga{3.2}$$

for every  $r_i, r_2$  in  $\mathbb{R}$  and  $g_1, g_2$  in  $\mathcal{H}$ . Hence,

$$C(r_1g_1 + r_2g_2) = u_{r_1g_1 + r_2g_2} - u_0 = r_1u_{g_1} + r_2u_{g_2} + (1 - r_1 - r_2)u_0 - u_0 =$$
  
=  $r_1(u_{g_1} - u_0) + r_2(u_{g_2} - u_0) = r_1C(g_1) + r_2C(g_2),$ 

i.e., the operator C is linear.

Now to check the continuity of C, we note that since  $\Gamma_1$  has positive measure, Poincaré inequality implies that the bilinear form  $a(\cdot, \cdot)$  is coercive on  $V_0$ , i.e., there exists  $\lambda_0 > 0$  such that

$$a(v,v) \ge \lambda_0 \|\nabla v\|_H^2, \quad \forall v \in V_0.$$

$$(3.3)$$

We have

$$(\dot{u}_g(t) - \dot{u}_0(t), v)_H + a(u_g(t) - u_0(t), v) = (g(t), v)_H, \quad \forall v \in V_0,$$

and, in particular, for  $v = u_g(t) - u_0(t)$ ,

$$\begin{aligned} &\frac{1}{2} \frac{\mathrm{d}}{\mathrm{d}t} \Big( \|u_g(t) - u_0(t)\|_H^2 \Big) + \lambda_0 \Big\| \nabla \big( u_g(t) - u_0(t) \big) \Big\|_H^2 \leq \\ & \leq \big( g(t), u_g(t) - u_0(t) \big)_H \leq \frac{1}{2\lambda_0} \|g(t)\|_{V_0'}^2 + \frac{\lambda_0}{2} \| \nabla \big( u_g(t) - u_0(t) \big) \|_{V_0}^2, \end{aligned}$$

where the dual norm is given by

 $\|v\|_{V_0'}^2 = \sup \{(v, \varphi)_H : \varphi \in V_0, \ \|\varphi\|_{V_0} \le 1\}.$ 

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This yields

$$\begin{aligned} \|\nabla C(g)\|_{\mathcal{H}} &\leq \frac{1}{\lambda_0} \Big[ \int_0^T \|g(t)\|_{V_0'}^2 \,\mathrm{d}t \Big]^{1/2}, \\ \sup_{0 \leq t \leq T} \|C(g)(t)\|_H &\leq \frac{1}{\sqrt{\lambda_0}} \Big[ \int_0^T \|g(t)\|_{V_0'}^2 \,\mathrm{d}t \Big]^{1/2}, \end{aligned}$$

and going back to the equation, we get

$$\left[\int_{0}^{T} \left\|\frac{\mathrm{d}}{\mathrm{d}t} \left(C(g)(t)\right)\right\|_{V_{0}'}^{2} \mathrm{d}t\right]^{1/2} \leq \frac{2}{\lambda_{0}} \left[\int_{0}^{T} \|g(t)\|_{V_{0}'}^{2} \mathrm{d}t\right]^{1/2}$$

Hence the operator

$$C \colon L^2 \big( 0, T; V_0' \big) \to \big\{ v \in L^2(0, T; V_0) \cap L^\infty(0, T; H) : \dot{v} \in L^2(0, T; V_0') \big\}$$

is actually continuous. As a consequence, the bilinear form  $\pi(\cdot, \cdot)$  is symmetric,

continuous and coercive on  $\mathcal{H} \times \mathcal{H}$ , since  $\mathcal{H} \subset L^2(0,T;V'_0)$ . To complete the argument, we choose v = C(h) in (3.1) and  $v = p_g$  in (1.4) with g = 0 and g = h to obtain, after integrating in t, the equalities

$$-\left(\dot{p}_g, C(h)\right)_{\mathcal{H}} + \int_0^T a\left(p_g(t), C(h)(t)\right) \mathrm{d}t = \left(u_g - z_d, C(h)\right)_{\mathcal{H}}$$

and

$$(\dot{u}_h - \dot{u}_0, p_g)_{\mathcal{H}} + \int_0^T a (u_h(t) - u_0(t), p_g(t)) dt = (h, p_g)_{\mathcal{H}}.$$

Thus

$$\int_0^T \frac{\mathrm{d}}{\mathrm{d}t} \big( p_g(t), C(h)(t) \big)_H \,\mathrm{d}t + (h, p_g)_{\mathcal{H}} = \big( u_g - z_d, C(h) \big)_{\mathcal{H}},$$

and because  $p_g(T) = 0$  and C(h)(0) = 0, we deduce  $J'(g) = mg + p_g$ .

To show that  $g \mapsto J(g)$  is strictly convex, one makes use of (1.7) and (3.2) to check that

$$(1-\theta)J(g_2) + \theta J(g_1) - J((1-\theta)g_1 + \theta g_2) = = \frac{1}{2}\theta(1-\theta) [||u_{g_1} - u_{g_2}||_{\mathcal{H}}^2 + m||g_1 - g_2||_{\mathcal{H}}^2],$$

for every  $\theta$  in [0, 1] and any  $g_1, g_2$  in  $\mathcal{H}$ .

Similarly, to study the optimal control problem (1.10), denote by 
$$u_{0\alpha}$$
 the so-  
lution  $u_{g\alpha}$  of the parabolic variational equality either (1.5) or equivalently (2.5)  
corresponding to  $g = 0$ , and define the (linear) operator  $C_{\alpha}: \mathcal{H} \to L^2(0,T;V)$ ,  
given by  $C_{\alpha}(g) := u_{g\alpha} - u_{0\alpha}$ . We have

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**Proposition 3.2.** With the previous notation and assumptions, the functional (1.8) can be expressed as

$$J_{\alpha}(g) = \frac{1}{2}\pi_{\alpha}(g,g) - \ell_{\alpha}(g) + \frac{1}{2}\|z_d - u_{0\alpha}\|_{\mathcal{H}}^2, \quad \forall g \in \mathcal{H},$$

where  $\pi_{\alpha}(g,h) := (C_{\alpha}(g), C_{\alpha}(h))_{\mathcal{H}} + m(g,h)_{\mathcal{H}}$  is a symmetric, continuous and coercive bilinear form on  $\mathcal{H}$  and  $\ell_{\alpha}(g) := (C_{\alpha}(g), z_d - u_{0\alpha})_{\mathcal{H}}$  is a linear continuous functional on  $\mathcal{H}$ . Moreover,  $J_{\alpha}$  is strictly convex and its Gateaux derivative of  $J_{\alpha}$  is given by  $\langle J'_{\alpha}(g), h \rangle = (u_g - z_d, C_{\alpha}(g))_{\mathcal{H}} + m(g,h)_{\mathcal{H}}$ . Furthermore, as a consequence, the optimal control problem (1.10) has a unique minimizer  $\hat{g}_{\alpha}$  in  $\mathcal{H}$ , i.e.,  $J_{\alpha}(\hat{g}_{\alpha}) \leq J_{\alpha}(g)$ , for every g in  $\mathcal{H}$ , and any solution  $\bar{g}_{\alpha}$  of the equation  $J'(\bar{g}_{\alpha}) = 0$  is indeed a minimizer. Also if  $p_{g\alpha}$  is the adjoint state defined by the parabolic variational equality with a terminal condition

$$\begin{cases} p_{g\alpha} \in L^2(0,T;V), & with \quad p_{g\alpha}(T) = 0 \quad and \quad \dot{p}_{g\alpha} \in L^2(0,T;V') \\ such that \quad -\langle \dot{p}_{g\alpha}(t), v \rangle + a_\alpha(p_{g\alpha}(t), v) = (u_{g\alpha} - z_d, v)_H, \ \forall v \in V, \end{cases}$$
(3.4)

then  $J'_{\alpha}(g) = mg_{\alpha} + p_{g_{\alpha}}$  for every g in  $\mathcal{H}$  and  $J'_{\alpha}(\hat{g}_{\alpha}) = m\hat{g}_{\alpha} + p_{\hat{g}_{\alpha}} = 0$ .

*Proof.* The calculations are similar to the previous proposition. We remark that the boundary conditions for the adjoint state  $p_{q\alpha}$  are

$$-\partial_n p_{g\alpha}(t) = \alpha p_{g\alpha}$$
 on  $\Gamma_1$  and  $\partial_n p_{g\alpha}(t) = 0$  on  $\Gamma_2$ .

for almost every t in ]0, T[. Moreover, we assume  $\alpha > 0$  so that the coerciveness (3.3) becomes

$$a_{\alpha}(v,v) \ge \lambda_1 \min\{1,\alpha\} [\|\nabla v\|_H^2 + \|v\|_H^2], \quad \forall v \in V,$$
(3.5)

Indeed, by contradiction one can show that  $a_1(v, v) \ge c_1 ||v||_H^2$  for every v in V, which implies (3.5). The continuity of  $a(\cdot, \cdot)$  in V uses the continuity of the trace in  $H^1(\Omega)$ , namely, for some  $\Lambda_1 > 0$  one has

$$a_{\alpha}(u,v) \leq \Lambda_1 \max\{1,\alpha\} \|u\|_V \|v\|_V, \quad \forall v \in V,$$

$$(3.6)$$

which depends on  $\alpha > 0$ .

The operator  $C_{\alpha}$  actually maps the space  $L^2(0,T;V')$  into the space

$$\left\{ v \in L^2(0,T;V) \cap L^\infty(0,T;H) : \dot{v} \in L^2(0,T;V') \right\}$$

and the estimates

$$\begin{aligned} \|\nabla C_{\alpha}(g)\|_{\mathcal{H}} &\leq \frac{1}{\lambda_{1}} \Big[ \int_{0}^{T} \|g(t)\|_{V'}^{2} \,\mathrm{d}t \Big]^{1/2}, \\ \sup_{0 \leq t \leq T} \|C_{\alpha}(g)(t)\|_{H} &\leq \frac{1}{\sqrt{\lambda_{1}}} \Big[ \int_{0}^{T} \|g(t)\|_{V'}^{2} \,\mathrm{d}t \Big]^{1/2}, \\ \Big[ \int_{0}^{T} \Big\| \frac{\mathrm{d}}{\mathrm{d}t} \big( C_{\alpha}(g)(t) \big) \Big\|_{V_{0}'}^{2} \,\mathrm{d}t \Big]^{1/2} &\leq \frac{2}{\lambda_{1}} \Big[ \int_{0}^{T} \|g(t)\|_{V'}^{2} \,\mathrm{d}t \Big]^{1/2} \end{aligned}$$

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are independent of  $\alpha > 1$ , but

$$\left[\int_{0}^{T} \left\|\frac{\mathrm{d}}{\mathrm{d}t} \left(C_{\alpha}(g)(t)\right)\right\|_{V'}^{2} \mathrm{d}t\right]^{1/2} \leq \frac{1+\alpha}{\lambda_{1}} \left[\int_{0}^{T} \|g(t)\|_{V'}^{2} \mathrm{d}t\right]^{1/2}$$

is depends on  $\alpha$ . Certainly, also one deduces

$$\alpha \int_0^T |C_\alpha(g)(t)|^2_{L^2(\Gamma_1)} \,\mathrm{d}t \le \|g\|_{L^2(0,T;V')} \|C_\alpha(g)\|_{L^2(0,T;V)},$$

which is uniformly bounded in  $\alpha > 1$ . On the other hand, note that the functions b and q (or  $\psi$ ) intervene to estimate  $u_{0\alpha}$  and  $\dot{u}_{0\alpha}$ .

To show that  $g \mapsto J_{\alpha}(g)$  is strictly convex, one show that

$$(1-\theta)J_{\alpha}(g_{2}) + \theta J_{\alpha}(g_{1}) - J_{\alpha}((1-\theta)g_{1} + \theta g_{2}) =$$
  
=  $\frac{1}{2}\theta(1-\theta)[||u_{g_{1}\alpha} - u_{g_{2}\alpha}||_{\mathcal{H}}^{2} + m||g_{1} - g_{2}||_{\mathcal{H}}^{2}],$ 

for every  $\theta$  in [0,1] and any  $g_1, g_2$  in  $\mathcal{H}$ .

Remark that one has nice estimates for the affine application  $g \mapsto u_{g\alpha}$ , namely

$$\begin{split} \|\nabla u_{g_{1}\alpha} - \nabla u_{g_{2}\alpha}\|_{\mathcal{H}} &\leq \frac{1}{\lambda_{1}} \|g_{1} - g_{2}\|_{L^{2}(0,T;V')}, \\ \sup_{0 \leq t \leq T} \|u_{g_{1}\alpha}(t) - u_{g_{2}\alpha}(t)\|_{H} \leq \frac{1}{\sqrt{\lambda_{1}}} \|g_{1} - g_{2}\|_{L^{2}(0,T;V')}, \\ \|\dot{u}_{g_{1}\alpha} - \dot{u}_{g_{2}\alpha}\|_{L^{2}(0,T;V'_{0})} &\leq \frac{2}{\lambda_{1}} \|g_{1} - g_{2}\|_{L^{2}(0,T;V')}, \\ \|\dot{u}_{g_{1}\alpha} - \dot{u}_{g_{2}\alpha}\|_{L^{2}(0,T;V')} \leq \frac{1+\alpha}{\lambda_{1}} \|g_{1} - g_{2}\|_{L^{2}(0,T;V')}, \\ \|u_{g_{1}\alpha} - u_{g_{2}\alpha}\|_{L^{2}(0,T;L^{2}(\Gamma_{1})} \leq \frac{1}{\sqrt{\lambda_{1}\alpha}} \|g_{1} - g_{2}\|_{L^{2}(0,T;V')}, \end{split}$$

and similarly, for the adjoint state mapping  $g \mapsto p_{g\alpha}$ , one obtain estimates as above replacing  $u_{g_i\alpha}$  with  $p_{g_i\alpha}$ .

On the other hand,  $u_{g_1\alpha} - u_{g_2\alpha}$  is the unique solution of a parabolic variational equality (1.5) with q = 0, b = 0 and  $g = g_1 - g_2$ , i.e.,  $(\partial_t - \Delta)(u_{g_1\alpha} - u_{g_2\alpha}) = g$  in  $L^2(\Omega \times ]0, T[)$  with homogeneous mixed (Robin on  $\Gamma_1$  and Neumann on  $\Gamma_2$ ) boundary conditions. Hence, regularity results implies that  $u_{g_1\alpha} - u_{g_2\alpha}$  belongs to  $L^2(0,T;H^2(\Omega)) \cap H^1(0,T;L^2(\Omega))$ . Similar arguments apply to  $u_{g_1} - u_{g_2}$ , i.e.,  $(\partial_t - \Delta)(u_{g_1} - u_{g_2}) = g$  in  $L^2(\Omega \times ]0,T[)$  with homogeneous mixed (Dirichlet on  $\Gamma_1$  and Neumann on  $\Gamma_2$ ) boundary conditions. Note that some difficulties due to the mixed boundary conditions do arrives, e.g., see Grisvard [9], but our interest is on the asymptotic behavior as  $\alpha$  becomes infinite.

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## 4 Asymptotic Estimates

First one needs to obtain estimates on  $u_{g\alpha}$  and  $p_{g\alpha}$  uniformly in  $\alpha > 1$  and any given g.

**Proposition 4.1.** Under the previous assumptions one has the estimate

$$\begin{aligned} u_{g\alpha} \|_{L^{\infty}(0,T;H)} &+ \|u_{g\alpha}\|_{L^{2}(0,T;V)} + \\ &+ \sqrt{(\alpha-1)} \|u_{g\alpha} - b\|_{L^{2}(\Gamma_{1}\times]0,T[)} \leq C(1 + \|g_{\psi}\|_{L^{2}(0,T;V')}), \end{aligned}$$

$$(4.1)$$

for every  $\alpha > 1$  and any g in  $\mathcal{H}$ , where the constant C depends only on the norms  $\|\dot{u}_g\|_{L^2(0,T;V')}, \|\nabla u_g\|_{L^2(0,T;H)}$ , and the coerciveness constant  $\lambda_1$  in (3.5). Moreover, as  $\alpha \to \infty$  one has  $u_{g\alpha} \to u_g$  strongly in  $L^2(0,T;V) \cap L^{\infty}(0,T;H)$  and  $\dot{u}_{g\alpha} \to \dot{u}_g$  in norm  $L^2(0,T;V'_0)$ .

*Proof.* First note that  $V_0 \subset V$  is a continuous (non dense) inclusion and the norms  $||v||_{V_0} = ||\nabla v||_H$  is equivalently to  $||v||_V = \sqrt{||v||_{V_0} + ||v||_H}$  on  $V_0$ .

Let  $\varphi$  be a function in  $L^2(0, T; V)$  such that  $\dot{\varphi}$  belongs to  $L^2(0, T; V')$ ,  $\varphi(0) = v_b$  and  $\varphi = b$  on  $\Gamma_1$ , e.g., an extension of b and  $v_b$  such as  $\psi$  in (2.2). Now, on the equality (1.5) defining  $u_{g\alpha}$  take  $v = u_{g\alpha}(t) - \varphi(t) := z_{g\alpha}(t)$  to get

$$\langle \dot{u}_{a\alpha}(t), z_{g\alpha}(t) \rangle + \left( \nabla u_{g\alpha}(t), \nabla z_{g\alpha}(t) \right)_{H} + \alpha \langle u_{g\alpha}(t), z_{g\alpha}(t) \rangle_{\Gamma_{1}} = = (g(t), z_{g\alpha}(t))_{H} - \langle q(t), z_{g\alpha}(t) \rangle_{\Gamma_{2}} + \alpha \langle b, z_{g\alpha}(t) \rangle_{\Gamma_{1}},$$

and because  $\varphi = b$  on  $\Gamma_1$  one deduces

$$\frac{1}{2} \frac{\mathrm{d}}{\mathrm{d}t} \| z_{g\alpha}(t) \|_{H}^{2} + \| \nabla z_{g\alpha}(t) \|_{H}^{2} + \alpha \| z_{g\alpha}(t) \|_{L^{2}(\Gamma_{1})}^{2} = (g(t), z_{g\alpha}(t))_{H} - (q(t), z_{g\alpha}(t))_{L^{2}(\Gamma_{2})} - \langle \dot{\varphi}(t), z_{g\alpha}(t) \rangle - (\nabla u_{g}, \nabla z_{g\alpha})_{H},$$
(4.2)

which together with coerciveness (3.5) and the condition  $z_{g\alpha}(0) = 0$  yield the bound (4.1). By means of estimate (4.1), there exists a sequence  $\alpha_n \to \infty$  and z in  $L^2(0,T;V) \cap L^{\infty}(0,T;H)$  such that  $z_{g\alpha_n} \to z$  weakly in  $L^2(0,T;V)$  and weakly\* in  $L^{\infty}(0,T;H)$ , and z = 0 on  $\Gamma_1$ , i.e., z belongs to  $L^2(0,T;V_0)$ .

Hence, note that  $a_{\alpha}(u, v) = a(u, v)$  and  $L_{g\alpha}(t, v) = L_g(t, v)$  if u belongs to  $V_0$ and v belongs to  $V_0$ , and take v in  $V_0$  in the equations (1.4) and (1.5) defining  $u_g$ and  $u_{g\alpha}$  to obtain  $\langle \dot{z}_{g\alpha}, v \rangle + a(z_{g\alpha}, v) = 0$ , for every  $v \in V_0$ . Therefore,  $\dot{z}_{g\alpha_n} \to \dot{z}$ weakly in  $L^2(0, T; V'_0)$  and because  $z_{g\alpha}(0) = 0$  and z = 0 on  $\Gamma_1$ , one deduces z = 0 in  $L^2(0, T; V)$ .

Thus, as  $\alpha \to \infty$  one has  $z_{g\alpha} \to 0$  weakly in  $L^2(0,T;V)$  and weakly<sup>\*</sup> in  $L^{\infty}(0,T;H)$ . It is clear that the inclusion  $V_0 \subset V$  is continuous and because the norm of V restricted to  $V_0$  is equivalent to the norm of  $V_0$ , Hahn-Banach Theorem implies that any element  $\vartheta$  of  $V'_0$  can be extended to an element in V' preserving its norm, in particular  $\dot{u}_g$  can be extended to be an element in  $L^2(0,T;V')$ . Then, take  $\varphi = u_g$  in the equality (4.2) and considering  $\dot{u}_g$  an element in  $L^2(0,T;V')$ , one deduces that the convergence of  $u_{g\alpha}$  toward  $u_g$  is indeed strongly in  $L^2(0,T;V) \cap L^{\infty}(0,T;H)$ . Moreover,  $z_{g\alpha} \to 0$  in norm  $L^2(\Gamma \times ]0,T[)$  and  $\dot{z}_{g\alpha} \to 0$  in norm  $L^2(0,T;V')$ .

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**Proposition 4.2.** Under the previous assumptions one has the estimate

$$\begin{aligned} \|p_{g\alpha}\|_{L^{\infty}(0,T;H)} + \|p_{g\alpha}\|_{L^{2}(0,T;V)} + \\ + \sqrt{(\alpha-1)} \|p_{g\alpha}\|_{L^{2}(\Gamma_{1}\times]0,T[)} \le C(1+\|u_{g\alpha}\|_{L^{2}(0,T;V')}), \end{aligned}$$
(4.3)

for every  $\alpha > 1$  and any g in  $\mathcal{H}$ , where the constant C depends only on the norms  $\|z_d\|_{\mathcal{H}}, \|\dot{p}_g\|_{L^2(0,T;V')}, \|\nabla p_g\|_{L^2(0,T;H)}$ , and the coerciveness constant  $\lambda_1$  in (3.5). Moreover, as  $\alpha \to \infty$  one has  $p_{g\alpha} \to p_g$  strongly in  $L^2(0,T;V) \cap L^{\infty}(0,T;H)$  and  $\dot{p}_{g\alpha} \to \dot{p}_g$  in norm  $L^2(0,T;V'_0)$ .

*Proof.* Note that even when  $b \neq 0$  the (Robin) boundary condition of  $p_g$  and  $p_{g\alpha}$  on  $\Gamma_1$  does not involve b directly. Certainly, the norm  $||u_{g\alpha}||_{L^2(0,T;V')}$  is bounded by  $||u_{g\alpha}||_{L^2(0,T;H)}$ , which is uniformly bounded in  $\alpha$ .

The technique used in Proposition 4.1 applies for the adjoint states  $p_{g\alpha}$  and  $p_g$ . Perhaps the only point to remark is the convergence as  $\alpha \to \infty$ . Indeed, one needs to make use of the weak (and later strong) convergence  $u_{g\alpha} \to u_g$  in  $L^2(0,T;V')$ , which is deduced for the convergence in  $L^2(0,T;H)$ .

# 5 Optimal Control Problems

We are now ready to consider the distributed control problems (1.9) and (1.10). Our purpose is to establish

**Theorem 5.1.** Let assumptions (2.1) and (2.2) be hold, and  $\hat{g}$  and  $\hat{g}_{\alpha}$  be the minimizers in  $\mathcal{H}$  of problems (1.9) and (1.10), respectively. Then, as the parameter  $\alpha \to \infty$ , the minimizers  $\hat{g}_{\alpha} \to \hat{g}$  strongly in  $\mathcal{H}$ . Moreover the corresponding optimal state and adjoint state satisfy  $(u_{\hat{g}_{\alpha}\alpha}, \dot{u}_{\hat{g}_{\alpha}\alpha}) \to (u_{\hat{g}}, \dot{u}_{\hat{g}})$  and  $(p_{\hat{g}_{\alpha}\alpha}, \dot{p}_{\hat{g}_{\alpha}\alpha}) \to (p_{\hat{g}}, \dot{p}_{\hat{g}})$  strongly in  $L^2(0, T; V) \times L^2(0, T; V'_0)$ .

*Proof.* We make several steps. First, be means of the estimate (4.1) in Proposition 4.1 one has

 $\|u_{0\alpha}\|_{\mathcal{H}} \le C, \quad \forall \alpha > 1,$ 

for some constant C. Now, from the inequality  $J(\hat{g}_{\alpha}) \leq J(0)$  we deduce

 $\|\hat{g}_{\alpha}\|_{\mathcal{H}} + \|u_{\hat{g}_{\alpha}\alpha}\|_{\mathcal{H}} \le C, \quad \forall \alpha > 1$ 

for some constant independent of  $\alpha > 1$ .

Again, estimate (4.1) in Proposition 4.1 and estimate (4.2) in Proposition 4.2 yield

$$\begin{split} \|u_{\hat{g}_{\alpha}\alpha}\|_{L^{2}(0,T;V)} + \|\dot{u}_{\hat{g}_{\alpha}\alpha}\|_{L^{2}(0,T;V'_{0})} + \\ &+ \sqrt{(\alpha-1)} \|u_{\hat{g}_{\alpha}\alpha} - b\|_{L^{2}(0,T;L^{2}(\Gamma_{1}))} \leq C, \quad \forall \alpha > 1 \end{split}$$

and

$$\begin{split} \|p_{\hat{g}_{\alpha}\alpha}\|_{L^{2}(0,T;V)} + \|\dot{p}_{\hat{g}_{\alpha}\alpha}\|_{L^{2}(0,T;V'_{0})} + \\ &+ \sqrt{(\alpha-1)} \|p_{\hat{g}_{\alpha}\alpha}\|_{L^{2}(0,T;L^{2}(\Gamma_{1}))} \leq C, \quad \forall \alpha > 1. \end{split}$$

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Hence, there exist  $\bar{g}$  in  $\mathcal{H}$ ,  $\hat{u}$  and  $\hat{p}$  in  $L^2(0,T;V_0)$  with  $\dot{\hat{u}}$  and  $\dot{\hat{p}}$  in  $L^2(0,T;V'_0)$ such that, for a convenient subsequence as  $\alpha \to \infty$  we has  $\hat{g}_{\alpha} \rightharpoonup \bar{g}$  weakly in  $\mathcal{H}$ ,  $u_{\hat{g}_{\alpha}\alpha} \rightharpoonup \hat{u}$  weakly in  $L^2(0,T;V)$ ,  $\dot{u}_{\hat{g}_{\alpha}\alpha} \rightharpoonup \hat{u}$  weakly in  $L^2(0,T;V'_0)$ ,  $p_{\hat{g}_{\alpha}\alpha} \rightharpoonup \hat{p}$ weakly in  $L^2(0,T;V)$ ,  $\dot{p}_{\hat{g}_{\alpha}\alpha} \rightharpoonup \dot{\hat{p}}$  weakly in  $L^2(0,T;V'_0)$ .

By taking v in  $V_0$  in the parabolic variational equality (2.5) and letting  $\alpha \to \infty$  we deduce that  $\hat{u}$  solves parabolic variational equality (2.4), and by uniqueness  $\hat{u} = u_{\hat{g}}$ . In particular  $u_{\hat{g}_{\alpha}\alpha} \rightharpoonup u_{\hat{g}}$  weakly in  $L^2(0,T;V'_0)$ . Thus, by taking v in  $V_0$  in the parabolic variational equality defining the adjoint state  $p_{\hat{g}_{\alpha}\alpha}$  in Proposition 3.2 and letting  $\alpha \to \infty$  we deduce that  $\hat{p} = p_{\bar{g}}$ . On the other hand, taking limit in the equality  $m\hat{g}_{\alpha} + p_{\hat{g}_{\alpha}\alpha} = 0$  we deduce  $m\bar{g} + p_{\bar{g}} = 0$ . Thus, by using Proposition 3.1, this proves that  $\bar{g}$  is a minimizer for the control problem (1.9), and by uniqueness  $\hat{g} = \bar{g}$ .

At this point, we have

$$(\hat{g}_{\alpha}, u_{\hat{g}_{\alpha}\alpha}, \dot{u}_{\hat{g}_{\alpha}\alpha}, p_{\hat{g}_{\alpha}\alpha}, \dot{p}_{\hat{g}_{\alpha}\alpha}) \rightharpoonup (\hat{g}, u_{\hat{g}}, \dot{u}_{\hat{g}}, p_{\hat{g}}, \dot{p}_{\hat{g}})$$

weakly in the corresponding spaces, initially for a convenient subsequence as  $\alpha \to \infty$ , but in view of the uniqueness of the limit, the weak convergence whole as  $\alpha \to \infty$ .

To prove the strong convergence we use the weak semicontinuity of the norm and the optimality of  $\hat{g}, \hat{g}_{\alpha}$ , namely,

$$J(\hat{g}) = \frac{1}{2} \|u_{\hat{g}} - z_d\|_{\mathcal{H}}^2 + \frac{m}{2} \|\hat{g}\|_{\mathcal{H}}^2 \le \liminf_{\alpha \to \infty} \left[\frac{1}{2} \|u_{\hat{g}_{\alpha}\alpha} - z_d\|_{\mathcal{H}}^2 + \frac{m}{2} \|\hat{g}_{\alpha}\|_{\mathcal{H}}^2\right] \le \limsup_{\alpha \to \infty} \left[\frac{1}{2} \|u_{\hat{g}_{\alpha}\alpha} - z_d\|_{\mathcal{H}}^2 + \frac{m}{2} \|\hat{g}_{\alpha}\|_{\mathcal{H}}^2\right] \le \limsup_{\alpha \to \infty} J_{\alpha}(g),$$

for any g in  $\mathcal{H}$ . In view of Proposition 4.1,  $u_{g\alpha} \to u_g$  strongly in  $L^2(0,T;V)$  as  $\alpha \to \infty$ , which implies that

$$\limsup_{\alpha \to \infty} J_{\alpha}(g) = \lim_{\alpha \to \infty} \left[ \frac{1}{2} \| u_{g\alpha} - z_d \|_{\mathcal{H}}^2 + \frac{m}{2} \| g \|_{\mathcal{H}}^2 \right] = J(g).$$

By taking infimum on g, all the above inequalities become equalities and therefore

$$\frac{1}{2} \|u_{\hat{g}_{\alpha}\alpha} - z_d\|_{\mathcal{H}}^2 + \frac{m}{2} \|\hat{g}_{\alpha}\|_{\mathcal{H}}^2 \to \frac{1}{2} \|u_{\hat{g}} - z_d\|_{\mathcal{H}}^2 + \frac{m}{2} \|\hat{g}\|_{\mathcal{H}}^2.$$

This and the weak convergence imply that  $(\hat{g}_{\alpha}, u_{\hat{g}_{\alpha}\alpha}) \rightarrow (\hat{g}, u_{\hat{g}})$  strongly in  $\mathcal{H} \times \mathcal{H}$ , as  $\alpha \to \infty$ .

Finally, if  $z_{\alpha} = u_{\hat{g}_{\alpha}\alpha} - u_{\hat{g}}$  then we deduce

$$\int_0^T \left[ \langle \dot{z}_{\alpha}(t), z_{\alpha}(t) \rangle + a_1(z_{\alpha}(t), z_{\alpha}(t)) + (\alpha - 1) \int_{\Gamma_1} |z_{\alpha}(x, t)|^2 \, \mathrm{d}x \right] \, \mathrm{d}t \le \\ \le \int_0^T \left[ \langle \hat{g}_{\alpha} - \dot{u}_{\hat{g}}, z_{\alpha} \rangle - a(u_{\hat{g}}, z_{\alpha}) - \int_{\Gamma_2} q(x, t) z_{\alpha}(x, t) \, \mathrm{d}x \right] \, \mathrm{d}t.$$

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Since  $z_{\alpha} \to 0$  weakly in  $L^2(0,T;V)$  and  $\hat{g}_{\alpha} \to \hat{g}$  strongly in  $\mathcal{H}$ , we obtain  $u_{\hat{g}_{\alpha}\alpha} \to u_{\hat{g}}$  strongly in  $L^2(0,T;V)$ , as  $\alpha \to \infty$ . Now, going back to the equation one has

$$\langle \dot{z}_{\alpha}(t), v \rangle + a(z_{\alpha}(t), v) = \langle \hat{g}_{\alpha} - \hat{g}, v \rangle.$$

Now, taking sup for v in  $V_0$  with  $||v_0||_{V_0} \leq 1$  and integrating in ]0, T[ one obtains the strong convergence of the time derivative. Similarly,  $(p_{\hat{g}_{\alpha}\alpha}, \dot{p}_{\hat{g}_{\alpha}\alpha}) \to (p_{\hat{g}}, \dot{p}_{\hat{g}})$ strongly in  $L^2(0, T; V) \times L^2(0, T; V'_0)$ , as  $\alpha \to \infty$ . This completes the proof.  $\Box$ 

Also we have

**Proposition 5.2.** If  $\alpha_2 \ge \alpha_1 \ge \alpha_0 > 0$  then there exists a constant  $C = C_{\alpha_0}$  such that for every g in  $\mathcal{H}$  one has

$$\|u_{g\alpha_1} - u_{g\alpha_2}\|_{L^2(0,T;V)} \le C_{\alpha_0}(\alpha_2 - \alpha_1) \|b - u_{g\alpha_2}\|_{L^2(0,T;H^{-1/2}(\Gamma_1))}, \quad (5.1)$$

and

$$\|p_{g\alpha_1} - p_{g\alpha_2}\|_{L^2(0,T;V)} \le C_{\alpha_0}(\alpha_2 - \alpha_1) \big( \|p_{g\alpha_2}\|_{L^2(0,T;H^{-1/2}(\Gamma_1))} + \\ + \|b - u_{g\alpha_2}\|_{L^2(0,T;H^{-1/2}(\Gamma_1))} \big),$$

$$(5.2)$$

i.e., the dependency in  $\alpha$  is Lipschitz continuous.

*Proof.* For a fixed g and  $\alpha_2 \ge \alpha_1 \ge \alpha_0 > 0$  set  $z = u_{g\alpha_2} - u_{g\alpha_1}$  to obtain from the equation (1.5) with  $\alpha_i$  the identity

$$\langle \dot{z}(t), v \rangle + a_{\alpha_1}(z(t), v) = (\alpha_2 - \alpha_1) \int_{\Gamma_1} (b - u_{g\alpha_2}) v \, \mathrm{d}\gamma, \quad \forall v \in V.$$

By taking v = z(t) and by means of the inequalities

$$\left|\int_{0}^{T} \mathrm{d}t \int_{\Gamma_{1}} (b - u_{g\alpha_{2}}) z \,\mathrm{d}\gamma\right| \leq C_{0} \|b - u_{g\alpha_{2}}\|_{L^{2}(0,T;H^{-1/2}(\Gamma_{1}))} \|z\|_{L^{2}(0,T;V)}$$

and

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$$a_{\alpha}(v,v) \ge \lambda(\alpha_0) \|v\|_V^2, \quad \forall v \in V, \ \alpha \ge \alpha_0,$$

we deduce the desired estimate with  $C_{\alpha_0} = C_0/\lambda(\alpha_0)$ .

Similarly, for a fixed g and  $\alpha_2 \ge \alpha_1 \ge \alpha_0 > 0$  set  $w = p_{g\alpha_2} - p_{g\alpha_1}$  to obtain from the equation (3.5) with  $\alpha_i$  the identity

$$\langle \dot{w}(t), v \rangle + a_{\alpha_1}(w(t), v) = (\alpha_1 - \alpha_2) \int_{\Gamma_1} p_{g\alpha_2} v \,\mathrm{d}\gamma + (u_{g\alpha_2} - u_{g\alpha_1}, v)_H,$$

for every v in V. By taking v = w(t) and in view of the estimate (5.1), we conclude.

Under some more restrict assumption we have monotonicity on  $\alpha$ 

**Proposition 5.3.** Let us assume the data b constant on  $\Gamma_1$ ,  $v_b \leq b$  on  $\Omega$ ,  $g \leq 0$  in  $\Omega \times ]0, T[$  and  $q \geq 0$  on  $\Gamma_2 \times ]0, T[$ . Then  $u_{g\alpha} \leq u_g \leq b$  for every  $\alpha > 0$ . Moreover, if  $0 < \alpha_1 \leq \alpha_2$  then  $u_{g\alpha_1} \leq u_{g\alpha_2} \leq u_g \leq b$  in  $\Omega \times ]0, T[$ . Furthermore, if  $b \leq z_d$  in  $\Omega \times ]0, T[$  then  $p_{g\alpha_1} \leq p_{g\alpha_2} \leq p_g \leq 0$  in  $\Omega \times ]0, T[$ , for every  $\alpha_2 \geq \alpha_1 > 0$ .

*Proof.* First, the maximum principle implies that  $u_{g\alpha} \leq b$ . Indeed, if  $z = (u_{g\alpha} - b)$  then we have

$$\langle \dot{z}(t), z^+(t) \rangle + a \bigl( z(t), z^+(t) \bigr) + \alpha \int_{\Gamma_1} z(t) z^+(t) \, \mathrm{d}\gamma =$$
$$= \bigl( g(t), z^+(t) \bigr) - \int_{\Gamma_2} q(t) z^+(t) \, \mathrm{d}\gamma$$

after using the fact that b is constant, which implies  $z^+ = 0$ .

Similarly, if  $w = u_{g\alpha_2} - u_{g\alpha_1}$  with  $\alpha_2 > \alpha_1$  then we get

$$\langle \dot{w}(t), w^+(t) \rangle + a_{\alpha_1} (w(t), w^+(t)) + (\alpha_2 - \alpha_1) \int_{\Gamma_1} (b - u_{g\alpha_2}(t) z^+(t) \, \mathrm{d}\gamma = 0,$$

which yields  $w \leq 0$ , i.e.,  $u_{g\alpha_2} \leq u_{g\alpha_1}$ .

Finally, if  $y = u_{g\alpha} - u_g$  then we obtain

$$\langle \dot{y}(t), y^+(t) \rangle + a \big( y(t), y^+(t) \big) + \alpha \int_{\Gamma_1} (b - u_{g\alpha}(t) y^+(t) \, \mathrm{d}\gamma = 0,$$

which yields  $y \leq 0$ , i.e.,  $u_{g\alpha} \leq u_g$ .

The estimate on the adjoint state follows from a comparison with the solution r of the parabolic variational equality with terminal condition

$$\begin{cases} r \in L^{2}(0,T;V), \quad r(T) = 0 \text{ and } \dot{r} \in L^{2}(0,T;V') \\ \text{such that} \quad -\langle \dot{r}(t), v \rangle + a(r(t), v) = (b - z_{d}, v)_{H}, \quad \forall v \in V. \end{cases}$$
(5.3)

Indeed, if  $b \leq z_d$  in  $\Omega \times ]0, T[$  then the maximum principle (as above) yields  $p_g \leq r \leq 0$ . Next, similarly to the state u with b = 0, one deduces that  $p_{g\alpha_1} \leq p_{g\alpha_2} \leq p_g \leq r \leq 0$  in  $\Omega \times ]0, T[$ , for every  $\alpha_2 \geq \alpha_1 > 0$ .

Certainly, the maximum principle yields  $u_{g_1} \leq u_{g_2}$  and  $u_{g_1\alpha} \leq u_{g_2\alpha}$  if  $g_1 \leq g_2$ , but a priori, it is not clear when the minimizers satisfy  $\hat{g} \geq \hat{g}_{\alpha}$  to deduce the monotonicity  $u_{g_{\alpha_1}\alpha_1} \leq u_{g_{\alpha_2}\alpha_2} \leq u_{g_{\alpha}} \leq b$ .

# 6 Final Comments

Variational inequalities was popular in the 70's, most of the main techniques for parabolic variational inequalities can be found in various classic books, e.g., Bensoussan and Lions [5], among other.

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It is well known that the regularity of the mixed problem is problematic when both portions of the boundary  $\Gamma_1$  and  $\Gamma_2$  have a nonempty intersection, e.g. see the book Grisvard [9]. Recently, sufficient conditions (on the data) to obtain a  $H^2$  regularity for a (elliptic) mixed boundary conditions are given in Bacuta et al. [3], see also Azzam and Kreyszig [1], among others.

Numerical analysis of a parabolic PDE with mixed boundary conditions (Dirichlet and Neumann) is studied in Babuska and Ohnimus [2], while a parabolic control problem with Robin boundary conditions is considered in Chrysafinos et al. [7] and Bergounioux and Troltzsch [6].

The state equation, i.e., a parabolic PDE with mixed boundary conditions (Robin and Neumann) has been discussed in Ben Belgacem et al. [4] and Tarzia [12].

Certainly, there are several possible extensions, e.g., a state equation of the form

$$\partial_t u - \operatorname{div}(A(x,t)\nabla u) + b(t,x)u = f$$
 in  $\Omega \times ]0,T[,$ 

with mixed boundary conditions. A carefully analysis is necessary, but the main techniques used to let  $\alpha \to \infty$  in the parabolic variational inequality seems to be very well adaptable to more general situations.

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