Gate-Bias Induced R_{ON} Instability in p-GaN Power HEMTs

Alessandro Chini, Nicolò Zagni, Member, IEEE, Giovanni Verzellesi, Senior Member, IEEE, Marcello Cioni, Student Member, IEEE, Giovanni Giorgino, Maria Concetta Nicotra, Maria Eloisa Castagna, and Ferdinando Iucolano

Abstract—In this letter, we investigate the on-resistance (R_{ON}) instability in p-GaN power HEMTs induced by a positive or negative gate bias (V_{GB}) , following the application of a quasi-static initialization voltage (V_{GP}) of opposite sign. The transient behavior of this instability was characterized at different temperatures in the 90–135 °C range. By monitoring the resulting drain current transients, the activation energy as well as time constants of the processes are characterized. Not trivially, both R_{ON} increase/decrease were found to be thermally activated and with same activation energy. We attribute the thermal activation of both R_{ON} increase/decrease to the charging/discharging of hole traps present in the AlGaN barrier in the region below the gate.

Index Terms—pGaN HEMTs, Dynamic ON-Resistance, Instability, NBTI, PBTI, Barrier Traps.

I. INTRODUCTION

 \mathbf{C} EVERAL studies are present in the literature that involve the Dinvestigation of the stability of normally-off p-GaN high electron mobility transistors (HEMTs) after applying gate bias stress [1]–[9]. These studies mostly focus on the instability of the threshold voltage (V_T) [2], [4]–[10] caused by dynamic effects occurring at (or near) the device region under the gate [11]. On the other hand, a detailed analysis of the consequences of gate biasing on device ON-resistance (R_{ON}) has not been presented yet, despite $R_{\rm ON}$ instability can be a more serious concern from the application standpoint (i.e., at

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A. Chini, and N. Zagni with the Department of Engineering "Enzo Ferrari", Università di Modena e Reggio Emilia, 41125, Modena, Italy. (*Corresponding author*: N. Zagni. E-mail: <u>nicolo.zagni@unimore.it</u>).

M. Cioni, and G. Giorgino are with the Department of Engineering "Enzo Ferrari", Università di Modena e Reggio Emilia, 41125, Modena, Italy and with STMicroelectronics, 95121, Catania, Italy.

G. Verzellesi is with Department of Sciences and Methods for Engineering (DISMI) and EN&TECH Center, University of Modena and Reggio Emilia, 42122, Reggio Emilia, Italy.

M. C. Nicotra, M. E. Castagna, and F. Iucolano are with STMicroelectronics, 95121, Catania, Italy.

large overdrive) than the more widely investigated $V_{\rm T}$ drift [7]. Recently, a preliminary characterization of the effects of negative/positive gate bias on $R_{\rm ON}$ in devices of a previous generation compared to the ones analyzed in this letter was presented in [12], which lacked however the physical interpretation of the observed instabilities.

In this letter, we investigate the influence of gate bias on the $R_{\rm ON}$ instability and characterize its time- and temperaturedependent behavior. First, we characterize $R_{\rm ON}$ after applying a quasi-static (50 min) gate bias ($V_{\rm GP}$) of different magnitudes between -6 and +6 V at a baseplate temperature of 90 °C. Then, we characterize the drain current transient response at different temperatures in the 75–135 °C range by applying baseline biases ($V_{\rm GB}$) with different magnitudes and by periodically sensing $R_{\rm ON}$. The obtained transients allow characterizing the time constants (τ) and activation energy ($E_{\rm A}$) of the processes involved. The experiments reveal a set of features of $R_{\rm ON}$ instability which we attribute to charging/discharging of hole traps present in the AlGaN barrier in the region below the gate.

II. EXPERIMENTAL CHARACTERIZATION

The devices under test (DUT's) are 650-V, packaged, normallyoff, p-GaN power HEMTs with a Schottky gate contact. Typical drain current (I_D) and gate current (I_G) vs gate-to-source bias (V_{GS}) characteristics at a drain-to-source bias (V_{DS}) of =50 mV and T = 90 °C are reported in Fig. 1(a), and (b). Figure 1(c) shows the R_{ON} variation occurring after applying a constant V_{GS} bias (V_{GP}) for 50 minutes with different magnitudes. R_{ON} is computed as the ratio between $V_{DS} = 50$ mV and I_D measured at $V_{GS} = 6$ V. Interestingly, little R_{ON} variation was observed after applying $V_{GP} \ge 1.5$ V, whereas R_{ON} increased linearly with decreasing V_{GP} below 1.5 V, resulting in a $\approx 17\%$ increase after $V_{GP} = -6$ V with respect to its minimum value, see Fig. 1(c). R_{ON} time evolution was investigated by applying to the gate

terminal the voltage waveforms depicted by applying to the gate terminal the voltage waveforms depicted in Fig. 2(a). After the 50-minutes initialization bias at V_{GP} , V_{GS} is settled to a baseline bias (V_{GB}) while periodically sampling (i.e., every 50 ms) I_D at $V_{GM} = 6$ V. A short, 2-ms V_{GM} pulse is applied to limit possible influence on the state of the device. This technique allows to acquire dependable data for wider time and temperature windows than the one employed in [12]. By employing a combination of different (V_{GP} , V_{GB}), several Drain Current Transients (DCTs) are extracted, from which the R_{ON} instability transients can be derived. The initialization phase with positive

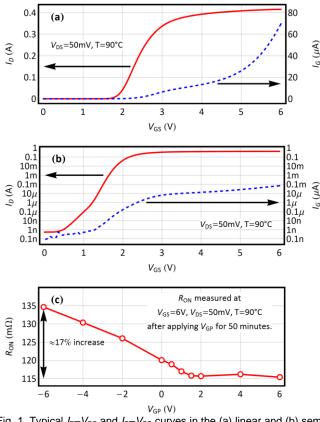


Fig. 1. Typical $I_{\rm D}-V_{\rm GS}$ and $I_{\rm G}-V_{\rm GS}$ curves in the (a) linear and (b) semilog scale obtained at $V_{\rm DS} = 50$ mV and T = 90 °C. (c) Observed $R_{\rm ON}$ variation obtained after applying a $V_{\rm GP}$ bias for 50 minutes with varying magnitude. Maximum $R_{\rm ON}$ variation after $V_{\rm GP} = -6$ V is ≈17% of its minimum value obtained at $V_{\rm GP} \ge 1.5$ V.

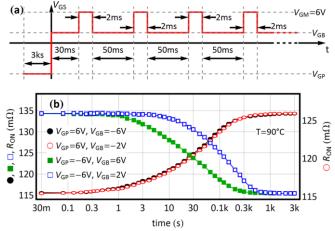


Fig. 2. (a) Voltage waveform applied at the gate terminal for the characterization of I_D/R_{ON} transients. V_{GP} is first applied for 3 ks (i.e., 50 minutes), then V_{GB} is applied and DUT's R_{ON} is periodically sensed at $V_{GM} = 6$ V for 2 ms every 50 ms, except for the first sensing done after 30 ms. (b) Results of the characterization in terms of R_{ON} vs time for four different (V_{GP} , V_{GB}) combinations (see legend).

(negative) V_{GP} ensures that the traps in the barrier (responsible for the instability, as discussed in Sec. III) are close to being fully occupied (empty) prior to the subsequent R_{ON} transient characterization. This comes as a consequene of the chosen V_{GP} and V_{GB} , which are always of opposite sign. Figure 2(b) shows the results of the characterization for a selection of the (V_{GP} , V_{GB}) combinations, namely: (+6, -6), (+6, -2), (-6, +6), and (-6, +2) V. In each case, an exponential-like evolution of I_{D} , and

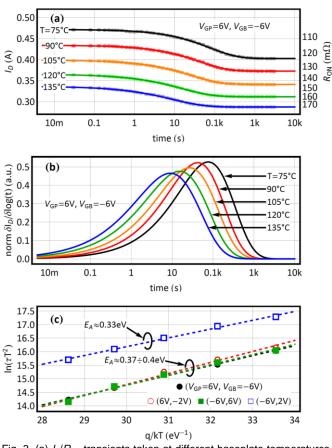


Fig. 3. (a) I_D/R_{ON} transients taken at different baseplate temperatures, *T*, see plot, at (V_{GP} , V_{GB}) = (+6, -6) V. (b) $\partial I_D/\partial [\log(t)]$ corresponding to (a) from which *r*'s are extracted. (c) Arrhenius plots built from the extracted *r*'s for different (V_{GP} , V_{GB}) combinations, see legend. *E*_A's extracted from the linear fitting of the plots are also indicated.

therefore R_{ON} , is obtained. We observe that when applying a negative (positive) V_{GB} after a positive (negative) V_{GP} , R_{ON} increases (decreases). Note that the first three combinations lead to a similar time evolution, whereas a slower transient is obtained for $(V_{\text{GP}}, V_{\text{GB}}) = (-6, +2)$ V. We mention that although the first data point is acquired at 30 ms after the initialization phase is completed, no significant R_{ON} variation occurs for shorter times. This is deduced from the fact that transients taken for equal and opposite $(V_{\text{GP}}, V_{\text{GB}})$ are within the same R_{ON} range. For instance, R_{ON} transient obtained at T = 90 °C with (V_{GP} , V_{GB}) = (-6, +6) V is bound between \approx (135, 115) m Ω and the one obtained with $(V_{\text{GP}}, V_{\text{GB}}) = (+6, -6)$ V is bound between $\approx (115, -6)$ 135) m Ω , see Fig. 2(b). Additional R_{ON} transients acquired by applying $V_{GB} = -6/+6$ V sequentially for three consecutive cycles (not shown) also reveal that R_{ON} stays within the bounds identified previously, further indicating that no additional effects take place for very short times.

DCTs were further characterized at different temperatures to extract the activation energy (E_A) of the processes leading to R_{ON} instability. DCTs obtained at T = 75-135 °C for (V_{GP} , V_{GB}) = (-6, +6) V are shown in Fig. 3(a). As it can be seen, the exponential behavior of I_D vs time is maintained, while the associated time constants (τ 's) decrease with increasing T. τ 's are determined as the instants corresponding to peaks in the $\partial I_D/\partial [\log(t)]$ signals [13]. Figure 3(b) shows the normalized $\partial I_D/\partial [\log(t)]$ curves, from which the τ 's are readily extracted

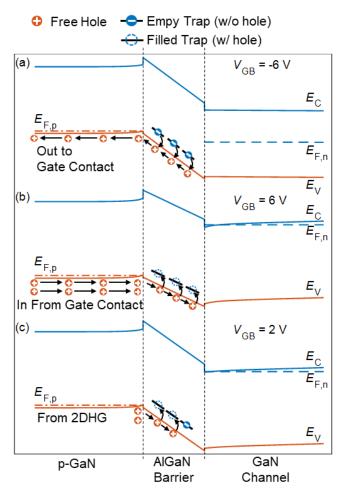


Fig. 4. Sketch of the band diagrams at (a)–(c) $V_{GB} = (-6, 6, 2)$ V illustrating the processes leading to R_{ON} instability. At $V_{GB} = -6$ V, see (a), holes are emitted from barrier traps and collected by the gate contact leaving negatively ionized charges which increase R_{ON} . This mechanism also occurs at $V_{GB} = -2$ V. At $V_{GB} = 6$ V, see (b), holes are injected by the reverse biased Schottky gate contact that get trapped into barrier traps decreasing R_{ON} . At $V_{GB} = 2$ V, see (c), holes are supplied by the 2DHG at the pGaN/AlGaN interface. These holes also get trapped into barrier traps thus decreasing R_{ON} but with longer r's with respect to (b).

and then used to build the Arrhenius plots reported in Fig. 3(c) for the (V_{GP} , V_{GB}) conditions previously discussed. Interestingly, the linear fitting of the Arrhenius plots corresponding to the DCTs with (V_{GP} , V_{GB}) = (+6, -6), (+6, -2), (-6, +6) V yield similar signatures with same E_A 's in the 0.37–0.4 eV range. On the other hand, the DCT at (V_{GP} , V_{GB}) = (-6, +2) V shows higher τ 's compared to the previous conditions for all *T*, and the extracted E_A is slightly lower, ≈ 0.33 eV.

III. DISCUSSION

The observed features of the $R_{\rm ON}$ transients can be attributed to charge/discharge of hole traps present in the AlGaN barrier in the region below the gate. Figure 4 shows a sketch of the band diagrams under different $V_{\rm GB}$ that we use as an aid to the discussion. At $V_{\rm GB} = -6/-2$ V, see Fig. 4(a), trapped holes in the barrier are emitted from traps and are collected by the gate contact, leaving behind negatively ionized charges which increase $R_{\rm ON}$. $E_A \approx 0.37-0.4$ eV extracted in Fig. 3(b) in this condition corresponds to the energy difference between the trap

level and the AlGaN valence band edge. At $V_{GB} = 6$ V, see Fig. 4(b), holes injected by the reverse biased Schottky gate contact get trapped by the negatively ionized barrier traps, which decreases R_{ON} . *T*-activation of trapping with the same activation energy $E_A \approx 0.37-0.4$ eV characterizing the hole emission process is not a trivial finding. A possible explanation is that under low hole injection conditions and relatively high trap concentration in the AlGaN barrier, hole density *p* is about proportional to $\exp[-(E_T-E_V)/kT]$. As a result, also the hole capture rate, that is proportional to *p* is thermally activated with the same $E_A = E_T-E_V \approx 0.37-0.4$ eV [14]. Considering the reported Mg energy level in GaN in the range of 0.16–0.26 eV [11], [15], [16] and the AlGaN/GaN valence band offset [17], the extracted E_A range is compatible with Mg out-diffused in the barrier during epitaxial growth [18].

Finally, at $V_{\text{GB}} = 2$ V, Fig. 4(c), shows that holes are negligibly injected by the gate contact, which are thus slowly supplied to the barrier traps by the two-dimensional hole gas (2DHG) present at the pGaN/AlGaN heterojunction [6]. These holes also get trapped into barrier traps thus decreasing R_{ON} similarly to the previous case but with longer τ 's due to the less effective hole injection mechanism within the AlGaN barrier. The injection mechanism also affects the *T*-dependence of R_{ON} decrease in this case. Being limited also by the hole injection mechanism however, the extracted activation energy ($E_A \approx 0.33$ eV) does not correspond to the trap energy level as in the previous cases.

IV. CONCLUSIONS

We investigated the on-resistance (R_{ON}) instability in p-GaN power HEMTs induced by gate biases. Results reveal that the activation energy (E_A) associated to the R_{ON} instabilities is to a first approximation unaffected by the applied gate bias, although the time constants (τ 's) depend on it. The thermal activation of the R_{ON} transients points to a charging/discharging mechanism of hole traps present in the AlGaN barrier in the region below the gate.

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