

Photometric IGM tomography with Subaru/HSC: the large-scale structure of Ly α emitters and IGM transmission in the COSMOS field at $z \sim 5$

Koki Kakiichi,^{1*} Joseph F. Hennawi,^{1,2} Yoshiaki Ono,³ Akio K. Inoue,^{4,5} Masami Ouchi,^{6,3,7}
Richard S. Ellis,⁸ Romain A. Meyer,⁹ and Sarah I. Bosman⁹

¹*Department of Physics, Broida Hall, University of California, Santa Barbara, CA 93106-9530, USA*

²*Leiden Observatory, Leiden University, Niels Bohrweg 2, 2333 CA Leiden, Netherlands*

³*Institute for Cosmic Ray Research, The University of Tokyo, 5-1-5 Kashiwanoha, Kashiwa, Chiba 277-8582, Japan*

⁴*Waseda Research Institute for Science and Engineering, Faculty of Science and Engineering, Waseda University, 3-4-1, Okubo, Shinjuku, Tokyo 169-8555, Japan*

⁵*Department of Physics, School of Advanced Science and Engineering, Faculty of Science and Engineering, Waseda University, 3-4-1, Okubo, Shinjuku, Tokyo 169-8555, Japan*

⁶*National Astronomical Observatory of Japan, 2-21-1 Osawa, Mitaka, Tokyo 181-8588, Japan*

⁷*Kavli Institute for the Physics and Mathematics of the Universe (WPI), University of Tokyo, Kashiwa, Chiba 277-8583, Japan*

⁸*Department of Physics and Astronomy, University College London, Gower Street, London WC1E 6BT, UK*

⁹*Max Planck Institut für Astronomie, Königstuhl 17, D-69117, Heidelberg, Germany*

Accepted XXX. Received YYY; in original form ZZZ

ABSTRACT

We present a novel technique called “photometric IGM tomography” to map the intergalactic medium (IGM) at $z \approx 4.9$ in the COSMOS field. It utilizes deep narrow-band (NB) imaging to photometrically detect faint Ly α forest transmission in background galaxies across the Subaru/Hyper-Suprime Cam (HSC)’s 1.8 sq. deg field of view and locate Ly α emitters (LAEs) in the same cosmic volume. Using ultra-deep HSC images and Bayesian spectral energy distribution fitting, we measure the Ly α forest transmission at $z \approx 4.9$ along a large number (140) of background galaxies selected from the DEIMOS10k spectroscopic catalogue at $4.98 < z < 5.89$ and the SILVERRUSH LAEs at $z \approx 5.7$. We photometrically measure the mean Ly α forest transmission and achieve a result consistent with previous measurements based on quasar spectra. We also measure the angular LAE-Ly α forest cross-correlation and Ly α forest auto-correlation functions and place an observational constraint on the large-scale fluctuations of the IGM around LAEs at $z \approx 4.9$. Finally, we present the reconstructed 2D tomographic map of the IGM, co-spatial with the large-scale structure of LAEs, at a transverse resolution of $11 h^{-1} \text{cMpc}$ across $140 h^{-1} \text{cMpc}$ in the COSMOS field at $z \approx 4.9$. We discuss the observational requirements and the potential applications of this new technique for understanding the sources of reionization, quasar radiative history, and galaxy-IGM correlations across $z \sim 3 - 6$. Our results represent the first proof-of-concept of photometric IGM tomography, offering a new route to examining early galaxy evolution in the context of the large-scale cosmic web from the epoch of reionization to cosmic noon.

Key words: methods: observational – intergalactic medium – dark ages, reionization, first stars – large-scale structure of Universe

1 INTRODUCTION

Cosmography, i.e. the science of mapping the Universe, is a fundamental pillar of astronomy. Mapping the distribution of objects on the sky has been pivotal for the discovery of the large-scale structure of the Universe. Our modern cosmological model is largely based on the maps of cosmic microwave background fluctuations (e.g. Planck Collaboration et al. 2020), the large-scale distribution of galaxies (e.g. eBOSS Collaboration: Alam et al. 2021), and gravitational lensing (e.g. DES Collaboration: Abbott et al. 2022). Mapping the structure of the intergalactic medium (IGM) with 21-cm tomography has great promise in advancing our understanding of the epoch of reionization and cosmic dawn (e.g. Pritchard & Loeb 2012). However, while sig-

nificant progress has been made in searching for the 21-cm power spectrum (Mertens et al. 2020; Trott et al. 2020; Abdurashidova et al. 2022) and the global signal (Bowman et al. 2018; Singh et al. 2022), there are still many challenges to overcome before IGM tomography can be achieved with the 21-cm line.

Meanwhile, the Ly α forest remains the best probe of the IGM available to date (e.g. Becker et al. 2015a; McQuinn 2016 for reviews). Recent measurements of effective optical depth indicate an ending of reionization as late as $z \sim 5.3 - 5.7$ (Becker et al. 2015b; Eilers et al. 2018; Bosman et al. 2022). The presence of long Gunn-Peterson troughs extending $\sim 10 - 100 h^{-1}$ comoving Mpc (cMpc) (Becker et al. 2015b; Zhu et al. 2022) and transmission spikes (Barnett et al. 2017; Yang et al. 2020) in the same redshift range indicates large spatial variation in the IGM opacity at the tail end of reionization. Simulations suggest the large-scale fluctuations of Ly α forest transmission could be caused by the UV background fluctuations from

* E-mail: kakiichi@ucsb.edu (KK)

galaxies (Becker et al. 2015b; D’Aloisio et al. 2018; Davies et al. 2018) or luminous rare sources such as AGN (Chardin et al. 2015, 2017; Meiksin 2020), thermal fluctuations in the IGM (D’Aloisio et al. 2015; Keating et al. 2018), islands of neutral hydrogen due to the late end of reionization (Kulkarni et al. 2019; Keating et al. 2020; Nasir & D’Aloisio 2020), and/or the spatially-varying distribution of self-shielded absorbers which modulate the mean free path of the ionizing radiation (Davies & Furlanetto 2016; D’Aloisio et al. 2018). However, without directly observing the sources of ionizing radiation, it is difficult to understand how the interplay of these various physical processes and how they shape the physical state of the IGM during the reionization process.

Establishing the direct spatial correlation between galaxy populations and Ly α forest transmission of the IGM is key to understanding how reionization proceeded. Previous ground-based surveys have mapped the distribution of galaxies both photometrically (Becker et al. 2018; Kashino et al. 2020; Christenson et al. 2021; Ishimoto et al. 2022) and spectroscopically (Kakiichi et al. 2018; Meyer et al. 2019, 2020; Bosman et al. 2020) along sightlines to luminous $z \geq 6$ quasars where high quality Ly α forest spectra are available. By using Subaru/Hyper-Suprime Cam (HSC) narrow-band (NB) imaging in a total of six quasar fields, Becker et al. (2018); Christenson et al. (2021); Ishimoto et al. (2022) find that $\sim 20 h^{-1} \text{cMpc}$ scale galaxy underdensities (overdensities) at $z \simeq 5.7$ correlate with opaque (transmissive) regions of the IGM on $\sim 50 - 100 h^{-1} \text{cMpc}$ scale. These observations favour the scenario where ionizing radiation from galaxies drives the large-scale UV background fluctuations at the tail end of reionization and/or completely neutral islands still exist in the IGM at $z < 5.7$. The spectroscopic survey using Keck/DEIMOS and VLT/MUSE (Kakiichi et al. 2018; Meyer et al. 2020) supports a similar picture. Surveying a total of eight quasar fields, Meyer et al. (2020) find a large scale excess transmission of Ly α forest around $z \simeq 5.8$ galaxies on scales of $\sim 10 - 40 h^{-1} \text{cMpc}$ at $\sim 2 - 3\sigma$. By modelling the observed galaxy-Ly α forest cross-correlation signal, they interpreted that this excess transmission is caused by the UV background fluctuations driven by the faint unseen population of galaxies clustered around luminous galaxies, with an average Lyman continuum (LyC) escape fraction of $\langle f_{\text{esc}} \rangle \simeq 14\%$ at $z \simeq 5.8$.

Recent fully-coupled cosmological radiation hydrodynamic simulations also show the excess transmission in the large-scale galaxy-Ly α forest cross-correlation and suggest that the cross-correlation signal contains important information about the timing of reionization (Garaldi et al. 2022). The large-scale UV background fluctuations in the Ly α forest have long been recognised to contain valuable information about the nature of LyC sources (e.g. host halo mass) through their clustering properties (Pontzen 2014; Gontcho A Gontcho et al. 2014; Meiksin & McQuinn 2019; Wolfson et al. 2022). In summary, establishing the spatial correlation between galaxies and Ly α forest at $5 < z < 7$ offers a smoking gun test of the reionization process and is one of the key science goals of ongoing JWST quasar field surveys (ID 2078, PI: Wang et al. 2021; ID 1243, PI: Lilly et al. 2017, see also Kashino et al. 2022). However, due to the rarity of high-redshift quasars, the connection between galaxies and the IGM in these surveys will remain limited to one-dimensional skewers.

Ultimately we seek to map both galaxies and the IGM in three dimensions. At intermediate redshifts $z \sim 2 - 3$, deep spectroscopic samples of background galaxies have enabled the construction of 3D Ly α forest tomographic maps of the IGM (Lee et al. 2014a,b, 2018; Newman et al. 2020; Horowitz et al. 2021) and the measurements of the galaxy-Ly α forest cross-correlation (Steidel et al. 2010; Chen et al. 2020). However, extending this approach to higher red-

shifts is extremely challenging because the diminishing Ly α forest transmission demands a much larger investment of telescope time. Spectroscopically detecting the UV continua and Ly α forest transmission of $z \simeq 5 - 6$ galaxies would require 30-m class telescopes such as the Thirty-Meter Telescope (TMT), Giant Magellan Telescope (GMT), and Extremely Large Telescope (ELT) (Japelj et al. 2019).

An alternative approach for IGM tomography is to utilize ultra-deep NB imaging to detect the Ly α forest transmission at a fixed redshift against the backdrop of more distant galaxies. As the throughput of an imager is much higher than a typical spectrograph, this photometric measurement of the Ly α forest transmission can be more sensitive than one based on spectroscopy. For the $\sim 60\%$ throughput of the HSC imager (cf. $\sim 10 - 20\%$ throughput of a typical spectrograph), the NB measurement of the Ly α forest with the Subaru 8.2-m can be comparable to undertaking a spectroscopic IGM survey with a 14-20 m telescope. The typical NB filter width is $\simeq 100 \text{\AA}$ which corresponds to a line-of-sight distance of $\sim 30 h^{-1} \text{cMpc}$ at $z \sim 5 - 6$. This matches the scale of fluctuations in the Ly α forest transmission seen by previous quasar field surveys (Becker et al. 2018; Kakiichi et al. 2018; Kashino et al. 2019; Meyer et al. 2019, 2020; Christenson et al. 2021; Ishimoto et al. 2022). Since up to a few hundred background galaxies can be identified using extant spectroscopic catalogues and NB-selected Ly α emitters (LAEs) in well-studied extragalactic fields, photometric IGM tomography can provide a $\times 100$ increase in the number of galaxy-Ly α sightline pairs compared to quasar surveys - a huge boost in statistical power. In an earlier article, Kakiichi et al. (2022) outlined the strategy for photometric IGM tomography which provides a path forward to map to map the Ly α forest transmission of the IGM and Ly α emitting galaxies in the same cosmic volume using only imaging data.

In this paper, we apply this photometric IGM tomography technique to the well-studied extragalactic COSMOS field and present measurements of the LAE-Ly α forest cross-correlation and the auto-correlation of the Ly α forest at $z \simeq 4.9$. The wealth of deep multi-wavelength imaging and spectroscopic data makes COSMOS an ideal field to demonstrate the method. We provide the first large-scale 2D tomographic map of the IGM at $z \simeq 4.9$ across the 1.8 deg^2 ($\sim 140 h^{-1} \text{cMpc}$ in diameter) field of view of Subaru/HSC, enabling us to directly visualise the spatial connection between galaxies and the IGM. After the reionization process is complete, we expect that the galaxy-Ly α forest cross-correlation will evolve from positive (Kakiichi et al. 2018; Meyer et al. 2020; Garaldi et al. 2019) owing to the large-scale UV background fluctuations and/or ionized bubbles to a negative (i.e. anti-correlation) signal (e.g. Turner et al. 2017; Nagamine et al. 2021) due to the increasing impact of gas overdensities around galaxies at lower redshifts. Our study at $z \simeq 4.9$ will provide a clue for how the cross-correlation evolves across cosmic time.

In Section 2 we describe the imaging data used in this analysis. Section 3 describes the galaxy catalogues and the selection for the photometric IGM tomography. Section 4 presents the method to estimate the Ly α forest transmission along background galaxies using a Bayesian spectral energy distribution (SED) fitting framework. In Sections 5-8, we present our main results including the measurements of mean Ly α forest transmission (Section 5), LAE-Ly α forest cross-correlation (Section 6), auto-correlation of the Ly α forest (Section 7), and the reconstruction of the 2D IGM tomographic map (Section 8). In Section 9 we discuss the requirement to improve the photometric IGM tomography and various science applications. We summarise our results and conclusions in Section 10. Throughput this paper we assume cosmological parameters $(\Omega_m, \Omega_\Lambda, \Omega_b, h, \sigma_8, n_s) =$

(0.3089, 0.6911, 0.0486, 0.6774, 0.8159, 0.9667) (Planck Collaboration et al. 2016). We use cMpc (pMpc) to indicate distances in comoving (proper) units. All magnitudes in this paper are quoted in the AB system (Oke & Gunn 1983).

2 DATA

We use the public release of a Subaru HSC NB718 image from CHORUS DR1 (Inoue et al. 2020) and broad-band (BB) *grizy* and NB816 images from HSC-SSP DR3 (Aihara et al. 2022) in the ultra-deep layer of COSMOS field (tract 9813). These images cover an area of approximately $1.67 \times 1.67 \text{ deg}^2$ centred at (RA, DEC) = (10h01m00s, +2d14m00s). Co-added images are retrieved from the public data release website.^{1,2} We use NB718 as a foreground NB filter to measure the Ly α forest transmission at $z \simeq 4.9$ (the central wavelength of 7170.5 Å corresponds to Ly α redshift of $z_{\text{Ly}\alpha} = 4.898$) covering the range of $4.85 < z < 4.94$ and z - and y -bands to measure the UV continua of the background galaxies.

We mask regions contaminated by artefacts (bright star haloes, ghosts, blooming, channel-stop, dip). Since we measure the residual transmitted fluxes in the NB718 image, particular care is needed for this filter as artefacts could potentially cause false positive detections. To address this we first flag all the masked pixels reported by CHORUS PDR1 (Inoue et al. 2020). This includes the pixels with the `hscPipe` flags: `pixelflags_bright_object=True` (pixels affected by bright objects) or `pixelflags_saturatedcenter=True` (pixels affected by count saturation). Pixels affected by the haloes of bright stars are also masked. After visual inspection of the NB718 image with the CHORUS PDR1 mask overlaid, we find that there are still some regions affected by the outer ghosts of bright stars, which extend to approximately ~ 270 arcsec in radius. To mask these, we conservatively follow the procedure adopted by HSC-SSP DR3 (Aihara et al. 2022). We select bright stars with $G < 10$ mag from GAIA DR3 catalogue³ in the footprint of tract 9813. We then mask regions inside 320 arcsec radius around these Gaia stars. 320 arcsec corresponds to the size of the outermost ghost identified by Aihara et al. (2022). They find that the outer ghost is significant for a star brighter than ~ 7 mag and the inner ghost is dominant at $\sim 7 - 9$ mag. We apply the same mask for the broad-band images.

The limiting magnitudes of the NB and BB images are estimated using synthetic apertures randomly distributed in the blank sky regions of the image. For NB718 we use the published limiting magnitude map from the CHORUS PDR1. For the BB images, we retrieve the synthetic apertures located in the empty regions of the sky (hereafter sky objects) from HSC database and perform photometry with a fixed 1.5'' aperture. As the sensitivity varies across the field of view, we estimate the limiting magnitudes for each patch. In each patch, there are typically ~ 40 sky objects and we compute the limiting magnitude from the standard deviation of the photometric measurements of the fluxes for the sky objects. Figure 1 shows the limiting magnitudes of NB718, z , and y -bands for each patch in the field.

The NB limiting magnitudes vary by about ~ 0.18 mag. The median depths of the NB and BB images are summarised in Tables 1 and 2.

A rule-of-thumb for the required depth for IGM tomography is given by Kakiichi et al. (2022). Assuming the flat UV continuum slope of a background galaxy, the NB magnitude needs to reach

$$m_{\text{NB}} = m_{\text{BB}} - 2.5 \log_{10} e^{-\tau_{\text{eff}}(z)} \approx m_{\text{BB}} + \tau_{\text{eff}}(z), \quad (1)$$

to detect the Ly α forest transmission with an effective optical depth $\tau_{\text{eff}}(z)$. Here, the NB magnitude corresponds to NB718 and the BB magnitude corresponds to z -band which covers the UV continuum of a background galaxy. Assuming the effective optical depth of $\tau_{\text{eff}}(z) = 1.5$ at $z = 4.9$ (Becker et al. 2013; Eilers et al. 2018; Bosman et al. 2022), the required NB718 depth for a $z = 25.0$ mag background source is 26.5 mag. The existing NB718 depth meets this requirement at $> 3\sigma$ (Table 1). For a fainter background source with $z = 26.5$ mag, the existing NB718 depth still has a $\sim 1\sigma$ sensitivity to the mean Ly α forest transmission. The HSC imaging of the COSMOS field thus has sufficient sensitivity for photometric IGM tomography.

2.1 Photometry

We measure the *grizy* and NB718 photometry from HSC-SSP DR3 and CHORUS PDR1 data using a fixed 1.5'' aperture for the background sources. The zero points of all photometric bands is 27 mag/DN according to the HSC-SSP data release. We ignore a few percent level correction arising from aperture corrections during the photometric calibration stage. This is negligible compared with the other photometric errors described below. We assign the photometric error of an object based on the limiting magnitude of the patch where the object is located.

3 CATALOGUES

The redshift range for the background sources is chosen such that the Ly α forest range between Ly β ($\lambda_{\beta} = 1026 \text{ \AA}$) and Ly α ($\lambda_{\alpha} = 1216 \text{ \AA}$) lines is covered by the NB718 filter. The lower and upper redshifts are set by $z_{\text{min}} = \lambda_{\text{NB,max}}/\lambda_{\alpha} - 1$ and $z_{\text{max}} = \lambda_{\text{NB,min}}/\lambda_{\beta} - 1$ where $\lambda_{\text{NB,max}}$, $\lambda_{\text{NB,min}}$ are the maximum and minimum wavelengths of the filter. We define the minimum and maximum wavelengths as a range where the NB718 filter transmission is $> 50\%$. The appropriate background source redshift for $z = 4.9$ IGM tomography is thus $4.98 < z < 5.89$.

To locate foreground galaxies at $z = 4.9$ in the same redshift slice corresponding to that for which our IGM transmission is being measured, we employ the SILVERRUSH catalogue of $z \simeq 4.9$ LAEs (VER20210224, Ono et al. 2021).

To identify background galaxies at $4.98 < z < 5.89$, we use both the SILVERRUSH catalogue of $z \simeq 5.7$ LAEs (Ono et al. 2021) and the spectroscopic redshift (spec- z) catalogue compiled with HSC-SSP DR3 (Aihara et al. 2022). For the latter, we find that DEIMOS10k (Hasinger et al. 2018) was the primary source of the background galaxies as we will describe below. Thus in the remainder of the paper, we refer the background galaxies derived from the catalogues to as the LAE and DEIMOS10k samples, respectively.

3.1 Spectroscopic redshift catalogue

The spec- z catalogue associated with HSC-SSP DR3 is a compilation of public spectroscopic redshifts from numerous previous redshift surveys including 2dFGRS (Colless et al. 2003), 3D-HST (Skelton et al. 2014; Momcheva et al. 2016), 6dFGRS (Jones et al. 2004,

¹ HSC-SSP DR3: https://hsc-release.mtk.nao.ac.jp/doc/index.php/data-access_pdr3/

² CHORUS DR1: <https://hsc-release.mtk.nao.ac.jp/doc/index.php/chorus/> The ancillary data including PSF FWHM and limiting magnitudes for each patch is also downloaded from here.

³ <https://gea.esac.esa.int/archive/>

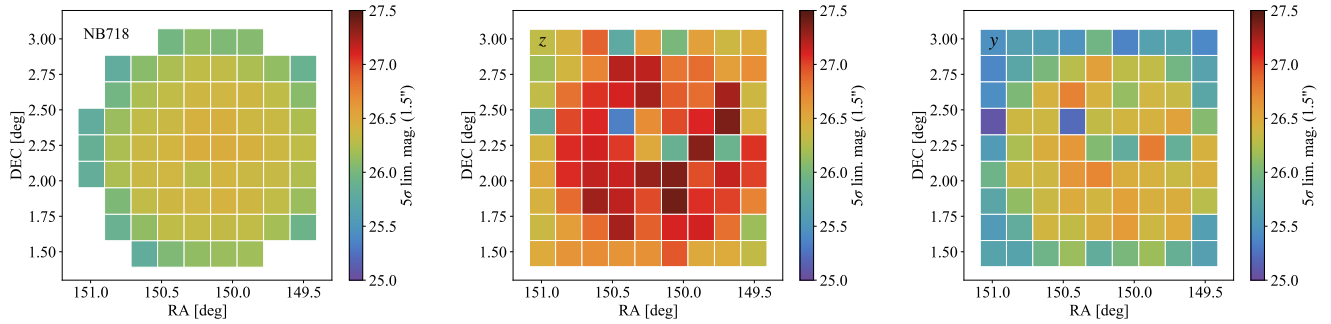


Figure 1. 5σ limiting magnitudes of the NB718, z, and y-band images with fixed $1.5''$ aperture in the COSMOS field (tract 9813).

Table 1. Summary of the foreground NB images

Filter	$\text{Ly}\alpha$ redshift	bkg. source redshift	5σ depth ($1.5''$) [AB mag]	3σ depth ($1.5''$) [AB mag]	1σ depth ($1.5''$) [AB mag]	Ref
NB718	4.90 ($4.85 < z_{\text{Ly}\alpha} < 4.94$)	$4.98 < z < 5.89$	26.30	26.86	28.05	Inoue et al (2020)

^b computed from the median of random sky objects in masked region (i.e. excluding near bright stars and artefacts) in each patch of tract 9813.

Table 2. Summary of the BB images from HSC-SSP DR3 in the UD-COSMOS field (tract 9813)

Median 5σ depth ($1.5''$) [mag]					Reference
<i>g</i>	<i>r</i>	<i>i</i>	<i>z</i>	<i>y</i>	
27.85	27.39	27.22	26.86	26.23	Aihara et al (2022)

2009), C3R2 DR2 (Masters et al. 2017, 2019), DEEP2 DR4 (Davis et al. 2003; Newman et al. 2013), DEEP3 (Cooper et al. 2011, 2012), DEIMOS10k (Hasinger et al. 2018), FMOS-COSMOS (Silverman et al. 2015; Kashino et al. 2019), GAMA DR2 (Liske et al. 2015), LEGA-C DR2 (Straatman et al. 2018), PRIMUS DR1 (Coil et al. 2011; Cool et al. 2013), SDSS DR16 (Ahumada et al. 2020), SDSS IV QSO catalog (Pâris et al. 2018), UDSz (Bradshaw et al. 2013; McLure et al. 2013), VANDELS DR1 (Pentericci et al. 2018), VIPERS PDR1 (Garilli et al. 2014), VVDS (Le Fèvre et al. 2013), WiggleZ DR1 (Drinkwater et al. 2010), and zCOSMOS DR3 (Lilly et al. 2009).

Spectroscopic sources are matched to the HSC photometry by position, thus the catalogue only includes the objects detected by HSC-SSP DR3. Using the CAS search, we retrieve the HSC-SSP DR3 spec-z catalogue after it was cross-matched using the `object_id`. We re-measure the fixed $1.5''$ aperture photometry in *grizy* and NB718 for all the objects to ensure consistent measurements of the fluxes across all bands.

Using the spec-z catalogue, we search for background source candidates at $4.98 < z < 5.89$ in the UD-COSMOS field (tract 9813). The catalogue contains a spec-z flag (`specz_flag_homogeneous=True` for secure and `False` for insecure) after homogenising the quality flags⁴ of spectroscopic redshifts from the above surveys. Selecting only objects with `specz_flag_homogeneous=True`, we find 236 candidates in the required redshift range, of which 60 belong to DEIMOS10k (Hasinger et al. 2018) and 176 belong to 3D-HST (v4.1.5, Momcheva et al. 2016). The majority of bright candidates with $z \lesssim 25.5$ comes from DEIMOS10k whereas fainter candidates are mostly from 3D-HST.

In order to visually confirm the spectroscopic redshifts, we downloaded the original DEIMOS10k spectra from NASA/IPAC Infrared Science Archive (IRSA)⁵, and the 3D-HST spectra from MAST archive⁶. For 54 of the 60 DEIMOS10k sources, we confirmed an emission line feature (mostly single $\text{Ly}\alpha$ line, one Lyman break only, one quasar). We rejected 6 objects because either (1) no published spectrum is available (DEIMOS10k ID: L234173) or (2) we could not visually confirm the reported redshift (L420065, L430951, L378903, C563716, L442206).

Out of the remaining 54 DEIMOS10k sources, we removed 11 residing in the masked regions. Furthermore, in order to secure a reliable UV continuum detection in each background sources, we applied a 5σ detection cut in the *z*-band using the limiting magnitude appropriate for the relevant patch. 3 candidates fail to meet this criterion in the *z*-band of HSC-SSP DR3. We also require a 3σ non-detection in *g*-band in order to reject low-redshift interlopers. This leads to the removal of a further 4 sources. As a result we finally have 36 DEIMOS10k background sources for our IGM tomography. We present the postage stamp image and DEIMOS spectrum of a representative background source in Figure 2.

For the 3D-HST sources, the catalogued redshifts are determined from either photometric and/or grism spectroscopic data. The ACS/G800L grism spectra cover $\text{Ly}\alpha$ in our desired redshift range. We downloaded the 3D-HST catalogue and find that all the relevant sources have only photometric redshifts; we could not find any sources with a convincing $\text{Ly}\alpha$ line or Lyman break in the grism spectra. In principle, we can use background objects with photo-*z*'s whose 95% confidence interval (i.e. `z_BEST_L95`, `z_BEST_U95` from `COSMOS_3DHST.v4.1.5.ZBEST.DAT`) lie within our required range. This would ensure that their $\text{Ly}\alpha$ forest region is appropriately covered by the NB718 filter. However, since it is unclear how catastrophic photo-*z* errors might affect the quality of our IGM tomography, we decided to remove all the 3D-HST objects from our final background sources in this paper. Nonetheless, in future work, it will be inter-

⁴ https://hsc-release.mtk.nao.ac.jp/doc/index.php/catalog-of-spectroscopic-redshifts__pdr3/

⁵ <https://irsa.ipac.caltech.edu/data/COSMOS/overview.html>

⁶ <https://archive.stsci.edu/prepds/3d-hst/>

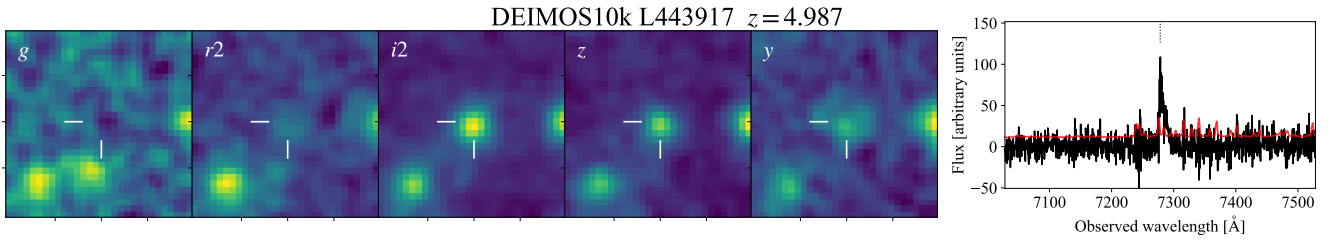


Figure 2. An example background galaxy from the DEIMOS10k catalogue. **(Left):** Postage stamps show 5×5 arcsec cutouts of *grizY* images around the object marked with a white crosshairs. **(Right):** DEIMOS 1D spectrum (flux: black, red: noise). The vertical dotted line indicates the Ly α line.

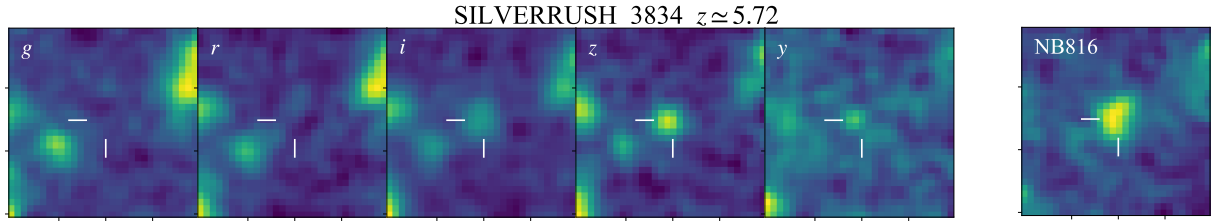


Figure 3. An example background LAE from the SILVERRUSH catalogue. **(Left):** Postage stamp showing 5×5 arcsec cutouts of *grizY* images with the object marked with a white cross. **(Right):** NB816 image. The NB colour excess is clearly detected.

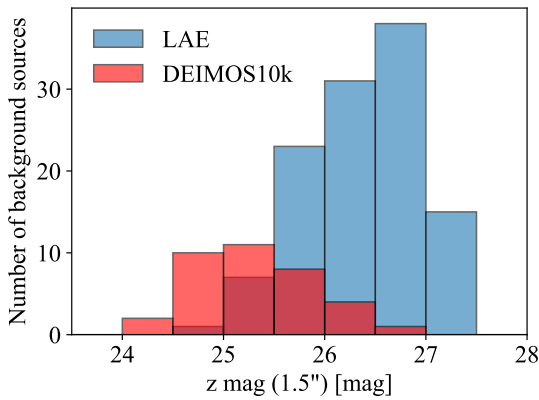


Figure 4. Magnitude distribution of background sources for NB718 IGM tomography at $z \approx 4.9$. Our final background source catalogue contains 151 sources in total (red: 36 spec- z objects from DEIMOS10k, blue: 115 LAE objects from SILVERRUSH).

esting to examine the utility of the photo- z background sources for IGM tomography.

To summarise, our final catalogue of background spec- z sources contains 36 objects from DEIMOS10k. Their z -band magnitudes and distribution in the UD-COSMOS field are shown in Figures 4 and 5, respectively. The z -band magnitudes range from 24.2 to 26.9 with the median SNR ≈ 19.5 . Because the previous DEIMOS spectroscopic campaigns are focused near the central region of the COSMOS field, their distribution reflects their survey footprints.

3.2 LAE catalogue

We use the LAE catalogue from Ono et al. (2021) constructed as part of the SILVERRUSH programme. The catalogue is based on the data from CHORUS survey (Inoue et al. 2020) and from the HSC-SSP internal data release of S18A which is basically identical to the Public Data Release 2 (Aihara et al. 2019). To ensure ho-

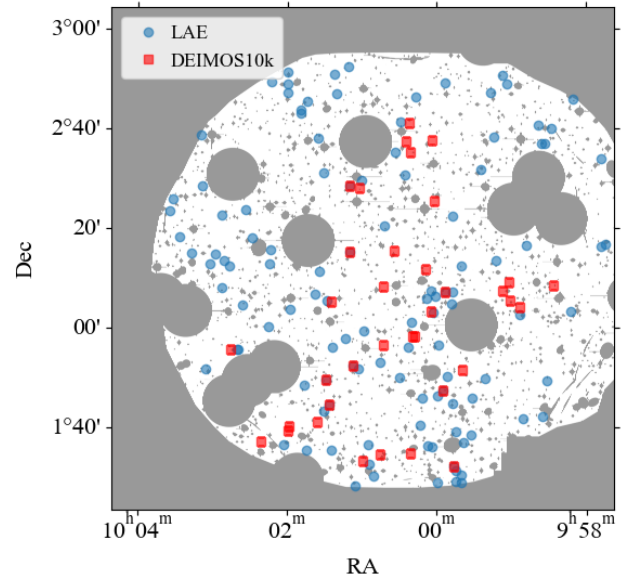


Figure 5. Sky distribution of background sources for NB718 IGM tomography (blue circles: $z = 5.7$ LAEs from SILVERRUSH, red squares: spec- z from DEIMOS 10k). Masked regions and those outside the field-of-view are indicated by the gray shaded regions.

mogeneous photometric measurements for the IGM tomography, we re-measure the fixed $1.5''$ aperture photometry at the coordinates of the SILVERRUSH LAEs using HSC-SSP DR3 *grizy* and CHORUS NB718 images. We use the SILVERRUSH catalogue to select both background and foreground LAEs at $z = 5.7$ and $z = 4.9$ located by NB816 and NB718 colour excess, respectively.

Table 3. Average physical properties of foreground LAEs.

Redshift	$\log_{10}\langle L_{\alpha} \rangle (2.0'')$ [erg s^{-1}]	$\langle M_{\text{UV}} \rangle (2.0'')$ [AB mag]
$z = 4.89$	42.62	-20.09

3.2.1 Background LAE selection: $z \approx 5.7$

To select $z \approx 5.7$ background LAEs, we draw a sample from the SILVERRUSH catalogue (Ono et al. 2021) which applies a NB colour excess $i - \text{NB816} \geq 1.2$ and 5σ detection in NB816. Such a NB816 colour excess can locate LAEs at a redshift $z = 5.726$ with an $\Delta z \approx 0.1$ accuracy (Ono et al. 2021) sufficient to ensure that the $\text{Ly}\alpha$ forest transmission is covered by the NB718 filter. Details of the LAE catalogue construction are described in Ono et al. (2021). The SILVERRUSH catalogue contains 378 $z \approx 5.7$ LAEs in the UD-COSMOS field. This double NB technique (Kakiichi et al. 2022) allows us to efficiently assemble a large number of background sources for IGM tomography.

In addition to the standard NB selection, we require a 5σ detection in z -band. This criterion is met by 176 objects out of 378 LAEs. We also removed 35 further objects which reside in the masked regions. While Ono et al. (2021) already applied masks in constructing the original catalogue, our revised mask in NB718 is more conservative. As before, we also require a 3σ non-detection in the g -band to avoid possible low-redshift interlopers; this removes a further 26 objects.

Thus, our final background LAE catalogue contains 115 objects. We visually inspected all of these objects in HSC DR3 *grizy* and NB816 images. The z -band magnitudes and the spatial distribution of the background LAEs are shown in Figures 4 and 5. The z -band magnitudes range from 24.7 to 27.2 with the median SNR ≈ 7.4 . For comparison, the i -band magnitudes of the background LAEs, which cover the $\text{Ly}\alpha$ forest flux, $\text{Ly}\alpha$ emission line, and UV continuum rewards of $\text{Ly}\alpha$ line, are much fainter than the z -band magnitudes, ranging from 28.0 to 29.7 with the median SNR ≈ 6.8 . The background LAEs are typically fainter than the DEIMOS10k sample, but distributed more evenly across the entire UD-COSMOS field as they are selected homogeneously via NB816 colour excess. We present an example postage stamp of a background LAE in Figure 3.

3.2.2 Foreground LAE selection: $z \approx 4.9$

In order to cross-correlate foreground LAEs with the $\text{Ly}\alpha$ forest transmission, we also use the HSC NB718 data to select LAEs at $z \approx 4.9$. The SILVERRUSH catalogue applies the selection criteria: $ri - \text{NB718} > 0.7$ and $r - i > 0.8$ and $ri - \text{NB718} > (ri - \text{NB718})_{3\sigma}$ and $g > g_{2\sigma}$ where ri is calculated by the linear combination of the fluxes in r - and i -bands, f_r and f_i , following $f_{ri} \approx 0.3f_r + 0.7f_i$ and the 2σ and 3σ subscripts denote 2σ and 3σ limiting magnitudes (Ono et al. 2021). Further detail is described in Ono et al. (2021). This gives 280 $z \approx 4.9$ LAEs in the UD-COSMOS field. We remove 17 objects lying in our updated masked regions. For these foreground LAEs, unlike the background galaxies, we do not apply any z -band detection cut and use all 263 NB718-selected LAEs for our subsequent analysis. The average luminosities of the foreground LAEs are summarised in Table 3.

4 IGM $\text{Ly}\alpha$ FOREST TRANSMISSION

A measurement of the IGM $\text{Ly}\alpha$ forest transmission T_{IGM} using the foreground NB718 filter requires us to infer the intrinsic spectral

energy distribution (SED) of each background galaxy in the absence of any IGM absorption. The NB-integrated $\text{Ly}\alpha$ forest transmission is then determined by the ratio between the observed and intrinsic NB fluxes,

$$T_{\text{IGM}} = \frac{f_{\text{NB}}^{\text{obs}}}{f_{\text{NB}}^{\text{intr}}}. \quad (2)$$

There are several ways to perform this measurement. One obvious way to estimate $f_{\text{NB}}^{\text{intr}}$, similar to the approach employed by Mawatari et al. (2017), is to first fit a SED to each background galaxy using broad-band photometry redward of the $\text{Ly}\alpha$ emission line $> 1216 \text{ \AA}$, and then extrapolate the continuum to the relevant rest-frame range of the $\text{Ly}\alpha$ forest between 1026 \AA and 1216 \AA covered by the foreground NB718 filter. While intuitive, it is difficult to rigorously propagate the photometric errors and systematic uncertainties of the SED modelling into the final measurement of T_{IGM} . Also, it is hard to quantify likely degeneracies between the SED parameters (e.g. UV continuum slope, or age and dust attenuation law) and the $\text{Ly}\alpha$ forest transmission.

A better way is to *simultaneously fit both the galaxy SED and the $\text{Ly}\alpha$ forest transmission of the IGM in a fully Bayesian framework*. This allows a rigorous propagation of photometric and systematic errors in the T_{IGM} estimate for each background galaxy and characterises the full posteriors including the degeneracy with the SED model parameters.

4.1 Bayesian SED fitting framework

We apply a Bayesian SED fitting framework to measure T_{IGM} . We forward model the observed photometric fluxes in narrow- and broad-band filters using realistic HSC filter transmission curves $t_{\text{NB}}(\nu)$ and $t_{\text{BB}}(\nu)$ including the CCD quantum efficiency, the transmittance of the dewar window and the Primary Focus Unit of the HSC. For our IGM tomography, NB = NB718 and BB = z, y since we use NB718 to measure the $\text{Ly}\alpha$ forest transmission and z - and y -bands to constrain the intrinsic galaxy SED.

We denote a model SED of a background galaxy by $L_{\nu}(\nu_e|\Theta)$ (in unit of $\text{erg s}^{-1} \text{ Hz}^{-1}$) where ν_e is the emitted frequency at the rest-frame of the galaxy and Θ is a set of the SED parameters. We assume a power-law SED with $L_{\nu} = L_{\nu}(1500\text{\AA})(\nu_e/\nu_{1500})^{-(2+\beta)}$ where $L_{\nu}(1500\text{\AA})$ and ν_{1500} represent the luminosity and frequency at 1500 \AA respectively, and β is the UV continuum slope. Thus our SED parameters are $\Theta = \{M_{\text{UV}}, \beta\}$. At the known redshift of the object (either from spectroscopy or NB detection of $\text{Ly}\alpha$), the observed flux is $f_{\nu}(\nu|\Theta) = (1+z)L_{\nu}[\nu_e = \nu(1+z)|\Theta]/(4\pi D_L(z)^2)$ where ν is the observed frequency and $D_L(z)$ is the luminosity distance.

As the foreground NB718 filter covers a portion of $\text{Ly}\alpha$ forest of a background galaxy, the observed NB718 flux is attenuated by $e^{-\tau_{\alpha}}$ where τ_{α} is the $\text{Ly}\alpha$ optical depth of the IGM, therefore,

$$f_{\text{NB}}(T_{\text{IGM}}, \Theta) = \frac{\int e^{-\tau_{\alpha}} f_{\nu}(\nu|\Theta) t_{\text{NB}}(\nu) d\nu}{\int t_{\text{NB}}(\nu) d\nu} \approx T_{\text{IGM}} f_{\text{NB}}^{\text{intr}}(\Theta). \quad (3)$$

We define the NB-integrated $\text{Ly}\alpha$ forest transmission of the IGM as $T_{\text{IGM}} = \int e^{-\tau_{\alpha}} t_{\text{NB}}(\nu) d\nu / \int t_{\text{NB}}(\nu) d\nu$. In the absence of the IGM, the NB718 flux is $f_{\text{NB}}^{\text{intr}}(\Theta) = \int f_{\nu}(\nu|\Theta) t_{\text{NB}}(\nu) d\nu / \int t_{\text{NB}}(\nu) d\nu$. The BB fluxes redward of $\text{Ly}\alpha$ are not affected by the IGM. They are thus modelled as

$$f_{\text{BB}}(\Theta) = \frac{\int f_{\nu}(\nu|\Theta) t_{\text{BB}}(\nu) d\nu}{\int t_{\text{BB}}(\nu) d\nu}. \quad (4)$$

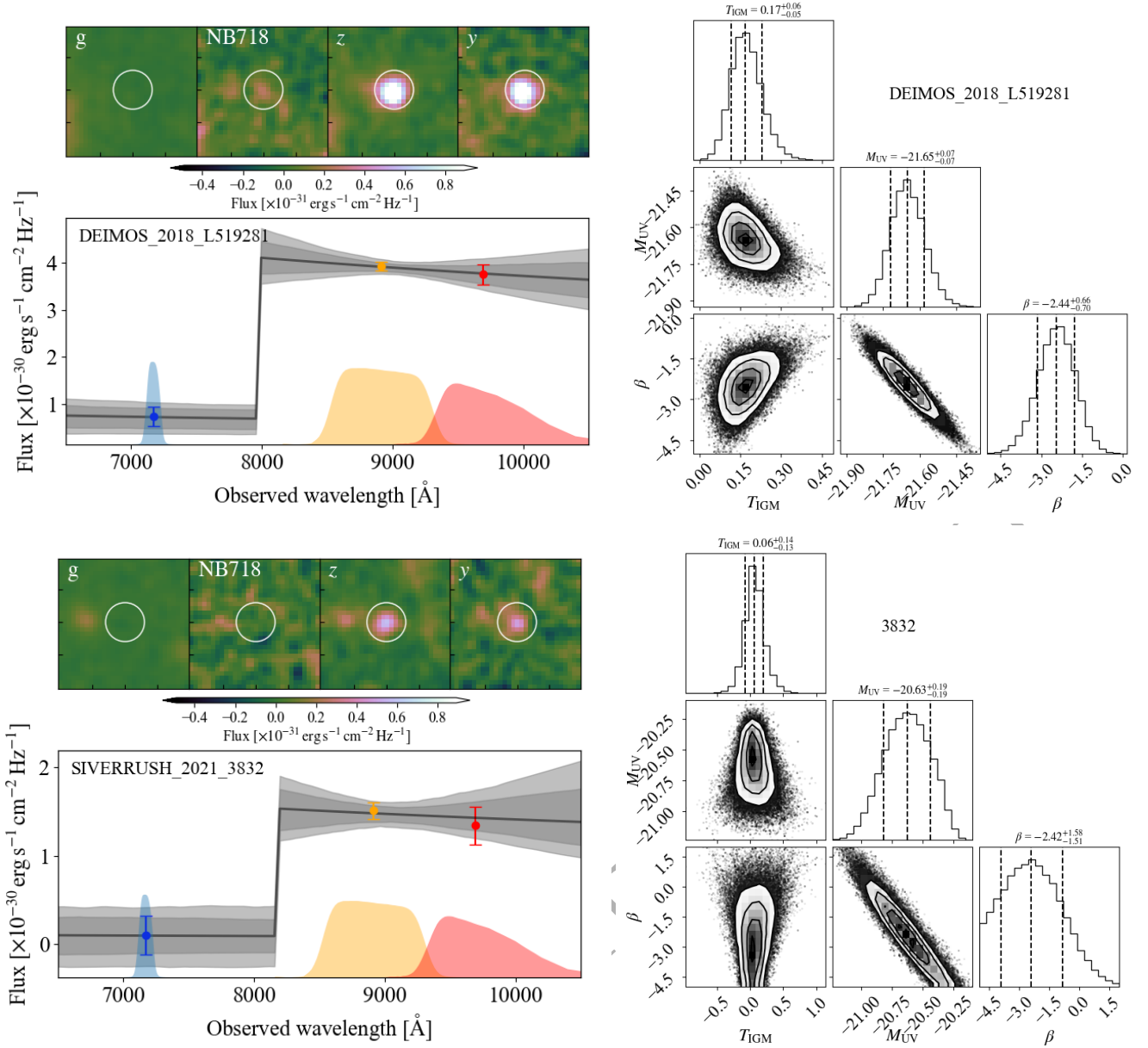


Figure 6. Representative examples illustrating results from the Bayesian SED fitting framework for the DEIMOS10k (top panels) and LAE (bottom panels) samples. The examples show a case for the detection of the transmitted Ly α forest flux in NB718 (e.g. DEIMOS_2018_519281) and one for a non-detection (e.g. SILVERRUSH_2021_7140). **(Left):** The best-fit power-law SEDs (black solid) with the 14 – 86% and 5 – 95% confidence intervals (dark and light grey shaded regions) are overlaid on the measured NB718 (blue), z (yellow), and y (red) band fluxes of the background source. The wavelength coverage of each NB718, z and y-band filter is indicated by the transparent filled curves. $5'' \times 5''$ postage stamps show the images of g, NB718, z, and y fluxes smoothed by a Gaussian kernel with a standard deviation of one pixel. The colourbar scales are the same for all images. **(Right):** MCMC corner plots of the key parameters (T_{IGM} , M_{UV} , β).

We assume the observed photometric noise follows a Gaussian distribution and that noise levels in the various filters do not correlate with one another. Therefore, the likelihood can be written as the sum of Gaussian likelihoods,

$$\ln \mathcal{L} = -\frac{1}{2} \left[\frac{f_{\text{NB}}^{\text{obs}} - f_{\text{NB}}(T_{\text{IGM}}, \Theta)}{\sigma_{\text{NB}}} \right]^2 - \frac{1}{2} \sum_{\text{BB}=z,y} \left[\frac{f_{\text{BB}}^{\text{obs}} - f_{\text{BB}}(\Theta)}{\sigma_{\text{BB}}} \right]^2. \quad (5)$$

Using Bayes theorem, we can express the posterior as a product of

prior $P(T_{\text{IGM}}, \Theta)$ and likelihood \mathcal{L} ,

$$P(T_{\text{IGM}}, \Theta | f_{\text{NB}}^{\text{obs}}, f_{\text{BB}}^{\text{obs}}) \propto P(T_{\text{IGM}}, \Theta) \mathcal{L}(f_{\text{NB}}^{\text{obs}}, f_{\text{BB}}^{\text{obs}} | T_{\text{IGM}}, \Theta). \quad (6)$$

Thus the measurement of the Ly α forest transmission T_{IGM} along each background galaxy is given by the marginalized posterior over the SED parameters Θ ,

$$P(T_{\text{IGM}} | f_{\text{NB}}^{\text{obs}}, f_{\text{BB}}^{\text{obs}}) = \int P(T_{\text{IGM}}, \Theta | f_{\text{NB}}^{\text{obs}}, f_{\text{BB}}^{\text{obs}}) d\Theta. \quad (7)$$

We implement this Bayesian SED fitting framework using a Markov Chain Monte Carlo method `emcee` (Foreman-Mackey et al.

2013). We use flat priors in the range of $-100 < T_{\text{IGM}} < 100$, $-30 < M_{\text{UV}} < -15$, and $-5 < \beta < 2$ as our default. We justify the very wide range of flat priors (instead of imposing a physical range between 0 and 1 for individual measurements for T_{IGM}) in Section 5.

When fitting the observed SEDs, we find occasional cases (13% in the background LAE sample and 0% in the DEIMOS10k sample) where the marginalized posterior peaks at an unphysically large value $T_{\text{IGM}} > 1$. This indicates there may be an unaccounted systematic error in our data. In order to assess and remove such objects, we compute the probability that the estimated T_{IGM} is greater than 1 using the posterior, i.e. $P(T_{\text{IGM}} > 1 | f_{\text{NB}}^{\text{obs}}, f_{\text{BB}}^{\text{obs}}) = \int_1^{\infty} P(T_{\text{IGM}} | f_{\text{NB}}^{\text{obs}}, f_{\text{BB}}^{\text{obs}}) dT_{\text{IGM}}$. We then flag objects with $P(T_{\text{IGM}} > 1 | f_{\text{NB}}^{\text{obs}}, f_{\text{BB}}^{\text{obs}}) > 50\%$. This choice is motivated by the fact that a Gaussian posterior centred at $T_{\text{IGM}} > 1$ gives $> 50\%$ probability that T_{IGM} is greater than 1, suggesting an unmodelled systematic error while SED fitting. We then visually check the original images and confirm such cases originate from systematic errors such as the under-subtraction of the sky background in NB718 and/or contamination from nearby objects in the photometric aperture. We apply this validation procedure⁷ in both the DEIMOS10k and LAE samples. We find no such cases (out of 36) in the DEIMOS10k catalogue. However, 11 cases (out of 115) are flagged in the LAE sample. Here, our visual inspection confirms likely contamination from nearby objects or a diffuse NB718 image compared to that in the BB z -band, suggesting that the region may be affected by the faint halo or ghost of a nearby bright star or by incorrect sky subtraction. Thus these 11 objects are removed and we use the remaining 104 background LAEs for our subsequent analysis.

4.2 Individual IGM Ly α forest transmission measurements

By applying the Bayesian SED fitting framework, we measure the Ly α forest transmission of the IGM along sightlines to individual background galaxies. The method simultaneously returns the constraints on the Ly α forest transmission T_{IGM} and the SED parameters (M_{UV} and β) for each source. In Figure 6 we show two representative examples of the derived constraints on these physical parameters for the DEIMOS10k and LAE samples. When the transmitted Ly α forest flux is detected in NB718 as demonstrated by the DEIMOS10k sample (e.g. DEIMOS_2018_L519281), we can clearly measure the Ly α forest transmission. In the case of a non-detection in NB718 (e.g. SILVERRUSH_2021_3832) the derived constraint on T_{IGM} is consistent with zero within the photometric uncertainty. The table of the MCMC results for the full DEIMOS10k and LAE samples is available as the supplementary online material.

Figure 7 shows the distribution of the estimated Ly α forest transmission. The SED fitting provides $|\delta T_{\text{IGM}}/T_{\text{IGM}}| \sim 49\%$ and 76% determinations of the Ly α forest transmission for the DEIMOS10k and LAE samples with median $\delta T_{\text{IGM}} \sim 0.10$ and 0.42 , where δT_{IGM} and T_{IGM} are given by the standard deviation and mean of the posterior. This includes errors from both photometric noise and the uncertainty in the intrinsic UV continuum. Defining the signal-to-noise ratio (SNR) to be the inverse of the relative error $\text{SNR} = |\delta T_{\text{IGM}}/T_{\text{IGM}}|^{-1}$, the median SNR of the individual T_{IGM}

⁷ Note that while one can similarly flag objects with $T_{\text{IGM}} < 0$ to avoid systematics due to over-subtracted sky background, non-detection of NB flux from faint background sources will also result in the value centred on $T_{\text{IGM}} = 0$. This means that removing these objects will unphysically bias the result to a larger mean Ly α forest transmission. We thus decided not to remove these objects from our analysis.

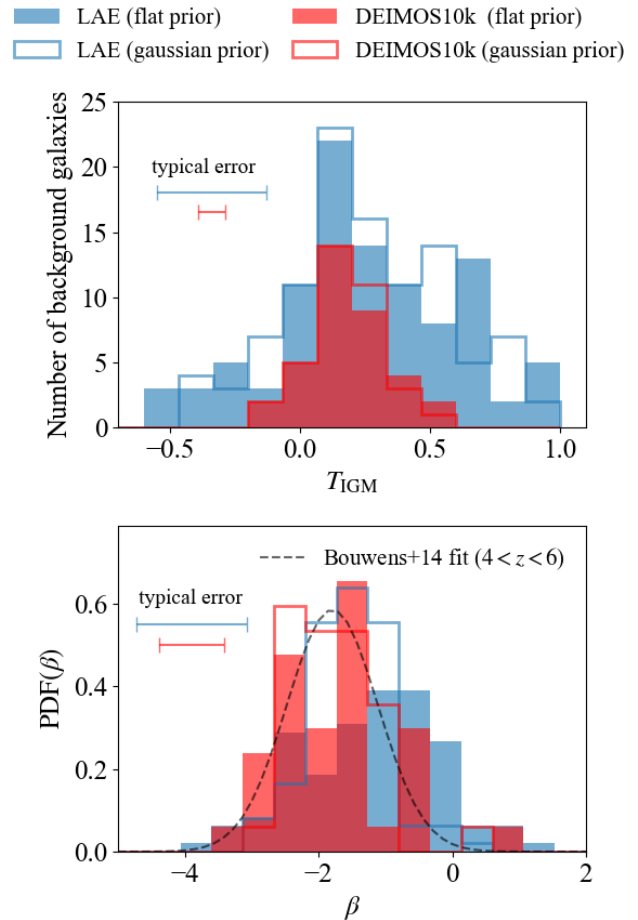


Figure 7. (Top): Distribution of estimated T_{IGM} values from the mean of each posterior of individual background galaxies (blue: LAEs, red: DEIMOS10k) using a flat prior (filled histogram) and Gaussian prior (step histogram) on the β UV slope. The median standard deviation of the posteriors is indicated as the typical error of the measurement. (Bottom): As above but for the UV slope β . The Gaussian fit to the distribution of β slopes from Bouwens et al. (2014) is indicated by the dashed line.

measurements is thus $\text{SNR} \approx 2.0$ and 1.3 for DEIMOS10k and LAE samples, respectively.

The uncertainty in T_{IGM} for the LAEs is larger than the DEIMOS10k sample. Since the LAE sample is typically fainter than the DEIMOS10k sample, it is more severely affected by photometric noise because of (i) the reduced contrast between the UV continuum level (z -band) and the Ly α forest flux (NB718) and (ii) a less precise determination of the UV continuum slope ($z - y$ colour).

The uncertainty in the UV continuum slope introduces a degeneracy in the estimate of the Ly α forest transmission. As discussed in Kakiichi et al. (2022), for power-law spectra the uncertainty in the continuum slope enters as $T_{\text{IGM}}^{\text{estimated}} = (\lambda_{\text{NB}}/\lambda_{\text{BB}})^{\beta_{\text{true}} - \beta_{\text{temp}}} T_{\text{IGM}}^{\text{true}}$ where β_{true} and β_{temp} are the continuum slopes of true and template spectra and the central wavelengths of the filters are $\lambda_{\text{NB}} = 7170.5 \text{ \AA}$ for NB718 and $\lambda_{\text{BB}} = 8912.6 \text{ \AA}$ for the z -band. The typical relative error $\delta \epsilon_{\text{cont}} \equiv |T_{\text{IGM}}^{\text{true}} - T_{\text{IGM}}^{\text{estimated}}|/T_{\text{IGM}}^{\text{true}}$ is then estimated as:

$$\langle \delta \epsilon_{\text{cont}} \rangle = \int |1 - (\lambda_{\text{NB}}/\lambda_{\text{BB}})^{\beta_{\text{true}} - \beta_{\text{temp}}}| P(\beta_{\text{true}} | \bar{\beta}, \sigma_{\beta}) d\beta_{\text{true}}, \quad (8)$$

resulting in $\langle \delta \epsilon_{\text{cont}} \rangle \approx 11\%$ (21%) error for $\beta_{\text{temp}} = -1.8$ (-1.8 ± 1) and assuming that true β slopes follows a Gaussian PDF with mean

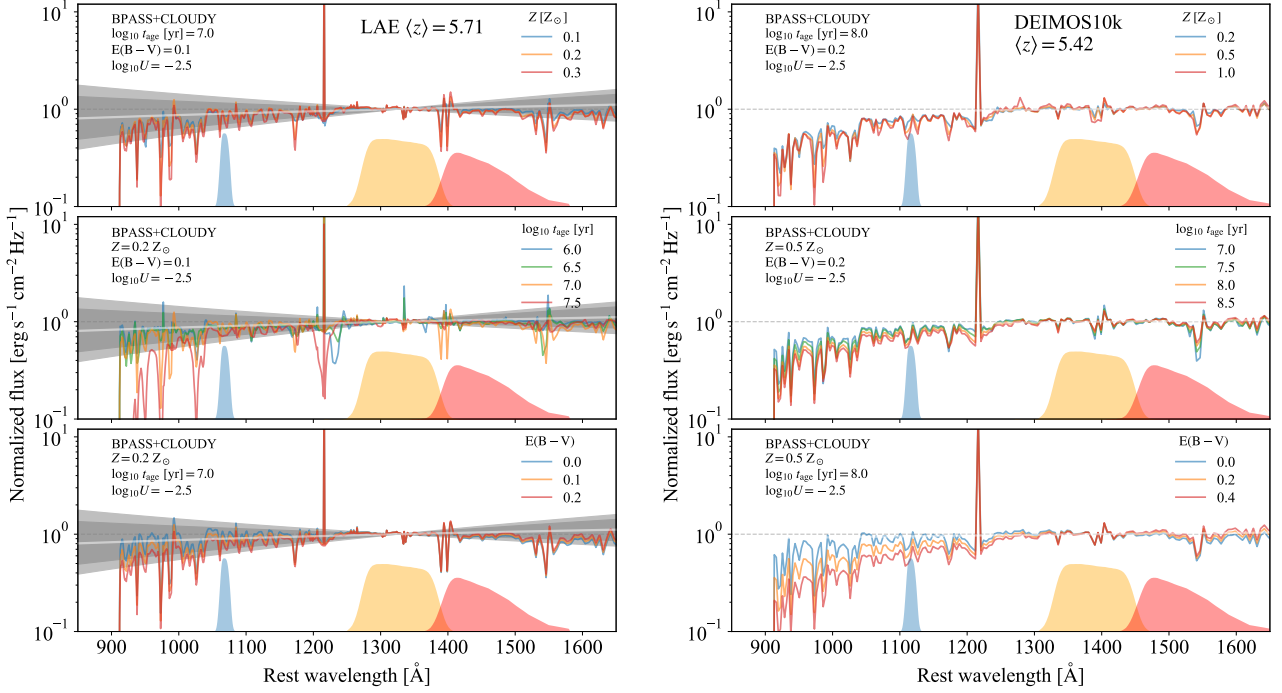


Figure 8. Comparison of the intrinsic galaxy SED without the IGM absorption from BPASS+CLOUDY and power-law fits to the LAE (left) and DEIMOS10k (right) samples. The dark (light) gray regions indicate the 16 – 84 (5 – 95) percentiles of the range of the best-fit power-law models with the white line indicating the population-averaged β slope. The default parameters of the BPASS+CLOUDY model are stellar age $t_{\text{age}} = 10$ Myr, metallicity $Z = 0.20Z_{\odot}$, ionization parameter $\log_{10} U = -2.5$ and Calzetti (2001) dust attenuation $E(B - V) = 0.1$. The filter transmission curves for NB718 (blue), z (yellow), and y (red) bands are indicated by the shaded regions.

$\bar{\beta} = -1.8$ and $\sigma_{\beta} = 0.7$ consistent with Bouwens et al. (2014). In comparison, the relative error due to the photometric noise in the Ly α forest transmission can be estimated by approximating $T_{\text{IGM}} \approx f_{\text{NB}}/f_z$ as

$$\delta\epsilon_{\text{phot}} = \sqrt{(\delta f_{\text{NB718}}/f_{\text{NB718}})^2 + (\delta f_z/f_z)^2}, \quad (9)$$

where δf_{NB718} and δf_z are the limiting fluxes in NB718 and z . The median relative error of the DEIMOS10k and LAE samples are $\sim 46\%$ and $\sim 89\%$. These estimates indicate that for our current HSC depth, the photometric noise dominates the continuum slope uncertainty.

This point is further reinforced by the fact that the choice of flat or Gaussian priors on the continuum slope has little impact on the resulting distribution of individual T_{IGM} (Figure 7). A two-sample Kolmogorov-Smirnov test indicates the difference is not statistically significant with p -values of $p = 0.392$ and 0.999 for the LAE and DEIMOS10k samples respectively.

4.2.1 Population synthesis vs power-law SEDs

We now test whether a power-law spectrum is sufficient to accurately represent the intrinsic galaxy spectrum for the purposes of IGM tomography. We compare the power-law spectrum with the intrinsic galaxy spectrum model without the Ly α forest absorption from the stellar population synthesis code BPASS (v2.2.1, Eldridge et al. 2017; Stanway & Eldridge 2018) processed with the photoionization code CLOUDY (c17.01 Ferland et al. 2017). We generate a grid of model spectra with metallicities, $Z = 0.1, 0.2, 0.3, 0.5, 1.0Z_{\odot}$ and stellar ages, $\log_{10} t_{\text{age}}/\text{yr} = 6.0, 6.5, 7.0, 7.5, 8.0, 8.5$, with a Salpeter in-

tial mass function with the upper mass limit set to $300M_{\odot}$ including binary stars. We assume an instantaneous starburst to model the LAEs and a continuous star formation history to model the Lyman-break galaxies (LBGs) in the DEIMOS10k sample. We assume the stellar population is spherically surrounded by gas with an electron density $n_e = 200 \text{ cm}^{-3}$ and ionization parameter $\log_{10} U = -2.5$ (e.g. Davies et al. 2021; Reddy et al. 2023). We then apply the Calzetti (2001) dust attenuation curve with $E(B - V) = 0.0, 0.1, 0.2, 0.4$ to the BPASS+CLOUDY outputs to model the intrinsic galaxy spectra.

In Figure 8 (left) we compare the LAE power-law spectra with a range of the best-fit UV continuum slopes with intrinsic spectra calculate from the BPASS+CLOUDY model with varying metallicities, ages, and dust extinctions. All the spectra are normalized at $\sim 1330 \text{ \AA}$ corresponding to the rest-frame wavelength coverage of the z -band. For a typical range of metallicities $Z \sim 0.1-0.3Z_{\odot}$, stellar ages $t_{\text{age}} \sim 1-30$ Myr, and dust extinction $E(B - V) \sim 0.0-0.3$ for LAEs (e.g. Ono et al. 2010; Guaita et al. 2011; Nakajima et al. 2012; Hagen et al. 2014; Trainor et al. 2016, see also reviews by Hayes 2019; Ouchi et al. 2020), the power-law template approximates the continuum shape of the BPASS+CLOUDY spectra at $\sim 1000-1600 \text{ \AA}$ very well. The relative error in the estimated T_{IGM} due to adopting different SEDs, $\delta\epsilon_{\text{SED}} \equiv |(T_{\text{IGM}}^{\text{power-law}} - T_{\text{IGM}}^{\text{BPASS+CLOUDY}})/T_{\text{IGM}}^{\text{BPASS+CLOUDY}}|$, is given by

$$\delta\epsilon_{\text{SED}} = |1 - f_{\text{NB}}^{\text{intr, BPASS+CLOUDY}}/f_{\text{NB}}^{\text{intr, power-law}}|. \quad (10)$$

Comparing the power-law template with median β slope and the BPASS+CLOUDY spectrum with typical LAE parameters of $Z = 0.20Z_{\odot}$, $\log_{10} t_{\text{age}}/\text{yr} = 7.0$, and $E(B - V) = 0.10$, the relative error is $\delta\epsilon_{\text{SED}} \sim 6\%$ (~ 15 and 17% for $E(B - V) = 0.2$ and $\log_{10} t_{\text{age}}/\text{yr} = 7.5$ respectively). This is much smaller than the error

from photometric noise and be captured by the continuum slope error budget of the power-law template. We conclude the power-law SED fit is sufficient to predict the intrinsic flux at the Ly α forest region for the current HSC depth.

We expect a similar uncertainty for the DEIMOS10k sample, which should primarily consist of LBGs. Figure 8 (right) shows the same comparison normalized at $\sim 1400 \text{ \AA}$ corresponding to the z -band coverage at the mean redshift of the DEIMOS10k sample. LBGs typically span a range of dust extinction $E(B - V) \sim 0.0 - 0.4$ (de Barros et al. 2014; Reddy et al. 2016), metallicities $Z \sim 0.2 - 1.0 Z_{\odot}$ (Steidel et al. 2014), and stellar ages $t_{\text{age}} \sim 50 - 500 \text{ Myr}$ (Stark et al. 2009; Curtis-Lake et al. 2013; de Barros et al. 2014). Being more mature star-forming galaxies than LAEs, their older stellar ages or higher dust extinctions introduce a downward trend towards shorter wavelengths. For typical LBG parameters of $Z = 0.5 Z_{\odot}$, $\log_{10} t_{\text{age}}/\text{yr} = 8.0$, and $E(B - V) = 0.2$, the impact on the estimated T_{IGM} between the BPASS+CLOUDY and power-law SEDs is $\delta\epsilon_{\text{SED}} \sim 27\%$ (~ 39 and 31% for $E(B - V) = 0.4$ and $\log_{10} t_{\text{age}}/\text{yr} = 8.5$ respectively). While larger than for the LAEs, this is still within the photometric error for the individual measurements of T_{IGM} . However, if this range of stellar ages and dust extinctions is representative of the true intrinsic spectra of background spec- z LBGs, it could introduce a systematic bias in stacked measurements. The use of power-law SEDs for background spec- z sample could systematically underestimate the measured Ly α forest transmission because it predicts an intrinsic UV continuum level larger than the true value. To quantify and eliminate the possible impact of this limitation, we would need to extend our analysis (which currently only uses z - and y -bands) including near-infrared data to better constrain the ages and dust extinctions of background galaxies. We discuss this strategy further in Section 9, but in the following analysis we only discuss the effect of this potential bias.

Another possible systematic arising from the assumption of a power-law spectrum is the presence of stellar photospheric and interstellar absorption lines in the Ly α forest region of a background galaxy. Prominent absorption lines between Ly β and Ly α lines ($1026 - 1216 \text{ \AA}$) include C II $\lambda 1036$, S IV/Fe II $\lambda 1063$, N II $\lambda 1084$, N I $\lambda 1134$, C III $\lambda 1176$, Si II $\lambda 1190, 1193$, Si III $\lambda 1207$ (Reddy et al. 2016). Previous spectroscopic IGM tomographic surveys mitigated this issue by masking the $\pm 2 - 5 \text{ \AA}$ regions around each absorption line (Lee et al. 2014b, 2018; Newman et al. 2020). For $z \sim 5.7$ background LAEs, the NB718 filter covers the rest-frame wavelength between 1050 and 1082 \AA , which coincides with the S IV/Fe II $\lambda 1063$ line. According to a stacked galaxy spectrum (Newman et al. 2020), the typical rest-frame equivalent width of the absorption line is $\text{EW}_{\text{abs}} \sim 0.8 \text{ \AA}$. The effect of an absorption line on the NB718 flux is thus $(1 + z)\text{EW}_{\text{abs}}/\Delta\lambda_{\text{NB718}} \approx 4.8\%$ reduction in flux integrated over the NB718 filter. Accordingly, the resulting bias in the estimated T_{IGM} is only

$$\delta\epsilon_{\text{abs}} = (1 + z)\text{EW}_{\text{abs}}/\Delta\lambda_{\text{NB718}}, \quad (11)$$

i.e. 4.8% , which is negligible compared to the other sources of error discussed above. For the DEIMOS10k sample, the spectroscopic redshifts span the range of $4.98 < z < 5.89$. As the contamination in NB718 flux by the absorption lines is expected to be randomized, the effect of absorption lines on the overall DEIMOS10k sample should be negligible.

5 MEAN LY α FOREST TRANSMISSION

5.1 Estimating mean Ly α forest transmission

The mean Ly α forest transmission is simply the mean in a representative volume V , $\bar{T}_{\text{IGM}} = V^{-1} \int T_{\text{IGM}}(\mathbf{x}) dV$ where \mathbf{x} is a 3D spatial position in the Universe. For IGM tomography, we can only sample $T_{\text{IGM}}(\mathbf{x})$ along sightlines to background galaxies. For each background galaxy at a location θ_i , we sample $T_{\text{IGM}}(\theta_i) = T_{\text{IGM},i}$ where $T_{\text{IGM},i}$ is the measured Ly α forest transmission integrated along the line of sight to an i th galaxy over the width of the NB filter. Since the background galaxies are uncorrelated with the foreground IGM structure, providing random sampling of $T_{\text{IGM}}(\mathbf{x})$, this is equivalent to performing the Monte Carlo integration of the mean Ly α forest transmission, i.e.

$$\bar{T}_{\text{IGM}} = \frac{1}{N_{\text{bg}}} \sum_{i=1}^{N_{\text{bg}}} T_{\text{IGM},i}. \quad (12)$$

Since we have noisy measurements of $T_{\text{IGM},i}$ characterised by the posterior $P(T_{\text{IGM},i} | f_{\text{NB},i}^{\text{obs}}, f_{\text{BB},i}^{\text{obs}})$ (equation 7) for a set of background galaxies $i = 1, \dots, N_{\text{bg}}$, our estimate of the mean Ly α forest transmission is also noisy. Thus, to characterise the uncertainty in the mean Ly α forest transmission, we need to know the posterior probability of \bar{T}_{IGM} , that is, $P(\bar{T}_{\text{IGM}} | \{f_{\text{NB},i}^{\text{obs}}, f_{\text{BB},i}^{\text{obs}}\}_{i=1, \dots, N_{\text{bg}}})$. Assuming individual measurements of $T_{\text{IGM},i}$ are statistically independent, the joint posterior probability is simply the product of all the individual posteriors, i.e. $P(\{T_{\text{IGM},i}\}_{i=1, \dots, N_{\text{bg}}} | \{f_{\text{NB},i}^{\text{obs}}, f_{\text{BB},i}^{\text{obs}}\}_{i=1, \dots, N_{\text{bg}}}) = \prod_{i=1}^{N_{\text{bg}}} P(T_{\text{IGM},i} | f_{\text{NB},i}^{\text{obs}}, f_{\text{BB},i}^{\text{obs}})$. Formally, the posterior probability of the mean Ly α forest transmission can then be written in terms of the convolution of the individual posteriors of $T_{\text{IGM},i}$ of all background galaxies (e.g. Sivia & Skilling 2006, Section 3.6),

$$P(\bar{T}_{\text{IGM}} | \{f_{\text{NB},i}^{\text{obs}}, f_{\text{BB},i}^{\text{obs}}\}_{i=1, \dots, N_{\text{bg}}}) = \int \prod_{i=1}^{N_{\text{bg}}} dT_{\text{IGM},i} P(T_{\text{IGM},i} | f_{\text{NB},i}^{\text{obs}}, f_{\text{BB},i}^{\text{obs}}) \delta_{\text{D}} \left(\bar{T}_{\text{IGM}} - \frac{1}{N_{\text{bg}}} \sum_{i=1}^{N_{\text{bg}}} T_{\text{IGM},i} \right), \quad (13)$$

where $\delta_{\text{D}}(x)$ is the Dirac delta function. Numerically, it is easy to generate the posterior of \bar{T}_{IGM} , i.e. $P(\bar{T}_{\text{IGM}} | \{f_{\text{NB},i}^{\text{obs}}, f_{\text{BB},i}^{\text{obs}}\}_{i=1, \dots, N_{\text{bg}}})$, by computing the histogram of many \bar{T}_{IGM} 's using random draws of $T_{\text{IGM},i}$ from $P(T_{\text{IGM},i} | f_{\text{NB},i}^{\text{obs}}, f_{\text{BB},i}^{\text{obs}})$.

Equation (13) illustrates an important, but subtle point on the choice of prior on $T_{\text{IGM},i}$. When the data has no constraining power, i.e. $P(T_{\text{IGM},i} | f_{\text{NB},i}^{\text{obs}}, f_{\text{BB},i}^{\text{obs}}) \rightarrow P(T_{\text{IGM},i})$, the posterior of \bar{T}_{IGM} becomes a convolution of multiple priors of $T_{\text{IGM},i}$. Assuming any prior with mean \bar{T}_{prior} and standard deviation σ_{prior} for all the individual measurements, by the virtue of central limit theorem, the posterior of \bar{T}_{IGM} approaches a Gaussian distribution with mean \bar{T}_{prior} and standard deviation $\sigma_{\bar{T}_{\text{IGM}}} = \sigma_{\text{prior}}/\sqrt{N_{\text{bg}}}$. In our case, if we were to impose a flat prior between 0 and 1 (i.e. $\bar{T}_{\text{prior}} = 0.5$ and $\sigma_{\text{prior}} = 0.29$) for the 104 individual measurements, the end result would be a Gaussian posterior on \bar{T}_{IGM} with mean $\bar{T}_{\text{prior}} = 0.5$ and $\sigma_{\bar{T}_{\text{IGM}}} = 0.29/\sqrt{104} = 0.028$ even for a completely uninformative dataset. This demonstrates that a reasonable prior on individual measurements propagates into an unreasonably tight constraint on the mean Ly α forest transmission. While imposing a flat prior with $0 < T_{\text{IGM}} < 1$ seems an innocent assumption, when we are interested in the mean we should use a maximally non-informative prior on the individual measurement to avoid an artificial constraint on the final

estimate of \bar{T}_{IGM} . We thus use a flat prior on individual measurement allowing a very wide range between $-100 < T_{\text{IGM}} < 100$.

Given the full posterior probability of \bar{T}_{IGM} , it is natural to take our best estimate of the mean Ly α forest transmission as the expectation value $\langle \bar{T}_{\text{IGM}} \rangle = \int \bar{T}_{\text{IGM}} P(\bar{T}_{\text{IGM}} | \{f_{\text{NB},i}^{\text{obs}}, f_{\text{BB},i}^{\text{obs}}\}_{i=1,\dots,N_{\text{bg}}}) d\bar{T}_{\text{IGM}}$, which simplifies as the average of the expectation values of individual measurements of $T_{\text{IGM},i}$ (see Appendix A for derivation),

$$\langle \bar{T}_{\text{IGM}} \rangle = \frac{1}{N_{\text{bg}}} \sum_{i=1}^{N_{\text{bg}}} \langle T_{\text{IGM},i} \rangle, \quad (14)$$

where $\langle T_{\text{IGM},i} \rangle = \int T_{\text{IGM},i} P(T_{\text{IGM},i} | f_{\text{NB},i}^{\text{obs}}, f_{\text{BB},i}^{\text{obs}}) dT_{\text{IGM},i}$. Similarly, the variance of the mean Ly α forest transmission, $\sigma_{\bar{T}_{\text{IGM}}}^2 = \int (\bar{T}_{\text{IGM}} - \langle \bar{T}_{\text{IGM}} \rangle)^2 P(\bar{T}_{\text{IGM}} | \{f_{\text{NB},i}^{\text{obs}}, f_{\text{BB},i}^{\text{obs}}\}_{i=1,\dots,N_{\text{bg}}}) d\bar{T}_{\text{IGM}}$, is given by

$$\sigma_{\bar{T}_{\text{IGM}}}^2 = \frac{1}{N_{\text{bg}}^2} \sum_{i=1}^{N_{\text{bg}}} \text{Var}[T_{\text{IGM},i}]. \quad (15)$$

This again simply follows from the sum of the variances of individual measurements, $\text{Var}[T_{\text{IGM},i}] = \int (T_{\text{IGM},i} - \langle T_{\text{IGM},i} \rangle)^2 P(T_{\text{IGM},i} | f_{\text{NB},i}^{\text{obs}}, f_{\text{BB},i}^{\text{obs}}) dT_{\text{IGM},i}$. Note that because the final variance $\sigma_{\bar{T}_{\text{IGM}}}^2$ scales like $1/N_{\text{bg}}$ times the average variance of the individual measurement, the final error scales as $\sigma_{\bar{T}_{\text{IGM}}} \propto 1/\sqrt{N_{\text{bg}}}$. All these quantities are easy to compute using the MCMC sample of $T_{\text{IGM},i}$ of the individual posteriors from the Bayesian SED fitting framework.

The central limit theorem guarantees that for a large number of background galaxies $P(\bar{T}_{\text{IGM}} | \{f_{\text{NB},i}^{\text{obs}}, f_{\text{BB},i}^{\text{obs}}\}_{i=1,\dots,N_{\text{bg}}})$ approaches a Gaussian distribution. Thus the expectation value $\langle \bar{T}_{\text{IGM}} \rangle$ is equivalent to the maximum *a posteriori* estimate of the mean Ly α forest transmission, $\bar{T}_{\text{IGM}}^{\text{MP}}$, which is given by

$$\bar{T}_{\text{IGM}}^{\text{MP}} = \arg \max_{\bar{T}_{\text{IGM}}} P(\bar{T}_{\text{IGM}} | \{f_{\text{NB},i}^{\text{obs}}, f_{\text{BB},i}^{\text{obs}}\}_{i=1,\dots,N_{\text{bg}}}). \quad (16)$$

Figure 9 shows the full posterior of \bar{T}_{IGM} directly computed from random realizations of individual measurements, which explicitly confirms the equivalence between the expectation value and the maximum *a posteriori* estimate of \bar{T}_{IGM} , i.e. $\bar{T}_{\text{IGM}}^{\text{MP}} = \langle \bar{T}_{\text{IGM}} \rangle$. As the full posterior is Gaussian, the variance (equation 15) completely characterises the total error in the estimated mean Ly α forest transmission as $\langle \bar{T}_{\text{IGM}} \rangle \pm \sigma_{\bar{T}_{\text{IGM}}}$ at one sigma level. This error is fully propagated, including both photometric noise and UV continuum uncertainties, from the Bayesian SED fitting procedure for the individual measurements of $T_{\text{IGM},i}$.

5.2 Results

Using the DEIMOS10k and LAE samples, the mean Ly α forest transmission at $z \approx 4.9$ are estimated to be

$$\langle \bar{T}_{\text{IGM}} \rangle = 0.179 \pm 0.028 \quad (\text{for DEIMOS10k}) \quad (17)$$

and

$$\langle \bar{T}_{\text{IGM}} \rangle = 0.293 \pm 0.045 \quad (\text{for LAE}) \quad (18)$$

This represents the first $\sim 15\%$ photometric measurement of the mean Ly α forest transmission of the IGM using background galaxies. The statistical precision is comparable to the $\sim 8\%$ determinations based using quasar spectra (Becker et al. 2013, 2015b; Eilers

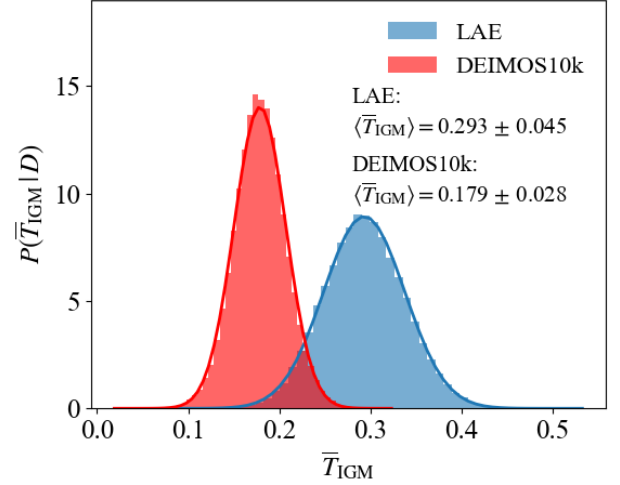


Figure 9. Full posterior of the mean Ly α forest transmission \bar{T}_{IGM} , $P(\bar{T}_{\text{IGM}} | \{f_{\text{NB},i}^{\text{obs}}, f_{\text{BB},i}^{\text{obs}}\}_{i=1,\dots,N_{\text{bg}}})$, for the LAE (blue histogram) and DEIMOS10k (red histogram) samples. The histogram is computed using the mean Ly α forest transmission from 1,000,000 random realizations of a set of $T_{\text{IGM},i}$ from the individual posteriors. The blue and red solid curves are the Gaussian distributions with the expectation value and variance computed via (14) and (15). Note that both the expectation value and the maximum *a posteriori* estimate of \bar{T}_{IGM} are equivalent.

et al. 2018; Yang et al. 2020; Bosman et al. 2018, 2022). Our precision arises from a large number of background galaxies roughly consistent with the $\propto 1/\sqrt{N_{\text{bg}}}$ scaling law, from which we expect $\delta T_{\text{IGM}}/T_{\text{IGM}} \sim 0.49/\sqrt{36} (0.76/\sqrt{104}) = 8\% (7\%)$ error on the mean transmission for the DEIMOS10k (LAE) sample. Note that the error budget quoted above using equation (15) takes into account only the photometric noise and UV continuum uncertainties, but not the error from cosmic or patch-to-patch variance. To quantify its impact, we empirically estimate the total error using Jackknife resampling (described in Section 6). We find that the Jackknife errors are $\sigma_{\text{JK}} = 0.024$ and 0.047 for the DEIMOS10k and LAE samples respectively, comparable to our analytic estimate of the (photometric noise + continuum) error. Thus, the additional error from cosmic variance is not significant.

To visually confirm that the measured Ly α forest transmission is truly representative of the physical value at $z \approx 4.9$, we applied a sigma-clipped mean stacking of the NB718 and BB images of the DEIMOS10k and LAE sample in Figure 10. The signal is clearly detected in NB718 as well as via a clear UV continuum detection in z . The signal is also detected in the median stack. The non-detection in mean g stack reinforces that the detected signal originates from the Ly α forest transmission towards the background galaxies and not due to low-redshift interlopers.

5.2.1 Comparison with the literature

In Figure 11 we compare our mean Ly α forest transmission with the measurements using quasars (Becker et al. 2013; Eilers et al. 2018; Bosman et al. 2022) and galaxy spectra (Thomas et al. 2017, 2020, 2021) in the literature. Using high signal-to-noise quasar spectra, the former measures the mean Ly α forest transmission in bins of $50 h^{-1} \text{cMpc}$ length. This is comparable to the line-of-sight comoving length $34 h^{-1} \text{cMpc}$ of the NB718 filter. Thomas et al. (2017, 2020, 2021) used a large spectroscopic sample of galaxies from the VAN-

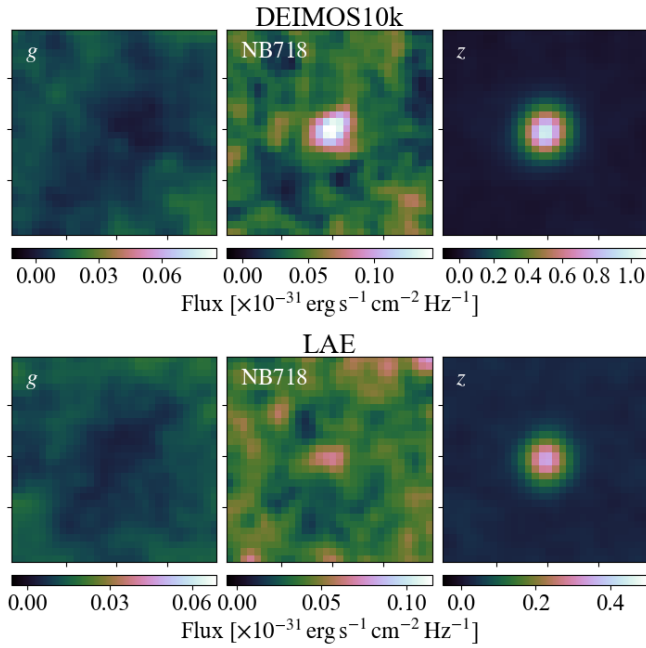


Figure 10. $5'' \times 5''$ cutout images of the sigma-clipped mean stack of g (left), NB718 (middle), z (right) images for the DEIMOS10k (top) and LAE (bottom) samples.

DELS and VUDS surveys and measured the Ly α forest transmission from the rest-frame 1070 – 1170 Å region of background galaxies by using various IGM templates in their spectral fitting method (see also [Monzon et al. \(2020\)](#) who used a stacking method).

Our measurement using the DEIMOS10k sample is in excellent agreement with the quasar studies, in particular with the latest high signal-to-noise measurement of $\langle \bar{T}_{\text{IGM}} \rangle = 0.171 \pm 0.014$ based on the XQR-30 quasar sample ([Bosman et al. 2022](#)). Although our measurement using background LAEs is slightly higher, it is still broadly in agreement with the previous values from the quasar- and galaxy-based measurements.

Our conclusion disagrees with the claim by [Thomas et al. \(2020\)](#) who argued that photometric data is insufficient to constrain the Ly α forest transmission. This is because their SED fitting used broad-band photometry to determine the Ly α forest transmission which is contaminated by the Ly α emission line and UV continuum and covers a region below the Ly β line depending on the redshift of a background galaxy. This renders the resulting measurement of the Ly α forest transmission uncertain and is a likely source of their $\sim 20\%$ discrepancy between their photometric and spectroscopic measurements. In contrast, our method uses a NB filter precisely covering the appropriate Ly α forest region enabling a clean photometric measurement of the transmitted Ly α forest flux. We argue that there is no fundamental limitation to the photometric approach when a carefully chosen combination of a NB filter and background galaxy redshifts is used.

5.3 Systematics: low-redshift interlopers and SEDs

There is a $\sim 2\sigma$ tension in our measurements of the mean Ly α forest transmission between the DEIMOS10k and LAE samples. The statistical significance is calculated from the difference between the posterior means divided by the quadrature sum of the errors ([Lemos et al. 2021](#)), $(0.293 - 0.179) / \sqrt{0.045^2 + 0.028^2} = 2.15\sigma$. While the

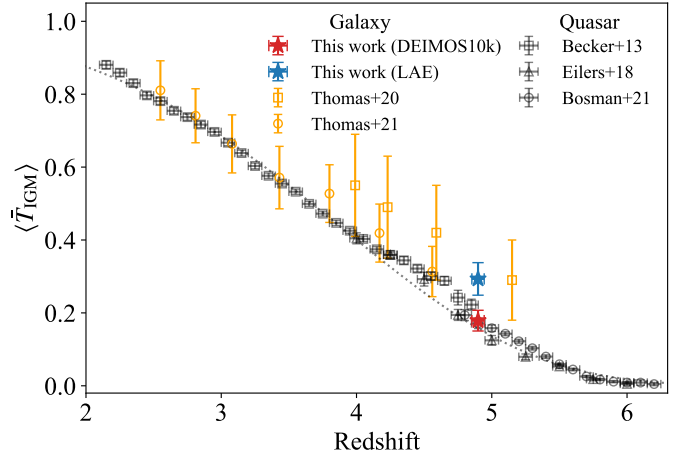


Figure 11. Comparison of the mean Ly α forest transmission measured in this work using the DEIMOS10k (red) and LAE (blue) samples with the previous measurements based on quasar ([Becker et al. 2013](#); [Eilers et al. 2018](#); [Bosman et al. 2022](#)) and galaxy spectra ([Thomas et al. 2020, 2021](#)).

discrepancy is statistically insignificant, it could indicate systematic effects outside our statistical error budget.

5.3.1 Low-redshift interlopers

The first possibility is contamination by the low-redshift interlopers in the background $z \approx 5.7$ LAE sample. A low-redshift interloper could bias the result by introducing a fictitious transmissive sightline. Possible interlopers in the LAE sample selected by a NB816 excess include low-redshift galaxies with strong emission lines from $z \approx 0.25$ H α , $z \approx 0.63$ [O III] $\lambda 5008$, and $z \approx 1.19$ [O II] $\lambda \lambda 3727, 3729$, as well as slightly lower redshift AGN at $z \approx 4.28$ with C IV $\lambda 1549$ emission (e.g. [Ouchi et al. 2008](#); [Shibuya et al. 2018](#); [Sobral et al. 2018](#)). Their rest-frame optical or UV continua could be mistaken in the NB718 filter as the transmitted Ly α forest flux at $z \approx 4.9$, artificially increasing the estimated mean Ly α forest transmission. The effect of a low-redshift interloper can be written as

$$\langle \bar{T}_{\text{IGM}} \rangle = (1 - f_{\text{bg.int}}) \langle \bar{T}_{\text{IGM}} \rangle^{\text{true}} + f_{\text{bg.int}} \langle \bar{T}_{\text{IGM}} \rangle^{\text{bg.int}}, \quad (19)$$

where $f_{\text{bg.int}}$ is the contamination rate by low-redshift interlopers in the background LAE sample and $\langle \bar{T}_{\text{IGM}} \rangle^{\text{bg.int}}$ is the fictitious mean Ly α forest transmission along the sightlines of the interlopers. The interloper fraction in the LAE selection is typically $\sim 20\%$ ([Shibuya et al. 2018](#)). Assuming that $\langle \bar{T}_{\text{IGM}} \rangle^{\text{bg.int}} = 0.7$ and the measured mean Ly α forest transmission from DEIMOS10k sample is the true value $\langle \bar{T}_{\text{IGM}} \rangle^{\text{true}} = 0.179$, the observed value using NB-selected background LAEs will become $\langle \bar{T}_{\text{IGM}} \rangle = (1 - 0.2) \times 0.179 + 0.2 \times 0.7 = 0.283$, being consistent with the measurement from our LAE sample. This would resolve the tension between DEIMOS10k and LAE samples.

In order to explore this further, we cross-matched our background LAE catalogue with the spectroscopic catalogue compiled with HSC-SSP DR3 ([Aihara et al. 2022](#)) containing 70,358 objects in the UD-COSMOS field (tract 9813). We find no interloper in our background LAE catalogue while 10 objects are spectroscopically-confirmed to be at $z \approx 5.7$. We also applied a stricter g -band non-detection cut compared to our default 3σ threshold and find that for a 1σ (2σ) threshold the estimated mean Ly α forest transmission becomes $\langle \bar{T}_{\text{IGM}} \rangle = 0.277 \pm 0.055$ (0.298 ± 0.048) for LAE sample; the values

are consistent with our result from the 3σ threshold within 1σ error. Thus we find no obvious evidence for low-redshift interlopers in our LAE sample.

Note that as we require $> 5\sigma$ detection in the z -band (< 26.86 mag), potential interlopers, if any, need to have red $g - z \gtrsim 2.0$ colours and be fainter than $29.60 - 28.84$ mag ($1 - 2\sigma$) in g -band. Such interlopers could be low-redshift dusty red galaxies or Balmer break galaxies with $H\alpha$, [O III], or [O II] doublet emission lines at $z \approx 0.25, 0.63, 1.19$. Ultimately, spectroscopic follow-up of the background LAEs would be necessary to fully reject systematic bias from low-redshift interlopers. Following up a random subset would determine the interloper fraction $f_{\text{bg.int}}$ and by measuring the mean Ly α forest transmission along the interlopers, one can also determine $\langle \bar{T}_{\text{IGM}} \rangle^{\text{bg.int}}$. Then the observed Ly α forest transmission using the parent photometric background LAE sample, $\langle \bar{T}_{\text{IGM}} \rangle^{\text{obs}}$, could be statistically corrected via

$$\langle \bar{T}_{\text{IGM}} \rangle^{\text{corrected}} = \frac{\langle \bar{T}_{\text{IGM}} \rangle^{\text{obs}} - f_{\text{bg.int}} \langle \bar{T}_{\text{IGM}} \rangle^{\text{bg.int}}}{1 - f_{\text{bg.int}}}. \quad (20)$$

5.3.2 SED templates

As discussed in Section 4.2.1, model galaxy SEDs may introduce a systematic error in T_{IGM} . The effect is expected to be larger in the DEIMOS10k sample which contains more mature galaxies with dust or more complex stellar populations for which the UV continuum cannot precisely be determined by only z and y -band photometry. If the intrinsic continua were systematically overestimated by $\sim 30\%$ (for example, because of an underestimated $E(B - V)$), then a more flexible galaxy SED template would lead to $\langle \bar{T}_{\text{IGM}} \rangle = 0.233 \pm 0.036$ relaxing the tension between the DEIMOS10k and LAE samples to $\sim 1\sigma$. However, this explanation would weaken the agreement between both our measures and those determined using quasars.

Our choice of the Ly α forest wavelength range ($1026 - 1216 \text{ \AA}$) is more generous compared to the $1040 - 1190 \text{ \AA}$ range commonly used for spectroscopic IGM tomography (Lee et al. 2014b, 2018; Newman et al. 2020). Although the effect of absorption lines is small (Section 4.2.1), the wider range means that a power-law continuum might neglect the effect of Ly $\beta + \text{C II } \lambda 1036$ and Ly $\alpha + \text{Si III } \lambda 1207$ absorption lines. Also, neutral gas in the circumgalactic medium of a background galaxy could contribute to the Ly α absorption blueward of the line centre (Rudie et al. 2013; Kakiichi et al. 2018; Bassett et al. 2021). We can test this effect by adopting a more restrictive range of $1040 - 1190 \text{ \AA}$ and limiting the redshift range of background galaxies to $5.072 < z < 5.841$ for the DEIMOS10k sample. We find $\langle \bar{T}_{\text{IGM}} \rangle = 0.174 \pm 0.035$, consistent with our main result. Thus contamination by absorption lines is negligible and cannot explain the tension.

5.3.3 Cosmic variance

Finally, cosmic variance may cause the tension between our two subsamples. Although the Jackknife resampling error should include the effects of cosmic variance, the small sample size may underestimate the effect. However, since both the DEIMOS10k and LAE samples probe similar regions of the sky in the COSMOS field, we consider cosmic variance an unlikely source of the tension.

In conclusion, we believe that the combination of a modest contribution from low-redshift interlopers and some variation in the SED templates is the mostly likely source of the 2σ tension between the DEIMOS10k and LAE samples. This can only be tested and corrected

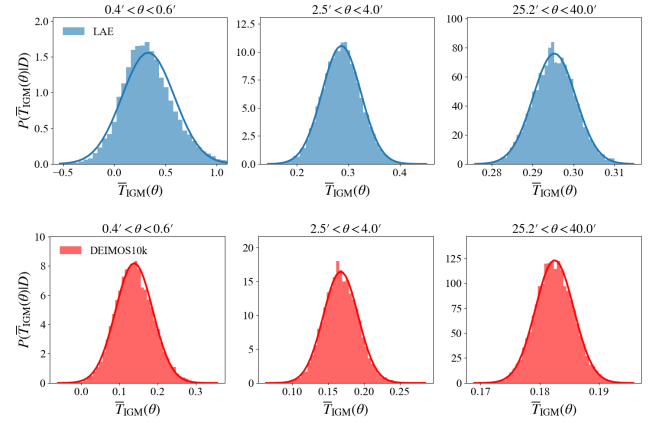


Figure 12. The full posterior of the mean Ly α forest transmission around $z \approx 4.9$ LAEs, $P(\bar{T}_{\text{IGM}}(\theta) | \{f_{\text{NB},i}^{\text{obs}}, f_{\text{BB},i}^{\text{obs}}\}_{i=1,\dots,N_{\text{bg}}})$, using background $z \approx 5.7$ LAE (blue histogram) and DEIMOS10k (red histogram) samples. The histogram is computed via $\bar{T}_{\text{IGM}}(\theta)$ from 10,000 random realizations of a set of $T_{\text{IGM},i}$ from the individual posteriors. The innermost (left), middle (middle), outermost (right) angular bins are shown. The blue and red solid curves are the Gaussian distributions with the expectation value and variance computed via equations (22) and (23). The equivalence between the expectation value and the maximum a posteriori estimate of the mean Ly α forest transmission around LAEs is verified.

by spectroscopic follow-up of the background LAEs and inclusion of near-infrared photometry to better constrain the background galaxy SEDs.

6 LAE-LY α FOREST CROSS-CORRELATION

6.1 Estimating the mean Ly α forest transmission around LAEs

The NB718 dataset can be used to reveal $z \approx 4.9$ LAEs in the same redshift slice as the IGM Ly α forest transmission and thus how the transmission varies as a function of angular separation θ from the foreground LAEs. As before, we can use a Monte Carlo sampling to evaluate the angular mean Ly α forest transmission around LAEs, $\bar{T}_{\text{IGM}}(\theta) = V(\theta)^{-1} \int T_{\text{IGM}}(x) dV(\theta)$. By angular averaging the transmission for all pairs of foreground LAEs and background galaxy sightlines, we obtain

$$\bar{T}_{\text{IGM}}(\theta) = \frac{1}{N_{\text{pair}}(\theta)} \sum_{j=1}^{N_{\text{fg}}} \sum_{i=1}^{N_{\text{bg}}} T_{\text{IGM},i} \mathcal{I}(|\theta - \theta_{ij}|), \quad (21)$$

where $N_{\text{pair}}(\theta) = \sum_{j=1}^{N_{\text{fg}}} \sum_{i=1}^{N_{\text{bg}}} \mathcal{I}(|\theta - \theta_{ij}|)$ is the number of pairs in each angular bin and $\mathcal{I}(|\theta - \theta_{ij}|)$ is an indicator function equal to unity when the angular separation $\theta_{ij} = |\theta_i - \theta_j|$ between i and j is in the angular bin specified by θ with width $\Delta\theta$, i.e. $\mathcal{I}(|\theta - \theta_{ij}|) = 1$ if $|\theta - \theta_{ij}| < \Delta\theta$, and zero otherwise. The sums run over all foreground LAEs $j = 1, \dots, N_{\text{fg}}$ and all background galaxy sightlines $i = 1, \dots, N_{\text{bg}}$.

Again as each background galaxy sightline provides a noisy measurement of $T_{\text{IGM},i}$, analogous to the argument made for the mean Ly α forest transmission, the full posterior probability of the mean Ly α forest transmission around LAEs, $P(\bar{T}_{\text{IGM}}(\theta) | \{f_{\text{NB},i}^{\text{obs}}, f_{\text{BB},i}^{\text{obs}}\}_{i=1,\dots,N_{\text{bg}}})$, can be expressed in terms of the posteriors of individual $T_{\text{IGM},i}$ measurements, which can be numerically computed by randomly sampling the individual posteriors. The expectation value of the angular mean Ly α forest transmission

around LAEs is given by (see Appendix A)

$$\langle \bar{T}_{\text{IGM}}(\theta) \rangle = \frac{1}{N_{\text{pair}}(\theta)} \sum_{j=1}^{N_{\text{fg}}} \sum_{i=1}^{N_{\text{bg}}} \langle T_{\text{IGM},i} \rangle \mathcal{I}(|\theta - \theta_{ij}|). \quad (22)$$

The variance of the estimated mean Ly α forest transmission around LAEs at each angular bin is given by

$$\text{Var}[\bar{T}_{\text{IGM}}(\theta)] = \frac{1}{N_{\text{pair}}(\theta)^2} \sum_{j=1}^{N_{\text{fg}}} \sum_{i=1}^{N_{\text{bg}}} \text{Var}[T_{\text{IGM},i}] \mathcal{I}(|\theta - \theta_{ij}|). \quad (23)$$

Figure 12 verifies that the full posterior $P(\bar{T}_{\text{IGM}}(\theta) | \{f_{\text{NB},i}^{\text{obs}}, f_{\text{BB},i}^{\text{obs}}\}_{i=1, \dots, N_{\text{bg}}})$ follows a Gaussian distribution, which can be fully characterised by an expectation value (22) and variance (23). The maximum *a posteriori* estimation,

$$\bar{T}_{\text{IGM}}^{\text{MP}}(\theta) = \arg \max_{\bar{T}_{\text{IGM}}(\theta)} P(\bar{T}_{\text{IGM}}(\theta) | \{f_{\text{NB},i}^{\text{obs}}, f_{\text{BB},i}^{\text{obs}}\}_{i=1, \dots, N_{\text{bg}}}), \quad (24)$$

is therefore equivalent to the expectation value $\langle \bar{T}_{\text{IGM}}(\theta) \rangle$. The error estimated by the square root of the variance (23) in the mean Ly α forest transmission around LAEs scales as $\propto 1/\sqrt{N_{\text{pair}}(\theta)}$ and includes the full uncertainties from photometric noise and continuum error in our individual T_{IGM} measurements from the Bayesian SED fitting framework.

6.2 The LAE-Ly α forest cross-correlation function

While the “mean Ly α forest transmission around LAEs” is well defined given the observed distribution of LAEs, in order to examine the angular “cross-correlation”, which is given by,

$$\omega_{\text{g}\alpha}(\theta) = \langle \bar{T}_{\text{IGM}}(\theta) \rangle / \langle \bar{T}_{\text{IGM}} \rangle - 1, \quad (25)$$

we must quantify whether the excess probability of finding galaxies in the environment of high or low Ly α forest transmission is statistically significant. An additional uncertainty arises from the Poisson sampling of foreground galaxies in the survey footprint, including the effect of mask regions and the edge of the field-of-view.

To understand the scatter, we generate a random galaxy catalogue. We populate N_{rand} objects at random locations $\{\theta_i^{\text{rand}}\}_{i=1, \dots, N_{\text{rand}}}$ within the survey footprint excluding the actual masked regions. We then repeat the measurement of the mean Ly α forest transmission around the random objects using the actual background sightlines,

$$\langle \bar{T}_{\text{IGM}}^{\text{random}}(\theta) \rangle = \frac{1}{N_{\text{pair}}(\theta)} \sum_{j=1}^{N_{\text{rand}}} \sum_{i=1}^{N_{\text{bg}}} \langle T_{\text{IGM},i} \rangle \mathcal{I}(|\theta - \theta_{ij}^{\text{rand}}|), \quad (26)$$

where $\theta_{ij}^{\text{rand}}$ is the angular separation between i -th random galaxy position and the observed location of j -th background galaxy sightline, $\theta_{ij}^{\text{rand}} = |\theta_i^{\text{rand}} - \theta_j|$. This de-correlates the real cross-correlation signal between LAEs and Ly α forest transmission and should converge to the mean Ly α forest transmission, $\langle \bar{T}_{\text{IGM}}^{\text{random}}(\theta) \rangle \approx \langle \bar{T}_{\text{IGM}} \rangle$.

One can also de-correlate by randomly shuffling the observed values of $T_{\text{IGM},i}$ among the background galaxy sightlines but keeping their angular locations fixed,

$$\langle \bar{T}_{\text{IGM}}^{\text{shuffle}}(\theta) \rangle = \frac{1}{N_{\text{pair}}(\theta)} \sum_{j=1}^{N_{\text{fg}}} \sum_{i=1}^{N_{\text{bg}}} \langle T_{\text{IGM},i}^{\text{shuffle}} \rangle \mathcal{I}(|\theta - \theta_{ij}|), \quad (27)$$

where $\langle T_{\text{IGM},i}^{\text{shuffle}} \rangle$ is a random draw from a set of real measurements $\{\langle T_{\text{IGM},1} \rangle, \langle T_{\text{IGM},2} \rangle, \dots, \langle T_{\text{IGM},N_{\text{bg}}} \rangle\}$ without replacement. The shuffled approach is convenient as it does not require us to

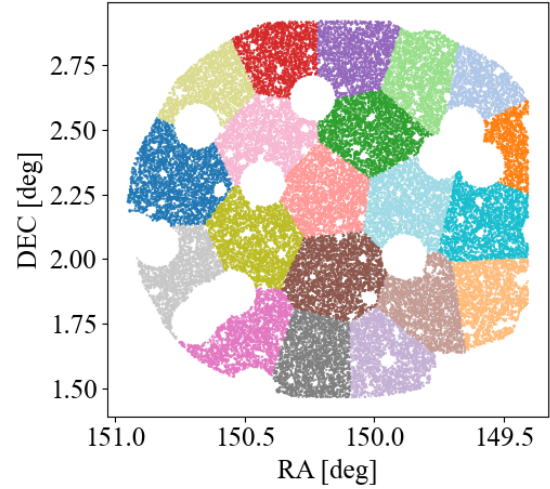


Figure 13. Angular regions (coloured patches) used to compute the Jackknife covariance matrix. Our survey footprint of the UD-COSMOS field are subdivided into $N_{\text{JK}} = 20$ regions.

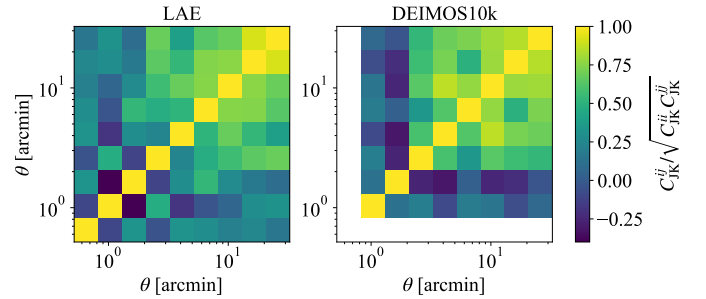


Figure 14. Correlation coefficients of the Jackknife covariance matrix for the LAE (left) and DEIMOS10k (right) samples. The covariance matrix at the innermost bin for the DEIMOS10k sample is not determined due to the small sample size in that bin.

model the galaxy selection function. This serves to verify that the observed LAE-Ly α forest cross-correlation is uncontaminated by the particular distribution of background galaxies on the sky. Both the shuffled and random measurements should be statistically identical, $\langle \bar{T}_{\text{IGM}}^{\text{shuffle}}(\theta) \rangle \approx \langle \bar{T}_{\text{IGM}}^{\text{random}}(\theta) \rangle$, and should converge to the mean Ly α forest transmission $\langle \bar{T}_{\text{IGM}} \rangle$ within the statistical error.

We estimate the error on the cross-correlation using the Jackknife estimator (e.g. Norberg et al. 2009). The Jackknife covariance matrix is given by

$$\text{Cov}_{\text{JK}} [\omega_{\text{g}\alpha}(\theta), \omega_{\text{g}\alpha}(\theta')] = \frac{N_{\text{JK}} - 1}{N_{\text{JK}}} \sum_{k=1}^{N_{\text{JK}}} \left[\omega_{\text{g}\alpha}^k(\theta) - \bar{\omega}_{\text{g}\alpha}^{\text{JK}}(\theta) \right] \left[\omega_{\text{g}\alpha}^k(\theta') - \bar{\omega}_{\text{g}\alpha}^{\text{JK}}(\theta') \right], \quad (28)$$

where

$$\bar{\omega}_{\text{g}\alpha}^{\text{JK}}(\theta) = \frac{1}{N_{\text{JK}}} \sum_{k=1}^{N_{\text{JK}}} \omega_{\text{g}\alpha}^k(\theta), \quad (29)$$

is the average of the cross-correlation functions from Jackknife resampling. Jackknife regions are obtained using the k-means cluster-

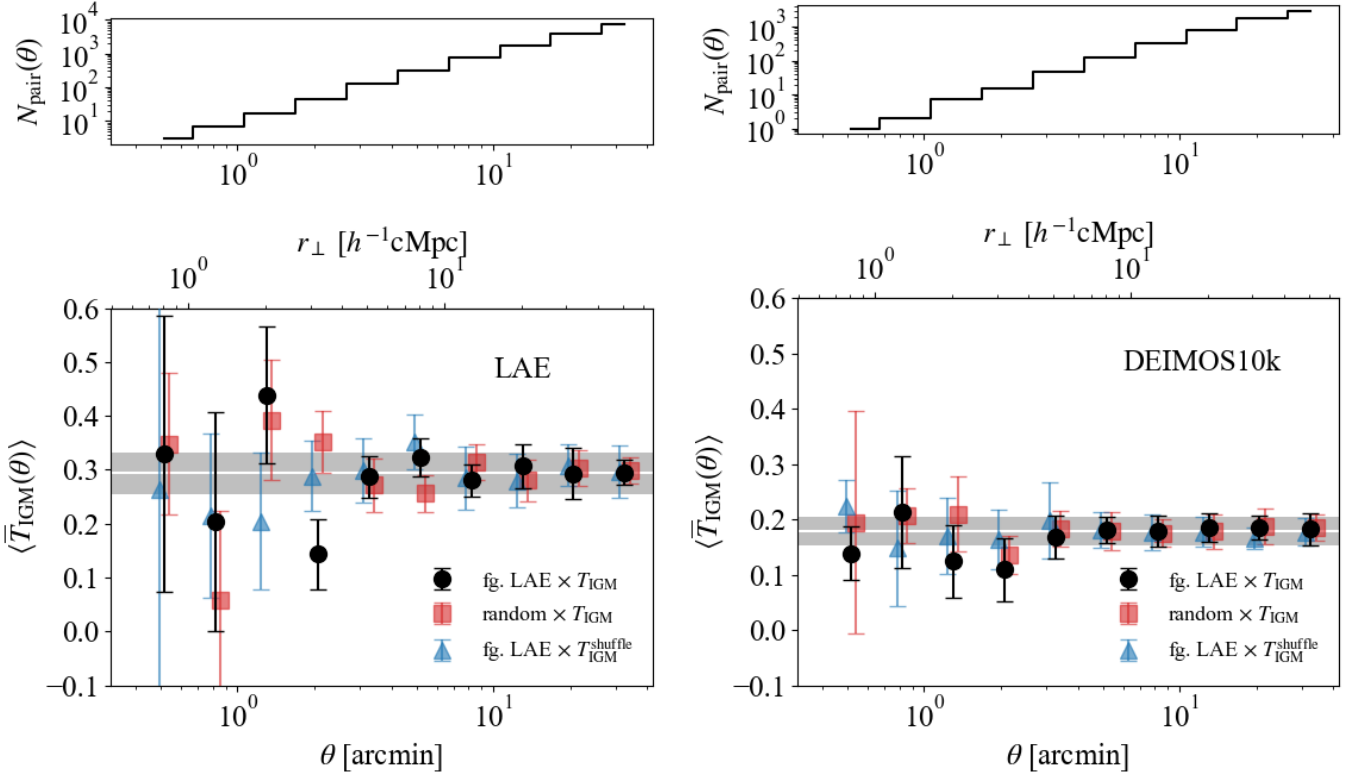


Figure 15. Mean Ly α forest transmission profiles around $z = 4.9$ LAEs (black circles) using the $z = 5.7$ LAE (left) and DEIMOS10k (right) samples as background sources. The same profiles around random foreground objects (red squares) and around foreground LAEs but using shuffled T_{IGM} along the background sources (blue triangles) are offset by ± 0.02 dex offset along the x-axis for clarity. The gray shaded region indicates the mean Ly α forest transmission and its 1σ error. Top panels indicate the number of foreground LAE - background sightline pairs for the LAE and DEIMOS10k samples respectively.

ing algorithm⁸ (Kwan et al. 2017) on a random galaxy catalogue with $N_{\text{rand}} = 50,000$. This algorithm subdivides the observed survey area into N_{JK} regions of a roughly equal area as shown in Figure 13. To compute the Jackknife covariance, we omit foreground LAEs and T_{IGM} along background galaxy sightlines located in each Jackknife region at a time and compute N_{JK} Jackknife re-sampled cross-correlation functions $\omega_{g\alpha}^k(\theta)$, $k = 1, \dots, N_{\text{JK}}$, using the remaining objects. We use the same procedure to compute the Jackknife error for the other summary statistics (the mean Ly α forest transmission and Ly α forest auto-correlation function) in this paper.

The correlation coefficients of the Jackknife covariance matrix are shown in Figure 14. The covariance matrix of the innermost bins of the DEIMOS10k sample could not be determined due to the small sample size. For both LAE and DEIMOS10k samples, there are significant off-diagonal correlations between angular bins as the same sightlines contribute multiple foreground LAE-background sightline pairs.

6.3 Result

Figure 15 shows the angular mean Ly α forest transmission around LAEs at $z \approx 4.9$. We find no excess transmission or absorption in the NB-integrated Ly α forest around the LAEs. The result is consistent with the global mean $\langle \bar{T}_{\text{IGM}} \rangle$ within the 2σ error both for the LAE and DEIMOS10k samples. We compare our result with the random

and shuffled measurements using the same number of foreground LAEs and background galaxies in the real data. Both measurements show similar fluctuations with the observed values, confirming that our result is consistent with no spatial correlation.

6.3.1 Correcting for contamination by low-redshift interlopers

Figure 15 shows a mean offset between $\langle \bar{T}_{\text{IGM}}(\theta) \rangle$ measured using the LAE and DEIMOS10k samples. As discussed in Section 5.3, this is likely caused by the low-redshift interlopers in both the foreground and background LAE samples (Grasshorn Gebhardt et al. 2019; Farrow et al. 2021). Since the distribution of any low-redshift interlopers would be random relative to structures in the tomographic slice of interest, the interlopers will dilute the observed cross-correlation. Assuming foreground and background contamination fractions $f_{\text{fg.int}}$ and $f_{\text{bg.int}}$, the observed angular mean Ly α forest transmission around the foreground LAEs can be expressed as (see Appendix B)

$$\begin{aligned} \langle \bar{T}_{\text{IGM}}(\theta) \rangle = & (1 - f_{\text{fg.int}})(1 - f_{\text{bg.int}})\langle \bar{T}_{\text{IGM}}(\theta) \rangle^{\text{true}} \\ & + f_{\text{fg.int}}(1 - f_{\text{bg.int}})\langle \bar{T}_{\text{IGM}} \rangle^{\text{true}} + f_{\text{bg.int}}\langle \bar{T}_{\text{IGM}} \rangle^{\text{bg.int}}, \end{aligned} \quad (30)$$

where $\langle \bar{T}_{\text{IGM}} \rangle^{\text{true}}$ is the true mean Ly α forest transmission and $\langle \bar{T}_{\text{IGM}} \rangle^{\text{bg.int}}$ is the fictitious mean Ly α forest transmission measured along the low-redshift interlopers in the background LAE sample. The second and third terms indicate contaminations from the low-redshift interlopers in the foreground and background LAE samples.

⁸ https://github.com/esheldon/kmeans_radec

Following our definition of the observed LAE-Ly α forest cross-correlation $\omega_{g\alpha}(\theta) = \langle \bar{T}_{\text{IGM}}(\theta) \rangle / \langle \bar{T}_{\text{IGM}} \rangle - 1$, we can similarly find that low-redshift interlopers dilute the cross-correlation amplitude by

$$\omega_{g\alpha}(\theta) = \frac{(1 - f_{\text{fg.int}})(1 - f_{\text{bg.int}})\langle \bar{T}_{\text{IGM}} \rangle^{\text{true}}}{(1 - f_{\text{bg.int}})\langle \bar{T}_{\text{IGM}} \rangle^{\text{true}} + f_{\text{bg.int}}\langle \bar{T}_{\text{IGM}} \rangle^{\text{bg.int}}} \omega_{g\alpha}^{\text{true}}(\theta). \quad (31)$$

The offset in the mean Ly α forest transmission around LAEs between the background LAE and DEIMOS10k samples can be explained by this effect. As in Section 5.3, we set $f_{\text{fg.int}} = f_{\text{bg.int}} = 0.2$ both for foreground and background LAE samples. The contamination fraction for the DEIMOS10k sample is $f_{\text{bg.int}} = 0.0$ as all are confirmed spectroscopically. We assume that the fictitious mean Ly α forest transmission along the interlopers in the background LAE sample is $\langle \bar{T}_{\text{IGM}} \rangle^{\text{bg.int}} = 0.7$. As before, LAE interlopers can explain the offset between the background LAE and DEIMOS10k samples.

These interlopers depress the observed LAE-Ly α forest cross-correlation by $\omega_{g\alpha}(\theta) \approx 0.40 \omega_{g\alpha}^{\text{true}}(\theta)$ and $\omega_{g\alpha}(\theta) \approx 0.80 \omega_{g\alpha}^{\text{true}}(\theta)$ for background LAE and DEIMOS10k samples, respectively. The cross-correlation measurement using the DEIMOS10k sample is also affected because of the interloper contamination in the foreground LAEs.

All the terms in the damping pre-factor in Equation 31 can be determined and statistically corrected *a posteriori* by spectroscopic follow up of a random subset of the foreground and background LAE samples as described in Section 5.3.

6.3.2 Limit on the IGM fluctuations around LAEs

Figure 16 shows the observed LAE-Ly α forest cross-correlation at $z \approx 4.9$ after correcting for possible low-redshift interloper contamination. To place an empirical constraint, we assume a simple power-law form,

$$\omega_{g\alpha}^{\text{model}}(\theta) = A_0(\theta/\theta_0)^{-\gamma}, \quad (32)$$

where the fluctuations are characterised by the amplitude A_0 at angular distance θ_0 and the power-law slope γ . We assume a Gaussian likelihood with the measured Jackknife covariance matrix and a fixed slope of $\gamma = 0.5$. The resulting 3σ bounds are shown in Figure 16. The 3σ lower and upper limits are

$$-0.29 \left(\frac{r_{\perp}}{10 h^{-1} \text{cMpc}} \right)^{-0.5} < \omega_{g\alpha}^{\text{model}} < 0.07 \left(\frac{r_{\perp}}{10 h^{-1} \text{cMpc}} \right)^{-0.5} \quad (33)$$

for the DEIMOS10k sample, and

$$-0.58 \left(\frac{r_{\perp}}{10 h^{-1} \text{cMpc}} \right)^{-0.5} < \omega_{g\alpha}^{\text{model}} < 0.40 \left(\frac{r_{\perp}}{10 h^{-1} \text{cMpc}} \right)^{-0.5} \quad (34)$$

for the LAE sample. The derived lower and upper limits are consistent for both the LAE and DEIMOS10k samples. While the bound from the DEIMOS10k sample is slightly shifted to the negative cross-correlation, this is likely due to an underestimated Jackknife error at small angular bins. The directly propagated error (equation 23) from the Bayesian SED fitting framework indicate the error from photometric noise and UV continuum uncertainty in the DEIMOS10k sample at the inner bins should be larger than the empirical estimate from the Jackknife method.

Our result indicates that the angular fluctuations of the Ly α forest transmission around LAEs should be $\leq 58\%$ at $10 h^{-1} \text{Mpc}$ relative

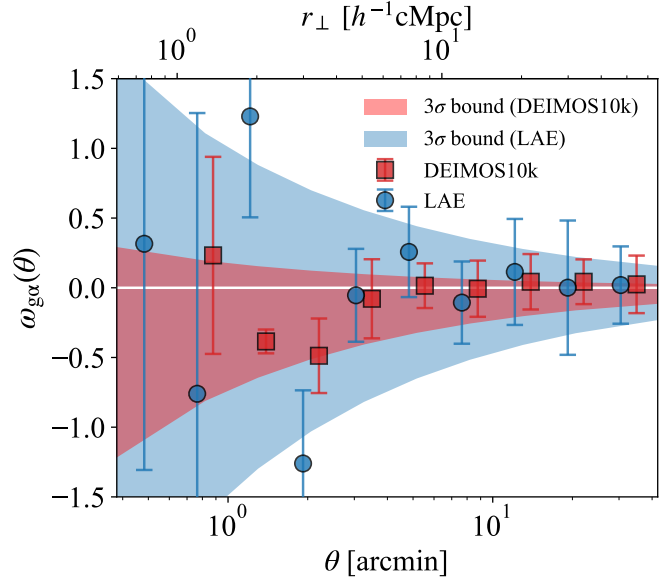


Figure 16. Observed LAE-Ly α forest cross-correlation function for the DEIMOS10k (red squares) and LAE (blue circles) samples after correcting for likely interloper contamination. The derived 3σ lower and upper limits of the cross-correlation assuming the power-law with slope $\gamma = 0.5$ are shown with the red and blue shaded regions for the DEIMOS10k and LAE sample, respectively. The Jackknife covariance matrices scaled by the interloper correction factors are used to estimate the error.

to the global mean at $z \approx 4.9$. The physical interpretation of the result will be discussed in a companion paper (Kakiichi et al in prep).

7 LY α FOREST AUTO-CORRELATION

7.1 Estimating the Ly α forest auto-correlation

We now turn our attention to examine the spatial fluctuations of Ly α forest transmission. These are expected to spatially correlate across different sightlines due to large-scale fluctuations of the IGM. Unlike the measurement of the 3D Ly α forest auto-correlation from spectra (e.g. Slosar et al. 2011), photometric IGM tomography measures the angular auto-correlation of Ly α forest transmission integrated over the line-of-sight width of the NB filter ($\approx 34 h^{-1} \text{cMpc}$).

In order to estimate the Ly α forest angular auto-correlation function, using the pairs of NB-integrated Ly α forest transmission measurements, we first compute, at each angular bin,

$$T_{\text{IGM}} T_{\text{IGM}}(\theta) = \frac{1}{N_{\text{pair}}(\theta)} \sum_{i=1}^{N_{\text{bg}}} \sum_{j>i}^{N_{\text{bg}}} T_{\text{IGM},i} T_{\text{IGM},j} \mathcal{I}(|\theta - \theta_{ij}|), \quad (35)$$

where $N_{\text{pair}}(\theta) = \sum_{i=1}^{N_{\text{bg}}} \sum_{j>i}^{N_{\text{bg}}} \mathcal{I}(|\theta - \theta_{ij}|)$ is the number of pairs in each angular bin. For independent measurements of $T_{\text{IGM},i}$, the expectation value of the Ly α forest angular auto-correlation is given by

$$\langle T_{\text{IGM}} T_{\text{IGM}}(\theta) \rangle = \frac{1}{N_{\text{pair}}(\theta)} \sum_{i=1}^{N_{\text{bg}}} \sum_{j>i}^{N_{\text{bg}}} \langle T_{\text{IGM},i} \rangle \langle T_{\text{IGM},j} \rangle \mathcal{I}(|\theta - \theta_{ij}|). \quad (36)$$

The error is computed from the Jackknife covariance matrix. The

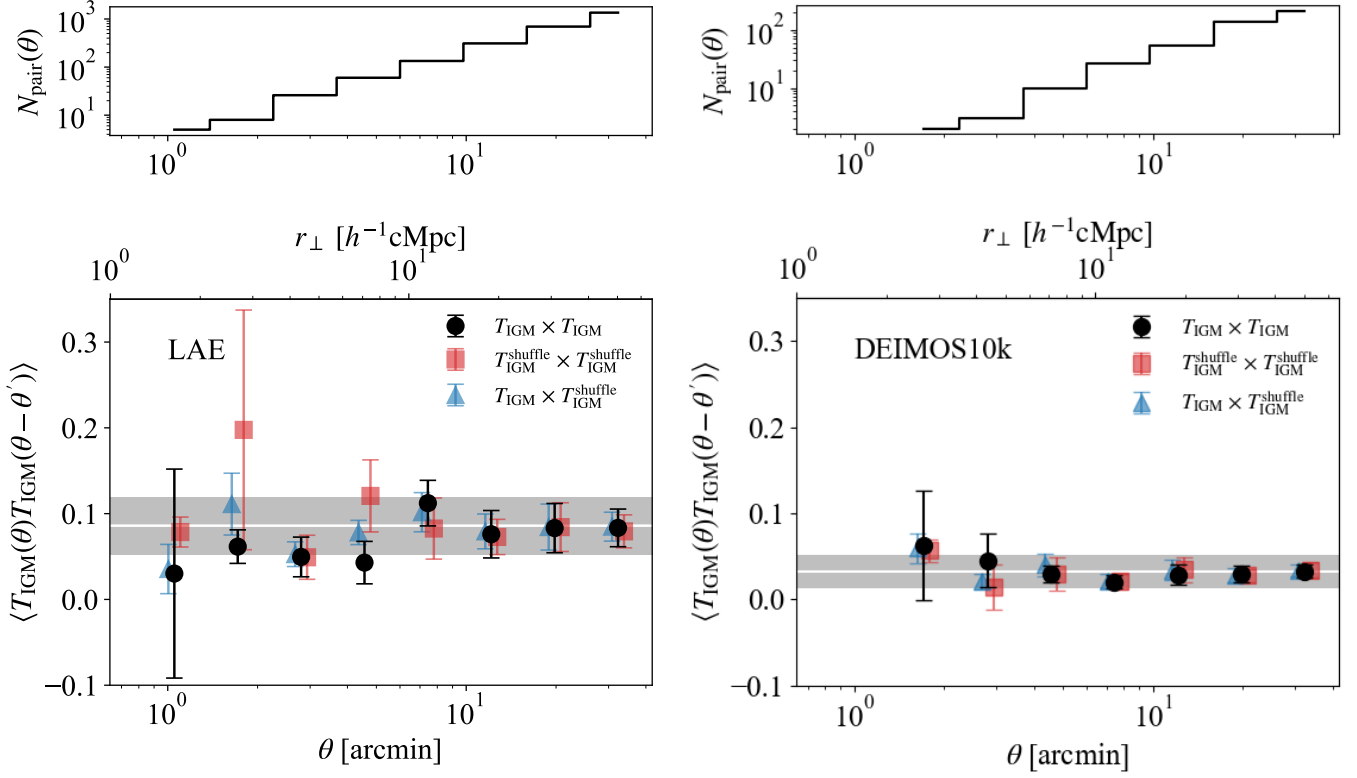


Figure 17. Ly α forest transmission auto-correlation function at $z = 4.9$ using the $z = 5.7$ LAE (left) and DEIMOS10k (right) samples. Values computed using only shuffled background sources (red squares) and the cross-correlation between data and shuffled samples (blue triangles) are shown offset ± 0.02 dex along the x-axis for clarity. The gray shaded region indicates the estimate in the case of no correlation based on the mean Ly α forest transmission and its 1σ error. Top panels indicate the number of sightline pairs for the LAE and DEIMOS10k samples.

Ly α forest auto-correlation function is estimated by

$$\omega_{\alpha\alpha}(\theta) = \langle T_{\text{IGM}} T_{\text{IGM}}(\theta) \rangle / \langle \bar{T}_{\text{IGM}} \rangle^2 - 1. \quad (37)$$

We can understand the scatter of $\langle T_{\text{IGM}} T_{\text{IGM}}(\theta) \rangle$ in the absence of any spatial correlation. As the Ly α forest auto-correlation has no complication from the window function or the survey geometry, $\langle T_{\text{IGM}} T_{\text{IGM}}(\theta) \rangle$ should be equal to $\langle \bar{T}_{\text{IGM}} \rangle^2$ if there is no spatial correlation. To test this, we de-correlate the observed correlation by shuffling either one or both of the measured values of T_{IGM} between the observed locations, i.e.

$$\langle T_{\text{IGM}} T_{\text{IGM}}^{\text{shuffle}}(\theta) \rangle = \frac{1}{N_{\text{pair}}(\theta)} \sum_{i=1}^{N_{\text{bg}}} \sum_{j>i}^{N_{\text{bg}}} \langle T_{\text{IGM},i} \rangle \langle T_{\text{IGM},j}^{\text{shuffle}}(\theta) \rangle \mathcal{I}(|\theta - \theta_{ij}|), \quad (38)$$

or

$$\langle T_{\text{IGM}}^{\text{shuffle}}(\theta) T_{\text{IGM}}^{\text{shuffle}}(\theta) \rangle = \frac{1}{N_{\text{pair}}(\theta)} \sum_{i=1}^{N_{\text{bg}}} \sum_{j>i}^{N_{\text{bg}}} \langle T_{\text{IGM},i}^{\text{shuffle}}(\theta) \rangle \langle T_{\text{IGM},j}^{\text{shuffle}}(\theta) \rangle \mathcal{I}(|\theta - \theta_{ij}|). \quad (39)$$

Using the real set of $\{T_{\text{IGM},i}\}_{i=1,\dots,N_{\text{bg}}}$, we generated a randomized set of T_{IGM} values keeping the angular positions of the sightlines the same. This artificially de-correlates the possible correlation. If there is no systematic, this should approach $\approx \langle \bar{T}_{\text{IGM}} \rangle^2$.

7.2 Result

Figure 17 shows the observed auto-correlation function of the Ly α forest transmission at $z \approx 4.9$. The observed auto-correlation is consistent with the square of the mean and the shuffled results within 2σ error, indicating the observed signal is consistent with no auto-correlation. This null detection can be interpreted as the observed limit on the Ly α forest transmission fluctuations at $z \approx 4.9$.

7.2.1 Correcting for the contamination by low-redshift interlopers

Similar to the angular mean Ly α forest transmission around LAEs, Figure 17 shows an offset between $\langle T_{\text{IGM}} T_{\text{IGM}}(\theta) \rangle$ measured using the background LAE and DEIMOS10k samples, which is likely caused by the low-redshift interlopers. The effect of the lower-redshift interlopers in $\langle T_{\text{IGM}} T_{\text{IGM}}(\theta) \rangle$ can be expressed as (see Appendix B)

$$\begin{aligned} \langle T_{\text{IGM}} T_{\text{IGM}}(\theta) \rangle &= (1 - f_{\text{bg.int}})^2 \langle T_{\text{IGM}} T_{\text{IGM}}(\theta) \rangle^{\text{true}} + \\ &2(1 - f_{\text{bg.int}}) f_{\text{bg.int}} \langle T_{\text{IGM}} \rangle^{\text{true}} \langle T_{\text{IGM}} \rangle^{\text{bg.int}} + (f_{\text{bg.int}} \langle T_{\text{IGM}} \rangle^{\text{bg.int}})^2. \end{aligned} \quad (40)$$

The second and third terms indicate contaminations from cross-correlation between low-redshift interlopers and true background galaxies and the auto-correlation of low-redshift interlopers, assuming there is no spatial correlation. In terms of the Ly α forest angular auto-correlation function, the true auto-correlation function $\omega_{\alpha\alpha}^{\text{true}}(\theta) = \langle T_{\text{IGM}} T_{\text{IGM}}(\theta) \rangle^{\text{true}} / \langle \bar{T}_{\text{IGM}} \rangle^{\text{true}}^2 - 1$ is diluted by the

interlopers such that

$$\omega_{\alpha\alpha}(\theta) = \left[\frac{(1 - f_{\text{bg.int}}) \langle \bar{T}_{\text{IGM}} \rangle^{\text{true}}}{(1 - f_{\text{bg.int}}) \langle \bar{T}_{\text{IGM}} \rangle^{\text{true}} + f_{\text{bg.int}} \langle \bar{T}_{\text{IGM}} \rangle^{\text{bg.int}}} \right]^2 \omega_{\alpha\alpha}^{\text{true}}(\theta). \quad (41)$$

Again, all factors can be determined *a posteriori* using spectroscopic follow-up of the background galaxy sample. Assuming $f_{\text{bg.int}} = 0.20$ and $\langle \bar{T}_{\text{IGM}} \rangle^{\text{bg.int}} = 0.7$ for the background LAE sample and taking $\langle \bar{T}_{\text{IGM}} \rangle^{\text{true}}$ to be the value from the DEIMOS10k sample can explain the observed offset in $\langle T_{\text{IGM}} T_{\text{IGM}}(\theta) \rangle$. This corresponds to the damping of $\omega_{\alpha\alpha}(\theta) = 0.39 \omega_{\alpha\alpha}^{\text{true}}(\theta)$ for the observed Ly α forest angular auto-correlation function from the LAE sample. There is no damping factor for the DEIMOS10k sample as the interloper contamination for the spectroscopically confirmed sample is zero.

8 IGM TOMOGRAPHIC MAP

8.1 Reconstruction method

Finally, we present a reconstructed tomographic map of the IGM. This is arguably the most unique aspect of photometric IGM tomography, since it enables us to directly visualise the large-scale structures of the IGM and galaxies in the same cosmic volume. To accomplish this we use the Nadaraya-Watson estimator for the 2D tomographic map of the IGM Ly α forest transmission fluctuations (Kakiichi et al. 2022),

$$\Delta T_{\text{IGM}}^{2\text{D}}(\theta) = \frac{\sum_{i=1}^{N_{\text{bg}}} K_R(\theta - \theta_i) (T_{\text{IGM},i} - \langle \bar{T}_{\text{IGM}} \rangle)}{\sum_{i=1}^{N_{\text{bg}}} K_R(\theta - \theta_i)}, \quad (42)$$

where $K_R(\theta) = (2\pi R^2)^{-1/2} \exp[-\theta^2/(2R^2)]$ is a Gaussian kernel with a smoothing length R . In practice, we create a 2D map on pixelised map of 4096×4096 pixels. Following a similar argument as in previous sections, given the posterior of Ly α forest transmission, the expectation value of the 2D tomographic map is given by

$$\langle \Delta T_{\text{IGM}}^{2\text{D}}(\theta) \rangle = \frac{\sum_{i=1}^{N_{\text{bg}}} K_R(\theta - \theta_i) (\langle T_{\text{IGM},i} \rangle - \langle \bar{T}_{\text{IGM}} \rangle)}{\sum_{i=1}^{N_{\text{bg}}} K_R(\theta - \theta_i)}. \quad (43)$$

At each point θ , the estimator is simply the weighted sum of (independent) individual $T_{\text{IGM},i}$ measurements. Thus the variance can be computed as

$$\text{Var}[\Delta T_{\text{IGM}}^{2\text{D}}(\theta)] = \frac{\sum_{i=1}^{N_{\text{bg}}} K_R(\theta - \theta_i)^2 \text{Var}[T_{\text{IGM},i}]}{\left[\sum_{i=1}^{N_{\text{bg}}} K_R(\theta - \theta_i) \right]^2}. \quad (44)$$

This includes errors from photometric noise and the UV continuum uncertainty. We define the SNR map as the ratio between the observed Ly α forest transmission fluctuations and the standard deviation,

$$\text{SNR}(\theta) = \frac{|\langle \Delta T_{\text{IGM}}^{2\text{D}}(\theta) \rangle|}{\sqrt{\text{Var}[\Delta T_{\text{IGM}}^{2\text{D}}(\theta)]}}. \quad (45)$$

8.2 Result

In Figure 18 we show the reconstructed 2D tomographic map of the Ly α forest transmission at $z \simeq 4.9$. We only apply the map reconstruction to the background LAE sample because their distribution spans the entire field of view. The smoothing length is chosen as the mean

inter-sightline separation, $R = 1/\sqrt{\Sigma_{\text{LAE}}} \simeq 0.118$ deg (7.1') where Σ_{LAE} is the surface number density of the background LAEs. We can visually see large-scale fluctuations of the Ly α forest transmission with a median contrast of $|\langle \Delta T_{\text{IGM}}^{2\text{D}}(\mathbf{x}) \rangle| \sim 0.1$. The typical standard

deviation in the reconstructed map is $\sqrt{\text{Var}[\Delta T_{\text{IGM}}^{2\text{D}}(\mathbf{x})]} \sim 0.14$. We find the mean SNR of the reconstructed map as $\langle \text{SNR}(\mathbf{x}) \rangle = 0.71$, indicating that our tomographic map is still noisy with contributions from photometric errors and the UV continuum uncertainty.

There are several tentative regions of transmissive and opaque transmission in the map located at (RA, DEC) \simeq (150.2°, 2.45°) and (RA, DEC) \simeq (150.0°, 1.95°) respectively, with the peak SNR \simeq 1.5 – 2.0. The Ly α forest transmission in these regions is reasonably coherent. We require deeper NB imaging data to secure the statistical significance. If confirmed, these opaque and transmissive regions of the IGM may represent a protocluster and a highly ionized region of the IGM by the enhanced UV background.

Although there is a higher transmissive region towards the edge at (RA, DEC) \simeq (150.8°, 1.80°), as there is no background galaxy here this is likely an artefact from the map reconstruction method. At the edge of the field of view, the reconstruction method is more affected by boundary effects. Although our estimator corrects for boundary effects by incorporating the sightline density, the estimator is more sensitive to the T_{IGM} values of individual background galaxies, whereas at the centre of the field smoothing corrects outlier values of T_{IGM} .

Figure 19 overlays the distribution of the $z = 4.9$ LAEs on the reconstructed tomographic map of the Ly α forest transmission. This represents the highest redshift 2D tomographic map of the IGM with the galaxy distribution at the present time and demonstrates the potential of the NB tomographic technique to examine the galaxy-IGM connection closer to the reionization epoch.

We briefly examine the spatial correlation between the $z = 4.9$ LAE distribution and the reconstructed Ly α forest transmission map. In order to apply the spatial correlation analysis at the map level, we first reconstruct LAE density map at the same smoothing scale using the Gaussian kernel density estimator with the boundary correction,

$$n_{\text{LAE}}(\theta) = \frac{\sum_{i=1}^{N_{\text{bg}}} w_i K_R(\theta - \theta_i)}{\int m(\theta') K_R(\theta - \theta') d\theta'}, \quad (46)$$

where $m(\theta')$ represents the mask and the denominator is the correction factor $C^{-1} = \int m(\theta') K_R(\theta - \theta') d\theta'$ for boundary effects including masked regions around bright stars. We introduce weights w_i , where $w_i = 1$ for the ordinary LAE density field and $w_i = L_{\alpha,i}$ for the Ly α luminosity-weighted LAE density field. The galaxy density fluctuation map is then $\delta_{\text{LAE}}(\mathbf{x}) = n_{\text{LAE}}/\bar{n}_{\text{LAE}} - 1$ where the mean density \bar{n}_{LAE} is computed by excluding the masked regions.

In Figure 20 we show the comparison of the 2D tomographic map with the LAE density and Ly α luminosity-weighted LAE density fields at the same smoothing length. Using all unmasked regions, the Pearson correlation coefficient ρ indicates that there is a negligible correlation between the reconstructed 2D tomographic map and the (luminosity-weighted) LAE density map with $\rho = 0.09$ (0.06). Within the precision of existing photometric data, it appears that $z \simeq 4.9$ LAEs do not occupy extreme Ly α transmissive or opaque regions of the IGM. This is consistent with the two-point cross-correlation analysis between LAE and Ly α forest transmission.

To avoid confusion from boundary effects, we have also limited the map-level analysis within the region with a low correction factor $C < 2.0$. We then find that the Pearson correlation coefficient becomes $\rho = 0.22$ for the LAE density field (0.20 for luminosity-weighted

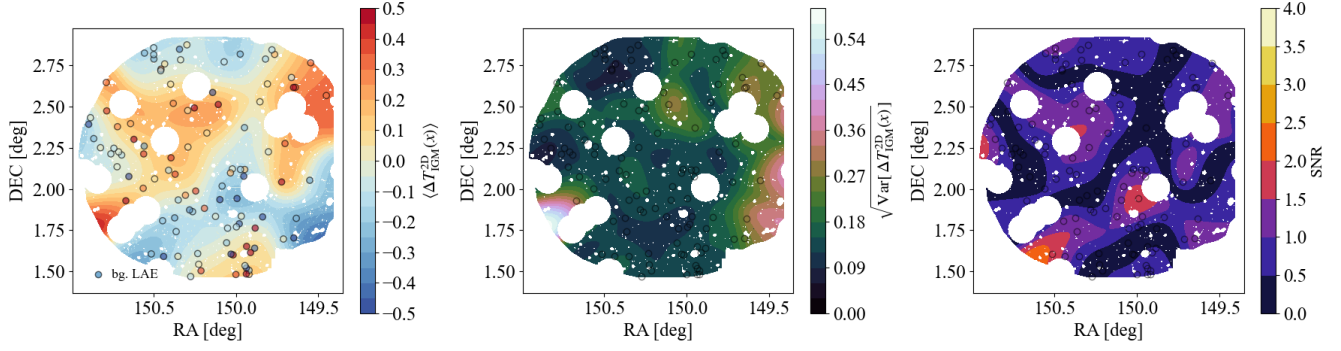


Figure 18. (Left): Reconstructed 2D tomographic map of the Ly α forest transmission of the IGM at $z \approx 4.9$ (colour map). The locations of the background ($z \approx 5.7$) LAEs are shown with coloured circles with the colours indicating the measured difference between the Ly α forest transmission along the sightline and the global mean. (Middle): The standard deviation of the reconstructed map including the uncertainties from the photometric and continuum errors in the Bayesian SED fitting. The locations of the background LAE sightlines are shown with open circles. (Right): The signal-to-noise ratio of the reconstructed map. The masked regions are left blank.

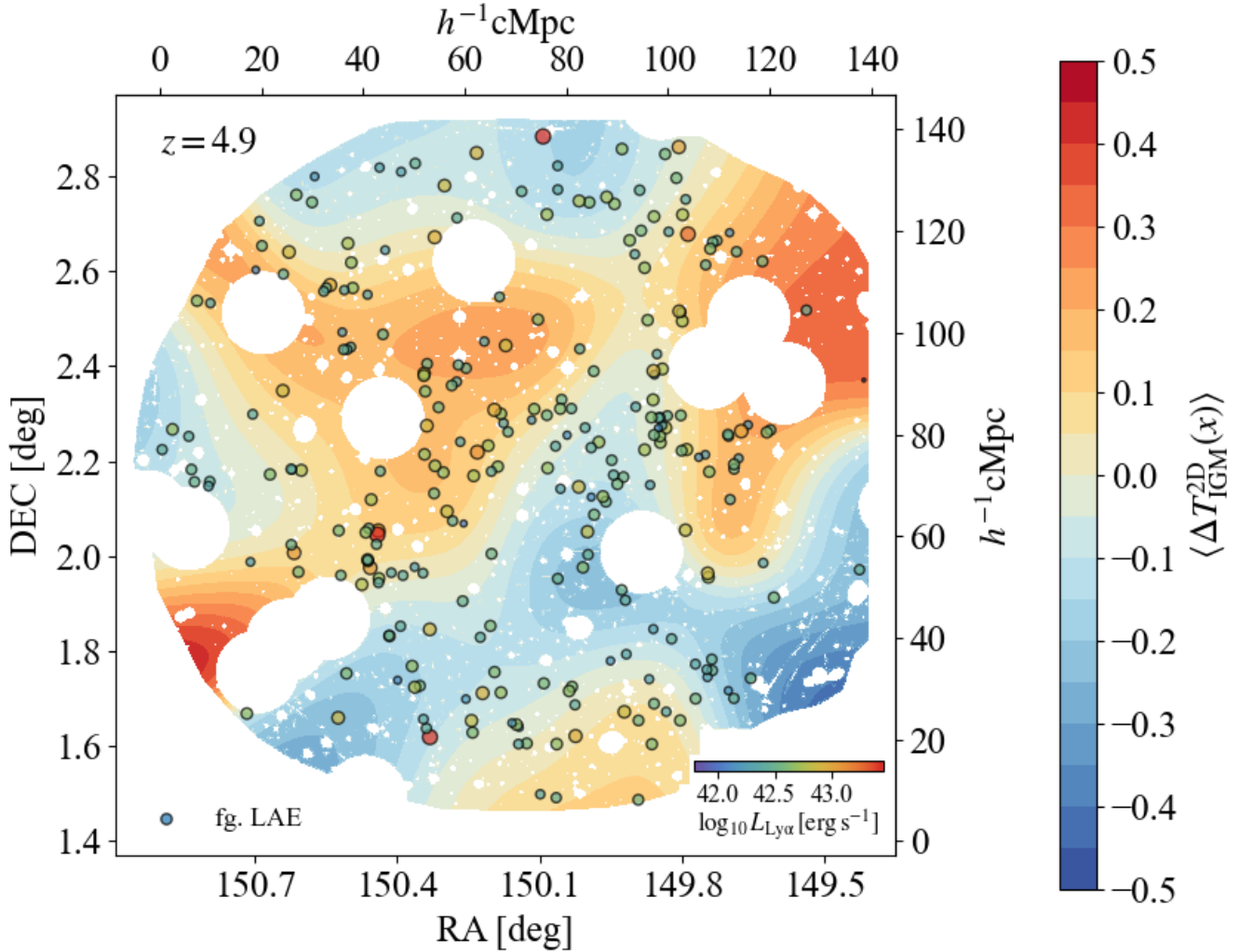


Figure 19. Reconstructed 2D tomographic map of the Ly α forest transmission of the IGM (coloured map) overlaid with the distribution of the $z = 4.9$ LAEs (coloured circles). The angular resolution of the reconstructed map is 0.12 deg corresponding to $11 h^{-1} \text{ cMpc}$. The colour of each circle indicates the Ly α luminosity of the LAE. Masked regions are left blank. Although the map is dominated by photometric noise, it represents the first IGM tomographic map co-spatial to the large-scale structure of LAEs in the same cosmic volume close to the end of cosmic reionization.

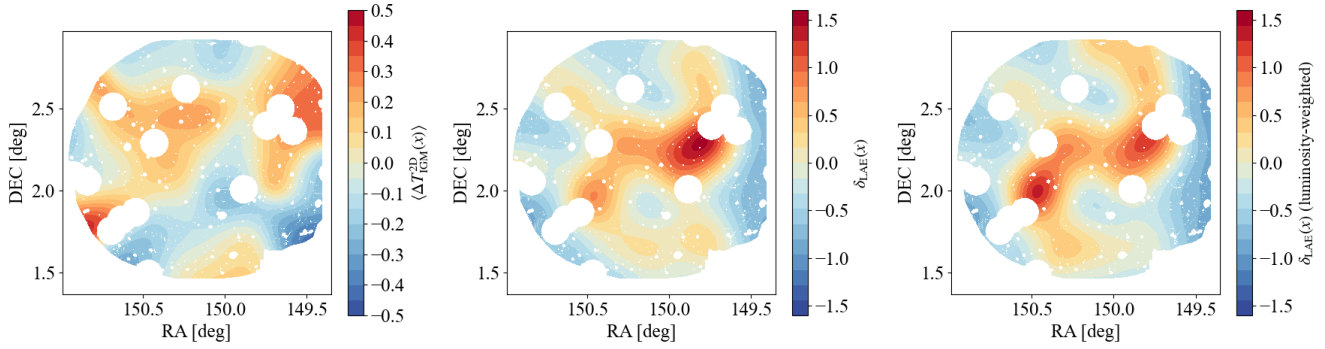


Figure 20. Fluctuations of the reconstructed 2D $\text{Ly}\alpha$ forest tomographic map (left), LAE number density field centre, and $\text{Ly}\alpha$ luminosity-weighted LAE number density field (right) with the angular resolution of 0.12 deg ($11 h^{-1} \text{ cMpc}$) at $z \approx 4.9$.

field), indicating a possible weak positive correlation between the LAE number density and the $\text{Ly}\alpha$ forest transmission map of the IGM. Although this weak correlation is intriguing, as noted above, our average SNR of the map is still low to conclude.

8.3 Systematics

Unlike the cross- and auto-correlation functions between LAEs and $\text{Ly}\alpha$ forest transmission, the reconstruction of the 2D tomographic map demands a higher purity of the background galaxy sample. A fictitious $\text{Ly}\alpha$ forest transmission due to a low-redshift interloper would produce a fake transmissive IGM region. While having many background galaxies within a smoothing length of the reconstruction reduces the effect of interlopers, we have not found a way to statistically correct for this effect as is possible statistically using a spectroscopic subset in the case of the correlation functions. It is difficult to quantify the level of contamination in each transmissive or opaque region of the IGM identified with photometric IGM tomography without directly confirming the redshifts of all background galaxies spectroscopically.

9 DISCUSSION

9.1 Improving photometric IGM tomography

9.1.1 Extremely-deep NB imaging and spectroscopic campaign

As discussed in Section 4.2, our estimate of T_{IGM} from individual background galaxies is dominated by photometric noise. This error propagates into our measurements of LAE- $\text{Ly}\alpha$ forest cross-correlation (Section 6) and $\text{Ly}\alpha$ forest auto-correlation functions (Section 7) and the reconstruction of the 2D tomographic map of the IGM (Section 8). As the measured $\text{Ly}\alpha$ forest transmission depends on the contrast between the foreground NB flux and the BB flux of a background galaxy, the noise scales approximately as $\delta T_{\text{IGM}} \approx \delta f_{\text{NB}}/f_{\text{BB}}$. To improve the signal-to-noise ratio of photometric IGM tomography, we therefore require (i) deeper imaging in the foreground NB filter and/or (ii) enlarge the sample of bright (spectroscopically-confirmed) background galaxies for which we can more accurately measure T_{IGM} given a great contrast between the NB and BB filters.

While the current NB718 depth (26.86 mag, 3σ) provides 3σ sensitivity to a typical mean value of the $\text{Ly}\alpha$ forest transmission when using 25.2 mag (z -band) background sources, the majority of our background galaxies are fainter. Secure ($> 3\sigma$) detection of the $\text{Ly}\alpha$

forest transmitted flux along individual sightlines is currently only possible for rare bright background galaxies ($M_{\text{UV}} \lesssim -21.4 \text{ mag}$). Extremely-deep NB imaging reaching 27.6 mag (28.2 mag) at 3σ depth would allow us to detect typical $\text{Ly}\alpha$ forest transmitted fluxes to fainter 26.0 – 26.4 mag (26.6 – 27.0 mag) background sources at $2 - 3\sigma$ significance level which comprises $\approx 50\%$ (90%) of our background LAE+DEIMOS10k sample. Based on the existing depth of NB718 from CHORUS PDR1 after $t_{\text{exp}} = 7.7$ hour exposures (Inoue et al. 2020), and assuming a factor of $\propto t_{\text{exp}}^{-1/2}$ reduction of photometric noise, such an extremely deep observation would require a total exposure of ≈ 28 (100) hours in NB718. As the reconstructed tomographic map of the IGM is dominated by the photometric noise, the improvement in the NB image quality by longer integration will directly increase the signal-to-noise ratio of the IGM tomographic map. While this may seem a significant investment of the telescope time, given a large number of potential science applications of photometric IGM tomography as we will discuss later, such an investment would be of great interest.

Concerning an increase in the number of bright spectroscopically-confirmed background sources, while we used a large compilation of spectroscopic catalogues, previous surveys have focused primarily on the central part of the COSMOS field, providing only 36 background objects with suitable spectroscopic redshifts for our IGM tomography. According to the Bouwens et al. (2021) UV luminosity function, there should be numerous star-forming galaxies brighter than $m_{\text{UV}} < 25.5 \text{ mag}$ ($M_{\text{UV}} \lesssim -21$) in the appropriate redshift range ($4.98 < z < 5.89$) with the surface density of $\Sigma_g \approx 726 \text{ deg}^{-2}$. This corresponds to a total of ≈ 1280 sources that can be in principle accessed across the HSC's 1.76 deg^2 field. Uncovering this population would provide a large boost in the number of background galaxies (cf. the surface density of our background LAEs of $\Sigma_{\text{LAE}} \approx 71.8 \text{ deg}^2$). The bright UV continua will ensure $\sim 2-3\sigma$ detection of the $\text{Ly}\alpha$ forest transmission with the current NB718 depth. There are a number of photometric catalogues with dropout selection (e.g. GOLDRUSH: Harikane et al. 2022) and photometric redshifts (COSMOS2020: Weaver et al. 2022) in the COSMOS field. We expect at least 10–20% of such UV continuum selected objects will show observable $\text{Ly}\alpha$ emission (Stark et al. 2010, 2011; Mallery et al. 2012; Cassata et al. 2015; Arrabal Haro et al. 2018; Kusakabe et al. 2020). A wide-field multi-object spectroscopic (MOS) follow-up campaign in the ultra-deep HSC footprint of the COSMOS field can locate $\approx 128 - 256$ background sources ($\Sigma_{\text{specz}} \approx 72.6 - 145.2 \text{ deg}^2$), providing $2-3\times$ increase in the total background sample (which is currently dominated by $z \approx 5.7$ LAEs with typical $\sim 26.4 \text{ mag}$ continua). As discussed in Section 5.3, including a subset of the background LAEs in the

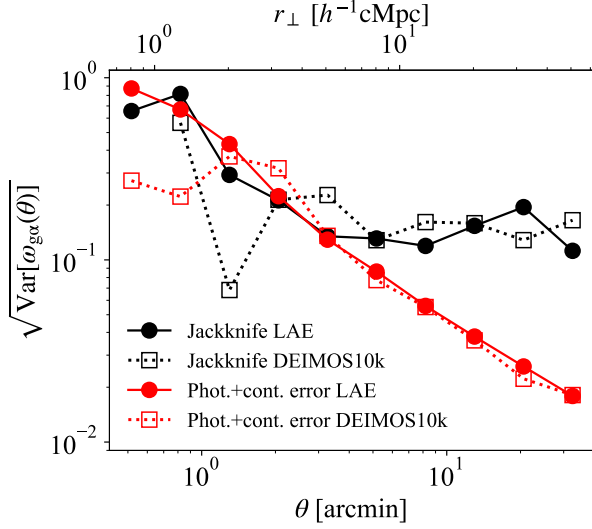


Figure 21. Comparison of the error budget in the cross-correlation estimated from the Jackknife covariance matrix (black) and analytic variance (red) for the LAE (filled circles) and DEIMOS10k (open squares) samples. The analytic variance includes the error from photometric noise and continuum slope uncertainty and follows the expected scaling $\sim ((\text{typical } T_{\text{IGM}} \text{ error}) / \langle T_{\text{IGM}} \rangle) / \sqrt{N_{\text{pairs}}(\theta)}$ as the number of pairs per bin increases.

spectroscopic follow-up campaign is also important as it allows us to statistically correct for the systematic bias in the correlation functions by low-redshift interlopers. A follow-up spectroscopic survey using wide-field MOS instruments such as Keck/DEIMOS and upcoming Subaru/Prime Focus Spectrograph (PFS) and VLT/MOONS is required to improve the significance and angular resolution of the photometric IGM tomography.

Alternatively, it should also be possible to uncover large numbers of bright star-forming galaxies with secure redshifts using rest-frame optical emission line such as $H\beta + [\text{O III}]$ and $H\alpha$ using a wide-field redshift survey with the NIRCcam wide-field slitless spectrograph (WFSS) on board *JWST*. Recently Sun et al. (2022b,a) (see also Matthee et al. 2022a) suggest that $\sim 88\%$ of bright $z \sim 6$ star-forming galaxies show strong $H\beta + [\text{O III}]$ and $H\alpha$ emission lines detectable with a shallow (~ 20 min) integration. If this holds true at $4.98 < z < 5.89$, the shallow wide-field NIRCcam/WFSS survey tiling the HSC COSMOS field can uncover a factor of $\sim 4 - 9$ larger number of background galaxies ($\Sigma_{\text{JWST}} \approx 638 \text{ deg}^{-2}$) than a ground-based wide-field MOS survey targeting $\text{Ly}\alpha$ lines.

9.1.2 Reducing the errors & systematics in the statistical analysis

For the statistical measurement of the correlation functions, there is patch-to-patch variance in addition to the photometric noise and UV continuum uncertainty. In Figure 21 we compare the Jackknife variance with our propagated errors from photometric noise and UV continuum uncertainty (Equation 23) from the Bayesian SED fitting framework in the angular mean $\text{Ly}\alpha$ forest transmission around LAEs $\langle T_{\text{IGM}}(\theta) \rangle$.

We find that the photometric noise is the dominant source of uncertainties at $\theta < 3 \text{ arcmin}$ ($< 5 h^{-1} \text{ cMpc}$). The Jackknife variance sometimes underestimates the error due to the small number of pairs in the inner angular bins. The photometric noise is comparable for both the LAE and DEIMOS10k samples because, while the individ-

ual error in T_{IGM} is larger in the LAE sample, the larger sample size reduces the error in the cross-correlation function.

At larger radii $\theta > 3 \text{ arcmin}$ ($> 5 h^{-1} \text{ cMpc}$), the Jackknife error becomes larger than the propagated error from photometric+continuum errors, indicating that the patch-to-patch variance in the field becomes the dominant uncertainty. While the photometric error is the dominant uncertainty in individual T_{IGM} , the error scale as $\propto \delta T_{\text{IGM}} / \sqrt{N_{\text{pairs}}(\theta)}$ where δT_{IGM} is the typical photometric uncertainty in individual T_{IGM} 's. The patch-to-patch variance may arise from cosmic variance or systematics such as imperfect sky subtraction or a coherent error in the photometric colours across the field. We believe that the current dominant source of the patch-to-patch variance is systematics. We checked this by comparing the Jackknife variances between the LAE and DEIMOS10k samples. As the former is sampled from the larger area, if the large-scale patch-to-patch variance is due to cosmic variance, we expect the Jackknife variance to decrease. However, this is not the case. We have also compared the Jackknife variance using a smaller number of angular bins to reduce the photometric error. The photometric error decreases as expected. In the both cases, the large-scale Jackknife variance remains roughly constant, suggesting that the observed excess of the Jackknife variance compared to the photometric+continuum error is likely caused by systematics such as sky background subtraction or possibly reflected lights within the optical units during the NB imaging. An unaccounted coherent change in the photometric colours (e.g. $\text{NB718} - z$ colour) across the field, e.g. due to imperfect PSF matching or Galactic dust extinction corrections, could also be a cause of the systematics. One possible way to quantify the systematics is to create an artificial IGM tomographic map using sky objects assuming constant $\text{Ly}\alpha$ forest transmission. While this artificial IGM map should be uniform across the field by construction, a coherent systematic change in the colours may introduce spatial fluctuations, which can be characterised by measuring the two-point correlation functions or similar statistics. More careful data reduction, background subtraction, and photometric calibration will be required to reduce the error in the large angular bins. We emphasise that these issues can be resolved with additional procedures during data reduction and calibration steps.

If the large-scale Jackknife variance is caused by the cosmic variance, we would benefit by enlarging the survey area to increase the sampling of the large-scale modes. Including other pointings such as the SXDS field would increase the constraining power on the large-scale correlation functions. Whether such an investment is worthwhile depends on the theoretically expected scale of the LAE- $\text{Ly}\alpha$ forest cross-correlation and $\text{Ly}\alpha$ forest auto-correlation functions (see companion paper, Kakiichi et al in prep). Further theoretical studies on photometric IGM tomography are necessary in order to understand the physical information contained on the different scales of the correlation functions.

9.1.3 Combining multi-wavelength data

Our present analysis employed only two broad-band filters (HSC z and y) to constrain the SEDs of background galaxies. The SED uncertainties (i.e continuum slope and SED template) can be mitigated by including multi-wavelength datasets available in the COSMOS field. Inclusion of near-infrared data such as UltraVISTA *JHK* (McCracken et al. 2012) will provide a longer baseline to better characterise the intrinsic SEDs. Constraining the dust and age of the background galaxies would eliminate the $\sim 6 - 27\%$ error in the measured T_{IGM} from SED fitting (Section 4.2). The soon-available

NIRCam imaging with F115W, F150W, F277W, F444W filters from COSMOS-Web (Kartaltepe et al. ID 1727, PI: 2021, see also Casey et al. 2022) can precisely determine the rest UV-to-optical SEDs of background galaxies.

9.2 Science applications

9.2.1 LyC escape fraction and the nature of ionizing sources

The statistical analysis of the LAE-Ly α forest angular cross-correlation from photometric IGM tomography provides a measure of the gas overdensity, temperature, and UV background fluctuations of the IGM around foreground LAEs. The angular cross-correlation is the NB-averaged version of the underlying galaxy-Ly α forest 3D cross-correlation that can be measured from the spectroscopic galaxy survey in quasar fields (Kakiichi et al. 2018; Meyer et al. 2020). We can thus adopt a similar approach to constrain the population-averaged LyC escape fraction from the LAE-Ly α forest cross-correlation. We will present the analysis in a companion paper (Kakiichi et al in prep).

The observed limit on the LAE-Ly α forest cross-correlation could also limit the contribution of bright LAEs to the UV background. Matthee et al. (2022b); Naidu et al. (2022) claimed that bright LAEs with $L_{\text{Ly}\alpha} \gtrsim 10^{42} \text{ erg s}^{-1}$ could contribute significantly to the total ionizing budget at $z \gtrsim 4$. Such bright LAEs may produce a large proximity zone which could be observed by photometric IGM tomography. The measurement of LAE-Ly α forest cross-correlation can be used to test this scenario. We will discuss the implication of our result on the relative contribution of bright and faint galaxies to the total ionizing budget in a following paper (Kakiichi et al in prep).

9.2.2 Search for ionized bubbles and protoclusters

Besides the statistical analysis, the IGM tomography can also be used to search for ionized bubbles and/or protoclusters at high redshifts. While our present analysis focused on NB718 filter ($z \approx 4.9$), one can perform tomographic analyses at any redshift where NB filters are available, including NB527 ($z \approx 3.31$) and NB816 ($z \approx 5.72$) (see Kakiichi et al. 2022, for the full list), assuming adequate deep NB imaging and background galaxies are available.

Ionized bubbles or the high UV background regions are expected to show transmissive Ly α forest with effective optical depth of $\tau_{\text{eff}} \approx 2$ at $z \approx 5.7$ (Davies et al. 2018; Keating et al. 2020). This corresponds to a NB-BB magnitude decrement of $m_{\text{NB816}} - m_{\text{UV}} = -2.5 \log_{10} e^{-\tau_{\text{eff}}} \approx 2.2 \text{ mag}$. With a NB816 depth of 27.7 mag, it is possible to search for such regions using 25.5 mag background sources at $5.81 < z < 6.86$.

The present $2 - 3\sigma$ depth of 26.89 – 27.33 mag of the NB816 image already allows us to search for extreme transmissive IGM regions with $\tau_{\text{eff}} \sim 1 - 2$ ($m_{\text{NB816}} - m_{\text{UV}} \approx 1.1 - 2.2 \text{ mag}$). One such region ($\tau_{\text{eff}} \approx 1.2$) has been already discovered serendipitously by Bosman et al. (2020) around a quasar using NB816 imaging. While current cosmological simulations (Davies et al. 2018; Keating et al. 2020) do not predict such extreme values, they are performed under the assumption that reionization is driven by a large population of relatively faint galaxies. Highly transmissive regions may exist if the contribution from quasars (Chardin et al. 2017) or luminous galaxies with high LyC escape fraction and ionizing photon production efficiency (e.g. Endsley et al. 2021; Topping et al. 2022; Marques-Chaves et al. 2022) are important.

Similarly, one can conduct a search for protoclusters as coherently opaque regions of the IGM in the tomographic map (Lee et al. 2016; Newman et al. 2020). Indeed, Mawatari et al. (2017) used the photometric IGM tomographic technique to examine the protocluster region with NB497 ($z \approx 3.1$). The $10 - 40 h^{-1} \text{ cMpc}$ scales of coherently strong Ly α absorption (Cai et al. 2016) is shown to be associated with an overdensity at $z \sim 2 - 3$ (Cai et al. 2017; Shi et al. 2021) whose scale matches closely with the line-of-sight width of a NB filter. LAEs in the tomographic slice permit an immediately confirmation of whether the coherently opaque IGM region is associated with a galaxy overdensity without a separate spectroscopic or dedicated imaging campaign. This provides an opportunity to test the relation between galaxy overdensities and the IGM environments (e.g. Momose et al. 2021; Newman et al. 2022) at higher redshifts without a need of extremely-deep spectroscopy of background galaxies.

9.2.3 Quasar light-echoes and past AGN activity

The relatively short variability timescale of a quasar $t_{\text{Q}} \sim 10^{6-8} \text{ yr}$ compared to the light crossing time of the IGM tomographic map means that the radiation from the quasar could leave an imprint on the ionization state of the IGM well after the non-thermal activity ceases (Adelberger 2004). Schmidt et al. (2019); Kakiichi et al. (2022) examined the prospect of examining the lifetime/past AGN activity of an active quasar using this light echo signal. Our photometric IGM tomography demonstrates that searching for quasar light echoes is possible if an appropriate quasar field is targeted.

In principle, the search for quasar light-echoes is not limited to the region around a luminous quasar, but can be applied to any region around a massive galaxy where quasar activity might have occurred during its previous $\sim 10^8 \text{ yr}$. Searches for fossil light-echoes around LAEs, or LBGs at redshift within the region of influence of the NB tomographic slice, could be used to constrain the past luminous ionizing activity of a source as a function of the travel time between the source and a position of the IGM. Bosman et al. (2020) used the photometric detection of the Ly α forest transmission in the NB816 filter located slightly foreground of $z \approx 5.8$ quasar to show that the quasar was active for at least $\sim 2 \times 10^4 \text{ yr}$ in the past. Bowler et al. (2015, 2020); Ono et al. (2018); Harikane et al. (2022) suggest that the double-power law luminosity function at $z \gtrsim 4$ may be a sign of inefficient quenching by AGN feedback acting on luminous galaxies. Finding a lack of highly transmissive regions around luminous galaxies would imply that these systems could not release a significant ionizing radiation by quasar activity during their history.

9.2.4 Correlation with direct LyC, He II, and C IV emitters

Multiple HSC NB imaging available in the COSMOS field provides a NB photometric measure of LyC leakage and the identification of strong He II and C IV emission lines for $z \approx 4.9$ LAEs via the CHORUS survey (Inoue et al. 2020). The correlation of these populations with photometric IGM tomographic map will have a number of applications.

For example, direct search for LyC emission along the line-of-sight of galaxies is increasingly difficult at $z \gtrsim 3.5$ as the IGM becomes increasingly opaque on average. Bassett et al. (2021, 2022) showed that any LyC detection could be biased towards the rare transmissive regions of the IGM and the bias introduced by assuming an average IGM transmission is severe at $3 < z < 5$. While LyC leaking candidates have reported at this redshift range (e.g. Shapley et al. 2016;

Vanzella et al. 2018; Ji et al. 2020; Meštrić et al. 2020; Prichard et al. 2022; Rivera-Thorsen et al. 2022; Marques-Chaves et al. 2022), the uncertain IGM transmission value against LyC photons make it difficult to convert the observed LyC flux to an absolute LyC escape fraction and sometimes results in unphysical values $f_{\text{esc}} > 1$. Fletcher et al. (2019) also noted the possibility of spatial variations in the inferred LyC escape fraction due to the uncertain IGM transmission. Photometric IGM tomography provides an useful independent measure of the IGM transmission and could be used to alleviate the issue of uncertain IGM LyC opacities along the lines-of-sight to high-redshift galaxies.

Combining the direct detection of LyC leakage and the indirect measurement from galaxy-Ly α forest cross-correlation, one can test the relative contributions of luminous and faint galaxies to the ionising budget. The former measures the ionising contribution of galaxies above the detection limit, while the latter provides a population-averaged LyC escape fraction including those below the detection limit. NB measurements of C IV and He II from LAEs are also valuable for examining hard ionising sources like AGN and X-ray binaries. The IGM tomographic map enables us to connect the properties of these populations with their large-scale IGM environments.

10 CONCLUSIONS

We present a novel technique called photometric IGM tomography to map the large-scale structure of the IGM at $z \sim 5$ in the COSMOS field. The technique utilizes ultra-deep NB718 imaging to detect the Ly α forest transmission along sightlines to various background galaxies including spectroscopically-confirmed DEIMOS10k sources and NB816-selected LAE catalogue from SILVERRUSH. In this paper, we describe a science verification of this new technique using public HSC data including HSC-SSP DR3 and CHORUS PDR1.

We have developed a Bayesian SED fitting framework to measure the Ly α forest transmission along background galaxies. This allows us to accurately propagate the error from photometric noise and uncertainty from the assumed SED template into the final measurement of the Ly α forest transmission. At the current imaging depths, photometric noise from NB718 imaging dominates the total error in the estimated Ly α forest transmission. Uncertainties from the UV continuum slopes and SED templates are subdominant in the present analysis.

Using a total sample of 140 background sources, we have photometrically measured the NB718-integrated mean Ly α forest transmission at $z \approx 4.9$. We find that our result is consistent with the previous measurement using quasar spectra, demonstrating that an accurate photometric NB measurement of the Ly α forest transmission is practical. We argue that the most likely systematic is contamination from low-redshift interlopers in the NB-selected LAE sample, which if not taken into account, would cause a fictitious Ly α forest transmission. This may explain an offset we see in the measured values between spectroscopically-confirmed background sources (DEIMOS10k) and NB-selected background sources (SILVERRUSH $z \approx 5.7$ LAEs). Fortunately, this can be corrected statistically provided that the low-redshift interloper fraction is known e.g. from spectroscopic follow-up of a subset of the population.

We developed a method for measuring the angular LAE-Ly α forest cross-correlation and the Ly α forest auto-correlation functions. Our method incorporates the individual posteriors from the Bayesian SED fitting framework to measure the angular correlation functions consistent with the propagated error including photometric noise and SED uncertainties. Low-redshift interlopers in our foreground

($z \approx 4.9$) and background ($z \approx 5.7$) LAE samples are again a main systematic uncertainty which can also be corrected statistically using a partial spectroscopic follow-up of the parent LAE samples. Applying the technique to the present data, we did not detect any angular LAE-Ly α forest cross-correlation and auto-correlation of the Ly α forest at $z \approx 4.9$. Our result is consistent with no Ly α forest fluctuations around LAEs and should be below 58% at $10 h^{-1} \text{cMpc}$ compared to the global mean transmission at $z \approx 4.9$. We will discuss the physical implications of this limit in a companion paper (Kakiichi et al in prep).

Finally, we presented a reconstructed 2D tomographic map of the IGM at $z \approx 4.9$, co-spatial with the distribution of foreground LAEs in the same cosmic volume in the COSMOS field. The map embraces a field $140 h^{-1} \text{cMpc}$ in diameter with a transverse spatial resolution of $\approx 11 h^{-1} \text{cMpc}$. While the current map is still dominated by photometric noise, it represents the most detailed tomographic map close to the end of cosmic reionization. The ability of photometric IGM tomography to map both the large-scale structures of galaxies and the IGM across a large region of the sky is extremely appealing, allowing us to apply the technique for many science cases including constraints on LyC escape fractions, the nature of ionizing sources through the UV background and thermal fluctuations of the IGM, and the searches for ionized bubbles, protoclusters, and quasar light-echoes.

Photometric IGM tomography can be embedded in traditional NB imaging and wide-field spectroscopic surveys and is applicable at all redshifts from $z \sim 2$ to 6 where NB filters are available, thus making it possible to examine the co-evolution of galaxies and the cosmic web during the first few billion years of cosmic history. We argue that the technique can be improved through extremely-deep NB imaging and large spectroscopic follow-up campaigns. Although this would require a large, but nonetheless practical, investment of telescope time with existing 8-10 m telescopes, we make the case that such an investment is worthwhile. Combining multi-wavelength dataset including the UltraVISTA JHK and *Spitzer*/IRAC imaging and soon available *JWST*/NIRCam imaging in the COSMOS field will enable us to better control the systematic uncertainty arising from the intrinsic SEDs of background galaxies. Future Subaru/PFS surveys will greatly increase the number of spectroscopically-confirmed background galaxies. Furthermore, a large-scale *JWST* spectroscopic survey tiling the COSMOS field would push the redshift frontier closer to the reionization epoch. The wide-field capability of the photometric IGM tomography is highly complementary to surveys based on using extremely-deep spectra with ELT/MOSAICS and TMT/WFOS as their small fields of view require interesting target regions to be pre-selected, e.g. from photometric IGM tomographic maps. As such, photometric IGM tomography has great potential to uncover the physics of galaxy-cosmic web connection in the early Universe in the coming decade.

ACKNOWLEDGEMENTS

KK thanks Harley Katz, Fred Davies, and K-G Lee for discussions and constructive comments. We thank the referee for carefully reading the manuscript. This project has received funding from the European Research Council (ERC) under the European Union's Horizon 2020 research and innovation programme (grant agreement No 885301). RSE acknowledges financial support from ERC Advanced Grant FP7/669253. RAM acknowledges support from the ERC Advanced Grant 740246 (Cosmic Gas). This work is supported by the World Premier International Research Center Initiative (WPI Initia-

tive), MEXT, Japan, as well as KAKENHI Grant-in-Aid for Scientific Research (A) (20H00180 and 21H04467) through the Japan Society for the Promotion of Science (JSPS). This work was supported by the joint research program of the Institute for Cosmic Ray Research (ICRR), University of Tokyo.

This paper is based on data collected at the Subaru Telescope and retrieved from the HSC data archive system, which is operated by the Subaru Telescope and Astronomy Data Center (ADC) at NAOJ. Data analysis was in part carried out with the cooperation of Center for Computational Astrophysics (CfCA), NAOJ. We are honored and grateful for the opportunity of observing the Universe from Maunakea, which has the cultural, historical and natural significance in Hawaii. The Hyper Suprime-Cam (HSC) collaboration includes the astronomical communities of Japan and Taiwan, and Princeton University. The HSC instrumentation and software were developed by the National Astronomical Observatory of Japan (NAOJ), the Kavli Institute for the Physics and Mathematics of the Universe (Kavli IPMU), the University of Tokyo, the High Energy Accelerator Research Organization (KEK), the Academia Sinica Institute for Astronomy and Astrophysics in Taiwan (ASIAA), and Princeton University. Funding was contributed by the FIRST program from the Japanese Cabinet Office, the Ministry of Education, Culture, Sports, Science and Technology (MEXT), the Japan Society for the Promotion of Science (JSPS), Japan Science and Technology Agency (JST), the Toray Science Foundation, NAOJ, Kavli IPMU, KEK, ASIAA, and Princeton University. This paper makes use of software developed for Vera C. Rubin Observatory. We thank the Rubin Observatory for making their code available as free software at <http://pipelines.lsst.io/>.

DATA AVAILABILITY

All the original HSC data and spectroscopic catalogues including DEIMOS10k are available online through the HSC-SSP website and NASA/IPAC IRSA (see the links listed in this paper). The SILVER-RUSH catalogue is available at <http://cos.icrr.u-tokyo.ac.jp/rush.html>. The table of the measured Ly α forest transmission along our background galaxies is available through online supplementary material.

REFERENCES

- Abbott T. M. C., et al., 2022, *Phys. Rev. D*, **105**, 023520
 Abdurashidova Z., et al., 2022, *ApJ*, **925**, 221
 Adelberger K. L., 2004, *ApJ*, **612**, 706
 Ahumada R., et al., 2020, *ApJS*, **249**, 3
 Aihara H., et al., 2019, *PASJ*, **71**, 114
 Aihara H., et al., 2022, *PASJ*, **74**, 247
 Alam S., et al., 2021, *Phys. Rev. D*, **103**, 083533
 Arrabal Haro P., et al., 2018, *MNRAS*, **478**, 3740
 Barnett R., Warren S. J., Becker G. D., Mortlock D. J., Hewett P. C., McMahon R. G., Simpson C., Venemans B. P., 2017, *A&A*, **601**, A16
 Bassett R., Ryan-Weber E. V., Cooke J., Meštrić U., Kakiichi K., Prichard L., Rafelski M., 2021, *MNRAS*, **502**, 108
 Bassett R., et al., 2022, *MNRAS*, **511**, 5730
 Becker G. D., Hewett P. C., Worseck G., Prochaska J. X., 2013, *MNRAS*, **430**, 2067
 Becker G. D., Bolton J. S., Lidz A., 2015a, *Publ. Astron. Soc. Australia*, **32**, e045
 Becker G. D., Bolton J. S., Madau P., Pettini M., Ryan-Weber E. V., Venemans B. P., 2015b, *MNRAS*, **447**, 3402

- Becker G. D., Davies F. B., Furlanetto S. R., Malkan M. A., Boera E., Douglass C., 2018, *ApJ*, **863**, 92
 Bosman S. E. I., Fan X., Jiang L., Reed S., Matsuoka Y., Becker G., Haehnelt M., 2018, *MNRAS*, **479**, 1055
 Bosman S. E. I., Kakiichi K., Meyer R. A., Gronke M., Laporte N., Ellis R. S., 2020, *ApJ*, **896**, 49
 Bosman S. E. I., et al., 2022, *MNRAS*, **514**, 55
 Bouwens R. J., et al., 2014, *ApJ*, **793**, 115
 Bouwens R. J., et al., 2021, *AJ*, **162**, 47
 Bowler R. A. A., et al., 2015, *MNRAS*, **452**, 1817
 Bowler R. A. A., Jarvis M. J., Dunlop J. S., McLure R. J., McLeod D. J., Adams N. J., Milvang-Jensen B., McCracken H. J., 2020, *MNRAS*, **493**, 2059
 Bowman J. D., Rogers A. E. E., Monsalve R. A., Mozdzen T. J., Mahesh N., 2018, *Nature*, **555**, 67
 Bradshaw E. J., et al., 2013, *MNRAS*, **433**, 194
 Cai Z., et al., 2016, *ApJ*, **833**, 135
 Cai Z., et al., 2017, *ApJ*, **839**, 131
 Calzetti D., 2001, *PASP*, **113**, 1449
 Casey C. M., et al., 2022, arXiv e-prints, p. arXiv:2211.07865
 Cassata P., et al., 2015, *A&A*, **573**, A24
 Chardin J., Haehnelt M. G., Aubert D., Puchwein E., 2015, *MNRAS*, **453**, 2943
 Chardin J., Puchwein E., Haehnelt M. G., 2017, *MNRAS*, **465**, 3429
 Chen Y., et al., 2020, *MNRAS*, **499**, 1721
 Christenson H. M., Becker G. D., Furlanetto S. R., Davies F. B., Malkan M. A., Zhu Y., Boera E., Trapp A., 2021, *ApJ*, **923**, 87
 Coil A. L., et al., 2011, *ApJ*, **741**, 8
 Colless M., et al., 2003, arXiv e-prints, pp astro-ph/0306581
 Cool R. J., et al., 2013, *ApJ*, **767**, 118
 Cooper M. C., et al., 2011, *ApJS*, **193**, 14
 Cooper M. C., et al., 2012, *MNRAS*, **419**, 3018
 Curtis-Lake E., et al., 2013, *MNRAS*, **429**, 302
 D'Aloisio A., McQuinn M., Trac H., 2015, *ApJ*, **813**, L38
 D'Aloisio A., McQuinn M., Davies F. B., Furlanetto S. R., 2018, *MNRAS*, **473**, 560
 Davies F. B., Furlanetto S. R., 2016, *MNRAS*, **460**, 1328
 Davies F. B., Becker G. D., Furlanetto S. R., 2018, *ApJ*, **860**, 155
 Davies R. L., et al., 2021, *ApJ*, **909**, 78
 Davis M., et al., 2003, in Guhathakurta P., ed., Society of Photo-Optical Instrumentation Engineers (SPIE) Conference Series Vol. 4834, Discoveries and Research Prospects from 6- to 10-Meter-Class Telescopes II. pp 161–172 (arXiv:astro-ph/0209419), doi:10.1117/12.457897
 Drinkwater M. J., et al., 2010, *MNRAS*, **401**, 1429
 Eilers A.-C., Davies F. B., Hennawi J. F., 2018, *ApJ*, **864**, 53
 Eldridge J. J., Stanway E. R., Xiao L., McClelland L. A. S., Taylor G., Ng M., Greis S. M. L., Bray J. C., 2017, *Publ. Astron. Soc. Australia*, **34**, e058
 Endsley R., Stark D. P., Charlot S., Chevallard J., Robertson B., Bouwens R. J., Stefanon M., 2021, *MNRAS*, **502**, 6044
 Farrow D. J., et al., 2021, *MNRAS*, **507**, 3187
 Ferland G. J., et al., 2017, *Rev. Mex. Astron. Astrofis.*, **53**, 385
 Fletcher T. J., Tang M., Robertson B. E., Nakajima K., Ellis R. S., Stark D. P., Inoue A., 2019, *ApJ*, **878**, 87
 Foreman-Mackey D., Hogg D. W., Lang D., Goodman J., 2013, *PASP*, **125**, 306
 Garaldi E., Gnedin N. Y., Madau P., 2019, *ApJ*, **876**, 31
 Garaldi E., Kannan R., Smith A., Springel V., Pakmor R., Vogelsberger M., Hernquist L., 2022, *MNRAS*, **512**, 4909
 Garilli B., et al., 2014, *A&A*, **562**, A23
 Gontcho A Gontcho S., Miralda-Escudé J., Busca N. G., 2014, *MNRAS*, **442**, 187
 Grasshorn Gebhardt H. S., et al., 2019, *ApJ*, **876**, 32
 Guaita L., et al., 2011, *ApJ*, **733**, 114
 Hagen A., et al., 2014, *ApJ*, **786**, 59
 Harikane Y., et al., 2022, *ApJS*, **259**, 20
 Hasinger G., et al., 2018, *ApJ*, **858**, 77
 Hayes M., 2019, *Saas-Fee Advanced Course*, **46**, 319
 Horowitz B., et al., 2021, arXiv e-prints, p. arXiv:2109.09660

- Inoue A. K., et al., 2020, *PASJ*, 72, 101
- Ishimoto R., et al., 2022, *MNRAS*, 515, 5914
- Japelj J., et al., 2019, *A&A*, 632, A94
- Ji Z., et al., 2020, *ApJ*, 888, 109
- Jones D. H., et al., 2004, *MNRAS*, 355, 747
- Jones D. H., et al., 2009, *MNRAS*, 399, 683
- Kakiichi K., et al., 2018, *MNRAS*, 479, 43
- Kakiichi K., Schmidt T., Hennawi J., 2022, *MNRAS*, 516, 582
- Kartaltepe J., et al., 2021, COSMOS-Webb: The Webb Cosmic Origins Survey, JWST Proposal. Cycle 1, ID. #1727
- Kashino D., et al., 2019, *ApJS*, 241, 10
- Kashino D., Lilly S. J., Shibuya T., Ouchi M., Kashikawa N., 2020, *ApJ*, 888, 6
- Kashino D., Lilly S. J., Matthee J., Eilers A.-C., Mackenzie R., Bordoloi R., Simcoe R. A., 2022, arXiv e-prints, p. arXiv:2211.08254
- Keating L. C., Puchwein E., Haehnelt M. G., 2018, *MNRAS*, 477, 5501
- Keating L. C., Weinberger L. H., Kulkarni G., Haehnelt M. G., Chardin J., Aubert D., 2020, *MNRAS*, 491, 1736
- Kulkarni G., Keating L. C., Haehnelt M. G., Bosman S. E. I., Puchwein E., Chardin J., Aubert D., 2019, *MNRAS*, 485, L24
- Kusakabe H., et al., 2020, *A&A*, 638, A12
- Kwan J., et al., 2017, *MNRAS*, 464, 4045
- Le Fèvre O., et al., 2013, *A&A*, 559, A14
- Lee K.-G., et al., 2014a, *ApJ*, 795, L12
- Lee K.-G., et al., 2014b, *ApJ*, 795, L12
- Lee K.-G., et al., 2016, *ApJ*, 817, 160
- Lee K.-G., et al., 2018, *ApJS*, 237, 31
- Lemos P., et al., 2021, *MNRAS*, 505, 6179
- Lilly S. J., et al., 2009, *ApJS*, 184, 218
- Lilly S. J., Bordoloi R., Kashino D., Simcoe R. A., 2017, Exploring the End of Cosmic Reionization, JWST Proposal. Cycle 1, ID. #1243
- Liske J., et al., 2015, *MNRAS*, 452, 2087
- Mallery R. P., et al., 2012, *ApJ*, 760, 128
- Marques-Chaves R., et al., 2022, *MNRAS*, 517, 2972
- Masters D. C., Stern D. K., Cohen J. G., Capak P. L., Rhodes J. D., Castander F. J., Paltani S., 2017, *ApJ*, 841, 111
- Masters D. C., et al., 2019, *ApJ*, 877, 81
- Matthee J., Mackenzie R., Simcoe R. A., Kashino D., Lilly S. J., Bordoloi R., Eilers A.-C., 2022a, arXiv e-prints, p. arXiv:2211.08255
- Matthee J., et al., 2022b, *MNRAS*, 512, 5960
- Mawatari K., et al., 2017, *MNRAS*, 467, 3951
- McCracken H. J., et al., 2012, *A&A*, 544, A156
- McLure R. J., et al., 2013, *MNRAS*, 428, 1088
- McQuinn M., 2016, *ARA&A*, 54, 313
- Meiksin A., 2020, *MNRAS*, 491, 4884
- Meiksin A., McQuinn M., 2019, *MNRAS*, 482, 4777
- Mertens F. G., et al., 2020, *MNRAS*, 493, 1662
- Meštrić U., et al., 2020, *MNRAS*, 494, 4986
- Meyer R. A., Bosman S. E. I., Kakiichi K., Ellis R. S., 2019, *MNRAS*, 483, 19
- Meyer R. A., et al., 2020, *MNRAS*, 494, 1560
- Momcheva I. G., et al., 2016, *ApJS*, 225, 27
- Momose R., et al., 2021, *ApJ*, 909, 117
- Monzon J. S., Prochaska J. X., Lee K.-G., Chisholm J., 2020, *AJ*, 160, 37
- Nagamine K., et al., 2021, *ApJ*, 914, 66
- Naidu R. P., et al., 2022, *MNRAS*, 510, 4582
- Nakajima K., et al., 2012, *ApJ*, 745, 12
- Nasir F., D'Aloisio A., 2020, *MNRAS*, 494, 3080
- Newman J. A., et al., 2013, *ApJS*, 208, 5
- Newman A. B., et al., 2020, *ApJ*, 891, 147
- Newman A. B., et al., 2022, *Nature*, 606, 475
- Norberg P., Baugh C. M., Gaztañaga E., Croton D. J., 2009, *MNRAS*, 396, 19
- Oke J. B., Gunn J. E., 1983, *ApJ*, 266, 713
- Ono Y., Ouchi M., Shimasaku K., Dunlop J., Farrah D., McLure R., Okamura S., 2010, *ApJ*, 724, 1524
- Ono Y., et al., 2018, *PASJ*, 70, S10
- Ono Y., et al., 2021, *ApJ*, 911, 78
- Ouchi M., et al., 2008, *ApJS*, 176, 301
- Ouchi M., Ono Y., Shibuya T., 2020, *ARA&A*, 58, 617
- Pâris I., et al., 2018, *A&A*, 613, A51
- Pentericci L., et al., 2018, *A&A*, 616, A174
- Planck Collaboration et al., 2016, *A&A*, 594, A13
- Planck Collaboration et al., 2020, *A&A*, 641, A6
- Pontzen A., 2014, *Phys. Rev. D*, 89, 083010
- Prichard L. J., et al., 2022, *ApJ*, 924, 14
- Pritchard J. R., Loeb A., 2012, *Reports on Progress in Physics*, 75, 086901
- Reddy N. A., Steidel C. C., Pettini M., Bogosavljević M., 2016, *ApJ*, 828, 107
- Reddy N. A., Topping M. W., Sanders R. L., Shapley A. E., Brammer G., 2023, arXiv e-prints, p. arXiv:2303.11397
- Rivera-Thorsen T. E., Hayes M., Melinder J., 2022, *A&A*, 666, A145
- Rudie G. C., Steidel C. C., Shapley A. E., Pettini M., 2013, *ApJ*, 769, 146
- Schmidt T. M., Hennawi J. F., Lee K.-G., Lukić Z., Oñorbe J., White M., 2019, *ApJ*, 882, 165
- Shapley A. E., Steidel C. C., Strom A. L., Bogosavljević M., Reddy N. A., Siana B., Mostardi R. E., Rudie G. C., 2016, *ApJ*, 826, L24
- Shi D. D., Cai Z., Fan X., Zheng X. Z., Huang Y.-H., Xu J., 2021, *ApJ*, 915, 32
- Shibuya T., et al., 2018, *PASJ*, 70, S15
- Silverman J. D., et al., 2015, *ApJS*, 220, 12
- Singh S., et al., 2022, *Nature Astronomy*,
- Sivia D., Skilling J., 2006, *Data Analysis: A Bayesian Tutorial*. Oxford science publications, OUP Oxford, <https://books.google.co.jp/books?id=LYMSDAAAQBAJ>
- Skelton R. E., et al., 2014, *ApJS*, 214, 24
- Slosar A., et al., 2011, *J. Cosmology Astropart. Phys.*, 2011, 001
- Sobral D., Santos S., Matthee J., Pablino-Afonso A., Ribeiro B., Calhau J., Khostovan A. A., 2018, *MNRAS*, 476, 4725
- Stanway E. R., Eldridge J. J., 2018, *MNRAS*, 479, 75
- Stark D. P., Ellis R. S., Bunker A., Bundy K., Targett T., Benson A., Lacy M., 2009, *ApJ*, 697, 1493
- Stark D. P., Ellis R. S., Chiu K., Ouchi M., Bunker A., 2010, *MNRAS*, 408, 1628
- Stark D. P., Ellis R. S., Ouchi M., 2011, *ApJ*, 728, L2
- Steidel C. C., Erb D. K., Shapley A. E., Pettini M., Reddy N., Bogosavljević M., Rudie G. C., Rakic O., 2010, *ApJ*, 717, 289
- Steidel C. C., et al., 2014, *ApJ*, 795, 165
- Stratman C. M. S., et al., 2018, *ApJS*, 239, 27
- Sun F., et al., 2022a, arXiv e-prints, p. arXiv:2209.03374
- Sun F., et al., 2022b, *ApJ*, 936, L8
- Thomas R., et al., 2017, *A&A*, 597, A88
- Thomas R., et al., 2020, *A&A*, 634, A110
- Thomas R., et al., 2021, *A&A*, 650, A63
- Topping M. W., Stark D. P., Endsley R., Plat A., Whittler L., Chen Z., Charlott S., 2022, arXiv e-prints, p. arXiv:2208.01610
- Trainor R. F., Strom A. L., Steidel C. C., Rudie G. C., 2016, *ApJ*, 832, 171
- Trott C. M., et al., 2020, *MNRAS*, 493, 4711
- Turner M. L., Schaye J., Crain R. A., Rudie G., Steidel C. C., Strom A., Theuns T., 2017, *MNRAS*, 471, 690
- Vanzella E., et al., 2018, *MNRAS*, 476, L15
- Wang F., et al., 2021, A SPectroscopic survey of biased halos In the Reionization Era (ASPIRE): A JWST Quasar Legacy Survey, JWST Proposal. Cycle 1, ID. #2078
- Weaver J. R., et al., 2022, *ApJS*, 258, 11
- Wolfson M., Hennawi J. F., Davies F. B., Oñorbe J., 2022, arXiv e-prints, p. arXiv:2208.09013
- Yang J., et al., 2020, *ApJ*, 904, 26
- Zhu Y., et al., 2022, *ApJ*, 932, 76
- de Barros S., Schaerer D., Stark D. P., 2014, *A&A*, 563, A81

APPENDIX A: DERIVATION OF THE ESTIMATORS

For clarity, we explicitly show the derivation of the estimator and the variance. The mean Ly α forest transmission is given

by $\langle \bar{T}_{\text{IGM}} \rangle = \int \bar{T}_{\text{IGM}} P(\bar{T}_{\text{IGM}} | \{f_{\text{NB},i}^{\text{obs}}, f_{\text{BB},i}^{\text{obs}}\}_{i=1,\dots,N_{\text{bg}}}) d\bar{T}_{\text{IGM}}$. By substituting Equation 13, we have

$$\begin{aligned} \langle \bar{T}_{\text{IGM}} \rangle &= \int \left(\frac{1}{N_{\text{bg}}} \sum_{i=1}^{N_{\text{bg}}} T_{\text{IGM},i} \right) P(T_{\text{IGM},i} | f_{\text{NB},i}^{\text{obs}}, f_{\text{BB},i}^{\text{obs}}) \prod_{i=1}^{N_{\text{bg}}} dT_{\text{IGM},i}, \\ &= \frac{1}{N_{\text{bg}}} \sum_{i=1}^{N_{\text{bg}}} \int T_{\text{IGM},i} P(T_{\text{IGM},i} | f_{\text{NB},i}^{\text{obs}}, f_{\text{BB},i}^{\text{obs}}) dT_{\text{IGM},i}, \\ &= \frac{1}{N_{\text{bg}}} \sum_{i=1}^{N_{\text{bg}}} \langle T_{\text{IGM},i} \rangle. \end{aligned}$$

The variance of the mean Ly α forest transmission is given by $\sigma_{\bar{T}_{\text{IGM}}}^2 = \int (\bar{T}_{\text{IGM}} - \langle \bar{T}_{\text{IGM}} \rangle)^2 P(\bar{T}_{\text{IGM}} | \{f_{\text{NB},i}^{\text{obs}}, f_{\text{BB},i}^{\text{obs}}\}_{i=1,\dots,N_{\text{bg}}}) d\bar{T}_{\text{IGM}}$. By substituting Equation 13 and the above result, we have

$$\begin{aligned} \sigma_{\bar{T}_{\text{IGM}}}^2 &= \\ &= \frac{1}{N_{\text{bg}}^2} \int \left(\sum_{i=1}^{N_{\text{bg}}} T_{\text{IGM},i} - \langle T_{\text{IGM},i} \rangle \right)^2 P(T_{\text{IGM},i} | f_{\text{NB},i}^{\text{obs}}, f_{\text{BB},i}^{\text{obs}}) \prod_{i=1}^{N_{\text{bg}}} dT_{\text{IGM},i}. \end{aligned}$$

Since all the individual measurements are uncorrelated, the cross-terms are zeros, i.e.

$$\begin{aligned} \iint dT_{\text{IGM},i} dT_{\text{IGM},j} P(T_{\text{IGM},i} | f_{\text{NB},i}^{\text{obs}}, f_{\text{BB},i}^{\text{obs}}) P(T_{\text{IGM},j} | f_{\text{NB},j}^{\text{obs}}, f_{\text{BB},j}^{\text{obs}}) \\ \times (T_{\text{IGM},i} - \langle T_{\text{IGM},i} \rangle) (T_{\text{IGM},j} - \langle T_{\text{IGM},j} \rangle) = 0. \end{aligned}$$

Thus, we obtain

$$\begin{aligned} \sigma_{\bar{T}_{\text{IGM}}}^2 &= \\ &= \frac{1}{N_{\text{bg}}^2} \sum_{i=1}^{N_{\text{bg}}} \int (T_{\text{IGM},i} - \langle T_{\text{IGM},i} \rangle)^2 P(T_{\text{IGM},i} | f_{\text{NB},i}^{\text{obs}}, f_{\text{BB},i}^{\text{obs}}) dT_{\text{IGM},i}. \end{aligned}$$

The calculation for the angular averaged Ly α forest transmission around galaxies is identical to that for the estimator and the variance for the mean Ly α forest transmission for each angular bin. Thus, we obtain the estimator,

$$\begin{aligned} \langle \bar{T}_{\text{IGM}}(\theta) \rangle &= \frac{1}{N_{\text{pair}}(\theta)} \sum_{j=1}^{N_{\text{fg}}} \sum_{i=1}^{N_{\text{bg}}} \mathcal{I}(|\theta - \theta_{ij}|) \\ &\times \int T_{\text{IGM},i} P(T_{\text{IGM},i} | f_{\text{NB},i}^{\text{obs}}, f_{\text{BB},i}^{\text{obs}}) dT_{\text{IGM},i}, \end{aligned}$$

and the variance,

$$\begin{aligned} \text{Var}[\bar{T}_{\text{IGM}}(\theta)] &= \frac{1}{N_{\text{pair}}(\theta)^2} \sum_{j=1}^{N_{\text{fg}}} \sum_{i=1}^{N_{\text{bg}}} \mathcal{I}(|\theta - \theta_{ij}|) \\ &\times \int (T_{\text{IGM},i} - \langle T_{\text{IGM},i} \rangle)^2 P(T_{\text{IGM},i} | f_{\text{NB},i}^{\text{obs}}, f_{\text{BB},i}^{\text{obs}}) dT_{\text{IGM},i}. \end{aligned}$$

APPENDIX B: INTERLOPER CONTAMINATION

The contamination by low-redshift interlopers in the foreground and background LAE samples dilutes the angular mean Ly α forest transmission around the foreground LAEs. When we measure the observed angular mean Ly α forest transmission around LAEs using the background LAE sample, one can be decomposed the estimator into four different contributions: (1) true foreground LAE - true

background LAE pairs (fg-bg pairs), (2) foreground interloper - true background LAE pairs (fg.int-bg pairs), (3) true foreground LAE - background interloper pairs (fg-bg.int pairs), and (4) foreground interloper - background interloper pairs (fg.int-bg.int pairs),

$$\begin{aligned} \bar{T}_{\text{IGM}}(\theta) &= \frac{1}{N_{\text{pair}}(\theta)} \times \\ &\left(\sum_{\text{fg-bg pairs}}^{N_{\text{fg}}^{\text{true}} N_{\text{bg}}^{\text{true}}} T_{\text{IGM},i}^{\text{true}} \mathcal{I}(|\theta - \theta_{ij}|) + \sum_{\text{fg.int-bg pairs}}^{N_{\text{fg}}^{\text{int}} N_{\text{bg}}^{\text{true}}} T_{\text{IGM},i}^{\text{true}} \mathcal{I}(|\theta - \theta_{ij}|) \right. \\ &\left. + \sum_{\text{fg-bg.int pairs}}^{N_{\text{fg}}^{\text{true}} N_{\text{bg}}^{\text{int}}} T_{\text{IGM},i}^{\text{bg.int}} \mathcal{I}(|\theta - \theta_{ij}|) + \sum_{\text{fg.int-bg.int pairs}}^{N_{\text{fg}}^{\text{int}} N_{\text{bg}}^{\text{int}}} T_{\text{IGM},i}^{\text{bg.int}} \mathcal{I}(|\theta - \theta_{ij}|) \right), \end{aligned}$$

where $N_{\text{pairs}}(\theta)$ is the total number of observed pairs per angular bin $N_{\text{pairs}}(\theta)$, $N_{\text{fg}}^{\text{true}}$ and $N_{\text{bg}}^{\text{true}}$ are the number of true foreground and background LAEs, $N_{\text{fg}}^{\text{int}}$ and $N_{\text{bg}}^{\text{int}}$ are the number of foreground and background interlopers, and $T_{\text{IGM},i}^{\text{true}}$ and $T_{\text{IGM},i}^{\text{bg.int}}$ are the measured Ly α forest transmission along true background LAEs and background LAE interlopers. The expectation value of the angular averaged Ly α forest transmission around LAEs is then given by

$$\begin{aligned} \langle \bar{T}_{\text{IGM}}(\theta) \rangle &= \frac{N_{\text{pair}}^{\text{fg-bg pairs}}(\theta)}{N_{\text{pair}}(\theta)} \langle \bar{T}_{\text{IGM}}^{\text{true}}(\theta) \rangle \\ &+ \frac{N_{\text{pair}}^{\text{fg.int-bg pairs}}(\theta)}{N_{\text{pair}}(\theta)} \langle \bar{T}_{\text{IGM}}^{\text{true}}(\theta) \rangle_{\text{fg.int-bg pairs}} \\ &+ \frac{N_{\text{pair}}^{\text{fg-bg.int pairs}}(\theta)}{N_{\text{pair}}(\theta)} \langle \bar{T}_{\text{IGM}}^{\text{bg.int}}(\theta) \rangle_{\text{fg-bg.int pairs}} \\ &+ \frac{N_{\text{pair}}^{\text{fg.int-bg.int pairs}}(\theta)}{N_{\text{pair}}(\theta)} \langle \bar{T}_{\text{IGM}}^{\text{bg.int}}(\theta) \rangle_{\text{fg.int-bg.int pairs}}, \end{aligned}$$

where $N_{\text{pairs}}^{\text{fg-bg pairs}}(\theta)$, $N_{\text{pairs}}^{\text{fg.int-bg pairs}}(\theta)$, $N_{\text{pairs}}^{\text{fg-bg.int pairs}}(\theta)$, and $N_{\text{pairs}}^{\text{fg.int-bg.int pairs}}(\theta)$ are the number of foreground LAE - background LAE, foreground LAE interloper - background LAE, foreground LAE - background LAE interloper, and foreground LAE interloper - background LAE interloper pairs per bin.

The total number of observed pairs depends on the angular correlation function of the observed populations $\omega_{\text{fg-bg}}^{\text{obs}}(\theta)$,

$$N_{\text{pairs}}(\theta) = N_{\text{fg}}^{\text{obs}} N_{\text{bg}}^{\text{obs}} \left[1 + \omega_{\text{fg-bg}}^{\text{obs}}(\theta) \right] \frac{2\pi\theta\Delta\theta}{\Omega_{\text{survey}}} \quad (\text{B1})$$

where Ω_{survey} is the survey area and $\Delta\theta$ is the angular bin size. Similarly, the number of true foreground LAE - true background LAE pairs etc is given by

$$N_{\text{pairs}}^{\text{fg-bg pairs}}(\theta) = N_{\text{fg}}^{\text{true}} N_{\text{bg}}^{\text{true}} \left[1 + \omega_{\text{fg-bg}}(\theta) \right] \frac{2\pi\theta\Delta\theta}{\Omega_{\text{survey}}},$$

$$N_{\text{pairs}}^{\text{fg.int-bg pairs}}(\theta) = N_{\text{fg}}^{\text{int}} N_{\text{bg}}^{\text{true}} \left[1 + \omega_{\text{fg.int-bg}}(\theta) \right] \frac{2\pi\theta\Delta\theta}{\Omega_{\text{survey}}},$$

$$N_{\text{pairs}}^{\text{fg-bg.int pairs}}(\theta) = N_{\text{fg}}^{\text{true}} N_{\text{bg}}^{\text{int}} \left[1 + \omega_{\text{fg-bg.int}}(\theta) \right] \frac{2\pi\theta\Delta\theta}{\Omega_{\text{survey}}},$$

$$N_{\text{pairs}}^{\text{fg.int-bg.int pairs}}(\theta) = N_{\text{fg}}^{\text{int}} N_{\text{bg}}^{\text{int}} \left[1 + \omega_{\text{fg.int-bg.int}}(\theta) \right] \frac{2\pi\theta\Delta\theta}{\Omega_{\text{survey}}}.$$

Defining the interloper fractions in the observed foreground and background LAE samples to be $f_{\text{fg.int}} = N_{\text{fg.int}}^{\text{int}}/N_{\text{fg}}^{\text{obs}}$ and $f_{\text{bg.int}} = N_{\text{bg.int}}^{\text{int}}/N_{\text{bg}}^{\text{obs}}$, the observed angular cross-correlation function from the

Table A1. The measured values of the Ly α forest transmission T_{IGM} , UV magnitude M_{UV} , and the UV continuum slope β from the Bayesian SED fitting framework for all background sources in the DEIMOS10k sample. The full table and the machine-readable file are available as online supplementary material.

ID	Object name	RA	DEC	T_{IGM}	M_{UV}	β	Note
1	DEIMOS_2018_L222036	10h02m20.94s	+01d37m06.78s	+0.049 ^{+0.196} _{-0.194}	-20.65 ^{+0.13} _{-0.13}	-1.56 ^{+1.65} _{-1.67}	
2	DEIMOS_2018_L416105	10h02m45.66s	+01d55m35.91s	+0.531 ^{+0.178} _{-0.176}	-21.18 ^{+0.12} _{-0.12}	-1.31 ^{+1.07} _{-1.07}	
3	DEIMOS_2018_L244697	10h01m59.64s	+01d39m16.81s	+0.045 ^{+0.064} _{-0.064}	-21.45 ^{+0.02} _{-0.02}	-1.69 ^{+0.70} _{-0.70}	
4	DEIMOS_2018_L254463	10h01m58.94s	+01d40m10.67s	+0.098 ^{+0.049} _{-0.049}	-21.62 ^{+0.02} _{-0.02}	-2.39 ^{+0.63} _{-0.62}	
5	DEIMOS_2018_L263567	10h01m35.65s	+01d41m08.01s	+0.070 ^{+0.139} _{-0.141}	-20.44 ^{+0.18} _{-0.18}	-2.81 ^{+1.47} _{-1.49}	
...							

Table A2. Same as Table A1 but for background sources ($z \approx 5.7$ LAEs) from the SILVERRUSH catalogue. The full table and the machine-readable file are available as online supplementary material.

ID	Object name	RA (J2000)	DEC (J2000)	T_{IGM}	M_{UV}	β	Note
1	SILVERRUSH_2021_15477	09h57m44.50s	+02d16m39.00s	+3.004 ^{+1.693} _{-1.648}	-19.47 ^{+0.32} _{-0.32}	-1.33 ^{+2.27} _{-2.32}	94% outlier
2	SILVERRUSH_2021_14862	09h57m47.81s	+02d16m18.53s	+1.431 ^{+0.826} _{-0.798}	-19.83 ^{+0.27} _{-0.27}	-0.87 ^{+2.04} _{-2.16}	63% outlier
3	SILVERRUSH_2021_5731	09h57m48.23s	+02d33m56.51s	+0.604 ^{+0.401} _{-0.381}	-20.57 ^{+0.27} _{-0.26}	-1.15 ^{+2.12} _{-2.17}	
4	SILVERRUSH_2021_18336	09h58m11.04s	+02d45m50.21s	+0.621 ^{+0.481} _{-0.455}	-20.33 ^{+0.28} _{-0.28}	-1.47 ^{+2.21} _{-2.22}	
5	SILVERRUSH_2021_10859	09h58m12.36s	+02d03m09.22s	-0.006 ^{+0.352} _{-0.368}	-20.31 ^{+0.22} _{-0.21}	-0.52 ^{+1.77} _{-1.82}	
...							

foreground and background LAE samples has contributions from true and interloper pairs,

$$\begin{aligned} \omega_{\text{fg-bg}}^{\text{obs}}(\theta) &= (1 - f_{\text{fg.int}})(1 - f_{\text{bg.int}})\omega_{\text{fg-bg}}(\theta) \\ &+ f_{\text{fg.int}}(1 - f_{\text{bg.int}})\omega_{\text{fg.int-bg}}(\theta) \\ &+ f_{\text{bg.int}}(1 - f_{\text{fg.int}})\omega_{\text{fg-bg.int}}(\theta) \\ &+ f_{\text{bg.int}}f_{\text{fg.int}}\omega_{\text{fg.int-bg.int}}(\theta). \end{aligned}$$

Since both interlopers and true LAEs in the two different NB-selected samples are located at different redshifts, all these angular correlation functions should be zeros. The objects in foreground and background LAE samples are not spatially correlated to each other. Thus,

$$\begin{aligned} \langle \bar{T}_{\text{IGM}}^{\text{true}}(\theta) \rangle_{\text{fg.int-bg pairs}} &\approx \langle \bar{T}_{\text{IGM}}^{\text{true}} \rangle, \\ \langle \bar{T}_{\text{IGM}}^{\text{bg.int}}(\theta) \rangle_{\text{fg-bg.int pairs}} &\approx \langle \bar{T}_{\text{IGM}}^{\text{bg.int}} \rangle, \\ \langle \bar{T}_{\text{IGM}}^{\text{bg.int}}(\theta) \rangle_{\text{fg.int-bg.int pairs}} &\approx \langle \bar{T}_{\text{IGM}}^{\text{bg.int}} \rangle. \end{aligned}$$

Therefore, we find that the effect of the interloper contamination on the observed angular-averaged Ly α forest transmission profile is expressed as

$$\begin{aligned} \langle \bar{T}_{\text{IGM}}(\theta) \rangle &= (1 - f_{\text{fg.int}})(1 - f_{\text{bg.int}})\langle \bar{T}_{\text{IGM}}^{\text{true}}(\theta) \rangle \\ &+ f_{\text{fg.int}}(1 - f_{\text{bg.int}})\langle \bar{T}_{\text{IGM}}^{\text{true}} \rangle \\ &+ f_{\text{bg.int}}(1 - f_{\text{fg.int}})\langle \bar{T}_{\text{IGM}}^{\text{bg.int}} \rangle \\ &+ f_{\text{fg.int}}f_{\text{bg.int}}\langle \bar{T}_{\text{IGM}}^{\text{bg.int}} \rangle. \end{aligned}$$

As the observed LAE-Ly α forest cross-correlation is defined with respect to the observed mean Ly α forest transmission which is also affected by the interlopers in the background LAE sample, i.e.

$$\langle \bar{T}_{\text{IGM}}^{\text{obs}} \rangle = (1 - f_{\text{bg.int}})\langle \bar{T}_{\text{IGM}}^{\text{true}} \rangle + f_{\text{bg.int}}\langle \bar{T}_{\text{IGM}}^{\text{bg.int}} \rangle, \text{ we find}$$

$$\begin{aligned} \omega_{\text{g}\alpha}^{\text{obs}}(\theta) &= \frac{\langle \bar{T}_{\text{IGM}}^{\text{obs}}(\theta) \rangle}{\langle \bar{T}_{\text{IGM}}^{\text{obs}} \rangle} - 1 \\ &= \frac{(1 - f_{\text{fg.int}})(1 - f_{\text{bg.int}})\langle \bar{T}_{\text{IGM}}^{\text{true}} \rangle}{(1 - f_{\text{bg.int}})\langle \bar{T}_{\text{IGM}}^{\text{true}} \rangle + f_{\text{bg.int}}\langle \bar{T}_{\text{IGM}}^{\text{bg.int}} \rangle} \omega_{\text{g}\alpha}^{\text{true}}(\theta), \end{aligned}$$

$$\text{where } \omega_{\text{g}\alpha}^{\text{true}}(\theta) = \langle \bar{T}_{\text{IGM}}^{\text{true}}(\theta) \rangle / \langle \bar{T}_{\text{IGM}}^{\text{true}} \rangle - 1.$$

APPENDIX C: TABLES OF THE ESTIMATED LY α FOREST TRANSMISSION

Tables A1 and A2 show the first five rows of the tables of the estimated Ly α forest transmission along all the background sources used in this paper. The full tables and the corresponding machine-readable files are available as online supplementary materials.