1	Local seismicity around the Chain Transform Fault at the
2	Mid-Atlantic Ridge from OBS observations
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19 Summary

20 Seismicity along transform faults provides important constraints for our understanding of the 21 factors that control earthquake ruptures. Oceanic transform faults are particularly informative 22 due to their relatively simple structure in comparison to their continental counterparts. The seismicity of several fast-moving transform faults has been investigated by local networks, 23 24 but as of today there been few studies of transform faults in slow spreading ridges. Here we 25 present the first local seismicity catalogue based on event data recorded by a temporary 26 broadband network of 39 ocean bottom seismometers located around the slow-moving Chain 27 Transform Fault (CTF) along the Mid-Atlantic Ridge (MAR) from March 2016 to March 2017. We locate 972 events in the area by simultaneously inverting for a 1-D velocity model 28 informed by the event P- and S-arrival times. We refine the depths and focal mechanisms of 29 30 the larger events using deviatoric moment tensor inversion. Most of the earthquakes are located along the CTF (700) and Romanche transform fault (94) and the MAR (155); a 31 32 smaller number (23) can be observed on the continuing fracture zones or in intraplate locations. The ridge events are characterised by normal faulting and most of the transform 33 events are characterised by strike slip faulting, but with several reverse mechanisms that are 34 35 likely related to transpressional stresses in the region. CTF events range in magnitude from 1.1 to 5.6 with a magnitude of completeness around 2.3. Along the CTF we calculate a b-36 value of 0.81 ± 0.09 . The event depths are mostly shallower than 15 km below sea level 37 38 (523), but a small number of high-quality earthquakes (16) are located deeper, with some (8) located deeper than the brittle-ductile transition as predicted by the 600°C-isotherm from a 39 40 simple thermal model. The deeper events could be explained by the control of seawater infiltration on the brittle failure limit. 41

42 Keywords: Seismicity, Atlantic Ocean, Mid-ocean ridge processes, Oceanic transform and
43 fracture zone processes, Seismicity and tectonics.

45 Introduction

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47 The factors that dictate the location, size and style of earthquake faulting are fundamental to our understanding of hazard and hazard mitigation (e.g., Slemmons et al., 1986). Oceanic 48 49 transform faults (OTF), such as those that connect the many fragmented sections of the slow 50 spreading Mid-Atlantic Ridge, are an ideal place to investigate the earthquake cycle. While 51 not a great risk to humans themselves due to their often remote locations, OTFs are thought 52 to be relatively simple in terms of thermal structure, fault zone geometry, slip rate and 53 rheology in comparison to more hazardous continental counterparts such as the San Andreas 54 fault (Behn et al., 2007). Therefore, the study of earthquakes associated with OTFs provides 55 important constraints to inform hazard assessments in continental areas. Ocean lithosphere is 56 characterised by relatively homogenous composition, with mafic to ultra-mafic lithologies. In addition, deformation tends to be localised in a narrow zone, roughly 20 to 30 km wide (Fox 57 58 & Gallo 1984) with nearly vertically oriented faults on which total motions should roughly correspond to ridge spreading rates. Finally, the region of potential seismic slip is also well-59 predicted by the region shallower than the 600°C isotherm predicted by simple thermal 60 models such as the halfspace cooling model (e.g., Abercrombie & Ekström, 2001). 61

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The vast majority of earthquakes along OTFs cannot be detected, due to their general inaccessibility. Being in the middle of the ocean, the closest landmasses and, thus, seismic land stations are often located at teleseismic distances. Previous studies of Atlantic transform earthquakes have used teleseismic data (e.g., Engeln et al., 1986; Bergman & Solomon, 1988; Abercrombie & Ekström, 2001). Smaller events are not recorded teleseismically, and the magnitude of completeness offered by global monitoring networks is high, around 4.4 (Parnell-Turner et al., 2022). This affects the observation of microseismicity around the

brittle-ductile transition. Furthermore, the use of a non-local velocity model can influence theaccuracy of hypocentre depth localisations.

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Over the last decade the temporary deployment of dense ocean bottom seismometer (OBS) 73 74 networks has enabled more detailed study of local seismicity along Pacific OTFs, e.g., along the Gofar and Discovery OTFs (McGuirre et al., 2012) and the Blanco OTF (Kuna et al., 75 76 2019). Spreading rates in the Pacific are mostly classified as intermediate or even fast (> 60 77 mm/yr), whereas the MAR is a slow-spreading ridge, which affects the rate and depth range, 78 and potentially other aspects of the local seismicity. The Chain and Romanche, as well as other slow-slipping OTFs in the Atlantic, are likely variable in structure, potentially 79 containing areas of very thin crust and highly serpentinized mantle (Detrick et al., 1993; 80 81 Marjanovic et al., 2020; Gregory et al., 2021) and teleseismic earthquakes can be observed in the shallow mantle beneath (Abercrombie & Ekström, 2001). While there have been studies 82 83 using OBS networks to investigate ridge segments or ridge-transform intersections in the Atlantic (e.g., Toomey at al., 1985; Grevemeyer et al., 2013, Gregory et al., 2021; Yu et al., 84 85 2021), local studies of Atlantic OTFs have been sparse.

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From March 2016 to March 2017, a temporary network of 39 seismic ocean bottom 87 seismometers (OBS) was deployed around the Chain transform fault (CTF) on the MAR and 88 89 continuing eastward fracture zone (CFZ) as part of the PI-LAB ("Passive Imaging of the Lithosphere-Asthenosphere Boundary") project and EURO-LAB (Experiment to Unearth the 90 Rheological Oceanic Lithosphere-Asthenosphere Boundary) (Fig. 1; Agius et al., 2018, 2021; 91 92 Harmon et al., 2018, 2020, 2021, 2022; Hicks et al., 2020; Bogiatzis et al., 2020; Wang et al., 2019, 2020; Rychert et al., 2016, 2021; Saikia et al., 2020, 2021; Leptokaropoulos et al., 93 2021, 2023). The CTF zone has a length of 300 km and varies in width between 7 km and 20 94

95 km with an African Plate half-spreading rate of 18.2mm/yr and a South American Plate half spreading rate of 15.7mm/yr (DeMets et al., 1994; Harmon et al., 2018). The aim of the 96 97 experiment was to achieve a better understanding of the definition of the tectonic plate (Fischer et al., 2010, 2020; Rychert & Shearer, 2009; Rychert et al., 2005, 2007, 2010, 2018a 98 99 2018b, 2020; Tharimena et al., 2016, 2017a, 2017b). The deployment occurred on oceanic crust from its formation at the MAR to older plate ages away from the ridge; the setup was 100 101 chosen to facilitate the observation of thickening oceanic lithosphere relative to increasing plate age. However, the experiment also offers an opportunity to study nearby plate boundary 102 103 seismicity. During the same time, the International Seismic Centre (ISC) recorded 40 earthquakes in the region, of which 23 are located in the vicinity of the Romanche transform 104 fault (RTF), 7 close to the CTF and 8 around the MAR; 2 earthquakes were found on the 105 106 inactive part of the CFZ.

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108 In this study, we present the first year-long catalogue of local seismicity around the equatorial MAR near the CTF, together with a 1D velocity model of the oceanic crust and upper mantle. 109 110 The catalogue provides event times, hypocentre locations and magnitude information. We 111 investigate small-scale seismicity variations along the CTF, the adjacent CFZ and the adjoining sections of the MAR. We also compute earthquake-magnitude distribution at CTF 112 and RTF. Furthermore, we investigate the seismicity in relation to the depth of the brittle 113 114 ductile region predicted by a simple thermal model (Harmon et al., 2018). 115 116 Method 117 Detection and location 118

The data were recorded by the 39 PI-LAB OBS stations during their one-year deployment 120 121 (Fig. 1, Table 1). The stations are mostly distributed in two lines north and south of the CTF 122 and eastern CFZ continuation with an average spacing of 40 km around the active part. On the CTF, the typical minimum station distance to the events is around 50 km. Before working 123 124 with the continuous data, we checked the time stability at each station. Due to timing issues 125 L23D was excluded from the study. We use a Butterworth bandpass filter with corner 126 frequencies of 4Hz and 18Hz, which enhances the signal-to-noise ratio on OBS in a marine 127 environment. The a priori temporary set of picks is created using automated picks from a 128 short-time average divided by long time average (STA/LTA) detector. We use a 1D velocity model that is based on CRUST1.0 (Laske et al., 2012), which includes an average regional 129 crustal thickness of around 8 km (Fig. 2A), plus a 2.90 km thick water layer, which 130 131 corresponds to the depth of the shallowest OBS location (S11D). Deeper than 35km, it is based on ak135 (Kennett et al., 1995). Here, ST and LT time windows of 0.2s and 10s are 132 133 used, together with thresholds of 4 (trigger on) and 1.5 (trigger off). A minimum of 3 P- and S-wave picks are set as a detection requirement. This results in 1492 a priori events with 134 11399 P- and 13948 S-picks. 135

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The P- and S- picks are refined manually and iteratively in combination with an update on the 137 1-D velocity model and the event locations. Based on the picks the average v_P/v_S ratio is 1.73 138 139 (Fig. 2B). Dividing the picks by station and event origin region (Suppl. Fig. S1) reveals deviations from that value. Events from the RTF have a v_P/v_S ratio of ~1.68 at all stations; 140 events from the western CTF have a v_P/v_S ratio of ~1.76 at stations located north of the 141 western CFZ. Event locations are computed using the NonLinLoc package (Lomax et al., 142 2000). The package performs a probabilistic, non-linear global-search inversion using the list 143 of picked arrival times and a regional velocity model as input. Following the probabilistic 144

145	approach of Tarantola & Valette (1982) it calculates the hypocentre location using the
146	maximum likelihood of a probability density function of model parameters in the temporal
147	and spatial domain (event origin time and hypocentre location).
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149	The velocity model is updated using VELEST3.1 (Kissling et al., 1994), which
150	simultaneously inverts for 1-D P- and S-wave velocities and hypocentral parameters. The
151	velocity model has a substantial influence on the depth accuracy, whereas lateral localisation
152	of most earthquakes can be constrained even if more general models are used. Further details
153	can be found in Appendix 1.
154	
155	Given this we perform an additional test to verify the robustness of the depth estimates. We
156	compare NonLinLoc solutions with different maximum search depth limits (15, 20, 25, 30
157	km, as well 200 km, which we will label as "no limit"). When we used the "no limit" setting

158 of 200 km, in some instances NonLinLoc finds some very deep values. Therefore, we chose

159 different depth limits to evaluate the range of possibilities. Here, we present the greatest

160 depth that produces a reliable result, but also present the shallower solutions, to show the

162 NonLinLoc testing with different maximum search depth limits. Although the lateral

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locations are typically the same regardless of the depth limit, we find a few solutions ($\sim 15\%$)

uncertainties. We use the following steps to establish the preferred solution from the

164 where the lateral location changes significantly with different depth limits. These outliers are

discarded for the remainder of the investigation; we then choose the location that has the

166 lowest root mean square (RMS) error if it is outside of error from the remaining location

solutions (see Suppl. Fig. S2 for an example). For locations where results using multiple

depth cut-offs are within error of each other and shallower than 30 km we use the solution

169 with the lowest RMS error. In cases where the lowest RMS error is reached at > 30 km depth

from the sea surface and the solution is within error of a shallower solution, we favour a 170 171 shallower solution, but we mark the quality as marginal. We use 30 km depth as a cut-off 172 because it corresponds to the approximate depth of the $900 - 1000^{\circ}$ isotherms from simple thermal models (Harmon et al., 2018) based on the approximate length of the transform (900 173 km), the averaged half spreading rate (20 mm/yr) and the mantle potential (1350° C), which 174 175 may be the limit of predicted brittle deformation (Molnar, 2020; Kohli et al., 2021). Although 176 locations deeper than 30 km are generally unlikely, we will also present the deep solutions 177 since deeper hypocentres have been found in other locations (McGuire et al., 2012; Kuna et 178 al., 2019; Yu et al., 2021). To ensure the robustness of the computed depths are unbiased by 179 heterogeneities due to long raypaths, we also carried out tests using only the data from the closest five stations with which an azimuthal gap below 180 degrees was still maintained. We 180 did this on a random sampling of 6 events ranging from M_L of 3.3 to 5.6. The locations 181 excluding longer paths were all within error of the original results. 182

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184 Moment tensor inversion

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186 For larger events, it is possible to carry out full moment tensor (MT) inversions, which provides additional constraints on the depth accuracy (e.g., Braunmiller & Nábělek, 2008). 187 We use the Grond software tool (Heimann et al., 2018), which performs a Bayesian 188 189 bootstrap-based inversion of the time-domain deviatoric MT. We only use the vertical component, which we found generally increases MT solution stability. The data are filtered 190 between 12.5 s and 25 s. We fix our epicentral locations at those determined from the 191 192 location inversion, for solution stability. In testing we find that if we relax this constraint the majority of epicentral locations are within the calculated uncertainties of one another. 193 Differences arise in less than 5% of the solutions; these are more unstable due to noisy data, 194

thus resulting in a lack of stations with clear onset signals. Even in these cases the solutions 195 are never more than 5 kilometres away from each other. We estimate the quality of these 196 197 measurements based on the following parameters (Suppl. Figs. S3, S4): 1) the number of stations for which the theoretical waveforms fit the data is high, we assign quality level "MT 198 best", respectively. 2) If the fits are generally poorer and/or the number of stations with good 199 fits is lower than four, we categorise the event as "MT fair". 3) The stacked radiation patterns 200 201 from different ensemble subsets of stations agree with each other (good agreement for best, overall agreement with potential outliers for fair). All solutions show a depth constraint that 202 203 is well defined. In all cases where an MT solution exists, we use this depth estimate.

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205 Magnitud	les
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We calculate local magnitudes for all the located earthquakes using amplitude measurements 207 208 of the maxima on all components in a window of 150 s after the initial P-peak. After demeaning and detrending the data, we remove the instrument response and simulate the 209 response of a Wood-Anderson seismometer (e.g., Bakun & Joyner, 1984; Hutton & Boore, 210 1987; Abercrombie, 1996). The data are then filtered using a high-pass filter with a corner 211 frequency of 4 Hz to reduce oceanic noise. To account for further noise we introduce a 212 signal-to-noise threshold of 1.3 for the zero-to-peak amplitude calculations that are used to 213 214 compute local magnitudes. To account for changes in recorded amplitude with hypocentral distance due to geometric spreading and anelastic attenuation the event magnitudes are 215 adjusted using the correction term: 216

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$$\text{Log}_{10}(A) - M_{ISC} = -1.46 \text{ Log}_{10}(r/100) - 0.0016 (r - 100) - 6.2$$
 (eq. 1)

220	Here, A is the recorded amplitude, M_{ISC} is the magnitude as derived by the International
221	Seismological Centre (ISC), and r is the hypocentral distance. The term is derived in a similar
222	way as done in the work of Abercrombie et al. (1996), using an empirical comparison to our
223	MT solutions (Fig. 3).
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225	Earthquake magnitude-frequency distribution
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227	The earthquake magnitude-frequency distribution can be described by the b -value, which is
228	the gradient of the Gutenberg-Richter relationship $(Log{N(M)} = a - bM; Gutenberg \&$
229	Richter, 1944). The relationship shows that the number of events (N) above a certain
230	magnitude (M) can be expressed by two positive constants, representing the overall seismic
231	activity (a) and the relative occurrence of small to large magnitude earthquakes (b) .
232	
233	To calculate the b -value and associated uncertainties we use the maximum likelihood method

234 (Aki 1965; Utsu 1965; Bender 1983) in combination with a Kolmogorov-Smirnov test. The

catalogue is truncated at a magnitude of 2.3, which represents the magnitude of

236 completeness, M_C , along the CTF. The solution is accepted if the magnitude distribution and

- 237 the straight-line gradient b reach a similarity above M_C above a predetermined significance
- level, in our case chosen to be rigorous at 20% (Schlaphorst et al., 2016).
- 239

240 **Results**

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242 *Earthquake locations*

244	In total, 972 earthquakes could be located laterally (Fig. 4). Constraining depths with a larger
245	network spacing of around 40 km poses challenges, but we were able to constrain depths for
246	928 events. 118 could be constrained by applying an MT inversion (Fig. 4B; Suppl. Fig. S6).
247	Of the remaining event locations 650 are classified as good (Fig. 4C) and 150 are classified
248	as marginal quality (Fig. 4D). The majority of the MT constrained events (88) are located on
249	the CTF and CFZ. The lateral location uncertainty is smaller than the depth uncertainty with
250	a median value of $\tilde{y} = 2.86$. Moment tensor solutions show the lowest depth uncertainties
251	with a median values of $\tilde{z} = 5.42$ km in comparison with the median NonLinLoc solutions for
252	good (\tilde{z} = 7.96 km) and marginal (\tilde{z} = 11.87 km) categories (Fig. 4B, C, D). NonLinLoc
253	hypocentres are better constrained inside the network with the following median values of \tilde{y}
254	= 2.39 km and \tilde{z} = 8.36 km for CTF in contrast to \tilde{y} = 15.09 km and \tilde{z} = 24.81 km for RTF.
255	Nearly all the events are found on or close to the OTFs (793) and the MAR (155). Only a
256	small number (10) are located on the inactive continuing CFZs. Another 13 can be found off
257	these features. Of the events located on or close to an OTF, 700 are close to the CTF, of
258	which the majority are concentrated towards the eastern side, particularly clustered towards
259	the ridge-transform intersection. A further 94 can be found around the RTF.

The earthquake depths determined using NonLinLoc, range from the seafloor to 43 km depth beneath the sea surface. Most (716) are located shallower than 15 km depth. Events on the MAR, away from the ridge-transform intersection, are mostly (36 out of 46) shallower than 10 km depth. Events on the MAR near the transform are deeper, with 12 events between 15 km and 30 km depth categorized as good. 60 events in the entire catalogue were located between 30 and 40 km depth, of which only 4 are classified as good and 56 as marginal. 14 are located at depths greater than 40 km and all are categorised as marginal (Fig. 4D).

271	Most of the MT depths are well-constrained (96 out of 118; Figs. 5, 6). Most solutions (63)
272	are characterised by probability density functions (pdfs) with a clear singular peak. Some
273	solutions (35) have a more complex distribution, but the dominant peak is shallow (<20 km).
274	We find 20 events where a deep $(20 - 49 \text{ km})$ MT solution exists that fits the waveforms
275	equally well or better than a shallower solution (<20 km). We present both depths for
276	completeness (Fig. 7). The deep solutions are likely unrealistic, given the expected maximum
277	depth of brittle deformation (Molnar, 2020; Kohli et al., 2021), and so, we favour the
278	shallower solution. In general, the deepest events are located along the CTF and RTF, while
279	the events beneath the MAR are shallower than 12 km below sea level. There is only one
280	event on the MAR that is deeper than 12 km. There is one intraplate earthquake with an MT
281	solution located on the African plate near the CFZ, with a depth of 18 km.
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283	MT inversion of 118 events in our catalogue shows a mix of strike slip, normal and reverse

283 faults across the region (Fig. 5). Of the 88 focal mechanisms located on the CTF, the majority 284 285 shows right lateral strike-slip movement (39), mostly dipping to the north. Towards the eastern end of the CTF, a combination of strike-slip and reverse faulting is observed (18). 286 287 Several reverse faulting mechanisms can be observed in some locations along the CTF as 288 well (20). Nearly all the 29 MAR events exhibit predominantly normal fault mechanisms with a minor strike-slip component in some cases. Furthermore, a rotation of the normal 289 290 faulting axis can be observed close to ridge-transform intersections (Fig. 5; Suppl. Fig. S8). 291 The one intraplate event on the CFZ has a moment tensor which shows normal faulting with 292 fault plane solutions with north-south strike orientations

296 We calculate local magnitudes from our catalogue (eq. 1) that range from 1.1 to 6.2. The local 6.2 magnitude is a clipped estimate of the Romanche event, that was a larger Mw 7.1 in 297 298 the global CMT (Dziewonski et al., 1981; Ekström et al., 2012) and in a detailed local study 299 of the event (Hicks et al., 2020). The local magnitude scale clips for earthquakes of this large 300 magnitude, and the event is also outside of our array with a small backazimuthal coverage. The next largest event in the catalogue is a M_L 5.6 located on the CTF. We note that our 301 302 relationship between amplitude and local magnitude scale (Fig. 3) resembles one derived for 303 the Azores (Gongóra et al., 2004), but also to the continental setting of southern California (Hutton & Boore, 1987). Similarly, Kuna et al. (2019) found that the southern California 304 305 relationship fit the data recorded along the Blanco transform. It is likely that the tails of the relationship curves are influenced by the number of distant events, since studies do not need 306 307 to consider the curve beyond their maximum event distance (e.g., Abercrombie, 1996; Baumbach et al., 2003). 308

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310 The moment magnitudes from MT inversion range from 3.0 to 6.2. This largest magnitude corresponds to the M_w 7.1 Romanche event (Hicks et al., 2020). The magnitude is vastly 311 underestimated for four main reasons: 1) the large azimuthal gap, 2) the distance to the 312 313 network, 3) the usage of a CTF based velocity model and magnitude calibration, and 4) the frequency range not being low enough for such large magnitude events. On the CTF three 314 M_w >5 events were located, the largest being a M_w 5.6. The largest earthquake that we 315 316 recorded on a MAR segment was a M_w 4.5, while the largest intraplate earthquake we recorded was a M_w 3.2. 317

321	The calculated <i>b</i> -value for the CTF is 0.81 ± 0.09 and for the RTF is 0.76 ± 0.22 (Fig. 8), so
322	indistinguishable given the uncertainties. The network configuration results in large distances
323	to the stations and azimuthal gaps for the events at the RTF, which leads to less well
324	constrained magnitude estimations. These factors combined with the lower number of
325	detected events leads to higher uncertainties for events at RTF. In addition, the magnitude of
326	completeness is larger outside the network at the RTF ($M_C = 3.4$) compared to inside the
327	network at CTF ($M_c = 2.3$). The RTF also has larger median local magnitudes ($M_L = 3.7$)
328	compared to the CTF (M_L = 2.4). Within the CTF the median local M_L values vary, with M_L =
329	2.5 in the west and $M_L = 2.3$ in the east.
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332	Discussion
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334	Earthquake locations
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336	The results generally show an expected pattern of event distribution. They are mostly located
337	close to the ridge-transform plate tectonic boundary between the African and South American
338	Plates. The sub-ridge events are located at depths shallower than 12 km beneath sea level.
339	These likely occur in the crust and the uppermost mantle assuming around 3 km average
340	water depth at the ridge and a 6-7 km average crustal thickness (Christeson et al., 2019),
341	which is similar to a previous, shorter-term local study that found ~4.6 to 5.9 km thick crust
342	to the west of the CTF using an OBS line perpendicular to the OTF direction (Marjanovic et
343	al., 2020). These can be explained through extension of the crust due to seafloor spreading.

Earthquakes at slightly deeper depths than the crust beneath the ridge may occur given that 344 geodynamic modelling suggests that lateral conductive cooling results in non-negligible 345 346 thickness of the mantle lithosphere at slow spreading centres (Parmentier & Morgan, 1990). 347 The deeper earthquakes observed along the transforms suggest that there is brittle deformation to greater depths than beneath the ridges. A deeper brittle-ductile transition 348 349 beneath the transform in comparison to the ridges is predicted by the simple thermal model 350 often used to describe a transform, created by averaging two opposing half-space cooling 351 models. Here we show this model assuming a mantle potential temperature of 1350°C and a thermal diffusivity of 10⁻⁶ mm²/s (Harmon et al., 2018). Brittle deformation in the mantle 352 lithosphere has also been observed in teleseismic estimates for CTF and RTF (Abercrombie 353 & Ekström, 2001). The strain rate can be another component in controlling the depth extent 354 355 of seismicity (Molnar, 2020).

356

357 Focal mechanisms

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359 The focal mechanisms along the plate boundary are generally consistent with the ridgetransform system in the region. Specifically, we observe mostly normal fault MT solutions 360 near the MAR segments, and mostly strike-slip solutions along the CTF and RTF consistent 361 with the global view from the global CMT (Dziewonski et al., 1981; Ekström et al., 2012). 362 363 The general pattern of north dipping focal mechanisms on the CTF and south dipping mechanisms on the RTF is in agreement with previous results that are based solely on 364 365 teleseismic observations (Abercrombie & Ekström, 2001; Hicks et al., 2020; Yu et al., 2021). The reason for the opposite dip directions is unclear but could be related to the symmetry of 366 the system and a transpressional stress regime that includes an anticlockwise rotation (Bonatti 367 et al., 1994; Maia et al., 2016). At both ridge-transform intersections we observe the rotation 368

of the strike of the normal fault focal mechanism from ridge parallel on the ridge to around 369 35° off that orientation towards the strike direction of the OTF. This occurs over roughly 10 370 371 km at the eastern end and around 20 km on the western end (Suppl. Fig. S8). This rotation 372 reflects the transition in stress regime as the plate boundary transitions from spreading to 373 transform fault tectonics (Fox and Gallo, 1984) as predicted by geodynamic modelling 374 (Morgan & Parmentier, 1984). This also results in the observed arcuate shaped scarps 375 bending in towards the transform faults in the ridge-transform intersection in the CTF 376 (Harmon et al. 2018).

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378 The presence of reverse fault MT solutions in some parts of the CTF is indicative of transpression in the strike slip system. The reverse fault mechanisms are typically associated 379 380 with topographic highs within the CTF, which have been interpreted as positive flower structures associated with restraining bends in the normal fault system (Harmon et al., 2018). 381 382 The four positive flower structures have visible sharp scarps in the bathymetry (Harmon et al., 2018), and taken with the MT solutions, suggests that these are active features and the 383 CTF is in active transpression today (Fig. 6). Transpression could be caused by a present-day 384 385 rotation and re-adjustment of the African-South American spreading system. Previously, it has been suggested that the relative plate motions in the region have undergone a rotation of 386 ~11° anticlockwise based on observations of transpressional features on the major OTFs to 387 the north, for example Romanche and St. Paul, (Bonatti et al., 1994; Maia et al., 2016). 388 Compressional mechanisms with fault planes roughly in SW-NE direction would be expected 389 for the reorganisation/transpression direction, which can be observed for some but not all the 390 391 moment tensor solutions. Still, our observations support these notions and suggest that the 392 change in motions is ongoing.

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396 The cluster of intraplate seismicity (5 events) observed on the CFZ on the African Plate at -397 11.5° longitude suggests there are internal stresses in the region of the earthquakes. The 398 epicentres are very close (~5 km) to a M5.2 event recorded in the USGS catalogue on 399 14/04/2016 (before the OBS network deployment), hinting towards an ongoing process. The 400 MT solution indicates a relatively deep event (18 km), suggesting strain is accumulating in 401 the deeper part of the lithosphere, which caused brittle deformation there. The near vertical 402 (or horizontal) fault plane and north-south strike direction are at odds with the expected 403 tectonic fabric of the roughly east-west trending CFZ, so it does not appear to be re-activation of the fracture zone (FZ). Another example of a large intraplate oceanic earthquake in the 404 405 Wharton Basin in the Indian Ocean was interpreted as FZ reactivation given that MT solutions were consistent with the FZ orientations (Yue et al., 2012), and the event may have 406 407 been related to the nearby Sumatran subduction zone. One explanation for our event is that the intraplate stresses are due to the aforementioned rotation of the relative plate motions 408 409 (Bonatti et al., 1994; Maia et al., 2016). It is also possible that strain caused by age 410 differences leading to variations in plate thickness and/or density across the FZ cause that 411 seismic activity (DeLong et al., 1977). However, we do not observe dip-slip faulting due to 412 topography. Another potential explanation of intraplate stress in the region may be due to 413 ocean mantle dynamics. Surface wave tomography and magnetotelluric imaging find a high velocity and high resistivity anomaly in this region centred at 12°W and 0°N (Fig. 4) with 414 415 shear velocities of greater than 4.6 km/s at depths of 30 to 50 km (Saikia et al., 2021). This 416 feature may be related to lithospheric heterogeneity and/or be the beginning of a lithospheric drip. We note that this feature is at shallower depths than the high velocity anomaly further 417 east at 10.5°W longitude (east of our research area) and 0°N latitude with shear velocities of 418

419 greater than 4.4 km/s and log₁₀ (resistivities) > 1.5 extending from 30 to 100 km depth that 420 was interpreted as a lithospheric drip (Harmon et al., 2020; 2021; Wang et al., 2020). The 421 lack of an observed receiver function phase from the lithosphere-asthenosphere boundary in 422 the latter location was interpreted as either the absence of melt ponding beneath the plate 423 and/or strong lithosphere-asthenosphere topography, again consistent with a drip (Rychert et 424 al., 2021). The observed vertical MT motion of the large event recorded in this region may be 425 consistent with the vertical stresses of the downwelling lithosphere (Fig. 5).

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427 Events beneath the $600^{\circ}C$ isotherm

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Some events beneath the CTF are located deeper than the predicted depth of the 600°C 429 430 isotherm from half-space cooling, i.e., a greater depth than the previously proposed limit of seismic slip at 600°C based on teleseismic observations (Abercrombie & Ekström, 2001) and 431 432 earthquake rupture experiments (Boettcher et al., 2007). Earthquakes deeper than expected based on isotherms predicted for an oceanic transform from the half-space cooling model 433 have also been observed at Blanco (Kuna et al., 2019), Gofar (McGuire et al., 2012) and 434 435 Romanche (Yu et al., 2021). These observations are typically interpreted as lithospheric cooling caused by hydrothermal circulation, which can deepen the depth of the brittle-ductile 436 437 transition (Roland et al., 2010). In some cases, the events were deeper than expected by 438 numerical modelling accounting for hydrothermal circulation (Roland et al., 2010; McGuire et al., 2012; Kuna et al., 2019). Alternatively, lateral variability in composition might have an 439 440 influence on the seismicity. A likely explanation for the CTF is that brittle and ductile deformation can occur over a broad range of temperatures $(300^{\circ}C - 900^{\circ}/1000^{\circ}C)$ and, 441 442 therefore, a broad range of depths, owing to variable seawater infiltration and grain sizes

along the fault. This is supported by observations of exhumed hydraulically altered mylonites 443 444 (Molnar, 2020; Kohli et al., 2021).

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Earthquake magnitude-frequency distribution 446

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467

448 The *b*-value of the entire catalogue of CTF events is lower than the global average value of 449 1.0 but in agreement with the results from a previous study of tidal triggering in the Mid-Atlantic segment just north of the CTF (Leptokaropoulos et al., 2021). The low b-value (0.81 450 451 \pm 0.09) also supports the notion that average coupling along the CTF is low and that much of the strain is accommodated aseismically. However, historical seismicity, which has resulted 452 in several large events (Mw > 6.5), suggests that coupling in discrete locations may be higher 453 454 (Shi et al., 2021). This agrees with results from Blanco Transform, where crustal *b*-values are similarly low ($b = 0.78 \pm 10$) and coupling is low but variable along the fault (Kuna et al., 455 2019). A study of different plate boundary settings found similarities between magnitude-456 457 frequency relationships of OTFs and continental TFs (Bird & Kagan, 2003). 458 459 Cumulative moment release 460 We compare the cumulative moment release over the 1-year of our array to examine the 461 462 average seismic coupling. We find that the cumulative moment release on the CTF was 4.9 x 10^{17} Nm. The release on the RTF was 5.7 x 10^{19} Nm. The release on the ridge segments had a 463 cumulative moment of 3.3×10^{16} Nm. The higher release on the Romanche is due to the M_w 464 7.1 event and its aftershock sequence. A more detailed investigation of this event is presented 465 in (Hicks et al., 2020). The predicted moment release per year for the thermal lithosphere 466 from Boettcher and Jordan, (2004) for the CTF is 5.6 x 10¹⁹ Nm yr⁻¹, which suggests only

~1% of the predicted moment release was released seismically over the year-long observation 468 period. Our results also support the notion that the CTF is relatively weak and not coupled. In 469 470 contrast, continental TFs show higher amounts of seismic coupling (Bird & Kagan, 2003). 471 Low coupling values of normally around 10-30% have been observed for OTFs around the world, although showing large variability (Boettcher & Jordan, 2016), and a consensus of a 472 473 general inverse correlation between seismic coupling and spreading rate has been found (e.g., 474 Kawasaki et al., 1985; Sobolev & Rundquist, 1999). For many OTFs in the Atlantic, higher values have been observed (e.g., Gibbs, Oceanographer, 15-20 OTF – Muller, 1983; Kane 475 OTF — Wilcock et al., 1990), which all have a present-day full spreading rate between 21 476 and 26 mm/yr (DeMets et al., 1994), slightly slower than the CTF. However, low values have 477 been found as well (e.g., Vema, Kane, Doldrums OTF - Muller, 1983), which, apart from the 478 479 Kane OTF, have spreading rates between 28 and 33 mm/yr. Uncertainties can arise due to shortness of the observation duration, as well as a large estimation of annual moment release 480 481 due to the choice of OTF dimensions (Muller, 1983; Wilcock et al., 1990). For the entire Romanche, the predicted value of cumulative moment release from the thermal lithosphere is 482 2.9 x 10¹⁹ Nm yr⁻¹, approximately half of the observed moment release. Large events like the 483 M_w 7.1 occur relatively rarely, approximately every 30 years on Romanche, indicating the 484 485 moment release during this time was particularly high. Given that these large events occur on different segments, and do not rupture the entire length of the transform fault, portions of the 486 487 RTF likely have high seismic coupling. This observation also highlights the temporal variability of moment release on OTFs. The MAR segments show a very low cumulative 488 moment (3.3 x 10^{16} Nm) compared to the OTFs (>=4.9 x 10^{17} Nm) in the region for the 489 490 duration of the deployment, which may suggest that the normal faulting also has a low 491 coupling during this time period.

492

495 We present an earthquake catalogue from a broadband ocean bottom seismic experiment 496 from March 2016 to March 2017 at the equatorial MAR centred around the CFZ, which provides one of the first detailed seismicity studies from a larger local network around a slow 497 498 moving OTF. Most of the events are located on the South American – African Plate boundary 499 as expected although we also recorded 10 events near the CFZ, east of the OTF and 13 500 additional intraplate earthquakes. Focal mechanisms on the MAR are predominantly 501 characterised as normal faulting and on the OTF they are predominantly strike slip as expected. A few (17) predominantly reverse faulting mechanisms on the transform are likely 502 caused by a transpressional stress regime related to current rotation of the ridge-transform 503 504 system. The focal mechanism of an intraplate event on the CFZ shows vertical displacement, 505 consistent with a hypothesized nearby lithospheric drip. The magnitudes range from 1.1 to 506 5.6 and we find a *b*-value of 0.81 ± 0.09 for the CTF. The earthquakes on the ridge are 507 limited to the crust and the shallowest mantle (< 12 km depth) while the events on the CTF are deeper, with high quality MT location depths of up to 26 km beneath the sea surface, 508 509 which is much deeper than the depth of the brittle-ductile transition predicted for simple 510 thermal models. A simple thermal model cannot explain these deeper events, and another 511 factor is required, for instance, lithospheric cooling by hydrothermal circulation.

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- 523 created the initial automated detection of events, assisted with the locations, and commented
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- 525 manuscript. J.M.K. co-managed the project and commented on the manuscript. R.E.A.
- 526 contributed to improving the event localization method and commented on the manuscript.
- 527

528 Data Availability

- 529 The continuous raw seismic waveform data from the PI-LAB OBS network is available to
 530 download from IRIS Data Management Center (https://doi.org/10.7914/SN/XS 2016).
- 531

532 Appendix 1 – Details on iterative pick refinement, model update and event localisation 533 in NonLinLoc

In step 1, P- and S-arrivals are readjusted using the vertical component and the Seismic

- 535 Analysis Code (SAC). This is done without the implementation of calculated theoretical
- arrival times derived from the automated picks in order to keep the reassessed picks unbiased.
- 537 Although different filters as well as the unfiltered data are used to evaluate potential onsets,

we use the same corner frequencies to determine the onsets to have consistent P- and Smoveouts and P-to-S travel time differences. The weighting of picks is assigned to reflect
their quality, using levels from 0 (highest) to 4 (lowest). The levels are associated with an
uncertainty in arrival time, using values of 0.1s, 0.2s, 0.5s, 0.8s, and 1.5s (Hicks et al., 2014).
Picks of quality level 4 are subsequently omitted from the locating process.

543

544 In step 2, the event locations are computed using the NonLinLoc package (Lomax et al., 545 2000). The package performs a probabilistic, non-linear global-search inversion using the list 546 of picked arrival times and a regional velocity model as input. Following the probabilistic approach of Tarantola & Valette (1982) it calculates the hypocentre location using maximum 547 likelihood of a probability density function of model parameters in the temporal and spatial 548 549 domain. The initial velocity model is given by CRUST1.0 resulting in an average regional crustal thickness of around 8 km (Fig. 2A), plus a 2.90 km thick water layer, which 550 551 corresponds to the depth of the shallowest OBS location (S11D). We do not include a sediment layer since it is supposed to be very thin in the area (Agius et al., 2018; Saikia et al., 552 2020). In this step, we also combine picks that are counted as multiple events by the 553 554 STA/LTA detector, thus reducing the overall number of events.

555

A minimum of 3 stations with clearly observable and coherent arrivals is required to identify the event. We use a search grid spacing of 1km in all three dimensions and an average v_P/v_S ratio of 1.73, based on picked P–S arrival time differences (Fig. 2B). The quality of the resulting location is based on temporal uncertainty (i.e., root mean square or RMS error) as well as spatial uncertainty. The latter is a combination of the size of the ellipsoidal approximation to the 1 σ confidence level of the location likelihood scatter and the overall shape of the scatter sample. Lateral localisation of most earthquakes can be constrained

relatively precisely, even if more general models are used, whereas the velocity model has alarger impact on depth localisation.

565

In step 3, the picks of events with low quality are refined. This includes all events that either 566 have large location uncertainty and/or where picks and theoretical arrival times show large 567 568 mismatches. Here we use the Seismic Data eXplorer (SDX; Hicks et al., 2014), additionally 569 including the two horizontal components and the hydrophone component. SDX updates theoretical onset time calculations of P- and S-phases automatically after each change of pick 570 571 placement. In general, this reduces the number of mispicks. There are, however, instances, 572 where clearly identifiable onsets do not match the theoretical arrivals. In most cases, this is 573 due to an inaccurate velocity model.

574

In step 4, the velocity model is updated, based on the updated set of picks. We use the program VELEST3.1 (Kissling et al., 1994), which simultaneously inverts for 1D P- and Swave velocities and hypocentral parameters. For this step, we concentrate on the picks of events on or close to the CTF. This ensures good azimuthal coverage with a gap of less than 180°. Also, regional heterogeneities are likely to distort a uniform velocity model if the area is chosen too large and diverse. We limit the selection further to events with an RMS of less than 1.2, a minimum number of P-picks of 5, and a minimum number of S picks of 3.

582

583 Steps 2, 3 and 4 are run iteratively until the model and arrival time picks stabilise. Here, three584 iterations proved to be sufficient.

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950 Tables

Tab. 1: PI-LAB network OBS station details. Stations indicated with an asterisk were notused. Except for L27A running time is equivalent to deployment time.

953		1 1	
954	Station Lon Lat Depth (m)	Running Time	Notes
955	I01D* -17.8855 1.2734 -4047	N/A	unrecovered
956	L02A -17.4085 -1.1667 -3499	06/03/2016 - 24/03/2017	did not level
957	S03D -17.0315 -2.4021 -3750	05/03/2016 - 25/03/2017	
958	I04D -16.1733 -2.1238 -3928	07/03/2016 - 24/03/2017	
959	L05A -15.4058 -1.8577 -4052	07/03/2016 - 21/03/2017	stopped after 150 days
960	S06D -14.4298 -1.6703 -3778	06/03/2016 - 19/03/2017	
961	I07D -14.0428 -1.5565 -3819	08/03/2016 - 22/03/2017	
962	L08D* -13.6409 -1.4493 -3357	08/03/2022 - 18/03/2017	only 1 horizontal channel
963	L09A -13.3185 -1.3569 -3378	08/03/2016 - 18/03/2017	1 flat horizontal channel
964	S10D -12.9697 -1.3180 -3015	07/03/2016 - 18/03/2017	-
965	S11D -12.4602 -1.1691 -2905	07/03/2016 - 18/03/2017	
966	I12D -10.7766 -0.8683 -4022	19/03/2016 - 17/03/2017	
967	L13D -9.5619 -0.5862 -4659	09/03/2016 - 16/03/2017	usable data until 23/12/16
968	I14D -7.9524 -0.3522 -4702	10/03/2016 - 10/03/2017	
969	S15D -6.6228 0.1814 -4927	09/03/2016 - 16/03/2017	rocking, some bad data
970	L16D -7.8953 0.8933 -4581	11/03/2016 - 15/03/2017	usable data until 24/02/16
971	S17D -8.5121 0.7422 -5205	09/03/2016 - 15/03/2017	
972	L18D -9.3765 0.5769 -4890	11/03/2016 - 14/03/2017	
973	S19D -9.9754 0.4809 -4607	10/03/2016 - 14/03/2017	
974	I20D -10.5352 0.3681 -4724	12/03/2016 - 13/03/2017	
975	L21D -11.0380 0.2364 -4625	11/03/2016 - 13/03/2017	
976	S22D -11.6799 0.1254 -4352	10/03/2016 - 13/03/2017	
977	L23D* -12.1478 0.0521 -4631	12/03/2016 - 12/03/2017	timing issues
978	S24D -12.7806 -0.1383 -4453	10/03/2016 - 13/03/2017	
979	L25D* -13.2230 -0.1745 -4207	12/03/2022 - 11/03/2017	disk failure
980	S26D -13.6260 -0.3434 -4216	11/03/2016 - 12/03/2017	
981	L27A* -13.9427 -0.4007 -3928	12/03/2016 - 23/08/2016	l flat horizontal channel
982	I28D -14.2684 -0.4918 -3711	13/03/2016 - 11/03/2017	
983	S29D -14.6272 -0.5597 -3626	11/03/2016 - 11/03/2017	
984	L30A -14.9467 -0.5880 -4003	13/03/2016 - 10/03/2017	
985	S31D -15.3188 -0.7141 -3408	11/03/2016 - 11/03/2017	
986	S32D -15.6470 -0.7968 -2967	11/03/2016 - 11/03/2017	
987	L33D -16.0152 -0.8747 -3919	13/03/2016 - 09/03/2017	
988	I34D -16.3485 -0.9579 -2964	14/03/2016 - 09/03/2017	
989	S35D* -16.6798 -1.0372 -3773	11/03/2016 - 11/03/2017	disk failure
990	I36D* -17.0306 -1.1170 -3938	N/A	unrecovered
991	L37D -14.9718 1.5657 -5054	19/03/2016 - 08/03/2017	
992	S38D -12.7623 1.9218 -4926	18/03/2016 - 08/03/2017	
993	L39D* -11.4904 2.0557 -4685	20/03/2016 - 07/03/2017	unrecovered
994			

996 Figures



Fig. 1: Map of the study area (see inset for global context), including the PI-LAB station 998 999 network. Some stations did not provide usable data for the entire timespan of the experiment 1000 (white triangles), some stations were only returning useful data for parts of the deployment 1001 (grey triangles); see details in Table 1. The thick solid line denotes the locations of the Mid 1002 Atlantic Ridge (MAR) and the Romanche and Chain Transform Faults and Fracture Zones 1003 (RTF, CTF, CFZ; adapted from Bird, 2003). Events from the ISC during the time of the deployment are shown by circles, events around the CTF and CFZ (white circles) have been 1004 used to calibrate the local event magnitude (see Fig. 3). In the global catalogue, all these 1005 1006 events have been assigned default depth values of 10km. Bathymetry data is taken from 1007 ETOPO1 (Amante & Eakins, 2009) with an additional higher resolution image around the CTF (Harmon et al., 2018). White thin lines show age contours in Ma (Müller et al, 2008). 1008 1009 Arrows depict the annual half-spreading rate (Harmon et al., 2018).



Fig. 2: (A) Initial (CRUST1.0) and final 1-D velocity-depth models. Note that CRUST1.0
only provides a P-velocity model. M2020 shows a model by Marjanovic et al. (2020),
adjusted to the same water column thickness of 2.9 km. (B) Wadati plot of P- and S- arrival

1015 time picks (7345 pairs) for all events. The red line shows the average v_P/v_S ratio of 1.73.

1016 About 97.1% (7130 pairs) fall within \pm 0.4 of that value, about 67.0% (4923 pairs in blue)

1017 fall within \pm 0.07. See Supplementary Figure S1 for detailed versions of the Wadati plot

1018 divided by station and event origin region.



Fig. 3: Magnitude calibration with MT inversion solutions. Note, that results from station
S15D (grey dots) are taken out of the calculation, since the instrument was located on a slope
rocking and produced erroneous data. The inset shows different correction curves, all shifted
to match the original setup by Richter (1935, 1958).



Fig. 4: Catalogue event location and magnitude results. (A) Map showing the epicentres
coloured by depth. The red circle indicates the location and rough extent of a high velocity
anomaly at 30 km depth (Saikia et al., 2021). Top and left panels show histograms by
projected longitude and latitude. The right panel shows a cross section along the MAR, with
event symbols representing different result quality levels corresponding to the symbols in the
other cross section panels. (B-D) Cross section along the CTF and CFZ. Depth uncertainties

- 1033 are shown in the blue histograms with the median value being depicted by vertical black
- 1034 lines. For clarity, error bars have been omitted from the figure but can be seen in
- 1035 Supplementary Figure S5. Events are divided by quality (B: events for which moment tensor
- 1036 inversion was possible; C: good locating results with NonLinLoc; D: marginal locating
- 1037 results with NonLinLoc). Isotherms (300 1200°C in steps of 300°C) from a half-space
- 1038 cooling model (using the lower half spreading rate from the African plate, 15.7mm/yr, which
- 1039 results in deeper isotherms) are also shown along the CTF. Also shown are depth histograms
- 1040 (left) and the final P-velocity model (right). See Figure 1 for further details. See
- 1041 Supplementary Table S1 for event details.
- 1042



Fig. 5: Map including all focal mechanisms (shown as MT inversion results in Fig. 4). The
mechanisms are placed apart from the OTF for better visibility. The shading is based on their
dominant mechanism according to rake angle. Red lines indicate mapped transpressive
flower-structures on the CTF (Harmon et al., 2018). Zoomed-in views of the regions around
the flower structures can be seen in Figure 6. See Figure 1 for further details, Supplementary
Table S2 for focal mechanism details, and Supplementary Figure S7 for an example of a
waveform fit.



1053 Fig. 6: Zoomed-in versions of flower structures on the CTF, including focal mechanisms,

- 1054 coloured by their rake angle. For further information see Figures 1 and 5.
- 1055



Fig. 7: Map showing all MT results that show a bimodal distribution with good shallow anddeep solutions, including shallow (circles) and deep (diamonds) in the cross section.



Fig. 8: Earthquake-magnitude distribution on the CTF (A) and RTF (B). Blue circles and
lines show the magnitude range used to calculate the Gutenberg-Richter relationship. The
result is shown (solid red line) with its uncertainties (dashed red lines).