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# Weak solutions for the Oseen system in 2D and when the given velocity is not sufficiently regular

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ABSTRACT

 $\mathbf{L}^{s}(\Omega)$  for s < n.

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#### 1. Introduction

This work is dedicated to the study of some existence aspect related to the Oseen problem in bounded domain  $\Omega \subset \mathbb{R}^n$ , n = 2, 3:

(O)  $-\Delta \boldsymbol{u} + \boldsymbol{v} \cdot \nabla \boldsymbol{u} + \nabla \pi = \boldsymbol{f}, \quad \nabla \cdot \boldsymbol{u} = h \text{ in } \Omega, \quad \boldsymbol{u} = \boldsymbol{g} \text{ on } \Gamma.$ 

In the 3-dimensional case, the existence of weak solutions  $(\boldsymbol{u}, \pi) \in \mathbf{H}^1(\Omega) \times L^2(\Omega)$ , regular solution in  $\mathbf{H}^2(\Omega) \times H^1(\Omega)$  and  $\mathbf{W}^{1,p}(\Omega) \times L^p(\Omega)$  (and intermediate Sobolev spaces) together with the analysis of the existence of very weak solutions in  $\mathbf{L}^p(\Omega) \times W^{-1,p}(\Omega)$  have been analyzed by the authors in [1], assuming  $\boldsymbol{v}$  a solenoidal field belonging to  $\mathbf{L}^s(\Omega)$  for  $s \geq 3$  (from now on, we will denote this solenoidal space by  $\mathbf{L}^s_{\sigma}(\Omega)$ ). However, the existence of solution for the 2-dimensional Oseen system has not been attacked in [1] because the "logical" assumption of considering the solenoidal field  $\boldsymbol{v} \in \mathbf{L}^2(\Omega)$  (in order to obtain weak solutions for (O)) poses some difficulties in the treatment of the convective term  $(\boldsymbol{v} \cdot \nabla)\boldsymbol{u}$ : On the one hand, it is not

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The aim of this work is twofold: proving the existence of solution  $(u, \pi) \in \mathbf{H}^1(\Omega) \times$ 

 $\mathbf{L}^{2}(\Omega)$  in bounded domains of  $\mathbb{R}^{2}$  and the whole plane for the Oseen problem (O)

for solenoidal vector fields v in  $\mathbf{L}^{2}(\Omega)$ , and analyzing the same problem in bounded

domains of  $\mathbb{R}^n$  for n = 2, 3 when h = 0, q = 0 and the solenoidal field v belongs to



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clear if the bilinear form associated is coercive and continuous. Some related results can be found in [2] for the scalar case (instead of considering a vector field solution  $\boldsymbol{u}$ , one considers a scalar unknown  $\theta$ ) and for  $\boldsymbol{g} = \boldsymbol{0}$ . On the other hand, when  $\Omega = \mathbb{R}^2$  an additional awkwardness appears because even if we can prove  $\nabla \boldsymbol{u} \in \mathbf{L}^2(\mathbb{R}^2)$ , it is not evident that  $\boldsymbol{u} \in \mathbf{L}^p(\mathbb{R}^2)$  (for any p). Giving a successful answer to both previous problems is our first aim.

Our second aim is to give a first answer to the question of the existence of solution for the Oseen problem (O) when  $\boldsymbol{v}$  only belongs to  $\mathbf{L}_{\sigma}^{s}(\Omega)$ , with s < n and n = 2, 3.

#### 2. Solutions for the Oseen problem in the 2-dimensional case

The existence of weak solutions in  $\mathbf{H}^{1}(\Omega)$  for Problem (O) in 2-dimensional domains is not known when a solenoidal field  $\boldsymbol{v}$  that only belongs to  $\mathbf{L}^{2}(\Omega)$  is considered. In this case, the term  $(\boldsymbol{v} \cdot \nabla)\boldsymbol{u}$  belongs only to  $\mathbf{L}^{1}(\Omega)$ . It is then not clear neither if the bilinear form associated to the Problem (O), with h = 0 and  $\boldsymbol{g} = \boldsymbol{0}$ :

$$a(\boldsymbol{u}, \boldsymbol{w}) = \int_{\Omega} \nabla \boldsymbol{u} \cdot \nabla \boldsymbol{w} \, d\boldsymbol{x} + \int_{\Omega} (\boldsymbol{v} \cdot \nabla) \boldsymbol{u} \cdot \boldsymbol{w} \, d\boldsymbol{x}$$

is coercive on the space  $\mathbf{V}(\Omega) = \{ \boldsymbol{w} \in \mathbf{H}_0^1(\Omega); \text{ div } \boldsymbol{w} = 0 \text{ in } \Omega \}$  nor if it is continuous on  $\mathbf{V}(\Omega) \times \mathbf{V}(\Omega)$ . In order to overcome this difficulty, we use the Hardy space  $\mathcal{H}^1(\mathbb{R}^2)$ . One equivalent definition of such a space (in the *n*-dimensional case) is [3]:

$$\mathcal{H}^1(\mathbb{R}^n) = \{ f \in L^1(\mathbb{R}^n), R_j f \in L^1(\mathbb{R}^n), 1 \le j \le n \} \text{ where } R_j = \frac{\partial}{\partial x_j} (-\Delta)^{-1/2}.$$

A partial study of the *BMO* spaces (Bounded Mean Oscillation) will be also necessary taking into account the duality between  $\mathcal{H}^1$  and the *BMO* (see [4]). Moreover, the *VMO*-space (Vanishing Mean Oscillator) is a subspace of the *BMO*: a function f in *BMO*( $\mathbb{R}^n$ ) is said to be in *VMO*( $\mathbb{R}^n$ ) if

$$\lim_{r \to 0} \sup_{\boldsymbol{x}_0 \in \mathbb{R}^n} \frac{1}{r^n} \int_{B(\boldsymbol{x}_0, r)} |f - \overline{f}| \mathrm{d}\boldsymbol{x} = 0, \quad \text{where } \overline{f} = \frac{1}{|B(\boldsymbol{x}_0, r)|} \int_{B(\boldsymbol{x}_0, r)} f.$$

It is also crucial the fact that  $H^1(\mathbb{R}^2) \hookrightarrow VMO(\mathbb{R}^2)$  (see [5]).

With these ingredients, we will prove one of the two main results of this work, namely Theorem 2.2 in bounded domains and Theorem 2.5 if  $\Omega = \mathbb{R}^2$ . In order to prove them, we use the following result:

**Lemma 2.1.** Assume  $v \in \mathbf{L}^2_{\sigma}(\Omega)$  and  $y \in H^1_0(\Omega)$ . Then  $(v \cdot \nabla)y \in H^{-1}(\Omega)$  and

$$\|(\boldsymbol{v}\cdot\nabla)\boldsymbol{y}\|_{H^{-1}(\Omega)} \le C \,\|\boldsymbol{v}\|_{\mathbf{L}^{2}(\Omega)} \|\nabla\boldsymbol{y}\|_{\mathbf{L}^{2}(\Omega)}.$$
(1)

Moreover, we have that

$$\langle \boldsymbol{v} \cdot \nabla \boldsymbol{z}, \boldsymbol{z} \rangle_{\mathbf{H}^{-1}(\Omega) \times \mathbf{H}_0^1(\Omega)} = 0 \quad \text{for all } \boldsymbol{z} \in \mathbf{H}_0^1(\Omega).$$
 (2)

**Proof.** Indeed, considering  $\boldsymbol{w} \in \mathbf{L}^2(\mathbb{R}^2)$  the extension of  $\boldsymbol{v}$  to  $\mathbb{R}^2$  given by:  $\boldsymbol{w} = \boldsymbol{v}$  in  $\Omega$ , and  $\boldsymbol{w} = \nabla \theta$  in  $\Omega' = \mathbb{R}^2 \setminus \overline{\Omega}$  where  $\theta$  is the solution of the following problem:

$$\Delta \theta = 0 ext{ in } \Omega', \quad rac{\partial heta}{\partial oldsymbol{n}} = -oldsymbol{v} \cdot oldsymbol{n} ext{ on } arGamma,$$

with  $\nabla \theta \in \mathbf{L}^2(\Omega')$  that satisfies

$$\|\nabla \theta\|_{\mathbf{L}^{2}(\Omega')} \leq C \|\boldsymbol{v} \cdot \boldsymbol{n}\|_{H^{-1/2}(\Gamma)} \leq C \|\boldsymbol{v}\|_{\mathbf{L}^{2}(\Omega)}$$

because  $\nabla \cdot \boldsymbol{v} = 0$  in  $\Omega$ . Moreover,  $\nabla \cdot \boldsymbol{v} = 0$  in  $\Omega$  implies  $\langle \boldsymbol{v} \cdot \boldsymbol{n}, 1 \rangle_{\Gamma} = 0$  and the existence of  $\theta$  is ensured by Theorem 3.1 [6]. Observe that  $\nabla \cdot \boldsymbol{w} = 0$  in  $\mathbb{R}^2$  because for  $\varphi \in \mathcal{D}(\mathbb{R}^2)$ :

$$\langle \nabla \cdot \boldsymbol{w}, \varphi \rangle = -\int_{\Omega} \boldsymbol{v} \cdot \nabla \varphi \, d\boldsymbol{x} - \int_{\Omega'} \nabla \theta \cdot \nabla \varphi \, d\boldsymbol{x} = \langle \boldsymbol{v} \cdot \boldsymbol{n}, \varphi \rangle_{\Gamma} - \langle \boldsymbol{v} \cdot \boldsymbol{n}, \varphi \rangle_{\Gamma} = 0,$$

and  $\|\boldsymbol{w}\|_{\mathbf{L}^2(\mathbb{R}^2)} \leq C \|\boldsymbol{v}\|_{\mathbf{L}^2(\Omega)}$ . On the other hand, we consider  $\tilde{y}$  the extension by zero of y that satisfies  $\tilde{y} \in H^1(\mathbb{R}^2)$ . Using Theorem II.2 point (2) or Theorem II.1 point (2) in [3], we can deduce that  $\boldsymbol{w} \cdot \nabla \tilde{y} \in \mathcal{H}^1(\mathbb{R}^2)$  and the bound

$$\|\boldsymbol{w}\cdot\nabla\widetilde{\boldsymbol{y}}\|_{\mathcal{H}^{1}(\mathbb{R}^{2})} \leq C\|\boldsymbol{w}\|_{\mathbf{L}^{2}(\mathbb{R}^{2})}\|\nabla\widetilde{\boldsymbol{y}}\|_{\mathbf{L}^{2}(\mathbb{R}^{2})} \leq C\|\boldsymbol{v}\|_{\mathbf{L}^{2}(\Omega)}\|\nabla\boldsymbol{y}\|_{\mathbf{L}^{2}(\Omega)}.$$

Now, we have to prove that  $\boldsymbol{v} \cdot \nabla y \in H^{-1}(\Omega)$  and  $\langle \boldsymbol{v} \cdot \nabla y, y \rangle_{H^{-1}(\Omega) \times H^1_0(\Omega)} = 0$ . Indeed, for  $\varphi \in \mathcal{D}(\Omega)$ 

$$\begin{aligned} \left| \int_{\Omega} \varphi \boldsymbol{v} \cdot \nabla y \, d\boldsymbol{x} \right| &= \left| \int_{\mathbb{R}^2} \widetilde{\varphi} \boldsymbol{w} \cdot \nabla \widetilde{y} \, d\boldsymbol{x} \right| &\leq \| \boldsymbol{w} \cdot \nabla \widetilde{y} \|_{\mathcal{H}^1(\mathbb{R}^2)} \| \widetilde{\varphi} \|_{BMO(\mathbb{R}^2)} \\ &\leq C \| \boldsymbol{v} \|_{\mathbf{L}^2(\Omega)} \| \nabla y \|_{\mathbf{L}^2(\Omega)} \| \widetilde{\varphi} \|_{H^1(\mathbb{R}^2)} \\ &\leq C \| \boldsymbol{v} \|_{\mathbf{L}^2(\Omega)} \| \nabla y \|_{\mathbf{L}^2(\Omega)} \| \varphi \|_{H^1(\Omega)} \end{aligned}$$

because  $H^1(\mathbb{R}^2) \hookrightarrow VMO(\mathbb{R}^2) \hookrightarrow BMO(\mathbb{R}^2)$ . In that way, as  $\mathcal{D}(\Omega)$  is dense in  $H^1_0(\Omega)$ , we can deduce that  $\boldsymbol{v} \cdot \nabla \boldsymbol{y} \in H^{-1}(\Omega)$  and estimate (1).

For the proof of (2), let us consider  $\mathbf{z}_k \in \mathcal{D}(\Omega)$  be such that  $\mathbf{z}_k \to \mathbf{z}$  in  $\mathbf{H}_0^1(\Omega)$ . Then,

$$\begin{aligned} |\langle \boldsymbol{v} \cdot \nabla \boldsymbol{z}, \boldsymbol{z} \rangle_{\mathbf{H}^{-1}(\Omega) \times \mathbf{H}_{0}^{1}(\Omega)} - \langle \boldsymbol{v} \cdot \nabla \boldsymbol{z}_{k}, \boldsymbol{z}_{k} \rangle_{\mathbf{H}^{-1}(\Omega) \times \mathbf{H}_{0}^{1}(\Omega)}| \\ & \leq |\langle \boldsymbol{v} \cdot \nabla(\boldsymbol{z} - \boldsymbol{z}_{k}), \boldsymbol{z} \rangle_{\mathbf{H}^{-1}(\Omega) \times \mathbf{H}_{0}^{1}(\Omega)}| + |\langle \boldsymbol{v} \cdot \nabla \boldsymbol{z}_{k}, (\boldsymbol{z}_{k} - \boldsymbol{z}) \rangle_{\mathbf{H}^{-1}(\Omega) \times \mathbf{H}_{0}^{1}(\Omega)}| \end{aligned}$$

Using (1) and the convergence of  $z_k$  to z in  $\mathbf{H}_0^1(\Omega)$ , both duality terms on the right-hand-side of the previous inequality tend to 0 when  $k \to +\infty$ .

Finally, from  $\langle \boldsymbol{v} \cdot \nabla \boldsymbol{z}_k, \boldsymbol{z}_k \rangle_{\mathbf{H}^{-1}(\Omega) \times \mathbf{H}^1_0(\Omega)} = 0$ , we can deduce (2).  $\Box$ 

**Theorem 2.2** (Existence of Weak Solution for (O)). Let  $\Omega$  be a Lipschitz bounded domain in  $\mathbb{R}^2$ . Let

$$\boldsymbol{f} \in \mathbf{H}^{-1}(\Omega), \quad \boldsymbol{v} \in \mathbf{L}^2_{\sigma}(\Omega), \quad h \in L^2(\Omega) \quad \text{and} \quad \boldsymbol{g} \in \mathbf{H}^{1/2}(\Gamma)$$

satisfy the compatibility condition

$$\int_{\Omega} h(\boldsymbol{x}) \, d\boldsymbol{x} = \int_{\partial \Omega} \boldsymbol{g} \cdot \boldsymbol{n} \, d\sigma.$$
(3)

Then, the problem (O) has a unique solution  $(\boldsymbol{u}, \pi) \in \mathbf{H}^1(\Omega) \times L^2(\Omega)/\mathbb{R}$ . Moreover, there exist some constants  $C_1 > 0$  and  $C_2 > 0$  such that:

$$\|\boldsymbol{u}\|_{\mathbf{H}^{1}(\Omega)} \leq C_{1}\left(\|\boldsymbol{f}\|_{\mathbf{H}^{-1}(\Omega)} + \left(1 + \|\boldsymbol{v}\|_{\mathbf{L}^{2}(\Omega)}\right)\left(\|\boldsymbol{h}\|_{L^{2}(\Omega)} + \|\boldsymbol{g}\|_{\mathbf{H}^{1/2}(\Gamma)}\right)\right),\tag{4}$$

$$\|\pi\|_{L^{2}(\Omega)/\mathbb{R}} \leq C_{2} \Big(\|\boldsymbol{f}\|_{\mathbf{H}^{-1}(\Omega)} + \Big(1 + \|\boldsymbol{v}\|_{\mathbf{L}^{2}(\Omega)}\Big) \Big(\|h\|_{L^{2}(\Omega)} + \|\boldsymbol{g}\|_{\mathbf{H}^{1/2}(\Gamma)}\Big)\Big),$$
(5)

where  $C_1 = C(\Omega)$  and  $C_2 = C_1 \left( 1 + \| \boldsymbol{v} \|_{\mathbf{L}^2(\Omega)} \right)$ .

**Proof.** Although some parts of this proof are identical to the proof made in [1], we include the whole argument here for completeness. In order to prove the existence of solution, first (using Lemma 3.3 in [7], for instance) we lift the boundary and the divergence data. Then, there exists  $u_0 \in \mathbf{H}^1(\Omega)$  such that  $\nabla \cdot u_0 = h$  in  $\Omega$ ,  $u_0 = g$  on  $\Gamma$  and:

$$\|\boldsymbol{u}_0\|_{\mathbf{H}^1(\Omega)} \le C \left( \|\boldsymbol{h}\|_{L^2(\Omega)} + \|\boldsymbol{g}\|_{\mathbf{H}^{1/2}(\Gamma)} \right).$$
(6)

Therefore, it remains to find  $(\boldsymbol{z}, \pi) = (\boldsymbol{u} - \boldsymbol{u}_0, \pi)$  in  $\mathbf{H}_0^1(\Omega) \times L^2(\Omega)$  such that:

$$-\Delta \boldsymbol{z} + \boldsymbol{v} \cdot \nabla \boldsymbol{z} + \nabla \boldsymbol{\pi} = \mathbf{F} \text{ and } \nabla \cdot \boldsymbol{z} = 0 \text{ in } \boldsymbol{\Omega}, \quad \boldsymbol{z} = \mathbf{0} \text{ on } \boldsymbol{\Gamma}.$$
(7)

being  $\mathbf{F} = \mathbf{f} + \Delta \mathbf{u}_0 - (\mathbf{v} \cdot \nabla) \mathbf{u}_0$ . From Lemma 2.1, we deduce that  $(\mathbf{v} \cdot \nabla) \mathbf{u}_0 \in \mathbf{H}^{-1}(\Omega)$ , then  $\mathbf{F} \in \mathbf{H}^{-1}(\Omega)$ . Since the space  $\mathcal{D}_{\sigma}(\Omega) = \{ \boldsymbol{\varphi} \in \mathcal{D}(\Omega); \nabla \cdot \boldsymbol{\varphi} = 0 \}$  is dense in the space  $\mathbf{V}(\Omega)$ , the previous problem is equivalent to:

Find 
$$\boldsymbol{z} \in \mathbf{V}(\Omega)$$
 such that:  $\forall \boldsymbol{\varphi} \in \mathbf{V}(\Omega)$   
$$\int_{\Omega} \nabla \boldsymbol{z} \cdot \nabla \boldsymbol{\varphi} \, d\boldsymbol{x} + \langle (\boldsymbol{v} \cdot \nabla) \boldsymbol{z}, \boldsymbol{\varphi} \rangle_{\mathbf{H}^{-1}(\Omega) \times \mathbf{H}_{0}^{1}(\Omega)} = \langle \boldsymbol{F}, \boldsymbol{\varphi} \rangle_{\mathbf{H}^{-1}(\Omega) \times \mathbf{H}_{0}^{1}(\Omega)}.$$

Now, using (2) by Lax–Milgram's Theorem, if we assume that  $\mathbf{F} \in \mathbf{H}^{-1}(\Omega)$ , then we can deduce the existence of a unique  $\mathbf{z} \in \mathbf{H}_0^1(\Omega)$  solution of (7) verifying:

$$\begin{aligned} \|\boldsymbol{z}\|_{\mathbf{H}^{1}(\Omega)} &\leq C \|\boldsymbol{F}\|_{\mathbf{H}^{-1}(\Omega)} \\ &\leq C \Big(\|\boldsymbol{f}\|_{\mathbf{H}^{-1}(\Omega)} + \Big(1 + \|\boldsymbol{v}\|_{\mathbf{L}^{2}(\Omega)}\Big) \Big(\|h\|_{L^{2}(\Omega)} + \|\boldsymbol{g}\|_{\mathbf{H}^{1/2}(\Gamma)}\Big)\Big), \end{aligned}$$
(8)

which added to estimate (6) makes (4). We can recover the pressure  $\pi$  thanks to the De Rham's Lemma (Lemma 6 in [1] and Corollary III.5.1 in [8]). Now,  $-\Delta \boldsymbol{z} + \boldsymbol{v} \cdot \nabla \boldsymbol{z} - \boldsymbol{F} \in \mathbf{H}^{-1}(\Omega)$  and:

$$\forall \boldsymbol{\varphi} \in \mathbf{V}(\Omega), \qquad \langle -\Delta \boldsymbol{z} + \boldsymbol{v} \cdot \nabla \boldsymbol{z} - \boldsymbol{F}, \boldsymbol{\varphi} \rangle_{\mathbf{H}^{-1}(\Omega) \times \mathbf{H}_{0}^{1}(\Omega)} = 0.$$

Thanks to De Rham's Lemma, there exists a unique  $\pi \in L^2(\Omega)/\mathbb{R}$  such that

$$-\Delta \boldsymbol{z} + \boldsymbol{v} \cdot \nabla \boldsymbol{z} + \nabla \pi = \mathbf{F}$$

with  $\|\pi\|_{L^2(\Omega)/\mathbb{R}} \leq C \|\nabla\pi\|_{\mathbf{H}^{-1}(\Omega)}$ . Finally, estimate (5) follows from the previous equation and estimate (8) for  $\boldsymbol{z}$ .  $\Box$ 

With the same procedure than in [1], we can prove strong and weak- $W^{1,p}(\Omega)$  regularity for (O) in the 2-dimensional bounded case. These results can be stated as follows:

**Theorem 2.3** (*Existence of Strong Solution For* (O)). Let p > 1,

$$f \in L^p(\Omega), \quad h \in W^{1,p}(\Omega), \quad v \in \mathbf{L}^s_{\sigma}(\Omega) \quad \text{and} \quad g \in \mathbf{W}^{2-1/p,p}(\Gamma)$$

satisfying the compatibility condition (3) with s = 2 if p < 2, s = p if p > 2 and  $s = 2 + \varepsilon$  ( $\varepsilon > 0$ ) if p = 2. Then, the unique solution of (O) given by Theorem 2.2  $(u, \pi)$  belongs to  $\mathbf{W}^{2,p}(\Omega) \times W^{1,p}(\Omega)$ , and there exists a constant C > 0 such that:

$$\begin{aligned} \|\boldsymbol{u}\|_{\mathbf{W}^{2,p}(\varOmega)} &+ \|\pi\|_{W^{1,p}(\varOmega)/\mathbb{R}} \leq C\left(1 + \|\boldsymbol{v}\|_{\mathbf{L}^{s}(\varOmega)}\right) \\ &\times \left(\|\boldsymbol{f}\|_{\mathbf{L}^{p}(\varOmega)} + \left(1 + \|\boldsymbol{v}\|_{\mathbf{L}^{s}(\varOmega)}\right)\left(\|h\|_{W^{1,p}(\varOmega)} + \|\boldsymbol{g}\|_{\mathbf{W}^{2-1/p,p}(\varGamma)}\right) \right) \end{aligned}$$

Theorem 2.4. Let

$$p > 1, \quad \boldsymbol{f} \in \mathbf{W}^{-1,p}(\Omega), \quad h \in L^p(\Omega), \quad \boldsymbol{v} \in \mathbf{L}^3_{\sigma}(\Omega) \quad \text{and} \quad \boldsymbol{g} \in \mathbf{W}^{1-1/p,p}(\Gamma)$$

satisfying the compatibility condition (3). Then, the problem (O) has a unique solution  $(\boldsymbol{u}, \pi) \in \mathbf{W}^{1,p}(\Omega) \times L^p(\Omega)/\mathbb{R}$ , and there exists a constant C > 0 such that:

$$\begin{aligned} \|\boldsymbol{u}\|_{\mathbf{W}^{1,p}(\Omega)} &+ \left(1 + \|\boldsymbol{v}\|_{\mathbf{L}^{3}(\Omega)}\right)^{\gamma} \|\boldsymbol{\pi}\|_{L^{p}(\Omega)/\mathbb{R}} \leq C \left(1 + \|\boldsymbol{v}\|_{\mathbf{L}^{3}(\Omega)}\right) \\ &\times \left(\|\boldsymbol{f}\|_{\mathbf{W}^{-1,p}(\Omega)} + \left(1 + \|\boldsymbol{v}\|_{\mathbf{L}^{3}(\Omega)}\right) \left(\|\boldsymbol{h}\|_{L^{p}(\Omega)} + \|\boldsymbol{g}\|_{\mathbf{W}^{1-1/p,p}(\Gamma)}\right)\right) \end{aligned}$$

with  $\gamma = 0$  if  $p \ge 2$  and  $\gamma = -1$  if p < 2.

If we treat the case of  $\Omega = \mathbb{R}^2$ , we have to introduce the Sobolev spaces:

$$W_0^{1,2}(\mathbb{R}^2) = \left\{ \varphi \in \mathcal{D}'(\mathbb{R}^2); \ \frac{\varphi}{w_1} \in L^2(\mathbb{R}^2), \ \nabla \varphi \in L^2(\mathbb{R}^2) \right\},$$
$$W_0^{2,2}(\mathbb{R}^2) = \left\{ \varphi \in \mathcal{D}'(\mathbb{R}^2); \ \frac{\varphi}{w_2} \in L^2(\mathbb{R}^2), \ \frac{\nabla \varphi}{w_1} \in L^2(\mathbb{R}^2), \ \nabla^2 \varphi \in L^2(\mathbb{R}^2) \right\}.$$

where  $w_1 = (1 + |\boldsymbol{x}|) \ln(2 + |\boldsymbol{x}|)$  and  $w_2 = (1 + |\boldsymbol{x}|)^2 \ln(2 + |\boldsymbol{x}|)$  (see Definition (7.1), p. 593 in [9]). We denote by  $W_0^{-1,2}(\mathbb{R}^2)$  the dual space of  $W_0^{1,2}(\mathbb{R}^2)$ . Recall [5] that the space  $W_0^{1,2}(\mathbb{R}^2)$  is densely embedded in  $VMO(\mathbb{R}^2)$ , and therefore  $\mathcal{H}^1(\mathbb{R}^2) = [VMO(\mathbb{R}^2)]' \hookrightarrow W_0^{-1,2}(\mathbb{R}^2)$ .

**Theorem 2.5** (*Case*  $\Omega = \mathbb{R}^2$ ). (i) Let

$$f = \operatorname{div} \mathbb{F}$$
 with  $\mathbb{F} \in \mathbf{L}^2(\mathbb{R}^2)$  and  $\mathbf{h} \in L^2(\mathbb{R}^2)$ .

Then, the problem (O) has a unique solution  $(\boldsymbol{u}, \pi)$  satisfying  $\boldsymbol{u} \in \mathbf{W}_0^{1,2}(\mathbb{R}^2)$  and  $\pi \in L^2(\mathbb{R}^2)$ , where  $\pi$  is unique and  $\boldsymbol{u}$  is unique up to an additive constant vector field.

(ii) Moreover, if

$$\boldsymbol{f} \in \mathcal{H}^1(\mathbb{R}^2)$$
 and  $\nabla h \in \mathcal{H}^1(\mathbb{R}^2),$ 

then

$$\nabla^2 \boldsymbol{u} \in \mathcal{H}^1(\mathbb{R}^2), \quad \nabla \pi \in \mathcal{H}^1(\mathbb{R}^2), \quad \nabla \boldsymbol{u} \in \boldsymbol{L}^{2,1}(\mathbb{R}^2) \quad and \quad \boldsymbol{u} \in L^{\infty}(\mathbb{R}^2), \tag{9}$$

being  $L^{2,1}(\mathbb{R}^2)$  is the Lorentz space of all measurable functions f satisfying

$$\int_0^\infty t^{-1/2} f^*(t) \, dt < +\infty$$

where the rearrangement function  $f^*$  is defined by  $f^*(t) = \sup\{s \in (0,\infty); \mu(\{\boldsymbol{x} \in \mathbb{R}^2; |f(\boldsymbol{x})| > s\}) > t\},$ for  $\mu$  the Lebesgue measure on  $\mathbb{R}^2$ .

**Proof.** (i) Existence: Let  $\chi \in W_0^{2,2}(\mathbb{R}^2)$  be the unique solution, up to a polynomial function of degree one, of  $\Delta \chi = h$  in  $\mathbb{R}^2$  (Theorem 9.6 in [9]). Then, we take  $u_h = \nabla \chi \in W_0^{1,2}(\mathbb{R}^2)$ . Problem (O) is then written as:

$$-\varDelta oldsymbol{z} + oldsymbol{v} \cdot 
abla oldsymbol{z} + 
abla \pi = oldsymbol{k}, \quad 
abla \cdot oldsymbol{z} = 0 \quad ext{in } \mathbb{R}^2,$$

with  $\mathbf{k} = \mathbf{f} + \Delta u_h - \mathbf{v} \cdot \nabla u_h$ . Because of  $(\mathbf{v} \cdot \nabla) u_h \in \mathcal{H}^1(\mathbb{R}^2) \hookrightarrow W_0^{-1,2}(\mathbb{R}^2)$ , by using Lax–Milgram's Lemma (as in the bounded case) we can deduce the existence of a solution  $\mathbf{z} \in \mathbf{W}_0^{1,2}(\mathbb{R}^2)$  with  $\nabla \cdot \mathbf{z} = 0$ , unique up to a constant vector of  $\mathbb{R}^2$ . Lax–Milgram's Lemma hypotheses are satisfied because, on the one hand, we know that the quotient norm  $\|\mathbf{z}\|_{\mathbf{W}_0^{1,2}(\mathbb{R}^2)/\mathbb{R}^2}$  is equivalent to that one defined as  $\|\nabla \mathbf{z}\|_{L^2(\mathbb{R}^2)}$ , and, on the other hand,  $(\mathbf{v} \cdot \nabla)\mathbf{z} \in \mathcal{H}^1(\mathbb{R}^2)$  for any  $\mathbf{z} \in \mathbf{W}_0^{1,2}(\mathbb{R}^2)$ . The pressure can be recovered by using Theorem 1 in [10].

(ii) Regularity: Assume that  $\mathbf{f} \in \mathcal{H}^1(\mathbb{R}^2)$  and  $\nabla h \in \mathcal{H}^1(\mathbb{R}^2)$  (which, in particular, implies that  $h \in L^{2,1}(\mathbb{R}^2)$ ). Therefore,

$$-\Delta \boldsymbol{u} + 
abla \pi = \boldsymbol{f} - \boldsymbol{v} \cdot 
abla \boldsymbol{u} \in \mathcal{H}^1(\mathbb{R}^2) \quad ext{and} \quad 
abla \cdot \boldsymbol{u} = h.$$

By using Theorem 3.14 in [10], one deduces (9).

### 3. The Oseen problem in bounded domains for a less regular v

The aim of this section is the analysis of the existence of solutions of (O) in a bounded domain (n = 2 or 3) when  $\boldsymbol{v} \in \mathbf{L}_{\sigma}^{s}(\Omega)$  for s < n. We analyze the case for  $\boldsymbol{f} \in \mathbf{H}^{-1}(\Omega)$ , h = 0 and  $\boldsymbol{g} = \mathbf{0}$ . Observe that the term  $(\boldsymbol{v} \cdot \nabla)\boldsymbol{u}$  can also be written as  $\nabla \cdot (\boldsymbol{u} \otimes \boldsymbol{v})$ . The proof of Theorem 3.2 (n = 2) applies directly from that one of Theorem 3.1 (n = 3).

## **Theorem 3.1.** Let $\Omega \subset \mathbb{R}^3$ a Lipschitz bounded domain,

$$\boldsymbol{f} \in \mathbf{H}^{-1}(\Omega), \quad h = 0, \quad \boldsymbol{g} = \mathbf{0} \quad \text{and} \quad \boldsymbol{v} \in \boldsymbol{L}_{\sigma}^{6/5+\alpha}(\Omega)$$

for any  $0 < \alpha \leq 9/5$ . Then, there exists a solution of (O) such that  $(\mathbf{u}, \pi) \in \mathbf{H}_0^1(\Omega) \times L^{q(\alpha)}(\Omega)/\mathbb{R}$  for  $q(\alpha) = (6(6+5\alpha))/(36+5\alpha)$  with the estimate:

$$\|\boldsymbol{u}\|_{\mathbf{H}^{1}(\Omega)} + \|\boldsymbol{\pi}\|_{L^{q(\alpha)}(\Omega)/\mathbb{R}} \leq C \left(1 + \|\boldsymbol{v}\|_{\mathbf{L}^{6/5+\alpha}(\Omega)}\right) \|\boldsymbol{f}\|_{\mathbf{H}^{-1}(\Omega)}$$
(10)

**Proof.** We approximate  $\boldsymbol{v}$  by  $\boldsymbol{v}_{\lambda} \in \boldsymbol{\mathcal{D}}_{\sigma}(\overline{\Omega})$  in the  $\mathbf{L}^{6/5+\alpha}(\Omega)$ -norm and look for the solution of the problem:

$$(O_{\lambda}) \quad -\Delta \boldsymbol{u}_{\lambda} + \nabla \cdot (\boldsymbol{u}_{\lambda} \otimes \boldsymbol{v}_{\lambda}) + \nabla \pi_{\lambda} = \boldsymbol{f} \quad \text{and} \quad \nabla \cdot \boldsymbol{u}_{\lambda} = 0 \quad \text{in } \boldsymbol{\Omega}, \quad \boldsymbol{u}_{\lambda} = \boldsymbol{0} \quad \text{on } \boldsymbol{I}$$

Taking  $u_{\lambda}$  as test function in  $(O_{\lambda})$ , we get the estimate:

$$\|\boldsymbol{u}_{\lambda}\|_{\mathbf{H}^{1}_{0}(\Omega)} \leq C(\Omega) \,\|\boldsymbol{f}\|_{\mathbf{H}^{-1}(\Omega)}.$$
(11)

By De Rham Theorem, there exists  $\pi_{\lambda} \in L^{2}(\Omega)$  (unique up to a constant) such that:

$$abla \pi_{\lambda} = oldsymbol{f} + arDeltaoldsymbol{u}_{\lambda} - 
abla \cdot (oldsymbol{u}_{\lambda} \otimes oldsymbol{v}_{\lambda}),$$

Moreover,  $\boldsymbol{u}_{\lambda} \otimes \boldsymbol{v}_{\lambda}$  belongs to a bounded set of  $\mathbb{L}^{q(\alpha)}(\Omega)$  with  $q(\alpha) = (6(6+5\alpha))/(36+5\alpha)$  and which implies that  $\nabla \cdot (\boldsymbol{u}_{\lambda} \otimes \boldsymbol{v}_{\lambda})$  belongs to a bounded set of  $\mathbf{W}^{-1,q(\alpha)}(\Omega)$ . Note that if  $0 < \alpha \leq 9/5$  then  $1 < q(\alpha) \leq 2$ . Using (11),

$$\begin{aligned} \|\nabla \pi_{\lambda}\|_{W^{-1,q(\alpha)}(\Omega)} &\leq C_{1}\left(1+C(\Omega)\right)\|f\|_{\mathbf{H}^{-1}(\Omega)} + \|\boldsymbol{u}_{\lambda} \otimes \boldsymbol{v}_{\lambda}\|_{\mathbb{L}^{q(\alpha)}(\Omega)} \\ &\leq C_{1}\left(1+C(\Omega)\right)\|f\|_{\mathbf{H}^{-1}(\Omega)} + C_{2}\|\boldsymbol{v}_{\lambda}\|_{\mathbf{L}^{6/5+\alpha}(\Omega)}\|\boldsymbol{u}_{\lambda}\|_{\mathbf{H}^{1}_{0}(\Omega)} \\ &\leq C(\Omega)\left(1+\|\boldsymbol{v}\|_{\mathbf{L}^{6/5+\alpha}(\Omega)}\right)\|f\|_{\mathbf{H}^{-1}(\Omega)} \end{aligned}$$
(12)

where  $C_1$  and  $C_2$  are the constant of the Sobolev embeddings  $\mathbf{H}^{-1}(\Omega) \hookrightarrow \mathbf{W}^{-1,q(\alpha)}(\Omega)$  and  $\mathbf{H}^1_0(\Omega) \hookrightarrow \mathbf{L}^6(\Omega)$ , respectively. Therefore, from (12) we obtain:

$$\inf_{K \in \mathbb{R}} \|\pi_{\lambda} + K\|_{L^{q(\alpha)}(\Omega)} \le C(\Omega) \left(1 + \|\boldsymbol{v}\|_{\mathbf{L}^{6/5 + \alpha}(\Omega)}\right) \|f\|_{\mathbf{H}^{-1}(\Omega)}$$

Now, it is necessary to take the limit when  $\lambda \to 0$ : We can extract a subsequence of  $(u_{\lambda})$  and  $(\pi_{\lambda} + C_{\lambda})$  (that will be called in the same way that the original one) such that:

 $u_{\lambda} \rightharpoonup u \quad \text{in } \mathbf{H}_{0}^{1}(\Omega), \qquad \pi_{\lambda} + C_{\lambda} \rightharpoonup \pi \quad \text{in } L^{q(\alpha)}(\Omega),$ 

where  $(u, \pi)$  is solution of (O) and satisfies (10).  $\Box$ 

**Theorem 3.2.** Let  $\Omega \subset \mathbb{R}^2$  a Lipschitz bounded domain,  $\mathbf{f} \in \mathbf{H}^{-1}(\Omega)$ , h = 0,  $\mathbf{g} = \mathbf{0}$  and  $\mathbf{v} \in \mathbf{L}^{1+\alpha}_{\sigma}(\Omega)$  with  $0 < \alpha \leq 1$ . Then, there exists a solution of (O) such that  $(\mathbf{u}, \pi) \in \mathbf{H}^{1}_{0}(\Omega) \times L^{q(\beta)}(\Omega)/\mathbb{R}$  for  $q(\beta) = 1 + \beta$ , for any  $0 < \beta < \alpha$ , with the estimate:

$$\|\boldsymbol{u}\|_{\mathbf{H}^{1}(\Omega)} + \|\boldsymbol{\pi}\|_{L^{q(\beta)}(\Omega)/\mathbb{R}} \leq C \left(1 + \|\boldsymbol{v}\|_{\mathbf{L}^{1+\alpha}(\Omega)}\right) \|\boldsymbol{f}\|_{\mathbf{H}^{-1}(\Omega)}$$

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