# On the logistic equation for the fractional $\boldsymbol{p}$-Laplacian 

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#### Abstract

We consider a Dirichlet problem for a nonlinear, nonlocal equation driven by the degenerate fractional $p$-Laplacian, with a logistic-type reaction depending on a positive parameter. In the subdiffusive and equidiffusive cases, we prove existence and uniqueness of the positive solution when the parameter lies in convenient intervals. In the superdiffusive case, we establish a bifurcation result. A new strong comparison result, of independent interest, plays a crucial role in the proof of such bifurcation result.


## KEYWORDS

bifurcation, comparison principle, fractional $p$-Laplacian, logistic equation

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## 1 | INTRODUCTION

The paper is devoted to the study of the following nonlinear elliptic equation of fractional order with Dirichlet-type condition:

$$
\left(P_{\lambda}\right) \begin{cases}(-\Delta)_{p}^{s} u=\lambda u^{q-1}-u^{r-1} & \text { in } \Omega \\ u>0 & \text { in } \Omega \\ u=0 & \text { in } \mathbb{R}^{N} \backslash \Omega .\end{cases}
$$

Here, $\Omega \subset \mathbb{R}^{N}(N \geqslant 2)$ is a bounded domain with $C^{1,1}$ boundary $\partial \Omega, s \in(0,1), p \geqslant 2$ are s.t. $p s<N$, and the leading operator is the degenerate fractional $p$-Laplacian, defined for all $u: \mathbb{R}^{N} \rightarrow \mathbb{R}$ smooth enough and $x \in \mathbb{R}^{N}$ by

$$
(-\Delta)_{p}^{s} u(x)=2 \lim _{\varepsilon \searrow 0} \int_{\{|x-y|>\varepsilon\}} \frac{|u(x)-u(y)|^{p-2}(u(x)-u(y))}{|x-y|^{N+p s}} d y
$$

[^0](which for $p=2$ reduces to the linear fractional Laplacian, up to a dimensional constant $C(N, s)>0)$. The reaction is of logistic type, with powers $1<q<r<p_{s}^{*}$, where $p_{s}^{*}=N p /(N-p s)$ denotes the critical exponent for fractional Sobolev spaces, and $\lambda>0$ is a parameter. Problem $\left(P_{\lambda}\right)$ is classified in three different cases, according to the principal exponent $q>1$ :
(a) subdiffusive, if $q<p<r$;
(b) equidiffusive, if $p=q<r$;
(c) superdiffusive, if $p<q<r$.

Logistic equations are widely studied mainly because of their important applications in mathematical biology. Indeed, the parabolic semilinear logistic equation describes the evolution and spatial distribution of a biological population in the presence of constant rates of reproduction and mortality (Verhulst's law), see [17]. This is the obvious reason why, in the study of logistic-type equations, authors are usually interested in positive solutions. More recently, evolutive systems involving logistic terms have been studied as a model for the biological phenomenon of chemotaxis [37], and existence of solutions in the presence of a parameter was studied in [1, 7]. Regarding the elliptic counterpart, it models an equilibrium distribution, see [10]. Several existence results for the equidiffusive case (b), combining variational and topological methods, can be found in $[2,3,36]$ (note that multiplicity often includes negative and nodal solutions). Bifurcation results for the superdiffusive case (c) can be found in [23] for the Dirichlet problem, and in [29] for the whole space.

Fractional order equations also have a close connection to mathematical biology. Indeed, since fractional elliptic operators model space diffusion via Lévy-type random motion with jumps, they can be effectively used to describe the movement of populations, see [4, 31]. Studies on logistic equations with several nonlocal operators of fractional order have appeared in recent years, including the square root of the Dirichlet Laplacian [8], the spectral Neumann fractional Laplacian [28], and the fractional Laplacian on the whole space [35].

The operator we consider here is both nonlinear and nonlocal. It represents the nonlinear generalization of the fractional Laplacian, and it can be seen as the gradient of the functional $u \mapsto[u]_{s, p}^{p} / p$ in the fractional Sobolev space $W^{s, p}\left(\mathbb{R}^{N}\right)$ (see Section 2), as first pointed out in [5]. The corresponding eigenvalue problem was studied in [26], which led to existence results for general nonlinear reactions via Morse theory in [18]. Due to the nature of the operator, regularity theory for weak solutions required a considerable effort as most of the usual techniques (including the Caffarelli-Silvestre extension method) do not apply here. For any $p>1$, Hölder continuity of weak solutions in the interior and up to the boundary was studied in [14] and [20], respectively.

In the degenerate case $p>2$, optimal interior Hölder regularity was proved in [6], while a weighted global Hölder regularity result was proved in [21] (the singular case $p \in(1,2)$ is still open). The result of [21] is the fractional counterpart of Lieberman's $C^{1, \alpha}$-regularity result for the classical $p$-Laplacian [25] and yields many applications, such as the equivalence of Sobolev and Hölder local minimizers of the energy functional [22], the existence of extremal constant sign solutions [16], and more recently a Brezis-Oswald-type weak comparison principle [19]. We also recall other interesting related results, such as the study of critical growth and singularity performed in [9] and the bifurcation results of [12, 32]. For further information, we refer the reader to the surveys [27, 30].

As far as we know, the present literature includes no specific study on the logistic equation for the fractional $p$-Laplacian. This paper aims at filling the gap, by presenting the following general result for the existence of solutions to problem $\left(P_{\lambda}\right)$ (in which $\hat{\lambda}_{1}>0$ denotes the principal eigenvalue of $(-\Delta)_{p}^{s}$ in $\Omega$ with Dirichlet conditions, see Equation (2.4)):

Theorem 1.1. Let $\Omega \subset \mathbb{R}^{N}$ be a bounded domain with $C^{1,1}$-boundary, $p \geqslant 2$, $s \in(0,1)$ s.t. $p s<N$, and $1<q<r<p_{s}^{*}$. Then, the following hold:
(a) if $q<p$, then for all $\lambda>0$ problem $\left(P_{\lambda}\right)$ has a unique solution $u_{\lambda}>0$, with $u_{\lambda}>u_{\mu}$ in $\Omega$ for all $\lambda>\mu>0$ and $u_{\lambda} \rightarrow 0$ as $\lambda \searrow 0$;
(b) if $q=p$, then for all $\lambda \in\left(0, \hat{\lambda}_{1}\right.$ ] problem $\left(P_{\lambda}\right)$ has no solution, while for all $\lambda>\hat{\lambda}_{1}\left(P_{\lambda}\right)$ has a unique solution $u_{\lambda}>0$, with $u_{\lambda}>u_{\mu}$ in $\Omega$ for all $\lambda>\mu>\hat{\lambda}_{1}$ and $u_{\lambda} \rightarrow 0$ as $\lambda \searrow \hat{\lambda}_{1}$;
(c) if $q>p$, then there exists $\lambda_{*}>0$ s.t. for all $\lambda \in\left(0, \lambda_{*}\right)$ problem $\left(P_{\lambda}\right)$ has no solution, while $\left(P_{\lambda_{*}}\right)$ has at least one solution $u_{*}>0$, and for all $\lambda>\lambda_{*}\left(P_{\lambda}\right)$ has at least two solutions $u_{\lambda}>v_{\lambda}>0$, with $u_{\lambda}>u_{\mu}$ in $\Omega$ for all $\lambda>\mu>\lambda_{*}$ and $u_{\lambda} \rightarrow u_{*}$ as $\lambda \searrow \lambda_{*}$.

More precise statements of the results above can be found in Theorems 3.1, 3.2, and 3.7. Our approach is variational, based on critical point theory and comparison-truncation arguments. For the sub- and equidiffusive cases, we apply direct minimization and the weak comparison result of [19] for uniqueness. In the superdiffusive case, we prove a bifurcation result and detect via the mountain pass theorem a second solution for all $\lambda>\lambda_{*}$.

We remark that our result is new even in the semilinear case $p=2$ (fractional Laplacian) and in the local case $s=1$ (classical $p$-Laplacian). Bifurcation theorems are proved in [8] for the superdiffusive logistic equation driven by the square root of the Laplacian, and in [23] for the classical p-Laplacian, but with no information about monotonicity, order between solutions, and convergence. Also, existence and uniqueness for the equidiffusive case with the fractional Laplacian are proved in [35].

A crucial role in our arguments is played by new strong minimum and comparison principles for weak sub- and supersolutions, including a Hopf-type property (see Theorems 2.6 and 2.7). Previous results of this type were proved in [13, 24], respectively, but our versions involve very general reactions and milder restrictions on the constants $p, s$ and can be of general interest, since they are applicable to a wide class of problems driven by the fractional $p$-Laplacian.

Structure of the paper: in Section 2, we recall some preliminary results (Section 2.1) and prove new minimum and comparison principles (Section 2.2); in Section 3, we deal with the logistic equation, distinguishing between the subdiffusive case (Section 3.1), the equidiffusive case (Section 3.2), and the superdiffusive case (Section 3.3).

Notation: For any $a \in \mathbb{R}, \nu>0$ we set $a^{\nu}=|a|^{\nu-1} a$. For any $A \subset \mathbb{R}^{N}$ we shall set $A^{c}=\mathbb{R}^{N} \backslash A$ and denote by $|A|$ the Lebesgue measure of $A$. For any two measurable functions $u, v: \Omega \rightarrow \mathbb{R}, u \leqslant v$ will mean that $u(x) \leqslant v(x)$ for a.e. $x \in \Omega$ (and similar expressions). The positive (resp., negative) part of $u$ is denoted as $u^{+}$(resp., $u^{-}$). Every function $u$ defined in $\Omega$ will be identified with its 0 -extension to $\mathbb{R}^{N}$. If $X$ is an ordered function space, then $X_{+}$will denote its non-negative order cone. For all $\nu \in[1, \infty],\|\cdot\|_{\nu}$ denotes the standard norm of $L^{\nu}(\Omega)\left(\right.$ or $L^{\nu}\left(\mathbb{R}^{N}\right)$, which will be clear from the context). Moreover, $C$ will denote a positive constant whose value may change case by case.

## 2 | PRELIMINARIES

Problem $\left(P_{\lambda}\right)$ falls into the following class of Dirichlet problems for the fractional $p$-Laplacian:

$$
\begin{cases}(-\Delta)_{p}^{s} u=f(x, u) & \text { in } \Omega  \tag{2.1}\\ u=0 & \text { in } \Omega^{c} .\end{cases}
$$

Here, $\Omega \subseteq \mathbb{R}^{N}(N \geqslant 2)$ is a bounded domain with $C^{1,1}$ boundary $\partial \Omega$, $s \in(0,1), p>1$ satisfy $p s<N$. Besides, the general reaction $f$ satisfies the following hypothesis:
$\mathbf{H} f: \Omega \times \mathbb{R} \rightarrow \mathbb{R}$ is a Carathéodory function s.t. for a.e. $x \in \Omega$ and all $t \in \mathbb{R}$

$$
|f(x, t)| \leqslant c_{0}\left(1+|t|^{r-1}\right) \quad\left(c_{0}>0, r \in\left(p, p_{s}^{*}\right)\right)
$$

In this section, we will collect some old and new properties of the solutions of problem (2.1).

## 2.1 | Variational formulation and properties of solutions

A variational theory for problem (2.1) was established in the recent literature (see, for instance, [16, 18, 22]). For the reader's convenience, we recall here some of its main features. First, for all measurable $u: \Omega \rightarrow \mathbb{R}$, we introduce the Gagliardo seminorm

$$
[u]_{s, p, \Omega}=\left[\iint_{\Omega \times \Omega} \frac{|u(x)-u(y)|^{p}}{|x-y|^{N+p s}} d x d y\right]^{\frac{1}{p}}
$$

setting $[u]_{s, p, \mathbb{R}^{N}}=[u]_{s, p}$. Then, we define the fractional Sobolev spaces

$$
\begin{gathered}
W^{s, p}(\Omega)=\left\{u \in L^{p}(\Omega):[u]_{s, p, \Omega}<\infty\right\}, \\
W_{0}^{s, p}(\Omega)=\left\{u \in W^{s, p}\left(\mathbb{R}^{N}\right): u=0 \text { in } \Omega^{c}\right\},
\end{gathered}
$$

the latter being a uniformly convex, separable Banach space with norm $\|u\|=[u]_{s, p}$, whose dual space is denoted by $W^{-s, p^{\prime}}(\Omega)$ (see [15]). The embedding $W_{0}^{s, p}(\Omega) \hookrightarrow L^{\nu}(\Omega)$ is continuous for all $v \in\left[1, p_{s}^{*}\right]$ and compact for all $\nu \in\left[1, p_{s}^{*}\right)$. We also recall from [20, Definition 2.1] the following special space:

$$
\widetilde{W^{s, p}}(\Omega)=\left\{u \in L_{\mathrm{loc}}^{p}\left(\mathbb{R}^{N}\right): \exists U \ni \Omega \text { s.t. } u \in W^{s, p}(U), \int_{\mathbb{R}^{N}} \frac{|u(x)|^{p-1}}{(1+|x|)^{N+p s}} d x<\infty\right\} .
$$

By [20, Lemma 2.3], we can define the fractional $p$-Laplacian as a nonlinear operator $(-\Delta)_{p}^{s}: \widetilde{W}^{s, p}(\Omega) \rightarrow W^{-s, p^{\prime}}(\Omega)$ by setting for all $u, v \in W_{0}^{s, p}(\Omega)$

$$
\left\langle(-\Delta)_{p}^{s} u, v\right\rangle=\iint_{\mathbb{R}^{N} \times \mathbb{R}^{N}} \frac{(u(x)-u(y))^{p-1}(v(x)-v(y))}{|x-y|^{N+p s}} d x d y
$$

(with the convention $a^{p-1}=|a|^{p-2} a$ established above). Such definition is equivalent to the one given in Section 1 as soon as $u$ is smooth enough (for instance, if $u \in S\left(\mathbb{R}^{N}\right)$ ).

Clearly $W_{0}^{s, p}(\Omega) \subset \widetilde{W}^{s, p}(\Omega)$. Also, whenever $u \in \widetilde{W}^{s, p}(\Omega)$ satisfies $u=0$ in $\Omega^{c}$, it is easily seen that $u \in W_{0}^{s, p}(\Omega)$. The restricted operator $(-\Delta)_{p}^{s}: W_{0}^{s, p}(\Omega) \rightarrow W^{-s, p^{\prime}}(\Omega)$ is continuous, maximal monotone, and enjoys the $(S)_{+}$-property, namely, whenever $\left(u_{n}\right)$ is a sequence in $W_{0}^{s, p}(\Omega)$ s.t. $u_{n} \rightharpoonup u$ in $W_{0}^{s, p}(\Omega)$ and

$$
\limsup _{n}\left\langle(-\Delta)_{p}^{s} u_{n}, u_{n}-u\right\rangle \leqslant 0,
$$

then $u_{n} \rightarrow u$ in $W_{0}^{s, p}(\Omega)$ (see [16, Lemma 2.1] for $p \geqslant 2$, with analogous argument for $p \in(1,2)$ ). For all $u \in W_{0}^{s, p}(\Omega)$, we have

$$
\begin{equation*}
\left\|u^{ \pm}\right\|^{p} \leqslant\left\langle(-\Delta)_{p}^{s} u, \pm u^{ \pm}\right\rangle . \tag{2.2}
\end{equation*}
$$

Another useful property, referred to as strict $T$-monotonicity, of $(-\Delta)_{p}^{s}$ is the following, which holds for any $p>1$ (see [26, proof of Lemma 9]):

Proposition 2.1. Let $u, v \in \widetilde{W}^{s, p}(\Omega)$ s.t. $(u-v)^{+} \in W_{0}^{s, p}(\Omega)$ satisfy

$$
\left\langle(-\Delta)_{p}^{s} u-(-\Delta)_{p}^{s} v,(u-v)^{+}\right\rangle \leqslant 0 .
$$

Then, $u \leqslant v$ in $\Omega$.
We say that $u \in \widetilde{W^{s, p}}(\Omega)$ is a (weak) supersolution of Equation (2.1) if $u \geqslant 0$ in $\Omega^{c}$ and for all $v \in W_{0}^{s, p}(\Omega)_{+}$

$$
\left\langle(-\Delta)_{p}^{s} u, v\right\rangle \geqslant \int_{\Omega} f(x, u) v d x
$$

and similarly we define a (weak) subsolution. Finally, $u \in W_{0}^{s, p}(\Omega)$ is a (weak) solution of Equation (2.1) if it is both a super- and a subsolution, that is, if for all $v \in W_{0}^{s, p}(\Omega)$

$$
\left\langle(-\Delta)_{p}^{s} u, v\right\rangle=\int_{\Omega} f(x, u) v d x .
$$

In such cases, we write that weakly in $\Omega$

$$
(-\Delta)_{p}^{s} u=(\geqslant, \leqslant) f(x, u) .
$$

From [9, Theorem 3.3], we have the following a priori bound on the solutions:
Proposition 2.2. Let $\mathbf{H}$ hold, $u \in W_{0}^{s, p}(\Omega)$ be a solution of Equation (2.1). Then, $u \in L^{\infty}(\Omega)$ with $\|u\|_{\infty} \leqslant C(\|u\|)$.
Classical nonlinear regularity theory does not apply to fractional order equations, whose solutions fail to be $C^{1}$ in general. Nevertheless, weighted Hölder continuity can replace higher smoothness in most cases. We set $\mathrm{d}_{\Omega}(x)=\operatorname{dist}\left(x, \Omega^{c}\right)$ for all $x \in \mathbb{R}^{N}$ and define the following space:

$$
C_{s}^{0}(\bar{\Omega})=\left\{u \in C^{0}(\bar{\Omega}): \frac{u}{\mathrm{~d}_{\Omega}^{s}} \text { has a continuous extension to } \bar{\Omega}\right\}
$$

a Banach space under the norm $\|u\|_{0, s}=\left\|u / \mathrm{d}_{\Omega}^{s}\right\|_{\infty}$. By [18, Lemma 5.1], the positive order cone $C_{s}^{0}(\bar{\Omega})_{+}$has a nonempty interior

$$
\operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)=\left\{u \in C_{s}^{0}(\bar{\Omega}): \inf _{\Omega} \frac{u}{\mathrm{~d}_{\Omega}^{s}}>0\right\}
$$

Similarly, for any $\alpha \in(0,1)$ we set

$$
C_{s}^{\alpha}(\bar{\Omega})=\left\{u \in C^{0}(\bar{\Omega}): \frac{u}{\mathrm{~d}_{\Omega}^{s}} \text { has a } \alpha \text {-Hölder continuous extension to } \bar{\Omega}\right\}
$$

a Banach space under the norm

$$
\|u\|_{\alpha, s}=\|u\|_{0, s}+\sup _{x, y \in \Omega, x \neq y} \frac{\left|u(x) / \mathrm{d}_{\Omega}^{s}(x)-u(y) / \mathrm{d}_{\Omega}^{s}(y)\right|}{|x-y|^{\alpha}} .
$$

By the Ascoli-Arzelà theorem, $C_{s}^{\alpha}(\bar{\Omega}) \hookrightarrow C_{S}^{0}(\bar{\Omega})$ with compact embedding for all $\alpha \in(0,1)$. From Proposition 2.2 , [20, Theorem 1.1], and [21, Theorem 1.1] we have the following weighted Hölder regularity result:

Proposition 2.3. Let $\mathbf{H}$ hold, $u \in W_{0}^{s, p}(\Omega)$ be a solution of Equation (2.1). Then, there exists $\alpha \in(0, s]$, independent of $u$, s.t. $u \in C^{\alpha}(\bar{\Omega})$. Besides, if $p \geqslant 2$, then $u \in C_{S}^{\alpha}(\bar{\Omega})$ and $\|u\|_{\alpha, s} \leqslant C(\|u\|)$.

Weighted Hölder continuity is known only for the degenerate case $p \geqslant 2$. This is the main reason why the next results, which use such type of regularity, are only stated for $p \geqslant 2$. From [19, Proposition 2.8], we have the following weak comparison principle under a special monotonicity assumption of Brezis-Oswald type:

Proposition 2.4. Let $\mathbf{H}$ hold, $p \geqslant 2$, and assume that

$$
t \mapsto \frac{f(x, t)}{t^{p-1}}
$$

is decreasing in $(0, \infty)$ for a.e. $x \in \Omega$. Let $u, v \in \operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right) \cap W_{0}^{s, p}(\Omega)$ be a subsolution and a supersolution, respectively, of Equation (2.1). Then, $u \leqslant v$ in $\Omega$.

The energy functional for problem (2.1) is defined by setting for all $u \in W_{0}^{s, p}(\Omega)$

$$
\Phi(u)=\frac{\|u\|^{p}}{p}-\int_{\Omega} F(x, u) d x
$$

where we have set for $\operatorname{all}(x, t) \in \Omega \times \mathbb{R}$

$$
F(x, t)=\int_{0}^{t} f(x, \tau) d \tau
$$

By classical results, we have $\Phi \in C^{1}\left(W_{0}^{s, p}(\Omega)\right)$, and $u \in W_{0}^{s, p}(\Omega)$ is a solution of Equation (2.1) iff $\Phi^{\prime}(u)=0$ in $W^{-s, p^{\prime}}(\Omega)$. Besides, by [18, Proposition 2.1] $\Phi$ satisfies a bounded $(P S)$-condition, namely, whenever $\left(u_{n}\right)$ is a bounded sequence
in $W_{0}^{s, p}(\Omega)$ s.t. $\left(\Phi\left(u_{n}\right)\right)$ is bounded in $\mathbb{R}$ and $\Phi^{\prime}\left(u_{n}\right) \rightarrow 0$ in $W^{-s, p^{\prime}}(\Omega)$, then $\left(u_{n}\right)$ has a convergent subsequence. In this connection, we recall from [22, Theorem 1.1] the following equivalence principle for Sobolev and Hölder local minimizers of $\Phi$ :

Proposition 2.5. Let $\mathbf{H}$ hold, $p \geqslant 2, u \in W_{0}^{s, p}(\Omega)$. Then, the following are equivalent:
(i) there exists $\rho>0$ s.t. $\Phi(u+v) \geqslant \Phi(u)$ for all $v \in W_{0}^{s, p}(\Omega) \cap C_{s}^{0}(\bar{\Omega}),\|v\|_{0, s} \leqslant \rho$;
(ii) there exists $\sigma>0$ s.t. $\Phi(u+v) \geqslant \Phi(u)$ for all $v \in W_{0}^{s, p}(\Omega),\|v\| \leqslant \sigma$.

Regarding the spectral properties of the fractional $p$-Laplacian, we refer the reader to [26]. We just recall that the eigenvalue problem is stated as

$$
\begin{cases}(-\Delta)_{p}^{s} u=\lambda u^{p-1} & \text { in } \Omega  \tag{2.3}\\ u=0 & \text { in } \Omega .\end{cases}
$$

The principal eigenvalue $\hat{\lambda}_{1}>0$ of Equation (2.3) is simple and isolated, with a unique positive eigenfunction $\hat{u}_{1} \in$ $\operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)$s.t. $\|u\|_{p}=1$, and both are defined as follows:

$$
\begin{equation*}
\hat{\lambda}_{1}=\inf _{u \in W_{0}^{S, p}(\Omega) \backslash\{0\}} \frac{\|u\|^{p}}{\|u\|_{p}^{p}}=\left\|\hat{u}_{1}\right\|^{p} . \tag{2.4}
\end{equation*}
$$

## 2.2 | Strong minimum and comparison principles

As mentioned in Section 1, a strong minimum principle and a Hopf-type lemma for the fractional $p$-Laplacian were proved in [13, Theorems 1.2, 1.5], while a strong comparison principle was obtained in [24, Theorem 1.1]. Nevertheless, the strong comparison principle of [24] does not fit with our purposes for two reasons: first, in the degenerate case $p>2$ it requires some special relations between the parameters $p$ and $s$ which, combined with the optimal Hölder continuity proved in [6], lead to the quite restrictive condition $s \leqslant 1 / p^{\prime}$; second, the result only ensures that the difference between the superand the subsolution is positive in $\Omega$, while we need to prove that such difference lies in $\operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)$.

Motivated by such difficulties, we present here a new pair of results, following an alternative approach based on the nonlocal superposition principle introduced in [21]. In view of future applications, we will prove such results for any $p>1$.

We begin with a strong minimum principle (including a Hopf-type boundary property):
Theorem 2.6. Let $\Omega \subset \mathbb{R}^{N}$ be a bounded domain with $C^{1,1}$ boundary, $p>1, s \in(0,1)$ s.t. $p s<N, g \in C^{0}(\mathbb{R}) \cap B V_{\text {loc }}(\mathbb{R})$, $u \in \widetilde{W}^{s, p}(\Omega) \cap C^{0}(\bar{\Omega}), u \not \equiv 0$ s.t.

$$
\begin{cases}(-\Delta)_{p}^{s} u+g(u) \geqslant g(0) & \text { weakly in } \Omega \\ u \geqslant 0 & \text { in } \mathbb{R}^{N} .\end{cases}
$$

Then,

$$
\inf _{\Omega} \frac{u}{\mathrm{~d}_{\Omega}^{s}}>0 .
$$

In particular, if $u \in C_{s}^{0}(\bar{\Omega})$, then $u \in \operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)$.
Proof. By Jordan's decomposition, we can find $g_{1}, g_{2} \in C^{0}(\mathbb{R})$ nondecreasing s.t. $g(t)=g_{1}(t)-g_{2}(t)$ for all $t \in \mathbb{R}$, and $g_{1}(0)=0$. So, we have weakly in $\Omega$

$$
\begin{aligned}
(-\Delta)_{p}^{s} u+g_{1}(u) & =(-\Delta)_{p}^{s} u+g(u)+g_{2}(u) \\
& \geqslant g(0)+g_{2}(0)=0 .
\end{aligned}
$$

Thus, without loss of generality we may assume that $g$ is nondecreasing and $g(0)=0$. In order to prove our assertion, we need a lower barrier for $u$. Let us consider the following torsion problem:

$$
\begin{cases}(-\Delta)_{p}^{s} v=1 & \text { in } \Omega  \tag{2.5}\\ v=0 & \text { in } \Omega^{c}\end{cases}
$$

By convexity, Equation (2.5) has a unique solution $v \in W_{0}^{s, p}(\Omega)$, which by [21, Lemma 2.3] satisfies $v \geqslant c \mathrm{~d}_{\Omega}^{s}$ in $\Omega$, for some $c>0$. By Proposition 2.3, we have $v \in C^{\alpha}(\bar{\Omega})$, in particular $v$ is continuous. So, since $u \not \equiv 0$, we can find $x_{0} \in \Omega, \rho, \varepsilon>0$, and $\eta_{0} \in(0,1)$ s.t. $\bar{B}_{\rho}\left(x_{0}\right) \subset \Omega$ and

$$
\begin{equation*}
\sup _{\bar{B}_{\rho}\left(x_{0}\right)} \eta_{0} v<\inf _{\bar{B}_{\rho}\left(x_{0}\right)} u-\varepsilon . \tag{2.6}
\end{equation*}
$$

Set for all $x \in \mathbb{R}^{N}, \eta \in\left(0, \eta_{0}\right]$

$$
w_{\eta}(x)= \begin{cases}\eta v(x) & \text { if } x \in \bar{B}_{\rho / 2}^{c}\left(x_{0}\right) \\ u(x) & \text { if } x \in \bar{B}_{\rho / 2}\left(x_{0}\right)\end{cases}
$$

First, by Equation (2.6) we have $w_{\eta} \leqslant u$ in $\bar{B}_{\rho}\left(x_{0}\right)$. Besides, by the nonlocal superposition principle [21, Proposition 2.6] we have $w_{\eta} \in \widetilde{W}^{s, p}\left(\Omega \backslash \bar{B}_{\rho}\left(x_{0}\right)\right)$ and weakly in $\Omega \backslash \bar{B}_{\rho}\left(x_{0}\right)$

$$
\begin{aligned}
(-\Delta)_{p}^{s} w_{\eta}(x) & =(-\Delta)_{p}^{s}(\eta v)(x)+2 \int_{\bar{B}_{\rho / 2}\left(x_{0}\right)} \frac{(\eta v(x)-u(y))^{p-1}-(\eta v(x)-\eta v(y))^{p-1}}{|x-y|^{N+p s}} d y \\
& \leqslant \eta^{p-1}+2 \int_{\bar{B}_{\rho / 2}\left(x_{0}\right)} \frac{(\eta v(x)-u(y))^{p-1}-(\eta v(x)-u(y)+\varepsilon)^{p-1}}{|x-y|^{N+p s}} d y
\end{aligned}
$$

where we have also used Equation (2.5) and again Inequality (2.6). Now, by continuity we can find $C_{\varepsilon}>0$, independent of $\eta \in\left(0, \eta_{0}\right]$, s.t. for all $x \in \Omega \backslash \bar{B}_{\rho}\left(x_{0}\right), y \in \bar{B}_{\rho / 2}\left(x_{0}\right)$

$$
(\eta v(x)-u(y))^{p-1}-(\eta v(x)-u(y)+\varepsilon)^{p-1} \leqslant-C_{\varepsilon}
$$

and $C_{\varepsilon} \rightarrow 0$ as $\varepsilon \searrow 0$. So, we have weakly in $\Omega \backslash \bar{B}_{\rho}\left(x_{0}\right)$

$$
(-\Delta)_{p}^{s} w_{\eta}(x) \leqslant \eta^{p-1}-2 \int_{\bar{B}_{\rho / 2}\left(x_{0}\right)} \frac{C_{\varepsilon}}{(\rho / 2)^{N+p s}} d y \leqslant \eta^{p-1}-\tilde{C}_{\varepsilon}
$$

with $\tilde{C}_{\varepsilon}>0$ independent of $\eta$. Choosing $\eta \in\left(0, \eta_{0}\right]$ small enough, we have weakly in $\Omega \backslash \bar{B}_{\rho}\left(x_{0}\right)$

$$
(-\Delta)_{p}^{S} w_{\eta}(x) \leqslant-\frac{\tilde{C}_{\varepsilon}}{2}
$$

Note that $g\left(w_{\eta}\right) \rightarrow 0$ uniformly in $\Omega \backslash \bar{B}_{\rho}\left(x_{0}\right)$ as $\eta \searrow 0$. So, for an even smaller $\eta \in\left(0, \eta_{0}\right]$ we have

$$
\begin{cases}(-\Delta)_{p}^{s} w_{\eta}+g\left(w_{\eta}\right) \leqslant 0 \leqslant(-\Delta)_{p}^{s} u+g(u) & \text { weakly in } \Omega \backslash \bar{B}_{\rho}\left(x_{0}\right) \\ w_{\eta} \leqslant u & \text { in }\left(\Omega \backslash \bar{B}_{\rho}\left(x_{0}\right)\right)^{c}\end{cases}
$$

We have $\left(w_{\eta}-u\right)^{+} \in \widetilde{W}^{s, p}\left(\Omega \backslash \bar{B}_{\rho}\left(x_{0}\right)\right)$ and, by the second inequality above, $\left(w_{\eta}-u\right)^{+}=0$ in $\left(\Omega \backslash \bar{B}_{\rho}\left(x_{0}\right)^{c}\right.$, hence $\left(w_{\eta}-u\right)^{+} \in W_{0}^{s, p}\left(\Omega \backslash \bar{B}_{\rho}\left(x_{0}\right)\right)$. So, we can employ such function to test the inequality above. We get

$$
\left\langle(-\Delta)_{p}^{s} w_{\eta}-(-\Delta)_{p}^{s} u,\left(w_{\eta}-u\right)^{+}\right\rangle \leqslant \int_{\Omega \backslash \bar{B}_{\rho}\left(x_{0}\right)}\left(g(u)-g\left(w_{\eta}\right)\right)\left(w_{\eta}-u\right)^{+} d x
$$

and the latter is negative by the monotonicity of $g$. By Proposition 2.1, we have $w_{\eta} \leqslant u$ in $\Omega \backslash \bar{B}_{\rho}\left(x_{0}\right)$. Combining with Inequality (2.6) we get in $\Omega$

$$
u \geqslant \eta v \geqslant \eta c \mathrm{~d}_{\Omega}^{s}
$$

hence the conclusion. In particular, if $u \in C_{S}^{0}(\bar{\Omega})$, then clearly we have $u \in \operatorname{int}\left(C_{S}^{0}(\bar{\Omega})_{+}\right)$.

With a similar technique, we prove a strong comparison principle:
Theorem 2.7. Let $\Omega \subset \mathbb{R}^{N}$ be a bounded domain with $C^{1,1}$ boundary, $p>1$, $s \in(0,1)$ s.t. $p s<N, g \in C^{0}(\mathbb{R}) \cap B V_{\text {loc }}(\mathbb{R})$, $u \in \widetilde{W}^{s, p}(\Omega) \cap C^{0}(\bar{\Omega}), v \in W_{0}^{s, p}(\Omega) \cap C^{0}(\bar{\Omega})$ s.t. $u \not \equiv v, K>0$ satisfy

$$
\begin{cases}(-\Delta)_{p}^{s} v+g(v) \leqslant(-\Delta)_{p}^{s} u+g(u) \leqslant K & \text { weakly in } \Omega \\ 0<v \leqslant u & \text { in } \Omega \\ u \geqslant 0 & \text { in } \Omega^{c}\end{cases}
$$

Then, $u>v$ in $\Omega$. In particular, if $u, v \in \operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)$, then $u-v \in \operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)$.
Proof. As in Theorem 2.6, we may assume $g$ nondecreasing. By continuity, we can find $x_{0} \in \Omega, \rho, \varepsilon>0$ s.t. $\bar{B}_{\rho}\left(x_{0}\right) \subset \Omega$ and

$$
\sup _{\bar{B}_{\rho}\left(x_{0}\right)} v<\inf _{\bar{B}_{\rho}\left(x_{0}\right)} u-\varepsilon
$$

Hence, for all $\eta \in(1,2)$ close enough to 1 we have

$$
\begin{equation*}
\sup _{\bar{B}_{\rho}\left(x_{0}\right)} \eta v<\inf _{\bar{B}_{\rho}\left(x_{0}\right)} u-\frac{\varepsilon}{2} \tag{2.7}
\end{equation*}
$$

Define $w_{\eta} \in \widetilde{W}^{s, p}\left(\Omega \backslash \bar{B}_{\rho}\left(x_{0}\right)\right)$ as in Theorem 2.6 , so by Inequality (2.7) we have $w_{\eta} \leqslant u$ in $\left(\Omega \backslash \bar{B}_{\rho}\left(x_{0}\right)\right)^{c}$. Applying nonlocal superposition as in the previous proof, we have weakly in $\Omega \backslash \bar{B}_{\rho}\left(x_{0}\right)$

$$
(-\Delta)_{p}^{s} w_{\eta} \leqslant \eta^{p-1}(-\Delta)_{p}^{s} v-C_{\varepsilon}
$$

for some $C_{\varepsilon}>0$ independent of $\eta \in(1,2)$. Further, we have weakly in $\Omega \backslash \bar{B}_{\rho}\left(x_{0}\right)$

$$
\begin{aligned}
(-\Delta)_{p}^{s} w_{\eta}+g\left(w_{\eta}\right) & \leqslant \eta^{p-1}(-\Delta)_{p}^{s} v+g\left(w_{\eta}\right)-C_{\varepsilon} \\
& \leqslant \eta^{p-1}\left((-\Delta)_{p}^{s} v+g(v)\right)+\left(g\left(w_{\eta}\right)-\eta^{p-1} g(v)\right)-C_{\varepsilon} \\
& \leqslant \eta^{p-1}\left((-\Delta)_{p}^{s} u+g(u)\right)+\left(g\left(w_{\eta}\right)-\eta^{p-1} g(v)\right)-C_{\varepsilon} \\
& \leqslant(-\Delta)_{p}^{s} u+g(u)+K\left(\eta^{p-1}-1\right)+\left(g\left(w_{\eta}\right)-\eta^{p-1} g(v)\right)-C_{\varepsilon}
\end{aligned}
$$

where we have used the hypothesis and the monotonicity of $g$. Since

$$
K\left(\eta^{p-1}-1\right)+\left(g\left(w_{\eta}\right)-\eta^{p-1} g(v)\right) \rightarrow 0
$$

uniformly in $\Omega \backslash \bar{B}_{\rho}\left(x_{0}\right)$ as $\eta \searrow 1$, we have for all $\eta>1$ close enough to 1

$$
\begin{cases}(-\Delta)_{p}^{s} w_{\eta}+g\left(w_{\eta}\right) \leqslant(-\Delta)_{p}^{s} u+g(u) & \text { weakly in } \Omega \backslash \bar{B}_{\rho}\left(x_{0}\right) \\ w_{\eta} \leqslant u & \text { in }\left(\Omega \backslash \bar{B}_{\rho}\left(x_{0}\right)\right)^{c}\end{cases}
$$

Testing with $\left(w_{\eta}-u\right)^{+} \in W_{0}^{s, p}\left(\Omega \backslash \bar{B}_{\rho}\left(x_{0}\right)\right)$, recalling the monotonicity of $g$, and applying Proposition 2.1 we get $u \geqslant w_{\eta}$ in $\Omega \backslash \bar{B}_{\rho}\left(x_{0}\right)$. So we have in $\Omega$

$$
u \geqslant \eta v>v
$$

hence the conclusion. In particular, if $u, v \in \operatorname{int}\left(C_{S}^{0}(\bar{\Omega})_{+}\right)$, then clearly

$$
\inf _{\Omega} \frac{u-v}{\mathrm{~d}_{\Omega}^{S}} \geqslant \inf _{\Omega} \frac{(\eta-1) v}{\mathrm{~d}_{\Omega}^{S}}>0
$$

so $u-v \in \operatorname{int}\left(C_{S}^{0}(\bar{\Omega})_{+}\right)$.

Remark 2.8. Both results above exhibit unexpected differences when compared to the corresponding local versions, that is, the case of the classical $p$-Laplacian. For example, according to Theorem 2.6, the strong minimum principle holds for non-negative supersolutions of the Dirichlet problem

$$
\begin{cases}(-\Delta)_{p}^{s} u+u^{\sigma}=0 & \text { in } \Omega \\ u=0 & \text { in } \Omega^{c}\end{cases}
$$

for any $\sigma>0$, while for $s=1$ the same is not true when $\sigma<p-1$ due to the possible presence of dead cores (see [34, p. 204]). Also, the strong comparison principle of Theorem 2.7 includes cases which are excluded in the local case (see [11, Example 4.1]). This is essentially due to the nonlocal nature of the operator.

## 3 | THE LOGISTIC EQUATION

In this section, we study problem $\left(P_{\lambda}\right)$ with $\Omega \subset \mathbb{R}^{N}(N \geqslant 2)$ bounded domain with a $C^{1,1}$ boundary, $p \geqslant 2$, $s \in(0,1)$ s.t. $p s<N$, and $1<q<r<p_{s}^{*}$. For all $\lambda>0, t \in \mathbb{R}$, we set

$$
\begin{gathered}
f_{\lambda}(t)=\lambda\left(t^{+}\right)^{q-1}-\left(t^{+}\right)^{r-1} \\
F_{\lambda}(t)=\int_{0}^{t} f_{\lambda}(\tau) d \tau=\lambda \frac{\left(t^{+}\right)^{q}}{q}-\frac{\left(t^{+}\right)^{r}}{r} .
\end{gathered}
$$

Note that $f_{\lambda}: \mathbb{R} \rightarrow \mathbb{R}$ satisfies hypotheses $\mathbf{H}$ as stated in Section 2. So we may set for all $u \in W_{0}^{s, p}(\Omega)$

$$
\begin{equation*}
\Phi_{\lambda}(u)=\frac{\|u\|^{p}}{p}-\int_{\Omega} F_{\lambda}(u) d x \tag{3.1}
\end{equation*}
$$

and deduce that $\Phi_{\lambda} \in C^{1}\left(W_{0}^{s, p}(\Omega)\right)$. As we will see, the positive critical points of $\Phi_{\lambda}$ coincide with the solutions of $\left(P_{\lambda}\right)$. In the following subsections, we separately study the different cases according to the position of $q$.

## 3.1 | The subdiffusive case

We assume $1<q<p<r<p_{s}^{*}$. In this case, we have the following global existence and uniqueness result (corresponding to case ( $a$ ) of Theorem 1.1):

Theorem 3.1. Let $1<q<p<r<p_{s}^{*}$. Then, for all $\lambda>0$ problem $\left(P_{\lambda}\right)$ has a unique solution $u_{\lambda} \in \operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)$, s.t. $u_{\lambda}-u_{\mu} \in \operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)$for all $\lambda>\mu>0$ and $u_{\lambda} \rightarrow 0$ in both $W_{0}^{s, p}(\Omega)$ and $C_{S}^{0}(\bar{\Omega})$ as $\lambda \searrow 0$.

Proof. Fix any $\lambda>0$. We will find the solution of $\left(P_{\lambda}\right)$ by direct minimization. First, we prove that the functional $\Phi_{\lambda}$ (defined in Equation (3.1)) is coercive. Indeed, since $q<r$, the mapping $F_{\lambda}$ is clearly bounded from above, that is, there
exists $C>0$ s.t. $F_{\lambda}(t) \leqslant C$ for all $t \in \mathbb{R}$. So, for all $u \in W_{0}^{s, p}(\Omega)$ we have

$$
\Phi_{\lambda}(u) \geqslant \frac{\|u\|^{p}}{p}-C|\Omega|,
$$

and the latter tends to $\infty$ as $\|u\| \rightarrow \infty$. Besides, by the compact embeddings $W_{0}^{s, p}(\Omega) \hookrightarrow L^{q}(\Omega), L^{r}(\Omega)$, it is easily seen that $\Phi_{\lambda}$ is sequentially weakly lower semicontinuous in $W_{0}^{s, p}(\Omega)$. So, there exists $u_{\lambda} \in W_{0}^{s, p}(\Omega)$ s.t.

$$
\begin{equation*}
\Phi_{\lambda}\left(u_{\lambda}\right)=\inf _{W_{0}^{s, p}(\Omega)} \Phi_{\lambda}=: m_{\lambda} . \tag{3.2}
\end{equation*}
$$

Besides, let $\hat{u}_{1} \in \operatorname{int}\left(C_{S}^{0}(\bar{\Omega})_{+}\right)$be defined by Equation (2.4). Then, for all $\tau>0$

$$
\Phi_{\lambda}\left(\tau \hat{u}_{1}\right)=\tau^{p} \frac{\left\|\hat{u}_{1}\right\|^{p}}{p}-\lambda \tau^{q} \frac{\left\|\hat{u}_{1}\right\|_{q}^{q}}{q}+\tau^{r} \frac{\left\|\hat{u}_{1}\right\|_{r}^{r}}{r},
$$

and the latter is negative for all $\tau>0$ small enough (recall that $q<p<r$ ). So, in Equation (3.2) we have $m_{\lambda}<0$, implying $u_{\lambda} \not \equiv 0$. From Equation (3.2), we deduce that $\Phi_{\lambda}^{\prime}\left(u_{\lambda}\right)=0$ in $W^{-s, p^{\prime}}(\Omega)$, that is, we have weakly in $\Omega$

$$
\begin{equation*}
(-\Delta)_{p}^{s} u_{\lambda}=f_{\lambda}\left(u_{\lambda}\right) \tag{3.3}
\end{equation*}
$$

By Proposition 2.3, we have $u_{\lambda} \in C_{s}^{\alpha}(\bar{\Omega})$. Besides, testing Equation (3.3) with $-u_{\lambda}^{-} \in W_{0}^{s, p}(\Omega)$ and applying Equation (2.2), we have

$$
\left\|u_{\lambda}^{-}\right\|^{p} \leqslant\left\langle(-\Delta)_{p}^{s} u_{\lambda},-u_{\lambda}^{-}\right\rangle=\int_{\Omega} f_{\lambda}\left(u_{\lambda}\right)\left(-u_{\lambda}^{-}\right) d x=0,
$$

so $u_{\lambda} \geqslant 0$. Now, Theorem 2.6 implies $u_{\lambda} \in \operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)$, so $u_{\lambda}$ solves $\left(P_{\lambda}\right)$.
Next, we prove uniqueness. Let $v_{\lambda} \in \operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)$be another solution of $\left(P_{\lambda}\right)$. We have for all $t>0$

$$
\frac{f_{\lambda}(t)}{t^{p-1}}=\lambda t^{q-p}-t^{r-p}
$$

and such mapping is decreasing in $(0, \infty)$. Applying Proposition 2.4 twice, we have $u_{\lambda}=v_{\lambda}$.
To see monotonicity, let $0<\mu<\lambda$, and $u_{\mu}, u_{\lambda} \in \operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)$be the solutions of $\left(P_{\mu}\right),\left(P_{\lambda}\right)$, respectively. We have weakly in $\Omega$

$$
(-\Delta)_{p}^{s} u_{\mu}<\lambda u_{\mu}^{q-1}-u_{\mu}^{r-1}
$$

so $u_{\mu}$ is a strict subsolution of $\left(P_{\lambda}\right)$. By Proposition 2.4 again, we have $u_{\mu} \leqslant u_{\lambda}$ in $\Omega$. This in turn implies that weakly in $\Omega$

$$
(-\Delta)_{p}^{s} u_{\mu}+u_{\mu}^{r-1}=\mu u_{\mu}^{q-1}<\lambda u_{\lambda}^{q-1}=(-\Delta)_{p}^{s} u_{\lambda}+u_{\lambda}^{r-1} .
$$

Since $g(t)=t^{r-1}$ is continuous and with locally bounded variation, we can apply Theorem 2.7 and see that $u_{\lambda}-u_{\mu} \in$ $\operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)$.

Finally, let $\left(\lambda_{n}\right)$ be a decreasing sequence in $(0, \infty)$ s.t. $\lambda_{n} \searrow 0$, and $u_{n} \in \operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)$be the solution of $\left(P_{\lambda_{n}}\right)$ for all $n \in \mathbb{N}$, that is, we have weakly in $\Omega$

$$
\begin{equation*}
(-\Delta)_{p}^{s} u_{n}=f_{\lambda_{n}}\left(u_{n}\right) . \tag{3.4}
\end{equation*}
$$

Since $q<p$ and $\left(\lambda_{n}\right)$ is decreasing, we can find $C>0$ s.t. for all $n \in \mathbb{N}, t \in \mathbb{R}$

$$
f_{\lambda_{n}}(t) t \leqslant C .
$$

Testing Equation (3.4) with $u_{n} \in W_{0}^{s, p}(\Omega)$, for all $n \in \mathbb{N}$ we have

$$
\left\|u_{n}\right\|^{p}=\left\langle(-\Delta)_{p}^{s} u_{n}, u_{n}\right\rangle=\int_{\Omega} f_{\lambda_{n}}\left(u_{n}\right) u_{n} d x \leqslant C|\Omega| .
$$

So, $\left(u_{n}\right)$ is a bounded sequence in $W_{0}^{s, p}(\Omega)$. By reflexivity and the compact embeddings $W_{0}^{s, p}(\Omega) \hookrightarrow L^{q}(\Omega), L^{r}(\Omega)$, we can pass to a subsequence s.t. $u_{n} \rightharpoonup u_{0}$ in $W_{0}^{s, p}(\Omega)$ and $u_{n} \rightarrow u_{0}$ in both $L^{q}(\Omega)$ and $L^{r}(\Omega)$. Testing Equation (3.4) with $\left(u_{n}-u_{0}\right) \in W_{0}^{S, p}(\Omega)$ and using Hölder's inequality, we have for all $n \in \mathbb{N}$

$$
\begin{aligned}
\left\langle(-\Delta)_{p}^{s} u_{n}, u_{n}-u_{0}\right\rangle & =\int_{\Omega}\left(\lambda_{n} u_{n}^{q-1}-u_{n}^{r-1}\right)\left(u_{n}-u_{0}\right) d x \\
& \leqslant \lambda_{1}\left\|u_{n}\right\|_{q}^{q-1}\left\|u_{n}-u_{0}\right\|_{q}+\left\|u_{n}\right\|_{r}^{r-1}\left\|u_{n}-u_{0}\right\|_{r}
\end{aligned}
$$

and the latter tends to 0 as $n \rightarrow \infty$. By the $(S)_{+}$-property of $(-\Delta)_{p}^{s}$, we have $u_{n} \rightarrow u_{0}$ in $W_{0}^{s, p}(\Omega)$. So, we can pass to the limit in Equation (3.4) as $n \rightarrow \infty$ and get weakly in $\Omega$

$$
(-\Delta)_{p}^{s} u_{0}=-u_{0}^{r-1}
$$

Testing with $u_{0} \in W_{0}^{s, p}(\Omega)$ we have

$$
\left\|u_{0}\right\|^{p}+\left\|u_{0}\right\|_{r}^{r}=0
$$

that is, $u_{0}=0$. Plus, we note that, by Equation (3.4) and Proposition 2.3, $\left(u_{n}\right)$ is bounded in $C_{s}^{\alpha}(\bar{\Omega})$, hence, passing to a further subsequence, $u_{n} \rightarrow 0$ in $C_{s}^{0}(\bar{\Omega})$. Recalling that $\lambda \mapsto u_{\lambda}$ is strictly increasing, we conclude that globally $u_{\lambda} \rightarrow 0$ in both $W_{0}^{s, p}(\Omega)$ and $C_{s}^{0}(\bar{\Omega})$, as $\lambda \searrow 0$.

## 3.2 | The equidiffusive case

Now, we assume $2 \leqslant q=p<r<p_{s}^{*}$, a case that does not differ too much from the previous one, except that the threshold for the parameter $\lambda$ turns out to be the principal eigenvalue $\hat{\lambda}_{1}>0$ defined in Equation (2.4). Our existence and uniqueness result (corresponding to case (b) of Theorem 1.1) is the following:

Theorem 3.2. Let $2 \leqslant q=p<r<p_{s}^{*}$. Then, for all $\lambda \in\left(0, \hat{\lambda}_{1}\right]$ problem $\left(P_{\lambda}\right)$ has no solution, while for all $\lambda>\hat{\lambda}_{1}$ problem $\left(P_{\lambda}\right)$ has a unique solution $u_{\lambda} \in \operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)$, s.t. $u_{\lambda}-u_{\mu} \in \operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)$for all $\lambda>\mu>\hat{\lambda}_{1}$ and $u_{\lambda} \rightarrow 0$ in both $W_{0}^{s, p}(\Omega)$ and $C_{s}^{0}(\bar{\Omega})$ as $\lambda \searrow \lambda_{1}$.

Proof. First, fix $\lambda \in\left(0, \hat{\lambda}_{1}\right]$. Assume that $u \in W_{0}^{s, p}(\Omega)_{+}$satisfies weakly in $\Omega$

$$
\begin{equation*}
(-\Delta)_{p}^{s} u=\lambda u^{p-1}-u^{r-1} \tag{3.5}
\end{equation*}
$$

Testing Equation (3.5) with $u \in W_{0}^{s, p}(\Omega)$ and applying Equation (2.4), we have

$$
0=\|u\|^{p}-\lambda\|u\|_{p}^{p}+\|u\|_{r}^{r} \geqslant\left(\hat{\lambda}_{1}-\lambda\right)\|u\|_{p}^{p}+\|u\|_{r}^{r} \geqslant\|u\|_{r}^{r}
$$

hence $u=0$. So $\left(P_{\lambda}\right)$ admits no solution.
Now, let $\lambda>\hat{\lambda}_{1}$, and define $\Phi_{\lambda}$ as in Equation (3.1). Arguing as in Theorem 3.1, we see that $\Phi_{\lambda}$ has a global minimizer $u_{\lambda} \in W_{0}^{s, p}(\Omega)_{+}$. Besides, let $\hat{u}_{1} \in \operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)$be as in Equation (2.4). Then, for all $\tau>0$ we have

$$
\begin{aligned}
\Phi_{\lambda}\left(\tau \hat{u}_{1}\right) & =\tau^{p}\left[\frac{\left\|\hat{u}_{1}\right\|^{p}}{p}-\lambda \frac{\left\|\hat{u}_{1}\right\|_{p}^{p}}{p}\right]+\tau^{r} \frac{\left\|\hat{u}_{1}\right\|_{r}^{r}}{r} \\
& =\tau^{p} \frac{\hat{\lambda}_{1}-\lambda}{p}+\tau^{r} \frac{\left\|\hat{u}_{1}\right\|_{r}^{r}}{r}
\end{aligned}
$$

and the latter is negative for $\tau>0$ small enough (as $p<r$ ). So, $u_{\lambda} \not \equiv 0$. The rest of the proof follows exactly as in Theorem 3.1.

## 3.3 | The superdiffusive case

In this final case, we assume $2 \leqslant p<q<r<p_{s}^{*}$ and define $\Phi_{\lambda}$ as in Equation(3.1). We will need a more accurate analysis. Let

$$
\Lambda=\left\{\lambda>0:\left(P_{\lambda}\right) \text { has a solution } u_{\lambda} \in \operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)\right\}
$$

In the following lemmas, we shall investigate the structure of the set $\Lambda$ and additional properties of solutions. We begin with a lower bound for $\Lambda$ :

Lemma 3.3. We have $\Lambda \neq \emptyset$ and $\lambda_{*}:=\inf \Lambda>0$.

Proof. Fix $\lambda>0$. As in the proof of Theorem 3.1, we find $u_{\lambda} \in W_{0}^{s, p}(\Omega)_{+}$s.t.

$$
\begin{equation*}
\Phi_{\lambda}\left(u_{\lambda}\right)=\inf _{W_{0}^{s, p}(\Omega)} \Phi_{\lambda}=: m_{\lambda} \tag{3.6}
\end{equation*}
$$

Let $\hat{u}_{1} \in \operatorname{int}\left(C_{S}^{0}(\bar{\Omega})_{+}\right)$be as in Equation (2.4), then we have

$$
\Phi_{\lambda}\left(\hat{u}_{1}\right)=\frac{\left\|\hat{u}_{1}\right\|^{p}}{p}-\lambda \frac{\left\|\hat{u}_{1}\right\|_{q}^{q}}{q}+\frac{\left\|\hat{u}_{1}\right\|_{r}^{r}}{r}
$$

which tends to $-\infty$ as $\lambda \rightarrow \infty$. So, for all $\lambda>0$ big enough we have $m_{\lambda}<0$ in Equation (3.6), hence $u_{\lambda} \neq 0$. As in Theorem 3.1 we see that $u_{\lambda} \in \operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)$and it solves $\left(P_{\lambda}\right)$. Thus, we have $\Lambda \neq \emptyset$.

We claim that there exists $\lambda_{0}>0$ s.t. for all $t \geqslant 0$

$$
\begin{equation*}
f_{\lambda_{0}}(t) \leqslant \hat{\lambda}_{1} t^{p-1} \tag{3.7}
\end{equation*}
$$

with $\hat{\lambda}_{1}>0$ defined by Equation (2.4). Indeed, since $p<q<r$ we have for any $\lambda>0$

$$
\lim _{t \searrow 0} \frac{f_{\lambda}(t)}{t^{p-1}}=0, \quad \lim _{t \rightarrow \infty} \frac{f_{\lambda}(t)}{t^{p-1}}=-\infty
$$

So, we can find $\delta \in(0,1)$ s.t. for all $t \in(0, \delta) \cup\left(\delta^{-1}, \infty\right)$ and all $\lambda \in(0,1]$

$$
f_{\lambda}(t) \leqslant \hat{\lambda}_{1} t^{p-1} .
$$

Now, set

$$
\lambda_{0}=\min \left\{\hat{\lambda}_{1} \delta^{q-p}, 1\right\}>0 .
$$

Then, for all $t \in\left[\delta, \delta^{-1}\right]$ we have

$$
f_{\lambda_{0}}(t)<\lambda_{0} t^{q-1} \leqslant \hat{\lambda}_{1} t^{p-1}
$$

hence Inequality (3.7) holds for all $t \geqslant 0$. We prove that inf $\Lambda \geqslant \lambda_{0}$, arguing by contradiction. Assume that for some $\lambda \in$ ( $0, \lambda_{0}$ ) problem $\left(P_{\lambda}\right)$ has a solution $u_{\lambda} \in \operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)$. Testing with $u_{\lambda} \in W_{0}^{s, p}(\Omega)$ and using Equation (3.7), we get

$$
\left\|u_{\lambda}\right\|^{p}=\int_{\Omega} f_{\lambda}\left(u_{\lambda}\right) u_{\lambda} d x<\int_{\Omega} f_{\lambda_{0}}\left(u_{\lambda}\right) u_{\lambda} d x \leqslant \hat{\lambda}_{1}\left\|u_{\lambda}\right\|_{p}^{p}
$$

against the characterization of $\hat{\lambda}_{1}$ in Equation (2.4).

Next, we prove that $\Lambda$ is a half-line and the mapping $\lambda \mapsto u_{\lambda}$ is strictly increasing:

Lemma 3.4. If $\lambda>\lambda_{*}$ then $\lambda \in \Lambda$. Besides, for all $\lambda>\mu>\lambda_{*}$, if $u_{\lambda}, u_{\mu} \in \operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)$are the solutions of $\left(P_{\lambda}\right)$, $\left(P_{\mu}\right)$ respectively, then $u_{\lambda}-u_{\mu} \in \operatorname{int}\left(C_{S}^{0}(\bar{\Omega})_{+}\right)$.

Proof. Fix $\lambda>\lambda_{*}$. Then, we can find $\mu \in \Lambda$ s.t. $\mu<\lambda$, and a solution $u_{\mu} \in \operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)$of $\left(P_{\mu}\right)$. We have weakly in $\Omega$

$$
\begin{equation*}
(-\Delta)_{p}^{s} u_{\mu}=f_{\mu}\left(u_{\mu}\right)<f_{\lambda}\left(u_{\mu}\right), \tag{3.8}
\end{equation*}
$$

that is, $u_{\mu}$ is a strict subsolution of $\left(P_{\lambda}\right)$. We use $u_{\mu}$ to truncate the reaction $f_{\lambda}$. Set for all $(x, t) \in \Omega \times \mathbb{R}$

$$
\hat{f}_{\lambda}(x, t)= \begin{cases}f_{\lambda}\left(u_{\mu}(x)\right) & \text { if } t \leqslant u_{\mu}(x) \\ f_{\lambda}(t) & \text { if } t>u_{\mu}(x)\end{cases}
$$

and

$$
\hat{F}_{\lambda}(x, t)=\int_{0}^{t} \hat{f}_{\lambda}(x, \tau) d \tau
$$

So $\hat{f}_{\lambda}: \Omega \times \mathbb{R} \rightarrow \mathbb{R}$ satisfies $\mathbf{H}$. Set for all $u \in W_{0}^{s, p}(\Omega)$

$$
\hat{\Phi}_{\lambda}(u)=\frac{\|u\|^{p}}{p}-\int_{\Omega} \hat{F}_{\lambda}(x, u) d x,
$$

then as in Section 2 it is seen that $\hat{\Phi}_{\lambda} \in C^{1}\left(W_{0}^{s, p}(\Omega)\right)$. Reasoning as in Theorem 3.1 we also see that $\hat{\Phi}_{\lambda}$ is coercive and sequentially weakly l.s.c., so there exists $u_{\lambda} \in W_{0}^{s, p}(\Omega)$ s.t.

$$
\hat{\Phi}_{\lambda}\left(u_{\lambda}\right)=\inf _{W_{0}^{s, p}(\Omega)} \hat{\Phi}_{\lambda} .
$$

As a consequence, we have $\hat{\Phi}_{\lambda}^{\prime}\left(u_{\lambda}\right)=0$ in $W^{-s, p^{\prime}}(\Omega)$, that is, weakly in $\Omega$

$$
\begin{equation*}
(-\Delta)_{p}^{s} u_{\lambda}=\hat{f}_{\lambda}(x, u) . \tag{3.9}
\end{equation*}
$$

Testing Equation (3.9) with $\left(u_{\mu}-u_{\lambda}\right)^{+} \in W_{0}^{s, p}(\Omega)_{+}$we get

$$
\begin{aligned}
\left\langle(-\Delta)_{p}^{s} u_{\lambda},\left(u_{\mu}-u_{\lambda}\right)^{+}\right\rangle & =\int_{\Omega} \hat{f}_{\lambda}\left(x, u_{\lambda}\right)\left(u_{\mu}-u_{\lambda}\right)^{+} d x \\
& =\int_{\Omega} f_{\lambda}\left(u_{\mu}\right)\left(u_{\mu}-u_{\lambda}\right)^{+} d x
\end{aligned}
$$

which along with Equation (3.8) gives

$$
\left\langle(-\Delta)_{p}^{s} u_{\mu}-(-\Delta)_{p}^{s} u_{\lambda},\left(u_{\mu}-u_{\lambda}\right)^{+}\right\rangle \leqslant 0 .
$$

By Proposition 2.1, we have $u_{\mu} \leqslant u_{\lambda}$ in $\Omega$. So, Equation (3.9) rephrases as

$$
(-\Delta)_{p}^{s} u_{\lambda}=f_{\lambda}\left(u_{\lambda}\right)
$$

weakly in $\Omega$, and moreover $u_{\lambda}>0$ in $\Omega$. As in Lemma 3.3 we see that $u_{\lambda} \in \operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)$and it solves $\left(P_{\lambda}\right)$, so $\lambda \in \Lambda$.
Finally, for all $\lambda>\mu>\lambda_{*}$ we have $u_{\lambda}, u_{\mu} \in W_{0}^{s, p}(\Omega) \cap C_{s}^{0}(\bar{\Omega})$ and

$$
\begin{cases}(-\Delta)_{p}^{s} u_{\mu}+u_{\mu}^{r-1}=\mu u_{\mu}^{q-1}<\lambda u_{\lambda}^{q-1}=(-\Delta)_{p}^{s} u_{\lambda}+u_{\lambda}^{r-1} & \text { weakly in } \Omega \\ 0<u_{\mu} \leqslant u_{\lambda} & \text { in } \Omega .\end{cases}
$$

By Theorem 2.7, we conclude that $u_{\lambda}-u_{\mu} \in \operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)$.
Note that in Lemma 3.4 we cannot use Proposition 2.4 to prove the monotonicity of $\lambda \mapsto u_{\lambda}$, as we did in sub- and equidiffusive cases: this is due to the fact that $t \mapsto f_{\lambda}(t) / t^{p-1}$ is not a decreasing mapping in $(0, \infty)$ (recall that $q>p$ ). The same reason prevents the use of Proposition 2.4 to prove uniqueness of the solution.

In fact, for $\lambda>\lambda_{*}$ we detect at least one more solution beside $u_{\lambda}$ :
Lemma 3.5. For all $\lambda>\lambda_{*}$ there exists a second solution $v_{\lambda} \in \operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)$of $\left(P_{\lambda}\right)$ s.t. $u_{\lambda}-v_{\lambda} \in \operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)$.
Proof. Fix $\lambda>\lambda_{*}$. As in Lemma 3.4 we pick $\mu \in \Lambda$ s.t. $\lambda_{*}<\mu<\lambda$, define $\hat{\Phi}_{\lambda} \in C^{1}\left(W_{0}^{s, p}(\Omega)\right)$, and find a global minimizer $u_{\lambda} \in \operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)$, which solves $\left(P_{\lambda}\right)$ and satisfies $u_{\lambda}-u_{\mu} \in \operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)$. Set now

$$
V=\left\{u_{\mu}+v: v \in \operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)\right\},
$$

which is an open set in $C_{s}^{0}(\bar{\Omega})$ containing $u_{\lambda}$. For all $x \in \Omega, t>u_{\mu}(x)$, we have

$$
\begin{aligned}
\hat{F}_{\lambda}(x, t) & =\int_{0}^{u_{\mu}(x)} f_{\lambda}\left(u_{\mu}(x)\right) d \tau+\int_{u_{\mu}(x)}^{t} f_{\lambda}(\tau) d \tau \\
& =F_{\lambda}(t)+\left[f_{\lambda}\left(u_{\mu}(x)\right) u_{\mu}(x)-F_{\lambda}\left(u_{\mu}(x)\right)\right],
\end{aligned}
$$

hence for all $u \in V \cap W_{0}^{s, p}(\Omega)$ (note that $u>u_{\mu}$ in $\Omega$ )

$$
\hat{\Phi}_{\lambda}(u)=\frac{\|u\|^{p}}{p}-\int_{\Omega} F_{\lambda}(u) d x-\int_{\Omega}\left[f_{\lambda}\left(u_{\mu}\right) u_{\mu}-F_{\lambda}\left(u_{\mu}\right)\right] d x=\Phi_{\lambda}(u)-C,
$$

with $C \in \mathbb{R}$ independent of $u$. So, recalling that $u_{\lambda}$ minimizes $\hat{\Phi}_{\lambda}$ over $W_{0}^{s, p}(\Omega)$, for all $u \in V \cap W_{0}^{s, p}(\Omega)$ we have

$$
\Phi_{\lambda}(u) \geqslant \Phi_{\lambda}\left(u_{\lambda}\right),
$$

that is, $u_{\lambda}$ is a local minimizer of $\Phi_{\lambda}$ in $C_{s}^{0}(\bar{\Omega})$. By Proposition $2.5, u_{\lambda}$ is as well a local minimizer of $\Phi_{\lambda}$ in $W_{0}^{s, p}(\Omega)$. To proceed with the proof, we need to perform a different truncation on the reaction. Set for all ( $x, t) \in \Omega \times \mathbb{R}$

$$
\tilde{f}_{\lambda}(x, t)= \begin{cases}f_{\lambda}(t) & \text { if } t \leqslant u_{\lambda}(x) \\ \lambda u_{\lambda}^{q-1}(x)-t^{r-1} & \text { if } t>u_{\lambda}(x)\end{cases}
$$

and as usual

$$
\tilde{F}_{\lambda}(x, t)=\int_{0}^{t} \tilde{f}_{\lambda}(x, \tau) d \tau .
$$

Clearly $\tilde{f}_{\lambda}: \Omega \times \mathbb{R} \rightarrow \mathbb{R}$ satisfies $\mathbf{H}$. So, we set for all $u \in W_{0}^{s, p}(\Omega)$

$$
\tilde{\Phi}_{\lambda}(u)=\frac{\|u\|^{p}}{p}-\int_{\Omega} \tilde{F}_{\lambda}(x, u) d x
$$

and thus define a functional $\tilde{\Phi}_{\lambda} \in C^{1}\left(W_{0}^{s, p}(\Omega)\right)$. We note that for all $(x, t) \in \Omega \times \mathbb{R}$ we have $\tilde{f}_{\lambda}(x, t) \leqslant f_{\lambda}(t)$ and hence $\tilde{F}_{\lambda}(x, t) \leqslant F_{\lambda}(t)$. This in turn implies for all $u \in W_{0}^{s, p}(\Omega)$

$$
\begin{equation*}
\tilde{\Phi}_{\lambda}(u) \geqslant \Phi_{\lambda}(u) . \tag{3.10}
\end{equation*}
$$

Since $u_{\lambda}$ is a local minimizer of $\Phi_{\lambda}$, we can find $\rho>0$ s.t. $\Phi_{\lambda}(u) \geqslant \Phi_{\lambda}\left(u_{\lambda}\right)$ for all $u \in B_{\rho}\left(u_{\lambda}\right)$, hence by Inequality (3.10)

$$
\tilde{\Phi}_{\lambda}(u) \geqslant \Phi_{\lambda}(u) \geqslant \Phi_{\lambda}\left(u_{\lambda}\right)=\tilde{\Phi}_{\lambda}\left(u_{\lambda}\right) .
$$

So, $u_{\lambda}$ is as well a local minimizer of $\tilde{\Phi}_{\lambda}$. Besides, fix $\varepsilon \in\left(0, \hat{\lambda}_{1}\right)$ (with $\hat{\lambda}_{1}>0$ defined by Equation (2.4)), then we can find $\delta>0$ s.t. for all $x \in \mathbb{R},|t| \leqslant \delta$

$$
\tilde{F}_{\lambda}(x, t) \leqslant F_{\lambda}(t) \leqslant \varepsilon \frac{\left(t^{+}\right)^{p}}{p} .
$$

Since $\Omega$ is bounded, we can find $\sigma>0$ s.t. $\|u\|_{\infty} \leqslant \delta$ for all $u \in C_{s}^{0}(\bar{\Omega}),\|u\|_{0, s} \leqslant \sigma$. Then, using also Equation (2.4), for all $u \in W_{0}^{s, p}(\Omega) \cap C_{s}^{0}(\bar{\Omega})$ with $0<\|u\|_{0, s} \leqslant \sigma$ we have

$$
\tilde{\Phi}_{\lambda}(u) \geqslant \frac{\|u\|^{p}}{p}-\int_{\Omega} \varepsilon \frac{\left(u^{+}\right)^{p}}{p} d x \geqslant\left(\hat{\lambda}_{1}-\varepsilon\right) \frac{\|u\|_{p}^{p}}{p}>0
$$

So, 0 is a strict local minimizer of $\tilde{\Phi}_{\lambda}$ in $C_{s}^{0}(\bar{\Omega})$. By Proposition 2.5 again, 0 is as well a local minimizer of $\tilde{\Phi}_{\lambda}$ in $W_{0}^{s, p}(\Omega)$. From Lemma 3.3, we know that $\Phi_{\lambda}$ is coercive in $W_{0}^{s, p}(\Omega)$, so by Inequality (3.10) $\tilde{\Phi}_{\lambda}$ is coercive as well. As recalled in Section 2 , $\tilde{\Phi}_{\lambda}$ then satisfies the ( $P S$ )-condition. Thus, we may apply the mountain pass theorem (see [33, Theorem 2.1]) and deduce the existence of $v_{\lambda} \in W_{0}^{S, p}(\Omega) \backslash\left\{0, u_{\lambda}\right\}$ s.t. $\tilde{\Phi}_{\lambda}^{\prime}\left(v_{\lambda}\right)=0$ in $W^{-s, p^{\prime}}(\Omega)$. So, we have weakly in $\Omega$

$$
\begin{equation*}
(-\Delta)_{p}^{S} v_{\lambda}=\tilde{f}_{\lambda}\left(x, v_{\lambda}\right) \tag{3.11}
\end{equation*}
$$

Testing Equation (3.11) with $-v_{\lambda}^{-} \in W_{0}^{s, p}(\Omega)$ and applying Equation (2.2) we have

$$
\left\|v_{\lambda}^{-}\right\|^{p} \leqslant\left\langle(-\Delta)_{p}^{s} v_{\lambda},-v_{\lambda}^{-}\right\rangle=\int_{\Omega} \tilde{f}_{\lambda}\left(x, v_{\lambda}\right)\left(-v_{\lambda}^{-}\right) d x=0
$$

so $v_{\lambda} \in W_{0}^{s, p}(\Omega)_{+} \backslash\{0\}$. Recalling the definition of $\tilde{f}_{\lambda}$ and testing Equation (3.11) with $\left(v_{\lambda}-u_{\lambda}\right)^{+} \in W_{0}^{s, p}(\Omega)$, we have

$$
\begin{aligned}
\left\langle(-\Delta)_{p}^{s} v_{\lambda},\left(v_{\lambda}-u_{\lambda}\right)^{+}\right\rangle & =\int_{\Omega} \tilde{f}_{\lambda}\left(x, v_{\lambda}\right)\left(v_{\lambda}-u_{\lambda}\right)^{+} d x \\
& \leqslant \int_{\Omega} f_{\lambda}\left(u_{\lambda}\right)\left(v_{\lambda}-u_{\lambda}\right)^{+} d x \\
& =\left\langle(-\Delta)_{p}^{s} u_{\lambda},\left(v_{\lambda}-u_{\lambda}\right)^{+}\right\rangle
\end{aligned}
$$

which by Proposition 2.1 implies $v_{\lambda} \leqslant u_{\lambda}$ in $\Omega$. So, Equation (3.11) rephrases as

$$
(-\Delta)_{p}^{s} v_{\lambda}=f_{\lambda}\left(v_{\lambda}\right)
$$

weakly in $\Omega$. Using Theorem 2.6 as in Theorem 3.1, we see that $v_{\lambda} \in \operatorname{int}\left(C_{S}^{0}(\bar{\Omega})_{+}\right)$and it solves $\left(P_{\lambda}\right)$. So we have

$$
\begin{cases}(-\Delta)_{p}^{s} v_{\lambda}+v_{\lambda}^{r-1}=\lambda v_{\lambda}^{q-1} \leqslant \lambda u_{\lambda}^{q-1}=(-\Delta)_{p}^{s} u_{\lambda}+u_{\lambda}^{r-1} & \text { weakly in } \Omega \\ v_{\lambda} \leqslant u_{\lambda} & \text { in } \Omega\end{cases}
$$

while $v_{\lambda} \not \equiv u_{\lambda}$. By Theorem 2.7, we have $u_{\lambda}-v_{\lambda} \in \operatorname{int}\left(C_{S}^{0}(\bar{\Omega})_{+}\right)$.

To complete the picture, we examine the limiting case $\lambda=\lambda_{*}$. In such case, we can prove existence of at least one solution, to which all principal solutions $u_{\lambda}$ converge:

Lemma 3.6. There exists a solution $u_{*} \in \operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)$of $\left(P_{\lambda_{*}}\right)$. Besides, if $u_{\lambda} \in \operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)$is the solution given in Lemma 3.4, then $u_{\lambda} \rightarrow u_{*}$ in both $W_{0}^{s, p}(\Omega)$ and $C_{S}^{0}(\bar{\Omega})$ as $\lambda \searrow \lambda_{*}$.

Proof. We prove a slightly more general assertion. Let $\left(\lambda_{n}\right)$ be a decreasing sequence s.t. $\lambda_{n} \searrow \lambda_{*}$, and denote by $u_{n} \in$ $\operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)$any solution of $\left(P_{\lambda_{n}}\right)$, then up to a subsequence $u_{n} \rightarrow u_{*}$ in both $W_{0}^{s, p}(\Omega)$ and $C_{s}^{0}(\bar{\Omega})$ as $n \rightarrow \infty$, being $u_{*} \in$ $\operatorname{int}\left(C_{S}^{0}(\bar{\Omega})_{+}\right)$a solution of $\left(P_{\lambda_{*}}\right)$. First, for all $n \in \mathbb{N}$ we have weakly in $\Omega$

$$
\begin{equation*}
(-\Delta)_{p}^{s} u_{n}=f_{\lambda_{n}}\left(u_{n}\right) \tag{3.12}
\end{equation*}
$$

Arguing as in the proof of Theorem 3.1, we find $u_{*} \in W_{0}^{s, p}(\Omega)_{+}$s.t. up to a subsequence $u_{n} \rightarrow u_{*}$ in both $W_{0}^{s, p}(\Omega)$ and $C_{s}^{0}(\bar{\Omega})$, hence we can pass to the limit in Equation (3.12) and get weakly in $\Omega$

$$
\begin{equation*}
(-\Delta)_{p}^{S} u_{*}=f_{\lambda_{*}}\left(u_{*}\right) \tag{3.13}
\end{equation*}
$$

We claim that $u_{*} \not \equiv 0$. Arguing by contradiction, assume that $u_{n} \rightarrow 0$ in both $W_{0}^{s, p}(\Omega)$ and $C_{s}^{0}(\bar{\Omega})$, hence in particular $u_{n} \rightarrow 0$ uniformly in $\Omega$. Then, for all $n \in \mathbb{N}$ big enough we have $0<u_{n} \leqslant 1$ in $\Omega$. Set for all $n \in \mathbb{N}$

$$
v_{n}=\frac{u_{n}}{\left\|u_{n}\right\|} \in W_{0}^{s, p}(\Omega) \cap \operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right) .
$$

The sequence $\left(v_{n}\right)$ is obviously bounded in $W_{0}^{s, p}(\Omega)$. By reflexivity and the compact embedding $W_{0}^{s, p}(\Omega) \hookrightarrow L^{p}(\Omega)$, passing to a subsequence we have $v_{n} \rightharpoonup v$ in $W_{0}^{\Omega, p}(\Omega), v_{n} \rightarrow v$ in $L^{p}(\Omega)$. Besides, by Equation (3.12), for all $n \in \mathbb{N}$ we have weakly in $\Omega$

$$
\begin{equation*}
(-\Delta)_{p}^{s} v_{n}=\lambda_{n} \frac{u_{n}^{q-1}}{\left\|u_{n}\right\|^{p-1}}-\frac{u_{n}^{r-1}}{\left\|u_{n}\right\|^{p-1}} . \tag{3.14}
\end{equation*}
$$

Consider the first term in the right-hand side of Equation (3.14). Since $0<u_{n} \leqslant 1$ in $\Omega$ and $p<q$, we have

$$
0<\frac{u_{n}^{q-1}}{\left\|u_{n}\right\|^{p-1}} \leqslant \frac{u_{n}^{p-1}}{\left\|u_{n}\right\|^{p-1}}=v_{n}^{p-1},
$$

so $\left(u_{n}^{q-1} /\left\|u_{n}\right\|^{p-1}\right)$ is bounded in $L^{p^{\prime}}(\Omega)$. Passing to a subsequence, we have $u_{n}^{q-1} /\left\|u_{n}\right\|^{p-1} \rightharpoonup w$ in $L^{p^{\prime}}(\Omega)$, hence a fortiori in $L^{1}(\Omega)$. By Hölder's inequality and the continuous embedding $W_{0}^{s, p}(\Omega) \hookrightarrow L^{q}(\Omega)$, we have

$$
\begin{aligned}
\|w\|_{1} & \leqslant \liminf _{n} \int_{\Omega} \frac{u_{n}^{q-1}}{\left\|u_{n}\right\|^{p-1}} d x \\
& \leqslant \limsup _{n} \frac{\left\|u_{n}\right\|_{q}^{q-1}|\Omega|^{\frac{1}{q}}}{\left\|u_{n}\right\|^{p-1}} \\
& \leqslant C \limsup _{n}\left\|u_{n}\right\|^{q-p}=0 .
\end{aligned}
$$

So, we get $w=0$, that is,

$$
\begin{equation*}
\frac{u_{n}^{q-1}}{\left\|u_{n}\right\|^{p-1}} \rightharpoonup 0 \text { in } L^{p^{\prime}}(\Omega) . \tag{3.15}
\end{equation*}
$$

An entirely similar argument proves that $\left(u_{n}^{r-1} /\left\|u_{n}\right\|^{p-1}\right)$ is bounded in $L^{p^{\prime}}(\Omega)$ and, up to a subsequence,

$$
\begin{equation*}
\frac{u_{n}^{r-1}}{\left\|u_{n}\right\|^{p-1}} \rightharpoonup 0 \text { in } L^{p^{\prime}}(\Omega) . \tag{3.16}
\end{equation*}
$$

Testing Equation (3.14) with $\left(v_{n}-v\right) \in W_{0}^{s, p}(\Omega)$ and using Hölder's inequality, we have for all $n \in \mathbb{N}$

$$
\begin{aligned}
\left\langle(-\Delta)_{p}^{s} v_{n}, v_{n}-v\right\rangle & =\int_{\Omega}\left[\lambda_{n} \frac{u_{n}^{q-1}}{\left\|u_{n}\right\|^{p-1}}-\frac{u_{n}^{r-1}}{\left\|u_{n}\right\|^{p-1}}\right]\left(v_{n}-v\right) d x \\
& \leqslant \lambda_{1}\left\|\frac{u_{n}^{q-1}}{\left\|u_{n}\right\|^{p-1}}\right\|_{p^{\prime}}\left\|v_{n}-v\right\|_{p}-\left\|\frac{u_{n}^{r-1}}{\left\|u_{n}\right\|^{p-1}}\right\|_{p^{\prime}}\left\|v_{n}-v\right\|_{p},
\end{aligned}
$$

and the latter tends to 0 as $n \rightarrow \infty$ by the relations above. By the $(S)_{+}-$property of $(-\Delta)_{p}^{s}$ we have $v_{n} \rightarrow v$ in $W_{0}^{s, p}(\Omega)$, hence $\|v\|=1$. On the other hand, testing Equation (3.14) with $v \in W_{0}^{s, p}(\Omega)$, we have for all $n \in \mathbb{N}$

$$
\left\langle(-\Delta)_{p}^{s} v_{n}, v\right\rangle=\int_{\Omega}\left[\lambda_{n} \frac{u_{n}^{q-1}}{\left\|u_{n}\right\|^{p-1}}-\frac{u_{n}^{r-1}}{\left\|u_{n}\right\|^{p-1}}\right] v d x .
$$

Passing to the limit as $n \rightarrow \infty$ and recalling Equations (3.15) and (3.16), we get $\|v\|^{p}=0$, a contradiction. Summarizing, $u_{*} \in W_{0}^{s, p}(\Omega)_{+} \backslash\{0\}$ and satisfies Equation (3.13). As in Lemma 3.3 we see that $u_{*} \in \operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)$solves $\left(P_{\lambda_{*}}\right)$.

Finally, taking into account the monotonicity property of Lemma 3.4, we conclude that globally $u_{\lambda} \rightarrow u_{*}$ in both $W_{0}^{s, p}(\Omega)$ and $C_{s}^{0}(\bar{\Omega})$, with monotone convergence, as $\lambda \searrow \lambda_{*}$, for some $u_{*} \in \operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)$solving $\left(P_{\lambda_{*}}\right)$.

Looking at the proof of Lemma 3.6 above, we can easily argue that, for any sequence $\left(\lambda_{n}\right)$ s.t. $\lambda_{n} \searrow \lambda_{*}$, the sequence of solutions ( $v_{\lambda_{n}}$ ) provided by Lemma 3.5 has a subsequence which converges to a solution of $\left(P_{\lambda_{*}}\right)$, which might differ from the global limit of $u_{\lambda}$.

Combining Lemmas 3.3-3.6, we obtain the following bifurcation result for the superdiffusive case (corresponding to case (c) of Theorem 1.1):

Theorem 3.7. Let $2 \leqslant p<q<r<p_{s}^{*}$. Then, there exists $\lambda_{*}>0$ with the following properties: for all $\lambda \in\left(0, \lambda_{*}\right)$ problem $\left(P_{\lambda}\right)$ has no solution; $\left(P_{\lambda_{*}}\right)$ has at least one solution $u_{*} \in \operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)$; and for all $\lambda>\lambda_{*}$ problem $\left(P_{\lambda}\right)$ has at least two solutions $u_{\lambda}, v_{\lambda} \in \operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)$s.t. $u_{\lambda}-v_{\lambda} \in \operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right), u_{\lambda}-u_{\mu} \in \operatorname{int}\left(C_{s}^{0}(\bar{\Omega})_{+}\right)$for all $\lambda>\mu>\lambda_{*}$, and $u_{\lambda} \rightarrow u_{*}$ in both $W_{0}^{s, p}(\Omega)$ and $C_{s}^{0}(\bar{\Omega})$ as $\lambda \searrow \lambda_{*}$.

Remark 3.8. For simplicity, we confined our study to the pure power logistic reactions. Nevertheless, most of our Theorem 3.7 can be extended to the following generalized logistic equation:

$$
\begin{cases}(-\Delta)_{p}^{s} u=\lambda f(x, u)-g(x, u) & \text { in } \Omega \\ u>0 & \text { in } \Omega \\ u=0 & \text { in } \Omega^{c},\end{cases}
$$

where $f, g: \Omega \times \mathbb{R} \rightarrow \mathbb{R}$ are Carathéodory mappings, both ( $p-1$ )-superlinear at $\infty$ and at 0 , satisfying a subcritical growth condition like $\mathbf{H}$, and jointly satisfying a pseudo-monotonicity condition (see [23] for the case of the $p$-Laplacian).

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## CONFLICT OF INTEREST

The authors declare no potential conflict of interests.

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