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PAPER

Kinetics of α hydrogen abstractions from thiols, sulfides and thiocarbonyl compounds†

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Hydrogen abstraction reactions involving organosulfur compounds play an important role in many industrial, biological and atmospheric processes. Despite their chemical relevance, little is known about their kinetics. In this work a group additivity model is developed that allows predicting the Arrhenius parameters for abstraction reactions of α hydrogen atoms from thiols, alkyl sulfides, alkyl disulfides and thiocarbonyl compounds by carbon-centered radicals at temperatures ranging from 300 to 1500 K. Rate coefficients for 102 hydrogen abstractions were obtained using conventional transition state theory within the high-pressure limit. Electronic barriers were calculated using the CBS-QB3 method and the rate coefficients were corrected for tunneling and hindered rotation about the transitional bond. Group additivity values for 46 groups are determined. To account for resonance and hyperconjugative stabilization in the transition state, 8 resonance corrections were fitted to a set of 32 reactions. The developed group additivity scheme was validated using a test set containing an additional 30 reactions. The group additivity scheme succeeds in reproducing the rate coefficients on average within a factor of 2.4 at 300 K and 1.4 at 1000 K. Mean absolute deviations of the Arrhenius parameters amount to, respectively, 2.5 kJ mol⁻¹ for E_a and 0.13 for log A , both at 300 and 1000 K. This work hence illustrates that the recently developed group additivity methods for Arrhenius parameters extrapolate successfully to hetero-element containing compounds.

1 Introduction

Despite the fact that sulfur compounds play an important role in many atmospheric, biological and chemical processes,^{1–3} large parts of their chemistry remain largely unknown. Recent studies have shown that reactions involving organosulfur compounds occur *via* complex radical chemistries making it hard to derive a reaction mechanism and extract reliable rate coefficients from experimental data only.^{4–10} Often, the decomposition of small organosulfur compounds is accompanied with the formation of large chain polysulfides,¹¹ which can result in deficient sulfur balances if those compounds remain undetected. Extended reaction networks accounting for the formation of heavier sulfur compounds can help in identifying unknown species. Recently, automatic reaction network generating tools

have become available that allow description of complex radical chemistries at the molecular level.¹² Radical reaction networks can easily contain up to thousands of reactions and for each reaction reliable rate coefficients need to be at hand. Since hydrogen abstraction reactions involving organosulfur compounds play an important role in the free-radical chemistry of many processes such as polymerization, combustion, pyrolysis and steam cracking of hydrocarbons,^{8,13} kinetic data for this reaction family are indispensable for reaction network generation.

Despite the increase in computational performance, *ab initio* methods still remain too demanding to calculate all the required rate coefficients, in particular for the larger species in the network. To keep the effort tractable, kinetic parameters are frequently obtained using kinetic correlations.^{14–20} The most popular correlation is without doubt the Bell–Evans–Polanyi relationship, which correlates the activation energy for reactions within the same reaction family to the reaction enthalpy.^{14,15} Many extensions to the Bell–Evans–Polanyi (BEP) relationship have been proposed in order to increase its accuracy. However, the introduction of additional parameters in the relationship limits its applicability in practice. Two such extended models are the intersecting parabolas (IP) method¹⁶ and the Blowers and Masel model.¹⁷ The major limitation of these BEP-type methods is that only activation energies can be estimated.

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† Electronic supplementary information (ESI) available: Rate coefficients, tunneling coefficients and symmetry numbers for all reactions presented in this work and B3LYP/6-311G(2d,d,p) geometries for all reactants, products and transition states. Guidelines for the determination of the symmetry numbers are provided. See DOI: 10.1039/c2cp41114h

Pre-exponential factors hence have to be obtained from other sources.

Alternative methods that can model both Arrhenius parameters rely on the additivity²¹ of enthalpy and entropy in the transition state or on the reaction class approach.^{22–33} Sumathi *et al.*^{27,28} proposed a method to obtain accurate kinetic data for hydrogen abstraction reactions using supergroups that encompass the reactive moiety of the transition state structure. The major advantage of these supergroups is that they can account for non-atom-centered contributions, *i.e.* contributions originating from interactions between non-adjacent groups, for example *cis/trans* or *gauche* interactions. Supergroups are larger than the commonly used Benson-like groups, resulting in the total amount of possible supergroups outnumbering the amount of their Benson analogues. Truong *et al.*²⁹ used reaction class transition state theory to predict rate coefficients for hydrogen abstraction reactions between methyl and alkanes³⁰ and between hydrogen and alcohols.^{32,33} In this approach the rate coefficient of a target reaction can be calculated by multiplying the rate coefficient of a reference reaction with a set of four correction factors accounting for effects of symmetry, tunneling, partition function and potential energy. In this work, the group additivity (GA) method as proposed by Saeys *et al.*²⁵ and further extended by Sabbe *et al.*²⁶ is used to model the kinetics of hydrogen abstractions involving organosulfur compounds. This method makes use of the additivity of enthalpy and entropy for reactants and transition state. This allows us to obtain the kinetic parameters as perturbations to a reference reaction. The group additivity values (GAVs) can be obtained from high-level quantum chemical calculations. The method has proven to be successful in predicting rate coefficients of addition reactions and hydrogen abstractions for hydrocarbons^{26,34} and H₂ additions, 1,2-hydrogen shifts and cyclization reactions for silicon-containing compounds.^{35–37}

The aim of this work is to extend the previously developed additivity schemes for Arrhenius parameters of hydrogen abstraction reactions between hydrocarbons to sulfur containing compounds.^{25,26} To this end, Arrhenius parameters are obtained from high-level quantum chemical calculations. Although hydrogen abstraction reactions have been extensively studied in literature,^{27,28,38–40} kinetic data for abstractions involving sulfur compounds are still scarce and only very limited attention has been given to the influence of neighboring S and C=S groups on the kinetics of a C–H–C type of hydrogen abstraction reactions.

Benchmark studies have assessed the influence of the level of theory, tunneling and hindered rotor corrections on the reaction kinetics of this reaction family. Coote⁴¹ studied geometries, barriers and reaction enthalpies for hydrogen abstraction reactions involving heteroatom containing hydrocarbons and showed that high-level composite procedures were in good agreement with experimental data. Barrier heights for an elaborate set of hydrogen abstractions were studied in the work of Zheng *et al.*^{42,43} In a recent study we assessed the influence of the level of theory, tunneling and 1-D hindered rotor (1-D HR) treatment for hydrogen abstractions involving hydrocarbons.⁴⁰ It was observed that the best agreement with experimental data for a large set of hydrogen abstractions was obtained using the CBS-QB3 method within the 1D-HR treatment of the internal

rotation about the transition state bond and corrected for tunneling effects according to the Eckart scheme.⁴⁴

In this work, the effect of an α S-atom or a thiocarbonyl group on the rate coefficients of hydrogen abstraction reactions is investigated. Due to their importance in combustion, pyrolysis and polymerization processes, the applicability of group additivity schemes for Arrhenius parameters is explored for this reaction family. The various steps that need consideration in order to extrapolate currently developed group additivity schemes to reactions involving hetero-elements are discussed. Arrhenius parameters for a set of 102 reactions are calculated using the CBS-QB3 composite method corrected for the 1-D hindered rotation about the transitional bond. Tunneling contributions were calculated according to the Eckart scheme.⁴⁴ 46 group additivity values are determined both for activation energies and pre-exponential factors. Eight corrections accounting for resonance in the transition state were determined by regression of group additively modeled Arrhenius parameters to kinetic data calculated for a set of 32 reactions. The accuracy of the group additivity method is assessed by comparison of CBS-QB3 data with group additively modeled data for a test set containing 30 additional reactions.

2 Computational methods

2.1 Rate coefficients

The rate coefficients presented in this work were calculated using classical transition state theory:

$$k_{\infty}(T) = \kappa_{\text{Eckart}}(T) \frac{k_{\text{B}}T}{h} \frac{q_{\ddagger}}{q_{\text{A}}q_{\text{B}}} e^{-\frac{\Delta E(0\text{ K})}{RT}} \quad (1)$$

In eqn (1) q represents the total partition function per unit volume, $\Delta E(0\text{ K})$ the electronic zero-point corrected reaction barrier and $\kappa_{\text{Eckart}}(T)$ the Eckart transmission coefficient accounting for quantum mechanical effects. The electronic barrier at 0 K is determined with the CBS-QB3 complete basis set method of Montgomery *et al.*⁴⁵ which is implemented in the *Gaussian 03* computational package.⁴⁶ Partition functions q are calculated at the B3LYP/6-311G(2d,d,p) level using a default scaling factor of 0.99. B3LYP geometries of all species considered in this work can be found in the ESI.† The rotation about the transitional bond is treated as a one-dimensional hindered internal rotor.^{47,48} In case the barrier for rotation is lower than 1 kJ mol⁻¹, the rotation is treated as a free rotor. Arrhenius parameters (E_{a} and $\log A$) were obtained by fitting to *ab initio* rate coefficients in the temperature range $T - 100\text{ K}$ to $T + 100\text{ K}$, with T the temperature of interest. Rate coefficients for all reactions considered in this study at temperatures ranging from 300–1500 K are presented in Tables S2–S5 of ESI.†

In this study, the accuracy of the GA method is assessed by comparing *ab initio* calculated rate coefficients and rate coefficients estimated by the group additivity method. As a measure of the deviation between both values we opted to define a factor ρ according to eqn (2).

$$\begin{cases} \rho = \frac{k_{\text{AI}}}{k_{\text{GA}}} & k_{\text{AI}} > k_{\text{GA}} \\ \rho = \frac{k_{\text{GA}}}{k_{\text{AI}}} & k_{\text{GA}} > k_{\text{AI}} \end{cases} \quad (2)$$

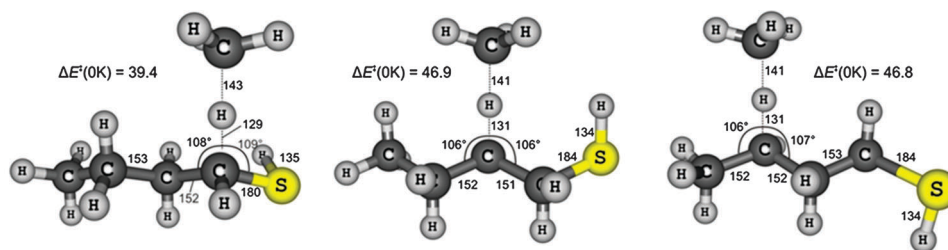
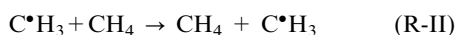


Fig. 2 B3LYP/6-311G(2d,d,p) transition state geometries for hydrogen abstractions from butane-1-thiol by methyl. Transition states for abstraction from the carbon atom in α , β and γ positions of the thiol group are shown. CBS-QB3 energy barriers $\Delta E^\ddagger(0\text{ K})$ are indicated in kJ mol^{-1} .

from hydrocarbons to organosulfur compounds only requires additional additivity values for hydrogen abstractions of carbon atoms having adjacent thiol and thiocarbonyl groups. In Section 3.3.2, the limited influence of secondary contributions and NNIs on the Arrhenius parameters will be further illustrated.

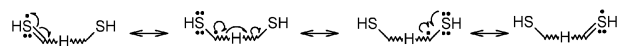
$\Delta\text{GAV}_{E_a}^\circ(C_1)$ is determined as the difference between the activation energy of the hydrogen abstraction by a given radical from methane and the activation energy of the reference reaction, *i.e.* the hydrogen abstraction by methyl from methane. $\Delta\text{GAV}^\circ(C_1)$ hence accounts for the influence of resonance and hyperconjugation resulting from groups centered on C_1 . Similarly, $\Delta\text{GAV}^\circ(C_2)$ includes the influence of resonance and hyperconjugation due to groups centered on C_2 . Therefore, the term $\Delta E_{a,\text{res}}^\circ$ in eqn (4) is a correction term that accounts for additional resonance and hyperconjugative stabilization in the transition state due to the simultaneous presence of ligands on the C_1 and C_2 groups. To capture these cross resonance and/or hyperconjugative stabilization effects in the transition state, Sabbe *et al.*²⁶ have introduced 4 corrections that allow us to account for every occurrence of a $\pi-\pi$, $\pi-\pi$, $\pi-\pi$, $\pi-\sigma_{\beta\text{C-H}}$ and $\sigma_{\beta\text{C-H}}-\sigma_{\beta\text{C-H}}$ interaction in the transition state for hydrogen transfer involving hydrocarbons. In this work, additional corrections are presented for cross-interactions involving α -sulfur atoms and thiocarbonyl groups.

The need for resonance correction terms is illustrated with an example, *i.e.* the hydrogen abstraction reaction from the C atom in methanethiol (CH_3SH) by the methylenethiol radical ($\text{C}^\bullet\text{H}_2\text{SH}$) (R-I).



The two ΔGAV° s required to estimate the Arrhenius parameters for this reaction are $\Delta\text{GAV}^\circ(C_1-\text{S})(\text{H})_2$ and $\Delta\text{GAV}^\circ(C_2-\text{S})(\text{H})_2$. These two ΔGAV° s account for the influence of a sulfur ligand on the C_1 and C_2 groups, respectively. The $\Delta\text{GAV}^\circ(C_1-\text{S})(\text{H})_2$ values are equal to the difference between the Arrhenius parameters obtained for the hydrogen abstraction from methane by $\text{C}^\bullet\text{H}_2\text{SH}$ (R-III) and those of the reference reaction (R-II). Similarly, the $\Delta\text{GAV}^\circ(C_2-\text{S})(\text{H})_2$ values are equal to the difference between the Arrhenius parameters of reaction (R-IV), which is the reverse of (R-III), and those of the reference reaction (R-II).

The spin density plots for the transition states of the reference reaction (R-II), reactions (R-III), (R-IV) and (R-I) shown in Fig. 3, illustrate the presence of resonance stabilization. In the transition state of the reference reaction (R-II), the spin is mainly located on the two carbon atoms that exchange the hydrogen atom. For the transition state of (R-III) and (R-IV) it is observed that the Mulliken atomic spin density on the sulfur atom amounts to 0.08. Since the spin delocalization is limited to the neighboring S-atom, this resonance effect on the Arrhenius parameters is captured within the ΔGAV° s. In the transition state of (R-I), spin delocalization indicates that both sulfur atoms are involved in delocalization of the unpaired electron in the transition state. Clearly, the spin density is distributed over the two carbon atoms exchanging the hydrogen atom and their neighboring S atoms. The electron delocalization in the transition state can be represented by the following resonance structures:



This electron delocalization effect cannot be captured within the ΔGAV° s, as it originates from cross-interactions caused by the simultaneous presence of a sulfur ligand on the C_1 and C_2 groups.

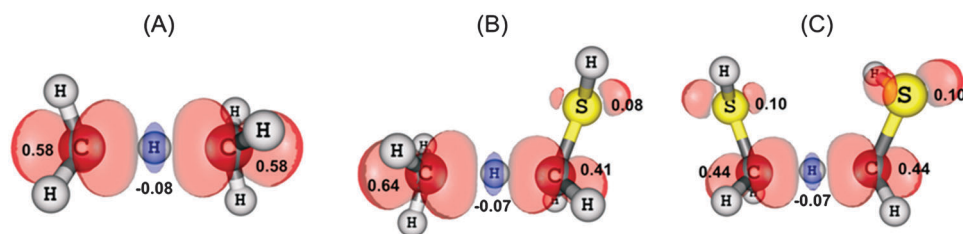


Fig. 3 Spin density plots and Mulliken atomic spin densities for the transition state of (A) hydrogen abstraction by methyl from methane (R-II), (B) hydrogen abstraction by methyl from methanethiol (R-III) and (R-IV), and (C) hydrogen abstraction by the methylenethiol radical from methanethiol (R-I) (B3LYP/6-311G(2d,d,p), contour value = 0.005).

Therefore, a correction factor $\Delta E_{a,\text{res}}^{\circ}$ needs to be added to eqn (4) in order to account for the additional stabilization due to the simultaneous presence of resonance in the abstracting and the formed radical.

As shown by Sabbe *et al.*,²⁶ a relationship similar to eqn (4) can be used to model the single-event pre-exponential factor $\log \tilde{A}$ of a target reaction:

$$\log \tilde{A}(T) = \log \tilde{A}_{\text{ref}}(T) + \sum_{i=1}^2 \Delta \text{GAV}_{\log \tilde{A}}^{\circ}(C_i) + \Delta \log \tilde{A}_{\text{res}}^{\circ} \quad (5)$$

$\Delta \log \tilde{A}_{\text{res}}^{\circ}$ in eqn (5) accounts for the influence of cross-interactions between resonance structures on C_1 and C_2 . The pre-exponential factor A is then obtained by adding the number of single events, n_e , to the single-event pre-exponential factor \tilde{A} obtained from the GA method:

$$\log A = \log \tilde{A} + \log n_e \quad (6)$$

n_e is analogous to the reaction path degeneracy and is given by the formula:⁵⁰

$$n_e = \frac{n_{\text{opt},\ddagger} \sigma_A \sigma_B}{n_{\text{opt},A} n_{\text{opt},B} \sigma_{\ddagger}} \quad (7)$$

In eqn (7) σ represents the total symmetry number of the molecule and n_{opt} a correction factor accounting for the mixing of optically active species. In analogy with the single-event pre-exponential factor \tilde{A} , single-event standard reaction entropies $\Delta_r S^{\circ}$ and single-event rate coefficients \tilde{k} are also used in this work.

2.3 Symmetry numbers

The total symmetry number of a molecule is equal to:

$$\sigma = \sigma_{\text{ext}} \prod_i \sigma_{\text{int},i} \quad (8)$$

with σ_{ext} and $\sigma_{\text{int},i}$ the external and internal symmetry numbers. External symmetry numbers were obtained from the *Gaussian 03* program. In case the reported experimental external symmetry number⁵¹ deviates from the one calculated by *Gaussian 03* (see Table S8 in ESI[†]), the experimental symmetry number was used to calculate the partition function. The internal symmetry number accounts for the symmetry of internal rotations. Methylene ($-\text{C}^{\bullet}\text{H}_2$) and *i*-propylene ($-\text{C}^{\bullet}(\text{CH}_3)_2$) rotating tops were assigned a symmetry number of 2, in line with previous work on hydrocarbons.²⁶

To account for nonsymmetrical torsion profiles, the various distinguishable structures can be calculated and their Boltzmann averaged contributions summed in order to obtain the total partition function.⁵¹ To avoid calculations for the various structures, the following approach was used: the higher energy conformers are neglected while for the energetically equivalent rotational minima the harmonic oscillator partition function of one of them is multiplied with a correction factor n_{conf} , which accounts for the mixing of the equivalent structures.⁵⁰ This is equivalent to dividing the total symmetry number in eqn (8) by a factor n_{conf} .⁵² For example, dihydrogen disulfide H_2S_2 has an external twofold axis while internal rotation around the S–S bond presents two distinct minima which are

each other's mirror image, having the S–H bonds nearly at right angles. To account for the two equivalent minima in H_2S_2 , Ercolani *et al.*⁵² proposed use of an apparent symmetry number of 1, instead of 2.

In analogy with previous work on hydrocarbons, n_{conf} is not included in the $\Delta \text{GAV}^{\circ}$ s.²⁶ Users of the GA method can opt to include these effects in the calculation of the pre-exponential factor by using an apparent symmetry number ($\sigma_{\text{app}} = \sigma/n_{\text{conf}}$) in eqn (7), rather than the symmetry number defined by eqn (8).

2.4 Tunneling

It has been shown for hydrogen abstractions that, especially at lower temperatures, tunneling corrections have a profound effect on the calculated rate coefficients and thereby influence the Arrhenius parameters.²⁹ Tunneling corrections are mainly temperature dependent and can differ significantly for the different reactions within a reaction family. Therefore, inclusion of the tunneling corrections in the $\Delta \text{GAV}^{\circ}$ s would lead to temperature dependent $\Delta \text{GAV}^{\circ}$ s. A group additivity model hence benefits from separating tunneling contributions from the calculated Arrhenius parameters. Alternatively, tunneling can be modeled explicitly using correlations with other known properties of the reaction as exemplified in the reaction class transition state theory of Truong *et al.*^{30,32,33} In the reaction class transition state theory, tunneling corrections for hydrogen abstractions are modeled by inclusion of a temperature dependent tunneling factor, which is taken as the ratio of the tunneling coefficients for a target reaction and the reference reaction.³⁰

In this work, the Eckart method is used to determine the tunneling contributions. This method requires information concerning the imaginary frequency and both the forward and reverse barrier heights. However, those values are not accessible using the group additivity method presented in this work. Since tunneling can only occur through the net electronic barrier and as within the same reaction family imaginary frequencies and barrier heights are strongly related, the activation energy of the exothermic reaction can be considered as one of the main factors determining the tunneling contribution. For hydrogen abstraction reactions of the C–H–C type, Sabbe *et al.*²⁶ have shown that tunneling contributions can be reproduced within a mean factor of deviation of 1.15 using a fourth order polynomial with temperature dependent coefficients:

$$\kappa(T) = 1 + \left(\frac{162}{T}\right)^3 E_{a,\text{exo}} + 2.71 \times 10^{-6} \exp\left(-\frac{T-300}{26}\right) E_{a,\text{exo}}^4 \quad (9)$$

with $E_{a,\text{exo}}$ the activation energy for the exothermic reaction in kJ mol^{-1} . The parameters in eqn (9) were obtained by fitting to tunneling coefficients for 55 reactions.²⁶ In Section 3.3.4 of this work the applicability of eqn (9) for the reactions studied in this work will be illustrated. Eckart tunneling coefficients for all reactions considered in this work at temperatures ranging from 300–1500 K are given in Table S7 of the ESI.[†]

3 Results and discussion

3.1 Rate coefficients and Arrhenius parameters

Tunneling coefficients, Arrhenius parameters and rate coefficients at 300 K for 24 hydrogen abstraction reactions by methyl are shown in Table 1. From the 24 reactions presented in Table 1 two sets of 23 ΔGAV 's (both for the forward and reverse reaction directions) will be obtained. In Table 2 kinetic parameters at 300 K for six sets of reactions are presented to study the effect of secondary contributions and non-nearest-neighbor interactions (NNIs) on the Arrhenius parameters. In Table 3, tunneling coefficients, Arrhenius parameters and rate coefficients at 300 K are presented for an additional set of 32 reactions from which the influence of resonance in the transition state on the Arrhenius parameters can be studied. These 32 reactions involve hydrogen abstractions between allylic, propargylic, thiocarbonylic radicals, their methyl substituted analogues, alkyl radicals and carbon-centered radicals having a sulfur-atom in the α -position of the radical site. Rate coefficients in the temperature range 300–1500 K and transition state geometries for all studied reactions can be found in the ESI.†

Validating calculated data with experiment is of primary importance. However, in particular for sulfur chemistry, experimental data are scarce and for the few experiments available, rate coefficients are often obtained by fitting to simplified reaction mechanisms or were measured at such low pressures that pressure dependence can no longer be ignored. However, the few data that could be retrieved agree well with the rate coefficients reported in this work. Arthur and Lee⁵³ report a rate coefficient at 500 K for the hydrogen abstraction from dimethyl sulfide by methyl of $4.1 \times 10^3 \text{ m}^3 \text{ mol}^{-1} \text{ s}^{-1}$. This agrees within a factor of 2 with the rate coefficient calculated in this work, *i.e.* $5.1 \times 10^3 \text{ m}^3 \text{ mol}^{-1} \text{ s}^{-1}$ (see ESI†). Ekwenchi *et al.*⁵⁴ studied hydrogen abstraction from an α -carbon atom in diethyl sulfide by ethyl and obtained a rate coefficient $k = 7.4 \times 10^6 \exp(-3452/T) \text{ m}^3 \text{ mol}^{-1} \text{ s}^{-1}$ or $1.3 \times 10^4 \text{ m}^3 \text{ mol}^{-1} \text{ s}^{-1}$ at 400 K. This rate coefficient is almost 3 orders of magnitude higher than for any reaction studied in this work. As the authors carefully reported their measurements,⁵⁴ we were able to reassess the reported rate coefficient and obtain $k = 6.2 \times 10^4 \exp(-3765/T) \text{ m}^3 \text{ mol}^{-1} \text{ s}^{-1}$ in the temperature range 383–443 K. Details of this reassessment can be found in the ESI.† This rate coefficient agrees within a factor of 5 with the calculated rate coefficients obtained for similar reactions, such as hydrogen abstraction by ethyl from a secondary α -C atom in ethanethiol or ethyl vinyl sulfide (reactions (40) and (41)). Estimated rate coefficients for hydrogen abstraction reactions by methyl from diethyl sulfide are reported by Zheng *et al.*⁹ At 1000 K our rate coefficients agree within one order of magnitude with their data.

The reactions studied in Tables 1 and 2 are hydrogen abstraction reactions by the methyl and ethyl radical from a carbon atom adjacent to a sulfur atom or a thiocarbonyl group. All reactions in Tables 1 and 2 are exothermic with the exception of the two abstraction reactions from thiophene (reactions (23) and (24)). The entropy change during reaction fluctuates around $0 \text{ J mol}^{-1} \text{ K}^{-1}$, as both the breaking and forming bonds during reaction are the same, *i.e.* a C–H bond.

At 300 K, the rate coefficients range between 3.7×10^{-24} up to $9.0 \times 10^2 \text{ m}^3 \text{ mol}^{-1} \text{ s}^{-1}$. The lowest rate coefficient is obtained for hydrogen abstraction from methane by the $\text{CH}_3\text{C}^\bullet(\text{SH})\text{C}(=\text{S})\text{H}$ radical (reaction (13)). The activation energy for this reaction amounts to 172 kJ mol^{-1} , which is the highest value obtained among all reactions studied in this work. The high barrier is caused by the stability of the $\text{CH}_3\text{C}^\bullet(\text{SH})\text{C}(=\text{S})\text{H}$ radical, resulting from (a) hyperconjugation, (b) π -conjugation with the adjacent thiocarbonyl group and (c) the α -effect due to the presence of a neighboring S atom. The lowest barriers are obtained for abstraction of hydrogen in the α position of a C=S group by methyl (reactions (5)–(7)) and for hydrogen abstraction from methane by the thiofen-2-yl radical (reaction (23)). Adjacent C=S groups significantly ease hydrogen abstraction, as in the resulting radical most spin will be evacuated to the sulfur atom.

At 300 K tunneling coefficients (Tables 1–3) are significant and range between 2 and 52. For hydrogen abstraction reactions by the methyl radical (Table 1), pre-exponential factors fluctuate between 10^5 and $10^7 \text{ m}^3 \text{ mol}^{-1} \text{ s}^{-1}$. This is somewhat higher than the values obtained for the reactions presented in Table 3 having a resonance stabilized transition state and for which the pre-exponential factor ranges around $10^{4.8}$. The higher pre-exponential factors obtained for hydrogen abstraction reactions involving methyl are due to the higher reaction path degeneracy.

The presented data show that the reaction enthalpy and the activation energy are strongly correlated. This can also be seen from Fig. 4, in which the activation energies for the reactions presented in Table 1 are shown as a function of the reaction enthalpy. A Blowers and Masel curve (eqn (10)) is fitted to the data:¹⁷

$$E_a = \left(\frac{w_b + w_f + \Delta_r H}{2} \right) \frac{(V_p - (w_b + w_f) + \Delta_r H)^2}{V_p^2 - (w_b + w_f)^2 + \Delta_r H^2} \quad (10)$$

Since typical C–H bond energies range around 400 kJ mol^{-1} , the parameter $w_b + w_f$ was set to 800 kJ mol^{-1} . A least-squares fit resulted in a V_p of 1114 kJ mol^{-1} , corresponding to an intrinsic barrier E_a^0 of 65.7 kJ mol^{-1} . The Blowers and Masel model succeeds in reproducing the calculated activation energies on average within 6 kJ mol^{-1} . However, maximum deviations can amount up to 21 kJ mol^{-1} .

For hydrogen abstractions by the methyl radical, barriers for rotation about the transition state bond are limited to 1 kJ mol^{-1} . Treating these rotations as free rotors significantly decreases the rate coefficients with at most a factor of 3 at 300 K up to a factor of 7 at 1500 K. For the reactions presented in Table 3 rotational barriers vary between 1 up to 13 kJ mol^{-1} . The higher barriers for rotation are due to steric effects and are in particular observed for reactions involving thiocarbonyl compounds.

3.2 Resonance interactions

The group additivity model for Arrhenius parameters of hydrogen abstraction reactions can be significantly improved by introducing correction factors that account for cross-resonance and/or hyperconjugative effects in the transition state.²⁶ To model resonance effects in the transition states of

Table 1 Tunneling coefficients κ (–), standard reaction enthalpies $\Delta_r H^\circ$ (kJ mol⁻¹), standard reaction entropies $\Delta_r S^\circ$ (J mol⁻¹ K⁻¹), pre-exponential factors $\log A$ (log(m³ mol⁻¹ s⁻¹)), activation energies E_a (kJ mol⁻¹) and rate coefficients k (m³ mol⁻¹ s⁻¹) at 300 K for the training set

| | | κ | Forward | | | | Reverse | | | |
|------|--|----------|--------------------|--------------------|----------|-------|----------------------|----------|-------|-----------------------|
| | | | $\Delta_r H^\circ$ | $\Delta_r S^\circ$ | $\log A$ | E_a | k | $\log A$ | E_a | k |
| (1) | <chem>CH3 + CH4 <=> CH4 + CH3</chem> | 49.1 | 0.0 | 0.0 | 6.171 | 69.7 | 5.2×10^{-5} | 6.171 | 69.7 | 5.2×10^{-5} |
| (2) | <chem>CH3 + -SH <=> CH4 + -S</chem> | 28.5 | -41.0 | -0.3 | 5.783 | 48.8 | 5.3×10^{-2} | 5.800 | 89.8 | 4.0×10^{-9} |
| (3) | <chem>CH3 + CH2SH <=> CH4 + CH2S</chem> | 17.6 | -46.6 | 1.4 | 5.458 | 41.2 | 3.2×10^{-1} | 5.384 | 87.8 | 2.1×10^{-9} |
| (4) | <chem>CH3 + CH(CH3)SH <=> CH4 + CH(CH3)S</chem> | 11.7 | -51.5 | 0.5 | 5.414 | 33.7 | 3.9 | 5.386 | 85.1 | 4.0×10^{-9} |
| (5) | <chem>CH3 + S=C=C <=> CH4 + S=C=C</chem> | 4.0 | -96.7 | -11.9 | 5.678 | 32.6 | 3.8 | 6.279 | 129.1 | 2.3×10^{-16} |
| (6) | <chem>CH3 + S=C=C <=> CH4 + S=C=C</chem> | 2.6 | -114.2 | -5.4 | 5.747 | 26.1 | 4.0×10 | 6.020 | 140.1 | 1.0×10^{-18} |
| (7) | <chem>CH3 + S=C=C <=> CH4 + S=C=C</chem> | 1.8 | -129.6 | -4.1 | 5.501 | 19.5 | 2.2×10^2 | 5.700 | 149.0 | 9.7×10^{-21} |
| (8) | <chem>CH3 + CH=CHSH <=> CH4 + CH=CHS</chem> | 26.6 | -3.2 | 1.5 | 5.980 | 52.6 | 1.7×10^{-2} | 5.892 | 55.6 | 4.0×10^{-3} |
| (9) | <chem>CH3 + CH=CHS <=> CH4 + CH=CHS</chem> | 30.3 | -48.0 | -2.4 | 6.065 | 63.6 | 2.9×10^{-4} | 6.185 | 111.5 | 1.7×10^{-12} |
| (10) | <chem>CH3 + CH=CHSH <=> CH4 + CH=CHS</chem> | 8.3 | -97.3 | -13.2 | 5.555 | 34.9 | 2.4 | 6.236 | 132.1 | 1.4×10^{-16} |
| (11) | <chem>CH3 + CH=C(CH3)SH <=> CH4 + CH=C(CH3)S</chem> | 5.5 | -102.3 | -0.3 | 5.322 | 28.4 | 1.3×10 | 5.331 | 130.6 | 2.0×10^{-17} |
| (12) | <chem>CH3 + S=C=C <=> CH4 + S=C=C</chem> | 2.7 | -140.5 | -23.6 | 5.602 | 27.1 | 1.9×10 | 6.803 | 167.4 | 1.1×10^{-22} |
| (13) | <chem>CH3 + S=C=C <=> CH4 + S=C=C</chem> | 2.0 | -150.2 | -15.5 | 5.491 | 22.1 | 8.2×10 | 6.274 | 172.1 | 3.7×10^{-24} |
| (14) | <chem>CH3 + CH#CHSH <=> CH4 + CH#CHS</chem> | 8.6 | -92.8 | -10.4 | 5.416 | 31.3 | 7.6 | 5.955 | 124.0 | 1.9×10^{-15} |
| (15) | <chem>CH3 + CH#CHSH <=> CH4 + CH#CHS</chem> | 5.5 | -98.3 | 1.3 | 5.424 | 25.1 | 6.0×10 | 5.341 | 124.5 | 2.4×10^{-16} |
| (16) | <chem>CH3 + HS-CH2-C6H5 <=> CH4 + HS-C6H5</chem> | 9.0 | -85.4 | -10.5 | 5.394 | 37.4 | 6.6×10^{-1} | 5.937 | 122.8 | 3.1×10^{-15} |
| (17) | <chem>CH3 + HS-CH(CH3)-C6H5 <=> CH4 + HS-C6H5</chem> | 6.7 | -86.2 | -7.3 | 5.098 | 31.7 | 2.4 | 5.473 | 117.8 | 5.7×10^{-15} |
| (18) | <chem>CH3 + =S <=> CH4 + HC=S</chem> | 18.7 | -54.6 | -3.6 | 6.145 | 39.7 | 3.1 | 6.321 | 94.2 | 1.5×10^{-9} |
| (19) | <chem>CH3 + CH=S <=> CH4 + CH=S</chem> | 12.9 | -56.0 | -2.0 | 5.716 | 35.8 | 3.8 | 5.809 | 91.7 | 8.4×10^{-10} |
| (20) | <chem>CH3 + HS-C=S <=> CH4 + HS-C=S</chem> | 7.1 | -70.5 | -5.2 | 5.644 | 30.4 | 1.5×10 | 5.905 | 100.8 | 1.5×10^{-11} |
| (21) | <chem>CH3 + CH=CHS <=> CH4 + CH=CHS</chem> | 17.7 | -68.1 | -1.2 | 5.678 | 37.6 | 2.3 | 5.738 | 105.7 | 3.7×10^{-12} |
| (22) | <chem>CH3 + CH#CHS <=> CH4 + CH#CHS</chem> | 7.7 | -86.4 | -8.2 | 5.662 | 32.2 | 8.4 | 6.107 | 118.7 | 2.0×10^{-14} |
| (23) | <chem>CH3 + S5 <=> CH4 + S5</chem> | 4.6 | 54.3 | 3.3 | 6.090 | 75.0 | 4.7×10^{-7} | 5.908 | 20.6 | 9.0×10^2 |
| (24) | <chem>CH3 + S5 <=> CH4 + S5</chem> | 9.1 | 44.1 | 2.5 | 6.005 | 72.6 | 2.0×10^{-6} | 5.867 | 28.4 | 7.1×10 |

hydrogen abstraction reactions involving hydrocarbons, four resonance corrections were introduced.²⁶ These four corrections account for the stabilization in the transition state due to the simultaneous presence of hyperconjugating and π -conjugating groups on the C₁ and C₂ atoms. To describe hydrogen abstractions from thiols, sulfides and thiocarbonyl compounds, additional corrections need to be introduced accounting for cross-interactions involving S- and CS-ligands.

To determine the resonance corrections needed to model hydrogen abstraction reactions from organosulfur compounds,

Arrhenius parameters and rate coefficients were gathered for the 32 reactions presented in Table 3. These reactions allow us to study the resonance stabilization caused by (a) cross-interactions of a sulfur ligand on the C₁ atom with hyperconjugating C–H bonds, π -sulfur atoms or π -conjugating systems on the C₂ atom (reactions (42)–(59)) and (b) cross-interactions of a thiocarbonyl group on the C₁ atom with hyperconjugating C–H bonds, α -sulfur atoms or π -conjugating systems on the C₂ atom (reactions (60)–(73)). Besides the 4 corrections previously introduced for hydrogen abstractions between hydrocarbons,

Table 3 Tunneling coefficients κ (–), standard reaction enthalpies $\Delta_r H^\circ$ (kJ mol⁻¹), standard reaction entropies $\Delta_r S^\circ$ (J mol⁻¹ K⁻¹), pre-exponential factors $\log A$ (log(m³ mol⁻¹ s⁻¹)), activation energies E_a (kJ mol⁻¹) and rate coefficients k (m³ mol⁻¹ s⁻¹) at 300 K for the 32 hydrogen abstraction reactions from which the effect of resonance in the transition state on the Arrhenius parameters is studied

| | | κ | Forward | | | | Reverse | | | |
|------|--|----------|--------------------|--------------------|----------|-------|-----------------------|----------|-------|-----------------------|
| | | | $\Delta_r H^\circ$ | $\Delta_r S^\circ$ | $\log A$ | E_a | k | $\log A$ | E_a | k |
| (42) | <chem>CC(C)S + C=C.[C] <=> CC(C)S + C=C</chem> | 38.1 | 35.1 | 15.4 | 5.629 | 88.4 | 6.1×10^{-9} | 4.844 | 53.5 | 1.2×10^{-3} |
| (43) | <chem>CC(C)S + C=C.[C] <=> CC(C)S + C=C</chem> | 21.9 | 48.1 | 5.6 | 5.047 | 92.6 | 1.7×10^{-10} | 4.462 | 44.5 | 1.1×10^{-2} |
| (44) | <chem>CC(C)S + C=C.[C] <=> CC(C)S + C=C</chem> | 35.6 | 29.5 | 17.1 | 5.162 | 78.2 | 1.2×10^{-7} | 4.285 | 48.8 | 2.1×10^{-3} |
| (45) | <chem>CC(C)S + C=C.[C] <=> CC(C)S + C=C</chem> | 20.7 | 42.5 | 7.3 | 5.048 | 81.8 | 1.2×10^{-8} | 4.371 | 39.4 | 6.3×10^{-2} |
| (46) | <chem>CC(C)S + C1=CC=CC=C1.[C] <=> CC(C)S + C1=CC=CC=C1</chem> | 44.4 | 19.8 | 13.9 | 5.642 | 78.7 | 3.6×10^{-7} | 4.927 | 59.0 | 1.9×10^{-4} |
| (47) | <chem>CC(C)S + C1=CC=CC=C1.[C] <=> CC(C)S + C1=CC=CC=C1</chem> | 36.2 | 14.2 | 15.6 | 5.124 | 67.6 | 7.8×10^{-6} | 4.317 | 53.4 | 3.5×10^{-4} |
| (48) | <chem>CC(C)S + C1=CC=CC=C1.[C] <=> CC(C)S + C1=CC=CC=C1</chem> | 30.0 | 9.3 | 14.6 | 4.897 | 57.7 | 2.0×10^{-5} | 4.137 | 48.4 | 1.4×10^{-3} |
| (49) | <chem>CC(C)S + C#C.[C] <=> CC(C)S + C#C</chem> | 51.5 | 17.4 | 1.6 | 5.612 | 75.5 | 1.5×10^{-6} | 5.532 | 58.1 | 1.3×10^{-3} |
| (50) | <chem>CC(C)S + C#C.[C] <=> CC(C)S + C#C</chem> | 32.3 | 30.0 | -5.3 | 4.728 | 77.8 | 4.7×10^{-8} | 5.005 | 47.8 | 1.5×10^{-2} |
| (51) | <chem>CC(C)S + C#C.[C] <=> CC(C)S + C#C</chem> | 39.4 | 11.8 | 3.3 | 5.266 | 64.9 | 3.4×10^{-5} | 5.096 | 53.1 | 2.6×10^{-3} |
| (52) | <chem>CC(C)S + C#C.[C] <=> CC(C)S + C#C</chem> | 26.7 | 24.4 | -3.6 | 4.294 | 66.9 | 1.1×10^{-6} | 4.481 | 42.4 | 3.1×10^{-2} |
| (53) | <chem>CC(C)S + C#C.[C] <=> CC(C)S + C#C</chem> | 44.7 | -38.9 | -11.3 | 5.086 | 55.0 | 1.3×10^{-3} | 5.672 | 93.8 | 9.0×10^{-10} |
| (54) | <chem>CC(C)S + CC(C)S.[C] <=> CC(C)S + CC(C)S</chem> | 49.0 | 0.0 | 0.0 | 4.765 | 61.5 | 5.4×10^{-5} | 4.765 | 61.5 | 5.4×10^{-5} |
| (55) | <chem>CC(C)S + CC(C)S.[C] <=> CC(C)S + CC(C)S</chem> | 38.3 | 5.6 | -1.7 | 4.356 | 57.7 | 7.4×10^{-5} | 4.446 | 52.1 | 8.6×10^{-4} |
| (56) | <chem>CC(C)S + CC(C)S.[C] <=> CC(C)S + CC(C)S</chem> | 26.3 | 10.5 | -0.8 | 4.251 | 54.1 | 1.7×10^{-4} | 4.296 | 43.7 | 1.2×10^{-2} |
| (57) | <chem>CC(C)S + CC(C)S.[C] <=> CC(C)S + CC(C)S</chem> | 29.5 | 0.0 | 0.0 | 4.018 | 47.4 | 1.6×10^{-3} | 4.018 | 47.4 | 1.6×10^{-3} |
| (58) | <chem>CC(C)S + CC(C)S.[C] <=> CC(C)S + CC(C)S</chem> | 39.0 | -13.9 | -16.8 | 4.574 | 50.3 | 2.4×10^{-3} | 5.451 | 64.2 | 6.9×10^{-5} |
| (59) | <chem>CC(C)S + CC(C)S.[C] <=> CC(C)S + CC(C)S</chem> | 30.5 | -6.6 | -13.9 | 4.473 | 47.3 | 4.9×10^{-3} | 5.200 | 54.0 | 1.8×10^{-3} |
| (60) | <chem>CC(C)S + C=C.[C] <=> CC(C)S + C=C</chem> | 33.1 | -20.6 | 3.8 | 5.364 | 62.6 | 9.1×10^{-5} | 5.163 | 83.1 | 1.5×10^{-8} |
| (61) | <chem>CC(C)S + C=C.[C] <=> CC(C)S + C=C</chem> | 20.8 | -7.6 | -6.0 | 4.932 | 62.1 | 2.6×10^{-5} | 4.932 | 69.5 | 1.3×10^{-6} |
| (62) | <chem>CC(C)S + C=C.[C] <=> CC(C)S + C=C</chem> | 16.7 | -1.0 | -17.8 | 4.106 | 59.7 | 8.2×10^{-6} | 5.023 | 60.6 | 4.7×10^{-5} |
| (63) | <chem>CC(C)S + C#C.[C] <=> CC(C)S + C#C</chem> | 31.4 | -38.3 | -10.0 | 5.453 | 55.8 | 1.6×10^{-3} | 5.959 | 93.9 | 1.2×10^{-9} |
| (64) | <chem>CC(C)S + C#C.[C] <=> CC(C)S + C#C</chem> | 24.6 | -25.7 | -16.9 | 4.406 | 55.1 | 1.5×10^{-4} | 5.269 | 80.5 | 4.0×10^{-8} |
| (65) | <chem>CC(C)S + C#C.[C] <=> CC(C)S + C#C</chem> | 18.0 | -16.3 | -19.5 | 4.315 | 51.1 | 4.5×10^{-4} | 5.308 | 67.1 | 6.9×10^{-6} |
| (66) | <chem>CC(C)S + C#C.[C] <=> CC(C)S + C#C</chem> | 29.4 | 0.0 | 0.0 | 4.847 | 69.5 | 1.5×10^{-6} | 4.847 | 69.5 | 1.5×10^{-6} |
| (67) | <chem>CC(C)S + C#C.[C] <=> CC(C)S + C#C</chem> | 20.4 | -17.5 | 6.5 | 5.116 | 58.0 | 2.0×10^{-4} | 4.787 | 75.6 | 8.1×10^{-8} |
| (68) | <chem>CC(C)S + C#C.[C] <=> CC(C)S + C#C</chem> | 19.6 | 0.0 | 0.0 | 4.740 | 62.2 | 1.5×10^{-5} | 4.740 | 62.2 | 1.5×10^{-5} |
| (69) | <chem>CC(C)S + C#C.[C] <=> CC(C)S + C#C</chem> | 15.8 | -33.0 | 7.9 | 4.783 | 46.1 | 8.2×10^{-3} | 4.381 | 79.2 | 5.8×10^{-9} |
| (70) | <chem>CC(C)S + CC(C)S.[C] <=> CC(C)S + CC(C)S</chem> | 10.1 | -55.7 | -11.7 | 4.557 | 41.2 | 2.3×10^{-2} | 5.142 | 96.7 | 1.9×10^{-11} |
| (71) | <chem>CC(C)S + CC(C)S.[C] <=> CC(C)S + CC(C)S</chem> | 6.3 | -73.1 | -5.2 | 4.824 | 30.8 | 1.8 | 5.081 | 103.8 | 6.1×10^{-13} |
| (72) | <chem>CC(C)S + C.[C] <=> CC(C)S + C</chem> | 4.3 | -81.2 | -32.3 | 4.512 | 29.7 | 8.8×10^{-1} | 6.172 | 110.7 | 3.1×10^{-13} |
| (73) | <chem>CC(C)S + C.[C] <=> CC(C)S + C</chem> | 3.9 | -62.3 | -25.6 | 4.037 | 18.6 | 2.3×10 | 5.349 | 80.7 | 7.2×10^{-9} |

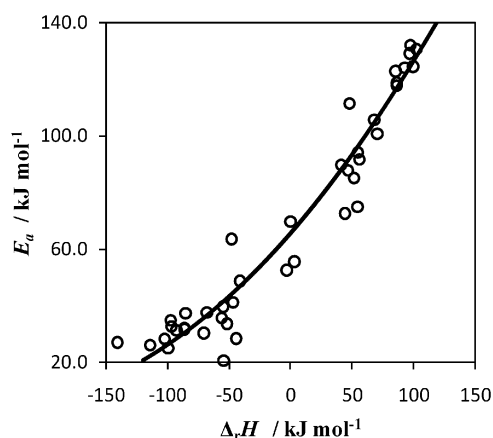


Fig. 4 Bell-Evans-Polanyi plot for the hydrogen abstraction reactions presented in Table 2. The full line corresponds to the Blowers and Masel model using $w_b + w_r = 800 \text{ kJ mol}^{-1}$ and $V_p = 1114 \text{ kJ mol}^{-1}$.

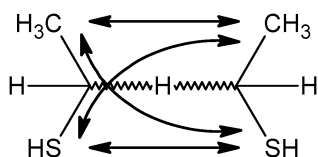


Fig. 5 Resonance interactions in the transition state of the symmetrical hydrogen abstraction reaction between ethylenethiol ($\text{CH}_3\text{C}^\bullet\text{HSH}$) and ethanethiol ($\text{CH}_3\text{CH}_2\text{SH}$).

one ligand on the C_2 atom, or *vice versa*. For example, $p_{\alpha S}-p_{\alpha S}$ accounts for the interaction of one sulfur atom bonded to the C_1 atom with one sulfur atom bonded to the C_2 atom. The number of corrections needed to describe the additional stabilization in the transition state due to cross-effects is obtained by counting all the cross interactions between ligands of the C_1 and C_2 groups. For example, the transition state for the hydrogen abstraction from methanethiol by the *i*-propyl radical (reaction (58)) is stabilized by the interaction of two hyperconjugating C–H bonds on the C_1 group with one α -sulfur atom on the C_2 group (see Fig. S7 of the ESI†). The additional transition state stabilization in reaction (58) is therefore equal to twice $p_{\alpha S}-\sigma_{\beta C-H}$. However, for more complex reactions resonance corrections can easily be overlooked and therefore as general aid eqn (11) and (12) can be applied:

$$n(X-X) = n_{X,C_1}n_{X,C_2} \quad (11)$$

$$n(X-Y) = n_{X,C_1}n_{Y,C_2} + n_{Y,C_1}n_{X,C_2} \quad (12)$$

with $X, Y = p_{\alpha S}, \sigma_{\beta C-H}, \pi_{C=C}, \pi_{\equiv}$ or $\pi_{C=S}$. $n(X-X)$ in eqn (11) is the number of resonance corrections needed to describe the cross-interactions caused by the simultaneous presence of the same X ligand on both the C_1 and C_2 atoms. $n(X-Y)$ in eqn (12) represents the number of cross-interactions between two different ligands, X and Y , on the C_1 and C_2 atoms. n_{X,C_1} and n_{X,C_2} represent the number of X ligands on C_1 and C_2 , respectively, while n_{Y,C_1} and n_{Y,C_2} are the number of Y ligands on C_1 and C_2 , respectively. The use of eqn (11) and (12) is illustrated for the symmetrical hydrogen abstraction reaction by ethylenethiol ($\text{CH}_3\text{C}^\bullet\text{HSH}$) from ethanethiol ($\text{CH}_3\text{CH}_2\text{SH}$). Fig. 5 illustrates that the transition state of this

reaction is stabilized by three types of cross-interactions: (a) between α -S atoms on both C_1 and C_2 ($p_{\alpha S}-p_{\alpha S}$), (b) between α -S atoms and a β C–H bond ($p_{\alpha S}-\sigma_{\beta C-H}$) and (c) between β C–H bonds on both C_1 and C_2 ($\sigma_{\beta C-H}-\sigma_{\beta C-H}$). The number of $p_{\alpha S}-p_{\alpha S}$ and $\sigma_{\beta C-H}-\sigma_{\beta C-H}$ corrections can be obtained from eqn (11), while eqn (12) accounts for the number of $p_{\alpha S}-\sigma_{\beta C-H}$ corrections required. The number of hyperconjugating C–H bonds on the C_1 and C_2 atoms amounts to 1 ($n_{\sigma_{\beta C-H},C_1} = n_{\sigma_{\beta C-H},C_2} = 1$). Also the number of S ligands on both primary groups is equal to 1 ($n_{p_{\alpha S},C_1} = n_{p_{\alpha S},C_2} = 1$). Therefore the total amount of corrections needed to describe the resonance in this transition state amounts to 2 $p_{\alpha S}-\sigma_{\beta C-H}$ corrections, 1 $p_{\alpha S}-p_{\alpha S}$ correction and 1 $\sigma_{\beta C-H}-\sigma_{\beta C-H}$ correction.

3.3 Group additivity model

3.3.1 Group additivity values. The data presented in Table 1 can be used to calculate ΔGAV° s for hydrogen abstraction reactions by/from carbon atoms having a sulfur or a thiocarbonyl ligand. The hydrogen abstraction from methane by methyl was chosen as the reference reaction. This ensures that the ΔGAV° s presented in this work can be combined with the ones reported in previous work for hydrogen abstraction reactions between hydrocarbons.²⁶ Arrhenius parameters at different temperatures for the reference reaction can be found in Table S1 of the ESI.† ΔGAV° s for both the pre-exponential factor and activation energy at 300 and 1000 K can be found in Table 4.

The derivation of the ΔGAV° s is straightforward. For example, $\Delta GAV_{E_a}^{\circ}(\text{CS}_2-(\text{H}))$ at 300 K can be obtained by subtracting the E_a of the reference reaction from the value obtained for reaction (18), *i.e.* $39.7-69.7 = -30.0 \text{ kJ mol}^{-1}$. Similarly, $\Delta GAV_{\log \tilde{A}}^{\circ}(\text{C}_2-(\text{S})(\text{H})_2)$ is obtained by subtracting the $\log \tilde{A}$ of the reference reaction from the $\log \tilde{A}$ obtained for reaction (18). The number of single events for the reference reaction equals 8, resulting in $\log \tilde{A}_{\text{ref}} = 6.171 - \log(8) = 5.267$, while n_e for reaction (18) is equal to 4, leading to $\log \tilde{A} = 6.145 - \log(4) = 5.543$. The $\Delta GAV_{\log \tilde{A}}^{\circ}(\text{CS}_2-(\text{H}))$ is the difference between both values and equals $5.543 - 5.267 = +0.276$. The reaction path degeneracy of 4 for reaction (18) is obtained by multiplying the number of equivalent hydrogen atoms in thioformaldehyde (=2) with a factor of two accounting for the two equivalent sites of the methyl radical which can abstract the hydrogen atom.

From Table 4 it can be seen that the ΔGAV° s are almost independent of temperature. The $\Delta GAV_{E_a}^{\circ}(C_1)$ s increase on average by 3 kJ mol^{-1} in going from 300 to 1000 K. The $\Delta GAV_{E_a}^{\circ}(C_2)$ s change at most by 1 kJ mol^{-1} and prove to be temperature independent. This can also be seen from Fig. 6 showing the temperature dependence of the C_1 and C_2 groups in the temperature range 300–1500 K. In general, the temperature dependence is small and restricted to 4 kJ mol^{-1} for $\Delta GAV_{E_a}^{\circ}$ and to 0.4 for $\Delta GAV_{\log \tilde{A}}^{\circ}$. Fig. 6 shows that the ΔGAV° s for the C_1 groups are slightly more dependent on the temperature than those obtained for the C_2 groups. The changes in E_a and $\log \tilde{A}$ for the C_2 contributions are limited to 1 kJ mol^{-1} and 0.1, respectively. A more pronounced temperature dependence is observed for the $C_1-(\text{S})(C_1)(C)$ group. This different behavior is caused by a free methyl rotor in the 3-sulfanylbuto-1-yn-3-yl radical,

Table 4 ΔGAV° s at 300 and 1000 K for abstractions of a hydrogen atom bonded to a carbon-centered group by a carbon-centered radical ($\Delta\text{GAV}^\circ_{\log \tilde{A}}$ in $\log(\text{m}^3 \text{mol}^{-1} \text{s}^{-1})$, $\Delta\text{GAV}^\circ_{E_a}$ in kJ mol^{-1})

$$\begin{array}{c} \text{X}_1 \\ | \\ \text{X}_2-\text{C}_1^\bullet \\ | \\ \text{X}_3 \end{array} + \text{H}-\begin{array}{c} \text{Y}_1 \\ | \\ \text{C}_2 \\ | \\ \text{Y}_2 \end{array}-\text{Y}_3 \rightarrow \begin{array}{c} \text{X}_1 \\ | \\ \text{X}_2-\text{C}_1-\text{H} \\ | \\ \text{X}_3 \end{array} + \begin{array}{c} \text{Y}_1 \\ | \\ \text{C}_2^\bullet \\ | \\ \text{Y}_2 \end{array}-\text{Y}_3$$

| | 300 K | | 1000 K | | 300 K | | 1000 K | |
|--|--------------------------------------|--------|------------------|--------|--------------------------------------|-------|------------------|-------|
| | $\log \tilde{A}$ | E_a | $\log \tilde{A}$ | E_a | $\log \tilde{A}$ | E_a | $\log \tilde{A}$ | E_a |
| Reference reaction | | | | | | | | |
| (1) $\text{C}^\bullet\text{H}_3 + \text{CH}_4$ | 5.267 | 69.7 | 6.509 | 84.4 | 5.267 | 69.7 | 6.509 | 84.4 |
| | $\Delta\text{GAV}^\circ(\text{C}_1)$ | | | | $\Delta\text{GAV}^\circ(\text{C}_2)$ | | | |
| | 300 K | | 1000 K | | 300 K | | 1000 K | |
| | $\log \tilde{A}$ | E_a | $\log \tilde{A}$ | E_a | $\log \tilde{A}$ | E_a | $\log \tilde{A}$ | E_a |
| Primary contributions | | | | | | | | |
| (2) $\text{C}_r(\text{S})(\text{H})_2$ | -0.672 | +20.1 | -0.570 | +21.1 | -0.564 | -20.9 | -0.542 | -20.8 |
| (3) $\text{C}_r(\text{S})(\text{C})(\text{H})$ | -0.485 | +18.1 | -0.247 | +20.4 | -0.110 | -28.5 | -0.053 | -28.1 |
| (4) $\text{C}_r(\text{S})(\text{C})_2$ | -1.086 | +15.5 | -0.724 | +18.8 | -0.155 | -36.0 | -0.072 | -35.5 |
| (5) $\text{C}_r(\text{CS})(\text{H})_2$ | +0.109 | +59.4 | +0.470 | +62.9 | -0.669 | -37.1 | -0.676 | -37.3 |
| (6) $\text{C}_r(\text{CS})(\text{C})(\text{H})$ | -0.151 | +70.5 | +0.232 | +74.1 | -0.123 | -43.6 | -0.088 | -43.5 |
| (7) $\text{C}_r(\text{CS})(\text{C})_2$ | -0.470 | +79.3 | -0.052 | +83.3 | -0.368 | -50.2 | -0.303 | -49.8 |
| (8) $\text{C}_{i,d}(\text{S})$ | +0.023 | -14.0 | +0.405 | -10.4 | +0.110 | -17.1 | +0.139 | -17.1 |
| (9) $\text{C}_{i,d}(\text{CS})$ | +0.316 | +41.8 | +0.645 | +45.0 | +0.497 | -6.1 | +0.611 | -5.3 |
| (10) $\text{C}_r(\text{S})(\text{C}_d)(\text{H})$ | +0.066 | +62.4 | +0.350 | +65.0 | -0.013 | -34.8 | +0.063 | -34.3 |
| (11) $\text{C}_r(\text{S})(\text{C}_d)(\text{C})$ | -0.840 | +60.9 | -0.487 | +64.2 | -0.247 | -41.3 | -0.160 | -40.7 |
| (12) $\text{C}_r(\text{S})(\text{CS})(\text{H})$ | +0.632 | +97.7 | +1.169 | +102.7 | +0.033 | -42.6 | +0.116 | -42.0 |
| (13) $\text{C}_r(\text{S})(\text{CS})(\text{C})$ | +0.103 | +102.4 | +0.651 | +107.5 | -0.078 | -47.6 | +0.019 | -46.9 |
| (14) $\text{C}_r(\text{S})(\text{C}_i)(\text{H})$ | -0.216 | +54.4 | +0.061 | +57.0 | -0.152 | -38.4 | -0.094 | -38.0 |
| (15) $\text{C}_r(\text{S})(\text{C}_i)(\text{C})$ | -0.829 | +54.8 | -0.220 | +61.0 | -0.145 | -44.6 | -0.055 | -44.0 |
| (16) $\text{C}_r(\text{S})(\text{C}_B)(\text{H})$ | -0.234 | +53.1 | +0.041 | +55.7 | -0.476 | -32.3 | -0.411 | -31.9 |
| (17) $\text{C}_r(\text{S})(\text{C}_B)(\text{C})$ | -0.698 | +48.1 | -0.331 | +51.6 | -0.771 | -38.0 | -0.692 | -37.5 |
| (18) $\text{CS}_r(\text{H})$ | +0.452 | +24.5 | +0.825 | +28.0 | +0.276 | -30.0 | +0.330 | -29.7 |
| (19) $\text{CS}_r(\text{C})$ | -0.060 | +22.0 | +0.387 | +26.3 | +0.147 | -33.9 | +0.187 | -33.8 |
| (20) $\text{CS}_r(\text{S})$ | +0.035 | +31.1 | +0.514 | +35.7 | +0.075 | -39.3 | +0.105 | -39.4 |
| (21) $\text{CS}_r(\text{C}_d)$ | -0.131 | +36.0 | +0.108 | +38.2 | +0.109 | -32.1 | +0.158 | -32.0 |
| (22) $\text{CS}_r(\text{C}_i)$ | +0.238 | +49.1 | +0.478 | +51.5 | +0.093 | -37.5 | +0.146 | -37.3 |
| (23) $\text{C}_{i,T}(\text{S})$ | +0.039 | -49.1 | +0.424 | -45.4 | +0.221 | +5.4 | +0.234 | +5.1 |
| (24) $\text{C}_{i,T}(\text{Cd})$ | -0.003 | -41.3 | +0.392 | -37.6 | +0.135 | +2.9 | +0.175 | +2.9 |
| Secondary contributions | | | | | | | | |
| (25) $\text{S}-(\text{C}_i(\text{H})_2)(\text{S})$ | 0 | 0 | 0 | 0 | +0.173 | -5.6 | +0.209 | -5.3 |

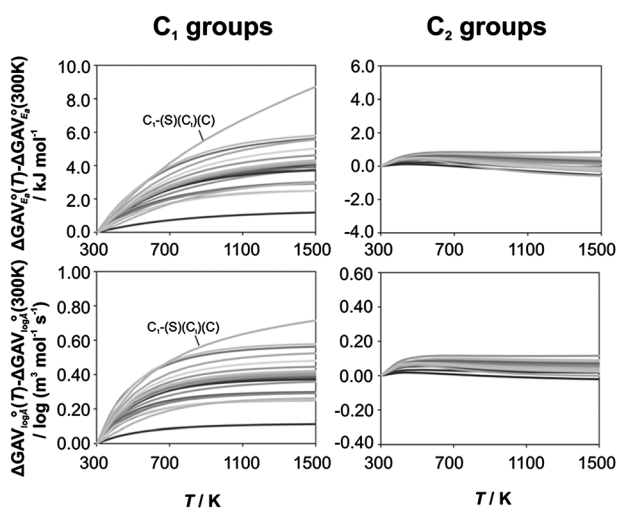


Fig. 6 Temperature dependence of the ΔGAV° s presented in Table 4.

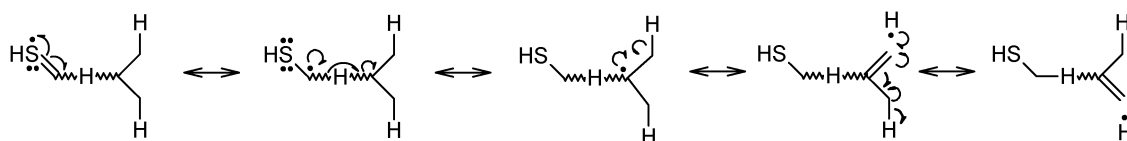
$\text{HC}\equiv\text{CC}^\bullet(\text{SH})\text{CH}_3$. The maximum error in the rate coefficient made by using the $\Delta\text{GAV}^\circ(300 \text{ K})$ s at different temperatures (up to 1500 K) amounts to a factor of 3. The limited temperature dependence illustrates that the ΔGAV° s can be used over a broad temperature range without a significant loss of accuracy.

3.3.2 The influence of secondary contributions and non-nearest-neighbor interactions. An elaborate study of the effect of secondary contributions and NNIs can be found in Table 2. In this table standard reaction enthalpies, single-event standard reaction entropies $\Delta_r S^\circ$ and Arrhenius parameters at 300 K are presented for 21 hydrogen abstraction reactions. Each reaction presented in Table 2 involves the hydrogen abstraction by a carbon-centered radical from a carbon atom having a neighboring sulfur atom, as illustrated by the reaction scheme presented at the top of this table. The 21 reactions are divided into six subsets according to the primary groups involved. Hydrogen abstractions within a subset of reactions have identical C_1 and C_2 groups but differ among each other by the Y and Z ligands (see the

reaction scheme at the top of Table 2). The influence of the Y ligands can be evaluated by comparing the Arrhenius parameters obtained for reactions (25)–(30) with reactions (3) and (8). The effect of different Z ligands is studied using reactions (31)–(41). The effect of secondary contributions and NNIs on the Arrhenius parameters can be quantified by studying the differences between the Arrhenius parameters for two reactions having identical C₁ and C₂ groups. Besides the influence of secondary contributions and NNIs on E_a and log \tilde{A} , also deviations of the tunneling coefficient, reaction enthalpy, $\Delta_r\tilde{S}^\circ$ and rate coefficient \tilde{k} were studied. The results of this study are shown in Table 5.

From Table 5 it can be seen that similar tunneling coefficients are obtained for reactions having the same C₁ and C₂ groups. Maximum deviations of the tunneling coefficients are limited to a factor of 1.25. Also for $\Delta_r\tilde{S}^\circ$ and Δ_rH° the differences between reactions having the same C₁ and C₂ groups are limited. The average deviations amount to -0.4 kJ mol^{-1} and $-1.7 \text{ J mol}^{-1} \text{ K}^{-1}$ for Δ_rH° and $\Delta_r\tilde{S}^\circ$, respectively. Maximum deviations are limited to 5 kJ mol^{-1} and $12 \text{ J mol}^{-1} \text{ K}^{-1}$. In general, good agreement is obtained between the Arrhenius parameters obtained within the same set of reactions, resulting in the fact that the \tilde{k} s are reproduced on average within a factor of 2.4. The mean absolute deviations of the log \tilde{A} and activation energy amount to 0.217 and 2.1 kJ mol^{-1} . This result shows that the influence of secondary contributions and NNIs on the studied hydrogen abstraction reactions is negligible. A truncated model including only primary contributions hence succeeds in accurately predicting Arrhenius parameters for hydrogen abstraction reactions involving organosulfur compounds.

The maximum deviation of the rate coefficients at 300 K amounts to a factor of 14.5 and is obtained for reaction (33) as compared to reaction (2). The large deviation is caused by an overestimation of the reaction barrier for the reverse reaction by 5.6 kJ mol^{-1} , *i.e.* the hydrogen abstraction from methane by HSSC•H₂. This overestimation of the reaction barrier is caused by a mesomeric effect: the sulfur atom in the β -position of the radical site helps to stabilize the radical site. This effect will mainly influence the barrier for the endothermic reaction. A similar, but smaller effect is also observed for reactions (35)–(37). For reactions (35)–(37) the stabilizing effect of the sulfur atom in the β -position of the radical site is less pronounced, as the C-centered radical site is also stabilized by hyperconjugation. The additional stabilization of the S-atom in the β -position of the radical site is therefore smaller and deviations are limited to approximately 2 kJ mol^{-1} . In general, additional ligands on the carbon atom bearing the radical site will reduce the influence of an S–S bond adjacent to the radical site.



To extend the applicability of the GA scheme to abstraction reactions involving C•H₂SSR type of radicals, one secondary ΔGAV° can be introduced accounting for the influence of an

S-atom in the β -position of the radical site, *i.e.* S–(C_i(H)₂)(S). The ΔGAV° s corresponding to this secondary contribution can be found at the bottom of Table 6. The data presented in Table 5 for reaction (33) show that for the forward reaction the influence on log \tilde{A} and E_a is limited to 0.027 and 1.2 kJ mol^{-1} , respectively. Therefore the $\Delta\text{GAV}^\circ(\text{S}-(\text{C}_i(\text{H})_2)(\text{S}))$ for the forward reaction can be set equal to 0. The contributions of the S–(C_i(H)₂)(S) group to the activation energy and log \tilde{A} for the reverse reaction amount to -5.6 kJ mol^{-1} and $+0.173$ at 300 K.

The data presented in Table 5 hence illustrate that secondary effects and NNIs have a limited influence on the reaction kinetics of hydrogen abstraction reactions. For all reactions studied, the maximum deviations of log \tilde{A} and E_a are limited to 0.67 and 6 kJ mol^{-1} , respectively. This result is in agreement with the observations of Saeys *et al.*²⁴ on the effect of secondary contributions and NNIs on the activation energies for radical addition reactions.

3.3.3 Resonance corrections. As explained in the methodology section, the resonance corrections account for stabilizing interactions in the transition state due to the simultaneous presence of groups on C₁ and C₂ and that cannot be included in the group-centered ΔGAV° s. In Table 3 Arrhenius parameters and rate coefficients were presented for 32 reactions for which stabilizing cross-interactions in the transition state are expected. The influence of those cross-interactions on the Arrhenius parameters can be quantified by comparing the *ab initio* calculated data with the GA modeled parameters using only the ΔGAV° s presented in Table 4.

$$\Delta E_{\text{res}}^\circ = E_a(T) - \left(E_{a,\text{ref}}(T) + \sum_{i=1}^2 \Delta\text{GAV}_{E_a}^\circ(C_i) \right) \quad (13)$$

$$\Delta \log \tilde{A}_{\text{res}}^\circ = \log \tilde{A}(T) - \left(\log \tilde{A}_{\text{ref}}(T) + \sum_{i=1}^2 \Delta\text{GAV}_{\log \tilde{A}}^\circ(C_i) \right) \quad (14)$$

For example, the Arrhenius parameters for the hydrogen abstraction reaction $\text{CH}_3\text{SH} + \text{CH}_3\text{C}\cdot\text{HCH}_3 \rightarrow \text{C}_2\cdot\text{HSH} + \text{C}_3\text{H}_8$ can be estimated from the data presented in Table 4 and ref. 26. The two primary groups involved in this reaction are the C₁–(C)₂(H) and C₂–(S)(H)₂ groups. The GA estimated activation energy at 300 K hence amounts to $69.7 + 6.5 - 20.9 = 55.3 \text{ kJ mol}^{-1}$. This is 5 kJ mol^{-1} higher than the calculated value which amounts to 50.3 kJ mol^{-1} (reaction (58)). This deviation between calculated and GA estimated activation energy is caused by resonance stabilization. For the transition state, the following resonance structures represent the electron delocalization:

These resonance structures describe the interaction of one of the lone electron pairs of the sulfur atom and two hyperconjugating C–H bonds with the partially filled orbital. As illustrated in Fig. S7

Table 5 Effect of secondary contributions and non-nearest-neighbor interactions on the tunneling coefficients κ , standard reaction enthalpies $\Delta_r H^\circ$ (kJ mol⁻¹), single-event standard reaction entropies $\Delta_r \tilde{S}^\circ$ (J mol⁻¹ K⁻¹), single-event pre-exponential factors $\log \tilde{A}$ (log(m³ mol⁻¹ s⁻¹)), activation energies E_a (kJ mol⁻¹) and single-event rate coefficients \tilde{k} (m³ mol⁻¹ s⁻¹) at 300 K for the reactions presented in Table 2. The reported values are the deviations between the calculated data for a target reaction and the corresponding data obtained for the smallest possible reaction having the same C₁ and C₂ groups (=training reaction, subscript t). The training reactions are indicated by —

$$\begin{array}{c} \text{X}_1 \\ | \\ \text{X}_2-\text{C}_1^\bullet \\ | \\ \text{X}_3 \end{array} + \text{H}-\begin{array}{c} \text{Y}_1 \\ | \\ \text{C}_2 \\ | \\ \text{Y}_2 \end{array} -\text{S}-\text{Z} \longrightarrow \text{X}_2-\begin{array}{c} \text{X}_1 \\ | \\ \text{C}_1-\text{H} \\ | \\ \text{X}_3 \end{array} + \begin{array}{c} \text{Y}_1 \\ | \\ \text{C}_2^\bullet \\ | \\ \text{Y}_2 \end{array} -\text{S}-\text{Z}$$

| | | Forward | | | | | Reverse | | | | |
|--------------------------|---|-------------------|--------------------------|----------------------------------|-------------------------|--------------|-------------------------|-------------------------|--------------|-------------------------|--|
| | | κ/κ_t | $\Delta\Delta_r H^\circ$ | $\Delta\Delta_r \tilde{S}^\circ$ | $\Delta \log \tilde{A}$ | ΔE_a | \tilde{k}/\tilde{k}_t | $\Delta \log \tilde{A}$ | ΔE_a | \tilde{k}/\tilde{k}_t | |
| Y-group | | | | | | | | | | | |
| (3) | $\text{CH}_3^\bullet + \text{HS}-\text{CH}_3 \rightleftharpoons \text{CH}_4 + \text{HS}-\dot{\text{C}}\text{H}_3$ | — | — | — | — | — | — | — | — | — | |
| (25)–(3) | $\text{CH}_3^\bullet + \text{HS}-\text{CH}_2\text{CH}_3 \rightleftharpoons \text{CH}_4 + \text{HS}-\dot{\text{C}}\text{H}_2\text{CH}_3$ | 1.09 | 1.0 | -1.3 | -0.173 | -0.7 | 0.96 | -0.111 | -1.7 | 1.66 | |
| (26)–(3) | $\text{CH}_3^\bullet + \text{HS}-\text{CH}_2\text{CH}_2\text{CH}_3 \rightleftharpoons \text{CH}_4 + \text{HS}-\dot{\text{C}}\text{H}_2\text{CH}_2\text{CH}_3$ | 1.05 | -0.2 | -0.9 | -0.067 | -1.3 | 1.54 | -0.024 | -1.2 | 1.62 | |
| (27)–(3) | $\text{CH}_3^\bullet + \text{HS}-\text{CH}(\text{CH}_3)_2 \rightleftharpoons \text{CH}_4 + \text{HS}-\dot{\text{C}}\text{H}(\text{CH}_3)_2$ | 1.14 | 1.4 | -1.7 | -0.243 | -0.4 | 0.77 | -0.158 | -1.8 | 1.65 | |
| (28)–(3) | $\text{CH}_3^\bullet + \text{HS}-\text{C}(\text{CH}_3)_3 \rightleftharpoons \text{CH}_4 + \text{HS}-\dot{\text{C}}(\text{CH}_3)_3$ | 1.25 | 4.7 | 1.8 | -0.264 | -0.6 | 0.86 | -0.364 | -5.3 | 4.53 | |
| (8) | $\text{CH}_3^\bullet + \text{HS}-\text{CH}=\text{CH}_2 \rightleftharpoons \text{CH}_4 + \text{HS}-\dot{\text{C}}\text{H}=\text{CH}_2$ | — | — | — | — | — | — | — | — | — | |
| (29)–(8) | $\text{CH}_3^\bullet + \text{HS}-\text{CH}=\text{CH}-\text{CH}_3 \rightleftharpoons \text{CH}_4 + \text{HS}-\dot{\text{C}}\text{H}=\text{CH}-\text{CH}_3$ | 0.96 | 0.2 | -4.9 | -0.319 | -0.6 | 0.57 | -0.062 | -0.7 | 1.11 | |
| (30)–(8) | $\text{CH}_3^\bullet + \text{HS}-\text{C}(\text{CH}_3)=\text{CH}_2 \rightleftharpoons \text{CH}_4 + \text{HS}-\dot{\text{C}}(\text{CH}_3)=\text{CH}_2$ | 1.05 | -4.1 | -8.0 | -0.671 | -4.4 | 1.29 | -0.260 | -0.3 | 0.65 | |
| Z-group | | | | | | | | | | | |
| (2) | $\text{CH}_3^\bullet + \text{HS}-\text{S}-\text{S}-\text{S} \rightleftharpoons \text{CH}_4 + \text{HS}-\dot{\text{S}}-\text{S}-\text{S}$ | — | — | — | — | — | — | — | — | — | |
| (31)–(2) | $\text{CH}_3^\bullet + \text{S}-\text{S}-\text{S} \rightleftharpoons \text{CH}_4 + \text{S}-\dot{\text{S}}-\text{S}$ | 0.82 | -5.2 | -2.4 | -0.051 | -2.6 | 2.11 | 0.070 | 2.5 | 0.35 | |
| (32)–(2) | $\text{CH}_3^\bullet + \text{S}-\text{S}-\text{S} \rightleftharpoons \text{CH}_4 + \text{S}-\dot{\text{S}}-\text{S}$ | 1.00 | -0.9 | -3.4 | 0.447 | 3.2 | 0.76 | 0.617 | 4.1 | 0.78 | |
| (33)–(2) | $\text{CH}_3^\bullet + \text{HS}-\text{S}-\text{S} \rightleftharpoons \text{CH}_4 + \text{HS}-\dot{\text{S}}-\text{S}$ | 1.02 | 4.4 | -3.9 | -0.027 | -1.2 | 1.53 | 0.173 | -5.6 | 14.5 | |
| (3) | $\text{CH}_3^\bullet + \text{HS}-\text{CH}_2\text{S} \rightleftharpoons \text{CH}_4 + \text{HS}-\dot{\text{C}}\text{H}_2\text{S}$ | — | — | — | — | — | — | — | — | — | |
| (34)–(3) | $\text{CH}_3^\bullet + \text{S}-\text{CH}_2\text{S} \rightleftharpoons \text{CH}_4 + \text{S}-\dot{\text{C}}\text{H}_2\text{S}$ | 0.85 | -5.3 | 1.1 | -0.049 | -2.8 | 2.28 | -0.106 | 2.6 | 0.24 | |
| (35)–(3) | $\text{CH}_3^\bullet + \text{HS}-\text{S}-\text{CH}_2\text{S} \rightleftharpoons \text{CH}_4 + \text{HS}-\dot{\text{S}}-\text{CH}_2\text{S}$ | 1.10 | 2.7 | 5.8 | -0.012 | -0.8 | 1.46 | -0.314 | -3.4 | 2.11 | |
| (36)–(3) | $\text{CH}_3^\bullet + \text{S}-\text{S}-\text{CH}_2\text{S} \rightleftharpoons \text{CH}_4 + \text{S}-\dot{\text{S}}-\text{CH}_2\text{S}$ | 1.09 | 0.9 | -11.7 | -0.130 | -0.2 | 0.88 | 0.457 | -1.3 | 5.17 | |
| (37)–(3) | $\text{CH}_3^\bullet + \text{S}-\text{S}-\text{S}-\text{CH}_2\text{S} \rightleftharpoons \text{CH}_4 + \text{S}-\dot{\text{S}}-\text{S}-\text{CH}_2\text{S}$ | 1.11 | 2.0 | 4.3 | -0.172 | -0.2 | 0.80 | -0.395 | -2.2 | 1.08 | |
| (38)–(3) | $\text{CH}_3^\bullet + \text{S}-\text{S}-\text{S} \rightleftharpoons \text{CH}_4 + \text{S}-\dot{\text{S}}-\text{S}$ | 1.04 | -1.2 | 1.6 | 0.355 | 2.8 | 0.78 | 0.271 | 4.0 | 0.40 | |
| (10) | $\text{CH}_3^\bullet + \text{HS}-\text{CH}=\text{CH}-\text{S} \rightleftharpoons \text{CH}_4 + \text{HS}-\dot{\text{C}}\text{H}=\text{CH}-\text{S}$ | — | — | — | — | — | — | — | — | — | |
| (39)–(10) | $\text{CH}_3^\bullet + \text{S}-\text{CH}=\text{CH}-\text{S} \rightleftharpoons \text{CH}_4 + \text{S}-\dot{\text{C}}\text{H}=\text{CH}-\text{S}$ | 0.93 | -2.5 | -4.5 | 0.038 | -0.4 | 1.20 | 0.267 | 2.0 | 0.77 | |
| (40) | $\text{---} + \text{HS}-\text{CH}_2\text{S} \rightleftharpoons \text{---} + \text{HS}-\dot{\text{C}}\text{H}_2\text{S}$ | — | — | — | — | — | — | — | — | — | |
| (41)–(40) | $\text{---} + \text{S}-\text{CH}_2\text{S} \rightleftharpoons \text{---} + \text{S}-\dot{\text{C}}\text{H}_2\text{S}$ | 0.85 | -5.3 | 1.1 | -0.115 | -3.0 | 2.15 | -0.171 | 2.3 | 0.23 | |
| MAD log <i>A</i> | | | | | | | | | | 0.217 | |
| MAD <i>E_a</i> | | | | | | | | | | 2.1 | |
| $\langle \rho \rangle$ | | | | | | | | | | 2.4 | |

of the ESI[†], the spin density is delocalized on the sulfur atom and on the hydrogen atoms of the hyperconjugating C–H bonds.

The effect of transition state resonance on both Arrhenius parameters for the reactions in Table 3 is presented in Table 6. The reactions are grouped according to the different types of

cross-interactions introduced in Section 2.2. Table 6 illustrates that, on average, resonance stabilization lowers the activation energies and pre-exponential factors by 15 kJ mol⁻¹ and 0.319. Similar deviations of the Arrhenius parameters are observed for reactions (42), (46) and (49), and, hence, one single correction ($p_{\alpha S}-\pi_{C=C/\equiv}$) can be introduced to account for the cross-interaction of an α -sulfur atom on C₁ with a vinyl, ethynyl or phenyl ligand on C₂. The smallest effects are obtained for reactions (58) and (59), whose transition states are stabilized by the $p_{\alpha S}-\sigma_{\beta C-H}$ cross-interaction of an α -sulfur atom and hyperconjugating C–H bonds. From the data obtained for reactions (58) and (59) and reactions (54)–(57), it can be seen that $|\Delta E_{a,res}^{\circ}|$ increases with an increasing number of methyl ligands. This increase amounts to approximately 2 kJ mol⁻¹ for each additional methyl ligand. The largest effect on the activation energy amounts up to 34 kJ mol⁻¹ and is obtained for reaction (68), *i.e.* the hydrogen abstraction of propanethial (CH₃CH₂C(=S)H) by the 1-thioxopropan-2-yl radical (CH₃C[•]HC(=S)H). The transition state of this reaction is stabilized both by the $\pi_{C=S}-\pi_{C=S}$ cross-interaction of two thiocarbonyl groups and by the $\pi_{C=S}-\sigma_{\beta C-H}$ cross-interaction of an adjacent thiocarbonyl group on the C₁ atom with a β C–H bond on the C₂ atom. The data in Table 6 clearly show that resonance stabilization also influences the pre-exponential factors. Due to the occurrence of resonance in the transition state, the relative motion of the two reactants in the transition state is hampered, leading to lower pre-exponential factors. In analogy with the activation energy, correction factors are therefore presented to model $\log \tilde{A}$.

From a least squares regression of the 32 Arrhenius parameters presented in Table 3, the 8 supplementary corrections required to account for cross-interactions in the transition state were derived. The results of this regression are presented in Table 7. The values for corrections $\pi_{C=C}-\pi_{C=C/\equiv}$, $\pi_{\equiv}-\pi_{\equiv}$, $\pi_{C=C/\equiv}-\sigma_{\beta C-H}$ and $\sigma_{\beta C-H}-\sigma_{\beta C-H}$ were taken from previous work.²⁶ The largest stabilizing effects are obtained for the cross-interaction of a thiocarbonyl group with another π -conjugating system (see Fig. S1–S3 of the ESI†). Activation energies can lower more than 20 kJ mol⁻¹ due to the cross-interaction of a thiocarbonyl group on C₁ and C₂. Also on $\log \tilde{A}$ an important influence is observed. Pre-exponential factors can decrease up to a factor of 3 due to the simultaneous presence of a thiocarbonyl group on C₁ and C₂. The smallest correction factors are those describing the cross-interaction of a hyperconjugating C–H bond with an α -sulfur atom ($p_{\alpha S}-\sigma_{\beta C-H}$) or a thiocarbonyl group ($\pi_{C=S}-\sigma_{\beta C-H}$). In particular, the contribution of $p_{\alpha S}-\sigma_{\beta C-H}$ to $\log \tilde{A}$ is close to 0 and can be neglected. Its contribution to the activation energy amounts to 2 kJ mol⁻¹ per interaction and can hence still be important, especially in cases where multiple corrections are needed to describe the occurring resonance. The data in Table 7 further illustrate that the correction factors are almost temperature independent. The deviations between the corrections obtained at 300 K and 1000 K amount to at most 0.6 kJ mol⁻¹ and 0.07 for E_a and $\log \tilde{A}$, respectively.

In Table 6 it is shown that by using the corrections presented in Table 7, the activation energies and pre-exponential factors can be reproduced within 0.9 kJ mol⁻¹ and 0.099, respectively. The application of the resonance correction method is illustrated

for one reaction, *i.e.* the hydrogen abstraction reaction by the propargyl radical from 2-propene-1-thiol (reaction (53)): HC≡C–CH[•]₂ + H₂C=CHCH₂SH → HC≡C–CH₃ + H₂C=CHC[•]HSH. The transition state of this reaction is presented in Fig. 7. In this figure it can be seen that the transition state is stabilized by the cross-interaction of (a) an ethynyl ligand on C₁ with an α -S atom on C₂ and (b) an ethynyl ligand on C₁ with an ethenyl ligand on C₂. Using the presented group additivity model the activation energy and pre-exponential factor of reaction (53) are hence calculated as:

$$\begin{aligned} E_a(300\text{ K}) &= E_{a,\text{ref}}(300\text{ K}) + \Delta GAV_{E_a}^{\circ}(C_1-(C_t)(H)_2) \\ &\quad + \Delta GAV_{E_a}^{\circ}(C_2-(S)(C_d)(H)) \\ &\quad + \pi_{C=C}-\pi_{C=C/\equiv} + p_{\alpha S}-\pi_{C=C/\equiv} \\ &= 69.7 + 34.6 - 34.8 - 9.8 - 8.6 \\ &= 51.1\text{ kJ mol}^{-1} \end{aligned}$$

$$\begin{aligned} \log A(300\text{ K}) &= \log \tilde{A}_{\text{ref}}(300\text{ K}) + \Delta GAV_{\log A}^{\circ}(C_1-(C_t)(H)_2) \\ &\quad + \Delta GAV_{\log A}^{\circ}(C_2-(S)(C_d)(H)) \\ &\quad + \pi_{C=C}-\pi_{C=C/\equiv} + p_{\alpha S}-\pi_{C=C/\equiv} + \log(2) \\ &= 5.267 - 0.083 - 0.013 - 0.180 - 0.206 \\ &\quad + 0.301 = 5.086 \end{aligned}$$

These values correspond well to the *ab initio* values of 55.0 kJ mol⁻¹ and 5.086 for E_a and $\log A$, respectively.

3.3.4 Tunneling. By using eqn (9) the Eckart transmission coefficients for all hydrogen abstraction reactions presented in Tables 1–3 are reproduced on average within 10%. The maximum deviation amounts to a factor of 2.7 and is obtained for reaction (62), *i.e.* the hydrogen abstraction by 3-methylbut-1-en-3-yl from ethanethial. In Fig. 8, a parity plot is shown in which the tunneling coefficients obtained with eqn (9) are presented *versus* those calculated using the Eckart scheme at 300 K and tunneling coefficients for all reactions are presented in Tables 1–3. At 300 K the Eckart tunneling coefficients are reproduced on average within a factor of 1.5. 90% of the tunneling coefficients are reproduced within a factor of 2. The agreement between modeled and calculated tunneling coefficients increases with increasing temperatures. At 1000 K the average deviation amounts to 3%, with a maximum deviation of 10%.

3.4 Validation

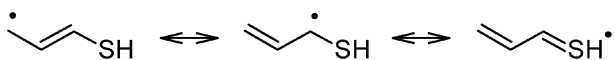
The presented ΔGAV° s are validated using a test set of *ab initio* data for 30 reactions. Arrhenius parameters and rate coefficients for the test set in the temperature range 300–1500 K can be found in the ESI.† In Table 8 the deviations of the transmission coefficients, pre-exponential factor, activation energy and rate coefficients between the *ab initio* calculated and GA modeled values are presented at 300 K. The MADs of $\log A$ and E_a amount to 0.132 and 2.5 kJ mol⁻¹, respectively. The average deviation of the calculated rate coefficients amounts to a factor 2.4, indicating that the presented GA model reproduces the calculated rate

Table 6 (continued)

| | | Resonance effect | | Remaining deviation | |
|------|--|--------------------------------------|----------------------------|---------------------|-----------------------|
| | | $\Delta \log \tilde{A}_{\text{res}}$ | $\Delta E_{\text{a, res}}$ | $\Delta \log A$ | ΔE_{a} |
| (68) | | -0.556 | -34.3 | 0.085 | -0.4 |
| (69) | | -0.526 | -32.8 | 0.066 | 0.8 |
| 7. | | | | | |
| (70) | | -0.449 | -11.5 | -0.097 | 1.0 |
| (71) | | -0.251 | -15.5 | 0.097 | -1.0 |
| 8. | | | | | |
| (72) | | -0.300 | -7.6 | -0.215 | -1.8 |
| (73) | | -0.387 | -18.3 | -0.130 | -0.8 |
| MAD | | 0.319 | 15.3 | 0.099 | 0.9 |

coefficients accurately. The agreement between *ab initio* and group additively estimated rate coefficients improves with increasing temperature and at 1000 K the rate coefficients are reproduced within a factor of 1.4. At 1000 K the largest deviation of the rate coefficients is limited to a factor of 3.

Table 8 (third column) also shows the ratio of the estimated transmission coefficient (using eqn (9)) and calculated Eckart transmission coefficient ($\kappa/\kappa_{\text{AI}}$). It is seen that the transmission coefficients are reproduced accurately. The average deviation between the estimated and calculated κ amounts to a factor of 1.4 at 300 K. The maximum deviation is limited to a factor of 2.8 and is obtained for reaction (85). The data in Table 8 further show that good agreement is obtained between the *ab initio* calculated kinetic parameters and the ones obtained by using the group additivity model. The largest deviation of $\log A$ is less than 0.5 while all activation energies are reproduced within 10 kJ mol⁻¹. The largest deviations of E_{a} are obtained for hydrogen abstraction reactions by radicals that are stabilized by resonance. In these cases, special attention needs to be given to the thermochemistry of the reaction. For example, reaction (75) leads to the formation of a 3-sulfanyl-1-en-3-yl radical:



The $\Delta_f H^\circ$ of this radical is obtained by calculating the $\Delta_f H^\circ$ s for the different resonance structures and then selecting the lowest value obtained.⁵⁵ Using the GAVs presented in literature,⁵⁵ the lowest $\Delta_f H^\circ$ is obtained for CH₂=CHC•HSH which is 11.9 kJ mol⁻¹ lower than the one obtained for C•H₂CH=CHSH. The $\Delta\text{GAV}^\circ(\text{C}_1-(\text{C}_d)(\text{H})_2)$ contribution used to estimate the activation energy for reaction (75) is obtained from a hydrogen abstraction by an allylic radical. It describes the enthalpy change between the C•-(C_d)(H)₂ group in the reactant and the C-(C_d)(H)₂ group in the transition state. Hence, $\Delta\text{GAV}^\circ(\text{C}_1-(\text{C}_d)(\text{H})_2)$ describes the enthalpy change going from C•H₂CH=CHSH to the transition state. However, as the C•H₂CH=CHSH resonance structure is not the most stable structure, an additional

contribution needs to be added, accounting for the enthalpy difference between CH₂=CHC•HSH and C•H₂CH=CHSH. By simply using $\Delta\text{GAV}^\circ(\text{C}_1-(\text{C}_d)(\text{H})_2)$, the GA model underestimates the barrier for reaction (75) by 9.5 kJ mol⁻¹, resulting in an overestimation of the rate coefficient at 300 K by a factor of 150. A correction for the enthalpy difference between the two resonance structures allows us to reproduce the barrier within 2.4 kJ mol⁻¹ and the rate coefficients within a factor of 1.5 (see bracketed values in Table 8 for reaction (75)). Problems of this nature can be avoided by enforcing thermodynamic consistency in the GA model. In this case, the Arrhenius parameters for the exothermic reaction are estimated using the GA scheme, while those for the endothermic reactions are subsequently obtained from thermodynamic consistency.

The reactions presented in Table 8 can be divided into five groups. The first group of reactions (reactions (74)–(76)) are hydrogen abstraction reactions having a sulfur atom in the β - or γ -position of the reactive center. For these three reactions the ΔGAV° s derived for hydrocarbons by Sabbe *et al.*²⁶ can be used. From Table 8 it is seen that using these ΔGAV° s the rate coefficients can be reproduced accurately, provided that thermodynamic consistency is enforced on reaction (75). The second group of reactions are hydrogen abstraction reactions by ethyl (reactions (40) and (77)–(80)). For these five reactions the GA model succeeds in reproducing all the rate coefficients within a factor of 4. The maximum deviations of $\log A$ and E_{a} are small and are restricted to 0.3 and 3 kJ mol⁻¹, respectively. For reactions (81)–(84), *i.e.* hydrogen abstraction reactions from dimethyl disulfide, larger deviations between calculated and GA data are observed. The barriers are on average reproduced within 2 kJ mol⁻¹, resulting in rate coefficients which are reproduced within a factor of 1.3 to 2.5. In case no additional correction for secondary effects in methyl disulfides was introduced, the reverse barriers would have been systematically overestimated by on average 5 kJ mol⁻¹ leading to an underestimation of the rate coefficients for the reverse reactions by almost a factor of 10 at 300 K. The fourth group of reactions, *i.e.* reactions (85)–(91), are hydrogen abstraction

Table 7 Correction factor to account for additional resonance and hyperconjugative stabilization in the transition state at 300, 600 and 1000 K (E_a in kJ mol^{-1}). *Italic values were taken from the work of Sabbe et al.*²⁶

| Correction | Corresponding structure | 300 K | | 600 K | | 1000 K | | Average | |
|---|-------------------------|------------------|-------|------------------|-------|------------------|-------|------------------|-------|
| | | $\log \tilde{A}$ | E_a | $\log \tilde{A}$ | E_a | $\log \tilde{A}$ | E_a | $\log \tilde{A}$ | E_a |
| (1) $\pi_{\text{C}=\text{C}}-\pi_{\text{C}=\text{C}}/\equiv$ | | -0.180 | -9.8 | -0.202 | -10.0 | -0.212 | -10.1 | -0.198 | -10.0 |
| (2) $\pi_{\equiv}-\pi_{\equiv}$ | | -0.067 | -6.2 | -0.003 | -5.7 | 0.008 | -5.5 | -0.021 | -5.8 |
| (3) $\pi_{\text{C}=\text{C}}/\equiv-\sigma_{\beta\text{C}-\text{H}}$ | | -0.061 | -3.4 | -0.052 | -3.3 | -0.053 | -3.4 | -0.055 | -3.4 |
| (4) $\sigma_{\beta\text{C}-\text{H}}-\sigma_{\beta\text{C}-\text{H}}$ | | -0.049 | -0.3 | -0.050 | -0.3 | -0.051 | -0.3 | -0.050 | -0.3 |
| (5) $p_{2s}-\pi_{\text{C}=\text{C}}/\equiv$ | | -0.206 | -8.6 | -0.207 | -8.6 | -0.217 | -8.7 | -0.210 | -8.6 |
| (6) $p_{2s}-p_{2s}$ | | -0.366 | -7.1 | -0.299 | -6.6 | -0.293 | -6.5 | -0.319 | -6.7 |
| (7) $p_{2s}-\sigma_{\beta\text{C}-\text{H}}$ | | 0.005 | -2.0 | 0.006 | -2.0 | 0.005 | -2.0 | 0.005 | -2.0 |
| (8) $\pi_{\text{C}=\text{S}}-\pi_{\text{C}=\text{C}}$ | | -0.278 | -18.4 | -0.292 | -18.6 | -0.306 | -18.8 | -0.292 | -18.6 |
| (9) $\pi_{\text{C}=\text{S}}-\pi_{\equiv}$ | | -0.184 | -11.1 | -0.181 | -11.1 | -0.188 | -11.2 | -0.184 | -11.1 |
| (10) $\pi_{\text{C}=\text{S}}-\pi_{\text{C}=\text{S}}$ | | -0.421 | -21.9 | -0.370 | -21.6 | -0.374 | -21.6 | -0.388 | -21.7 |
| (11) $\pi_{\text{C}=\text{S}}-p_{2s}$ | | -0.352 | -12.5 | -0.400 | -12.9 | -0.417 | -13.1 | -0.390 | -12.8 |
| (12) $\pi_{\text{C}=\text{S}}-\sigma_{\beta\text{C}-\text{H}}$ | | -0.085 | -5.8 | -0.099 | -6.0 | -0.108 | -6.1 | -0.097 | -6.0 |

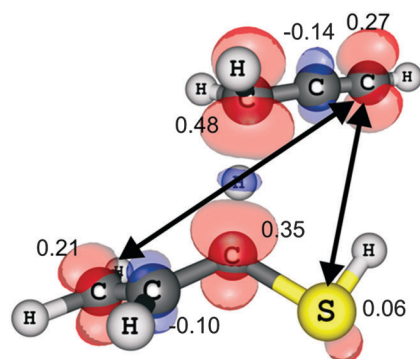


Fig. 7 Spin density plot and Mulliken atomic spin densities illustrating the resonance stabilization in the transition state for the hydrogen abstraction reaction by a propargyl radical from 2-propen-1-thiol (reaction (53)) (B3LYP/6-311G(2d,d,p), 0.006 isosurface).

reactions involving allylic and propargyl radicals. Five of the reactions studied in this group have transition state stabilization corrections ranging up to 15 kJ mol^{-1} and higher. It is seen that the GA model again succeeds in accurately reproducing the rate coefficients for these reactions; the rate coefficients are predicted within an average factor of 2.5, with a maximum of 5. The last group pertains to hydrogen abstractions from methanethiol (reactions (92)–(102)). Once more the GA model succeeds in reproducing the calculated rate coefficients for all reactions in this group on average within a factor of 2.8.

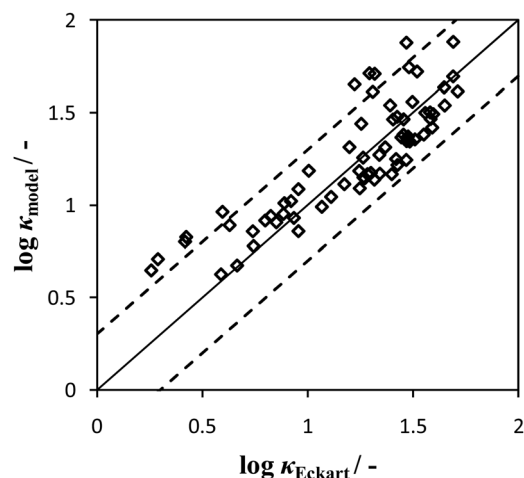


Fig. 8 Parity plot between Eckart tunneling coefficients at 300 K for the reactions presented in Tables 1–3 and those estimated using eqn (9). The two dashed lines indicate deviations by a factor of 2.

The use of the group additivity method is illustrated for two reactions, *i.e.* the abstraction of a hydrogen atom from methyl ethyl sulfide by ethyl (reaction (77)) and the hydrogen abstraction reaction from methanethiol by prop-2-en-1-thiol-1-yl (reaction (97)). The Arrhenius parameters at 300 K for the reference reaction, *i.e.* hydrogen abstraction from methane by methyl, can be found in Table 4 and amount to 5.267 and 69.7 kJ mol^{-1} , for, respectively, $\log \tilde{A}$ and E_a . The primary

Table 8 Comparison between *ab initio* and GA kinetic parameters at 300 K for a set of 30 reactions. ($\Delta \log A = \log A_{GA} - \log A_{AI}$ in $\log(\text{m}^3 \text{mol}^{-1} \text{s}^{-1})$, $\Delta E_a = E_{a,GA} - E_{a,AI}$ in kJ mol^{-1})

| | | Forward | | | | Reverse | | |
|-------|------------------------|----------------------|-----------------|--------------|-----------------|-----------------|--------------|-----------------|
| | | κ/κ_{AI} | $\Delta \log A$ | ΔE_a | k_{GA}/k_{AI} | $\Delta \log A$ | ΔE_a | k_{GA}/k_{AI} |
| (74) | | 1.29 | 0.461 | 2.6 | 1.4 | -0.350 | 0.8 | 0.4 |
| (75) | | 1.39 | -0.019 | 2.4 | 0.5 | 0.353 | -9.5 (2.4) | 151.7 (1.4) |
| (76) | | 1.24 | -0.064 | 1.6 | 0.6 | 0.143 | 5.2 | 0.2 |
| (40) | | 0.60 | -0.008 | 0.1 | 0.6 | -0.008 | 0.1 | 0.6 |
| (77) | | 0.70 | 0.106 | 3.1 | 0.3 | 0.163 | -2.2 | 2.7 |
| (78) | | 0.93 | -0.016 | 2.8 | 0.3 | -0.016 | 2.8 | 0.3 |
| (79) | | 0.83 | 0.251 | 0.7 | 1.2 | 0.251 | 0.7 | 1.2 |
| (80) | | 0.89 | 0.283 | 0.0 | 1.8 | 0.182 | -1.6 | 2.8 |
| (81) | | 0.85 | 0.066 | 1.2 | 0.6 | -0.002 | -2.0 | 2.0 |
| (82) | | 0.78 | 0.154 | 2.3 | 0.5 | 0.086 | -0.9 | 1.4 |
| (83) | | 0.90 | 0.019 | 2.2 | 0.4 | -0.048 | -1.0 | 1.3 |
| (84) | | 0.96 | 0.350 | 3.6 | 0.6 | 0.283 | 0.3 | 1.7 |
| (85) | | 0.37 | -0.035 | -5.7 | 3.6 | -0.035 | -5.7 | 3.6 |
| (86) | | 0.53 | -0.042 | -3.5 | 2.1 | -0.042 | -3.5 | 2.1 |
| (87) | | 0.86 | 0.045 | 3.5 | 0.2 | 0.045 | 3.5 | 0.3 |
| (88) | | 0.54 | -0.098 | -1.4 | 0.8 | -0.098 | -1.4 | 0.8 |
| (89) | | 0.71 | 0.246 | -0.2 | 1.5 | 0.246 | -0.2 | 1.4 |
| (90) | | 0.67 | 0.031 | -2.4 | 2.1 | 0.031 | -2.4 | 2.0 |
| (91) | | 1.14 | 0.087 | 5.0 | 0.2 | 0.087 | 5.0 | 0.2 |
| (92) | | 0.84 | 0.129 | 1.1 | 0.8 | 0.129 | 1.1 | 0.8 |
| (93) | | 0.67 | 0.119 | 2.2 | 0.4 | 0.119 | 2.2 | 0.4 |
| (94) | | 0.89 | 0.231 | 5.8 | 0.2 | 0.231 | 5.8 | 0.2 |
| (95) | | 2.05 | 0.078 | 1.9 | 1.2 | 0.078 | 1.9 | 1.2 |
| (96) | | 0.97 | 0.179 | 5.5 | 0.2 | 0.179 | 5.5 | 0.2 |
| (97) | | 0.57 | -0.189 | -3.0 | 1.3 | -0.189 | -3.0 | 1.3 |
| (98) | | 0.98 | 0.250 | 0.1 | 1.8 | 0.250 | 0.0 | 1.8 |
| (99) | | 0.79 | -0.187 | -3.8 | 2.6 | -0.187 | -3.8 | 2.5 |
| (100) | | 0.81 | -0.014 | -0.5 | 1.0 | -0.014 | -0.5 | 1.0 |
| (101) | | 1.79 | 0.077 | 2.5 | 0.8 | 0.077 | 2.5 | 0.8 |
| (102) | | 0.64 | 0.078 | -1.6 | 1.6 | 0.078 | -1.6 | 1.6 |
| | MAD log A | | | | | | | 0.132 |
| | MAD E _a | | | | | | | 2.5 |
| | $\langle \rho \rangle$ | | | | | | | 2.4 |

groups for reaction (77) are $C_1-(C)(H)_2$ and $C_2-(S)(C)(H)$. Their contributions to the activation energy amount to, respectively, +4.7 and -28.5 kJ mol^{-1} while their contributions to $\log \tilde{A}$ are, respectively, -0.866 and -0.110. The number of single events for this reaction amounts to four. This number originates from the two equivalent sides of the ethyl radical and from chirality in the transition state. The transition state of this reaction is stabilized by the cross-interaction of one hyperconjugating methyl group on the C_1 group with one sulfur atom and one methyl group on the C_2 group. The activation energy and pre-exponential factor are hence calculated as:

$$\begin{aligned} E_a(300 \text{ K}) &= E_{a,\text{ref}}(300 \text{ K}) + \Delta GAV_{E_a}^{\circ}(C_1-(C)(H)_2) \\ &+ \Delta GAV_{E_a}^{\circ}(C_2-(S)(C)(H)) \\ &+ \sigma_{\beta C-H} - \sigma_{\beta C-H} + p_{\alpha S} - \sigma_{\beta C-H} \\ &= 69.7 + 4.7 - 28.5 - 0.3 - 2.0 \\ &= 43.6 \text{ kJ mol}^{-1} \end{aligned}$$

$$\begin{aligned} \log A(300 \text{ K}) &= \log \tilde{A}_{\text{ref}}(300 \text{ K}) + \Delta GAV_{\log \tilde{A}}^{\circ}(C_1-(C)(H)_2) \\ &+ \Delta GAV_{\log \tilde{A}}^{\circ}(C_2-(S)(C)(H)) \\ &+ \sigma_{\beta C-H} - \sigma_{\beta C-H} + p_{\alpha S} - \sigma_{\beta C-H} + \log(4) \\ &= 5.267 - 0.866 - 0.110 - 0.049 + 0.005 \\ &+ 0.602 = 4.849 \end{aligned}$$

In order to obtain reliable tunneling correction for this reaction, the exothermic reaction direction has to be identified. Using the ΔGAV 's presented in Table 4 and the corrections in Table 7 the barrier for the reverse reaction is calculated as:

$$\begin{aligned} E_a(300 \text{ K}) &= E_{a,\text{ref}}(300 \text{ K}) + \Delta GAV_{E_a}^{\circ}(C_1-(S)(C)(H)) \\ &+ \Delta GAV_{E_a}^{\circ}(C_2-(C)(H)_2) \\ &+ \sigma_{\beta C-H} - \sigma_{\beta C-H} + p_{\alpha S} - \sigma_{\beta C-H} \\ &= 69.7 + 18.1 + 5.4 - 0.3 - 2.0 \\ &= 90.9 \text{ kJ mol}^{-1} \end{aligned}$$

As this barrier is higher than the one obtained for the forward reaction, it is concluded that hydrogen abstraction from ethanethiol by ethyl is exothermic. Substitution of $E_{a,\text{exo}} = 43.6 \text{ kJ mol}^{-1}$ in eqn (9) yields a transmission coefficient of 17.7 at 300 K, which corresponds within a factor of 2 to the calculated transmission coefficients of 25.1. The calculated E_a and $\log A$ amount to, respectively, 40.5 kJ mol^{-1} and 4.743. The GA model hence succeeds in estimating the rate coefficient for this reaction within a factor of 4.

The primary groups for reaction (97) are $C_1-(S)(C_d)(H)$ and $C_2-(S)(H)_2$. Their contributions to the activation energy amount to, respectively, +62.4 and -20.9 kJ mol^{-1} while their contributions to $\log \tilde{A}$ are, respectively, +0.066 and -0.564. The transition state is stabilized by the cross-interaction of one α -S atom and one $C=C$ ligand on the C_1 group with one α -S atom on the C_2 group. The number of single events for this reaction amounts to six corresponding to the three equivalents of hydrogen atoms in methanethiol and

chirality in the transition state. The activation energy and pre-exponential factor are hence calculated as:

$$\begin{aligned} E_a(300 \text{ K}) &= E_{a,\text{ref}}(300 \text{ K}) + \Delta GAV_{E_a}^{\circ}(C_1-(S)(C_d)(H)) \\ &+ \Delta GAV_{E_a}^{\circ}(C_2-(S)(H)_2) \\ &+ p_{\alpha S} - \pi_{C=C/\equiv} + p_{\alpha S} - p_{\alpha S} \\ &= 69.7 + 62.4 - 20.9 - 8.6 - 7.1 \\ &= 95.5 \text{ kJ mol}^{-1} \end{aligned}$$

$$\begin{aligned} \log A(300 \text{ K}) &= \log \tilde{A}_{\text{ref}}(300 \text{ K}) + \Delta GAV_{\log \tilde{A}}^{\circ}(C_1-(S)(C_d)(H)) \\ &+ \Delta GAV_{\log \tilde{A}}^{\circ}(C_2-(S)(H)_2) \\ &+ p_{\alpha S} - \pi_{C=C/\equiv} + p_{\alpha S} - p_{\alpha S} + \log(6) \\ &= 5.267 + 0.066 - 0.564 - 0.206 - 0.366 \\ &+ 0.778 = 4.975 \end{aligned}$$

The activation energy for the reverse reaction amounts to:

$$\begin{aligned} E_a(300 \text{ K}) &= E_{a,\text{ref}}(300 \text{ K}) + \Delta GAV_{E_a}^{\circ}(C_1-(S)(H)_2) \\ &+ \Delta GAV_{E_a}^{\circ}(C_2-(S)(C_d)(H)) \\ &+ p_{\alpha S} - \pi_{C=C/\equiv} + p_{\alpha S} - p_{\alpha S} \\ &= 69.7 + 20.1 - 34.8 - 8.6 - 7.1 \\ &= 39.4 \text{ kJ mol}^{-1} \end{aligned}$$

The reverse reaction is hence the exothermic reaction. Substitution of $E_{a,\text{exo}} = 39.4 \text{ kJ}$ in eqn (9) yields a transmission coefficient of 13.7 at 300 K. The calculated E_a and $\log A$ amount to, respectively, 98.5 kJ mol^{-1} and 5.165. The GA model succeeds in predicting the rate coefficient for this reaction within a factor of 1.3.

In Fig. 9 the accuracy of the GA model is compared to other models such as the one proposed by Blowers and Masel (eqn (10)) and the intersecting parabolas model. The IP model departs from the following relationship between the activation energy and reaction enthalpy:¹⁶

$$b = a(E_a - \Delta_r H^\circ)^{1/2} - E_a^{1/2} \quad (15)$$

with a and b two fitting parameters which relate to the force constants of the broken and formed hydrogen bonds and the hydrogen atom displacement during abstraction. The parity diagram shows the estimated activation energy as a function of the *ab initio* calculated activation energy for the 30 reactions presented in Table 6. It can be seen that the GA model outperforms the other models. As discussed above, the MAD between calculated and GA estimated activation energies amounts to approximately 3 kJ mol^{-1} . The MADs obtained with the Blowers and Masel model and the intersecting parabolas model amount to 8–9 kJ mol^{-1} . Large deviations are particularly obtained for reactions having resonance stabilized transition states. For example, both BEP type models overestimate the barriers for reaction (83) by 16 kJ mol^{-1} . Besides the less accurate estimates for the activation energy, these BEP-type models do not give access to pre-exponential factors.

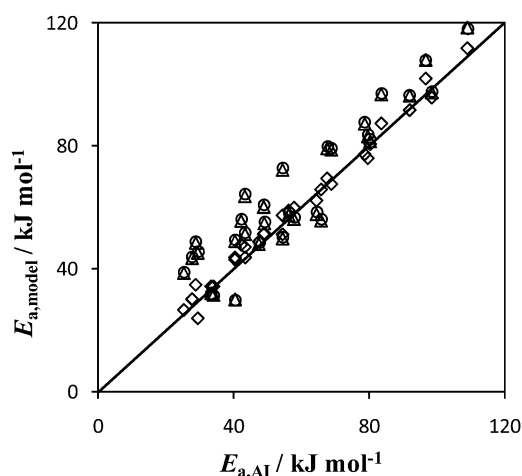


Fig. 9 Parity plot for the activation energies at 300 K calculated from CBS-QB3 rate coefficients and those obtained using the GA method presented in this work (\diamond), the Blowers and Masel model with $w_r + w_b = 800 \text{ kJ mol}^{-1}$ and $V_p = 1114 \text{ kJ mol}^{-1}$ and the intersecting parabolas model with $a = 1$ and $b = 16.3$ (\circ).

4 Conclusions

In this work a group additivity model is presented that allows us to estimate the Arrhenius parameters and rate coefficients for hydrogen abstraction reactions from carbon atoms in sulfides and thiocarbonyl compounds in the temperature range 300–1500 K. This work is an extension of previous work on modeling of the Arrhenius parameters for hydrogen abstractions involving hydrocarbons.²⁶

In order to derive and validate group additivity values, high-pressure limit rate coefficients were calculated using the CBS-QB3 method for more than 100 reactions. The rate coefficients were corrected for 1-D hindered rotation about the transition state bond. Tunneling contributions are modeled explicitly using an equation that correlates the tunneling contributions to the temperature and activation energy for the exothermic reaction.

From a training set containing 24 reactions, 46 ΔGAV° s were derived that can be used to model hydrogen abstraction reactions involving carbon atoms having a neighboring S or CS ligand. In general, the influence of secondary effects and NNIs on the Arrhenius parameters is limited. However, at low temperatures deviations of a few kJ mol^{-1} of the activation energy can lead to deviations of the rate coefficients up to one order of magnitude. In order to enhance the applicability of the presented GA method, one secondary contribution was introduced to account for hydrogen abstractions involving methyl disulfides. Resonance stabilization in the transition states was studied for a set of 32 reactions. As resonance stabilization can decrease the activation energy to 30 kJ mol^{-1} , 8 additional correction terms were introduced, which are able to describe this effect. The values for the 8 correction terms were obtained from a least squares regression. The introduction of corrections terms for transition state stabilization lowered the mean absolute deviation between *ab initio* and group additively estimated activation energies from 15.3 to 0.9 kJ mol^{-1} , hence significantly improving the accuracy of the GA model.

The developed group additivity scheme was validated using a test set containing an additional 30 reactions. The group additivity scheme succeeds in reproducing the rate coefficients on average within a factor of 2.4 at 300 K. The mean absolute deviations of the Arrhenius parameters amount to, respectively, 2.5 kJ mol^{-1} for E_a and 0.132 for $\log A$. The agreement between calculated and GA estimated data improves with increasing temperature. At 1000 K the rate coefficients are reproduced within an average factor of deviation of 1.4. It is hence concluded that the group additivity schemes, developed for elementary reactions between hydrocarbons, extrapolate successfully to compounds containing hetero-elements.

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