

# Nonlinear optics on silicon: towards integrated long-wavelength light sources

Bart Kuyken<sup>1,2</sup>, Xiaoping Liu<sup>3</sup>, Richard M. Osgood Jr.<sup>3</sup>, Roel Baets<sup>1,2</sup>, Günther Roelkens<sup>1,2</sup> and William M. J. Green<sup>4</sup>

<sup>1</sup>Photonics Research Group, Department of Information Technology, Ghent University – imec, Ghent B-9000, Belgium.

<sup>2</sup>Center for Nano- and Biophotonics (NB-Photonics), Ghent University, Ghent, Belgium.

<sup>3</sup>Microelectronics Sciences Laboratories, Columbia University, New York, NY 10027, USA.

<sup>4</sup>IBM Thomas J. Watson Research Center, 1101 Kitchawan Road, Yorktown Heights, NY 10598, USA.

<sup>‡</sup>Current address: OFS Labs, 19 Schoolhouse Road, Somerset, NJ 08873, USA.

[\\*Bart.Kuyken@intec.ugent.be](mailto:Bart.Kuyken@intec.ugent.be)

*The short-wave and mid-infrared wavelength region, often referred to as the molecular fingerprint region, has an enormous potential for spectroscopic applications. Silicon's transparency from 1.1  $\mu\text{m}$  up to 7  $\mu\text{m}$  could potentially be used to make a platform for a whole new class of compact devices for gas trace sensing, environmental monitoring etc. By using the large effective  $\chi^{(3)}$  nonlinearity, light generation by nonlinear mixing in this new wavelength region is demonstrated. We present both a silicon broadband supercontinuum spanning from the telecom window into the mid-infrared as well as a silicon-based widely tuneable optical parametric oscillator.*

## Introduction

Many molecules have distinct absorption features in the short-wave and mid-infrared wavelength range. These absorption features resulting from the molecular vibrational states of the molecules are specific and act as a fingerprint for these molecules. For now, silicon has been used primarily to construct planar lightwave circuits in the telecom wavelength range. However due to its transparency up to 7  $\mu\text{m}$ , silicon waveguides can be used as a platform to guide this mid-infrared light. This would enable to integrate complex functions on a compact silicon chip and enable a whole new set of applications for spectroscopic sensing. However, to probe these absorption features tuneable sources or broadband sources need be integrated on these chips. The integration of such sources with a silicon chip is not straightforward. However, the high intrinsic nonlinear  $\chi^{(3)}$  nonlinearity can be used to generate new wavelengths. Indeed by combining the high intensities obtained by the high confinement in silicon wire waveguides with the enormous nonlinearity of the material, low-threshold nonlinear effects can be obtained. The nonlinear interactions can be further enhanced by using dispersion engineered wire waveguides. In these dispersion engineered waveguides phase matching between the optical waves in the nonlinear optical process is satisfied. Here we demonstrate the use of nonlinear optics in silicon waveguides to generate a broadband supercontinuum light source as well as a widely tunable silicon based optical parametric oscillator.

## Nonlinear optics in silicon nanowire waveguides in the mid-infrared

At telecommunication wavelengths silicon suffers from two-photon absorption, which limits the efficiency of nonlinear processes in silicon wire waveguides. Given the 1.12 eV bandgap of silicon, this two-photon absorption process disappears at wavelengths

beyond 2.2  $\mu\text{m}$ . In this work, 900 nm wide and 220 nm thick silicon strip nanowire waveguides with an air top cladding are pumped close to 2.2  $\mu\text{m}$  to achieve efficient nonlinear interactions. The waveguide cross-section is shown in the inset of Figure 1. Using a cut back technique, the TE mode waveguide propagation loss is found to be approximately 2.5 dB/cm across the 2050-2450 nm wavelength range. In Figure 1 the higher order dispersion terms of the waveguide are shown. These waveguides are designed to have a negative second order dispersion and positive fourth order dispersion to achieve phase matching in the four wave mixing process.

Indeed, efficient phase matching in the four wave mixing occurs when [4]

$$\Delta k = \Delta k_{lin} + \Delta k_{nonlin} = 2k_{pump} - k_s - k_i - 2\gamma P = 0 \quad (1)$$

in which  $k_{pump}$ ,  $k_s$  and  $k_i$  are the linear propagation constants of the pump, signal and idler wave respectively. The term  $2\gamma P$ , in which  $\gamma$  is the effective nonlinear parameter of the waveguide and  $P$  is the peak power of the pump pulse, accounts for the self-phase and cross-phase modulation of the interacting waves. Using a Taylor expansion of the waveguide dispersion relation around  $\omega_{pump}$  and taking into account the conservation of energy in the four wave mixing process, this results in a phase matching condition

$$-\beta_2 \Delta\omega^2 - \frac{1}{12} \beta_4 \Delta\omega^4 - 2\gamma P = 0 \quad (2)$$

in which  $\Delta\omega$  is the frequency detuning between pump and signal (and also between pump and idler),  $\beta_2$  the second order and  $\beta_4$  the fourth order dispersion. By averaging the nonlinear susceptibility of bulk Si over the electric field of the fundamental TE polarized waveguide mode, the real part of the nonlinearity parameter is estimated to be  $\gamma = 150 (\text{W}\cdot\text{m})^{-1}$  in the silicon nanowire being used.

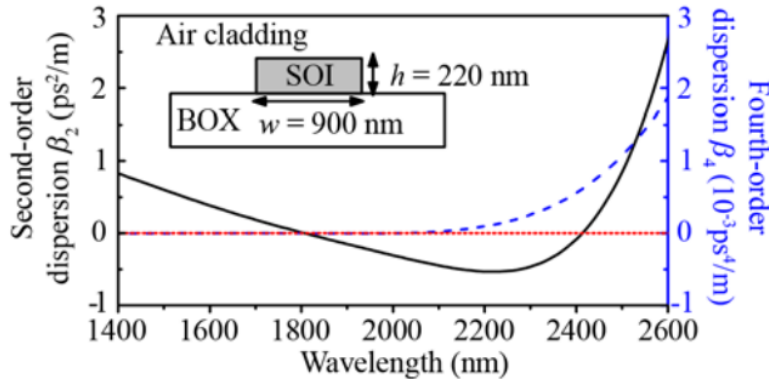


Fig. 1: Dispersion of the silicon photonic wire, simulated using a commercial finite element solver (RSoft FemSIM). The inset shows a cross section of the silicon wire waveguide.

## Supercontinuum generation

The supercontinuum generation experiments are conducted using a picosecond pulse train at a center wavelength of 2120 nm as the pump (Coherent Mira-OPO, FWHM = 2 ps, repetition rate = 76 MHz). Figure 2 illustrates the evolution of the waveguide output

spectrum as the input coupled peak pump power is gradually increased from 3.1 W (green trace) to 12.7 W (black trace). These spectra reveal that a number of different nonlinear processes ultimately combine to produce the broadband supercontinuum. At an input power of 3.1 W, a series of sidebands are generated in the vicinity of the pump at 2120 nm. Closest to the pump, two broad sidebands (labeled as MI(1)) are generated near wavelengths of 1990 nm and 2250 nm. Further away from the pump, a pair of narrowband peaks (labeled as MI(2)) appear at wavelengths of 1870 nm and approximately 2510 nm. Both the broad and narrow sideband pairs originate from modulation instability, i.e. the amplification of background noise at wavelengths for which the phase matching condition in Eq. (2) is satisfied. At a pump power of 7.9 W, several new spectral components are observed, peaked near 1700 nm and 1600 nm respectively. The term at 1700 nm is generated through cascaded four wave mixing (FWM), where the original MI(2) peak at 1890 nm serves as the degenerate pump and the input pulse at 2120 nm acts as the signal. At maximum power the supercontinuum was inspected with an FTIR, this revealed that the supercontinuum [1] spans from 1.53  $\mu\text{m}$  up to 2.55  $\mu\text{m}$ .

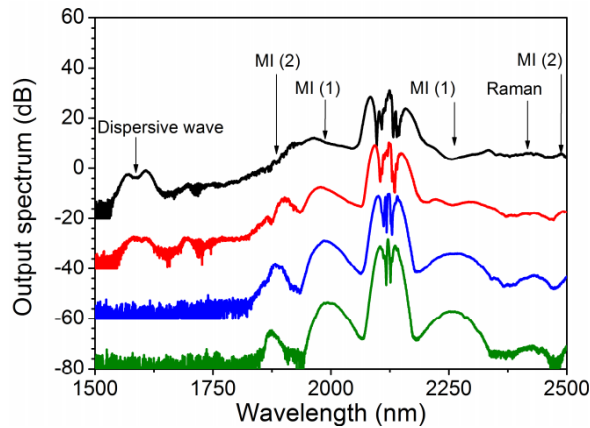


Fig. 2. Measured output spectrum for increasing values of coupled input peak power: 3.1 W (green), 4.3 W (blue), 7.9 W (red) and 12.7 W (black).

## Widely Tuneable Silicon Mid-Infrared Optical Parametric Oscillator

The broadband on-chip parametric gain demonstrated by the modulation instability sidebands [3] can be combined with optical feedback in order to achieve optical parametric oscillation over a broad wavelength range [4]. The synchronously pumped fiber optic loop configuration employed to construct the mid-IR OPO is illustrated in Figure 3. The mid-IR pumped photonic wire waveguide serves as the core gain element within the fiber loop. The cavity of the OPO is formed by a loop of standard single mode fiber, along with a variable free-space delay element, which facilitates temporal synchronization of the re-circulating amplified pulses with successive pump pulses. Due to the absence of special purpose wavelength multiplexers/demultiplexers within the mid-IR wavelength range of operation, a 90/10 coupler is used to combine the cavity pulses with successive pump pulses after each round-trip. Polarization controllers are used to align both the pump and cavity pulses to the quasi-TE-mode at the input of the photonic wire. The output of the silicon-fiber OPO is monitored by a mid-infrared

optical spectrum analyzer (OSA). The high dispersion of the single mode fiber feedback loop at wavelengths beyond 2000 nm allows for wavelength selective feedback. By adjusting the round-trip time using the free-space variable delay line, a particular temporal “slice” of the circulating amplified pulse is selectively synchronized with the pump pulse train. The dispersion in the fiber loop likewise ensures that only one particular wavelength within the silicon parametric gain spectrum can be synchronized with the pump pulses, and thus enables broadband tuneability of the oscillator wavelength via delay tuning. The energy of the output pulses as a function of the output wavelength for pump pulses with an energy of 48 pJ at 2175 nm is shown in Figure 4. It was demonstrated that the OPO is tuneable over a 75 nm-wide band centered around the gain peak at 2075 nm.

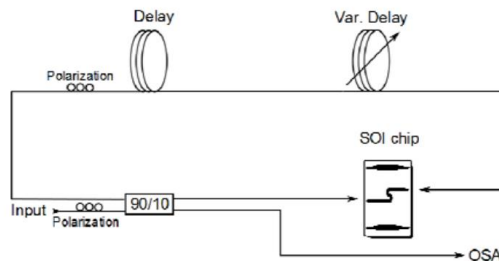


Fig 3. Configuration of the synchronously-pumped silicon-fiber OPO.

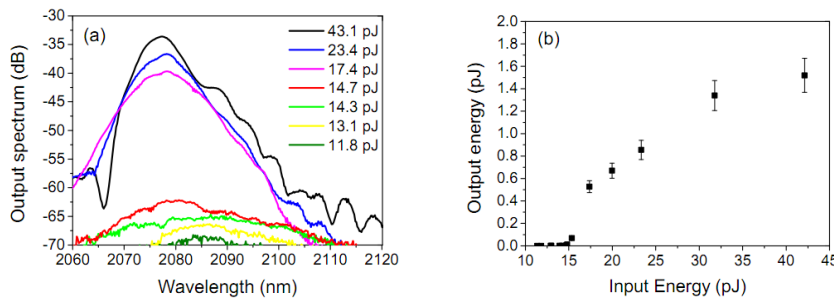


Fig 4 Output spectrum of the silicon-fiber OPO for different pump pulse energies, when the fiber cavity is tuned to an output wavelength of 2075 nm. (b) On-chip output pulse energy versus the coupled input pulse energy.

## Conclusion

Nonlinear optics has a huge potential for integrating tuneable or broadband sources in the mid-infrared on a silicon chip. A broadband supercontinuum spanning from the telecom wavelength range into the mid-infrared was demonstrated as well as a tuneable silicon-based optical parametric oscillator.

- [1] B. Kuyken, et al., “Mid-infrared to telecom-band supercontinuum generation in highly nonlinear silicon-on-insulator wire waveguides” ,*Optics Express*, **19**, p.20172-20181, 2011
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- [4] B. Kuyken, et al., “Widely Tuneable Silicon Mid-Infrared Optical Parametric Oscillator”, *Group IV Photonics*, 2011.