



ELSEVIER

Contents lists available at SciVerse ScienceDirect

## Journal of Magnetism and Magnetic Materials

journal homepage: [www.elsevier.com/locate/jmmm](http://www.elsevier.com/locate/jmmm)Nano-ilmenite FeTiO<sub>3</sub>: Synthesis and characterizationA.T. Raghavender<sup>a,\*</sup>, Nguyen Hoa Hong<sup>a</sup>, Kyu Joon Lee<sup>b</sup>, Myung-Hwa Jung<sup>b</sup>, Z. Skoko<sup>c</sup>, M. Vasilevskiy<sup>d</sup>, M.F. Cerqueira<sup>d</sup>, A.P. Samantilleke<sup>d</sup><sup>a</sup> Nanomagnetism Laboratory, Department of Physics and Astronomy, Seoul National University, Seoul 151–747, South Korea<sup>b</sup> Department of Physics, Sogang University, 1 Shinsu-dong, Mapo-gu, Seoul 121–742, South Korea<sup>c</sup> Department of Physics, Faculty of Science, University of Zagreb, Bijenicka C. 32, HR-10000, Zagreb, Croatia<sup>d</sup> Centro de Física, Universidade do Minho, Braga 4710-057, Portugal

## ARTICLE INFO

## Article history:

Received 29 March 2012

Received in revised form

10 September 2012

Available online 23 November 2012

## Keywords:

Nano-ilmenite

Structural

Optical

Magnetic properties

## ABSTRACT

In general, ilmenite FeTiO<sub>3</sub> is synthesized by solid-state reaction at very high pressure and high temperature. Synthesis of FeTiO<sub>3</sub> is not an easy task as the Fe<sup>2+</sup> ions are not stable. Therefore, it is really challenging to prepare this material. In this work nano-ilmenite FeTiO<sub>3</sub> was synthesized by the sol-gel method. Structural, optical and magnetic characterizations were performed. The bandgap of FeTiO<sub>3</sub> was determined to be 2.8 eV showing FeTiO<sub>3</sub> as suitable wide bandgap material for technological applications. The FeTiO<sub>3</sub> nanoparticles exhibit weak ferromagnetic properties at and below room temperature. The Néel temperature was observed to be around 52 K.

© 2012 Elsevier B.V. All rights reserved.

## 1. Introduction

Ilmenite FeTiO<sub>3</sub> is antiferromagnetic (AFM) and insulating but the intermediate compositions can be semiconducting and insulating [1–6]. FeTiO<sub>3</sub> is one of the most common mineral in the Earth crust [7]. FeTiO<sub>3</sub> is an interesting wide bandgap (2.58–2.9 eV) antiferromagnetic semiconductor [8,9] with potential applications in spintronics, optoelectronics, high temperature integrated circuits, chemical catalysts and photocatalysts etc. [10–13]. Most of the efforts have been devoted for designing ferromagnetic semiconductors operating at room temperature on homogeneous doping of semiconductors with magnetic impurities [14–16]. Furthermore, since the conductivity type (*p* or *n*) can be tuned by changing the doping density of Ti into hematite promises additional applications in electronic devices such as low-voltage varistors [17]. Ilmenite has rhombohedral structure with space group  $R\bar{3}$  and the Néel temperature is around 55 K [3]. In FeTiO<sub>3</sub>, Fe and Ti layers are arranged alternate, this reduces the symmetry to  $R\bar{3}$  with AFM coupling between the Fe layers with  $T_N=56\text{--}59\text{ K}$  [18,19]. Fujii et al. [20] have reported that FeTiO<sub>3</sub> thin films with  $R\bar{3}$  symmetry were ferrimagnetic at lower temperature, while  $R\bar{3}C$  symmetry was antiferromagnetic. Assuming that Ti<sup>4+</sup> is in B site and Fe<sup>3+</sup> is equally distributed between A

and B sites, the chemical composition can be written as

$$(Fe_x^{2+}, Fe_{1-x}^{3+})_{A\text{-site}}(Ti_x^{4+}, Fe_{1-x}^{3+})_{B\text{-site}}O_3 \quad (1)$$

It was reported that FeTiO<sub>3</sub> could be synthesized by using solid-state reaction above 1200 °C under vacuum [21–23]. Popular techniques that have been employed to synthesize FeTiO<sub>3</sub> in the recent year include ball milling [24], chemical reduction techniques [11], co-precipitation [25] etc. The structural, optical, electrical and magnetic properties of FeTiO<sub>3</sub> are strongly influenced by the preparation technique [26]. It is also observed that the synthesis of FeTiO<sub>3</sub> is made complicated by the instability of Fe<sup>2+</sup> in FeTiO<sub>3</sub>. In this work, we have successfully synthesized nano-ilmenite FeTiO<sub>3</sub>. In addition the magnetic, structural and optical properties of nano-ilmenite were characterized.

## 2. Experiment

Nano-ilmenite FeTiO<sub>3</sub> has been synthesized by using the sol-gel auto-ignition method [27]. The A.R grade citric acid (C<sub>6</sub>H<sub>8</sub>O<sub>7</sub>·H<sub>2</sub>O), ferric (III) nitrate nonahydrate (Fe(NO<sub>3</sub>)<sub>3</sub>·9H<sub>2</sub>O), titanium dioxide (TiO<sub>2</sub>) (≥ 99.99%) from Sigma-Aldrich were used as precursor materials. The molar ratio of metal nitrates to citric acid was taken as 1:3. Titanium nitrate was made by converting, titanium dioxide by dissolving in nitric acid. Ferric nitrate and citric acid were dissolved in minimum amount of de-ionized water to obtain a clear solution. The resulting two solutions were mixed together and stirred for about 30 min to get a uniform mixture solution. The pH was adjusted to 7 using ammonia; the

\* Corresponding author. Tel.: +82 2 880 6603; fax: +82 2 884 3002.

E-mail addresses: atraghavender@snu.ac.kr, raghavi9@gmail.com (A.T. Raghavender).

temperature of the solution was increased to 200 °C. After 4–5 h, the water molecules evaporated from the chemical solution. During evaporation, the solution became viscous and finally formed a very viscous brown/black gel. When all the remaining water was released from the mixture, the sticky mass began to bubble. During this time the temperature of the beaker was increased to 380 °C in order to increase the reaction rate. After several minutes, the gel automatically ignited and burnt with glowing flints. The decomposition reaction would not stop before the whole citrate complex was consumed. The auto-ignition was completed within a minute yielding the brown-colored ashes. The as-prepared powder sample was annealed at 500 °C for 10 h to get the final product. For transmission measurements, a thin layer of FeTiO<sub>3</sub> was coated on fluorine doped tin oxide (FTO) glass electrode using the doctor blade method [28]. For Raman measurements, FeTiO<sub>3</sub> powder was then pressed into pellets of diameter 10 mm and thickness 4 mm at a pressure of 40 MPa. The obtained pellets were sintered in air at 500 °C for 10 h.

The microstructure of the FeTiO<sub>3</sub> was studied by Raman scattering. Micro-Raman spectra have been measured in a frequency range of 50–750 cm<sup>-1</sup>, using the 488 nm excitation line of an Ar<sup>+</sup> laser, in the back scattering geometry, on a Jobin-Yvon T64000 spectrometer equipped with a liquid nitrogen cooled CCD detector. In order to avoid sample heating the incident sample laser power has been kept at 4 mW. The thicknesses of the films were measured using VEECO Dektak 150 surface profiler. The grain size and morphology were verified by using Hitachi SU-70, high resolution scanning electron microscope (HR-SEM) with an energy dispersive X-ray spectrometer (EDS) installed. EDS was performed to further confirm the composition of the prepared samples. The magnetization was measured by a commercial MPMS (Quantum Design) superconducting quantum interference device (SQUID) magnetometer in applied fields up to 70 kOe and in the temperature range of 5–300 K.

### 3. Results and discussions

It is well known that the positions and widths of Raman bands of materials are related to their vibrational and structural properties. According to literature, the Raman peaks of FeTiO<sub>3</sub>, in the studied range, are at 220, 250, 300, 410, 509, 610 and 650 cm<sup>-1</sup>. For the standard FeTiO<sub>3</sub>, the most intense mode should be 410 cm<sup>-1</sup>. Fig. 1 shows the Raman spectra of FeTiO<sub>3</sub> sample and Table 1 shows the result of the Raman fitting using Lorentzian functions for each band. Thus, we could get the wavenumber and the full-width at half-maximum (FWHM) of each mode. This sample shows the usual bands related with vibration modes of bonds related with FeTiO<sub>3</sub> (mentioned in the literature between 200 and 800 cm<sup>-1</sup>). According to Olivares et al. [29], the produced material is mainly based on iron oxide compounds as  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> (hematite: 226, 296, 410, 617 cm<sup>-1</sup>) and  $\gamma$ -Fe<sub>2</sub>O<sub>3</sub> (maghemite: 667 cm<sup>-1</sup>), as can be seen by the presence of the modes slightly shifted at 224, 336, 443, 612 (hematite) and 662 cm<sup>-1</sup> (maghemite) in the Raman spectra shown in Fig. 1. The presence of iron oxides, such as magnetite indicates a reducing atmosphere, and the observed iron oxides reddish, such as hematite, means that an oxidizing environment was used [30]. The Raman analysis observed for our nano-ilmenite samples are very well in agreement with Sharma et al. [31].

The morphology of FeTiO<sub>3</sub> was determined by HR-SEM. Fig. 2 represents the HR-SEM image of FeTiO<sub>3</sub> showing a uniform and a narrow distribution of the particles. EDS was performed to further confirm the compositions of the prepared samples. Fig. 3 shows that the prepared FeTiO<sub>3</sub> sample consists of Fe, Ti and O. The weight and atomic fractions of all elementary constituents in the synthesized FeTiO<sub>3</sub> nanoparticles determined by EDS are presented in Table 2.

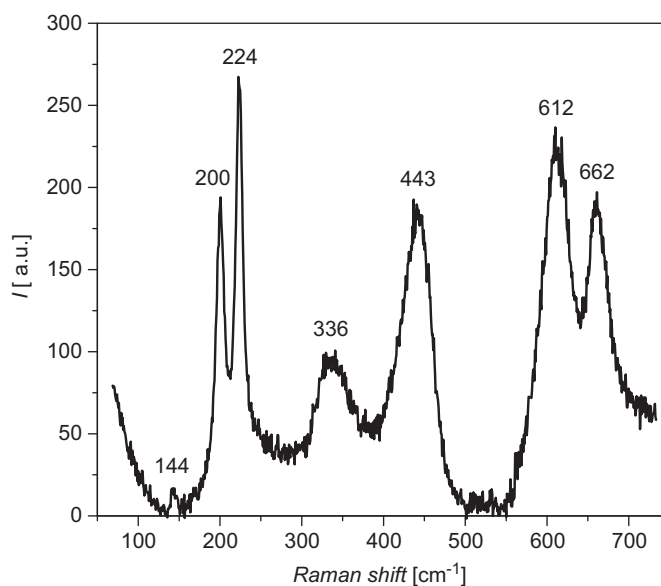


Fig. 1. Room temperature Raman spectra of nanocrystalline FeTiO<sub>3</sub>.

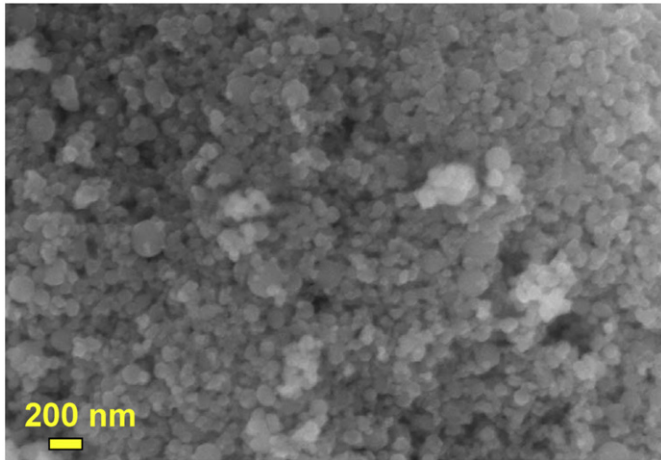
Fig. 4 shows the transmission spectrum of a typical FeTiO<sub>3</sub> nanoparticles coated FTO glass electrode with a film thickness of 1200 nm. The transmission is below 50% for most of the spectrum mainly due to the large thickness of the sample. However, the transmission is particularly weak in the visible part of the solar spectrum. This may be ascribed to the presence of high molar ratio of Fe atoms in the lattice of TiO<sub>2</sub>, resulting in increased visible light absorption. The optical bandgap  $E_g$  is determined from the extrapolated Tauc plot (Fig. 5) made by using the transmission spectrum. The estimated bandgap for the nanocrystalline FeTiO<sub>3</sub> from the Tauc plot (Fig. 5) was approximately 2.80 eV, which is comparable with reported bandgap values for antiferromagnetic FeTiO<sub>3</sub> (2.58–2.9 eV) elsewhere [32–34].

The M–H loops for FeTiO<sub>3</sub> were measured at different temperatures as of 10 K, 100 K and 300 K with an applied field between –70 and 70 kOe are shown in Fig. 6. The inset of Fig. 6 shows the expanded M–H at lower scale showing a small non-linear behavior. At 10 K, the sample shows large coercivity of 818 Oe, and at 100 K and 300 K with 16 Oe and 12 Oe. It is also observed that M–H loops do not saturate even up to 70 kOe magnetic field due to paramagnetic or antiferromagnetic nature of the material [34,35]. The remanence magnetization measured for FeTiO<sub>3</sub> at different temperatures was considerably small. This kind of magnetic behavior is attributed due to the presence of surface defects/ unsatisfied surface spins that develop in the crystal during the crystallization process [36]. Even though the  $M_r$  and  $H_c$  observed in our samples are very small they still show a decreasing trend with increasing temperature. This might be due to the flux trapped by the defects in nano-ilmenite FeTiO<sub>3</sub> that gets excluded as the temperature increases [37].

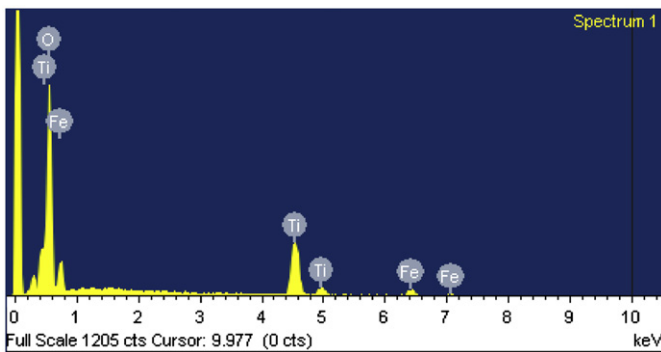
Fig. 7 shows the zero-field-cooled (ZFC) and field-cooled (FC) magnetization curves measured with applied fields 100 Oe, 500 Oe and 1000 Oe. As the temperature increases the magnetic moment in the FC curve decreases. However, as the temperature begins to increase from 2 K, the magnetic moment in the ZFC curves also decrease up to 45 K and beyond this point, ZFC curves reach a maximum position at 52 K and thereafter FC and ZFC merge together and then steadily decrease up to 300 K. The peak position at 52 K is assigned to be the Néel temperature of FeTiO<sub>3</sub>. The FC–ZFC diverges below the  $T_N \approx 52$  K which agrees with the magnetic transition temperature of FeTiO<sub>3</sub>. One can clearly see that all the

**Table 1**  
Wavenumber and FWHM of the FeTiO<sub>3</sub> Raman modes.

Sample	$\omega(\text{cm}^{-1})/\Gamma(\text{cm}^{-1})$	$\omega(\text{cm}^{-1})/\Gamma(\text{cm}^{-1})$	$\omega(\text{cm}^{-1})/\Gamma(\text{cm}^{-1})$	$\omega(\text{cm}^{-1})/\Gamma(\text{cm}^{-1})$	$\omega(\text{cm}^{-1})/\Gamma(\text{cm}^{-1})$	$\omega(\text{cm}^{-1})/\Gamma(\text{cm}^{-1})$	$\omega(\text{cm}^{-1})/\Gamma(\text{cm}^{-1})$
FeTiO <sub>3</sub>	144/12	200/12	224/11	336/60	443/50	612/50	662/40



**Fig. 2.** SEM image of nanocrystalline FeTiO<sub>3</sub> powder.

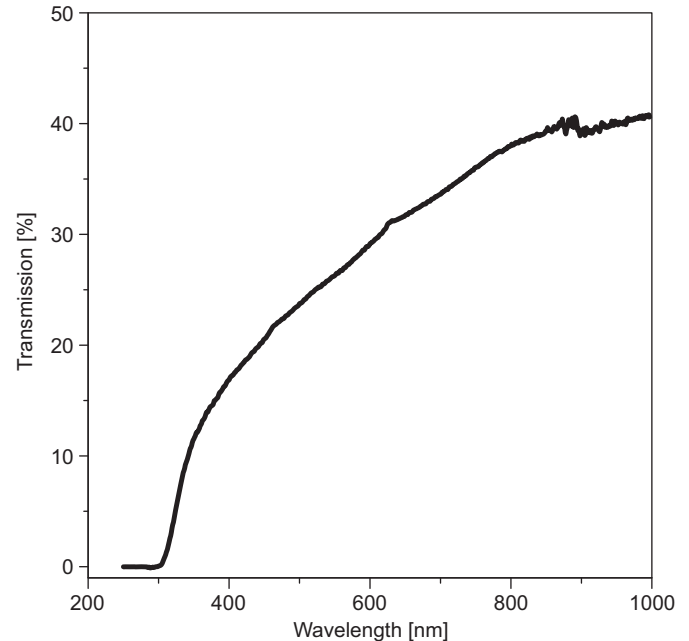


**Fig. 3.** EDS of nanocrystalline FeTiO<sub>3</sub> powder.

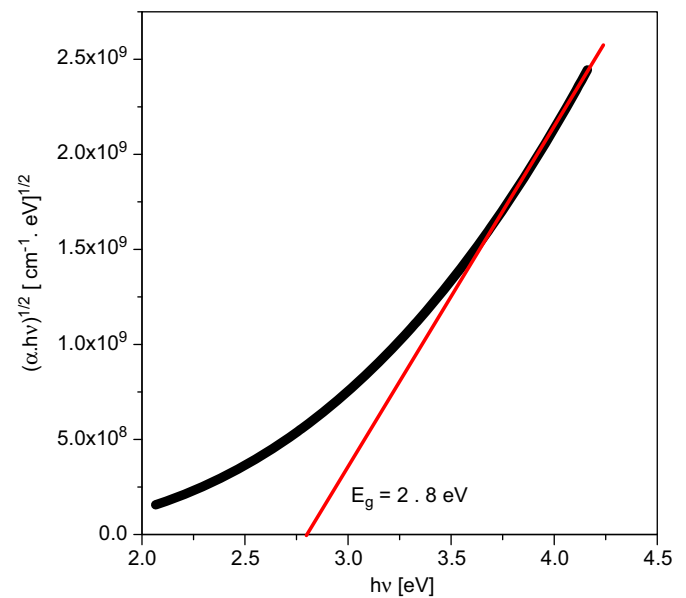
**Table 2**  
The elementary constituents of FeTiO<sub>3</sub> powder measured by EDS.

Elements	O	Ti	Fe	Total
Weight percentage (%)	36.45	28.24	35.31	100
Atomic percentage (%)	62.08	16.84	18.06	100

ZFC and FC curves show similar tendencies. The peak in the ZFC curves corresponds to the Néel temperature [38]. Both FC and ZFC curves resemble those reported in Refs. [18,39]. Above  $T_N=52$  K, Ilmenite is paramagnetic, while below  $T_N$  the Fe<sup>2+</sup> layers couple antiferromagnetically and are separated magnetically inert Ti<sup>4+</sup> layers [40]. The peak positions of ZFC and FC (Fig. 7) which are almost independent of the applied fields further confirm the AFM phase transition. Such kind of behavior was observed previously for CoO<sub>3</sub> nanoparticles [41]. These unusual magnetic properties were observed from the temperature dependence of the resistivity, magneto-elastic properties and the neutron diffraction properties [42,43]. Here we observe that the antiferromagnetic transition  $\approx 52$  K does not depend on the external applied fields, probably due to relatively strong antiferromagnetic forces [36]. Earlier, ZFC–FC that were done at much higher applied fields showed only Néel temperature with deviations in the ZFC–FC below 90 K [44,45].



**Fig. 4.** Transmission spectra of FeTiO<sub>3</sub> deposited on FTO glass substrate.



**Fig. 5.** Tauc plot of FeTiO<sub>3</sub>.

Hence, it is observed that the structural, optical and magnetic properties of FeTiO<sub>3</sub> were found to be influenced by the nanoscaled structures.

#### 4. Conclusions

Nano-ilmenite FeTiO<sub>3</sub> have been successfully synthesized using sol-gel method. The bandgap of FeTiO<sub>3</sub> was determined to be 2.8 eV.

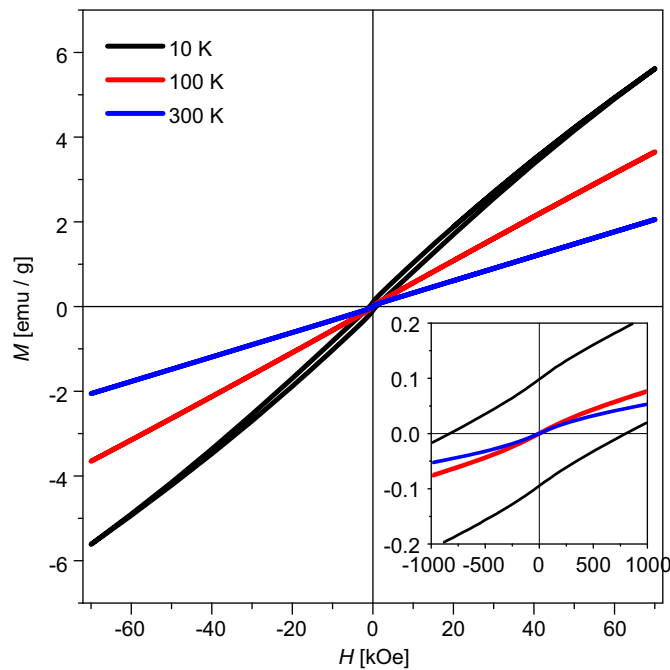


Fig. 6. Magnetization vs. magnetic field curve for nanocrystalline FeTiO<sub>3</sub> powder. The inset shows the plot for magnetic field in expanded scale.

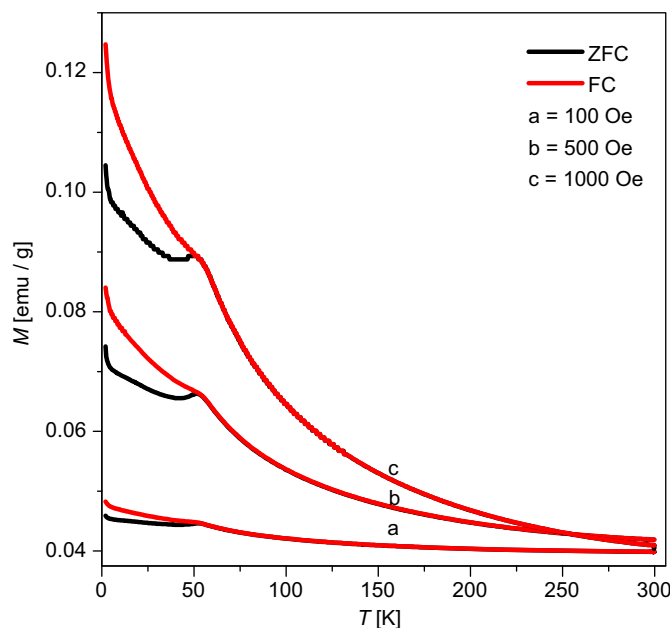


Fig. 7. Temperature dependent magnetization (ZFC-FC).

The synthesized FeTiO<sub>3</sub> nanoparticles showed weak ferromagnetic properties at and below 300 K. The M–H curves measured for all the temperature range showed the weak ferromagnetism linked with the onset of the antiferromagnetic state. The antiferromagnetic (AFM) transition  $T_N \approx 52$  K was determined from the magnetization versus temperature data. The  $T_N$  values are highly independent to the applied field further confirming the AFM nature of FeTiO<sub>3</sub>.

## Acknowledgments

A.T.R would like to thank, BK21 Frontier Physics Research Division, Department of Physics and Astronomy, Seoul National University,

Seoul, South Korea, for the Postdoctoral fellowship and EU-FP program of NRF (National Research Foundation), South Korea for financial support. M.F.C and A.P.S. would like to thank the FCT (Fundação para a Ciência e Tecnologia) for funding and M.H.J. would like to thank KRF grant (2012-0004082) through the pluriannual contract with CFUM, Ciencia 2007 Program and the European Commission through FP7-PEOPLE-2010-IRSES-NanoCIS (269279).

## References

- [1] C. Frandsen, B.P. Burton, H.K. Rasmussen, S.A. McEnroe, S. Morup, *Physical Review B* 8 (2010) 224423.
- [2] T. Nagata, *Nature* 172 (1953) 850.
- [3] Y. Ishikawa, S. Akimoto, *Journal of the Physical Society of Japan* 12 (1957) 1083.
- [4] R.M. Bozorth, D.E. Walsh, A.J. Williams, *Physical Review* 108 (1957) 157.
- [5] P. Robinson, R.J. Harrison, S.A. McEnroe, R.B. Hargraves, *Nature* 418 (2002) 517.
- [6] N.H. Butler, A. Bandopadhyay, R. Srinivasan, *Journal of Applied Physics* 93 (2003) 7882.
- [7] S.E. Haggerty, V. Sautter, *Science* 248 (1990) 993.
- [8] X. Tang, K. Hu, *Journal of Materials Science* 41 (2006) 8025.
- [9] F. Ye, A. Ohmori, *Surface and Coatings Technology* 160 (2002) 62.
- [10] T. Fujii, M. Kayano, Y. Takada, M. Nakanishi, J. Takada, *Solid State Ionics* 172 (2004) 289.
- [11] Z. Dai, P. Zhu, S. Yamamoto, A. Miyashita, K. Narum, H. Naramoto, *Thin Solid Films* 339 (1) (1999) 11.
- [12] F. Zhou, S. Kotru, R.K. Pandey, *Materials Letters* 57 (2003) 2104.
- [13] T. Dietl, H. Ohno, *MRS Bulletin* 28 (2003) 714.
- [14] H. Yamada, Y. Ogawa, Y. Ishii, H. Sato, M. Kawasaki, H. Akoh, Y. Tokura, *Science* 305 (2004) 646.
- [15] S. Kuroda, N. Nishizawa, K. Takita, M. Mitome, Y. Bando, K. Osuch, T. Dietl, *Nature Materials* 6 (2007) 440.
- [16] A.H. MacDonald, P. Schiffer, N. Samarth, *Nature Materials* 4 (2005) 195.
- [17] F. Zhou, S. Kotru, R.K. Pandey, *Thin Solid Films* 408 (2002) 33.
- [18] R. Pentcheva, H. Sadat Nabi, *Physical Review B* 77 (2008) 172405.
- [19] P.F. McDonald, A. Parasiris, R.K. Pandey, B.L. Gries, W.P. Kirk, *Journal of Applied Physics* 69 (1991) 1104.
- [20] T. Fujii, M. Kayano, Y. Takada, M. Nakanishi, J. Takada, *Journal of Magnetism and Magnetic Materials* 272–276 (2004) 2010.
- [21] C.H. Shomate, *Journal of the American Chemical Society* 68 (1946) 964.
- [22] B.F. Naylor, O.A. Cook, *Journal of the American Chemical Society* 68 (1946) 1003.
- [23] R.W. Grant, R.M. Housley, S. Geller, *Physical Review B* 5 (1972) 1700.
- [24] S. Ohara, K. Sato, Z. Tan, H. Shimoda, M. Ueda, T. Fukui, *Journal of Alloys and Compounds* 504 (2010) L17–L19.
- [25] S. Ya, S. Ge, W. Qiao, Y. Zuo, *Journal of Magnetism and Magnetic Materials* 322 (2010) 824.
- [26] N.C. Wilson, J. Muscat, D. Mkhonto, P.E. Ngoepe, M.M. Harrison, *Physical Review B* 71 (2005) 075202.
- [27] A.T. Raghavender, D. Pajić, K. Zadro, T. Mileković, P. Venkateshwar Rao, K.M. Jadhav, D. Ravinder, *Journal of Magnetism and Magnetic Materials* 316 (2007) 1.
- [28] A.T. Raghavender, A.P. Samanthilleke, Pedro Sa, B.G. Almeida, M.I. Vasilevskiy, N.H. Hong, *Materials Letters* 69 (2012) 59.
- [29] M. Olivares, M.C. Zuluaga, L.A. Ortega, X. Murelaga, A. Alonso-Olazarbal, M. Urteaga, L. Amundaray, I. Alonso-Martin, N. Etxebarria, *Journal of Raman Spectroscopy* 41 (2010) 1543.
- [30] L. Nodari, E. Marcuz, L. Maritan, C. Mazzoli, U. Russo, *Journal of the European Ceramic Society* 27 (2007) 4665.
- [31] Y.K. Sharma, M. Kharkwal, S. Uma, R. Nagarajan, *Polyhedron* 28 (2009) 579.
- [32] S.S. Sunkara, R.K. Pandey, *Ceramic Transactions* 60 (1995) 83.
- [33] Z. Dai, H. Naramoto, K. Narumi, S. Yamamoto, A. Miyashita, *Journal of Applied Physics* 85 (1999) 7433.
- [34] Y.H. Chen, *Journal of Non-Crystalline Solids* 357 (2011) 136.
- [35] T. Varga, A. Kumar, E. Vlahos, S. Denev, M. Park, T. Sanehira, Y. Wang, C.J. Fennie, S.K. Streiffer, X. Ke, P. Schiffer, V. Gopalan, J.F. Mitchell, *Physical Review B* 103 (2009) 047601.
- [36] R. Das, A. Jaiswal, S. Adyanthaya, P. Poddar, *Journal of Physical Chemistry C* 114 (2010) 12104.
- [37] X. Xu, A. Umezawa, G.W. Crabtree, *Physical Review B* 46 (1992) 11928.
- [38] V. Yu Galkin, W.A. Ortiz, E.J. Fawcett, *Journal of Physics: Condensed Matter* 9 (1997) L577.
- [39] T. Katsufuji, S. Mori, M. Masaki, Y. Maritomo, N. Yamamoto, H. Takagi, *Physical Review B* 64 (2001) 419.
- [40] H. Sadat Nabi, R.J. Harrison, R. Pentcheva, *Physical Review B* 81 (2010) 214432.
- [41] Y. Endoh, Y. Ishikawa, H. Ohno, *Journal of the Physical Society of Japan* 24 (1968) 263.
- [42] H.L. Alberts, J.A. Lourens, *Journal of Physics F: Metal Physics* 13 (1983) 873.
- [43] L. He, C. Chen, N. Wang, W. Zhou, L. Guo, *Journal of Applied Physics* 102 (2007) 103911.
- [44] K. Yoshii, H. Abe, *Journal of Solid State Chemistry* 165 (2002) 131.
- [45] K. Uusi-Eko, J. Malm, N. Imamura, H. Yamamuchi, M. Karppinen, *Materials Chemistry and Physics* 112 (2008) 1029.