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Research Paper

Approximation solution for steel concrete beam accounting high-order shear deformation using trigonometric-series

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ARTICLE INFO	ABSTRACT			
Article history :	Steel concrete beams have a reasonable structure in terms of using material and high load			
Received : 15 November 2022	carrying capacity. This paper deals with an approximate solution based on a trigonometric series for the static of steel concrete beams. The displacement field is based on the higher-			
Revised : 16 December 2022	order theory using Reddy's hypothesis. The governing equations are derived from			
Accepted : 18 December 2022	variation principles. An approximate solution based on the representation of displacement fields by trigonometric series is developed to solve the static problem of steel concrete			
	beams. In order to verify the accuracy of the present approximate solution, numerical results are compared with those of exact solutions using classical beam theory. The			
Keywords:	displacements and nominal stress distributions in the depth direction are obtained with			
Steel concrete beam	various high of beams. The present approximate approach can accurately predict the displacements and stresses of steel concrete beams.			
Spproximate solution	r			
Higher-order beam	F. ASMA & H. HAMMOUM (Eds.) special issue, 4 th International Conference on Sustainability in Civil Engineering ICSCE 2022, Hanoi, Vietnam, J. Mater. Eng. Struct. 9(4) (2022)			
Trigonometric series.				

1 Introduction

The steel-concrete beam is commonly used in civil engineering, which is composed of steel beams cast in situ with reinforced concrete slabs. This configuration utilizes sufficiently high capacities in the compression of concrete and steel

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tension. Thus, the steel-concrete beam has improved the stiffness and load capacity of the structure compared to those of the individual ones.

The literature reviews showed several works concerned with experimental and computational research on steel-concrete composite structures in civil and bridge engineering. P.V Phe and N.X. Huy [1] analysed the interfacial shear and peeling stresses of FRP-strengthened beams subjected to load and thermal effects by the finite difference method. Wang [2] studied the deflection of steel-concrete beams when there is a partial adhesion between steel and concrete. Chen Tao et al. [3] focused on the behaviour of steel-concrete beams with web openings under a negative moment. Jorge Luis Palomino et al. [4] studied the short-term performance of steel-concrete beams by developing a finite element model using plate elements. Pham Van Phe et al. [5] presented the computing of steel beams strengthened with composite panels by the finite element method based on the principle of stationary complementary energy. Sliseris and Korjakins [6] conducted an experimental study on high-performance concrete containing steel fibre beams by numerical method. Sugar Bükülmez Pınar and Celik Oguz [7] presented an experimental work on the fire resistance capacity of steel-concrete beams with a large openings ratio. Mahmoud Ashraf Mohamed [8] used finite element software to simulate steel-concrete beams with a nonlinear material.

There are many different computation beam theories depending on the members' dimensions, materials, and behavior that can be used when computing beams or plates. Classical beam theory can be used as presented on the issues of the strength of materials or structural mechanics[9], Timoshenko beam theory [10], and beam theory considering higher-order shear deformation [11]. L.T Ha and N.T.K Khue calculated the free oscillation of porous nanobeams using classical beam theory [12]. N.T Nhung et al. [13] investigated the stochastic free vibration of non-uniform beams using stochastic finite element method. Gianluca Ranzia and Alessandro Zona [14] presented a study on computing steel-concrete beams using both classical beam theory. Due to the neglect of shear strain in classical beam theory, many researchers have used higher-order beam theory. In which, the higher-order shear deformation term has been considered. Thus, increasing the accuracy of the observations. Atilla Ozutoka and Emrah Madenci [15] calculated multilayer composite beams using the finite element method by considering the higher-order shear deformation. T.D Hien et al. [16] developed the stochastic finite element method for higher-order beams involving the random field of elastic modulus. P.B Thang and L.L.Anh [17] analyzed the steel-concrete beam by using classical beam theory and Timoshenko beam theory considering the anterior of the observation between connectors and reinforced concrete slab. Fabrizio Gara et al. [18] proposed the finite element model steel-concrete beams considering the partial shear interaction between concrete and steel.

Although numerous types of research on steel-concrete beams have been presented in the literature, there are still a few works that consider the higher-order shear deformation of beams. Thus, it is necessary to research the computing method for steel-concrete beams, in which, the higher-order shear deformation is taken into account. Finally, the observations on internal forces and displacements are predicted close to the actual behaviour of the beams.

2 Governing equation of steel-concrete beam using Reddy's beam theory

Considering a schematic of the steel-concrete beam as shown in Fig. 1, a popular structure used in bridge construction.



Fig. 1 – Cross-section of steel-concrete beam

Following the classical beam theory, the reference coordinate is usually chosen at the centroid of the section. Then, the axial displacement will be neglected. In this study, when the theory of higher-order beams is applied, determining the neutral

axis is complex following the above. Hence, the reference axis is determined at the middle of the cross-section's height. Therefore, axial displacement exists. The displacement field of the steel-concrete beam under Reddy's theory [11] is determined with the origin located in the middle of the beam's height as follows:

$$u(x,z) = u_0(x) + z\psi(x) - \frac{4z^3}{3h^2} \left(\psi(x) + \frac{\partial w_0}{\partial x} \right); \quad w(x,z) = w_0(x)$$
(1)

where u_0, w_0, ψ are the displacement components at the mid-axis in the x, z directions, and the displacement parameter, respectively.

From the displacement field in Eq.(1), deformation can be calculated as follows:

$$\varepsilon_{x} = \frac{\partial u_{0}}{\partial x} + \frac{\partial \psi}{\partial x} z - \frac{4z^{3}}{3h^{2}} \left(\frac{\partial \psi}{\partial x} + \frac{\partial^{2} w_{0}}{\partial x^{2}} \right); \quad \gamma_{xz} = \psi \left(1 - \frac{4z^{2}}{h^{2}} \right) + \frac{\partial w_{0}}{\partial x} \left(1 - \frac{4z^{2}}{h^{2}} \right)$$
(2)

The deformation potential of the beam:

$$U = \frac{1}{2} \int_{\Omega} E(z) \left[\frac{\partial u_0}{\partial x} + \frac{\partial \psi}{\partial x} z - \frac{4}{3h^2} z^3 \left(\frac{\partial \psi}{\partial x} + \frac{\partial^2 w_0}{\partial x^2} \right) \right] d\Omega + \frac{1}{2} \int_{\Omega} G(z) \left[\left(1 - \frac{4z^2}{h^2} \right) \left(\psi + \frac{\partial w_0}{\partial x} \right) \right] dx$$
(3)

The beam's equilibrium position is determined according to the minimum potential energy principle:

$$\delta(U+\Pi) = 0 \tag{4}$$

The Eq. (4) can write as variations formulation:

$$\int_{\Omega} E(z) \left[\frac{\partial u_0}{\partial x} + \frac{\partial \psi}{\partial x} z - \frac{4}{3h^2} z^3 \left(\frac{\partial \psi}{\partial x} + \frac{\partial^2 w_0}{\partial x^2} \right) \right] \times \delta \left[\frac{\partial u_0}{\partial x} + \frac{\partial \psi}{\partial x} z - \frac{4}{3h^2} z^3 \left(\frac{\partial \psi}{\partial x} + \frac{\partial^2 w_0}{\partial x^2} \right) \right] d\Omega + \int_{\Omega} G(z) \left[\left(1 - \frac{4z^2}{h^2} \right) \left(\psi + \frac{\partial w_0}{\partial x} \right) \right] \delta \left[\left(1 - \frac{4z^2}{h^2} \right) \left(\psi + \frac{\partial w_0}{\partial x} \right) \right] dx - \int_{\Omega} q(x) \delta w_0(x) d\Omega = 0$$
(5)

Transforming and isolating the differential variables in Eq. (5) we have a system of differential equations of the beam:

$$\delta u_{0}: \quad \frac{\partial N_{x}}{\partial x} = 0 \quad \delta w_{0}: \quad \frac{\partial Q_{x}}{\partial x} - \frac{4}{h^{2}} \frac{\partial R_{x}}{\partial x} + \frac{4}{3h^{2}} \frac{\partial^{2} P_{x}}{\partial x^{2}} + q = 0$$

$$\delta \psi: \quad \frac{\partial M_{x}}{\partial x} - \frac{4}{3h^{2}} \frac{\partial P_{x}}{\partial x} - Q_{x} + \frac{4}{h^{2}} R_{x} = 0$$
(6)

In which the internal force components are defined as follows: ~

$$\begin{cases} N_x \\ M_x \\ P_x \end{cases} = \int_A E(z) \varepsilon_x \begin{cases} 1 \\ z \\ z^3 \end{cases} dA; \quad \begin{cases} Q_x \\ R_x \end{cases} = \int_A G(z) \gamma_{xz} \begin{cases} 1 \\ z^2 \end{cases} dA$$
(7)

Substituting the expression of displacement in Eq. (2) into Eq. (8), we have:

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$$\begin{cases}
N_{x} \\
M_{x} \\
P_{x}
\end{cases} = \int_{A} E(z) \begin{bmatrix}
\frac{\partial u_{0}}{\partial x} + \frac{\partial \psi}{\partial x} \left(z - \frac{4z^{3}}{3h^{2}}\right) \\
-\frac{4z^{3}}{3h^{2}} \frac{\partial^{2} w_{0}}{\partial x^{2}}
\end{bmatrix} \begin{cases}
1 \\
z \\
z^{3}
\end{cases} dA; \quad \begin{cases}
Q_{x} \\
R_{x}
\end{cases} = \int_{A} G(z) \begin{bmatrix}
\psi \left(1 - \frac{4z^{2}}{h^{2}}\right) \\
+ \frac{\partial w_{0}}{\partial x} \left(1 - \frac{4z^{2}}{h^{2}}\right) \\
+ \frac{\partial w_{0}}{\partial x} \left(1 - \frac{4z^{2}}{h^{2}}\right)
\end{bmatrix} \begin{cases}
1 \\
z^{2}
\end{cases} dA \quad (8)$$

Computing the integrals in Eq. (9) we obtain the internal force expressions under the displacement parameters as follows:

$$\begin{cases}
N_{x} \\
M_{x} \\
P_{x}
\end{cases} = \begin{bmatrix}
A_{1} & A_{2} & A_{3} \\
B_{1} & B_{2} & B_{3} \\
D_{1} & D_{2} & D_{3}
\end{bmatrix}
\begin{cases}
\frac{\partial u_{0}}{\partial x} \\
\frac{\partial \psi}{\partial x} \\
\frac{\partial^{2} w_{0}}{\partial x^{2}}
\end{cases}; \quad
\begin{cases}
Q_{x} \\
R_{x}
\end{cases} = \begin{bmatrix}
F_{1} & F_{2} \\
H_{1} & H_{2}
\end{bmatrix}
\begin{cases}
\psi \\
\frac{\partial w_{0}}{\partial x}
\end{cases}$$
(9)

In which the constants of the properties of a plane area of the cross-section are defined as follows:

)

$$\begin{cases} A_{1} \\ A_{2} \\ A_{3} \\ \end{pmatrix} = \int_{A} E(z) \begin{cases} 1 \\ z - \frac{4z^{3}}{3h^{2}} \\ -\frac{4z^{3}}{3h^{2}} \\ -\frac{4z^{3}}{3h^{2}} \\ \end{cases} dA; \begin{cases} B_{1} \\ B_{2} \\ B_{3} \\ -\frac{4z^{3}}{3h^{2}} \\ -\frac{4z^{3}$$

Analytical solution 3

Considering a beam with simple hinges at both ends, from the boundary conditions, we use the approximate solution based on a trigonometric series of Navier [11]:

$$u_0 = \sum_{i=1}^{\infty} U_i \cos(\alpha x); \quad w_0 = \sum_{i=1}^{\infty} W_i \sin(\alpha x); \quad \psi_0 = \sum_{i=1}^{\infty} \Psi_i \cos(\alpha x)$$
(11)

where
$$\alpha = \frac{i\pi}{L}$$

By substituting into Eq. (9), we observe:

$$\begin{cases}
N_{x} \\
M_{x} \\
P_{x}
\end{cases} = \sum_{i=1}^{\infty} \begin{bmatrix}
A_{1} & A_{2} & A_{3} \\
B_{1} & B_{2} & B_{3} \\
D_{1} & D_{2} & D_{3}
\end{bmatrix} \begin{bmatrix}
-\alpha U_{i} \\
-\alpha \Psi_{i} \\
-\alpha^{2}W_{i}
\end{bmatrix} \sin(\alpha x); \begin{cases}
Q_{x} \\
R_{x}
\end{cases} = \sum_{i=1}^{\infty} \begin{bmatrix}
F_{1} & F_{2} \\
H_{1} & H_{2}
\end{bmatrix} \begin{bmatrix}
\alpha \Psi_{i} \\
\alpha W_{i}
\end{bmatrix} \cos(\alpha x)$$
(12)

Substituting the internal force expressions from Eq. (12) into Eq. (6) and using the orthogonal property of the sine and cosine functions, Eq. (6) is altered as follows:

$$\begin{bmatrix} A_{1} & A_{2} & \alpha A_{3} \end{bmatrix} \begin{cases} U_{i} \\ \Psi_{i} \\ W_{i} \end{cases} = 0$$

$$\begin{bmatrix} \alpha^{3} \frac{4}{3h^{2}} D_{1} & \alpha^{3} \frac{4}{3h^{2}} D_{2} - \alpha^{2} F_{1} + \alpha^{2} \frac{4}{h^{2}} H_{1} & \alpha^{4} \frac{4}{3h^{2}} D_{3} - \alpha^{2} F_{2} + \alpha^{2} \frac{4}{h^{2}} H_{2} \end{bmatrix} \begin{cases} U_{i} \\ \Psi_{i} \\ W_{i} \end{cases} = 0$$

$$\begin{bmatrix} -\alpha^{2} B_{1} - \alpha^{2} \frac{4}{3h^{2}} D_{1} & -\alpha^{2} B_{2} - \alpha^{2} \frac{4}{3h^{2}} D_{2} & -\alpha^{3} B_{3} - \alpha^{3} \frac{4}{3h^{2}} D_{3} \\ -\alpha F_{1} + \alpha \frac{4}{h^{2}} H_{1} & -\alpha F_{2} + \alpha \frac{4}{h^{2}} H_{2} \end{bmatrix} \begin{cases} U_{i} \\ \Psi_{i} \\ W_{i} \end{cases} = 0$$

$$\begin{bmatrix} W_{i} \\ W_{i} \end{bmatrix} = 0$$

$$\begin{bmatrix} W_{i} \\ W_{i} \end{bmatrix} = 0$$

or abbreviation form:

$$\begin{bmatrix} S_{11} & S_{12} & S_{13} \\ S_{21} & S_{22} & S_{23} \\ S_{31} & S_{32} & S_{33} \end{bmatrix} \begin{bmatrix} U_i \\ \Psi_i \\ W_i \end{bmatrix} + \begin{bmatrix} 0 \\ q_i \\ 0 \end{bmatrix} = 0$$
(14)

where load parameter:

$$q_i = \frac{2}{L} \int_0^L q(x) \sin(\alpha x) dx \tag{15}$$

and S_{ii} is the parameter of stiffness

$$S_{11} = A_{1}; S_{12} = A_{2}; S_{13} = \alpha A_{3}$$

$$S_{21} = \alpha^{3} \frac{4}{3h^{2}} D_{1}; S_{22} = \alpha^{3} \frac{4}{3h^{2}} D_{2} - \alpha^{2} F_{1} + \alpha^{2} \frac{4}{h^{2}} H_{1};$$

$$S_{23} = \alpha^{4} \frac{4}{3h^{2}} D_{3} - \alpha^{2} F_{2} + \alpha^{2} \frac{4}{h^{2}} H_{2}$$

$$S_{31} = -\alpha^{2} B_{1} - \alpha^{2} \frac{4}{3h^{2}} D_{1}$$

$$S_{32} = -\alpha^{2} B_{2} - \alpha^{2} \frac{4}{3h^{2}} D_{2} - \alpha F_{1} + \alpha \frac{4}{h^{2}} H_{1}$$

$$S_{33} = -\alpha^{3} B_{3} - \alpha^{3} \frac{4}{3h^{2}} D_{3} - \alpha F_{2} + \alpha \frac{4}{h^{2}} H_{2}$$
(16)

By solving the system of equations (14) to obtain the displacement constants U_i, Ψ_i, W_i . So, the displacement is determined following Eq. (11). Based on those results, the deformation and stress of the beam can be calculated.

4 Numerical examples

To estimate the accuracy of the approximate solution based on a trigonometric series, we consider a simple span of steelconcrete beam, length L=15m, and cross-sectional dimensions are defined in Fig. 1.

Dimensions of the cross-section:

 $h = 1, 2m; h_1 = 0, 02m; h_3 = 0, 02m; h_4 = 0, 2m;$ $b_1 = 0, 3m; b_2 = 0, 014m; b_3 = 0, 3m; b_4 = 2, 5m;$

Applied uniform load: q=10kN/m.

We use the classical beam theory to estimate the simulation results as presented in the same material. Then, the beam deflection is observed of 3.025 mm (w=3,025(mm)). In addition to using the theoretical formula, the displacements can be checked by the finite element method. Specifically, in this present work, Sap2000 software is utilized, a beam element simulate the steel beam, and the concrete slab is simulated by a plate element. The simulation result shows the displacement in the middle of the beam is w=3.205(mm).

Using the trigonometric series in Eq. (11), the displacement observation in the middle of the beam is w=3.131(mm) when computing with the first 15 terms of the series. Thus, it shows that the displacement result computed according to the theory of higher-order beams is 3.3 % higher than that computed according to classical beam theory but lower than that calculated by finite element simulation by 2.36%. When computing the beam using classical theory, because of the neglection of the shear deformation, the observed displacement is lower. However, when simulating finite elements using a plate model for concrete slab, the result will be close to the actual behavior of the beam, the displacement results are larger than the two theoretical ones using the classical beam theory and higher-order beam theory.

For analyzing the influence of shear deformation on the internal force and displacement of the beam, we investigate the internal force and displacement of the beam, which the height varies from 0.8m to 1.5m, and the observations are compared with the observations according to the classical beam theory without considering the shear deformation. Table 1 presents the observations of the maximum deflection and normal stress at the mid-span and lower boundary fiber. The observations clearly show that when the height increases, the influence of the shear deformation increases, leading to an increase significantly of the deflection computed according to the higher-order shear deformation theory compared that computed according to the classical beam theory is higher than that computed according to the classical beam theory by 6.1%. In the case of the normal stress at the lower boundary fiber of the beam, it remained unchanged when computing the

normal stress according to the classical beam theory and the higher-order shear theory, although the difference increases as the height increases.

	Displacement (mm)			Norn	Normal stress (MPa)		
h(m)	Higher-order beam theory	Classical beam theory	Error (%)	Higher-order beam theory	Classical beam theory	Error (%)	
0,8	8,46	8,379	0,96%	43,822	43,813	0,02%	
0,9	6,341	6,252	1,40%	37,171	37,157	0,04%	
1,0	4,895	4,799	1,96%	31,93	31,911	0,06%	
1,1	3,874	3,772	2,63%	27,739	27,716	0,08%	
1,2	3,131	3,025	3,39%	24,313	24,342	0,12%	
1,3	2,576	2,467	4,23%	21,548	21,514	0,16%	
1,4	2,152	2,042	5,11%	19,225	19,186	0,20%	
1,5	1,822	1,711	6,09%	17,271	17,227	0,25%	

Table 1 – Comparison the displacement and normal stress at of mid-span.

5 Conclusions

The work presents an approach to compute the internal forces and displacements of steel-concrete beams by using the approximate solution based on a trigonometric series, in which, the higher-order shear deformation is considered. The comparison of the results based on the present approach and the results based on classical beam theory and finite element simulation shows the reliability of the approximate solution based on a trigonometric series. The observations from a steel-concrete beam with a height to span length ratio of 1/20 - 1/10 show higher displacement and stress based on the higher-order beam theory compared to those based on classical beam theory. As the beam height increases, the displacement increases significantly while the normal stress increases slightly. As a result, the effect of the shear deformation on computing the beam, which is height relatively higher than the span length, is clarified.

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