



UvA-DARE (Digital Academic Repository)

Telecloning of qudits via partially entangled states

Araneda, G.; Cisternas, N.; Delgado, A.

DOI

[10.1007/s11128-016-1348-2](https://doi.org/10.1007/s11128-016-1348-2)

Publication date

2016

Document Version

Final published version

Published in

Quantum Information Processing

License

CC BY

[Link to publication](#)

Citation for published version (APA):

Araneda, G., Cisternas, N., & Delgado, A. (2016). Telecloning of qudits via partially entangled states. *Quantum Information Processing*, 15(8), 3443-3458. <https://doi.org/10.1007/s11128-016-1348-2>

General rights

It is not permitted to download or to forward/distribute the text or part of it without the consent of the author(s) and/or copyright holder(s), other than for strictly personal, individual use, unless the work is under an open content license (like Creative Commons).

Disclaimer/Complaints regulations

If you believe that digital publication of certain material infringes any of your rights or (privacy) interests, please let the Library know, stating your reasons. In case of a legitimate complaint, the Library will make the material inaccessible and/or remove it from the website. Please Ask the Library: <https://uba.uva.nl/en/contact>, or a letter to: Library of the University of Amsterdam, Secretariat, Singel 425, 1012 WP Amsterdam, The Netherlands. You will be contacted as soon as possible.

Telecloning of qudits via partially entangled states

Gabriel Araneda¹ · Nataly Cisternas² ·
Aldo Delgado^{3,4,5}

Received: 11 February 2016 / Accepted: 12 May 2016 / Published online: 21 May 2016
© The Author(s) 2016. This article is published with open access at Springerlink.com

Abstract We study the process of quantum telecloning of d -dimensional pure quantum states using partially entangled pure states as quantum channel. This process efficiently mixes optimal universal symmetric cloning with quantum teleportation. It is shown that it is possible to implement universal symmetric telecloning in a probabilistic way using unambiguous state discrimination and quantum state separation schemes. It is also shown that other strategies, such as minimum error discrimination, lead to a decrease in the fidelity of the copies and that certain partially entangled pure states with maximal Schmidt rank lead to an average telecloning fidelity which is always above the optimal fidelity of measuring and preparation of quantum states. We also discuss the case of partially entangled pure states with non-maximal Schmidt rank. The results presented here are valid for arbitrary numbers of copies of a single-input qudit state of any dimension.

Keywords Quantum cloning · Quantum teleportation · Quantum telecloning · Quantum state discrimination

✉ Gabriel Araneda
gabriel.araneda-machuca@uibk.ac.at

¹ Institut für Experimentalphysik, Universität Innsbruck, Technikerstrasse 25, 6020 Innsbruck, Austria

² Van der Waals-Zeeman Institute for Experimental Physics, Universiteit van Amsterdam, Science Park 904, 1098 XH Amsterdam, The Netherlands

³ Center for Optics and Photonics, Universidad de Concepción, Casilla 4012, Concepción, Chile

⁴ MSI-Nucleus on Advanced Optics, Universidad de Concepción, Casilla 160-C, Concepción, Chile

⁵ Departamento de Física, Universidad de Concepción, Casilla 160-C, Concepción, Chile

1 Introduction

The encoding of information in quantum systems exhibits features which do not have classical counterparts, including quantum teleportation [1,2] and the no-cloning theorem [3]. Quantum teleportation enables quantum states to be transmitted from a sender, Alice, to a receiver, Bob, without sending a physical quantum system. The no-cloning theorem rules out the perfect cloning of unknown quantum states by the linearity of quantum mechanical transformations. It does allow approximate deterministic cloning [4–7] and perfect probabilistic cloning [8–10]. Approximate deterministic cloning is the creation of imperfect copies of an unknown quantum state with the maximal fidelity permitted by quantum mechanics. Perfect probabilistic cloning is the creation of exact copies of an unknown quantum state in a probabilistic fashion.

Quantum teleportation and approximate universal cloning can be combined to allow a sender to distribute approximate clones to many receiving parties. The sender may first generate all the required clones before teleporting them to the receivers. This strategy requires the use of as many bipartite maximally entangled states as there are clones to be broadcast, as well as the transmission of classical information to complete the teleportation stage. A new and more efficient approach, termed telecloning [11,12], allows simultaneous conveyance of all the clones by means of a single local measurement carried out by the sender, who publicly broadcasts the result of the measurement. Each receiving party then performs a local quantum operation which depends on the publicly broadcast measurement result, and thereafter the approximate clones are finally broadcast. This kind of scheme has been experimentally demonstrated using entangled photons [13], where three non-perfect but optimal copies were transmitted using maximally entangled bipartite pure states. This strategy is more efficient than the simple scheme described above since it requires only $O(\log_2 M)$ maximally entangled bipartite pure states, where M is the number of clones to be distributed.

Telecloning is based on the generation of a multipartite entangled state which is distributed among the parties participating in the telecloning process. In realistic situations, however, generation and distribution do not, in general, lead to a maximally entangled state. Here we study the performance of telecloning when implemented with a partially entangled pure state as a quantum channel. We focus on the errors introduced in the telecloning process and the degradation of the fidelity of the clones. We show that the problem of recovering high fidelity in the copies reduces to the problem of discrimination between non-orthogonal quantum states. We also propose a probabilistic and conclusive method to correct the errors caused by the non-maximal entanglement of the multipartite entangled state, using standard strategies of quantum state discrimination. This article is organised as follows: in Sect. 2, we briefly review the telecloning process. In Sect. 3, we show how the problem of telecloning via a partially entangled pure quantum state can be related to the problem of discrimination of non-orthogonal quantum states. In Sect. 4, we consider the case of a pure quantum channel of full Schmidt rank, and we combine unambiguous state discrimination with the telecloning scheme to correct the errors introduced by the non-maximal entanglement of the multipartite entangled state. We show that the fidelity of this process, averaged over the Hilbert space of the states to be cloned, is higher than the fidelity

of estimating quantum states for certain quantum channels. Also in this section, we study the combination with other discrimination schemes, namely minimum error and quantum state separation. In Sect. 5, we comment the case of a pure quantum channel with non-maximal Schmidt rank, and our results are summarised in Sect. 6.

2 Telecloning process

The linearity of quantum operations forbids the perfect cloning of unknown quantum states. Nevertheless, it is possible to consider an approximate cloning process which creates clones of unknown quantum states with the highest fidelity allowed by quantum mechanics. This optimal process is carried out by applying separate unitary transformations onto a set of quantum systems. We consider a universal and symmetric telecloning process [11, 12]. It is universal in the sense that the fidelities of the clones do not depend on the particular state undergoing the cloning process, and it is symmetric in that all of the clones have the same fidelity with respect to the input state. Additionally, each of the copies and the system to be cloned may be spatially separated.

Let us consider the case of creating M clones from a single copy of an arbitrary qudit of dimension d , denoted as $1 \rightarrow M$ telecloning [12]. The maximal fidelity for this kind of process is given by the optimal fidelity of the $1 \rightarrow M$ universal cloning, $F_{\text{opt}}^{1 \rightarrow M} = (2M + d - 1)/(M + Md)$. The input state to be teleported is the qudit state $|\psi\rangle_X = \sum_{j=0}^{d-1} \alpha_j |j\rangle_X$, where the set of states $\{|j\rangle\}$ is the computational basis for the system X . The output state is represented through the basis $|\phi_j\rangle$, which describes $N_0 = 2M - 1$ d -dimensional systems, where $M - 1$ of them are ancillary systems and the remaining M systems encode the clones. The elements of this basis are written in terms of the basis of normalised and symmetrised states $\{|\xi_k^M\rangle\}$. They are given by

$$|\phi_j\rangle_{AC} = \frac{\sqrt{d}}{\sqrt{d[M]}} \sum_{k=0}^{d[M]} P(j|\xi_k^M)_{PA} \otimes |\xi_k^M\rangle_C, \quad (1)$$

where $d[M] = (d + M - 1)!/M!(d - 1)!$. The index P represents the port qudit, which belongs to Alice, A represents the $M - 1$ ancillas, and C represents the M qudits which host the clones. A constructive procedure for the states $\{|\xi_k^M\rangle_C\}$ is detailed in [12]. The optimal quantum channel for the telecloning process is given by

$$|\xi\rangle_{PAC} = \frac{1}{\sqrt{d}} \sum_{j=0}^{d-1} |j\rangle_P \otimes |\phi_j\rangle_{AC}, \quad (2)$$

which is a maximally entangled state between the port qudit and the AC system. Therefore, the joint state of the total system (input state and quantum channel) is given by

$$|\psi\rangle_{XPAC} = |\psi\rangle_X \otimes |\xi\rangle_{PAC}. \quad (3)$$

This state can be cast as

$$|\psi\rangle_{XPAC} = \sum_{n,m=0}^{d-1} |\Phi_{nm}\rangle_{XP} \frac{1}{\sqrt{d}} \sum_{j=0}^{d-1} \omega^{jn} \alpha_j |\phi_{j\oplus m}\rangle_{AC}, \quad (4)$$

where $\omega = \exp(2\pi i/d)$ and the generalised Bell states $\{|\Phi_{nm}\rangle\}$ ($n, m = 0, \dots, d-1$) are given by

$$|\Phi_{nm}\rangle_{XP} = \frac{1}{\sqrt{d}} \sum_{k=0}^{d-1} \omega^{kn} |k\rangle_X |k \oplus m\rangle_P, \quad (5)$$

with $k \oplus m = (j + m) \bmod(d)$. Equivalently, the states of the separable basis of systems X and P are given in terms of the Bell states by the expression

$$|k\rangle_X |k \oplus m\rangle_P = \frac{1}{\sqrt{d}} \sum_{n=0}^{d-1} \omega^{-nk} |\Phi_{nm}\rangle_{XP}. \quad (6)$$

We now perform a measurement on the systems X and P in the $\{|\Phi_{nm}\rangle\}$ basis. The outcomes of this measurement are the eigenvalues associated with the indexes n and m , and they are transmitted by means of classical communications to the carriers of systems A and C . In order to recover the optimal copies, the carriers of these systems apply local reconstruction unitary operations, conditioned to the outcomes of the measurement. These operations are given by

$$U_{nm}^A = \sum_{j=0}^{d-1} \omega^{-jn} |j\rangle \otimes \langle j \oplus m|, \quad (7)$$

$$U_{nm}^C = \sum_{j=0}^{d-1} \omega^{jn} |j\rangle \otimes \langle j \oplus m|, \quad (8)$$

for the ancillas and copy systems, respectively. Therefore, the state in the AC system is $|\psi\rangle_{AC} = \sum_{k=0}^{d-1} \alpha_j |\phi_j\rangle_{AC}$, which contains optimal copies of the input state in each system C . The fidelity of each copy is optimal and given by $F_{\text{opt}}^{1 \rightarrow M}$.

3 Telecloning via partial entanglement

Now let us suppose that the quantum channel given by Eq. (2) is not a maximally entangled state between the port qudit and the AC system and, instead, is given by the state

$$|\tilde{\xi}\rangle_{PAC} = \sum_{j=0}^{d-1} c_j |j\rangle_P \otimes |\phi_j\rangle_{AC}, \quad (9)$$

where $\{c_j\}$ are d coefficients which define the channel and satisfy $\sum_{j=0}^{d-1} |c_j|^2 = 1$. We assume that the channel is written in its Schmidt decomposition so that all the c_j

are real, positive numbers. In this way, the state of the total system is given by

$$|\psi\rangle_{XPAC} = |\psi\rangle_X \otimes |\tilde{\xi}\rangle_{PAC}, \quad (10)$$

which can be cast as

$$|\psi\rangle_{XPAC} = \sum_{n,m=0}^{d-1} |\Phi_{nm}\rangle_{XP} \sum_{j=0}^{d-1} \alpha_j c_{j+m} \exp(-2\pi i j n/d) |\phi_{j+m}\rangle_{AC}. \quad (11)$$

If we perform a measurement in the generalised Bell basis Eq. (5), the state of the system is projected onto

$$|\psi\rangle_{nm}^{AC} = \frac{1}{\sqrt{P_m}} \sum_{j=0}^{d-1} \alpha_j c_{j+m} \exp(-2\pi i j n/d) |\phi_{j+m}\rangle_{AC}, \quad (12)$$

where m and n denote the generalised Bell state associated with the measurement result and P_m is the probability of projecting onto the $|\Phi\rangle_{nm}$ state. This probability depends only on m ,

$$P_m = \sum_{j=0}^{d-1} |\alpha_j|^2 |c_{j+m}|^2. \quad (13)$$

Since the values of n and m are known after this measurement, it is possible to apply the corresponding local reconstruction unitary operations U_{nm}^A and U_{nm}^C on the ancillas and copies systems. After these operations, the state of the system is described by

$$|\psi'\rangle_{nm}^{AC} = \frac{1}{\sqrt{P_m}} \sum_{j=0}^{d-1} \alpha_j c_{j+m} |\phi_j\rangle_{AC}, \quad (14)$$

and the associated density matrix can be written as

$$\rho_{nm}^{AC} = \frac{1}{\sqrt{P_m}} \sum_{j,j'=0}^{d-1} \alpha_j c_{j+m} \alpha_{j'} c_{j'+m} |\phi_j\rangle_{AC} \langle\phi_{j'}|_{AC}. \quad (15)$$

Let us study the fidelity in the simplest case in which two qudit copies ($M = 2$) are created. In this case, the vectors $|\phi_j\rangle_{AC}$ can be written explicitly as

$$|\phi_j\rangle_{AC_1C_2} = \sqrt{\frac{1}{2(d+1)}} \sum_{k=0}^{d-1} |k\rangle_A (|jk\rangle_{C_1C_2} + |kj\rangle_{C_1C_2}). \quad (16)$$

The reduced density matrix of one of the copies, e.g. C_2 , is then

$$\begin{aligned} \rho_{nm}^{C_2} &= \frac{1}{2(d+1)} \sum_{l=0}^{d-1} |l\rangle\langle l| \\ &+ \frac{(2+d)}{2(d+1)} \frac{1}{P_m} \sum_{j,j'=0}^{d-1} [\alpha_j c_{j+m} \alpha_{j'} c_{j'+m}] |j\rangle\langle j'|. \end{aligned} \tag{17}$$

Hence, the local fidelity of the clones $F_m^{1 \rightarrow 2, PE} = \langle \psi | \rho_{nm}^{C_2} | \psi \rangle$ with respect to the initial input state for a partially entangled channel is given by

$$F_m^{1 \rightarrow 2, PE} = \frac{1}{2(d+1)} + \frac{(2+d)}{2(d+1)} \frac{1}{P_m} \left(\sum_{k=0}^{d-1} |\alpha_k|^2 c_{k+m} \right)^2. \tag{18}$$

The average fidelity of the process for a fixed channel is the sum of all the possible fidelities weighted by the probability of obtaining the different results in the generalised Bell measurement, that is $F^{1 \rightarrow 2, PE} = \sum_{m=0}^{d-1} P_m F_m^{1 \rightarrow 2, PE}$, which takes the explicit form

$$F^{1 \rightarrow 2, PE} = \frac{1}{2(d+1)} + \frac{(2+d)}{2(d+1)} \sum_{m=0}^{d-1} \left(\sum_{k=0}^{d-1} |\alpha_k|^2 c_{k+m} \right)^2. \tag{19}$$

This fidelity is always smaller than the optimal universal cloning fidelity $F_{opt}^{1 \rightarrow 2}$. In the case of qubits, the average fidelity of telecloning using a partially entangled state reduces to

$$F_{d=2}^{1 \rightarrow 2, PE} = \frac{N_\theta^2}{12} \left[5|a|^4 c_1^2 + 5|b|^4 c_2^2 + |a|^2 |b|^2 (1 + 8c_1 c_2) \right], \tag{20}$$

where

$$N_\theta = \frac{\sqrt{2}}{\sqrt{a^2 c_1^2 + b^2 c_2^2}}. \tag{21}$$

As we can see from Eq. (19), the lack of maximal entanglement not only reduces the fidelity of the broadcast clones, but also makes the cloning process state dependent. As indicated by Eq. (17), this is due to the fact that the coefficients of the channel are introduced in the state of the clones. It is possible to improve the fidelities of the copies, even to optimal fidelity by using additional steps. In the following sections, we distinguish between the cases where the channel can be described by a state with maximal Schmidt rank and the cases in which it may not be described in this fashion.

4 Quantum channel of maximal Schmidt rank

In the previous section, we have shown that the use of a partially entangled channel leads to a telecloning process characterised by a suboptimal state-dependent fidelity.

Here, we will show that it is possible to achieve the optimal state-independent fidelity at the expense of making the telecloning process probabilistic.

This approach allows the coefficients of the state to be cloned to be decoupled from those of the channel. We can cast the state of the total system, Eq. (4), before the telecloning process, as

$$|\psi\rangle_{XPAC} = \frac{1}{d} \sum_{n,m=0}^{d-1} |\tilde{\Phi}_{nm}\rangle_{XP} \otimes U_{nm}^{-1} \sum_{j=0}^{d-1} \alpha_j |\phi_j\rangle_{AC}, \quad (22)$$

where the states $\{|\tilde{\Phi}_{nm}\rangle\}$ are defined by

$$|\tilde{\Phi}_{nm}\rangle_{XP} = \sum_{k=0}^{d-1} c_k \omega^{kn} |k \ominus m\rangle_X |k\rangle_P, \quad (23)$$

with $n \ominus m = (n - m) \bmod(d)$, and

$$U_{nm}^{-1} \sum_{j=0}^{d-1} \alpha_j |\phi_j\rangle_{AC} = \sum_{j=0}^{d-1} \omega^{-(j+m)n} \alpha_j |\phi_{j \oplus m}\rangle_{AC}. \quad (24)$$

The coefficients of the quantum channel are now transferred into the new set $\{|\tilde{\Phi}_{nm}\rangle\}$ of generalised Bell states. These are non-orthogonal and cannot be perfectly distinguished in a deterministic fashion. Therefore, the selection of the correct unitary U_{nm} is ambiguous and leads to errors in the clones.

In order to avoid this limitation and improve the performance of the telecloning process, we can use probabilistic processes such as quantum state discrimination [14, 15] or quantum state separation [16]. These processes allow the correct reconstruction operation U_{nm} to be chosen, or to improve the probability of choosing the correct reconstruction operation in a controlled way. It is also possible to use minimum error discrimination [17–19], which is a deterministic discrimination strategy. In Sects. 4.1, 4.2 and 4.3, we analyse the performance of the telecloning process when combined with each of these discrimination strategies.

4.1 Telecloning combined with unambiguous state discrimination

The set $\{|\tilde{\Phi}_{nm}\rangle\}$ of d^2 non-orthogonal, partially entangled states can be transformed into a set of separable states. This is done by applying the GXOR gate onto the XP system. This gate is defined as

$$\text{GXOR}_{PX}|n\rangle_P|m\rangle_X = |n\rangle_P|n \ominus m\rangle_X. \quad (25)$$

The GXOR gate is a generalisation of CNOT for qudits, and it can entangle and disentangle two qudit states. After applying this gate, the state given by Eq. (22) is transformed into the state

$$\begin{aligned} \text{GXOR}_{PX}|\psi\rangle_{PXC} &= \frac{1}{d} \sum_{n,m=0}^{d-1} \sum_{k=0}^{d-1} c_k \omega^{kn} |k\rangle_P |m\rangle_X \\ &\otimes U_{nm}^{-1} \sum_{j=0}^{d-1} \alpha_j |\phi_j\rangle_{AC}, \end{aligned} \tag{26}$$

where the systems X and P are factorisable. The X system is then projected into its computational basis by a projective measurement. The X system is projected onto the state $|m\rangle_X$ with probability

$$P_m = \sum_{k=0}^{d-1} |c_{k+m}|^2 |\alpha_k|^2. \tag{27}$$

The state of the total system, which depends on the results of the measurement on system X , becomes

$$\begin{aligned} |\psi_m\rangle_{XPAC} &= \frac{1}{\sqrt{P_m}} \frac{1}{d} \sum_{n=0}^{d-1} |\Psi_n\rangle_P |m\rangle_X \\ &\otimes \sum_{j=0}^{d-1} \omega^{-(j+m)n} \alpha_j |\phi_{j\oplus m}\rangle_{AC}, \end{aligned} \tag{28}$$

where the states

$$|\Psi_n\rangle_P = \sum_{k=0}^{d-1} c_k \omega^{nk} |k\rangle_P, \quad n = 0, \dots, d - 1 \tag{29}$$

are a set of d symmetric states. These are symmetric because they are defined by the successive action of the operator $Z = \sum_{j=0}^{d-1} \omega^j |j\rangle\langle j|$ over the seed state $|\Psi_0\rangle_P = \sum_{k=0}^{d-1} c_k |k\rangle_P$, namely $|\Psi_n\rangle_P = Z^n |\Psi_0\rangle_P$. Thus, the problem of discriminating among the d^2 states $|\tilde{\Phi}_{nm}\rangle$ of the bipartite system XP is reduced to distinguishing among the d states $|\Psi_n\rangle_P$ of the system P . These states are also mutually non-orthogonal, and so they cannot be deterministically distinguishable. We will now apply unambiguous state discrimination (USD) to the set $\{|\Psi_n\rangle_P\}$. It has been demonstrated that this discrimination scheme exists only for sets of linearly independent states [20]. In our case, this corresponds to a set of non-vanishing c_j coefficients, or equivalently, to a quantum channel of full Schmidt rank.

Now we apply the USD protocol to states $|\Psi_n\rangle_P$. This scheme corresponds to a unitary transformation U_{PX}^{USD} acting on the bipartite system PX followed by a von Neumann measurement on system X [21]. The unitary transformation U_{PX}^{USD} is given by

$$U_{PX}^{\text{USD}} |\Psi_n\rangle_P |m\rangle_X = \sqrt{pd} |u_n\rangle_P |m\rangle_X + \sqrt{1 - pd} |\chi_n\rangle_P |m \oplus 1\rangle_X, \tag{30}$$

where the set $\{|u_l\rangle_P\}$ is composed by d mutually orthogonal and distinguishable states of system P , the states $\{|\chi\rangle_P\}$ are d linearly dependent states, and pd is the probability

of successful discrimination of the states $|\Psi_n\rangle$. All the states $|\Psi_n\rangle$ are generated with the same probability $1/d$, so the discrimination probability is the same for each of them and is given by $p_d = d|c_{\min}|^2$ where c_{\min} is the channel coefficient with the smallest absolute value. The state $|m \oplus 1\rangle$ is orthogonal to the state $|m\rangle$; thus, it is possible to know when the discrimination process is successful by means of a projective measurement on system X . The states $\{|\chi_l\rangle\}$ are given by

$$|\chi_l\rangle = \frac{1}{\sqrt{d}} \frac{1}{\sqrt{1-p_d}} \sum_{m,n=0}^{d-1} \omega^{(l-m)n} \sqrt{c_n^2 - c_{\min}^2} |m\rangle. \tag{31}$$

These d states belong to a subspace of dimension $d-1$ and are thus linearly dependent, which forbids the possibility of applying a further stage of unambiguous discrimination. The states $\{|u_k\rangle\}$ turn out to be the Fourier transform of the computational basis [22], that is

$$|u_n\rangle = \mathcal{F}|n\rangle = \frac{1}{\sqrt{d}} \sum_{k=0}^{d-1} \omega^{kn} |k\rangle. \tag{32}$$

Therefore, the application of the inverse Fourier transform on the system P allows the states $|\Psi_n\rangle$ to be discriminated on the computational basis. First applying the discrimination unitary operation U and then the inverse Fourier transform, we find that the state of the joint system $PXAC$ is given by

$$\begin{aligned} |\tilde{\psi}_m\rangle = & \frac{\sqrt{p_d}}{d\sqrt{P_m}} \left(\frac{1}{\sqrt{d}} \sum_{n=0}^{d-1} |n\rangle_P U_{nm}^{-1} \sum_{j=0}^{d-1} \alpha_j |\phi_j\rangle_{AC} \right) |m\rangle_X \\ & + \frac{\sqrt{1-p_d}}{d\sqrt{P_m}} \left(\sum_{n=0}^{d-1} \mathcal{F}_P^{-1} |\chi_n\rangle_P U_{nm}^{-1} \sum_{j=0}^{d-1} \alpha_j |\phi_j\rangle_{AC} \right) |m \oplus 1\rangle_X, \end{aligned} \tag{33}$$

where we have defined $|\tilde{\psi}_m\rangle = \mathcal{F}_X^{-1} U_{PX}^{USD} |\psi_m\rangle_{XPAC}$. We now perform a projective measurement on system X . If the outcome of the measurement is the eigenvalue associated with the state $|m\rangle_X$, which occurs with probability p_d , then it is possible to conclusively discriminate the states of system P . We therefore obtain the values of n and m . These are transmitted by classical communications to the carriers of systems A and C , who separately apply the local reconstruction unitary operations U_{nm}^A on each ancilla system and U_{nm}^C on each system encoding a clone. These reconstruction operations are given by

$$U_{n,m}^A = \sum_{j=0}^{d-1} \omega^{-(j+m)n} |j\rangle \otimes |j \oplus m\rangle \tag{34}$$

and

$$U_{nm}^C = \sum_{j=0}^{d-1} \omega^{(j+m)n} |j\rangle \otimes |j \oplus m\rangle. \tag{35}$$

After these transformations, the copies encoded in the systems C are mixed states with optimal fidelity $F_{\text{opt}}^{1 \rightarrow M}$ with respect to the input qudit state. Thus, the combination of telecloning process via a partially entangled pure state and an unambiguous discrimination stage leads to a probabilistic telecloning process which produces optimal copies with probability $p_d = d|c_{\min}|^2$ for any values of M and d .

In the case that the discrimination process fails, with probability $1 - p_d$, the result of the measurement in the X system turns out to be $m \oplus 1$. The fidelity of the clones turns out to be non-optimal and depends on the state of system X after the first measurement (through m) and on the input state. Indeed, if the discrimination fails, the state of the AC system corresponds to the second term of Eq. (33), namely

$$|\tilde{\psi}_m\rangle^{\text{fail}} = \frac{1}{d\sqrt{P_m}} \left(\sum_{n=0}^{d-1} \mathcal{F}_P^{-1} |\chi_n\rangle_P U_{nm}^{-1} \sum_{j=0}^{d-1} \alpha_j |\phi_j\rangle_{AC} \right) |m \oplus 1\rangle_X. \tag{36}$$

Starting from Eq. (36), we can construct the density matrix for the final state in the case of a failed discrimination attempt and then obtain the reduced density matrix of one of the copies, namely C_2 , by calculating, in the case of the $1 \rightarrow 2$ telecloning,

$$\rho_{n,m}^{C_2,\text{fail}} = \text{Tr}_{C_1} \left[\text{Tr}_{PA} \left[|\tilde{\psi}_m\rangle^{\text{fail}} \langle \tilde{\psi}_m| \right] \right]. \tag{37}$$

The fidelity of the clones $F_m^{1 \rightarrow 2,\text{fail}} = \langle \psi | \rho_m^{C_2,\text{fail}} | \psi \rangle$ is given by

$$F_m^{1 \rightarrow 2,\text{fail}} = \frac{1}{2(d+1)} + \frac{1}{2(d+1)} \frac{1}{P_m} \sum_{j=0}^{d-1} |\alpha_{j+m}|^2 |\alpha_j|^2 \times \left\{ (2+d) (c_{j+m}^2 - c_{\min}^2) \right\}. \tag{38}$$

This depends on the result of the first measurement, given by the value of m , and on the state to be cloned, so that, in the case of failure, the process is neither universal nor optimal. The average failure fidelity is computed by integrating over all the possible input states and summing over all the possible outcomes of the first measurement on system X ,

$$F^{\text{fail}} = \int d\alpha \sum_{m=0}^{d-1} P_m F_m^{1 \rightarrow 2,\text{fail}}, \tag{39}$$

where $\int d\alpha$ represents integration over all pure d -dimensional qudits. Using the identity $\int d\psi |\psi_j|^2 |\psi_k|^2 = (\delta_{j,k} + 1)/d(d+1)$ and Eq. (38), we obtain

$$F^{\text{fail}} = \frac{1}{d}. \tag{40}$$

A simple cloning strategy consists of estimating the state to be cloned and then creating as many copies of the estimated state as needed. The fidelity of this process is given by

the optimal state estimation fidelity $F_{\text{est}} = 2/(d + 1)$. Thus, the average telecloning fidelity obtained in cases which the discrimination attempts fail is always smaller than that obtained when estimating.

Finally, as a measure of the quality of the total process obtained by concatenating telecloning to unambiguous state discrimination, it is possible to compute the total average fidelity F_{av} , which includes failures and successes in the discrimination stage. This average fidelity is obtained by adding the optimal fidelity of cloning $F_{\text{opt}}^{1 \rightarrow M}$ weighted by the optimal discrimination probability p_d and the average fidelity F^{fail} , given in Eq. (39), weighted by the failure probability $1 - p_d$:

$$F_{\text{av}} = F_{\text{opt}}^{1 \rightarrow M} p_d + F^{\text{(fail)}}(1 - p_d). \quad (41)$$

In the case of $1 \rightarrow 2$ telecloning of qudits, this fidelity is given by

$$F_{\text{av}}^{1 \rightarrow 2} = p_d \left(\frac{3 + d}{2 + 2d} \right) + (1 - p_d) \frac{1}{d}, \quad (42)$$

which depends on the quantum channel through p_d . We can compare this fidelity with the classical fidelity F_{set} for $1 \rightarrow 2$ cloning processes, i.e. the fidelity of the optimal measure-and-prepare cloner. For the $1 \rightarrow 2$ telecloning of qudits, we look for dimensions where

$$p_d \left(\frac{3 + d}{2 + 2d} \right) + (1 - p_d) \frac{1}{d} \geq \frac{2}{d + 1} \quad (43)$$

holds. In addition, the condition $p_d \leq 1$ also holds. This implies that $|c_{\text{min}}|^2 \leq 1/d$. Therefore, it is possible to outperform the classical cloning fidelity in every dimension d if the condition for the channel

$$|c_{\text{min}}|^2 \geq \frac{2}{d(d + 2)}, \quad (44)$$

is fulfilled. Hence, for certain quantum channels, even when including clones produced in cases with failed discrimination processes, it is possible to achieve a higher telecloning fidelity than in the classical case.

4.2 Telecloning combined with minimum error discrimination

Minimum error discrimination of quantum states [17–19, 23] is based on the minimisation of the average error when making guesses about a set of states. Operationally, following the same procedure as using unambiguous state discrimination, and after a projective measurement in the system X , the state of the PAC system is given by Eq. (28),

$$|\psi_m\rangle = \frac{1}{\sqrt{P_m}} \frac{1}{d} \sum_{n=0}^{d-1} |\Psi_n\rangle_P \sum_{j=0}^{d-1} \omega^{-n(j+m)} \alpha_j |\phi_{j \oplus m}\rangle_{AC}, \quad (45)$$

Now, in order to discriminate between the $|\Psi_n\rangle$ states, we apply the inverse Fourier transform over P ,

$$\mathcal{F}^{-1}|\Psi_n\rangle_P = \frac{1}{\sqrt{d}} \sum_{k=0}^{d-1} \sum_{j=0}^{d-1} c_k \omega^{k(n-j)} |j\rangle, \tag{46}$$

and subsequently, we projectively measure in the canonical basis. If the result is the one associated with the state $|n\rangle$, the state of the system is given by

$${}_P \langle n | \mathcal{F}_P^{-1} | \psi_m \rangle = \frac{1}{\sqrt{\tilde{P}_n}} \frac{1}{\sqrt{d}} \sum_{j=0}^{d-1} c_{j+m} \alpha_j \omega^{-n(j+m)} | \phi_{j \oplus m} \rangle_{AC}, \tag{47}$$

which is non-normalised, but with norm \tilde{P}_n . Knowing the values of n and m , we apply the reconstruction operations given by Eqs. (34) and (35). The new state of the AC system is given by

$$| \psi_{nm} \rangle = \frac{1}{\sqrt{\tilde{P}_n}} \frac{1}{\sqrt{P_m}} \frac{1}{\sqrt{d}} \sum_{j=0}^{d-1} c_{j+m} \alpha_j \omega^{-nm} | \phi_j \rangle_{AC}. \tag{48}$$

The probability of measuring the state associated with n' is given by $\tilde{P}_n = \frac{1}{P_m} \frac{1}{d} \sum_{j=0}^{d-1} |c_{j+m}|^2 |\alpha_j|^2$, and replacing Eq. (13), we get $\tilde{P}_n = \frac{1}{d}$.

As in the previous cases, it is possible to compute the fidelity of the copies in the $1 \rightarrow 2$ telecloning using Eqs. (16) and (48) to get first the density matrix of the $AC_1 C_2$ system, and then the reduced density matrix of one of the copies, namely C_2 , which is given by

$$\rho_{n,m}^{C_2} = \text{Tr}_{C_1} [\text{Tr}_A [| \psi_{n,m} \rangle_{A,C_1,C_2} \langle \psi_{n,m} |]]. \tag{49}$$

Hence, the fidelity of the copies, which is given by $F_m^{1 \rightarrow 2, \text{ME}} = \langle \psi | \rho_{n,m}^{C_2} | \psi \rangle$, has the explicit form

$$F_m^{1 \rightarrow 2, \text{ME}} = \frac{1}{2(d+1)} \frac{1}{d^3} + \frac{(2+d)}{2(d+1)} \frac{1}{P_m} \frac{1}{d^3} \left(\sum_{j=0}^{d-1} |\alpha_j|^2 c_{j+m} \right)^2. \tag{50}$$

This fidelity depends on the result of the first measurement, and this dependence is carried in the m indexes and in the input state, so that, after applying minimum error discrimination, the process is no longer universal, but still symmetric. Finally, the overall fidelity of the process is calculated by the weighted sum of all the possible measurements results, i.e. $\langle F^{\text{ME}} \rangle = \sum_{n,m=0}^{d-1} \tilde{P}_n P_m F_{nm}^2$, which is given explicitly by

$$\langle F^{\text{ME}} \rangle = \frac{1}{2(d+1)} \frac{1}{d^3} + \frac{(2+d)}{2(d+1)} \frac{1}{d^3} \sum_{m=0}^{d-1} \left(\sum_{j=0}^{d-1} |\alpha_j|^2 c_{j+m} \right)^2. \tag{51}$$

The last expressions correspond to $F^{1 \rightarrow 2, \text{PE}}/d^3$, which means that the overall fidelity of the process with minimum error discrimination is always worse than the process without any discrimination protocol. In the case of qubits, this fidelity is given by

$$\langle F^{\text{ME}} \rangle_{d=2}^{1 \rightarrow 2} = \frac{N_\theta^2}{96} \left[5|a|^4 c_1^2 + 5|b|^4 c_2^2 + |a|^2 |b|^2 (1 + 8c_1 c_2) \right]. \quad (52)$$

4.3 Telecloning combined with quantum state separation

The process of quantum state separation [16] consists of changing the separation between different states by a different and, in general, larger one. This separation is given by the inner product of the corresponding states. Therefore, the increase in this value between non-orthogonal states allows us to increase the chance of discrimination between them. This process is characterised by a unitary operation S and is a probabilistic process. In the case of a set of equidistant symmetric states, as in the case of the states given by Eq. (29), and considering that we want to symmetrically increase the separation of these states, the probability of this process is given by $p_{\text{SEP}} = c_{\text{min}}^2 / \tilde{c}_{\text{min}}^2$, where \tilde{c}_{min} is the minimal coefficient that we want to achieve. If we set the increase in the distance between the symmetric states on Eq. (29) so that after applying the suitable quantum state separation operation over system P they will be mutually orthogonal, it is then possible to discriminate and correctly choose the unitary reconstruction operations and at the end have optimal copies of the input state, i.e. with fidelity $F_{\text{opt}}^{1 \rightarrow 2}$, but with a limited probability $p_{\text{SEP}}^{\text{opt}} = c_{\text{min}}^2 / d$, which is always smaller than the probability achieved by unambiguous state discrimination.

However, it is also possible to change the separation of this state arbitrarily, not necessarily to orthogonality. If we apply this procedure over the system P , considering the states of Eq. (29), so that the new set of coefficients are denoted by \tilde{c}_k , and without other processes of discrimination, it is possible improve the fidelity of the copies, which in the case of $1 \rightarrow 2$ telecloning will be given by

$$F_m^{1 \rightarrow 2, \text{SEP}} = \frac{1}{2(d+1)} + \frac{(2+d)}{2(d+1)} \frac{1}{\tilde{P}_m} \left(\sum_{k=0}^{d-1} |\alpha_k|^2 \tilde{c}_{k+m} \right)^2, \quad (53)$$

where $\tilde{P}_m = \sum_{j=0}^{d-1} |\alpha_j|^2 |\tilde{c}_{j+m}|^2$. The probability of obtain this fidelity is given by p_{SEP} . The weighted fidelity over every possible outcome m is given by

$$\langle F^{\text{SEP}} \rangle^{1 \rightarrow 2} = \frac{1}{2(d+1)} + \frac{(2+d)}{2(d+1)} \sum_{m=0}^{d-1} \left(\sum_{k=0}^{d-1} |\alpha_k|^2 \tilde{c}_{k+m} \right)^2 \quad (54)$$

In the case of telecloning of qubits, this quantity reduces to

$$\langle F^{\text{SEP}} \rangle_{d=2}^{1 \rightarrow 2} = \frac{\tilde{N}_\theta^2}{12} \left[5|a|^4 \tilde{c}_1^2 + 5|b|^4 \tilde{c}_2^2 + |a|^2 |b|^2 (1 + 8\tilde{c}_1 \tilde{c}_2) \right] \quad (55)$$

where

$$\tilde{N}_\theta = \frac{\sqrt{2}}{\sqrt{a^2\tilde{c}_1^2 + b^2\tilde{c}_2^2}}. \tag{56}$$

If we use the quantum state separation process to orthogonalise in the qubit case, the fidelity achieved is the optimal 5/6 for any input state and the probability of success is $c_{\min}^2/2$, which is four times smaller than the probability of achieving the same fidelity when using unambiguous state discrimination.

5 Quantum channel of non-maximal Schmidt rank

In the case where $c_j = 0$, for some j , the states of Eq. (29) are linearly dependent and it is not possible to unambiguously distinguish between them, so it is necessary to resort to other discrimination schemes. A useful discrimination scheme in this scenario is maximum confidence state discrimination (MCD) [24]. This scheme allows us to discriminate between linearly dependent states, but with errors associated with the identification of some of them, even if we permit non-conclusive results in the discrimination measurements. Therefore, it is possible to construct a set of measurements which enables us to identify the states with the maximum possible confidence.

We apply, as in the USD case, the transformation of maximum confidence discrimination on the system X using the system P as ancilla system. The initial state of the ancilla system is $|m\rangle$. If this state remains unchanged, the discrimination process succeeds and it is possible to discriminate with maximum confidence. In any other case, the result is inconclusive. After applying the MCD unitary transformation U^{MC} , the state of the system is given by

$$\begin{aligned} |\tilde{\psi}\rangle = & \sqrt{1-p?} |m\rangle_X \sum_{n=0}^{d-1} |\tilde{u}_n\rangle_P U_{nm}^{-1} \sum_{j=0}^{d-1} \alpha_j |\phi_j\rangle_{AC} \\ & + \sqrt{p?} |m \oplus 1\rangle_X \sum_{n=0}^{d-1} |\tilde{\chi}_n\rangle_P U_{nm}^{-1} \sum_{j=0}^{d-1} \alpha_j |\phi_j\rangle_{AC}, \end{aligned} \tag{57}$$

where $|\tilde{\psi}\rangle = U_{XP}^{MC} |\psi_m\rangle_{XPAC}$. The states $|\tilde{u}\rangle_n$ are orthonormal states given by

$$|\tilde{u}_n\rangle = \frac{1}{\sqrt{N}} \sum_{k=0}^{N-1} e^{2\pi kn/d} |k\rangle_2, \tag{58}$$

and the states $|\tilde{\chi}\rangle$ are normalised non-orthogonal states, given by

$$|\tilde{\chi}_n\rangle = \sum_{k=0}^{N-1} \sqrt{\frac{c_k^2 - c_{\min}^2}{p?}} e^{2\pi kn/d} |k\rangle. \tag{59}$$

N is the number of non-vanishing coefficients c_j . As in the USD case, if we get the outcome m by measuring the system X , we will be able to perform an inverse Fourier transform and measure the system P on the computational basis, obtaining the outcome n , and then perform the transformations of Eqs. (34) and (35). The confidence of the process, i.e. the probability of correctly distinguishing the state n through the measurement outcome n , is N/d . Additionally, the probability of getting a non-conclusive outcome is $p? = 1 - Nc_{\min}^2$, where in this case c_{\min} is the non-vanishing coefficient with the smallest modulus.

6 Summary

We presented a general scheme to probabilistically teleclone qudit states via partially pure entangled channels as described in Eq. (9). After introducing these kinds of channels, the process is still symmetric, but no longer universal. However, it is still possible to recover the universality of the cloning process in the case of maximally ranked Schmidt channels, by using unambiguous state discrimination and turning the process into a probabilistic one. The same result could be achieved using quantum state separation, but with a much smaller probability. If instead of using these processes we use minimum error discrimination to try to improve the results, the overall fidelities are even worse than not using any strategy. In [25], the authors propose a similar probabilistic scheme, but just in the $1 \rightarrow 2$ case and maximal Schmidt rank using a *pseudo-control unitary operation* which corresponds to a special case of our unambiguous state discrimination strategy.

Acknowledgments Open access funding provided by University of Innsbruck and Medical University of Innsbruck. This work was supported by Millenium Scientific Initiative Grant No. RC130001. N.C. and G.A. acknowledge support from CONICYT. We thank M. W. van Mourik and M. G. Higgins for their valuable contribution to the improvement of the manuscript.

Open Access This article is distributed under the terms of the Creative Commons Attribution 4.0 International License (<http://creativecommons.org/licenses/by/4.0/>), which permits unrestricted use, distribution, and reproduction in any medium, provided you give appropriate credit to the original author(s) and the source, provide a link to the Creative Commons license, and indicate if changes were made.

References

1. Bennett, C.H., Brassard, G., Crepeau, C., Jozsa, R., Peres, A., Wootters, W.K.: Teleporting an unknown quantum state via dual classical and Einstein–Podolsky–Rosen channels. *Phys. Rev. Lett.* **70**, 1895 (1993)
2. Vaidman, L.: Teleportation of quantum states. *Phys. Rev. A* **49**, 1473 (1994)
3. Wootters, W.K., Zurek, W.H.: A single quantum cannot be cloned. *Nature* **299**, 802 (1982)
4. Bužek, V., Hillery, M.: Quantum copying: beyond the no-cloning theorem. *Phys. Rev. A* **54**, 1844 (1996)
5. Werner, R.F.: Optimal cloning of pure states. *Phys. Rev. A* **58**, 1827 (1998)
6. Keyl, M., Werner, R.F.: Optimal cloning of pure states, testing single clones. *J. Math. Phys.* **40**, 3283 (1999)
7. Alber, G., Delgado, A., Jex, I.: Optimal universal two-particle entanglement processes in arbitrary dimensional Hilbert spaces. *Quantum Inf. Comput.* **1**, 33 (2001)

8. Duan, L.M., Guo, G.C.: Probabilistic cloning and identification of linearly independent quantum states. *Phys. Rev. Lett.* **80**, 4999 (1998)
9. Jiménez, O., Bergou, J., Delgado, A.: Probabilistic cloning of three symmetric states. *Phys. Rev. A* **82**, 062307 (2010)
10. Jiménez, O., Roa, L., Delgado, A.: Probabilistic cloning of equidistant states. *Phys. Rev. A* **82**, 022328 (2010)
11. Murao, M., Jonathan, D., Plenio, M., Vedral, V.: Quantum telecloning and multiparticle entanglement. *Phys. Rev. A* **59**, 156 (1999)
12. Murao, M., Plenio, M.B., Vedral, V.: Quantum-information distribution via entanglement. *Phys. Rev. A* **61**, 032311 (2000)
13. Chiuri, A., Greganti, C., Paternostro, M., Vallone, G., Mataloni, P.: Experimental quantum networking protocols via four-qubit hyperentangled dicke states. *Phys. Rev. Lett.* **109**, 173604 (2012)
14. Chefles, A.: Quantum state discrimination. *Contemp. Phys.* **41**, 401 (2000)
15. Barnett, S.M., Croke, S.: Quantum state discrimination. *Adv. Opt. Photonics* **1**, 238 (2009)
16. Chefles, A., Barnett, S.: Quantum state separation, unambiguous discrimination and exact cloning. *J. Phys. A: Math. Gen.* **31**, 10097 (1998)
17. Helstrom, C.W.: *Quantum Detection and Estimation Theory*. Academic Press, New York (1976)
18. Holevo, A.S.: Statistical decision theory for quantum systems. *J. Multivar. Anal.* **3**, 337 (1973)
19. Yuen, H.P., Kennedy, R.S., Lax, M.: Optimum testing of multiple hypotheses in quantum detection theory. *IEEE Trans. Inf. Theory* **21**, 125 (1975)
20. Chefles, A.: Unambiguous discrimination between linearly independent quantum states. *Phys. Lett. A* **239**, 339 (1998)
21. Chefles, A., Barnett, S.M.: Optimum unambiguous discrimination between linearly independent symmetric states. *Phys. Lett. A* **250**, 223 (1998)
22. Roa, L., Delgado, A., Fuentes-Guridi, I.: Optimal conclusive teleportation of quantum states. *Phys. Rev. A* **68**, 022310 (2003)
23. Bae, J.: Structure of minimum-error quantum state discrimination. *New J. Phys.* **15**, 073037 (2015)
24. Croke, S., Andersson, E., Barnett, S., Gilson, C.R., Jeffers, J.: Maximum confidence quantum measurements. *Phys. Rev. Lett.* **96**, 070401 (2006)
25. Li, D., Shi, Z.: Probabilistic telecloning with partially entangled states. *Open Syst. Inf. Dyn.* **16**, 413 (2009)