



CO($J = 1-0$) Mapping Survey of 64 Galaxies in the Fornax Cluster with the ALMA Morita Array

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Abstract

We conduct a $^{12}\text{C}^{16}\text{O}(J = 1-0)$ (hereafter CO) mapping survey of 64 galaxies in the Fornax cluster using the Atacama Large Millimeter/submillimeter Array Morita array in cycle 5. CO emission is detected from 23 out of the 64 galaxies. Our sample includes dwarf, spiral, and elliptical galaxies with stellar masses of $M_{\text{star}} \sim 10^{6.3-11.6} M_{\odot}$. The achieved beam size and sensitivity are $15'' \times 8''$ and $\sim 12 \text{ mJy beam}^{-1}$ at the velocity resolution of $\sim 10 \text{ km s}^{-1}$, respectively. We study the cold gas (molecular and atomic gas) properties of 38 subsamples with $M_{\text{star}} > 10^9 M_{\odot}$ combined with literature H I data. We find that (1) the low star formation (SF) activity in the Fornax galaxies is caused by the decrease in the cold gas mass fraction with respect to stellar mass (hereafter, gas fraction) rather than the decrease of the SF efficiency from the cold gas; (2) the atomic gas fraction is more heavily reduced than the molecular gas fraction of such galaxies with low SF activity. A comparison between the cold gas properties of the Fornax galaxies and their environmental properties suggests that the atomic gas is stripped tidally and by the ram pressure, which leads to the molecular gas depletion with an aid of the strangulation and consequently SF quenching. Preprocesses in the group environment would also play a role in reducing cold gas reservoirs in some Fornax galaxies.

Unified Astronomy Thesaurus concepts: [Galaxy environments \(2029\)](#); [Galaxy clusters \(584\)](#); [Molecular gas \(1073\)](#)

1. Introduction

The cosmic star formation rate (SFR) density decreases in the last ~ 10 Gyr (e.g., Madau & Dickinson 2014). Studying how the star formation (SF) in galaxies has been quenched is essential to understand galaxy evolution. In particular, the SF

quenching due to the galaxy environment becomes important in the current universe since the large-scale structure of the universe develops with time (e.g., Peng et al. 2010; Darvish et al. 2016). Indeed, various observations find that the galaxies in dense regions are gas poor and passive (e.g., Dressler 1980; Gómez et al. 2003; Balogh et al. 2004; Hogg et al. 2004; Kauffmann et al. 2004; Baldry et al. 2006).

The environmental effects on the SF in the galaxies can be roughly classified into two categories: gravitational interactions and hydrodynamical interactions. One of the major differences between the two categories is that gravitational interactions act in the same way on all components of a galaxy (dark matter,

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stars, and gas), while hydrodynamic interactions only affect the gas component in the galaxy (circumgalactic medium and interstellar medium [ISM]). The former includes the galaxy–galaxy interaction (Merritt 1983) and the galaxy–cluster interaction (Byrd & Valtonen 1990). The multiple high-speed galaxy encounter is specifically called “galaxy harassment” (Moore et al. 1996, 1999), which is expected to be more important in the core than on the outskirts of the cluster. Although the slow galaxy–galaxy encounter is rare in the cluster core, it is important in the group environment. The hydrodynamical interactions include the ram pressure stripping (RPS) of ISM in galaxies by the intracluster medium (ICM; Gunn & Gott 1972) and the strangulation (Larson et al. 1980). The typical timescale of the ISM RPS is $\lesssim 500$ Myr (Abadi et al. 1999; Vollmer et al. 2001; Roediger & Brüggén 2006). In the process of the strangulation, the hot coronal gas of galaxies is stripped by the ram pressure and/or tidally but the SF continues until the remaining ISM is consumed. The typical timescale of the strangulation is estimated to be ~ 4 Gyr from the stellar metallicity difference between quiescent and star-forming galaxies (Peng et al. 2015).

It is essential to investigate the molecular gas (H_2) properties of cluster galaxies to understand the quenching processes in the galaxy clusters. This is because the H_2 molecule is the dominant molecular species in the universe and is the raw material for SF. However the H_2 molecule does not have a dipole moment and does not emit under cold environment with a temperature of ~ 10 K where stars generally form in the local universe. Astronomers have made use of $^{12}\text{C}^{16}\text{O}(J = 1 - 0)$ (hereafter CO) as a tracer of molecular gas since it is the most abundant molecule after H_2 in galaxies.

The Fornax cluster is the second nearest cluster from the Milky Way at a distance of 20 Mpc (Blakeslee et al. 2009), after the Virgo cluster at 16.5 Mpc (Mei et al. 2007). The total mass enclosed at the virial radius of 0.7 Mpc is estimated to be $5 \times 10^{13} M_\odot$ (Drinkwater et al. 2001). From the X-ray observation of its ICM, the Fornax cluster is a cool-core cluster, in which the cooling timescale at the cluster core is less than 1 Gyr, while the Fornax cluster is considered to be in the process of forming a cool core (i.e., “nascent” core, Burns et al. 2008) unlike the typical cool-core clusters that generally have a well-established core (Hudson et al. 2010). The member galaxies have also been studied in various projects in mostly optical wavelengths such as the Fornax Cluster Catalog (FCC; Ferguson 1989a), the Fornax spectroscopic survey (Drinkwater et al. 2000), the ACS Fornax Cluster Survey (Jordán et al. 2007), Next Generation Fornax Survey (Muñoz et al. 2015), the Fornax Deep Survey (Venhola et al. 2018), the Fornax3D project (Sarzi et al. 2018), and the Fornax Cluster VLT Spectroscopic Survey (Pota et al. 2018). Furthermore, they are observed in infrared (the Herschel Fornax Cluster Survey, Davies et al. 2013), H I, and radio continuum (Serra et al. 2019; Maccagni et al. 2020; Kleiner et al. 2021; Loni et al. 2021).

CO emission has also been searched for in the Fornax galaxies. Horellou et al. (1995) conducted CO single/multiple pointing observations toward 21 Fornax spiral and lenticular galaxies with the 15 m Swedish-ESO Submillimeter Telescope, whose beam size was $44''$, which corresponds to 4.3 kpc at 20 Mpc. They detected CO emission from 11 galaxies and found that CO emission of the Fornax galaxies is ~ 10 times lower than that of controlled field galaxies with similar far-infrared luminosity. They considered that

CO deficiency in the Fornax galaxies is attributed to tidal interactions rather than RPS provided by the weak X-ray emission and the high number density of galaxies in the cluster. Recently, the CO deficiency in the Fornax galaxies has been confirmed by the ALMA Fornax Cluster Survey (AlFoCS) project (Zabel et al. 2019). It observed 30 galaxies with stellar mass of $10^{8.5-11} M_\odot$ only with the main array of the Atacama Large Millimeter/submillimeter Array (ALMA) and detected CO emission from 15 galaxies. Among the 30 galaxies, 10 galaxies were also observed with the Mopra telescope. With the high-resolution ALMA observations ($2''-3'' \sim 0.2-0.3$ kpc at 20 Mpc), they also revealed that the CO-deficient galaxies with low stellar masses of $< 3 \times 10^9 M_\odot$ show disturbed CO morphologies. This suggests that the molecular gas in such less massive galaxies could be affected by the ram pressure from the ICM in the Fornax cluster.

We conduct CO mapping observations toward the 64 galaxies in the Fornax cluster using the Morita array (i.e., the Atacama Compact Array [ACA]) of the ALMA to measure the spatial distribution and the total flux in CO of the galaxies, making it by far the largest CO survey of the Fornax galaxies so far. With the data of one of the sample galaxies, NGC 1316 (i.e., Fornax A), obtained in the survey, we have already presented the complex distribution and kinematics of molecular gas (Morokuma-Matsui et al. 2019), the recurrent nuclear activity (Maccagni et al. 2020), and the physical properties of the multiphase gas in NGC 1316 (Maccagni et al. 2021). The cold gas properties of the Fornax A group have been presented in Kleiner et al. (2021) as well.

The structure of this paper is as follows. We briefly introduce the ALMA observations and data in Section 2. The ancillary data (stellar mass, SFR, and atomic gas mass) used in this study are described in Section 3. With the data set, we present the galaxy-integrated cold (molecular and atomic) gas properties of the Fornax galaxies as a function of the clustocentric distance, local galaxy number density, and accretion phase to the cluster that is defined on the projected phase-space diagram (PSD) in Section 4. Since the various properties of galaxies depend on their stellar masses and the mass segregation is often observed in galaxy clusters, we subtract the mass dependence of the quantities explored in this study before the investigation. We compare our results with previous studies and discuss the possible dominant quenching process in the Fornax cluster in Section 5. Finally, we summarize our study in Section 6.

We adopt the following parameters of the Fornax cluster throughout the paper: the coordinates of the cluster center (R.A., decl.) = (54°6, −35°5), the distance of 20 Mpc (Blakeslee et al. 2009), the virial radius of $2^\circ 0$ (~ 0.7 Mpc at 20 Mpc; Drinkwater et al. 2001), the velocity dispersion of 318 km s^{-1} , and the systemic velocity of 1442 km s^{-1} (Maddox et al. 2019).

2. ALMA Observations and Data

2.1. Sample

Our 64 sample galaxies meet either of the following conditions: (1) galaxies observed in the AlFoCS project (Zabel et al. 2019, 2020) or (2) galaxies in the vicinity of the Fornax cluster with archival H I data in the HyperLEDA (Makarov et al. 2014). The galaxies in condition (1) are selected based on the FCC (Ferguson 1989b) with stellar masses of $> 3 \times 10^8 M_\odot$ and Herschel detection (Fuller et al. 2014) or H I detection down to $\sim 3 \times 10^7 M_\odot$ (Waugh et al. 2002). The FCC number, coordinates, recessional kinematic LSR velocity, and morphology of the sample galaxies are presented in Table 1. The

Table 1
Basic Properties of Sample Fornax Galaxies

Name	FCC No.	R.A. (deg)	Decl. (deg)	V (kLSR) (km s ⁻¹)	Morphology
ESO 302-14	...	57.920417	-38.452222	854	IB(s)m
ESO 302-9	322	56.892167	-38.576528	970	SB(s)dm?
ESO 357-25	26	50.90525	-35.778389	1806	SAB0 ⁺ 0?
ESO 358-16	115	53.288327	-35.718527	1700	Sc
ESO 358-20	139	53.738958	-32.639833	1752	IB(s)m? pec
ESO 358-5	53	51.819208	-33.486361	1611	SAB(s)m pec?
ESO 358-51	263	55.385833	-34.888333	1707	S0/a?
ESO 358-G015	113	53.278542	-34.808111	1371	SBm? pec
ESO 358-G063	312	56.579208	-34.943556	1911	I0?
ESO 359-2	335	57.653042	-35.909333	1412	SB0 ⁺ -?
ESO 359-3	338	58.003833	-33.467639	1557	Sab? edge-on
ESO 359-5	...	58.514244	-36.063653	1370	dwarf
FCC 102	102	53.04474	-36.220809	1706	Irr
FCC 117	117	53.31106	-37.819638	1500	...
FCC 120	120	53.392593	-36.605914	887	S
FCC 177	177	54.197875	-34.739611	1544	S0 ⁺ 0? edge-on
FCC 198	198	54.427823	-37.208339	1500	...
FCC 206	206	54.556202	-37.29018	1385	dwarf
FCC 207	207	54.580292	-35.129083	1403	S0
FCC 261	261	55.339708	-33.769222	1475	...
FCC 302	302	56.300593	-35.570906	786	IB(s)m? edge-on
FCC 306	306	56.439158	-36.34653	868	S
FCC 316	316	56.756333	-36.437472	1529	dwarf
FCC 32	32	51.2185	-35.435444	1301	S0
FCC 332	332	57.45425	-35.945583	1308	S0
FCC 44	44	51.531017	-35.127491	1232	...
IC 1993	315	56.770042	-33.709861	1062	(R')SAB(rs)b
IC 335	153	53.879333	-34.447056	1602	S0 edge-on
MCG-06-08-024	90	52.784417	-36.290139	1796	SA0 ⁺ -?
MCG-06-09-023	282	55.688833	-33.9205	1208	S0
NGC 1310	...	50.264292	-37.101694	1788	SA(s)c?
NGC 1316	21	50.673825	-37.208227	1743	SAB0 ⁺ 0(s) pec
NGC 1316C	33	51.243583	-37.009472	1783	(R')SA0 ⁺ 0?
NGC 1317	22	50.684527	-37.103688	1924	SAB(r)a
NGC 1326	29	50.985	-36.464667	1343	(R)SB0 ⁺ (+r)
NGC 1326A	37	51.285405	-36.363949	1814	SB(s)m?
NGC 1326B	39	51.334782	-36.384954	982	SB(s)m? edge-on
NGC 1336	47	51.634125	-35.713556	1401	SA0 ⁺ -
NGC 1339	63	52.027417	-32.286111	1375	cD pec?
NGC 1341	62	51.993417	-37.15	1859	SAB(s)ab
NGC 1350	88	52.783833	-33.628639	1888	(R')SB(r)ab
NGC 1351	83	52.64575	-34.853944	1497	SA0 ⁺ -pec?
NGC 1351A	67	52.203	-35.178139	1336	SB(rs)bc? edge-on
NGC 1365	121	53.401548	-36.140402	1619	SB(s)b
NGC 1374	147	53.819125	-35.22625	1277	E
NGC 1375	148	53.820083	-35.265667	723	S0
NGC 1379	161	54.016458	-35.441194	1307	E
NGC 1380	167	54.114968	-34.976225	1860	SA0
NGC 1381	170	54.132	-35.295194	1707	SA0? edge-on
NGC 1386	179	54.192425	-35.999408	851	SB0 ⁺ (+s)
NGC 1387	184	54.23775	-35.506639	1285	SAB0 ⁺ -(s)
NGC 1399	213	54.620941	-35.450657	1408	E1 pec
NGC 1404	219	54.716333	-35.594389	1930	E2
NGC 1406	...	54.847083	-31.321417	1058	SB(s)bc? edge-on
NGC 1419	249	55.175458	-37.510833	1576	dE3
NGC 1425	...	55.547792	-29.893333	1493	SA(s)b
NGC 1427	276	55.580933	-35.392563	1371	E4
NGC 1427A	235	55.03875	-35.624444	2011	dE2
NGC 1436	290	55.9045	-35.853028	1370	ScII
NGC 1437A	285	55.75914	-36.273374	868	SdIII?
NGC 1437B	308	56.478542	-36.356972	1480	Sd(onedge)
NGC 1484	...	58.583875	-36.968889	1022	SB(s)b?
PGC 012625	...	50.52	-37.5875	1609	IB(s)m?
WOMBAT I	...	55.245185	-38.855266	812	S

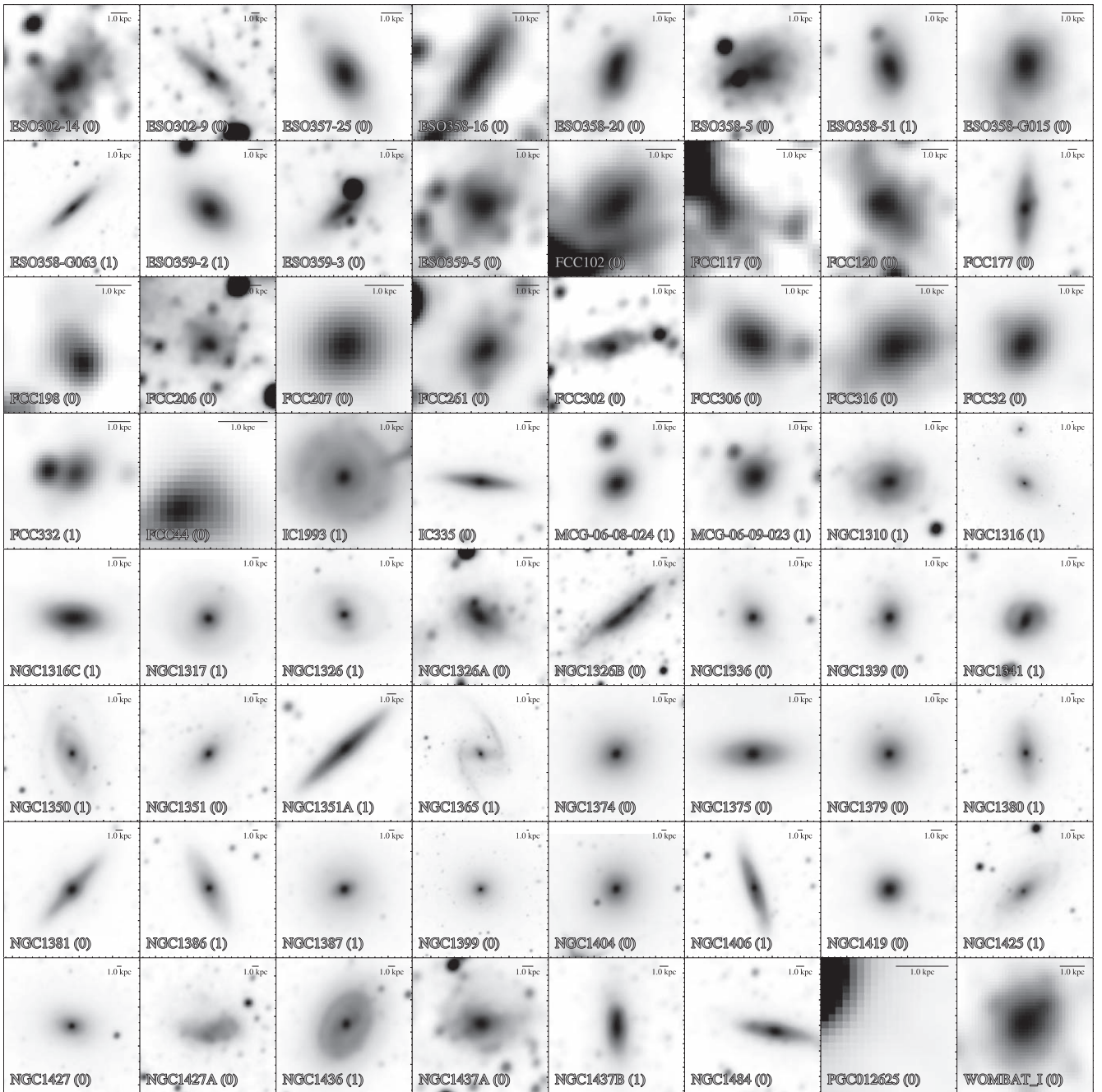


Figure 1. WISE $3.4 \mu\text{m}$ images of our sample galaxies (north is up and east is to the left). A bar on the upper right corner in each panel indicates 1 kpc. The integer in the brackets after the galaxy name indicates the detection (1) or nondetection (0) of the ALMA CO observation. FCC 117 and PGC 012625 are not clearly detected in the WISE band.

$3.4 \mu\text{m}$ images of our sample galaxies obtained with the Wide-field Infrared Survey Explorer (WISE; Wright et al. 2010) are shown in Figure 1. The sky distribution of our sample galaxies in addition to all of the FCC galaxies is presented in the left panel of Figure 2. We can see that our sample galaxies are evenly distributed over the region where the FCC galaxies are distributed. The relationship between stellar mass and SFR is presented in the right panel of Figure 2 (see Section 3.1 for the derivation of stellar mass and SFR). Most sample galaxies are on or below the main sequence of star-forming galaxies.

2.2. ALMA 7 m and Total Power Observations

ALMA CO observations with the 7 m (7M) and 12 m total power (TP) antennas arrays of the 64 Fornax galaxies were performed during cycle 5 (project code of 2017.1.00129.S and P.I. Kana Morokuma). The mapping area for the galaxies is basically determined so that the optical disk of galaxies is covered.

The 7M data were obtained from 2017 October 16 to December 21 for the scheduling block (SB) of NGC1316_a_03_7M and from 2017 November 30 to December 28 for the SB of ESO358-G_a_03_7M with 9–11 antennas

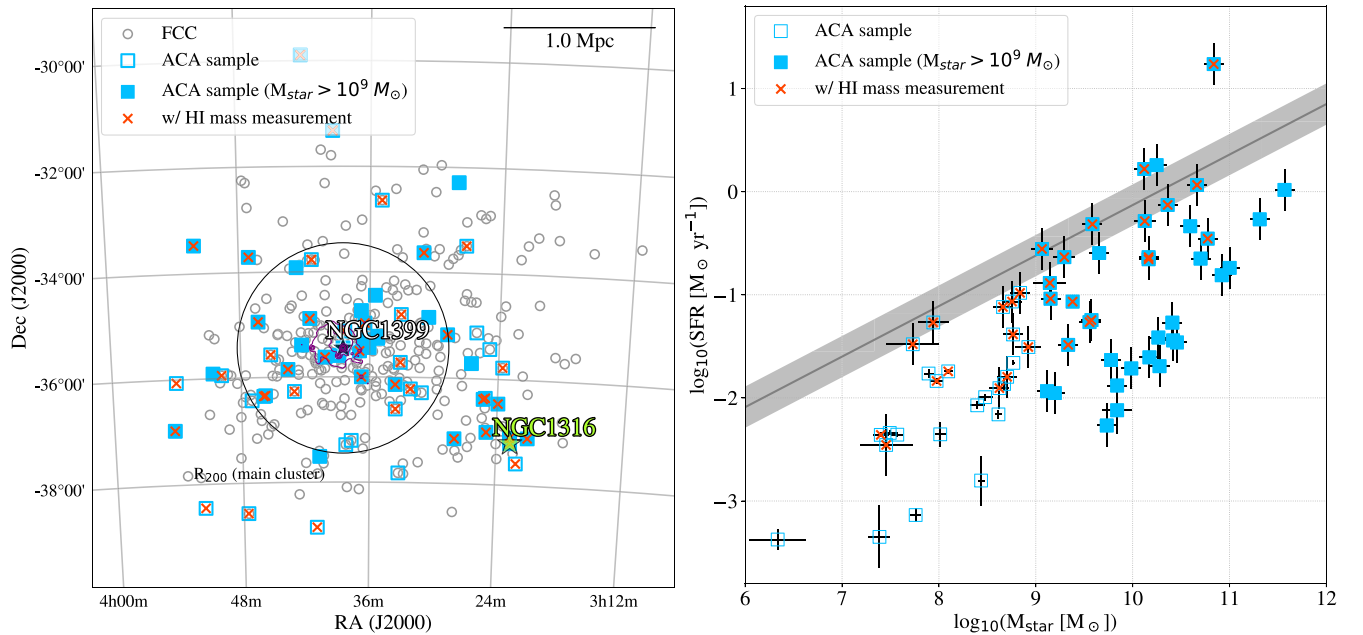


Figure 2. (Left) Sky distributions of Fornax cluster galaxies. Our entire ACA sample and those with H I mass measurements (including upper limits) are indicated with blue open squares or orange X marks, respectively. Those with $M_{\text{star}} > 10^9 M_{\odot}$ are indicated with blue filled squares. The gray open circles indicate the FCC galaxies (Ferguson 1989a). The purple contours indicate the X-ray distribution obtained with the Suzaku and XMM-Newton data (Murakami et al. 2011). Our ACA sample is distributed over the region where the FCC galaxies are distributed. (Right) Stellar mass vs. SFR relation of our sample galaxies. The gray solid line and gray shaded region indicate the main sequence of star-forming galaxies defined in Speagle et al. (2014) and its standard deviation, respectively.

Table 2
Observations for the 7M Data

ID	SB Name	Obs. Date in 2017	No. of EBs	Flux/Bandpass Cal.	Complex Gain Cal.
1	NGC1316_a_03_7M	Oct 16–Dec 21	33	J0006-0623, J0522-3627	J0334-4008, J0403-3605, J0424-3756
2	ESO358-G_a_03_7M	Nov 30–Dec 28	33	J0006-0623, J0522-3627	J0334-4008, J0403-3605, J0406-3826

whose baselines range 8.9–48.9 m. The 64 galaxies were divided into the two SBs, and each SB consists of 33 execution blocks (EBs). The maximum recoverable scales are $40''7\text{--}58''3$ ($\sim 3.9\text{--}5.7$ kpc at 20 Mpc) for NGC1316_a_03_7M and $55''2$ (~ 5.4 kpc) for ESO358-G_a_03_7M, respectively. The flux, bandpass, and complex gain calibrators used in the two SBs are presented in Table 2. The systemic velocity for each source was fixed to 1500 km s^{-1} to cover the entire-sample galaxies with a single correlator setup. The CO line was observed in one of the upper-sideband spectral windows (SPWs) with bandwidth and resolution of 1875 MHz (4902.5 km s^{-1}) and 1.128 MHz (2.952 km s^{-1}), respectively.

The on-the-fly observations with the TP array were carried out to account for zero spacing information between March and 2017 September with three or four antennas. There are 64 SBs, and each SB consists of one to 13 EBs (13 EBs for NGC 1365). The mapping area was set to cover the entire area observed with the 7M array. The correlator and SPW setup was the same as for the 7M observations.

2.3. ALMA Data Analysis with CASA

The obtained ALMA data were calibrated and imaged with the standard ALMA data analysis package, the Common Astronomy Software Applications (CASA; McMullin et al. 2007; Petry & CASA Development Team 2012). The absolute flux and gain fluctuations of the 7M data were calibrated with the ALMA Science Pipeline (version r40896 of Pipeline-

CASA51-P2-B) in the CASA 5.1.1 package. The resultant fluxes of the phase calibrator at each SPW are mostly consistent within the errors with the values reported in the other ALMA measurements of the same calibrator on the closest date to our observations. The 7M CO mosaic data cube for each source was generated with the TCLEAN task in CASA version 5.5 with standard options as follows: Briggs weighting with a robust parameter of 0.5, the auto-multithresh mask with standard value for 7M data provided in the CASA guides for the automasking, and niter of 10,000. The TP data were also reduced with the pipeline version r40896 of Pipeline-CASA51-P2-B in the CASA 5.1.1 package. The 7M and TP data were combined with the FEATHER task in CASA. The synthesized beam is roughly $15'' \times 8''$, and the values for each galaxy are presented in Table 3. The velocity width of the final fits cube is set to $\sim 10 \text{ km s}^{-1}$. The achieved sensitivity is 12 mJy beam^{-1} .

2.4. Molecular Gas Mass

CO emission is detected in 23 galaxies out of the 64 sample galaxies. Through inspecting the output masks, we adjust and decide an optimized set of configuration parameters of the Source Finding Application (SoFiA; Serra et al. 2015; Westmeier et al. 2021). We use the smooth and clip algorithm, which smoothes the cube with Gaussian kernels of different sizes, select voxels with a 5σ threshold in each smoothed cube as well as in the original cube, and then use the subset of voxels

Table 3
Parameters for ALMA Data Analysis

Name	SB ID ^a	$b_{\text{maj}}; b_{\text{min}}; b_{\text{P.A.}}$ (", ", °)	Region Shape ^b	Center ^c (deg, deg)	Width/ a^d (arcseconds)	Height/ b^e (arcseconds)	P.A. ^f (deg)	Velocity Range ^g (km s ⁻¹)
ESO 302-14	1	14.5; 7.7; 87.1	Ellipse	57.920417, -38.452222	35.0	17.2	67.9	701 - 1001
ESO 302-9	1	14.8; 7.7; 87.8	Ellipse	56.892167, -38.576528	70.5	27.8	128.6	821 - 1120
ESO 357-25	1	14.6; 7.7; 87.4	Ellipse	50.90525, -35.778389	27.9	10.4	117.2	1650 - 1950
ESO 358-16	2	14.4; 7.6; 89.1	Ellipse	53.288327, -35.718527	21.1	8.4	66.2	1540 - 1840
ESO 358-20	1	14.7; 7.7; 86.8	Ellipse	53.738958, -32.639833	40.4	22.6	55.0	1590 - 1890
ESO 358-5	1	14.8; 7.6; 87.5	Ellipse	51.819208, -33.486361	43.4	31.3	12.5	1450 - 1750
ESO 358-51	1 + 2	14.3; 7.6; 88.6	Rectangle	55.385833, -34.88972188888889	30.0	30.0	0.0	1650 - 1740
ESO 358-G015	2	14.4; 7.6; 88.9	Ellipse	53.278542, -34.808111	27.1	17.6	84.4	1210 - 1510
ESO 358-G063	2	14.4; 7.6; 88.9	Rectangle	56.579208, -34.943556	140.0	35.0	45.0	1750 - 2060
ESO 359-2	2	14.2; 7.6; 88.1	Rectangle	57.65331977777777, -35.909333	18.0	14.0	0.0	1410 - 1480
ESO 359-3	2	14.6; 7.7; 85.8	Ellipse	58.003833, -33.467639	56.0	29.1	44.0	1400 - 1700
ESO 359-5	2	14.6; 7.6; 87.6	Ellipse	58.514244, -36.063653	27.0	17.3	45.3	1210 - 1510
FCC 102	2	14.7; 7.7; -89.1	Ellipse	53.04474, -36.220809	19.1	11.7	59.5	1550 - 1850
FCC 117	2	14.6; 7.8; 88.5	Ellipse	53.31106, -37.819638	6.5	6.5	45.0	1340 - 1640
FCC 120	2	14.6; 7.7; -89.5	Ellipse	53.392593, -36.605914	21.4	8.2	132.9	731 - 1031
FCC 177	2	14.6; 7.7; -89.2	Ellipse	54.197875, -34.739611	63.2	35.7	82.1	1390 - 1690
FCC 198	2	14.5; 7.7; 88.5	Ellipse	54.427823, -37.208339	6.5	6.5	45.0	1340 - 1640
FCC 206	2	14.5; 7.7; 88.5	Ellipse	54.556202, -37.29018	49.2	26.2	-24.3	1230 - 1530
FCC 207	2	14.4; 7.6; 89.0	Ellipse	54.580292, -35.129083	15.0	11.5	-17.6	1250 - 1550
FCC 261	2	14.4; 7.6; 88.5	Ellipse	55.339708, -33.769222	64.2	22.2	49.2	1320 - 1620
FCC 302	2	14.4; 7.6; 88.8	Ellipse	56.300593, -35.570906	46.3	9.7	13.1	632 - 931
FCC 306	2	14.5; 7.8; 88.0	Ellipse	56.439158, -36.34653	18.8	13.9	-35.4	711 - 1011
FCC 316	2	14.3; 7.6; 88.6	Ellipse	56.756333, -36.437472	17.7	10.0	35.0	1370 - 1670
FCC 32	2	14.7; 7.7; -88.9	Ellipse	51.2185, -35.435444	22.0	18.9	72.8	1150 - 1450
FCC 332	2	14.3; 7.6; 88.8	Rectangle	57.45494444444444, -35.94669411111111	15.0	12.0	0.0	1260 - 1330
FCC 44	2	14.7; 7.7; -89.1	Ellipse	51.531017, -35.127491	11.8	8.3	33.6	1080 - 1380
IC 1993	1	14.8; 7.6; 88.0	Rectangle	56.770042, -33.71124988888889	115.0	115.0	0.0	1001 - 1130
IC 335	1	14.7; 7.6; 87.2	Ellipse	53.879333, -34.447056	79.7	35.9	4.2	1440 - 1740
MCG-06-08-024	2	14.6; 7.7; 88.4	Rectangle	52.78469477777777, -36.28986122222223	16.0	18.0	0.0	1780 - 1840
MCG-06-09-023	2	14.3; 7.6; 87.9	Rectangle	55.690083, -33.92147222222222	17.0	17.0	0.0	1210 - 1320
NGC 1310	1	14.8; 7.7; 88.0	Rectangle	50.264292, -37.101694	80.0	55.0	0.0	1700 - 1860
NGC 1316	1	15.2; 7.7; -89.3	Rectangle	50.6727, -37.2045	80.0	105.0	0.0	1410 - 2020
NGC 1316C	1	14.6; 7.7; 87.4	Rectangle	51.243583, -37.009472	45.0	20.0	0.0	1880 - 2030
NGC 1317	1	14.9; 7.6; 89.2	Rectangle	50.684527, -37.10313244444444	70.0	60.0	0.0	1820 - 2040
NGC 1326	1	14.9; 7.7; 89.2	Rectangle	50.985, -36.46605588888889	50.0	60.0	0.0	1190 - 1490
NGC 1326A	1	14.7; 7.7; 88.0	Ellipse	51.285405, -36.363949	62.3	27.7	12.7	1650 - 1950
NGC 1326B	1	14.9; 7.7; 88.9	Ellipse	51.334782, -36.384954	91.9	29.8	32.7	831 - 1130
NGC 1336	1	14.6; 7.7; 87.5	Ellipse	51.634125, -35.713556	76.2	52.5	109.2	1240 - 1540
NGC 1339	1	14.8; 7.6; 87.1	Ellipse	52.027417, -32.286111	75.2	59.2	49.9	1220 - 1520
NGC 1341	1	14.7; 7.7; 87.9	Rectangle	51.993417, -37.14861111111111	65.0	55.0	0.0	1810 - 1890
NGC 1350	1	15.4; 7.7; -89.2	Rectangle	52.783833, -33.628639	100.0	200.0	0.0	1670 - 2100
NGC 1351	1	14.8; 7.6; 87.7	Ellipse	52.64575, -34.853944	110.1	77.8	53.3	1340 - 1640
NGC 1351A	2	14.7; 7.6; -89.4	Rectangle	52.20355555555556, -35.17980566666667	100.0	30.0	45.0	1150 - 1500
NGC 1365	2	15.5; 7.6; -88.6	Rectangle	53.401548, -36.140402	320.0	450.0	10.0	1370 - 1880
NGC 1374	1	14.8; 7.6; 88.1	Ellipse	53.819125, -35.22625	78.3	74.5	63.0	1120 - 1420
NGC 1375	1	14.6; 7.6; 87.1	Ellipse	53.820083, -35.265667	54.1	27.0	0.2	572 - 871
NGC 1379	1	14.6; 7.6; 87.6	Ellipse	54.016458, -35.441194	90.3	82.7	-4.2	1150 - 1450
NGC 1380	2	14.7; 7.6; -89.3	Rectangle	54.11566244444444, -34.97678055555556	20.0	14.0	0.0	1550 - 2141
NGC 1381	1	14.7; 7.6; 87.5	Ellipse	54.132, -35.295194	81.8	41.3	51.8	1550 - 1850
NGC 1386	2	14.6; 7.7; 88.8	Rectangle	54.19298055555556, -35.99996355555556	40.0	70.0	25.0	632 - 1080
NGC 1387	2	14.4; 7.6; 89.2	Rectangle	54.23775, -35.506639	50.0	40.0	0.0	1190 - 1350
NGC 1399	1	14.8; 7.6; 88.5	Ellipse	54.620941, -35.450657	271.1	225.9	-15.5	1250 - 1550
NGC 1404	1	14.7; 7.6; 87.6	Ellipse	54.716333, -35.594389	115.0	88.6	72.4	1770 - 2070
NGC 1406	1	14.7; 7.6; 86.7	Rectangle	54.84791633333333, -31.32252811111111	50.0	140.0	20.0	861 - 1220
NGC 1419	1	14.5; 7.6; 86.8	Ellipse	55.175458, -37.510833	53.0	43.3	-7.4	1420 - 1720

Table 3
(Continued)

Name	SB ID ^a	$b_{\text{maj}}; b_{\text{min}}; b_{\text{P.A.}}$ ($''; ''; ^\circ$)	Region Shape ^b	Center ^c (deg, deg)	Width/ a^d (arcseconds)	Height/ b^e (arcseconds)	P.A. ^f (deg)	Velocity Range ^g (km s^{-1})
NGC 1425	1	15.3; 7.5; 88.9	Rectangle	55.54668088888889, −29.892221888888887	180.0	100.0	45.0	1290 – 1690
NGC 1427	1	14.7; 7.6; 87.6	Ellipse	55.580933, −35.392563	125.0	106.8	−19.1	1210 – 1510
NGC 1427A	2	14.4; 7.6; 89.2	Ellipse	55.03875, −35.624444	87.5	55.6	−18.7	1850 – 2151
NGC 1436	2	14.7; 7.7; − 89.0	Rectangle	55.90561111111111, −35.85386133333334	90.0	110.0	0.0	1250 – 1490
NGC 1437A	2	14.6; 7.7; 88.5	Ellipse	55.75914, −36.273374	52.7	44.4	17.5	711 – 1011
NGC 1437B	2	14.2; 7.7; 87.9	Rectangle	56.478542, −36.35836088888889	30.0	70.0	0.0	1410 – 1570
NGC 1484	1	14.6; 7.6; 87.0	Ellipse	58.583875, −36.968889	72.2	24.7	−13.9	871 – 1170
PGC 012625	1	14.8; 7.7; 89.0	Ellipse	50.52, −37.5875	34.0	17.0	100.0	1450 – 1750
WOMBAT I	1	14.7; 7.7; 87.6	Ellipse	55.245185, −38.855266	23.5	19.5	37.4	662 – 961

Notes.^a See Table 2.^b The region shape to generate the galactic CO spectra for calculating total molecular gas mass of galaxies. Rectangle and ellipse are adopted for the CO-detected and CO-nondetected galaxies.^c The region center for the galactic CO spectra.^d The width and semimajor axis for the region shapes of rectangle and ellipse, respectively.^e The height and semiminor axis for the region shapes of rectangle and ellipse, respectively.^f The position angle of the region calculated from the north in a counterclockwise direction.^g The velocity range to calculate the CO integrated intensity. The range was determined by eye for the CO detected sources; $\sim 300 \text{ km s}^{-1}$ is adopted for the CO-nondetected sources.

selected from any of the smoothed cubes or from the original cube. The kernels have FWHMs of 8 and 16 pixels (1 and $2 \times$ the major axis of the synthesis beam) in the R.A. and decl. directions, and 3 pixels (30 km s^{-1}) in the velocity direction. We used the reliability module, having verified that we have enough statistics to apply that method. The SoFiA reliability of 0.95 is required to rule out false detections. Then the detection mask is dilated with a maximum growing size of 10 pixels and three channels. By this stage, the targets have been separated into detections and nondetections. For nondetections, we further run the *optical finder* of SoFiA, which searches fluxes around input optical positions. We combine the final signal-to-noise ratio and moment maps to decide whether the detections are reliable. The final products for each source from SoFiA include a detection mask, moment maps, and integral measurements of the fluxes.

For CO-detected galaxies, their CO intensities are calculated with galactic total CO spectra, which are generated by averaging the spectra of the pixels enclosed by rectangles. Here we use the original data to generate the total CO spectra since the spectra generated using the SoFiA masks do not show the level of rms. The rectangle area to generate the total CO spectrum of each galaxy and the velocity range to calculate integrated intensities are determined by eye to cover the entire area where CO emission is detected. In case of nondetections, we calculate the upper limit of their CO integrated intensities with galactic total CO spectra. The galactic total CO spectrum of a CO-nondetected galaxy is generated with an ellipse whose major and minor axes are two times smaller than the ones used to calculate the total fluxes in the infrared and ultraviolet wavelengths (see Section 3.1). For the CO line profile, we assume boxcar function with the 30 velocity channels ($\sim 300 \text{ km s}^{-1}$). In Figure 3, we present the CO moment maps of the

CO-detected galaxies on the X-ray image (Murakami et al. 2011) and the stellar mass–SFR plane. We can see that CO-detected galaxies tend to avoid being distributed near the core of the Fornax cluster and are biased to relatively massive galaxies with stellar mass of $M_{\text{star}} > 10^9 M_{\odot}$. Moment 1 maps of most samples indicate rotation of molecular gas, whereas some also show noncircular motion (NGC 1365 in Gao et al. 2021; NGC 1316 in Morokuma-Matsui et al. 2019). The CO moment maps of the 23 CO-detected galaxies are presented in Figures 14–36 in the Appendix. Their individual CO distribution and kinematics will be discussed in a separate paper.

With the CO integrated intensities obtained above, we calculate molecular gas masses (M_{mol}) in the Fornax galaxies using a constant CO-to- H_2 conversion factor of the Milky Way of $\alpha_{\text{CO,MW}} = 4.36 M_{\odot} (\text{K km s}^{-1} \text{ pc}^2)^{-1}$ (Bolatto et al. 2013) as we did in our previous study on the Virgo cluster galaxies (Morokuma-Matsui et al. 2021). This is because we do not have a homogeneous estimation of the metallicity of our sample. Furthermore, we focus on the relatively massive systems with $M_{\text{star}} > 10^9 M_{\odot}$ in the analysis in the following sections. Note that $M_{\text{star}} \sim 10^9 M_{\odot}$ is near the turnover mass of the local mass–metallicity relation of galaxies (Andrews & Martini 2013) and the uncertainty of the metallicity dependence of the CO-to- H_2 conversion factor is expected to be small for the massive galaxies. The helium contribution to mass is already included in $\alpha_{\text{CO,MW}}$ by a factor of 1.36.

For a comparison with the AIFoCS project (Zabel et al. 2019, 2020), we find that molecular gas masses for the CO-detected objects in both the AIFoCS and this study (13 objects) are mostly consistent within a factor of ~ 1.4 , except for NGC 1365, whose M_{mol} estimation in this study is 3.7 times larger than the one in the AIFoCS. Note that the molecular gas mass in the AIFoCS data is recalculated with the constant CO-to- H_2

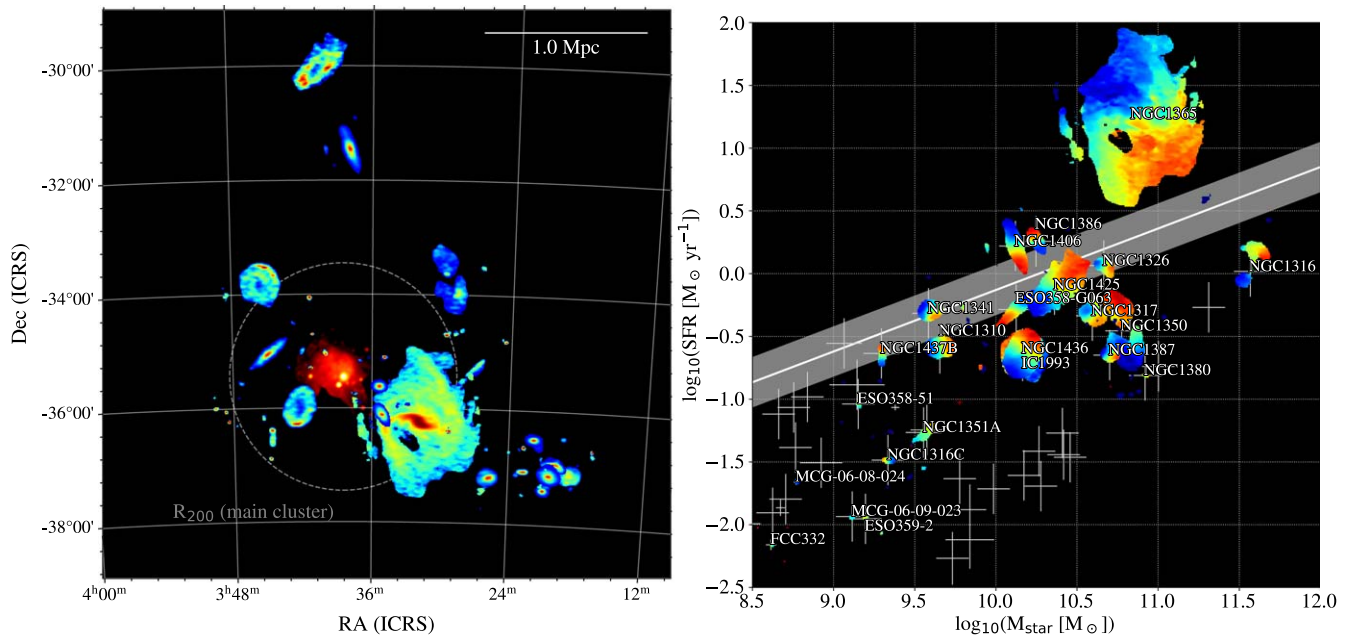


Figure 3. (Left) Sky distributions of 23 Fornax cluster galaxies with CO detections. The spatial size of CO data (moment 0) is enlarged by a factor of 30. The linear scaling is adopted for plotting the moment 0 map of the CO-detected galaxies except for NGC 1365, for which the logarithmic scaling is adopted. Background is the X-ray image obtained by combining the Suzaku and XMM-Newton data for the Fornax cluster (Murakami et al. 2011). (Right) Stellar mass vs. SFR relation of our sample galaxies. Moment 1 maps of the CO-detected galaxies are shown. The white solid line and shaded region indicate the star-forming main sequence of Speagle et al. (2014) and its standard deviation, respectively.

conversion factor of $\alpha_{\text{CO,MW}} = 4.36 M_{\odot} (\text{K km s}^{-1} \text{pc}^2)^{-1}$ (Bolatto et al. 2013) to match our data. If we compare the moment 0 maps of NGC 1365 with a large difference in the M_{mol} estimation between this study and the AIFoCS project, there would be missing flux in the 12 m array data, confirming the importance of the ACA data for calculating total fluxes of the extended nearby objects. Note that the maximum recoverable scales of the AIFoCS project are $20''8\text{--}23''3$ ($\sim 2.0\text{--}2.3$ kpc at 20 Mpc). The AIFoCS detected CO emission from FCC 207 and FCC 261, whereas we could not since M_{mol} values of these galaxies are below the sensitivity of our observation. There is a weak M_{mol} dependence in the M_{mol} estimation difference of the CO-detected objects where our estimation tends to be lower/higher for less-massive/massive objects compared with the AIFoCS estimation (Figure 4), although the dependence becomes insignificant if NGC 1365 is excluded ($r = 0.55$ and $p = 0.07$). For all the nondetected sources in this study, our upper limits are larger (mostly five times larger) than those estimated in the AIFoCS project, that is, a more conservative estimation in this study. This is partially due to the difference in the sensitivity and the assumed line width to calculate the upper limit of the CO intensity. We assumed 300 km s^{-1} for all the CO nondetections including dwarf and ETGs, whereas the AIFoCS assumed 50 km s^{-1} since their CO nondetections are all dwarf galaxies.

3. Ancillary Data

3.1. Stellar Mass and SFR

We adopted M_{star} and SFR values derived in the project $z = 0$ Multiwavelength Galaxy Synthesis (z0MGS; Leroy et al. 2019). Leroy et al. (2019) estimated M_{star} using the WISE 3.4 μm data and SFR using the WISE 22 μm data and far-ultraviolet data obtained with the Galaxy Evolution Explorer (GALEX; Martin et al. 2005) to account for both obscured and

unobscured components. M_{star} and SFR values of 46 out of 64 galaxies are found in the z0MGS catalog. For the rest of the 18 galaxies, we adopt the M_{star} and SFR values estimated by ourselves as follows.

We first measure the total fluxes in the WISE and GALEX bands of the 64 Fornax galaxies using the Comprehensive and Adaptable Aperture Photometry Routine (CAAPR²⁷; Clark et al. 2015, 2018). The CAAPR can remove the foreground and background objects with the Python Toolkit for SKIRT (Camps et al. 2015) and also remove the any large-scale background structure if desired. The aperture ellipse in each wavelength is determined to cover the pixels where the signal-to-noise ratio > 2 . The signal-to-noise ratio = 2 ellipse is then enlarged by a factor of 1.25, which is a default expansion factor in the CAAPR. The final aperture area is the smallest ellipse that contains all the ellipses in different wavelengths. The detailed description on the CAAPR is found in Clark et al. (2018). For MCG-06-08-024 and PGC 012625, we manually determined the photometry aperture to avoid including nearby stars. The parameters for the final ellipses obtained with the CAAPR for our sample are presented in Table 4. With the total fluxes of galaxies in the WISE and GALEX bands, the M_{star} and SFR values are calculated in the same manner adopted in the z0MGS. The relationship between M_{star} and SFR of our sample galaxies is presented in the right panel of Figure 2.

To assess our estimations of M_{star} and SFR, Figure 5 compares the M_{star} and SFR values estimated in this study and in the z0MGS for the 46 Fornax galaxies with the z0MGS data. Overall, our estimations of M_{star} and SFR are consistent with those of the z0MGS project within the errors. We find that the M_{star} estimations for two galaxies (ESO 302-9 and IC 1993) and the SFR estimations for two galaxies (NGC 1419 and ESO 358-5) are different from those in the z0MGS. Given that we

²⁷ <https://github.com/Stargrazer82301/CAAPR>

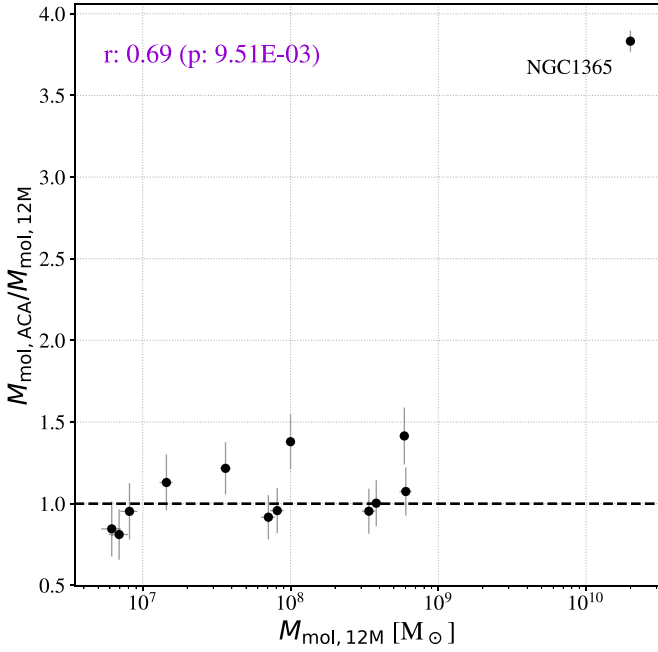


Figure 4. M_{mol} comparison between Zabel et al. (2019) ($M_{\text{mol},12\text{M}}$) and this study ($M_{\text{mol},\text{ACA}}$). The Spearman's correlation coefficient and its p -value are presented at the upper left corner in purple. The galaxies with a large M_{mol} value tend to have a high $M_{\text{mol},\text{ACA}}/M_{\text{mol},12\text{M}}$ ratio, suggesting the missing flux in the 12 m data.

adopted the same calibration methods of the z0MGS, the difference in the M_{star} and SFR values is mainly caused by the differences in the photometry processes such as the determination of the photometry/background areas and the subtraction of the foreground and background objects. Note that the M_{star} and SFR values of the massive galaxies analyzed in Section 4 are mostly from the z0MGS (37 out of 38 galaxies); thus the effects to obtained results due to the difference in the photometry processes are expected to be negligible.

3.2. Atomic Gas Mass

We make use of the atomic gas mass (M_{atom}) of the Fornax galaxies from Loni et al. (2021) and Kleiner et al. (2021), which conducted H I observations with the Australia Telescope Compact Array (ATCA) and the Meer Karoo Array Telescope (MeerKAT). For the galaxies whose H I masses are not listed in these studies, we utilize the H I fluxes that are mostly measured in the H I Parkes All Sky Survey (HIPASS) in the HyperLEDA catalog. H I data were found for 42 out of the 64 Fornax galaxies: 20 main-cluster galaxies from Loni et al. (2021), seven Fornax A group galaxies from Kleiner et al. (2021), and 15 galaxies from the Hyper-LEDA. M_{atom} is calculated based on the optically thin assumption as

$$M_{\text{atom}} = 1.36 \times 2.36 \times 10^5 D^2 \times \sum_i S_i \Delta v_i, \quad (1)$$

where the factor of 1.36 accounts for the helium contribution to mass, D is the luminosity distance to the galaxies in megaparsecs (20 Mpc), and $\sum_i S_i \Delta v_i$ is the total integrated flux in Jy km s^{-1} .

Table 4
Parameters for CAAPR^a

Name	Semimajor (arcseconds)	Semiminor (arcseconds)	P.A. (deg)
ESO 302-14	70.0	34.5	67.9
ESO 302-9	141.0	55.6	128.6
ESO 357-25	55.8	20.8	117.2
ESO 358-16	42.2	16.8	66.2
ESO 358-20	80.8	45.2	55.0
ESO 358-5	86.8	62.6	12.5
ESO 358-51	80.5	45.2	94.5
ESO 358-G015	54.2	35.1	84.4
ESO 358-G063	266.7	102.3	38.6
ESO 359-2	31.4	26.3	-35.0
ESO 359-3	111.9	58.2	44.0
ESO 359-5	54.0	34.7	45.3
FCC 102	38.1	23.4	59.5
FCC 117	13.0	13.0	45.0
FCC 120	42.9	16.4	132.9
FCC 177	126.5	71.4	82.1
FCC 198	13.0	13.0	45.0
FCC 206	98.5	52.4	-24.3
FCC 207	29.9	23.0	-17.6
FCC 261	128.4	44.4	49.2
FCC 302	92.6	19.3	13.1
FCC 306	37.6	27.8	-35.4
FCC 316	35.5	20.0	35.0
FCC 32	44.0	37.9	72.8
FCC 332	52.6	36.9	16.4
FCC 44	23.6	16.6	33.6
IC 1993	112.0	104.5	126.4
IC 335	159.5	71.7	4.2
MCG-06-08-024	22.0	18.3	45.0
MCG-06-09-023	84.5	64.3	-12.5
NGC 1310	128.8	101.0	5.1
NGC 1316	797.4	557.0	-34.9
NGC 1316C	84.4	45.2	-0.7
NGC 1317	177.5	167.3	-8.5
NGC 1326	236.6	180.2	-13.0
NGC 1326A	124.5	55.4	12.7
NGC 1326B	183.8	59.7	32.7
NGC 1336	152.4	105.0	109.2
NGC 1339	150.4	118.5	49.9
NGC 1341	134.5	109.4	51.4
NGC 1350	304.1	181.5	96.5
NGC 1351	220.3	155.6	53.3
NGC 1351A	125.9	36.5	40.5
NGC 1365	486.1	330.9	124.0
NGC 1374	156.7	149.0	63.0
NGC 1375	108.2	54.1	0.2
NGC 1379	180.5	165.4	-4.2
NGC 1380	329.1	240.4	98.6
NGC 1381	163.5	82.6	51.8
NGC 1386	221.1	194.2	133.2
NGC 1387	211.2	182.9	-27.7
NGC 1399	542.2	451.7	-15.5
NGC 1404	230.0	177.3	72.4
NGC 1406	226.0	181.8	123.5
NGC 1419	106.0	86.6	-7.4
NGC 1425	265.7	128.2	42.9
NGC 1427	250.0	213.7	-19.1
NGC 1427A	175.0	111.3	-18.7
NGC 1436	168.0	127.0	56.9
NGC 1437A	105.4	88.7	17.5
NGC 1437B	133.5	59.9	85.8
NGC 1484	144.3	49.4	-13.9
PGC 012625	68.0	34.0	100.0
WOMBAT I	46.9	39.0	37.4

Note.

^a The photometric aperture ellipses for MCG-06-08-024 and PGC 012625 are manually determined to prevent nearby stars from being included in the ellipses.

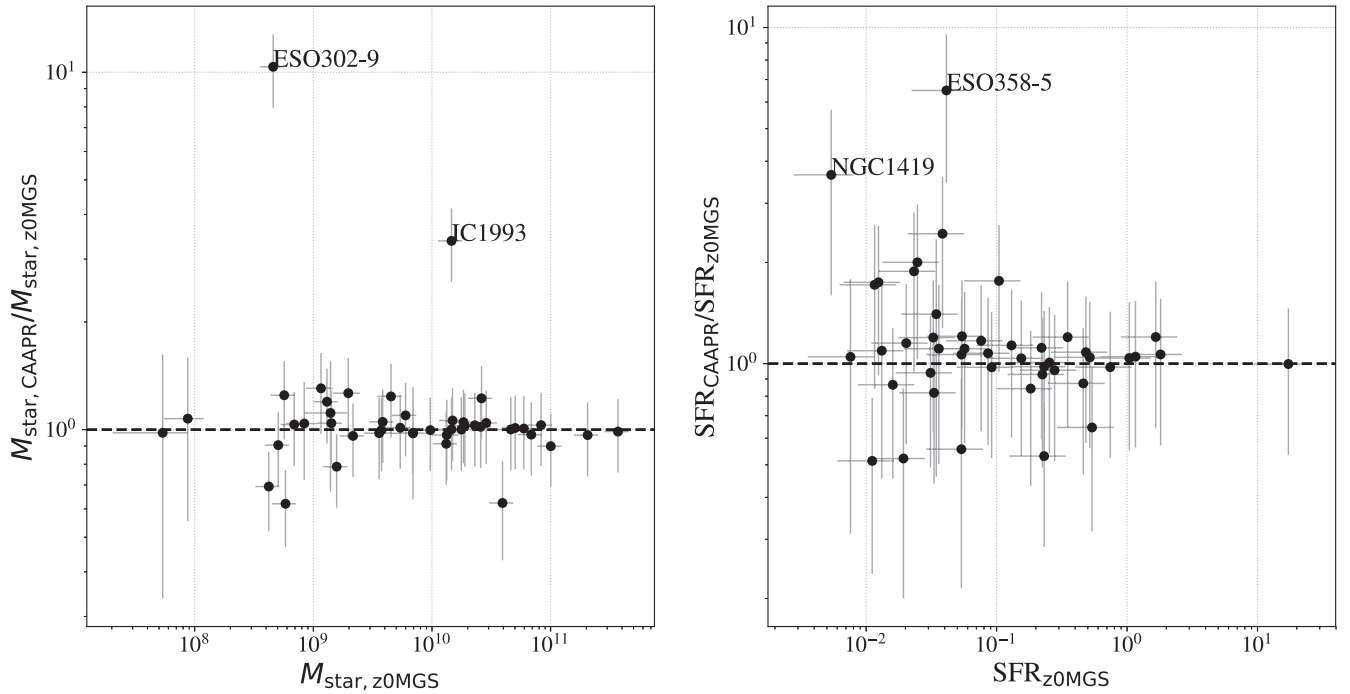


Figure 5. Comparison of M_{star} and SFR estimations in Leroy et al. (2019) ($M_{\text{star,z0MGS}}$ and $\text{SFR}_{\text{z0MGS}}$) and the ones in this study ($M_{\text{star,CAAPR}}$ and $\text{SFR}_{\text{CAAPR}}$).

4. Cold Gas and Star Formation Properties

In the following sections, we basically follow the analysis procedure adopted in Morokuma-Matsui et al. (2021) for the Virgo cluster galaxies. The comparison among different local clusters including the Fornax and Virgo will be presented in a separate paper (K. Morokuma-Matsui et al. 2022, in preparation). We present cold gas and SF properties of the Fornax galaxies focusing on the following key quantities (Table 5): (1) SFR, (2) specific SFR (sSFR = $\text{SFR}/M_{\text{star}}$), (3) molecular-gas mass to atomic gas mass ratio ($R_{\text{H}_2} = M_{\text{mol}}/M_{\text{atom}}$), (4) total gas mass to stellar mass ratio ($\mu_{\text{gas}} = M_{\text{gas}}/M_{\text{star}} = [M_{\text{mol}} + M_{\text{atom}}]/M_{\text{star}}$), (5) atomic gas mass to stellar mass ratio ($\mu_{\text{H I}} = M_{\text{atom}}/M_{\text{star}}$), (6) molecular gas mass to stellar mass ratio ($\mu_{\text{H}_2} = M_{\text{mol}}/M_{\text{star}}$), (7) SF efficiency (SFE) from total gas ($\text{SFE}_{\text{gas}} = \text{SFR}/M_{\text{gas}}$), (8) SFE from atomic gas ($\text{SFE}_{\text{H I}} = \text{SFR}/M_{\text{H I}}$), and (9) SFE from molecular gas ($\text{SFE}_{\text{H}_2} = \text{SFR}/M_{\text{mol}}$).

We consider the property distributions of two samples to be different if the p -value of the Kolmogorov–Smirnov (K-S) test is smaller than 0.05. We also assess the correlation between two quantities by Spearman’s rank-order correlation method. We consider that the correlation coefficient, r , is reliable if the p -value is smaller than 0.05.²⁸ Here we make use of `ks_2samp` and `spearmanr` in `Scipy` for the K-S test and the estimation of the correlation coefficient and p -value, respectively. The key quantities of the Fornax galaxies are compared on the best-effort basis, that is, all the galaxies with measurements (including upper limits) of the numerators and denominators of the key quantities (“best-effort sample”). However, for R_{H_2} , we limit the sample to the galaxies with detection in both CO and H I. Depending on the quantities to be compared, the number of galaxies in the following plots is different. The galaxies with the H I measurement are mostly H I detected and

Table 5
Key Quantities

Quantity	Unit	Description
SFR	$[M_{\odot} \text{ yr}^{-1}]$	Star formation rate
sSFR	$[\text{yr}^{-1}]$	$\text{SFR}/M_{\text{star}}$, specific SFR
R_{H_2}		$M_{\text{mol}}/M_{\text{atom}}$, molecular-gas mass to atomic gas mass ratio
μ_{gas}		$M_{\text{gas}}/M_{\text{star}} = [M_{\text{mol}} + M_{\text{atom}}]/M_{\text{star}}$, total gas mass to stellar mass ratio
$\mu_{\text{H I}}$		$M_{\text{atom}}/M_{\text{star}}$, atomic gas mass to stellar mass ratio
μ_{H_2}		$M_{\text{mol}}/M_{\text{star}}$, molecular gas mass to stellar mass ratio
SFE_{gas}	$[\text{yr}^{-1}]$	$\text{SFR}/M_{\text{gas}}$, star formation efficiency from total gas
$\text{SFE}_{\text{H I}}$	$[\text{yr}^{-1}]$	$\text{SFR}/M_{\text{atom}}$, star formation efficiency from atomic gas
SFE_{H_2}	$[\text{yr}^{-1}]$	$\text{SFR}/M_{\text{mol}}$, star formation efficiency from molecular gas

biased to star-forming main sequence (SFMS) galaxies (Figure 2 right). Therefore, the quantities with H I data, such as R_{H_2} , $\mu_{\text{H I}}$, μ_{gas} , SFE_{gas} , and $\text{SFE}_{\text{H I}}$, may not represent the majority of our sample galaxies. To assess the sample bias caused by this treatment, the key quantities are also compared for galaxies with all the measurements (“CO+H I-obs. sample”) or detections of H I and CO (“CO+H I-det. sample”) in Section 5.1, although the sample size becomes small. The summary of the subsamples is presented in Table 6.

We first compare the Fornax galaxies and the field galaxies in Section 4.1 and then investigate the dependence of the key quantities on three environmental parameters: clustocentric distance in Section 4.2, the distance to the fifth nearest galaxy in Section 4.3, and the accretion phase in the cluster in Section 4.4.

4.1. Comparison with Field Galaxies

To clarify the environmental effects on cluster galaxies independently of the mass-dependent galaxy evolution, we first

²⁸ We use the term “significant correlation” if its p -value is smaller than 0.05 and the term “stronger/weaker” for even a moderate correlation with $r \sim 0.5$ just in a relative sense, i.e., in a comparison of two r values.

Table 6
Numbers of the 64 Fornax Galaxies for the Subsamples

Subsample	With Data	With Detection
H ₂	64	23
H I	42	36
H ₂ and H I	42	16
H ₂ , <i>M</i> _{star9} ^a	38	21
H I, <i>M</i> _{star9}	23	20
H ₂ , <i>M</i> _{star9} and H I, <i>M</i> _{star9}	23 ^b	16 ^c

Notes.

^a Galaxies with $M_{\text{star}} > 10^9 M_{\odot}$, i.e., with $\Delta(\text{Field})$ measurements.

^b CO+H I-obs. sample.

^c CO+H I-det. sample.

determine the field relations, that is, the stellar mass dependence of the key quantities of the field star-forming galaxies (hereafter “field SFG relations”), by making use of the data of the extended GALEX Arecibo SDSS Survey (xGASS, 1179 galaxies; Catinella et al. 2018) for H I-related relations and the extended CO Legacy Database for GASS (xCOLD GASS, 532 galaxies; Saintonge et al. 2017) for H₂-related relations. The data of the xGASS and xCOLD GASS projects are used because their samples are selected only by stellar mass and redshift. For consistency with our Fornax data, the molecular gas mass of the xCOLD-GASS galaxies is recalculated with the constant $\alpha_{\text{CO,MW}}$ of $4.36 M_{\odot} (\text{K km s}^{-1} \text{pc}^2)^{-1}$, and the atomic gas mass of the xGASS galaxies is scaled with 1.36 to account for the helium. The field SFG relations are derived nonparametrically and by third-order polynomial fitting using the data of CO-detected xCOLD GASS galaxies and/or H I-detected xGASS galaxies. For the SFR-related relations, we adopted the SFMS derived based on the orthogonal distance regression using the data of $>10^6$ galaxies at $0 < z < 6$ in Speagle et al. (2014).

The comparison of the key quantities among the Fornax sample, the field (xGASS and xCOLD GASS) sample including both SF and passive galaxies, and the field SFG relations is shown in Figure 6. The contours in panels (a) and (b) indicate the entire field galaxies, and the contours in the rest of the panels indicate the field galaxies with H I or CO detection, that is, star-forming galaxies. The open triangles/inverse triangles indicate the field galaxies with lower/upper limits of the quantities. We can see that most of our Fornax galaxies (orange and blue symbols) are on or below the SFMS. There seems to be a tendency where all the gas fractions of the Fornax galaxies are lower than those of the field galaxies even when we limit to the galaxies with H I or CO detections. Although the R_{H_2} values of the Fornax galaxies seem to be similar to those of the field galaxies, some Fornax galaxies have higher R_{H_2} . In terms of the SFEs, most Fornax galaxies seem to have higher SFE_{gas} and $\text{SFE}_{\text{H I}}$ than the field galaxies, especially for the massive galaxies with $>10^{10} M_{\odot}$. Although a similar mass dependence is observed for SFE_{H_2} , the difference between the Fornax and field galaxies seems insignificant compared to SFE_{gas} and $\text{SFE}_{\text{H I}}$.

For both the Fornax and field samples, the deviation from the field SFG relations is calculated by

$$\Delta(\text{Field}) = \log_{10}(X_{\text{Fornax/Field}}(M_{\text{star}})/X_{\text{Field,SFG}}(M_{\text{star}})). \quad (2)$$

$X_{\text{Fornax/Field}}(M_{\text{star}})$ is a key quantity of the Fornax or field galaxies with a stellar mass of M_{star} . $X_{\text{Field,SFG}}(M_{\text{star}})$ is the field

SFG relation that is derived based on the third-order polynomial regression (except for SFR and sSFR where Speagle’s SFMS is used). Figure 7 shows the histograms of key quantities in the form of $\Delta(\text{Field})$ of the Fornax and field galaxies. According to the K-S test, the tendency seen in Figure 6 is confirmed quantitatively. In the Fornax cluster a larger fraction of galaxies tend to have lower SF activity, a smaller amount of atomic and molecular gas, and a higher SFE from the atomic or the total cold (atomic plus molecular) gas, compared to the field galaxies with the same stellar mass. Furthermore, the Fornax galaxies have a wider range of R_{H_2} and have a higher median value of R_{H_2} than the field galaxies. Although the upper/lower limits of the key quantities for the field galaxies are looser than those of the Fornax galaxies (Figure 6), we confirmed that the obtained result does not change even when limiting to the galaxies with H I or CO detection.

Both the μ_{H_2} and SFE_{H_2} values of galaxies have been claimed to positively correlate with the offset from the SFMS of the field galaxies,

$$\Delta(\text{MS}) = \log_{10}(\text{SFR}/\text{SFR}_{\text{MS}}) \quad (3)$$

(e.g., Saintonge et al. 2012; Genzel et al. 2015; Scoville et al. 2017; Ellison et al. 2020), where SFR_{MS} is the SFR of the SFMS of the field galaxies for a fixed stellar mass and redshift, that is, $\Delta(\text{MS}) = \Delta(\text{Field})_{\text{SFR}}$. These relations are claimed to be independent of the galaxy environment at least for field and group galaxies (Koyama et al. 2017). For denser environments, Morokuma-Matsui et al. (2021) found that the Virgo galaxies with low $\Delta(\text{MS})$ values offset to lower μ_{H_2} and higher SFE_{H_2} regimes with respect to the field and group relations.

Figure 8 shows the distribution of the Fornax and field galaxies with $M_{\text{star}} > 10^9 M_{\odot}$ on the plane of $\Delta(\text{MS})$ versus $\Delta(\text{Field})$ values of the cold gas mass fractions, R_{H_2} , or SFEs. The comparison to the original values (before the mass dependence is subtracted) is presented in Figure 37 in the Appendix. Here we limit the galaxies to those with H I or CO detections; this is because the observation sensitivity, that is, the strength of the upper/lower limits, is different between our ALMA observation and the ones for the field galaxies. The Fornax galaxies below the main sequence [$-1.5 < \Delta(\text{MS}) < -0.5$] tend to have smaller gas fractions and larger R_{H_2} ratio, SFE_{gas} , and $\text{SFE}_{\text{H I}}$ compared to field galaxies according to the K-S test. μ_{H_2} and μ_{gas} of the Fornax galaxies is lower than those of the field galaxies even on the main sequence [$-0.5 < \Delta(\text{MS}) < 0.5$].

To understand the properties of galaxies whose molecular and/or atomic gas is depleted, Figure 9 compares the offsets of the H I and H₂ gas fractions and SFEs from the values of the field galaxies for fixed stellar masses. The color of the symbols indicates $\Delta(\text{MS})$ of galaxies. Note that the sample is limited to the galaxies with both H I and H₂ measurements and the galaxies with large H I depletion are largely missed in this analysis. For a reference, the field galaxies with both CO and H I detections are shown as contours. There is no significant correlation between $\mu_{\text{H I}}$ and μ_{H_2} or between $\text{SFE}_{\text{H I}}$ and SFE_{H_2} of the Fornax galaxies. The Fornax galaxies with lower $\Delta(\text{MS})$ tend to have lower $\mu_{\text{H I}}$ and μ_{H_2} , while the location of galaxies on the $\text{SFE}_{\text{H I}}-\text{SFE}_{\text{H}_2}$ plot does not strongly depend on $\Delta(\text{MS})$.

For the comparison of the gas fractions, we can see that the Fornax galaxies tend to distribute at the third quadrant in the

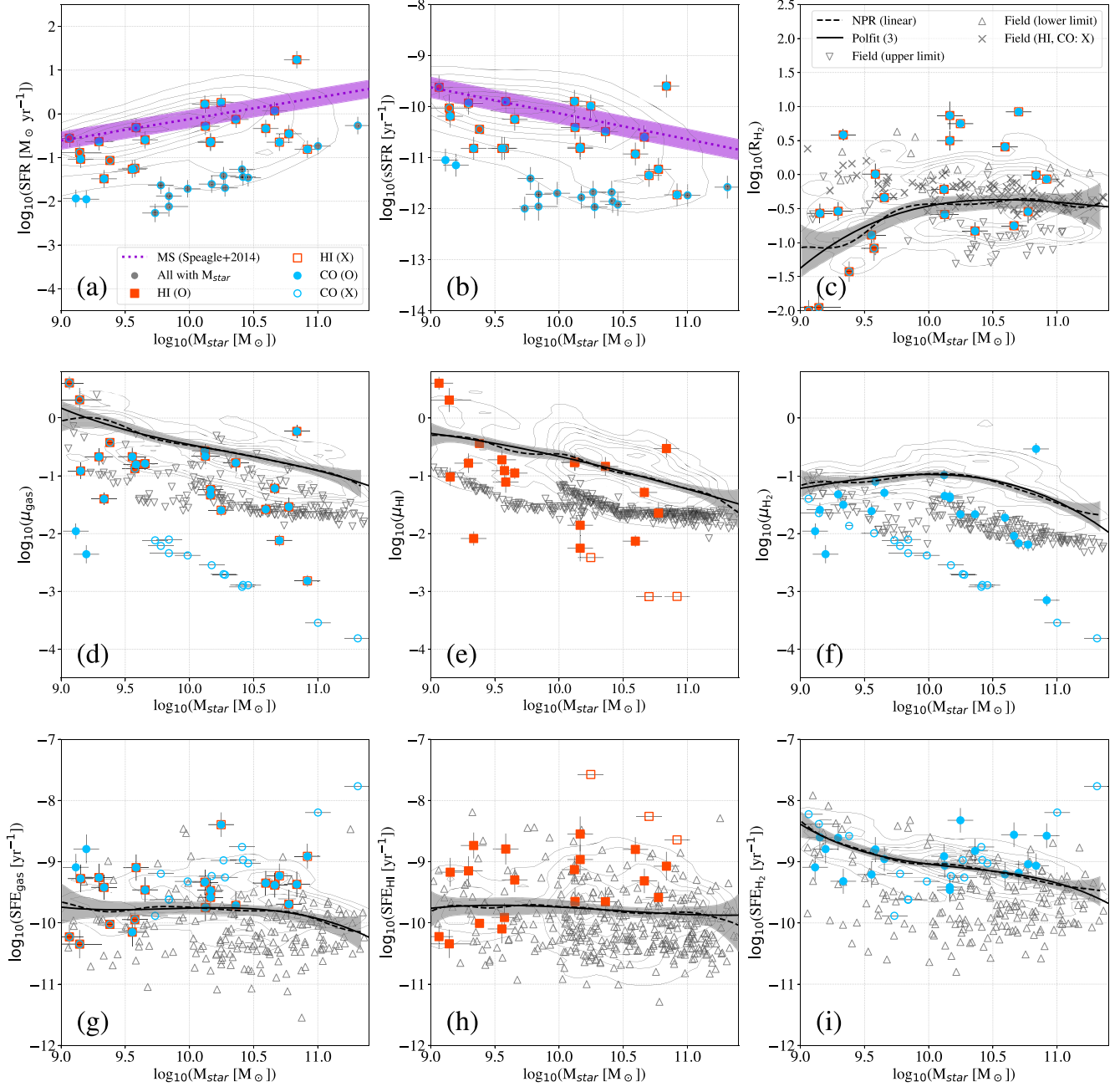


Figure 6. Comparison between the Fornax galaxies and field galaxies from the z0MGS (Leroy et al. 2019), xGASS (Catinella et al. 2018), and xCOLDGASS (Saintonge et al. 2017) projects: M_{star} dependences of (a) SFR, (b) sSFR, (c) R_{H_2} , (d) $\mu_{\text{gas}} = M_{\text{gas}}/M_{\text{star}}$, (e) $\mu_{\text{H I}} = M_{\text{H I}}/M_{\text{star}}$, (f) $\mu_{\text{H}_2} = M_{\text{H}_2}/M_{\text{star}}$, (g) $\text{SFE}_{\text{gas}} = \text{SFR}/M_{\text{gas}}$, (h) $\text{SFE}_{\text{H I}} = \text{SFR}/M_{\text{H I}}$, (i) $\text{SFE}_{\text{H}_2} = \text{SFR}/M_{\text{H}_2}$. The main sequence of star-forming galaxies defined in Speagle et al. (2014) is indicated as purple dotted lines, and the 0.2 dex scatter range is indicated as purple shading in panels (a) and (b). The field relations are estimated with CO- or H I-detected galaxies, i.e., star-forming galaxies, and indicated as dashed (with nonparametric fitting) and solid (with third-order polynomial fitting) lines in panels (c)–(i). Gray-shaded regions are the errors of the nonparametric fitting and estimated by the bootstrap method with a 95% confidence interval. The Fornax galaxies with CO/H I detection and nondetection are indicated as blue filled circles/orange filled squares and blue open circles/orange open squares, respectively. The contours indicate field star-forming galaxies used to derive the field relations. The lower/upper limits of the field galaxies are indicated as open triangles/inverted triangles. In panels (d) and (g), $M_{\text{H}_2}/M_{\text{star}}$ and $\text{SFR}/M_{\text{H}_2}$ are plotted for the galaxies without H I data, respectively. For the R_{H_2} panel, the field galaxies with H I and CO upper limits are indicated as X marks.

plot, that is, both the molecular and atomic gas is depleted in the Fornax galaxies compared to the field galaxies for their stellar masses. The dotted line indicates $\Delta(\text{Field})$, $\mu_{\text{H}_2} = \Delta(\text{Field})$, $\mu_{\text{H I}}$, that is, the galaxies on this line have the same degree of depletion in molecular and atomic gas. In the galaxies that distribute above the dotted line in the third quadrant, H I gas is more deficient than the H_2 gas compared to the field

galaxies. For some galaxies with the smallest $\mu_{\text{H I}}$, molecular gas is not depleted as much as atomic gas. Next, for the comparison of the SFE, most of the galaxies have comparable SFE_{H_2} to the field galaxies, while many galaxies have enhanced $\text{SFE}_{\text{H I}}$ compared to the field galaxies.

Hereafter, we use the $\Delta(\text{Field})$ values, that is, the values from which M_{star} dependences of field galaxies have been

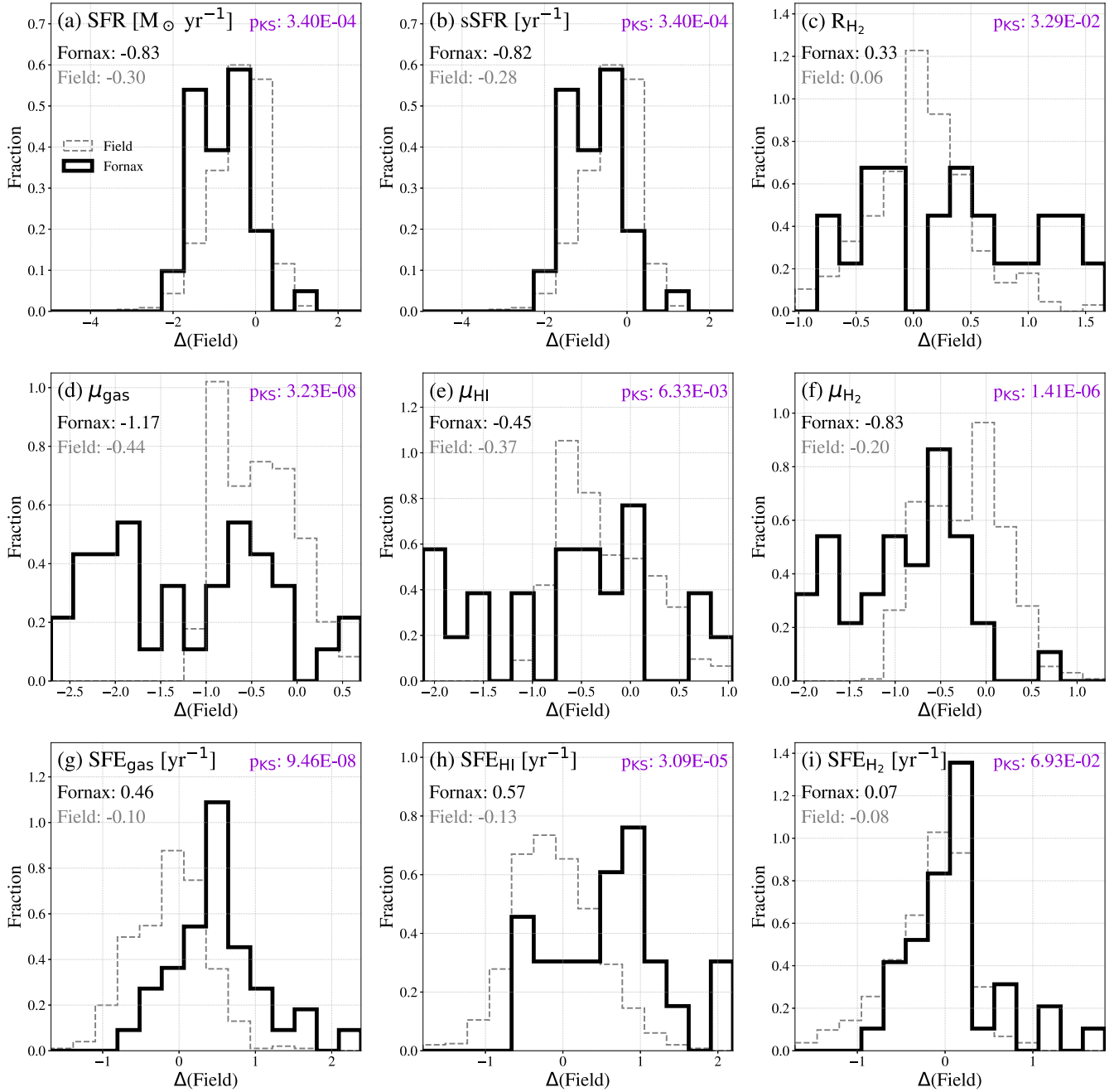


Figure 7. Comparison between the Fornax galaxies (solid line) and field galaxies (dashed line) with $M_{\text{star}} > 10^9 M_{\odot}$: histograms for the offsets from the field relations, $\Delta(\text{Field})$. Note that $\Delta(\text{Field})$ s of SFR and sSFR are calculated based on the SFMS definition in Speagle et al. (2014). The medians for the Fornax and field galaxies are indicated at the upper left corner. The p -value for the K-S test comparing the Fornax and field galaxies is indicated at the upper right corner. Our Fornax sample has lower SFR, sSFR, and gas fractions and higher R_{H_2} , SFE_{gas} , and SFE_{H_1} for their masses. SFE_{H_2} of our Fornax sample is not significantly different from the one of the field galaxies.

subtracted, in order to investigate the environmental effect without the M_{star} effect.

4.2. Dependence on the Projected Clustocentric Distance

The dependence of the galaxy properties on the projected clustocentric distance (R_{proj}) has been explored in various studies (e.g., Giovanelli & Haynes 1985; Solanes et al. 2001). R_{proj} is an easy indicator of the elapsed time since the galaxy is trapped by the cluster potential, that is, the accretion phase, as well as an indicator of the ICM density since the ICM density is higher at the cluster core than the outskirts to a first

approximation. The clustocentric radius of our sample is divided into six radial bins, with each bin spanning $0.5 R_{200}$, that is, $R_{\text{proj}}/R_{200} < 0.5$ (R_1), $0.5 \leq R_{\text{proj}}/R_{200} < 1.0$ (R_2), $1.0 \leq R_{\text{proj}}/R_{200} < 1.5$ (R_3), $1.5 \leq R_{\text{proj}}/R_{200} < 2.0$ (R_4), $2.0 \leq R_{\text{proj}}/R_{200} < 2.5$ (R_5), and $R_{\text{proj}}/R_{200} \geq 2.5$ (R_6), where R_{200} is the radius where the mean interior density is 200 times the critical density of the universe and we adopted the virial radius of $2^{\circ}0$ for R_{200} (Drinkwater et al. 2001). The most outskirts galaxy in our best-effort sample locates at $\sim 2.5R_{200}$.

Figure 10 shows the medians and first and third quartiles of the key quantities for galaxies in each category. There seems to be a radial trend in SFR, sSFR, and μ_{H_2} , where galaxies that are

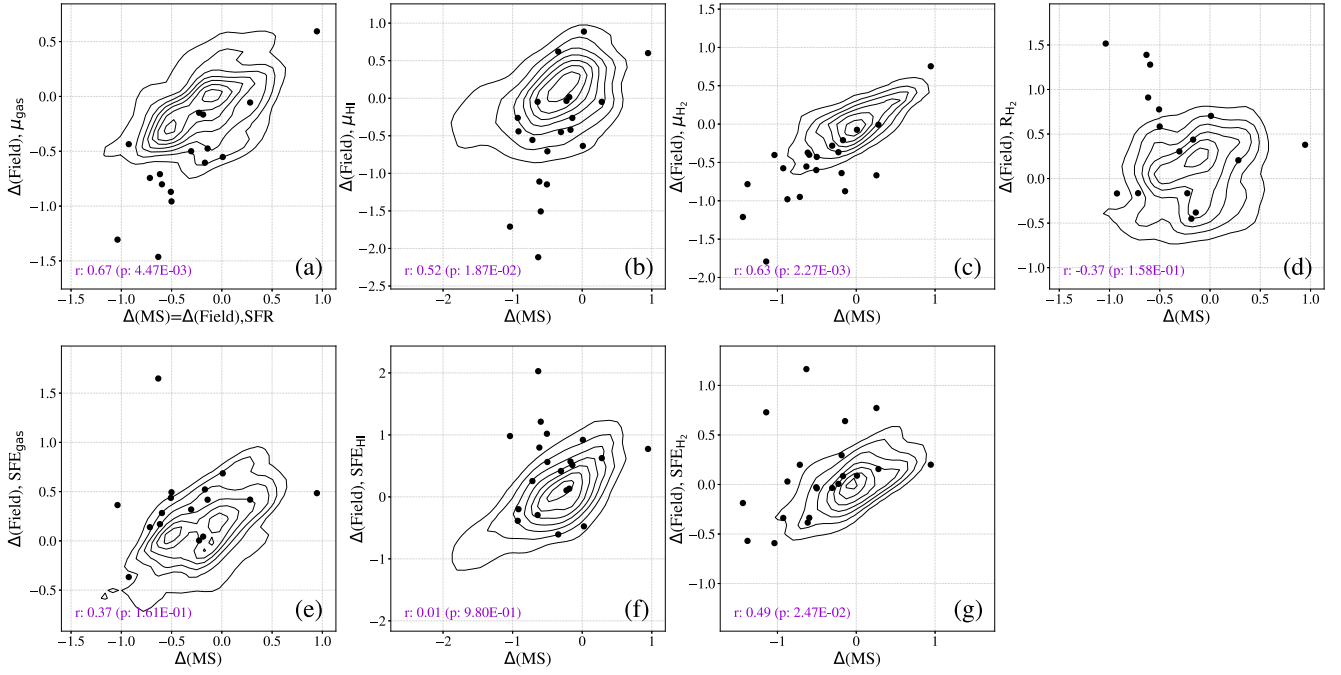


Figure 8. Comparison between the Fornax galaxies and field galaxies with $M_{\text{star}} > 10^9 M_{\odot}$: relationships between the offsets from the main sequence of star-forming galaxies [$\Delta(\text{MS}) = \Delta(\text{Field})_{\text{SFR}}$] with $\Delta(\text{Field})$ values of gas fractions, R_{H_2} , and SFEs. The contours indicate field galaxies with CO/H I detection, and filled black circles indicate the Fornax galaxies with CO/H I detection. The Spearman's rank-order correlation coefficient r and the p -value are shown on the lower left corner of each panel in purple. Our Fornax sample has lower gas fractions and higher R_{H_2} and SFEs than the field galaxies at $\Delta(\text{MS}) \sim -1$.

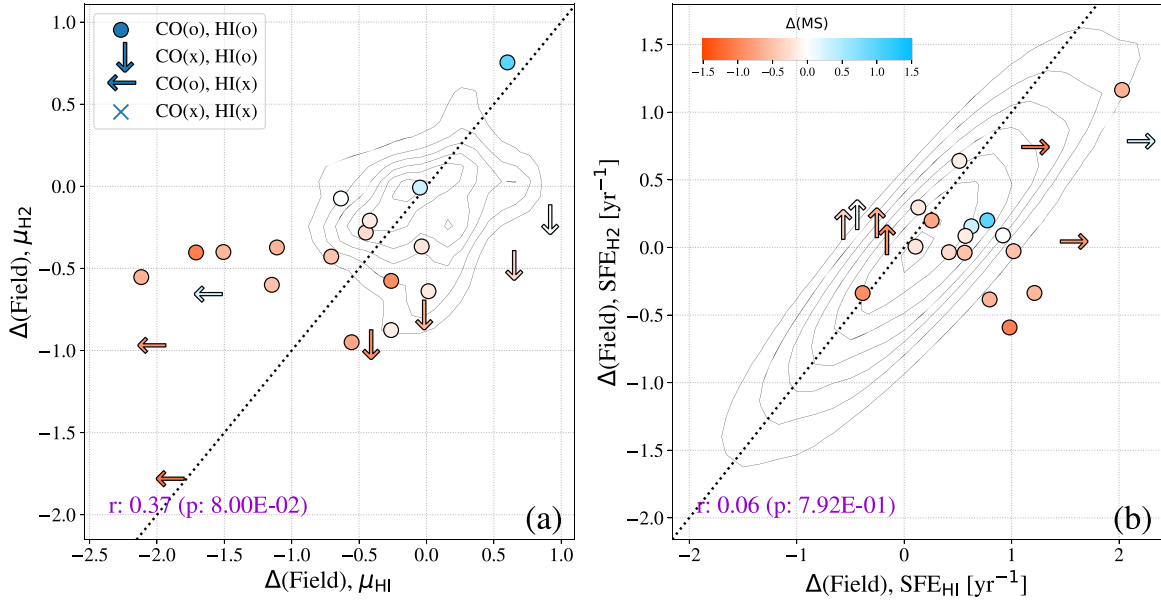


Figure 9. H I–H₂ comparison of the offsets from the field SF galaxies at fixed stellar masses of gas fraction (a) and SFE (b). The galaxies with both CO and H I detection are indicated as circles, and those with upper limits are indicated as arrows. The color of the symbol indicates $\Delta(\text{MS})$. The contours indicate the field galaxies with CO and H I detection. The Spearman's rank-order correlation coefficient r and the p -value are shown on the lower left corner of each panel in purple. Although there is no clear correlation between $\Delta(\text{Field}), \mu_{\text{H I}}$, and $\Delta(\text{Field}), \mu_{\text{H}_2}$ or $\Delta(\text{Field}), \text{SFE}_{\text{H I}}$, and $\Delta(\text{Field}), \text{SFE}_{\text{H}_2}$, our Fornax sample with H I and H₂ measurements tends to be more depleted in H I than in H₂ and tends to have higher $\text{SFE}_{\text{H I}}$ than field galaxies, whereas their SFE_{H_2} is comparable to that of the field galaxies.

nearer to the cluster center have a lower value compared to the galaxies at the cluster outskirts. However, only the trend of μ_{H_2} is further confirmed by the Spearman's rank-order correlation coefficient (0.5) and the p -value (< 0.05). A significant radial dependence is not seen in the other quantities. When compared to the field SF galaxies represented by a zero value in Figure 10, the key quantities of the Fornax galaxies are not

significantly different from those of the field galaxies except for SFR, sSFR, and μ_{H_2} .

4.3. Dependence on the Galaxy Number Density

Galaxy number density is also one of the important parameters to understand the dominant quenching mechanism in cluster galaxies. If the SF and cold gas properties of galaxies

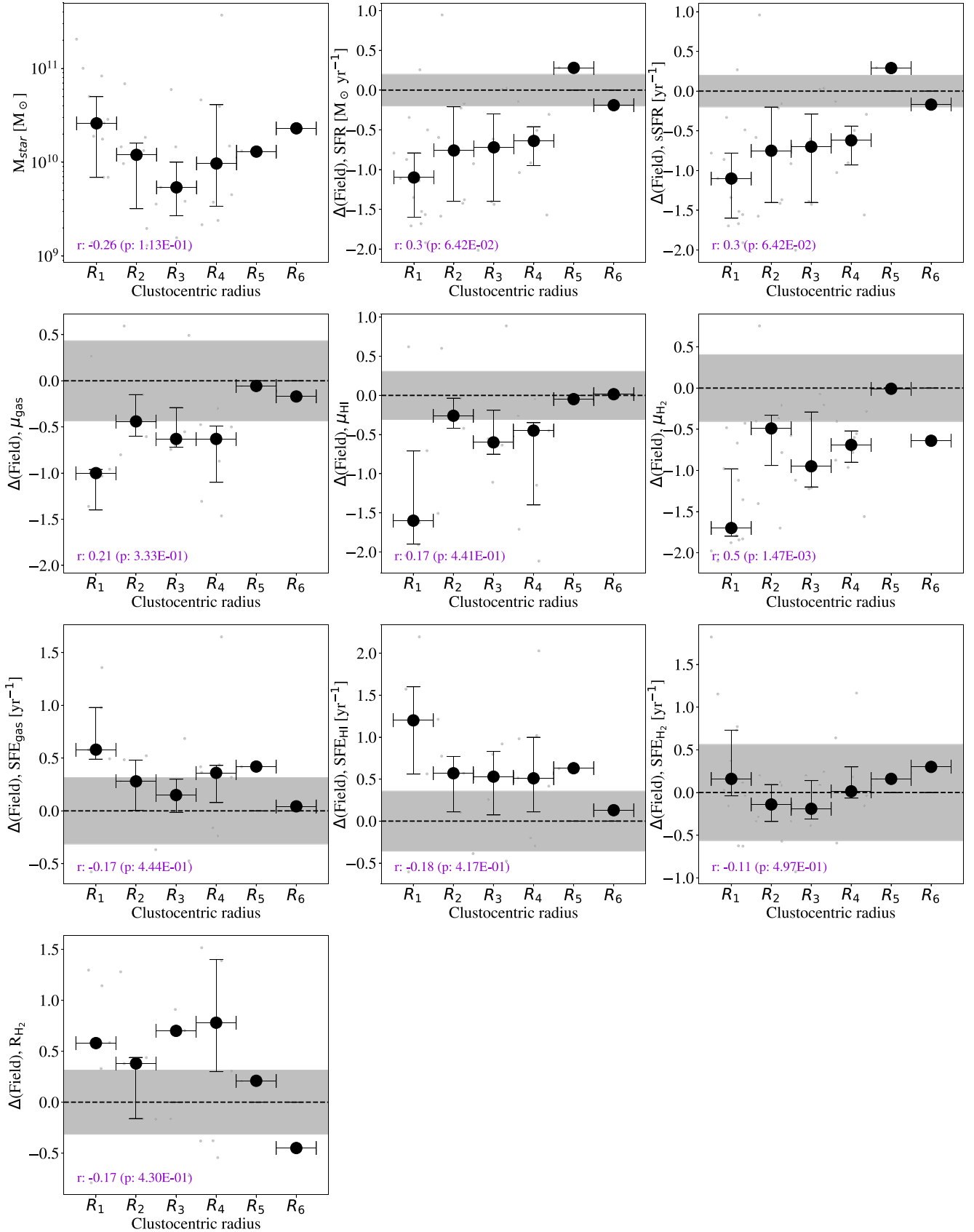


Figure 10. Radial variation of key quantities whose M_{star} dependence has been subtracted, i.e., $\Delta(\text{Field})$ values. The black filled circle and its upper/lower limits indicate the median and first/third quartiles, respectively. The individual data are indicated as gray filled circles. Zero of each key quantity indicates the values of the field galaxies (dashed line), and the gray-shaded regions indicate the 1σ of the $\Delta(\text{Field})$ values for the field galaxies. Note that the numbers of galaxies used to plot in each panel are different. The Spearman's rank-order correlation coefficient r and the p -value are shown on the lower left corner of each panel in purple. Only $\Delta(\text{Field}), \mu_{\text{H}_2}$ significantly correlates to the clustrocentric distance.

more strongly depend on the local galaxy number density than the other parameters of the parent clusters, galaxy–galaxy interactions are expected to be relatively important.

We adopt two distances to characterize the local galaxy number density: the distances to the fifth nearest galaxy (1) on the sky ($R_{5,\text{sky}}$) and (2) on the PSD ($R_{5,\text{PSD}}$). The vertical and horizontal axes of the PSD are, respectively, R_{proj}/R_{200} and $\Delta v_{\text{gal}}/\sigma_{\text{Fornax}}$, where Δv_{gal} and σ_{Fornax} are the velocity offset of a galaxy to the cluster center and the velocity dispersion of the Fornax cluster, respectively. Both values are dimensionless quantities. We assume a unity weight for both the x -axis and the y -axis of the PSD for calculating $R_{5,\text{PSD}}$. To calculate the distances to the fifth nearest galaxy, we utilize all the FCC galaxies with the information of their recession velocity in the NED database (NASA/IPAC Extragalactic Database 2019), resulting in 148 galaxies. $R_{5,\text{sky}}$ positively correlates to $R_{5,\text{PSD}}$ with a correlation coefficient of 0.32 (p -value of 8×10^{-5}). Although there are several ways to estimate the galaxy local density, the N th nearest-neighbor-based measures are considered to be the best probe of the internal density of a massive halo (Cooper et al. 2005; Muldrew et al. 2012), and N generally ranges from 3 to 10 (e.g., Dressler 1980; Gómez et al. 2003; Cooper et al. 2005; Muldrew et al. 2012).

The two distances work complementarily to evaluate the different galaxy–galaxy interactions. The tidal force between galaxies is inversely proportional to the cube of their physical distance. $R_{5,\text{sky}}$ is simply the projected distance between the galaxies on the sky, and it can be diluted by apparently close galaxy pairs whose distance along the line of sight is large. However, $R_{5,\text{PSD}}$ takes the relative line-of-sight velocity of the galaxy pairs into consideration in addition to the projected distance. However, $R_{5,\text{PSD}}$ would miss the effect of the galaxy harassment since galaxies with a large relative velocity could be neglected due to their large distance on the PSD.

Before comparing the two distances and the key quantities, $R_{5,\text{sky}}$ or $R_{5,\text{PSD}}$ needs to be compared to the projected clustocentric distance since there does exist radial tendency for some quantities (see Section 4.2). The results are shown in the Appendix. $R_{5,\text{sky}}$ and $R_{5,\text{PSD}}$ both positively correlate to the clustocentric distance with a correlation coefficient of 0.85 and 0.35, respectively, that is, it is denser near the center than at the outskirts of the cluster (Figure 38). These positive correlations suggest that the dependence of the key quantities on the clustocentric distance seen in Figure 10 might just reflect the dependence on the local galaxy number densities, and vice versa.

The comparison between the key quantities and $R_{5,\text{sky}}$ or $R_{5,\text{PSD}}$ of our sample is shown in Figures 11 and 39, respectively. $R_{5,\text{sky}}$ positively correlates to SFR, sSFR, and μ_{H_2} with correlation coefficients of ~ 0.4 . $R_{5,\text{PSD}}$ positively correlates to SFR, sSFR, and μ_{H_2} with correlation coefficients of ~ 0.4 and negatively correlates to SFE_{gas} with correlation coefficients of ~ -0.4 . Considering the stronger correlation between μ_{H_2} and the clustocentric distance than $R_{5,\text{sky}}$ or $R_{5,\text{PSD}}$ and the positive correlation between the clustocentric distance and $R_{5,\text{sky}}$ or $R_{5,\text{PSD}}$, the observed μ_{H_2} dependence on $R_{5,\text{sky}}$ or $R_{5,\text{PSD}}$ could just reflect the more intrinsic dependence of μ_{H_2} on the clustocentric distance. Following the same logic, the unique correlations between $R_{5,\text{PSD}}$ with SFE_{gas} indicates that SFE_{gas} intrinsically depends on the distances to the fifth nearest galaxy on the PSD. However, SFR and sSFR correlate to $R_{5,\text{sky}}$

and $R_{5,\text{PSD}}$ but do not correlate to the projected clustocentric distance.

4.4. Dependence on the Accretion Phase to the Cluster

The galaxy distribution on the PSD is useful to trace the accretion phase of galaxies in a cluster as suggested by previous studies (e.g., Jaffé et al. 2015; Yoon et al. 2017; Wang et al. 2021), where they show that H I gas removal in cluster galaxies can be explained by the RPS effects. The galaxies trapped by a cluster potential generally follow a “wedge” trajectory on the PSD, that is, oscillating in position and velocity, and eventually fall into the cluster core (see Figure 4 of Jaffé et al. 2015). In order to clarify the effects of the ram pressure and the other interactions with ICM or the cluster potential, the galaxies in our sample are divided into two groups according to their locations on the PSD of the cluster: (1) “virialized” (VIR) and “ram pressure stripping” (RPS) galaxies and (2) “recent infall” (RIF) galaxies. The classification is done so that the RPS and VIR galaxies are currently affected and have been affected by the ram pressure, respectively, while the RIF galaxies are unlikely to have been strongly affected by the ram pressure so far. It should be noted that the position on the PSD provides just a statistical probability that each galaxy is in a different accretion phase and is not univocally associated to the quenching mechanisms.

The VIR region is enclosed by the lines connecting the points of $(R_{\text{proj}}/R_{200}, |\Delta v_{\text{gal}}|/\sigma_{\text{Fornax}}) = (1.2, 0.0)$ and $(0.0, 1.5)$ (Jaffé et al. 2015; Yoon et al. 2017). The boundary of the RPS region is defined so that the gas in a typical $M_{\text{star}} \sim 10^9 M_{\odot}$ galaxy can be totally stripped, that is, even the gas at the galaxy center, where the anchoring force by the galaxy potential is the strongest, can be removed. Specifically, the boundary line is determined by the balance between the ram pressure,

$$P_{\text{ram}} = \rho_{\text{ICM}} v_{3\text{D,gal}}^2, \quad (4)$$

and the galaxy’s anchoring force per area,

$$\Pi_{\text{gal}} = 2\pi G \Sigma_{\text{star}} \Sigma_{\text{gas}}, \quad (5)$$

where ρ_{ICM} , $v_{3\text{D,gal}}$, G , $\Sigma_{\text{d,star}}$, and $\Sigma_{\text{d,gas}}$ are the ICM density distribution, 3D velocity of the galaxy ($v_{3\text{D,gal}} = \sqrt{3} v_{\text{obs,gal}}$), gravitational constant, surface mass density of stellar, and gas disks, respectively (Gunn & Gott 1972). The model parameters used to calculate the RPS boundary line are summarized in Table 7. The galaxies that are not in the VIR or RPS regions are classified as the RIF galaxies.

The galaxy distributions on the PSD that are color-coded for SFR, μ_{H_2} , and SFE_{H_2} are presented in Figure 12 (the one with color-coding for all the key quantities is presented in Figure 40 in the Appendix). The black solid and dotted lines indicate the boundary lines for the VIR and RPS regions, respectively. As seen in this figure, there is no RPS galaxy in our sample since the RPS region defined in this study is quite small. Note that the RPS region is wider for the less massive galaxies. We compare the cumulative histogram of the key quantities for the VIR and RIF galaxies in Figure 41. There seem to be differences between the two samples: compared to the RIF galaxies, the VIR galaxies have smaller median values of SFR, sSFR, and gas fractions and higher median values of SFE_{gas} and SFE_{H_2} . The K-S test confirms that only the difference in μ_{H_2} between VIR and RIF galaxies is significant.

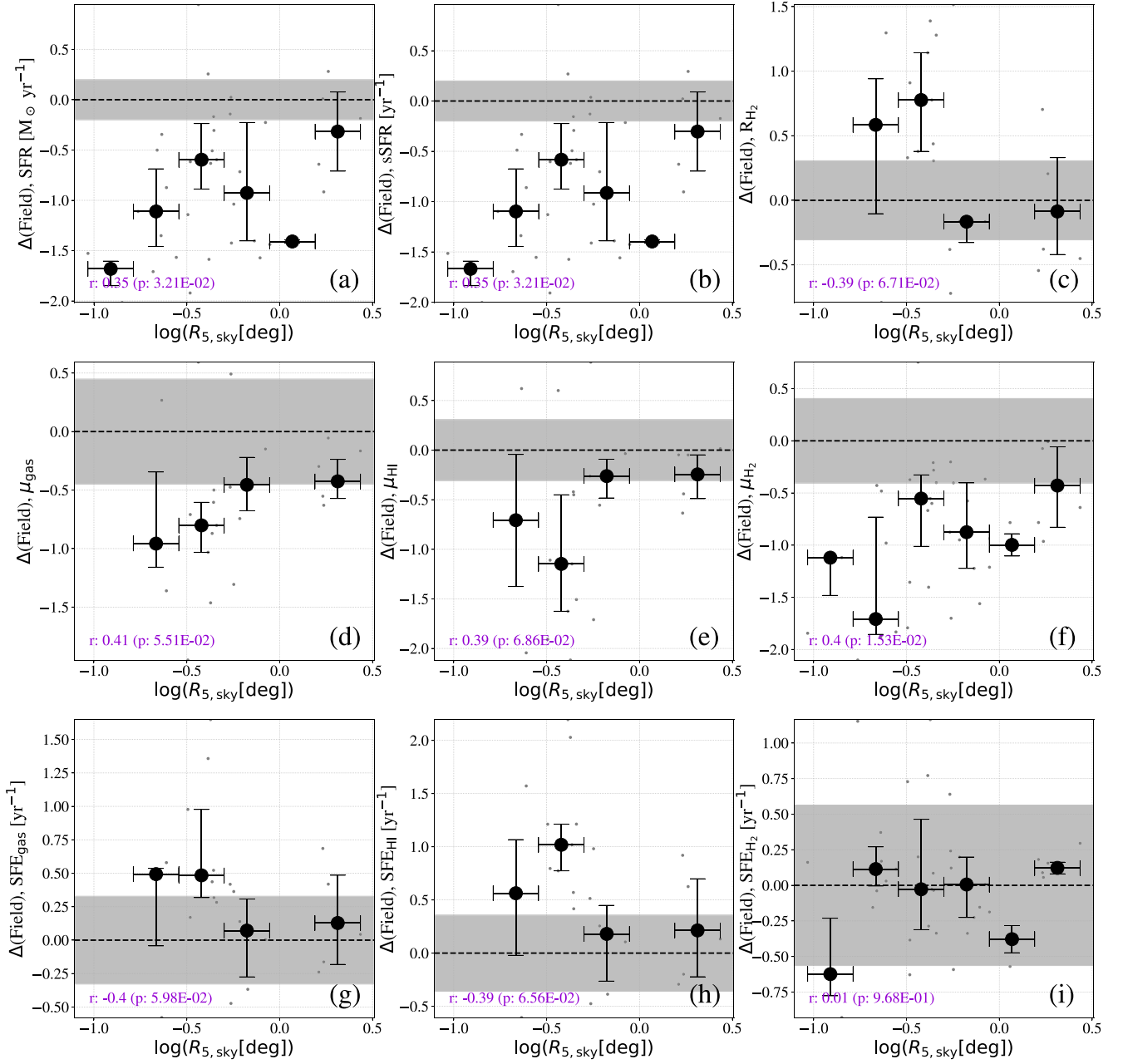


Figure 11. The relationships between $R_{5,\text{sky}}$ with the key quantities [$\Delta(\text{Field})$ values]. Symbols are the same as in Figure 10. $R_{5,\text{sky}}$ positively correlates to SFR, sSFR, μ_{gas} , and μ_{H_2} and negatively correlates to R_{H_2} .

We further study the variation of the key quantities with the product of $(|\Delta v_{\text{gal}}|/\sigma_{\text{Fomax}}) \times (R_{\text{proj}}/R_{200})$, which is claimed to be a measure of the accretion epoch. The galaxies with lower value of this quantity, that is, smaller clustocentric distance and/or smaller relative velocity with respect to the cluster, are considered to be accreted earlier. The result is shown in Figure 13. From the figure, $(|\Delta v_{\text{gal}}|/\sigma_{\text{Fomax}}) \times (R_{\text{proj}}/R_{200})$ positively correlates to SFR, sSFR, and μ_{H_2} whose correlation coefficients are among ~ 0.4 – 0.6 , that is, the ancient infallers tend to have lower SF activity and smaller molecular gas contents for their masses.

We find that $(|\Delta v_{\text{gal}}|/\sigma_{\text{Fomax}}) \times (R_{\text{proj}}/R_{200})$ positively correlates to the $R_{5,\text{sky}}$ or $R_{5,\text{PSD}}$ with correlation coefficients of 0.42 and 0.54, respectively, that is, the ancient infallers tend to reside in a denser region than the recent infallers (Figure 38 in the Appendix). The relationships between $R_{5,\text{sky}}$ and $R_{5,\text{PSD}}$ with SFR, sSFR, and μ_{H_2} would just reflect the more intrinsic

dependence of SFR, sSFR, and μ_{H_2} on the accretion phase, comparing the correlation coefficients. The null correlation between the accretion phase and μ_{gas} further supports the scenario that μ_{gas} intrinsically depends on the local galaxy number density.

5. Discussions

5.1. Sample Biases for the H I-Related Properties

As suggested at the beginning of Section 4, there may exist sample bias when we use the best-effort sample. In particular, the galaxies with H I measurements are biased to the recent infallers that are mostly SFMS galaxies with $\Delta(\text{MS}) > -1$ compared to the entire ALMA sample (Figures 2, 6, 13, and 40). Thus, we also perform all the analysis done in the previous sections for the galaxies with the measurements or detections of

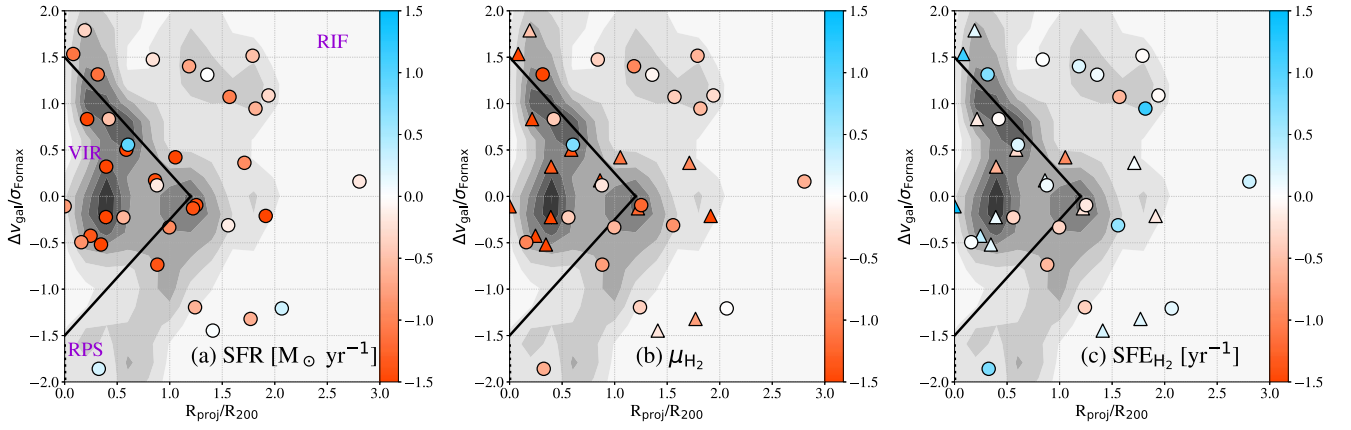


Figure 12. (a) $\Delta(\text{Field})_{\text{SFR}}$, (b) $\Delta(\text{Field})_{\mu_{\text{H}_2}}$, and (c) $\Delta(\text{Field})_{\text{SFE}_{\text{H}_2}}$ on the PSD. The gray contours indicate the FCC galaxies with the NED velocity (NASA/IPAC Extragalactic Database 2019). The triangle region enclosed by the solid black line is the VIR region. The boundary for the RPS region is indicated with the dotted line and is too close to the y-axis of the plot to be recognized in this plot. The rest of the region is the RIF region. The CO nondetections are indicated with triangles.

Table 7
ICM and Galaxy Parameters for the RPS Boundary Line

Parameters	Values	References
	ICM model ^a	
ρ_0	0.018 cm^{-3}	1
R_C	4.36 kpc	1
β	0.35	1
	Galaxy model ^b	
$M_{\text{d,star}}$	$10^9 M_\odot$	2
$M_{\text{d,gas}}$	$1.0 \times M_{\text{d,star}}$	3
$R_{\text{d,star}}$	0.9 kpc	4
$R_{\text{d,gas}}$	$1.7 \times R_{\text{d,star}}$	5

Notes.

^a Standard β -model of ICM gas density distribution (Cavaliere & Fusco-Femiano 1976): $\rho_{\text{ICM}}(r_{3\text{D}}) = \rho_0 \left[1 + \left(\frac{r_{3\text{D}}}{R_C} \right)^2 \right]^{-3\beta/2}$, where $r_{3\text{D}} = (\pi/2)R_{\text{proj}}$, ρ_0 , R_C , and β are the 3D distance from the cluster center, the number density at the center of the cluster, core radius, and the power of the distribution.

^b Exponential surface density distribution: $\Sigma_{\text{d,star or gas}}(r) = \frac{M_{\text{d,star or gas}}}{2\pi R_{\text{d,star or gas}}^2} \exp\left(-\frac{r}{R_{\text{d,star or gas}}}\right)$, where $M_{\text{d,star}}$ and $M_{\text{d,gas}}$ are the stellar and gas masses and $R_{\text{d,star}}$ and $R_{\text{d,gas}}$ are the scale lengths of the stellar and gas disks. We adopted an RPS boundary for the $M_{\text{star}} = 10^9 M_\odot$ galaxies where their entire gas (both atomic and molecular gas) is removed. The $R50$ for the $M_{\text{star}} \sim 10^9 M_\odot$ galaxies listed in the Extended Virgo Cluster Catalog (Kim et al. 2014) is converted to the scale length with an empirical relation between $R50$ and scale length of $R50 = 1.69 \times R_{\text{d,star}}$.

References. (1) Jones et al. (1997); (2) this study; (3) Hunt et al. (2015); (4) Morokuma-Matsui et al. (2021); (5) Cayatte et al. (1994).

CO and H I (the CO+H I-obs. or CO+H I-det. samples) to see if any of the obtained results changed. As a result, all the correlations with respect to SFR, sSFR, and μ_{H_2} seen in previous sections disappear in the case of CO+H I-obs. samples. However, SFE_{gas} correlates to $R_{5,\text{PSD}}$ ($r \sim -0.4$ and $p = 0.03$) with the CO+H I-obs. samples. When based on the CO+H I-det. samples, all of the correlation disappears except for the relation between R_{H_2} and $R_{5,\text{sky}}$ ($r \sim -0.5$ and $p = 0.04$).

Additionally, we investigate how the results change when assuming the upper limits of the atomic gas mass for the galaxies without the M_{atom} estimation. The H I data of our sample galaxies are retrieved from the unbiased wide-field surveys. Thus, it is a reasonable approximation to apply the

upper limits of the Fornax galaxies obtained in the surveys to other galaxies in the survey areas but not listed in the previous studies.

Loni et al. (2021) carried out the H I survey with the ATCA covering an area of 15 deg^2 roughly centered on NGC 1399, and the spatial/velocity resolutions and the H I mass sensitivity are $67'' \times 95''$ ($\sim 6 \times 9 \text{ kpc}$)/ 6.6 km s^{-1} and $\sim 4.7 \times 10^7 M_\odot$ (3σ over 300 km s^{-1}), respectively. Kleiner et al. (2021) observed the Fornax A group using the MeerKAT. The spatial/velocity resolutions and the H I mass sensitivity of the MeerKAT observations are $33''0 \times 29''2$ ($\sim 3.2 \times 2.8 \text{ kpc}$)/ 44.1 km s^{-1} and $2.9 \times 10^6 M_\odot$ (3σ over 300 km s^{-1} at the most sensitive area in case of the point source), respectively. The HIPASS covers the area where $\text{decl.} < 2^\circ$, that is, the Fornax cluster area is covered in this survey. Its spatial/velocity resolution and the H I mass sensitivity are $14'/18.0 \text{ km s}^{-1}$ and $4.7 \times 10^8 M_\odot$ (3σ over 300 km s^{-1} at the distance of 20 Mpc), respectively. All the H I mass sensitivities have been multiplied by 1.36 to account for helium.

As a simple test, we assume $4.7 \times 10^7 M_\odot$ as an upper limit of M_{atom} for the galaxies that are within the survey areas of Loni et al. (2021) and Kleiner et al. (2021) but without M_{atom} estimations in those studies and $4.7 \times 10^8 M_\odot$ as an upper limit of M_{atom} for the galaxies that are outside of the survey areas of Loni et al. (2021) and Kleiner et al. (2021). We obtain the qualitatively same results for the comparison of the key quantities between the field and the Fornax galaxies. However, the $\Delta(\text{Field})_{\mu_{\text{H I}}}$ [$\Delta(\text{Field})_{\text{SFE}_{\text{H I}}}$] becomes correlated to $\Delta(\text{Field})_{\mu_{\text{H}_2}}$ [$\Delta(\text{Field})_{\text{SFE}_{\text{H}_2}}$] with a correlation coefficient of 0.64 [0.41]. For a comparison between the key quantities and the environment parameters, we find that the μ_{gas} and $\mu_{\text{H I}}$ become correlated to all the environment parameters (clustocentric distance, $R_{5,\text{sky}}$, $R_{5,\text{PSD}}$, and the accretion phase), while SFE_{gas} and $\text{SFE}_{\text{H I}}$ remain independent of the environment parameters except for the clustocentric distance. The variable $\mu_{\text{H I}}$ most strongly correlates to the clustocentric distance (r of 0.55), unlike μ_{H_2} , which most strongly depends on the accretion phase. Furthermore, $\mu_{\text{H I}}$ and μ_{gas} ($\text{SFE}_{\text{H I}}$) of the VIR galaxies become significantly lower (higher) than those of the RIF galaxies.

These results suggest observational bias exists and that the weaker dependence of H I-related quantities compared with the H_2 ones on the environment properties when we use the best-effort sample would *not* simply indicate that the

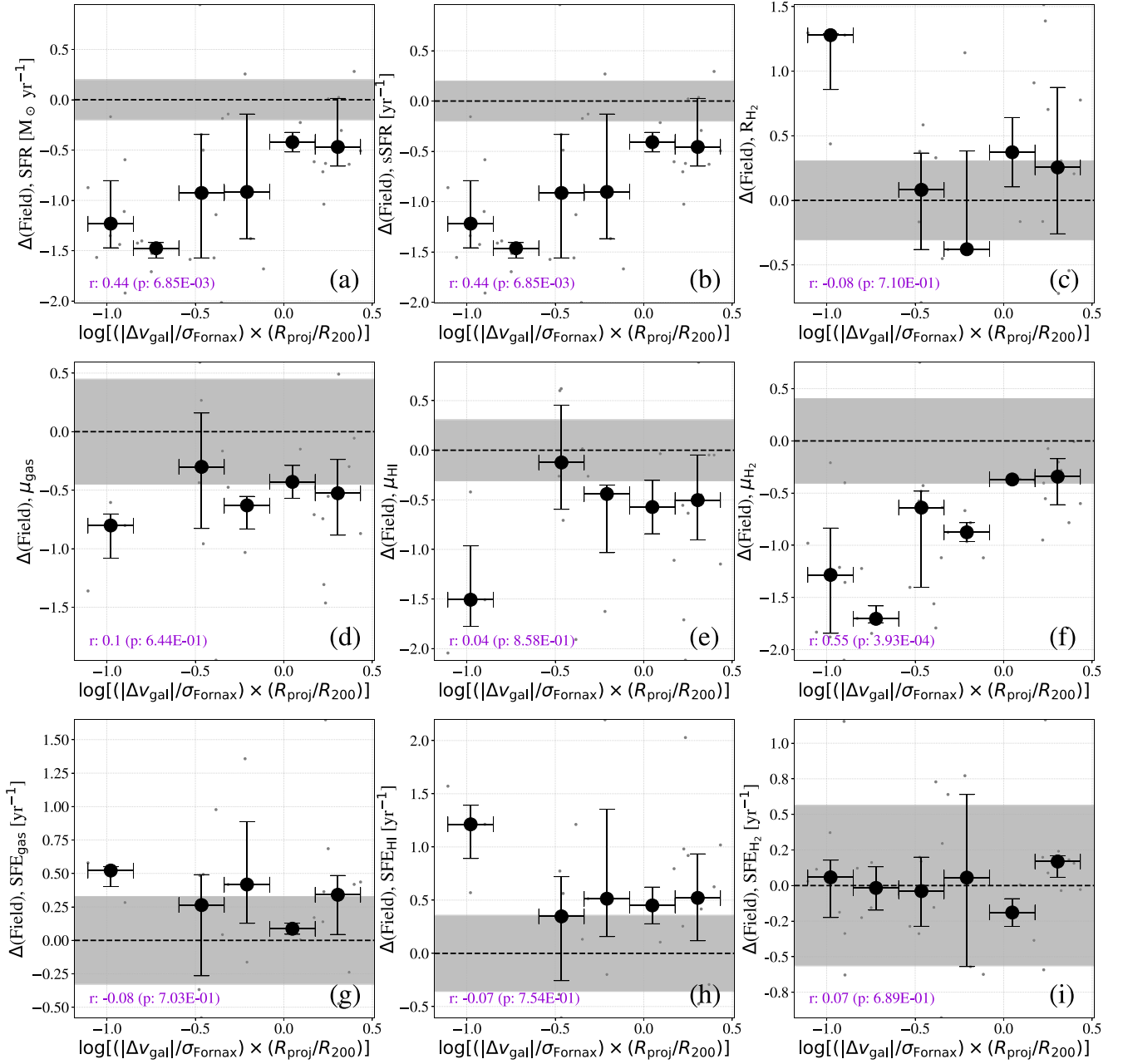


Figure 13. Relationship between the accretion phase, $(|\Delta v_{\text{gal}}|/\sigma_{\text{Fornax}}) \times (R_{\text{proj}}/R_{200})$, with the key quantities. Symbols are the same as in Figure 10. Among all the environmental parameters, SFR, sSFR, and μ_{H_2} most strongly depend on the accretion phase, where ancient infallers tend to have lower SFR, sSFR, and μ_{H_2} values than the recent infallers.

molecular gas in the Fornax galaxies is more affected by the cluster environment than their atomic gas. We speculate that the tendencies seen in the H_2 properties would be also observed in the H I properties if we could investigate the H I properties of the entire ALMA sample with M_{star} of $10^9 M_{\odot}$.

5.2. Comparison with Previous Studies on the Cold Gas Properties in the Fornax Galaxies

We find that our Fornax sample with $M_{\text{star}} > 10^9 M_{\odot}$ has lower gas fractions and higher SFE_{gas} and SFE_{HI} than the field galaxies for their stellar masses, while the difference in SFE_{H_2} between the Fornax and field galaxies is not significant. The deficiency of the cold gas reservoirs observed in this study is consistent with previous studies on the atomic and molecular

gas properties in the Fornax galaxies (Horellou et al. 1995; Zabel et al. 2019; Kleiner et al. 2021; Loni et al. 2021).

Zabel et al. (2020) have already reported that depletion time (the inverse number of SFE) of the Fornax galaxies is comparable to those of the field galaxies using ALMA CO and $\text{H}\alpha$ data obtained with the Multi-Unit Spectroscopic Explorer on the ESO Very Large Telescope (VLT). They also report that some Fornax dwarf galaxies have smaller depletion time (higher SFE) than the field galaxies.

We find a higher R_{H_2} in the Fornax galaxies than in the field galaxies and that the Fornax galaxies show a wider range of R_{H_2} than the field galaxies as claimed in Loni et al. (2021) and Kleiner et al. (2021). These studies discuss that the wide R_{H_2} range indicates that there are various galaxies in different stages

Table 8
Derived Physical Quantities of Sample Fornax Galaxies

Name	$\log M_{\text{star}}$ (M_{\odot})	M_{star} Method ^a	$\log \text{SFR}$ ($M_{\odot} \text{ yr}^{-1}$)	SFR Method ^b	$\log M_{\text{mol}}^c$ (M_{\odot})	$\log M_{\text{atom}}$ (M_{\odot})	M_{atom}	References ^d
ESO 302-14	7.94 ± 0.16	11	-1.27 ± 0.20	11	<7.54	8.92 ± 0.10		3
ESO 302-9	8.66 ± 0.10	11	-1.12 ± 0.20	12	<7.42	9.04 ± 0.08		3
ESO 357-25	8.70 ± 0.10	11	-1.80 ± 0.20	11	<7.66	8.12 ± 0.09		3
ESO 358-16	7.40 ± 0.08	4	-2.36 ± 0.02	4	<7.62	7.98 ± 0.18		1
ESO 358-20	8.92 ± 0.13	11	-1.51 ± 0.20	11	<7.64	7.40 ± 0.10		3
ESO 358-5	8.77 ± 0.10	11	-1.38 ± 0.20	11	<7.56	8.53 ± 0.19		3
ESO 358-51	9.15 ± 0.10	11	-1.04 ± 0.20	11	7.56 ± 0.07	8.13 ± 0.12		1
ESO 358-G015	8.62 ± 0.10	11	-1.91 ± 0.20	11	<7.60	8.09 ± 0.14		1
ESO 358-G063	10.12 ± 0.11	11	-0.29 ± 0.20	11	8.78 ± 0.02	9.36 ± 0.11		1
ESO 359-2	9.20 ± 0.10	12	-1.95 ± 0.20	13	6.84 ± 0.12
ESO 359-3	9.57 ± 0.10	11	-1.25 ± 0.20	11	<7.58	8.66 ± 0.11		3
ESO 359-5	7.98 ± 0.06	4	-1.84 ± 0.03	4	<7.43	8.83 ± 0.11		3
FCC 102	7.49 ± 0.07	4	-2.34 ± 0.02	4	<7.52	7.83 ± 0.12		1
FCC 117	6.34 ± 0.29	4	-3.37 ± 0.10	4	<7.68
FCC 120	7.57 ± 0.06	4	-2.36 ± 0.02	4	<7.63	7.74 ± 0.15		1
FCC 177	9.84 ± 0.15	12	-2.12 ± 0.23	13	<7.50
FCC 198	7.39 ± 0.11	1	-3.35 ± 0.30	1	<7.57
FCC 206	8.40 ± 0.07	4	-2.07 ± 0.03	4	<7.59
FCC 207	8.44 ± 0.03	1	-2.80 ± 0.24	1	<7.66	<7.13		1
FCC 261	8.48 ± 0.07	4	-1.99 ± 0.03	4	<7.68	<7.13		1
FCC 302	7.73 ± 0.27	11	-1.48 ± 0.20	11	<7.62	9.17 ± 0.11		1
FCC 306	7.89 ± 0.05	1	-1.76 ± 0.04	1	<7.79	8.17 ± 0.11		1
FCC 316	8.01 ± 0.05	1	-2.35 ± 0.12	2	<7.67
FCC 32	8.67 ± 0.03	1	-1.87 ± 0.06	2	<7.47
FCC 332	8.61 ± 0.03	5	-2.16 ± 0.04	5	6.82 ± 0.11	<7.13		1
FCC 44	7.76 ± 0.06	4	-3.14 ± 0.06	4	<7.56
IC 1993	10.16 ± 0.10	11	-0.66 ± 0.20	11	8.81 ± 0.01	8.31 ± 0.14		3
IC 335	9.99 ± 0.10	11	-1.71 ± 0.20	11	<7.61
MCG-06-08-024	8.77 ± 0.01	1	-1.66 ± 0.02	1	6.80 ± 0.12
MCG-06-09-023	9.11 ± 0.10	12	-1.94 ± 0.20	13	7.16 ± 0.09
NGC 1310	9.65 ± 0.10	11	-0.60 ± 0.20	11	8.36 ± 0.02	8.70 ± 0.01		2
NGC 1316	11.57 ± 0.10	11	0.02 ± 0.20	11	8.74 ± 0.03	7.85 ± 0.04		2
NGC 1316C	9.33 ± 0.10	11	-1.49 ± 0.20	11	7.83 ± 0.04	7.25 ± 0.07		2
NGC 1317	10.59 ± 0.10	11	-0.34 ± 0.20	11	8.87 ± 0.01	8.46 ± 0.00		2
NGC 1326	10.66 ± 0.10	11	0.06 ± 0.20	11	8.62 ± 0.02	9.38 ± 0.01		2
NGC 1326A	8.84 ± 0.10	11	-0.98 ± 0.20	11	<7.53	9.17 ± 0.03		2
NGC 1326B	9.06 ± 0.11	11	-0.56 ± 0.20	11	<7.67	9.67 ± 0.01		2
NGC 1336	9.78 ± 0.10	11	-1.63 ± 0.20	12	<7.56
NGC 1339	10.17 ± 0.10	11	-1.61 ± 0.20	11	<7.63
NGC 1341	9.58 ± 0.10	11	-0.32 ± 0.20	11	8.48 ± 0.01	8.48 ± 0.16		3
NGC 1350	10.77 ± 0.10	11	-0.46 ± 0.20	11	8.59 ± 0.06	9.13 ± 0.16		3
NGC 1351	10.27 ± 0.10	11	-1.41 ± 0.20	11	<7.56
NGC 1351A	9.55 ± 0.11	11	-1.27 ± 0.20	11	7.94 ± 0.08	8.83 ± 0.12		1
NGC 1365	10.84 ± 0.10	11	1.24 ± 0.20	11	10.30 ± 0.00	10.31 ± 0.12		1
NGC 1374	10.42 ± 0.10	11	-1.44 ± 0.20	11	<7.53
NGC 1375	9.84 ± 0.10	11	-1.88 ± 0.20	11	<7.74
NGC 1379	10.41 ± 0.10	11	-1.27 ± 0.20	11	<7.49
NGC 1380	10.92 ± 0.10	11	-0.81 ± 0.20	11	7.77 ± 0.03	<7.83		1
NGC 1381	10.28 ± 0.10	11	-1.69 ± 0.20	11	<7.56
NGC 1386	10.25 ± 0.10	11	0.26 ± 0.20	11	8.58 ± 0.02	<7.83		1
NGC 1387	10.70 ± 0.10	11	-0.65 ± 0.20	11	8.53 ± 0.01	<7.61		1
NGC 1399	11.31 ± 0.10	11	-0.27 ± 0.20	11	<7.50
NGC 1404	11.00 ± 0.10	11	-0.74 ± 0.20	11	<7.46
NGC 1406	10.12 ± 0.10	11	0.22 ± 0.20	12	9.13 ± 0.01	9.35 ± 0.09		3
NGC 1419	9.73 ± 0.10	11	-2.27 ± 0.21	11	<7.62
NGC 1425	10.36 ± 0.10	11	-0.13 ± 0.20	11	8.70 ± 0.03	9.52 ± 0.13		3
NGC 1427	10.46 ± 0.10	11	-1.46 ± 0.20	11	<7.57
NGC 1427A	9.14 ± 0.17	11	-0.89 ± 0.20	11	<7.50	9.46 ± 0.12		1
NGC 1436	10.16 ± 0.10	11	-0.64 ± 0.20	11	8.78 ± 0.02	7.91 ± 0.20		1
NGC 1437A	8.75 ± 0.10	11	-1.07 ± 0.20	11	<7.45	8.91 ± 0.10		1
NGC 1437B	9.29 ± 0.11	11	-0.64 ± 0.20	12	7.97 ± 0.05	8.51 ± 0.13		1
NGC 1484	9.38 ± 0.02	1	-1.07 ± 0.02	1	<7.51	8.94 ± 0.04		3
PGC 012625	7.45 ± 0.27	1	-2.46 ± 0.30	1	<7.67	8.52 ± 0.13		3

Table 8
(Continued)

Name	$\log M_{\text{star}}$ (M_{\odot})	M_{star} Method ^a	$\log \text{SFR}$ ($M_{\odot} \text{ yr}^{-1}$)	SFR Method ^b	$\log M_{\text{mol}}$ (M_{\odot})	$\log M_{\text{atom}}$ (M_{\odot})	M_{atom} References ^d
WOMBAT I	8.10 ± 0.04	4	-1.74 ± 0.02	4	< 7.71	8.33 ± 0.13	3

Notes.

^a Double-digit and single-digit numbers indicate the M_{star} values from Leroy et al. (2019) and those derived in this study, respectively. For the values from Leroy et al. (2019): 11, SSFRLIKE; 12, W4W1; 13, FIXEDUL. See Table 6 in Leroy et al. (2019) for details. For the values derived in this study, we determine M/L with sSFR values estimated as follows: 1, as 11 with SFR derived using far UV (FUV)+W4; 2, as 11 with SFR derived using NUV+W4; 3, as 12 with W4-to-W1 color; 4, 5, and 6, with Equation (24) in Leroy et al. (2019) using FUV-to-W1 color, NUV-to-W1 color, and W3-to-W1 color, respectively.

^b As in note *a*, double-digit and single-digit numbers indicate SFR values from Leroy et al. (2019) and those derived in this study, respectively. For the values from Leroy et al. (2019), SFR was determined using 11, FUV+W4; 12, NUV+W4; 13, W4. See Table 7 in Leroy et al. (2019) for details. For the values derived in this study, we determine SFR using 1, FUV+W4; 2, NUV+W4; 3, W4; 4, FUV; 5, NUV, and 6, W3.

^c The constant CO-to-H₂ conversion factor of $4.36 M_{\odot} (\text{K km s}^{-1} \text{ pc}^2)$ is adopted for all the galaxies.

^d References for the M_{atom} data, (1) Loni et al. (2021); (2) Kleiner et al. (2021); (3) Hyper-LEDA.

of the gas removal/consumption processes in the Fornax cluster.

5.3. Star Formation Quenching in the Fornax Galaxies

Amount of gas or conversion efficiency from gas to stars?—When limiting to the galaxies at the lower $\Delta(\text{MS})$ regime, we find that the CO- or H I-detected Fornax galaxies with $M_{\text{star}} > 10^9 M_{\odot}$ have lower gas fractions (μ_{gas} , $\mu_{\text{H I}}$ and μ_{H_2}) and higher SFE_{gas} , $\text{SFE}_{\text{H I}}$, and R_{H_2} than the field galaxies (Figure 8). SFE_{H_2} of the Fornax galaxies is higher than or comparable to that of the field galaxies at $\Delta(\text{MS}) \sim -1.0$, although the difference is statistically insignificant. In addition, even for the SFMS galaxies, μ_{H_2} in the Fornax galaxies is lower than that in the field galaxies. This suggests that the decrease in the cold gas reservoir is more important for the SF quenching in the Fornax galaxies than the decrease of SFE (our Fornax sample has higher SFE than the field galaxies!). Furthermore, the H I gas reservoir is more heavily reduced than the H₂ gas reservoir in the low $\Delta(\text{MS})$ Fornax galaxies.

Timescale of the gas removal—Our results also imply that the cold gas reservoir in the Fornax galaxies with $\Delta(\text{MS}) < 0$ is decreased (stripped or consumed) on a timescale shorter than the typical gas-depletion timescale by SF in the field SF galaxies of $M_{\text{gas}}/\text{SFR} \sim 3$ Gyr, that is, SFR does not have enough time to change according to the change in the amount of the cold gas. This is incompatible with the strangulation, whose typical timescale is claimed to be ~ 4 Gyr (Peng et al. 2015). A similar timescale for the gas depletion is estimated in a different way in Loni et al. (2021) based on the blind H I survey of the Fornax cluster. They find that the H I-deficient galaxies tend to distribute in the VIR region on the PSD, and they consequently found that the H I removal occurs on a timescale shorter than the cluster’s crossing time of $R_{200}/\sigma_{\text{Fornax}} \sim 2$ Gyr. However, for the galaxies in the cluster outskirts, they do not find a significant offset from the field relation between the $\Delta(\text{Field})_{\text{SFR}}$ and $\Delta(\text{Field})_{\mu_{\text{H I}}}$ (Figure 11 in Loni et al. 2021), suggesting that the SFR has had sufficient time to respond to a variation in H I mass as the field galaxies. In the same figure, they also find that some H I-nondetected Fornax galaxies would have lower $\mu_{\text{H I}}$ than the field galaxies for fixed $\Delta(\text{MS})$. The shortage of the H I gas can lead to the reduction of the H₂ gas reservoirs in galaxies without a significant enhancement in SFE_{H_2} .

Ram pressure gas stripping—One of the (most famous) fast processes removing gas from galaxies in a cluster is ISM RPS.

Schröder et al. (2001) discussed that the H I gas deficiency in the Fornax galaxies is caused by the RPS based on the H I gas survey toward 66 Fornax galaxies with the Parkes 64 m telescope. They find that the H I-deficient galaxies tend to have small velocity dispersion compared with the other galaxies, suggesting that they have more radial orbits than the other galaxies. Loni et al. (2021) detected H I emission from 16 Fornax galaxies located outside the VIR region, and they find that eight of the 16 galaxies show disturbed H I morphology with their $\sim 6 \text{ kpc} \times 9 \text{ kpc}$ ATCA beam. Among the eight H I disturbed galaxies, five galaxies (ESO 358-G063, NGC 1351A, NGC 1365, NGC 1427A, and NGC 1437B) have $> 10^9 M_{\odot}$ in our study. They discussed the RPS as the process responsible for the disturbed H I morphologies of ESO 358-G063 and NGC 1351A among the five galaxies. Thus, H I gas of some Fornax galaxies with $M_{\text{star}} > 10^9 M_{\odot}$ could be removed by the ram pressure from the ICM.

However, Horellou et al. (1995) consider that the ISM RPS is not as effective as in the Virgo cluster, given that the X-ray emission in the Fornax cluster is weaker than that of the Virgo cluster. In general, molecular gas is expected to be less affected by the ram pressure of ICM compared with atomic gas (Cortese et al. 2021, and references therein) since molecular gas resides within optical disks of galaxies, whereas atomic gas often distributes far out of the optical disks. However, Zabel et al. (2019) find that the low-mass Fornax galaxies with $< 3 \times 10^9 M_{\odot}$ have disturbed molecular gas morphology and such galaxies tend to have a molecular gas deficit. This may possibly suggest that some low-mass galaxies are affected by the ram pressure of the ICM. Zabel et al. (2021) find that the Fornax galaxies generally have a lower gas-to-dust ratio than the field galaxies due to the depletion of both the H I and H₂ gas even though the dust mass is also reduced compared to the field galaxies. They discuss various possible explanations for these findings including the more efficient stripping of H₂ compared to dust, more efficient enrichment of dust in the SF process, and altered ISM physics in the cluster environment. As one of the possible explanations, they suggest a combination of the dust formation and H₂ consumption by ongoing SF but inefficient H₂ formation owing to the H I gas depletion due to the strangulation or stripping.

Our data cannot completely reject the ISM RPS as a dominant physical process responsible for the reduction of molecular gas reservoirs in the Fornax galaxies studied here. The short-timescale atomic gas stripping by the ram pressure is rather compatible with our results. Galaxies with a low gas

fractions and high SFEs are naturally explained by RPS, provided that galaxies generally tend to have high SFE near the galactic center (a steeper radial gradient for earlier spiral galaxies; Villanueva et al. 2021) and that ram pressure preferentially strips outskirts components in galaxies (Bekki 2014; Cortese et al. 2016). Interestingly, NGC 1351A has a tentative elongated CO structure to the southeast direction that is consistent with the H I data in Loni et al. (2021), whereas there is no clear disturbance in the WISE 3.4 μm image (Figure 28). This would suggest the ram pressure strips not only H I but also H_2 gas in NGC 1351A. A detailed comparison of morphologies (and kinematics) of the stellar, atomic gas, and molecular gas components of the galaxy is required to confirm this possibility.

As another clue to the dominant quenching processes in the Fornax cluster, we find that $\Delta(\text{Field})_{\text{SFR}}$ and $\Delta(\text{Field})_{\mu_{\text{H}_2}}$ more strongly correlate to the accretion phase to the cluster than the local galaxy number density and the projected clustocentric distance. However, $\Delta(\text{Field})_{\mu_{\text{H I}}}$ most strongly depends on the projected clustocentric distance to which $\Delta(\text{Field})_{\text{SFR}}$ does not significantly correlate (Section 5.1). These results suggest that the reduction of H I gas reservoirs in galaxies does not directly lead to their SF quenching and is partially attributed to the RPS. The effect of RPS would blur the dependence of key quantities on the accretion phase since the galaxies that are being affected by ram pressure could have a small R_{proj}/R_{200} and a large $|\Delta v_{\text{gal}}|/\sigma_{\text{Fornax}}$, that is, distribute at the upper left corner on the $|\Delta v_{\text{gal}}|/\sigma_{\text{Fornax}}$ versus R_{proj}/R_{200} plot. The strongest dependence of SFR and μ_{H_2} on the accretion phase among all the environmental parameters explored here basically suggests the importance of long-timescale processes such as the interaction with the cluster potential and the strangulation for SF quenching in the Fornax galaxies. However, it is also possible that the reduction of H I gas due to the RPS drives the H_2 gas depletion and consequently SF quenching in the Fornax galaxies.

Tidal interaction—The tidal interaction is expected to strip gas in the galaxy outskirts as well as enhance SF activity in a galaxy by inducing gas inflow to the galaxy center (e.g., Noguchi & Ishibashi 1986; Noguchi 1988; Byrd & Valtonen 1990; Hernquist & Mihos 1995) and possibly accelerate the gas depletion in the galaxy along with the strangulation. We find that some Fornax galaxies with $\Delta(\text{MS}) < 0$ tend to have lower μ_{H_2} and higher SFE_{H_2} than the field galaxies. The depletion timescale of molecular gas ($1/\text{SFE}_{\text{H}_2}$) in some samples with $\Delta(\text{MS}) \lesssim -1.5$ is as short as $\lesssim 1$ Gyr (Figure 37). Since our sample is currently forming stars less actively, their cold gas reservoirs would have been already depleted when SF is enhanced by the tidal interaction or they would be in the later stage of the tidal interaction. Although the morphological analysis will be presented in a separate paper, there are galaxies with a ring-like gas structure (Figures 3, 18, and 34, 35), which is predicted in the later stage of the tidal interaction based on the numerical simulations (Noguchi & Ishibashi 1986).²⁹

We approximate the tidal radius of each galaxy as $r_{\text{tidal}} = R_{\text{proj}} \left(\frac{M_{\text{galaxy}}}{3M_{\text{cluster}}} \right)^{1/3}$ and compare it with the galaxy size. The tidal radius of a galaxy is defined as the point at which galactic stars/ISM escape the galaxy’s potential well and become part of the cluster. We adopt $M_{\text{galaxy}} = 10 \times M_{\text{star}}$,

$M_{\text{cluster}} = 5 \times 10^{13} M_{\odot}$ (Drinkwater et al. 2001) and the semimajor axis of the ellipse aperture for the WISE and GALEX photometry as the galaxy size (r_{CAAPR}). The $r_{\text{tidal}}/r_{\text{CAAPR}}$ ratio mostly ranges from zero to five, and the median value is 2.7. Four out of the 38 galaxies have $r_{\text{tidal}}/r_{\text{CAAPR}} < 1.0$. Given that the galaxies often have from a few to several times more extended H I disk than optical disk (e.g., Walter et al. 2008; Koribalski et al. 2018), H I gas in the galaxy outskirts would have been tidally stripped by the cluster potential.

The dependence of the key quantities on the $r_{\text{tidal}}/r_{\text{CAAPR}}$ ratio is also investigated. We find that $r_{\text{tidal}}/r_{\text{CAAPR}}$ positively correlates to μ_{H_2} with a correlation coefficient of 0.35. Although galaxies with high SFE_{H_2} tend to have smaller $r_{\text{tidal}}/r_{\text{CAAPR}}$, SFE_{H_2} does not significantly depend on the $r_{\text{tidal}}/r_{\text{CAAPR}}$ ($r \sim -0.3$ and $p = 0.06$). Considering $r_{\text{tidal}}/r_{\text{CAAPR}} \propto R_{\text{proj}}$ and the stronger dependence of μ_{H_2} on R_{proj} ($r = 0.5$), the correlation between $r_{\text{tidal}}/r_{\text{CAAPR}}$ and μ_{H_2} may just reflect the more intrinsic relation between R_{proj} and μ_{H_2} . The $\mu_{\text{H I}}$ and $\text{SFE}_{\text{H I}}$ of galaxies with M_{atom} measurements in the previous studies do not depend on $r_{\text{tidal}}/r_{\text{CAAPR}}$. However $\mu_{\text{H I}}$ and $\text{SFE}_{\text{H I}}$ correlate to $r_{\text{tidal}}/r_{\text{CAAPR}}$ with correlation coefficients of 0.42 and -0.32 , respectively, if we assume the M_{atom} upper limit as we did in Section 5.1. These correlation coefficients are smaller than and the same as those for the $R_{\text{proj}}-\mu_{\text{H I}}$ relation (0.55) and the $R_{\text{proj}}-\text{SFE}_{\text{H I}}$ relation (-0.32), respectively. Although the impact of the tidal effect from the cluster potential on galaxies is not so large in the current state, it may have boosted the SFEs and help to decrease gas fractions of some galaxies in the process of galaxy accretion to the cluster.

The caveats here are that we use $10 \times M_{\text{star}}$ instead of the halo mass and R_{proj} instead of the distance at the pericenter. To calculate r_{tidal} , we should use halo mass that is generally ~ 100 times larger than the stellar mass of galaxies depending on the halo mass (e.g., Behroozi et al. 2010, 2013). However, it is difficult to estimate the dark matter mass of cluster galaxies since it is predicted that substantial quantities of dark matter are tidally stripped before stellar stripping begins (galaxies that lose $\sim 80\%$ of their dark matter lose only $\sim 10\%$ of their stars; Smith et al. 2016). Therefore r_{tidal} estimated here should be considered to be a measure of the true tidal radius.

Preprocesses in the group environment—The preprocesses in the group environment would also be important for Fornax galaxies. Theoretical studies predict that many cluster galaxies (20%–40%) had been accreted as a group of galaxies (McGee et al. 2009; De Lucia et al. 2012; Benavides et al. 2020) and the preprocesses have nonnegligible impact on their evolution (but see also Berrier et al. 2009). The galaxy groups within clusters would be disassembled and virialized into the host cluster after the first pericentric passage depending on their mass concentrations (Taffoni et al. 2003). Kleiner et al. (2021) investigated the cold gas properties of the galaxies in the Fornax A group located at $\sim 2 R_{\text{vir}}$ from NGC 1399. The Fornax A group is considered to be falling into the main cluster for the first time. They find that most member galaxies they explored show signs of the preprocesses in the group environment such as H I tails, truncated H I disk, and H I deficiency. Loni et al. (2021) find that the possible members of the NGC 1365 subgroup (Drinkwater et al. 2001) also show H I deficiency.

²⁹ A ring-like structure in H α is also predicted in some galaxies that are affected by ram pressure (Bekki 2014).

Even within the virial radius of the Fornax cluster, Iodice et al. (2019a) identified three groups of galaxies: the old core, the clump on the north-northwest side of the cluster, and a group of galaxies that fell in more recently. Iodice et al. (2019b) find that the most bright early-type galaxies (ETGs) reside in the former two regions and there are intracluster baryons (diffuse light and globular clusters; Iodice et al. 2017) that are indicative of the tidal interaction. They discuss that the diffuse morphology of the intracluster lights (ICLs) and their relative locations to the bright ETGs suggest the gravitational interaction between galaxies rather than the galaxy–cluster interaction.

Combined with the previous studies, our results suggest that the atomic gas in the Fornax galaxies has been stripped tidally and/or by the ram pressure in a short time and that the (slightly) boosted SFEs due to the tidal interaction would accelerate the atomic and molecular gas depletion along with the strangulation. In the early epoch of the cluster formation, the preprocesses in the group environment would have already reduced a certain amount of the cold gas and SF activity in some galaxies in the core of the cluster (Iodice et al. 2019a).

6. Summary

We observed 64 Fornax galaxies with the ALMA Morita array in cycle 5, and CO emission is detected from 23 out of the 64 galaxies (Figure 3). Combined with the ancillary data of stellar mass, SFR, and atomic gas mass, we investigated atomic and molecular gas–related properties of the massive Fornax galaxies ($M_{\text{star}} > 10^9 M_{\odot}$). Our main results are as follows:

Fornax versus field (Section 4.1): Compared to field galaxies, the Fornax galaxies have lower SFR, sSFR, and both atomic and molecular gas fractions and higher SFE from atomic gas, suggesting that the reduction of the gas reservoirs is more essential than the SFE reduction for the quenching in the Fornax galaxies. In particular, the Fornax galaxies with low Δ (MS) tend to have low atomic and molecular gas fractions and high SFE from atomic gas compared with field galaxies. This suggests the atomic gas reduction processes occurred on a timescale shorter than the typical gas-depletion timescale by SF in field galaxies. The Fornax galaxies have a higher R_{H_2} median and a wider range of R_{H_2} compared with the field galaxies. This suggests that the H I gas is more depleted than the H_2 gas and that there are various galaxies in different stages of the gas stripping. However, the galaxies with H I measurements are biased to SF galaxies, and the obtained properties related to H I gas would not represent the common feature of our ALMA sample. The ongoing H I survey with MeerKAT will shed light on this problem (Serra et al. 2019).

Clustocentric distance (Section 4.2): The galaxies nearer to the cluster center have low μ_{H_2} (the Spearman rank-order correlation coefficient r of 0.5). The μ_{H_2} median of outskirt galaxies is already slightly lower than those of the field galaxies.

Local galaxy density (Section 4.3): The distance to the fifth nearest neighbor on the sky ($R_{5,\text{sky}}$) positively correlates to SFR, sSFR, and μ_{H_2} ($r \sim 0.4$). In the case of the distance to the fifth nearest neighbor on the PSD ($R_{5,\text{PSD}}$), there are correlations with SFR, sSFR, and μ_{H_2} with correlation coefficients of ~ 0.4 and a negative correlation with SFE_{gas} with a correlation coefficient of ~ -0.4 . SFR and sSFR more strongly depend on $R_{5,\text{sky}}$ or $R_{5,\text{PSD}}$ than the clustocentric

distance. However, μ_{H_2} dependence on the local density is expected to just reflect the dependence on the clustocentric distance given the positive correlation between $R_{5,\text{sky}}$ or $R_{5,\text{PSD}}$ with the clustocentric distance and the stronger dependence of μ_{H_2} on the clustocentric distance than on $R_{5,\text{sky}}$ and $R_{5,\text{PSD}}$.

Accretion phase (Section 4.4): The product of $(|\Delta_{\text{gal}}|/\sigma_{\text{Fornax}}) \times (R_{\text{proj}}/R_{200})$ positively correlates to SFR, sSFR, and μ_{H_2} (r of ~ 0.4 – 0.6), suggesting that galaxies that were accreted to the Fornax cluster at an earlier epoch tend to have lower SF activity and molecular gas contents. SFR, sSFR, and μ_{H_2} more strongly correlate to $(|\Delta_{\text{gal}}|/\sigma_{\text{Fornax}}) \times (R_{\text{proj}}/R_{200})$ than the other environment-related quantities.

SF quenching processes in Fornax (Section 5.3): Our results suggest that decreasing gas fractions is more important than decreasing SFEs for SF quenching in the Fornax galaxies with $M_{\text{star}} > 10^9 M_{\odot}$. In addition, the atomic gas is considered to be stripped tidally and by the ram pressure, which would lead to the depletion of molecular gas and eventually low SF activity in the galaxies. The galaxy–cluster interaction would accelerate molecular-gas deficiency by increasing SFE along with the strangulation, while preprocesses in the group environment would have reduced the molecular gas reservoirs of the Fornax galaxies in the early phase of the cluster formation. Detailed comparison of the morphology of the stellar and cold gas components in galaxies is required to shed light on the direct effect of the tidal interaction and the RPS on the molecular gas contents in the Fornax galaxies.

Further studies with larger statistics and sample homogeneity, as well as high resolution for cold gas, SFR, and stellar components, would be required to determine the dominant quenching process in the Fornax cluster.

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Facilities: GALEX, WISE, ALMA, ATCA, MeerKAT, Parkes.

Software: astropy (Astropy Collaboration et al. 2013), APLpy (Robitaille & Bressert 2012).

Appendix

In this section, supplementary plots are presented: Figures 14–36 show the CO moment maps of the 23 CO-detected galaxies out of the 64 sample galaxies; Figure 37 shows the dependence of the gas fractions, R_{H_2} , and SFE on the distance from the main sequence of star-forming galaxies; Figure 38 shows the relationships between the clustocentric distance with $R_{5,sky}$ and $R_{5,PSD}$ and the relationships between the accretion phase with $R_{5,sky}$ and $R_{5,PSD}$; Figure 39 shows the relationship between the $R_{5,PSD}$ with the key quantities; Figure 40 shows the PSD color coded for the key quantities explored in this study; Figure 41 compares the key quantities of the VIR and the RIF galaxies.

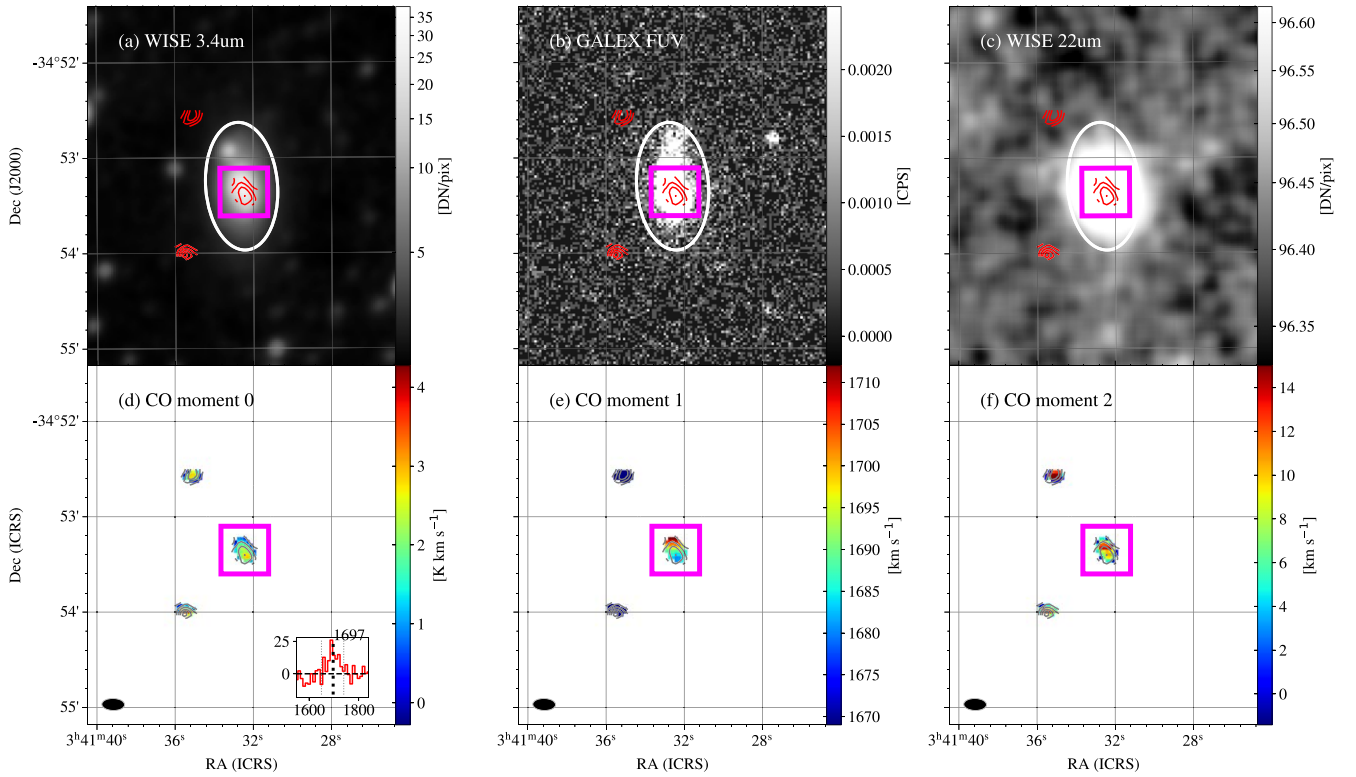


Figure 14. Ancillary images and CO moment maps of ESO 358-51: (a) WISE 3.4 μm , (b) GALEX FUV, (c) WISE 22 μm , (d) CO integrated-intensity map, (e) CO velocity map, and (f) CO velocity-dispersion map. The CO spectrum is also shown in panel (d). The vertical and horizontal axes of the CO spectrum plot are the brightness temperatures in units of millikelvin and the velocity in units of kilometers per second, respectively. The black and gray dotted lines on the CO spectrum panel indicate the galaxy systemic velocity and the velocity range in which the molecular gas mass is calculated, respectively. The red and gray contours indicate CO integrated-intensity maps. The aperture ellipse for the WISE and GALEX data to calculate M_{star} and SFR is indicated with a white solid line. The magenta line indicates the area to generate total CO spectra (shown in panel (d)) to calculate M_{mol} .

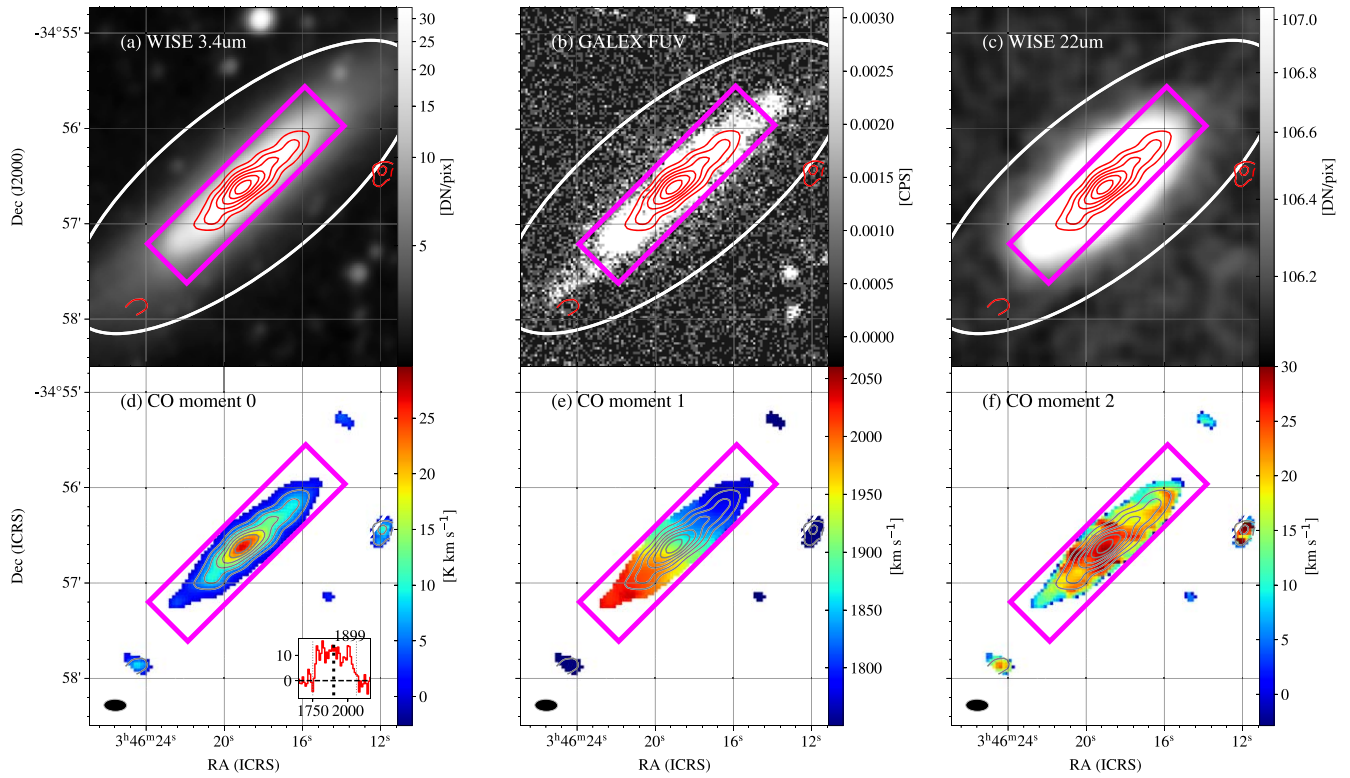


Figure 15. Same as Figure 14, but for ESO 358-G063.

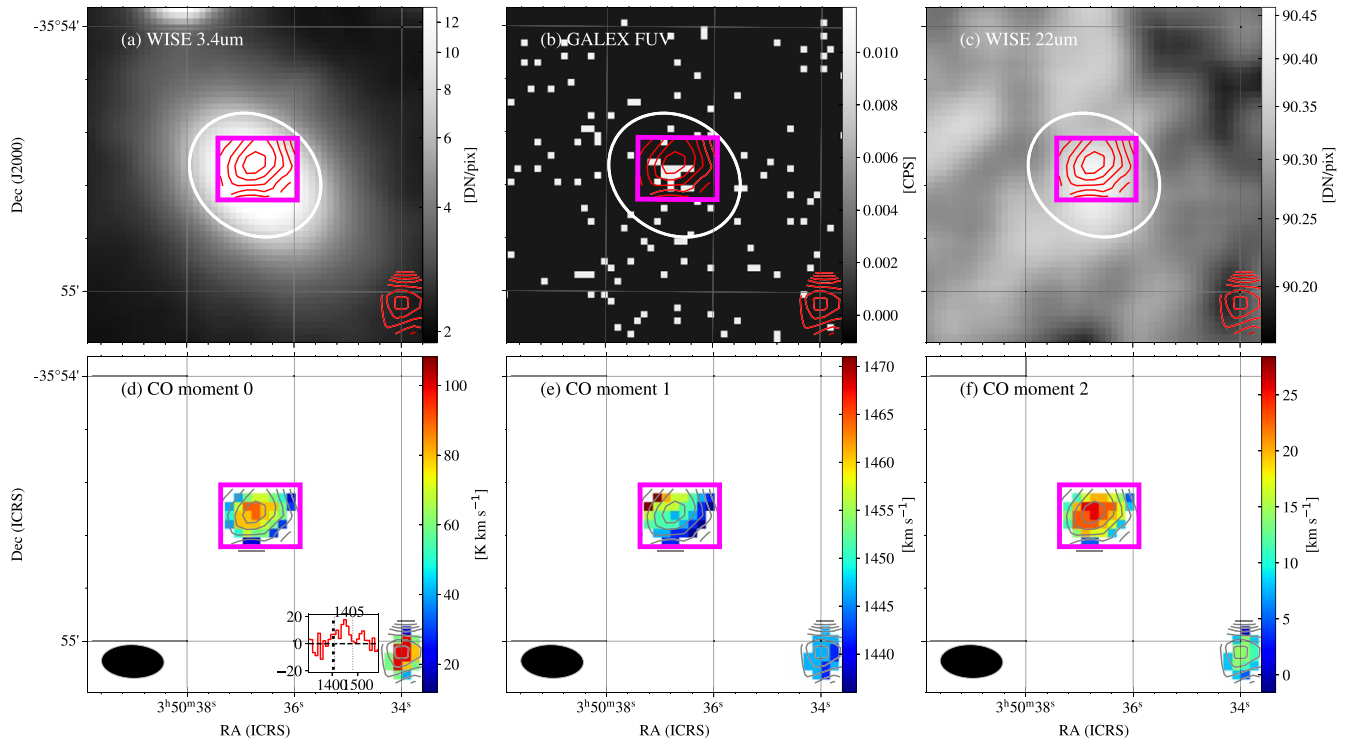


Figure 16. Same as Figure 14, but for ESO 359-2.

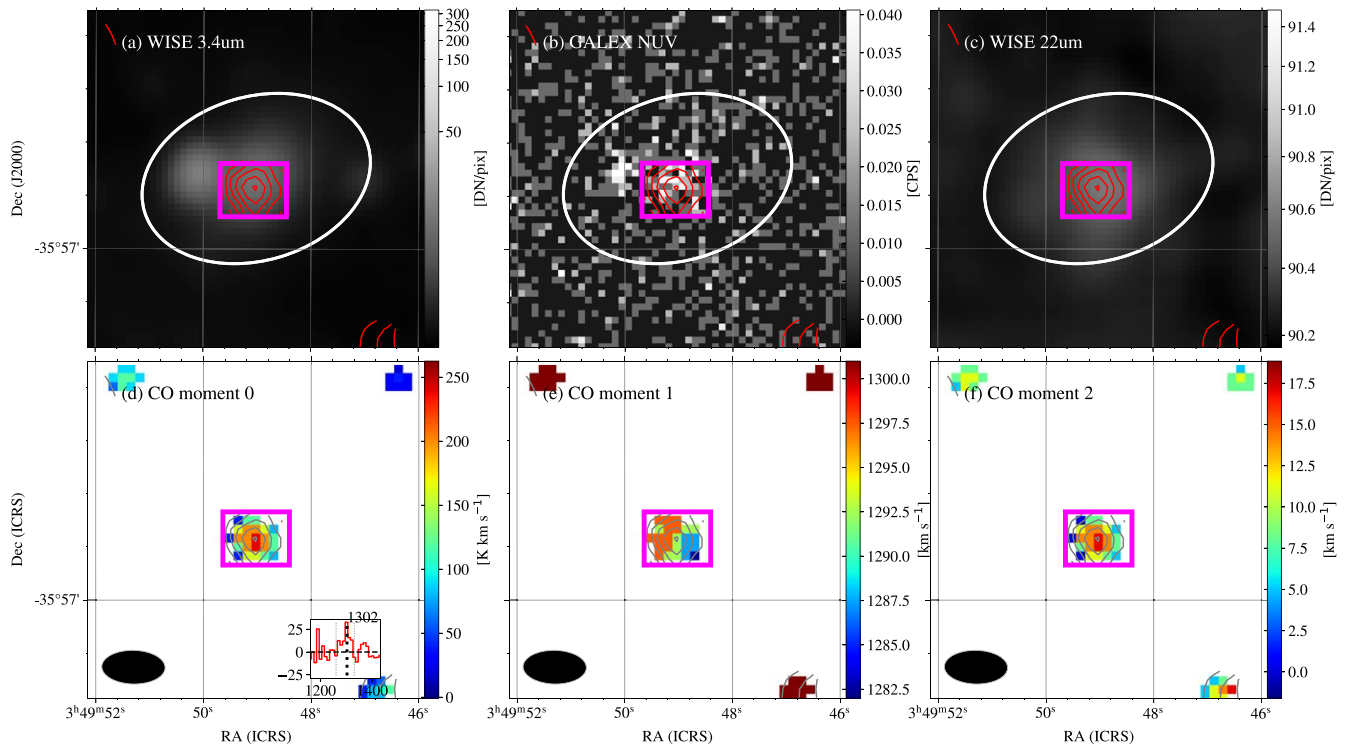


Figure 17. Same as Figure 14, but for FCC 332.

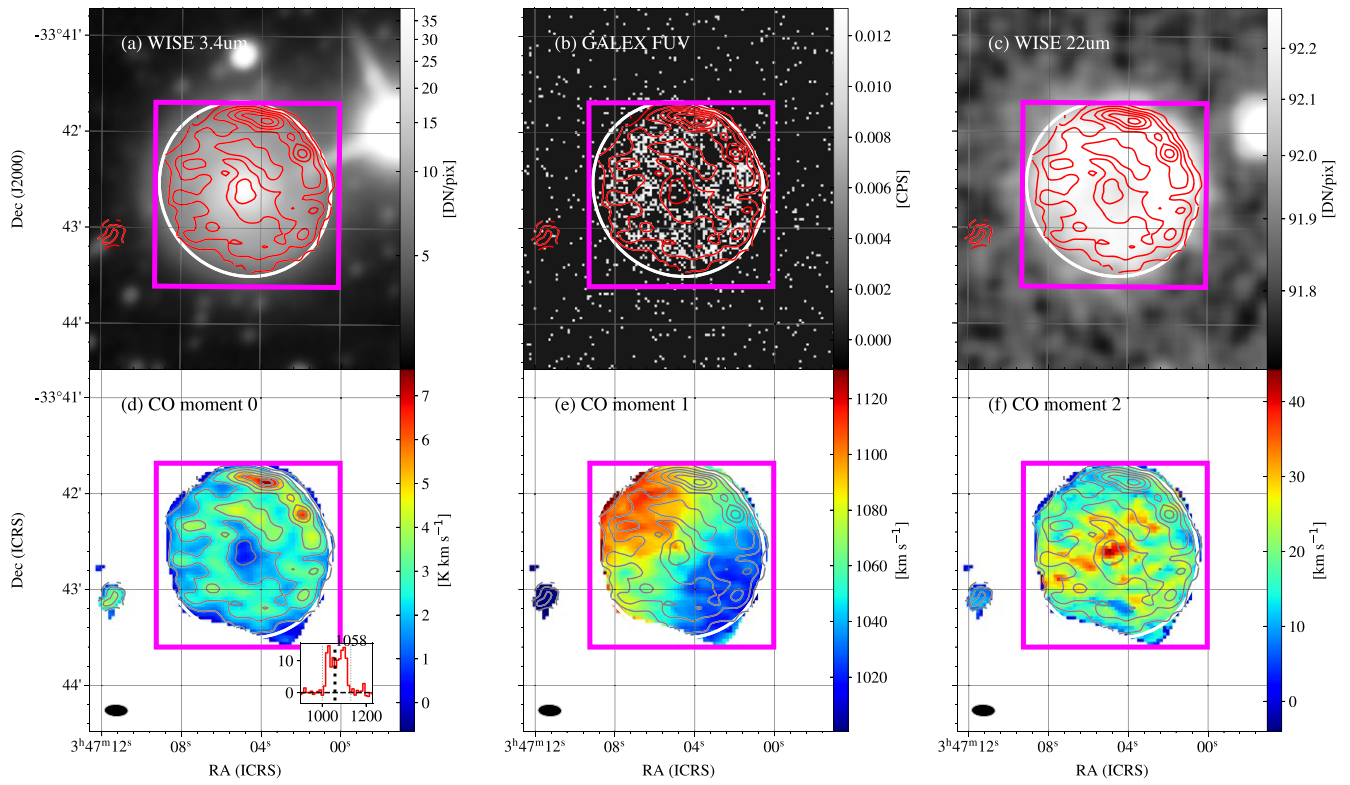


Figure 18. Same as Figure 14, but for IC 1993.

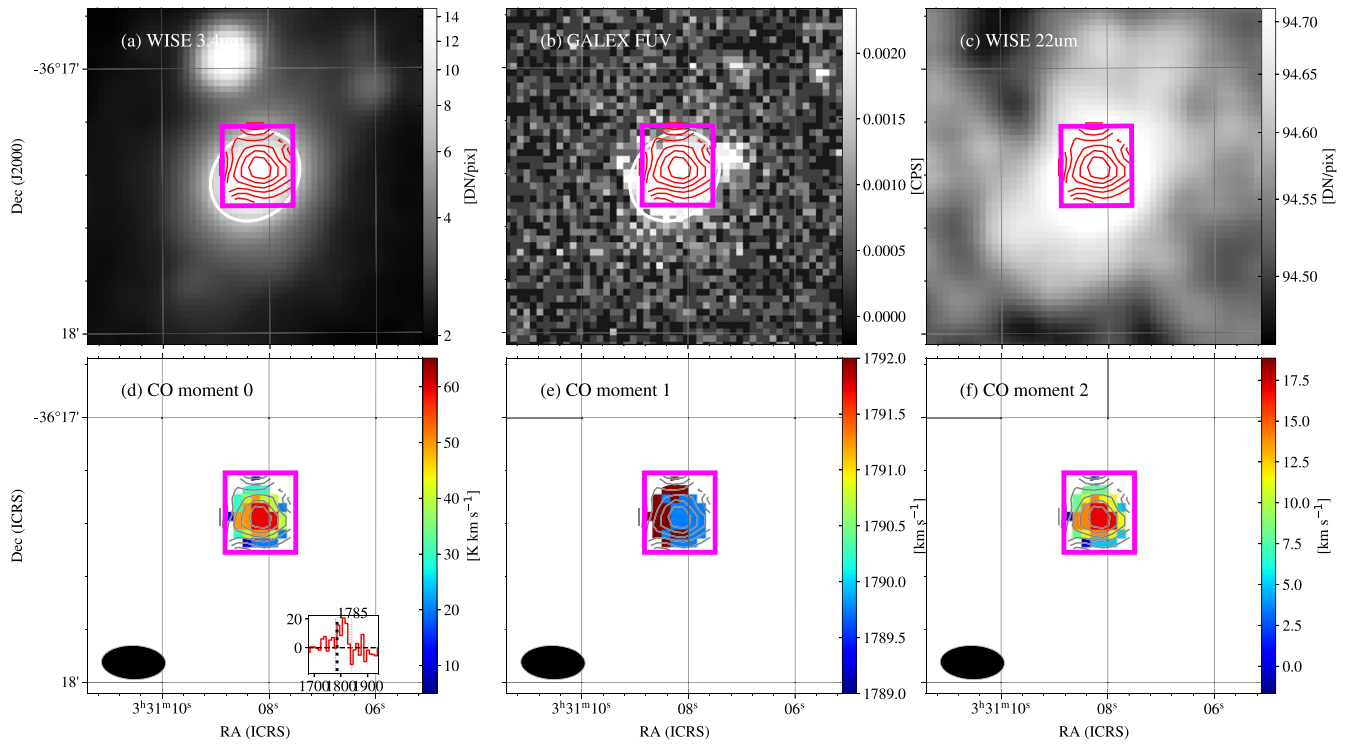


Figure 19. Same as Figure 14, but for MCG-06-08-024.

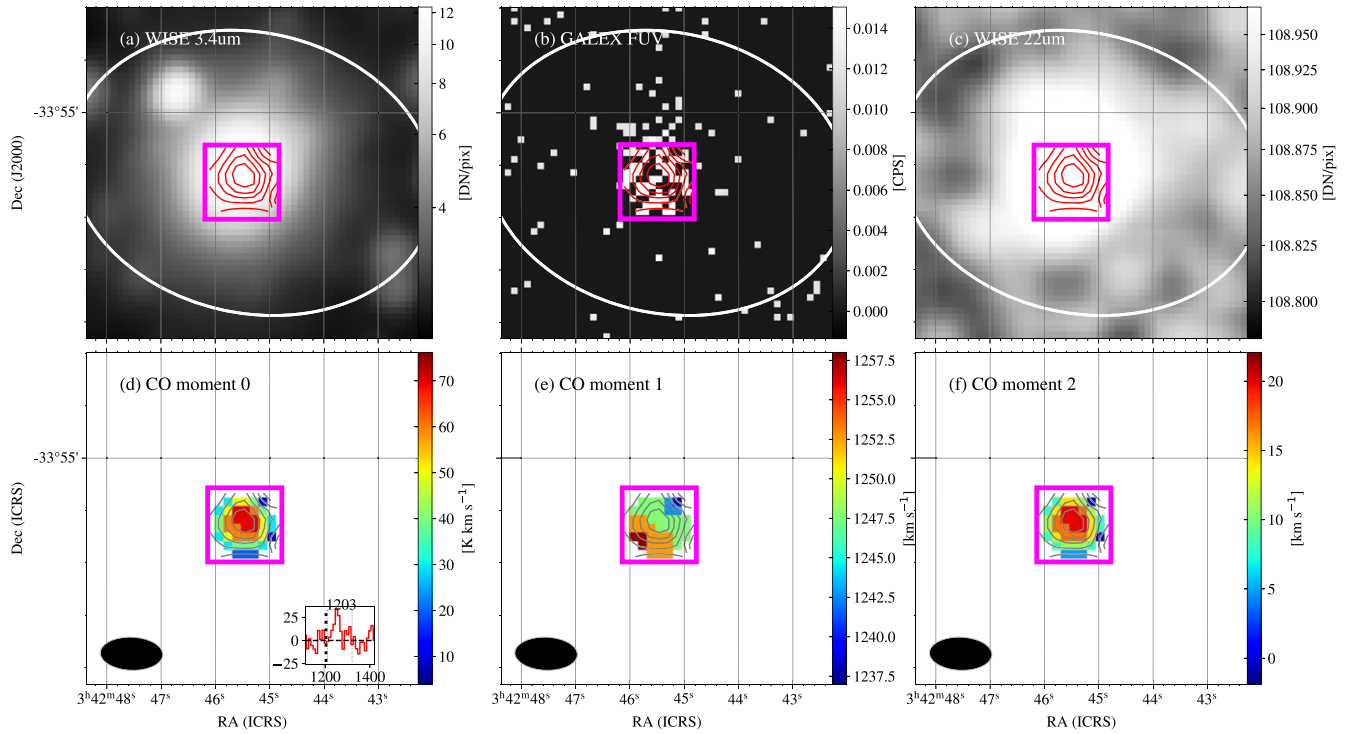


Figure 20. Same as Figure 14, but for MCG-06-09-023.

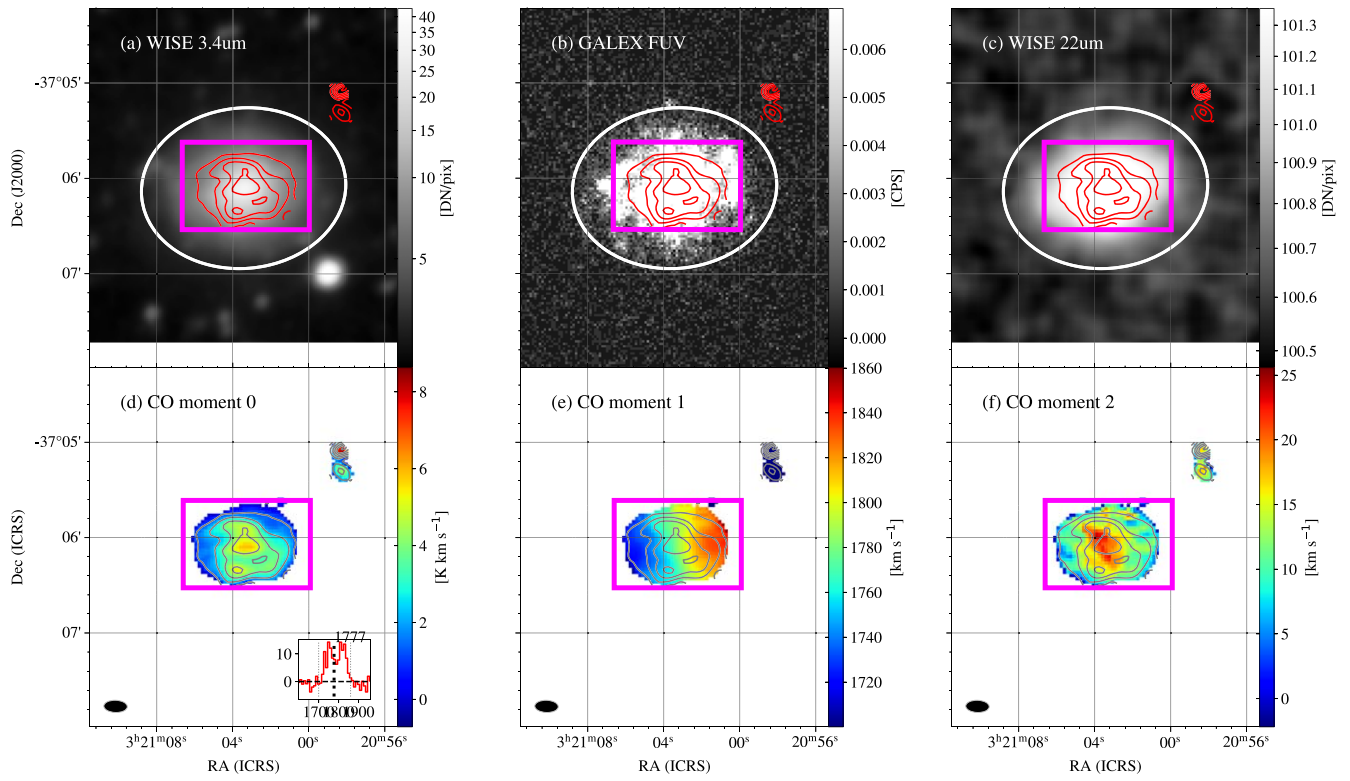


Figure 21. Same as Figure 14, but for NGC 1310.

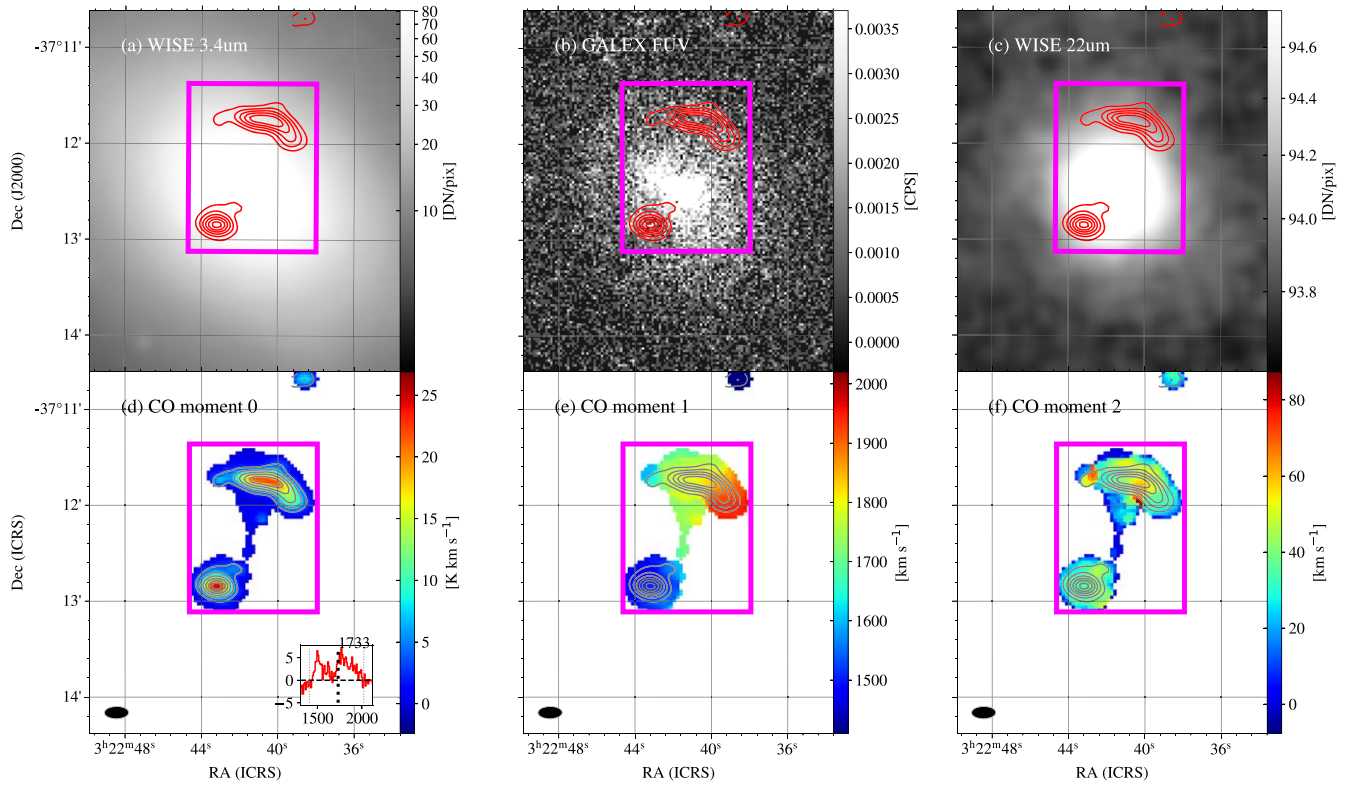


Figure 22. Same as Figure 14, but for NGC 1316.

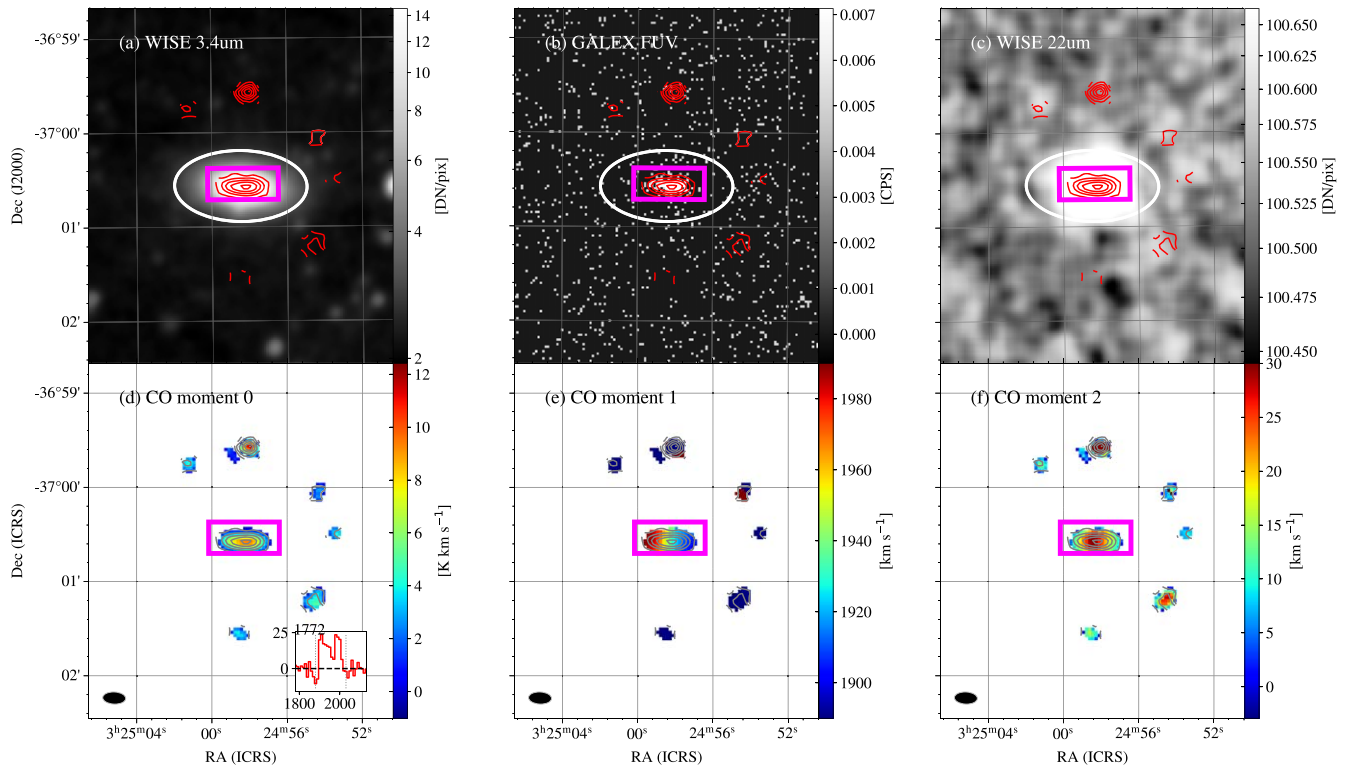


Figure 23. Same as Figure 14, but for NGC 1316C.

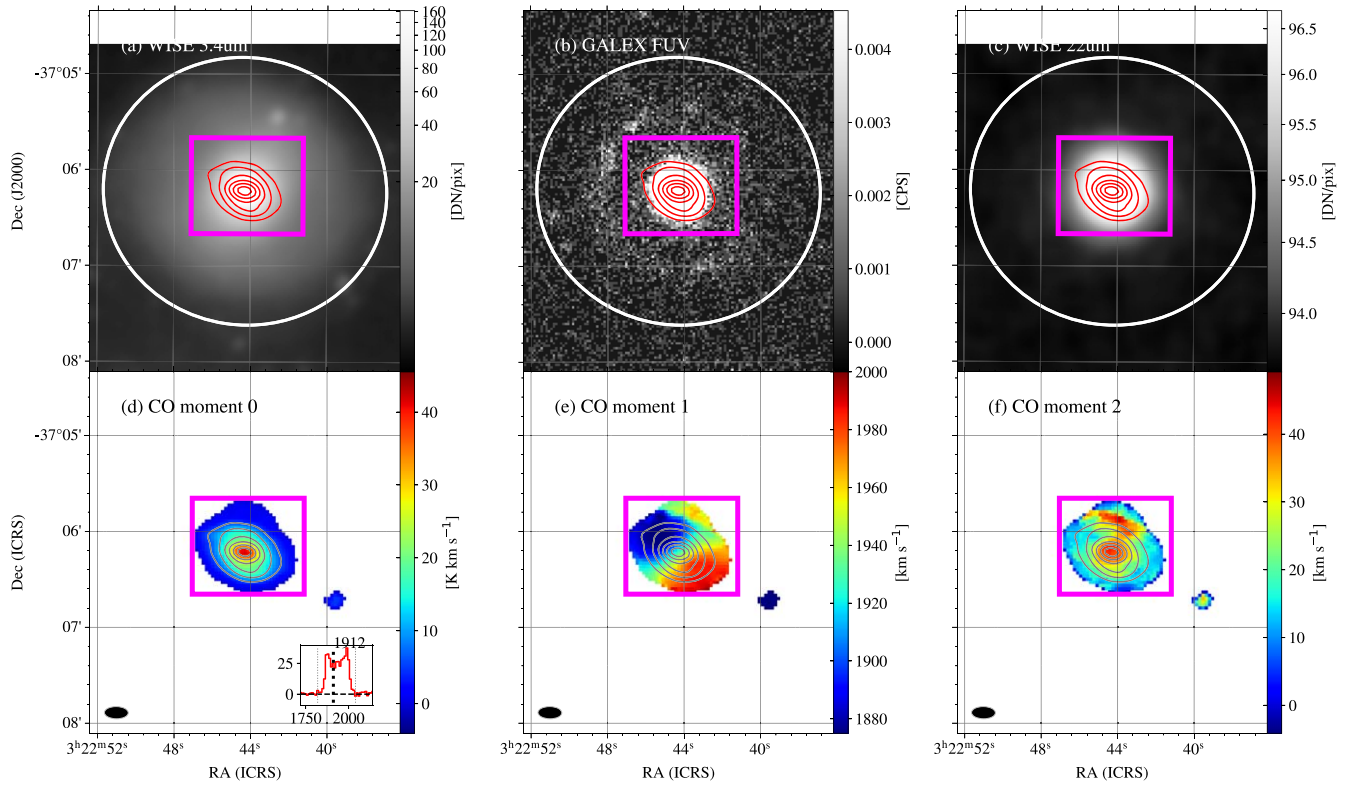


Figure 24. Same as Figure 14, but for NGC 1317.

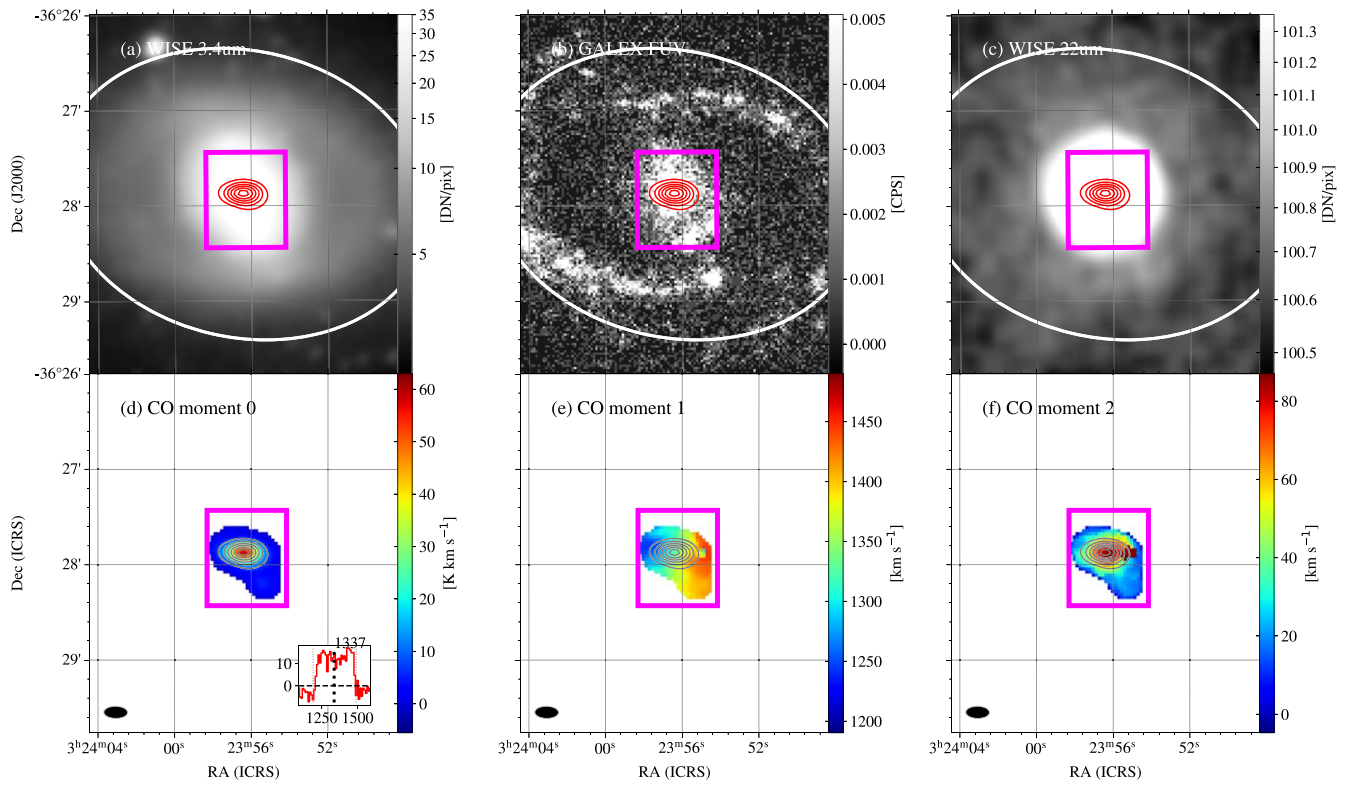


Figure 25. Same as Figure 14, but for NGC 1326.

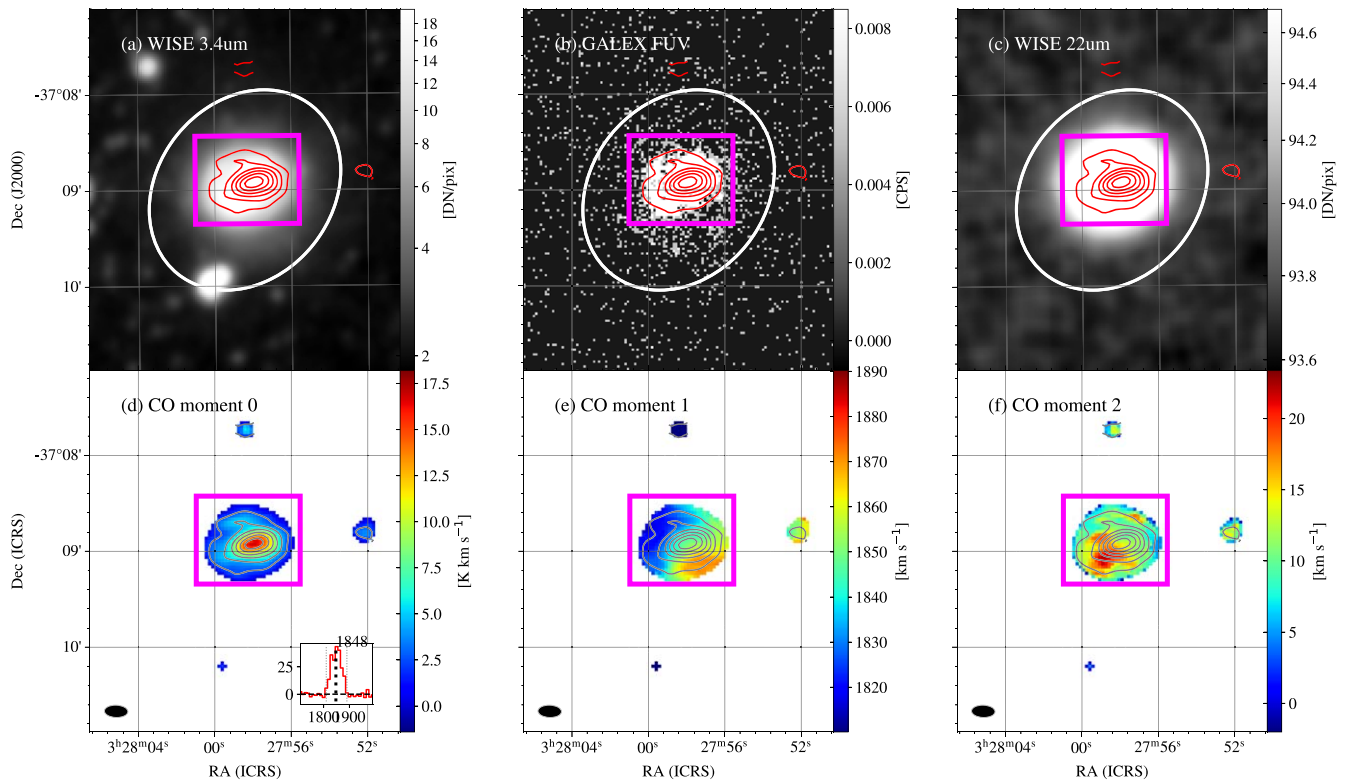


Figure 26. Same as Figure 14, but for NGC 1341.

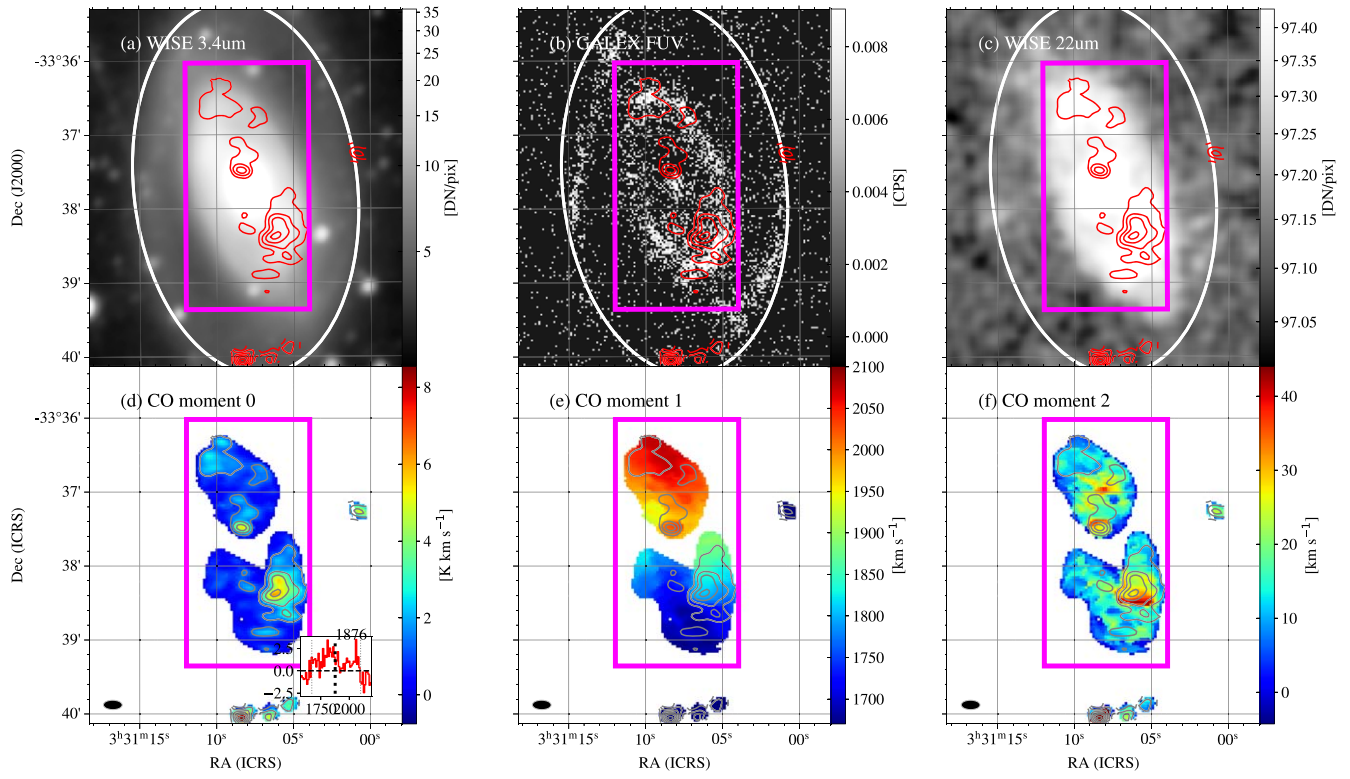


Figure 27. Same as Figure 14, but for NGC 1350.

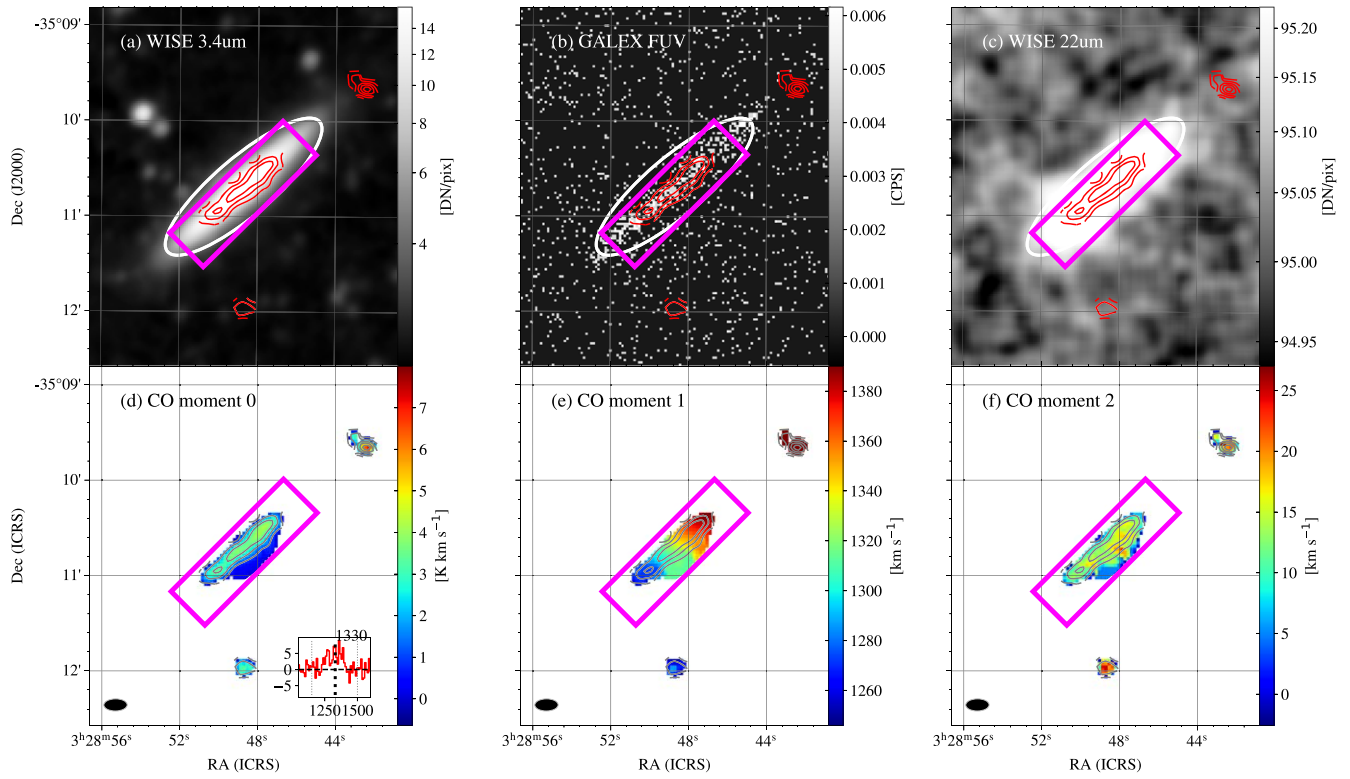


Figure 28. Same as Figure 14, but for NGC 1351A.

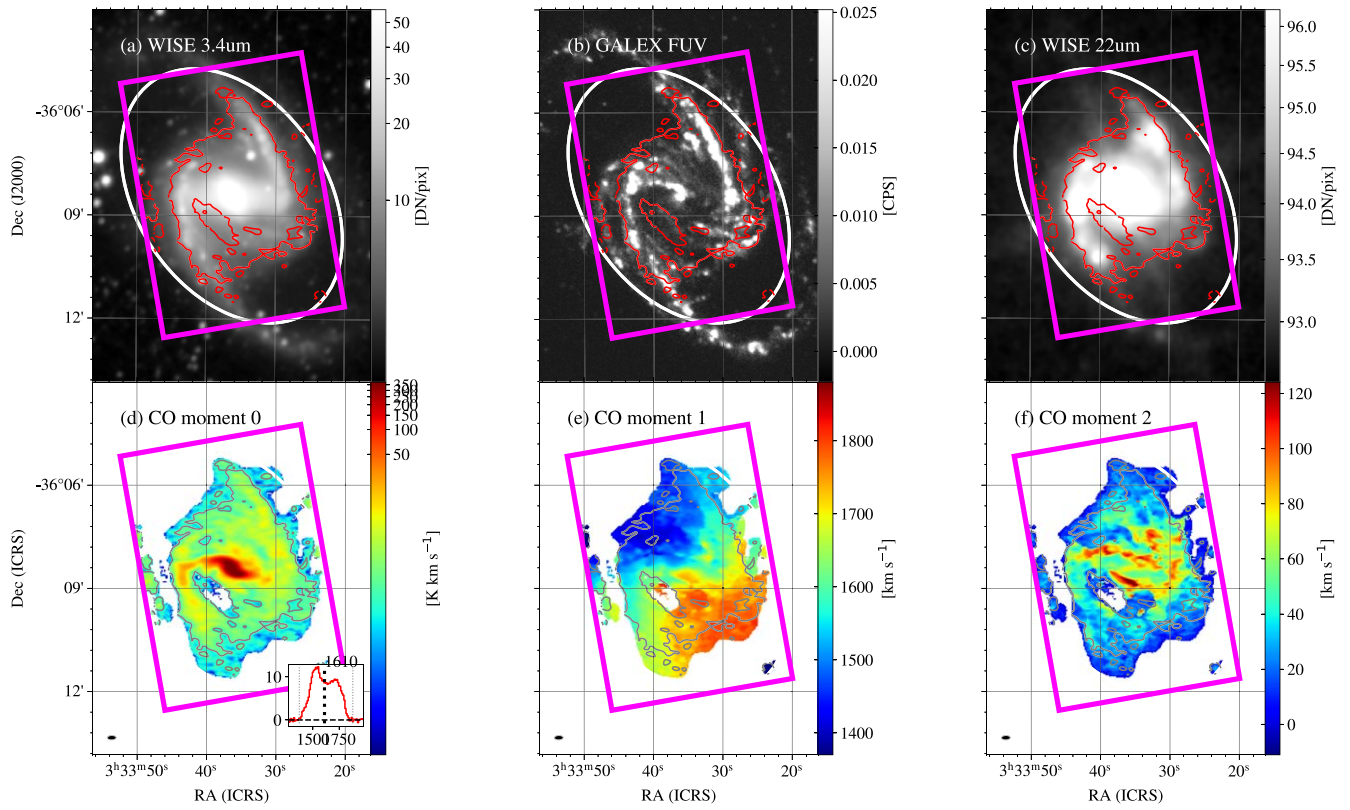


Figure 29. Same as Figure 14, but for NGC 1365.

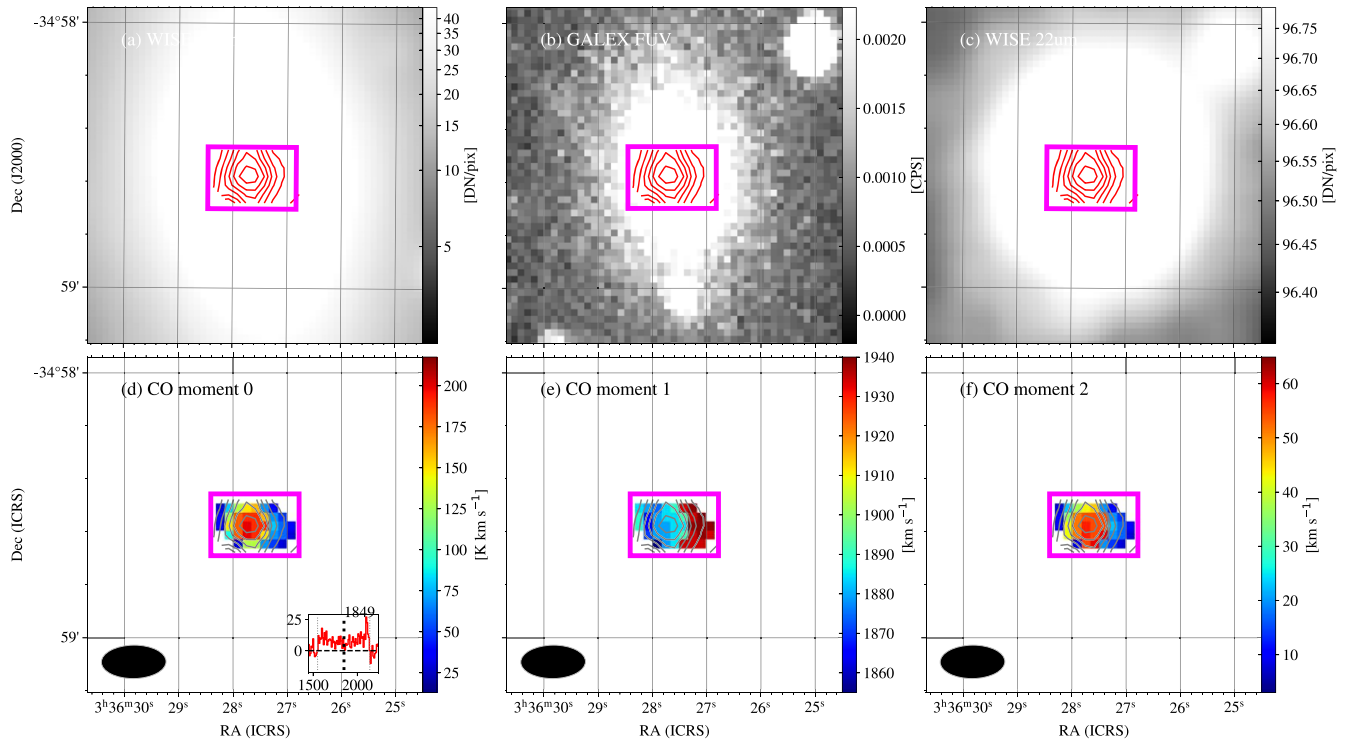


Figure 30. Same as Figure 14, but for NGC 1380.

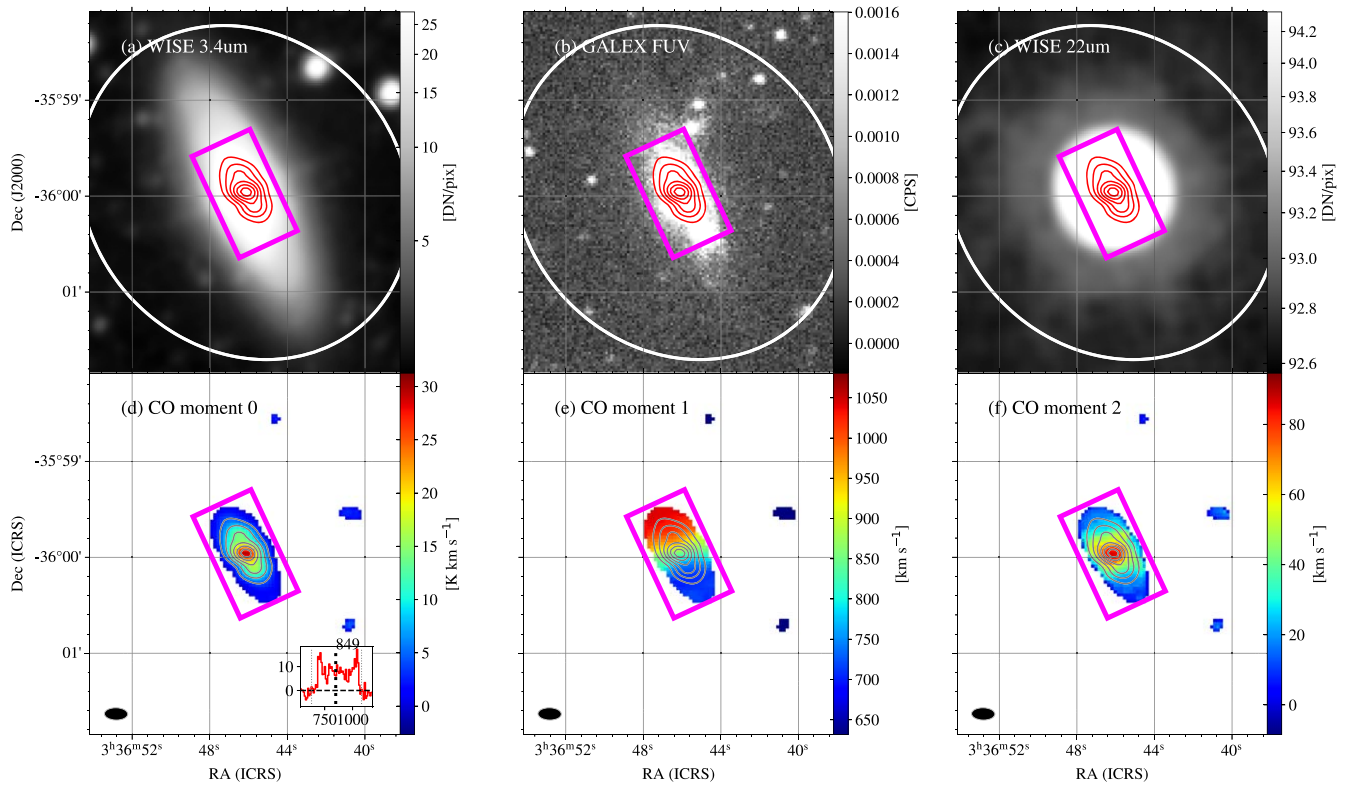


Figure 31. Same as Figure 14, but for NGC 1386.

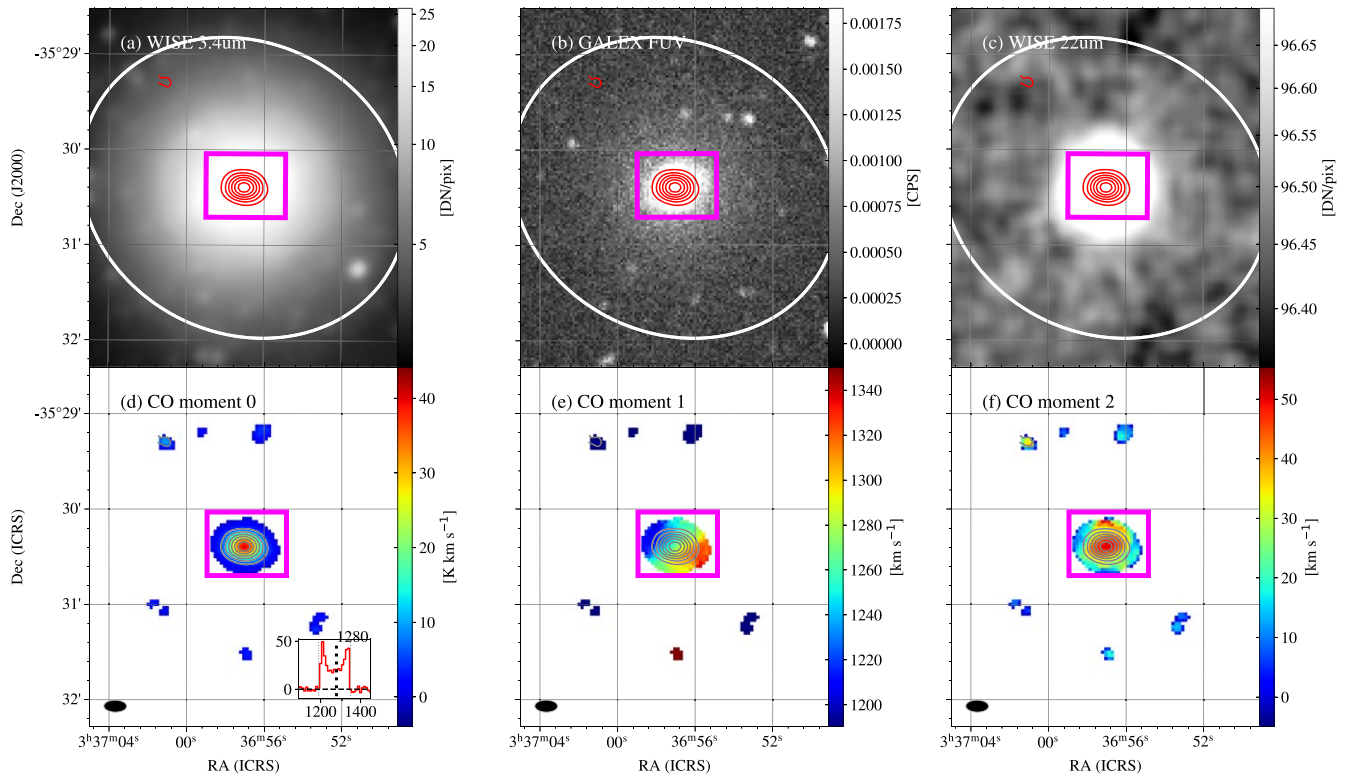


Figure 32. Same as Figure 14, but for NGC 1387.

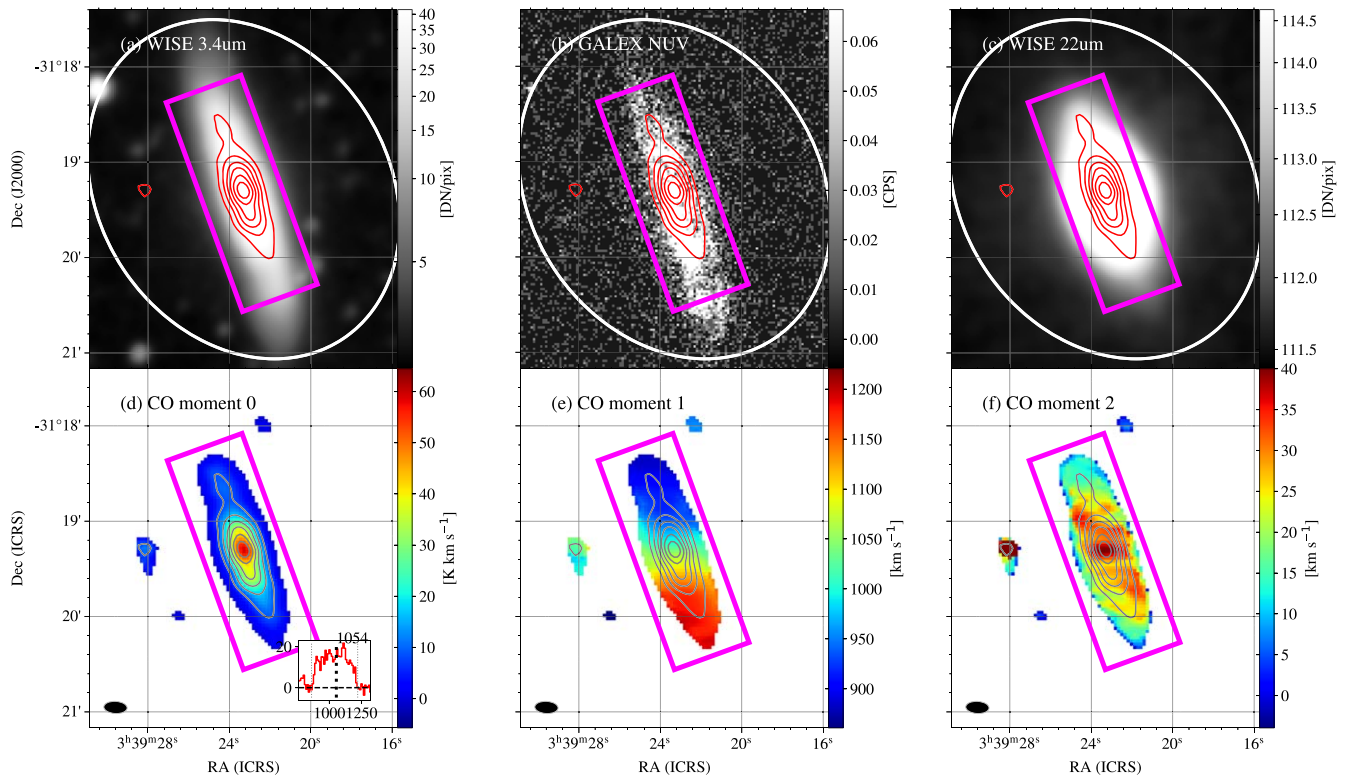


Figure 33. Same as Figure 14, but for NGC 1406.

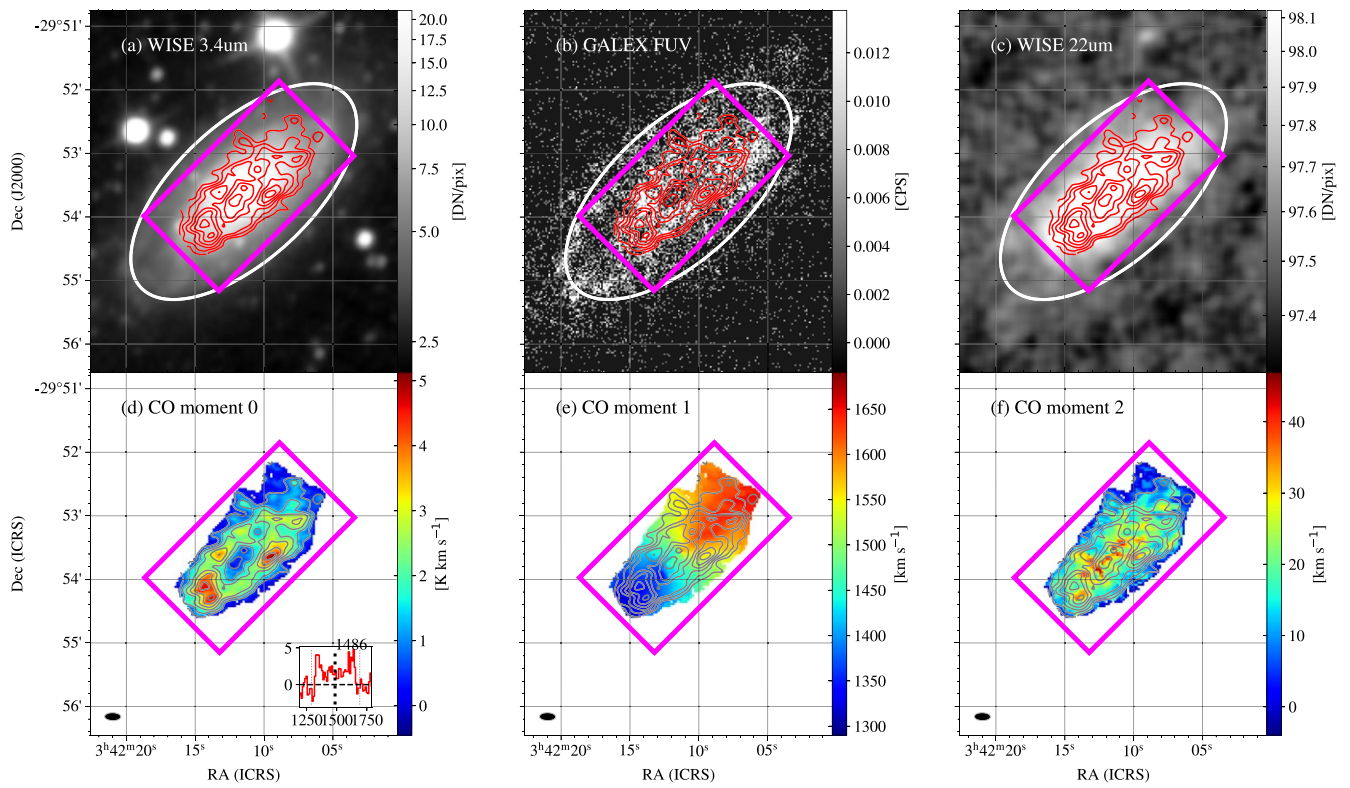


Figure 34. Same as Figure 14, but for NGC 1425.

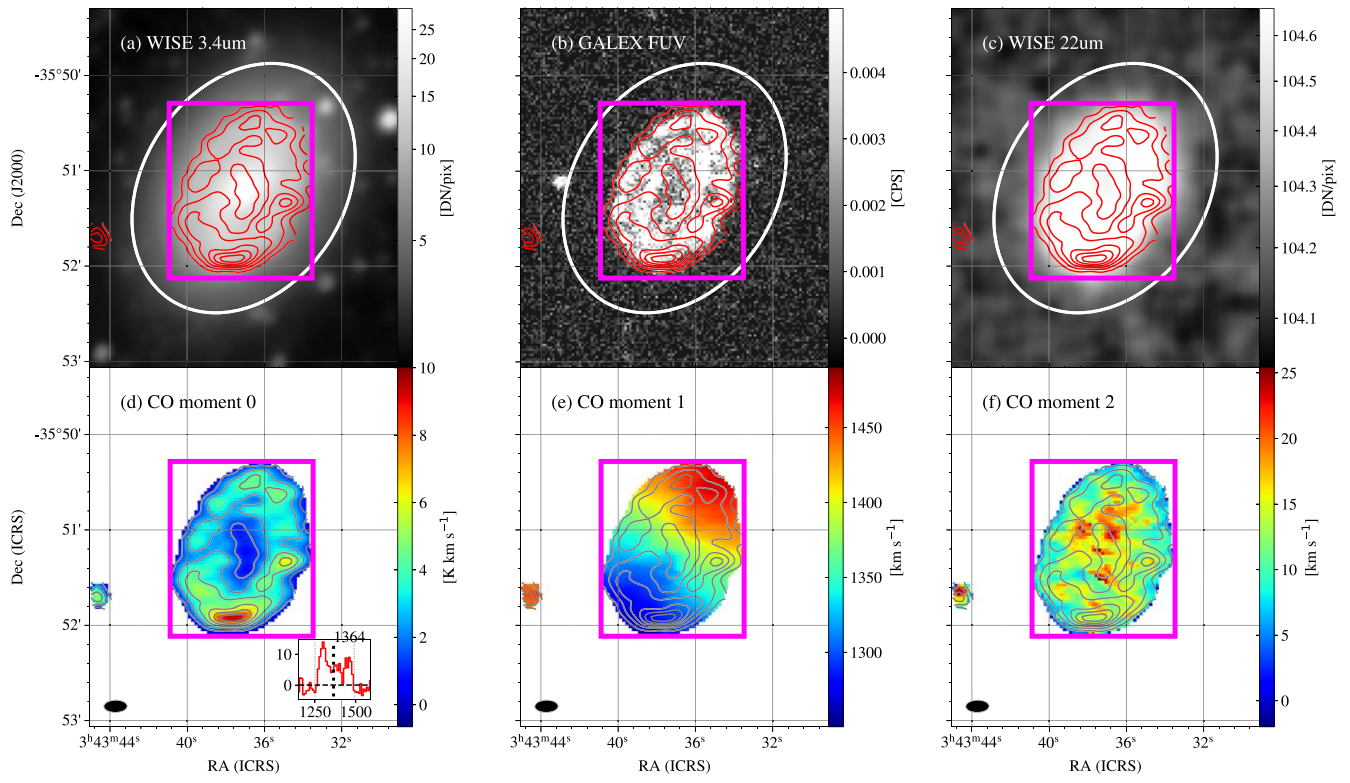


Figure 35. Same as Figure 14, but for NGC 1436.

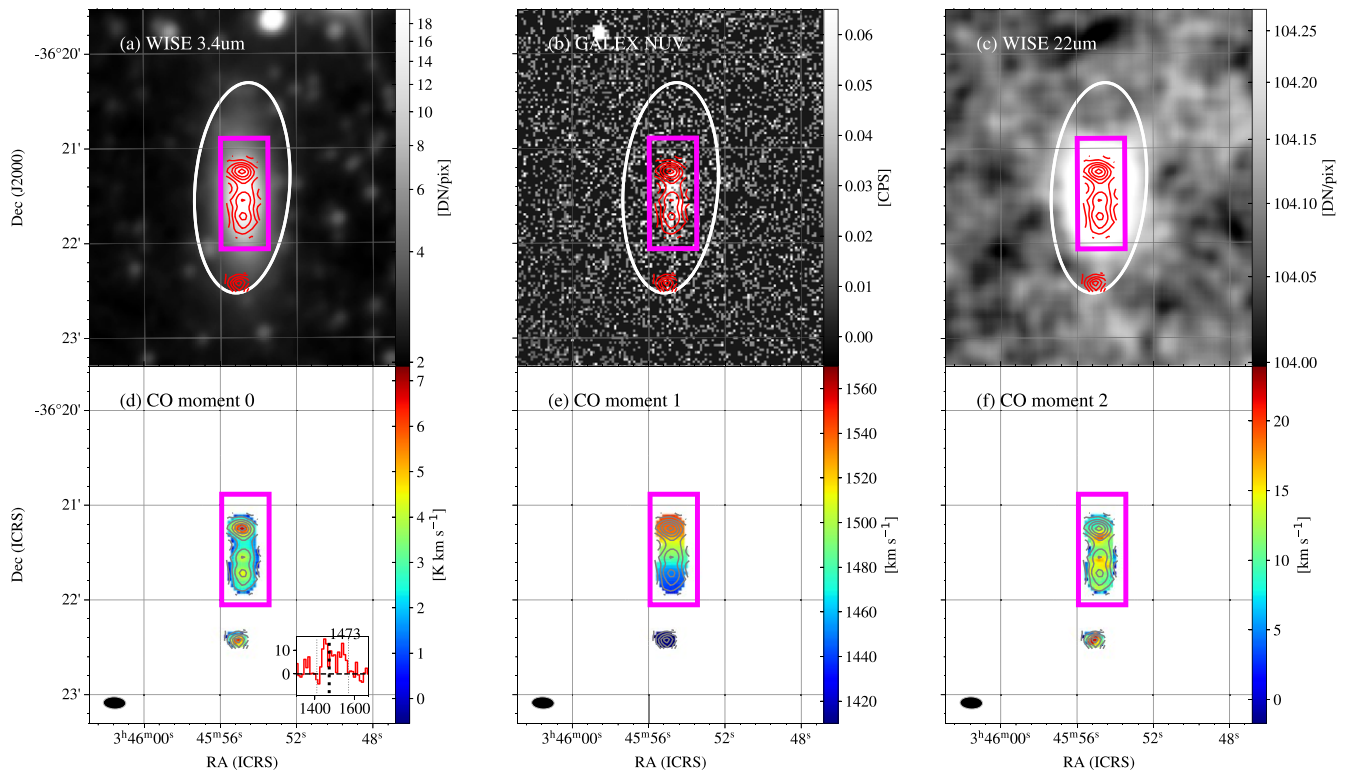


Figure 36. Same as Figure 14, but for NGC 1437B.

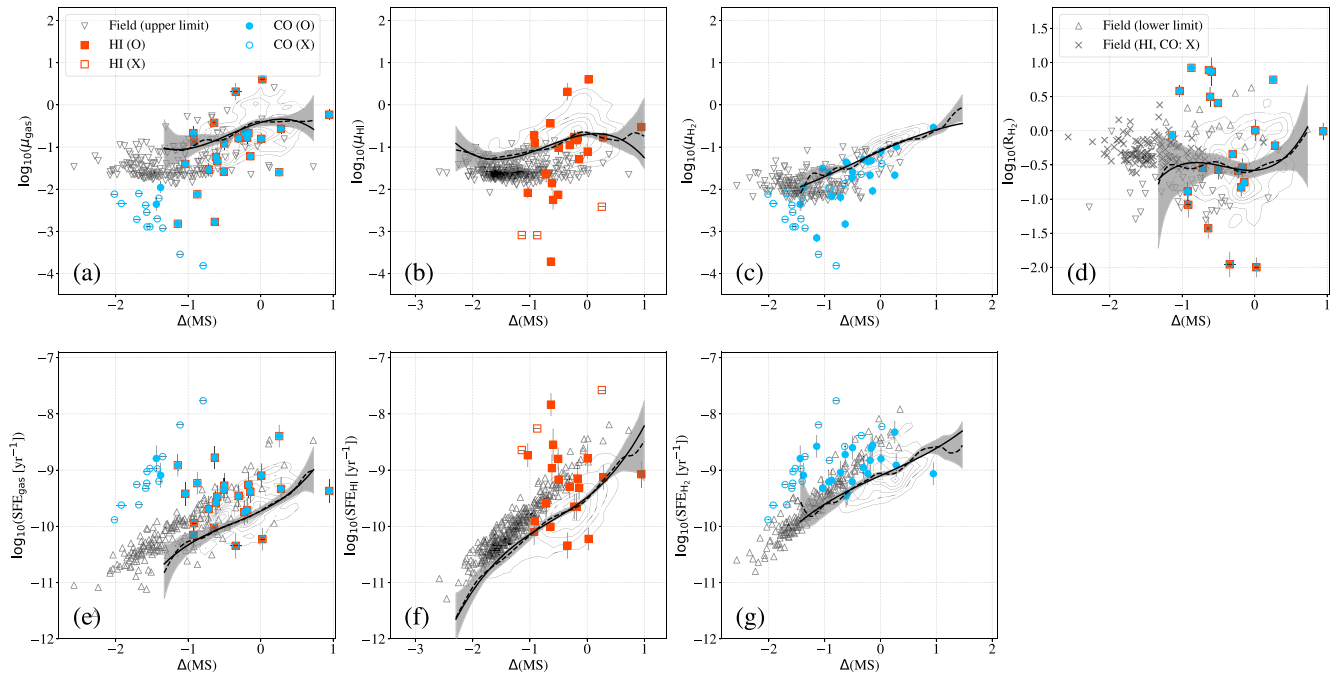


Figure 37. Comparison between the Fornax galaxies and field galaxies with $M_{\text{star}} > 10^9 M_{\odot}$: relationships between the offsets from the main sequence of star-forming galaxies with gas fractions, R_{H_2} , and SFEs. Symbols and lines are the same as in Figure 6. Our Fornax sample has lower gas fractions and higher R_{H_2} and SFEs than the field galaxies at $\Delta(\text{MS}) \sim -1$.

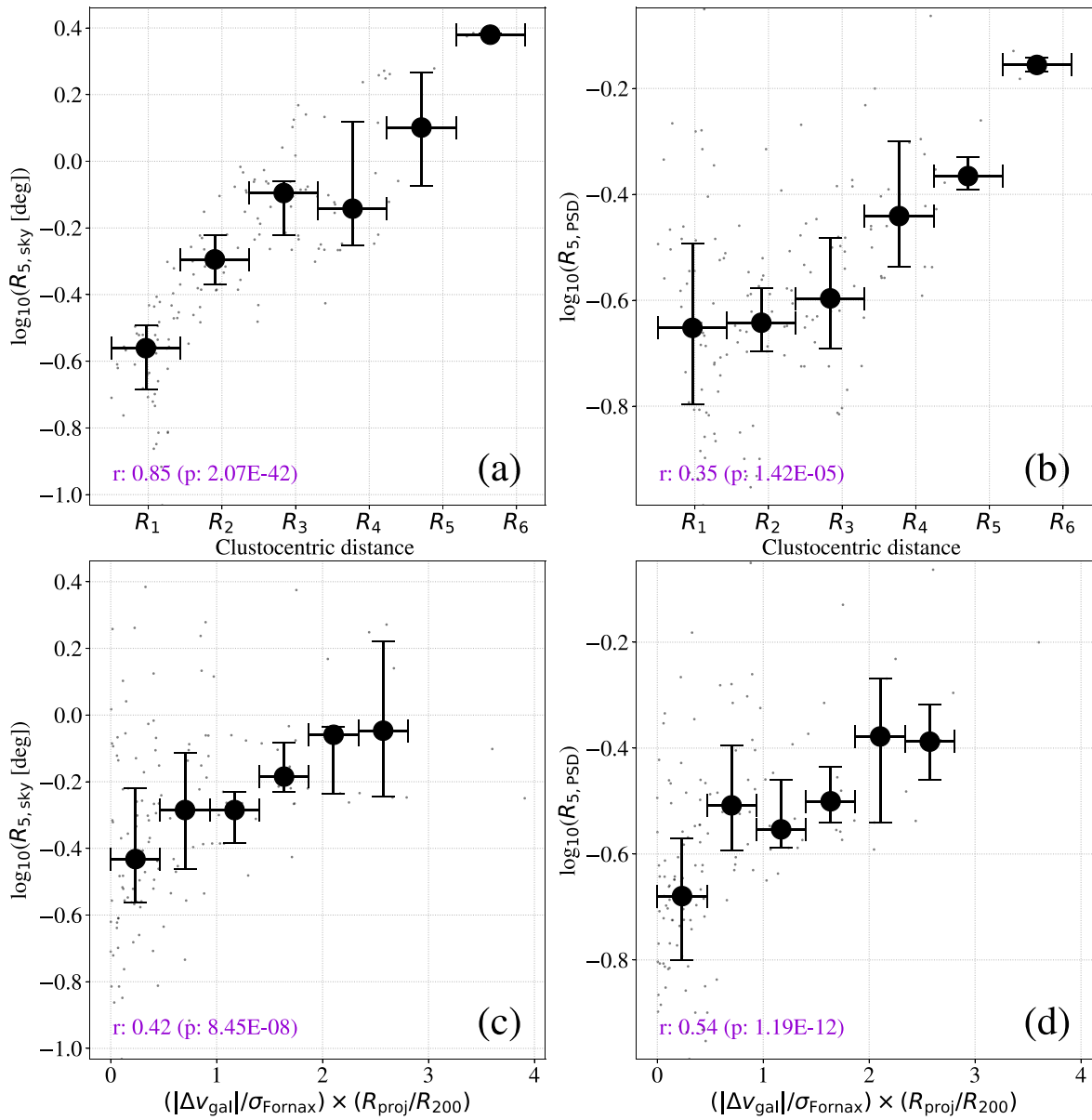


Figure 38. The relationship between the clustocentric distance with (a) $R_{5,\text{sky}}$ and (b) $R_{5,\text{PSD}}$ and between the accretion phase with (c) $R_{5,\text{sky}}$ and (d) $R_{5,\text{PSD}}$. The median values for each bin of the clustocentric distance are indicated with black filled circles. The upper and lower caps indicate the third and first quartiles of $R_{5,\text{sky}}$ or $R_{5,\text{PSD}}$ in each clustocentric distance or accretion phase bin, respectively. The Spearman's rank-order correlation coefficient r and the p -value are shown on the lower left corner of each panel in purple.

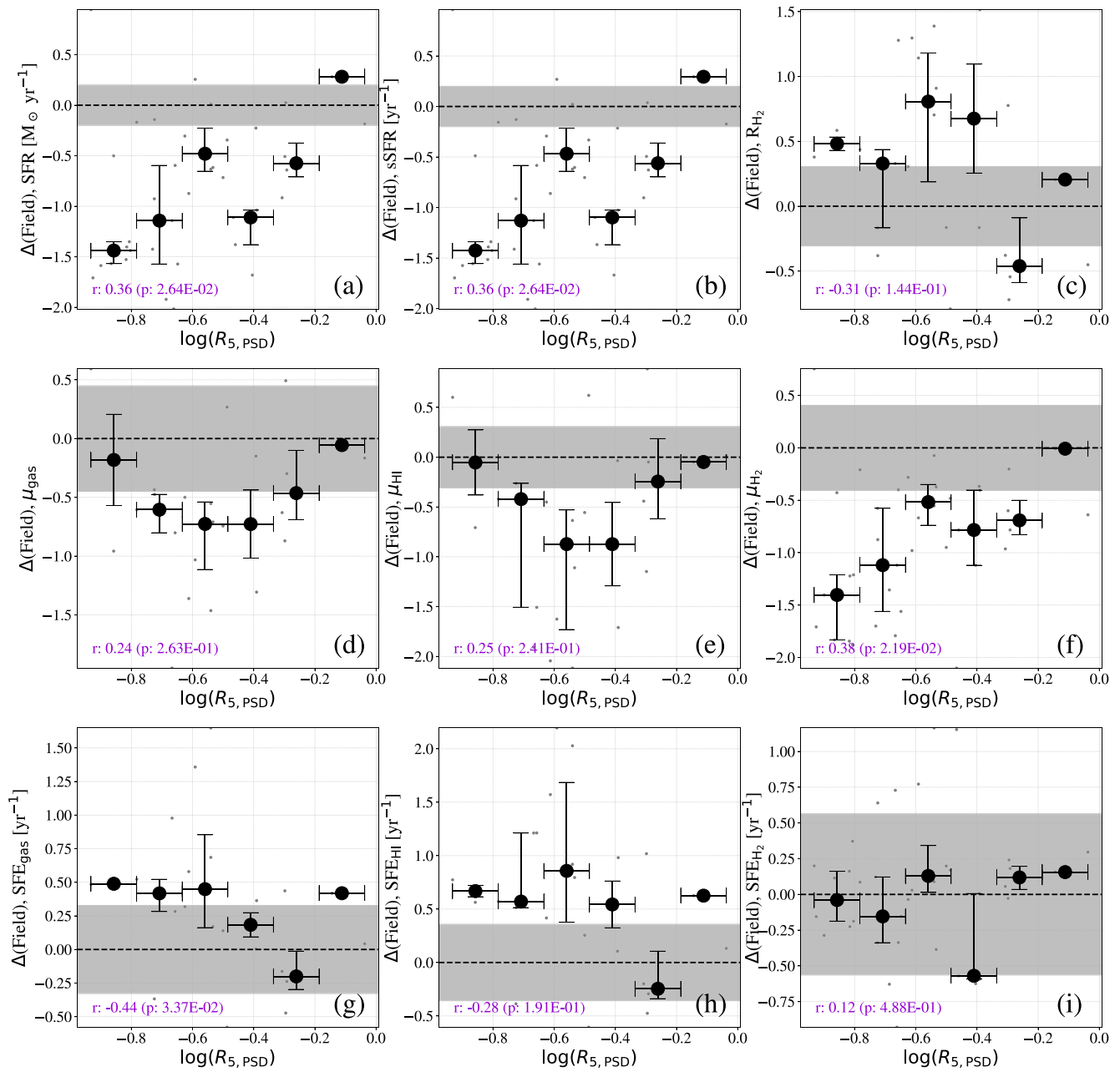


Figure 39. The relationships between $R_{5,PSD}$ with the key quantities [$\Delta(\text{Field})$ values]. Symbols are the same as in Figure 10. $R_{5,PSD}$ positively correlates to μ_{H_2} and negatively correlates to SFE_{gas} .

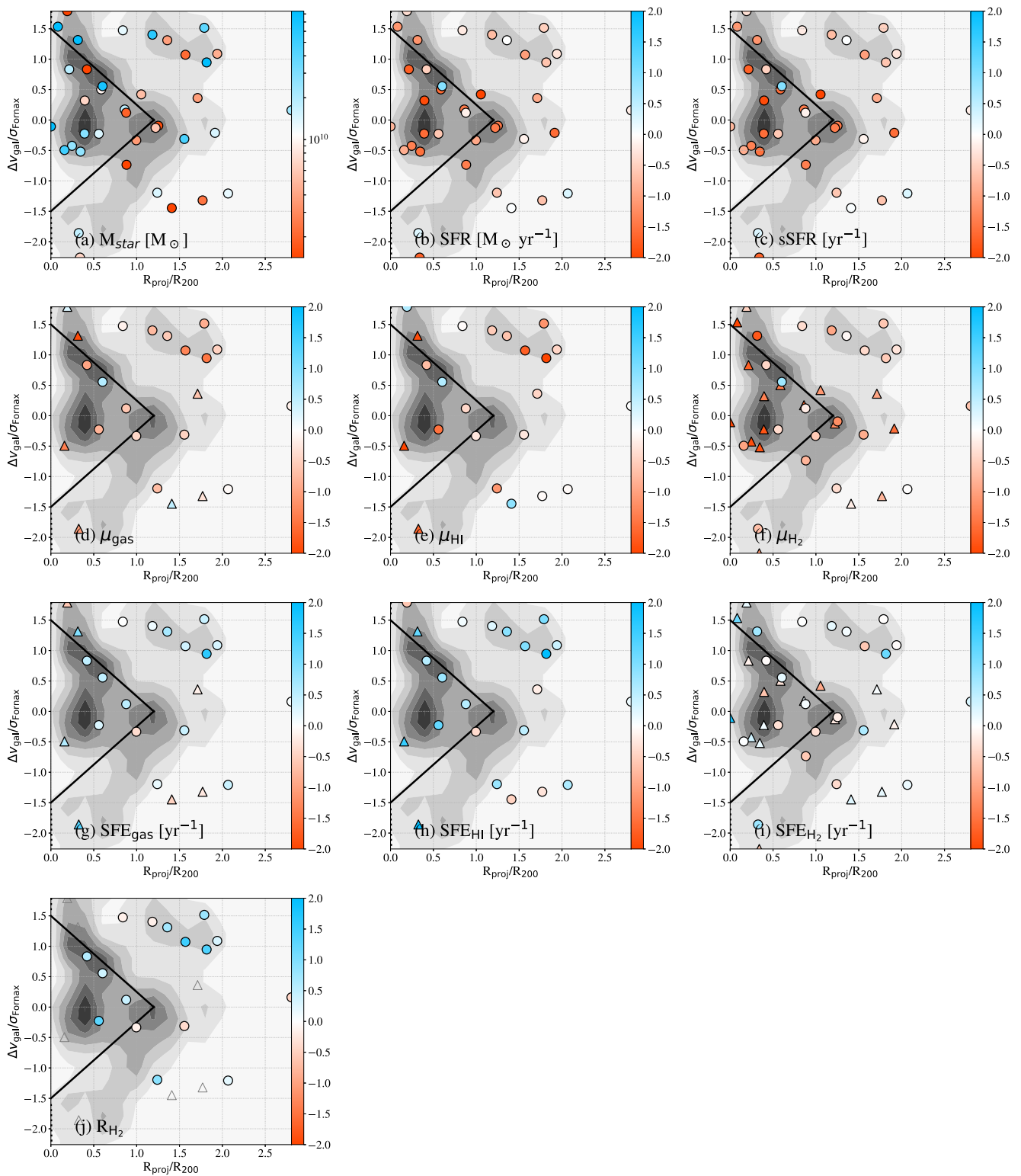


Figure 40. PSD of our sample galaxies in the Fornax cluster, color-coded for the key quantities defined in Section 4.

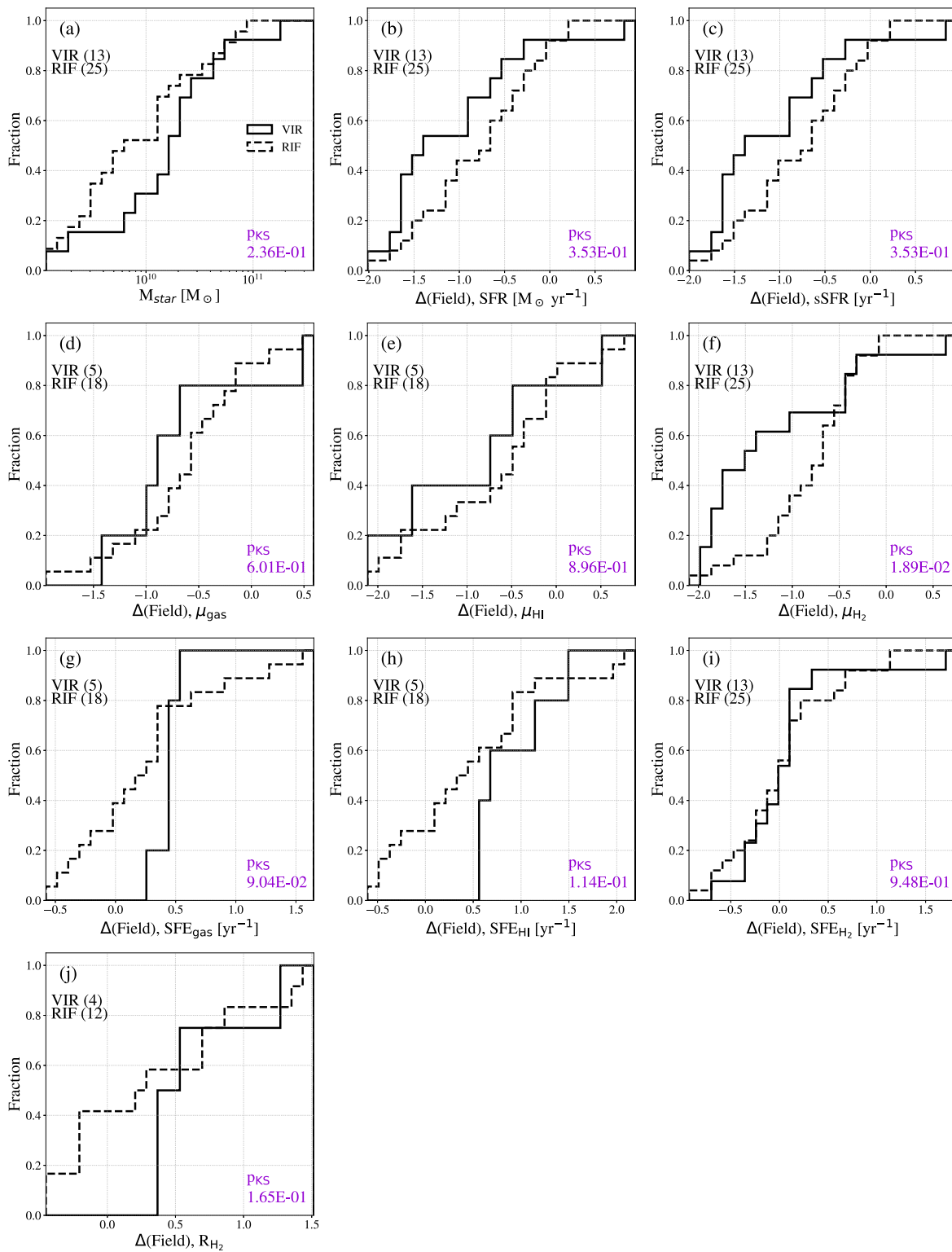


Figure 41. Cumulative histograms of stellar mass and the key quantities of galaxies in the VIR (solid line) and RIF (dashed line). The numbers of galaxies in the VIR and RIF regions are indicated at the upper left corner in each panel. The p -values of the K-S test are indicated in purple at the lower right corner.

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