



Complex Analysis

On the Lelong–Demainly numbers of plurisubharmonic currents

Sur les nombres de Lelong–Demainly des courants plurisousharmoniques

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ARTICLE INFO

Article history:

Received 25 March 2011

Accepted 28 March 2011

Available online 20 April 2011

Presented by Jean-Pierre Demailly

I dedicate this work to the martyrs of the Tunisian revolution, in particular to my colleague Hatem Bettaher.

ABSTRACT

In this Note we study the existence of the Lelong–Demainly number of a negative plurisubharmonic current with respect to a positive plurisubharmonic function on an open subset of \mathbb{C}^n . Then we establish some estimates of the Lelong–Demainly numbers of positive or negative plurisubharmonic currents.

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RÉSUMÉ

Dans cette Note, on étudie l'existence du nombre de Lelong–Demainly d'un courant négatif plurisousharmonique relativement à une fonction positive plurisousharmonique sur un ouvert de \mathbb{C}^n puis on donne quelques estimations des nombres de Lelong–Demainly des courants positifs ou négatifs plurisousharmoniques.

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Version française abrégée

L'existence des nombres de Lelong des courants positifs a été résolu par P. Lelong dans les années 1950 pour le cas des courants fermés, puis ce résultat a été étendu par Skoda au cas des courants positifs plurisousharmoniques. En revanche, il existe des courants négatifs plurisousharmoniques qui n'admettent pas de nombres de Lelong, on peut voir par exemple que $\log(|z_2|^2)[z_1 = 0]$ est un exemple de courant négatif plurisousharmonique de bidimension (1, 1) sur la boule unité de \mathbb{C}^2 qui n'admet pas de nombre de Lelong en 0.

Le principal objectif de cette note est de traiter le problème de l'existence des nombres de Lelong généralisés introduits par Demailly [1,2] d'un courant négatif plurisousharmonique T de bidimension (p, p) sur un ouvert Ω de \mathbb{C}^n ; pour cela on note $\text{PSH}(T, \Omega)$ l'ensemble des fonctions φ positives plurisousharmoniques semi-exhaustives dont le logarithme $\log \varphi$ est plurisousharmonique sur Ω , et telles que le produit extérieur $T \wedge (dd^c \varphi)^p$ soit bien défini. Le nombre de Lelong–Demainly de T relativement à un poids φ tel que $\varphi \in \text{PSH}(T, \Omega)$ est $\nu(T, \varphi) := \lim_{r \rightarrow 0^+} \nu(T, \varphi, r)$, où $\nu(T, \varphi, .)$ est la fonction définie par

$$\nu(T, \varphi, r) := \frac{1}{r^p} \int_{\{\varphi < r\}} T \wedge (dd^c \varphi)^p.$$

Le résultat principal de cette Note est :

Théorème 1. Soient T un courant négatif plurisousharmonique de bidimension (p, p) sur Ω et $\varphi \in \text{PSH}(T, \Omega)$. Si la fonction $t \mapsto \frac{\nu(dd^c T, \varphi, t)}{t}$ est intégrable au voisinage de 0, alors le nombre de Lelong–Demainly $\nu(T, \varphi)$ du courant T relativement à φ existe.

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1. Introduction

The existence of Lelong numbers of positive currents was proved by P. Lelong in the 1950's in the case of *closed* currents, then this result was extended by H. Skoda to the case of *positive plurisubharmonic* currents. However, there are *negative* plurisubharmonic currents which do not admit Lelong numbers, for example $-(-\log(|z_2|^2))^\epsilon [z_1 = 0]$ is a negative plurisubharmonic current of bidimension $(1, 1)$ on the unit ball of \mathbb{C}^2 , which admits no Lelong number at 0 for all $0 < \epsilon \leq 1$.

The main purpose of the first part of this Note is to study the existence of generalized Lelong numbers, introduced by Demainly [1,2], in the case of negative plurisubharmonic currents of bidimension (p, p) on an open set Ω of \mathbb{C}^n . In the second part, we give some proprieties of the Lelong–Demainly numbers of positive or negative plurisubharmonic currents. In particular, we prove that the Lelong–Demainly numbers do not depend to the system of coordinates. To this aim, we consider a non-negative plurisubharmonic function φ on Ω such that $\log \varphi$ is plurisubharmonic on $\{\varphi > 0\}$ and for every $r > 0$, $r_2 > r_1 > 0$, we set

$$\begin{aligned} B_\varphi(r) &:= \{z \in \Omega; \varphi(z) < r\}, \quad B_\varphi(r_1, r_2) := \{z \in \Omega; r_1 \leq \varphi(z) < r_2\} = B_\varphi(r_2) \setminus B_\varphi(r_1), \\ \beta_\varphi &:= dd^c \varphi = \frac{i}{\pi} \partial \bar{\partial} \varphi, \quad \alpha_\varphi := dd^c \log \varphi \text{ on } \{\varphi > 0\}. \end{aligned}$$

We assume that φ is semi-exhaustive, i.e. there exists $R = R(\varphi) > 0$ such that $B_\varphi(R)$ is relatively compact in Ω . If S is a positive (or negative) current of bidimension (p, p) on the set Ω then we denote by $\text{PSH}(S, \Omega)$ the set of non-negative semi-exhaustive plurisubharmonic functions φ such that $\log \varphi$ is also plurisubharmonic on Ω and the exterior product $S \wedge (dd^c \varphi)^p$ is well defined. Finally, we denote by $\mathcal{I}_S(\varphi) := \{k > 0; \varphi^k \in \text{PSH}(S, \Omega)\}$ for every $\varphi \in \text{PSH}(S, \Omega)$; in particular, if φ is C^2 then $\mathcal{I}_S(\varphi) \supset [2, +\infty[$ for every current S .

The Lelong–Demainly number of S relatively to the weight φ is $\nu(S, \varphi) := \lim_{r \rightarrow 0^+} \nu(S, \varphi, r)$ where $\nu(S, \varphi, \cdot)$ is the function defined on $]0, R(\varphi)[$ by $\nu(S, \varphi, r) := \frac{1}{r^p} \int_{B_\varphi(r)} S \wedge \beta_\varphi^p$. The classical Lelong number is given by the choice $\varphi(z) = \varphi_0(z) := |z|^2$.

Example 1. Let $S_\epsilon(z_1, z_2) = (|z_2|^{2\epsilon} - 1)[z_1 = 0]$ where $\epsilon > 0$. Then S_ϵ is a negative plurisubharmonic current of bidimension $(1, 1)$ on the unit ball \mathbb{B} of \mathbb{C}^2 , $\varphi_0 \in \text{PSH}(S_\epsilon, \mathbb{B})$ and $\mathcal{I}_{S_\epsilon}(\varphi_0) = \mathcal{I}_{dd^c S_\epsilon}(\varphi_0) =]0, +\infty[$.

Proof. A simple computation proves that $S_\epsilon \wedge dd^c(\varphi_0^k) = \frac{k^2}{\pi} (|z_2|^{2\epsilon} - 1) |z_2|^{2(k-1)} i dz_2 \wedge d\bar{z}_2 \wedge [z_1 = 0]$ and $dd^c S_\epsilon = \frac{\epsilon^2}{\pi} |z_2|^{2(\epsilon-1)} i dz_2 \wedge d\bar{z}_2 \wedge [z_1 = 0]$, hence $S_\epsilon \wedge dd^c(\varphi_0^k)$ is well defined for all $k > 0$. Therefore, $\mathcal{I}_{S_\epsilon}(\varphi_0) = \mathcal{I}_{dd^c S_\epsilon}(\varphi_0) =]0, +\infty[$. One can check easily that $\nu(S_\epsilon, \varphi_0^k, r) = 2k^2 (\frac{r^{\epsilon/k}}{\epsilon+k} - \frac{1}{k})$ and $\nu(dd^c S_\epsilon, \varphi_0, r) = 2\epsilon r^\epsilon$, hence $\nu(S_\epsilon, \varphi_0^k) = -2k$. \square

2. Main result

In the following, we will use a Lelong–Jensen formula proved by Demainly [1]. In [3], Toujani used an analogue of this formula to prove the existence of the directional Lelong–Demainly numbers of positive plurisubharmonic currents.

Lemma 1. (See [1] or [3].) Let S be a positive or negative plurisubharmonic current of bidimension (p, p) on Ω and $\varphi \in \text{PSH}(S, \Omega)$. Then, for all $0 < r_1 < r_2 < R(\varphi)$,

$$\begin{aligned} &\nu(S, \varphi, r_2) - \nu(S, \varphi, r_1) \\ &= \frac{1}{r_2^p} \int_{B_\varphi(r_2)} S \wedge \beta_\varphi^p - \frac{1}{r_1^p} \int_{B_\varphi(r_1)} S \wedge \beta_\varphi^p \\ &= \int_{B_\varphi(r_1, r_2)} S \wedge \alpha_\varphi^p + \int_{r_1}^{r_2} \left(\frac{1}{t^p} - \frac{1}{r_2^p} \right) dt \int_{B_\varphi(t)} dd^c S \wedge \beta_\varphi^{p-1} + \left(\frac{1}{r_1^p} - \frac{1}{r_2^p} \right) \int_0^{r_1} dt \int_{B_\varphi(t)} dd^c S \wedge \beta_\varphi^{p-1}. \end{aligned} \tag{1}$$

According to Lemma 1, if S is positive plurisubharmonic then $\nu(S, \varphi, \cdot)$ is a non-negative increasing function on $]0, R(\varphi)[$, so $\nu(S, \varphi) := \lim_{r \rightarrow 0^+} \nu(S, \varphi, r)$ exists.

It is well known that if S is a positive plurisubharmonic current, then for every $\varphi \in \text{PSH}(S, \Omega)$ the function $t \mapsto \frac{\nu(dd^c S, \varphi, t)}{t}$ is integrable in the neighborhood of 0. Throughout this Note, for every negative plurisubharmonic current T on Ω , we say that a function $\varphi \in \text{PSH}(T, \Omega)$ satisfies condition (C) if the function $t \mapsto \frac{\nu(dd^c T, \varphi, t)}{t}$ is integrable on a neighborhood of 0. In particular, if φ satisfies condition (C), we must have $\nu(dd^c T, \varphi) = 0$.

Now we state the main result concerning the case of negative plurisubharmonic currents.

Theorem 1. Let T be a negative plurisubharmonic current of bidimension (p, p) on Ω and $\varphi \in \text{PSH}(T, \Omega)$ satisfying condition (C). Then, the Lelong–Demainly number $\nu(T, \varphi)$ of the current T relatively to φ exists.

Proof. For every $0 < r < R(\varphi)$, we set

$$f(r) = \frac{1}{r^p} \int_{B_\varphi(r)} T \wedge \beta_\varphi^p + \frac{1}{r^p} \int_0^r dt \int_{B_\varphi(t)} dd^c T \wedge \beta_\varphi^{p-1} - \int_0^r \frac{dt}{t^p} \int_{B_\varphi(t)} dd^c T \wedge \beta_\varphi^{p-1}.$$

Thanks to condition (C), the function f is well defined and non-positive on $]0, R(\varphi)[$. Indeed,

$$f(r) = v(T, \varphi, r) + \int_0^r \left(\frac{t^p}{r^p} - 1 \right) \frac{v(dd^c T, \varphi, t)}{t} dt \leqslant 0$$

because the function $v(dd^c T, \varphi, .)$ is non-negative on $]0, R(\varphi)[$.

For $0 < r_1 < r_2 < R(\varphi)$ we set $A(r_1, r_2) := f(r_2) - f(r_1)$. The current T is negative plurisubharmonic, thus Lemma 1 gives

$$\begin{aligned} A(r_1, r_2) &= v(T, \varphi, r_2) - v(T, \varphi, r_1) + \frac{1}{r_2^p} \int_0^{r_2} t^{p-1} v(dd^c T, \varphi, t) dt \\ &\quad - \frac{1}{r_1^p} \int_0^{r_1} t^{p-1} v(dd^c T, \varphi, t) dt + \int_{r_1}^{r_2} \frac{v(dd^c T, \varphi, t)}{t} dt \\ &= \int_{B_\varphi(r_1, r_2)} T \wedge \alpha_\varphi^p \leqslant 0. \end{aligned} \tag{2}$$

Therefore, f is a non-positive decreasing function on $]0, R(\varphi)[$, and this implies the existence of the limit $\varrho := \lim_{r \rightarrow 0^+} f(r)$. The hypothesis of integrability of $v(dd^c T, \varphi, t)/t$ and the fact that $(t^p/r^p - 1)$ is uniformly bounded give

$$\lim_{r \rightarrow 0^+} \int_0^r \left(\frac{t^p}{r^p} - 1 \right) \frac{v(dd^c T, \varphi, t)}{t} dt = 0.$$

Therefore, $\varrho = \lim_{r \rightarrow 0^+} f(r) = \lim_{r \rightarrow 0^+} v(T, \varphi, r) = v(T, \varphi)$. \square

Remark 1. If T is a positive (resp. negative) plurisubharmonic current of bidimension (p, p) on Ω and $\varphi \in \text{PSH}(T, \Omega)$ (resp. satisfying condition (C)) then the extension $\widetilde{T \wedge \alpha_\varphi^p}$ of $T \wedge \alpha_\varphi^p$ by 0 over the compact set $\{\varphi = 0\}$ exists and we have

$$\int_{B_\varphi(r)} \widetilde{T \wedge \beta_\varphi^p} = f(r) - v(T, \varphi).$$

Indeed, it suffices to use Lemma 1 to prove that $\int_{B_\varphi(\epsilon, r)} T \wedge \beta_\varphi^p$ is uniformly bounded and then tend ϵ to 0 to prove the equality of the remark.

A natural question arises: does the existence of the Lelong–Demainly number of the negative plurisubharmonic current T implies the existence of $\widetilde{T \wedge \alpha_\varphi^p}$? So condition (C) will be necessary in Theorem 1.

In the following proposition, we give some properties of Lelong–Demainly numbers in both cases of negative or positive plurisubharmonic currents:

Proposition 1. Let T be a positive or negative plurisubharmonic current of bidimension (p, p) on Ω and $\varphi \in \text{PSH}(T, \Omega)$. Then

$$v(T, \varphi^k, r^k) = k^p \left[v(T, \varphi, r) + \int_0^r \frac{v(dd^c T, \varphi, t)}{t} \left(\frac{t^p}{r^p} - \frac{t^{kp}}{r^{kp}} \right) dt \right], \quad \forall k \in \mathcal{I}_T(\varphi), \forall r \in]0, R(\varphi)[. \tag{3}$$

In particular, if T is negative, then $v(T, \varphi)$ exists if and only if for all $k \in \mathcal{I}_T(\varphi)$, $v(T, \varphi^k)$ exists.

Furthermore, in both cases (with the assumption that $v(T, \varphi)$ exists if T is negative) we have

$$v(T, \varphi^k) = k^p v(T, \varphi).$$

The previous equality is due to Demailly in the case of closed positive currents.

Proof. Let $\epsilon > 0$ be sufficiently small. By replacing φ with $\varphi_\epsilon = \varphi + \epsilon$ and $\psi := \varphi^k$ with $\psi_\epsilon = \varphi_\epsilon^k$, when T is negative (respectively positive), equality (2) (resp. equality (1)) gives for $0 < r_1 < \epsilon < r < R(\varphi)$

$$\nu(T, \varphi_\epsilon, r) = \int_{B_{\varphi_\epsilon}(\epsilon, r)} T \wedge \alpha_{\varphi_\epsilon}^p - \frac{1}{r^p} \int_{\epsilon}^r t^{p-1} \nu(dd^c T, \varphi_\epsilon, t) dt + \int_{\epsilon}^r \frac{\nu(dd^c T, \varphi_\epsilon, t)}{t} dt \quad (4)$$

and

$$\begin{aligned} \nu(T, \psi_\epsilon, r^k) &= \int_{B_{\psi_\epsilon}(\epsilon^k, r^k)} T \wedge \alpha_{\psi_\epsilon}^p - \frac{1}{r^{kp}} \int_{\epsilon^k}^{r^k} t^{p-1} \nu(dd^c T, \psi_\epsilon, t) dt + \int_{\epsilon^k}^{r^k} \frac{\nu(dd^c T, \psi_\epsilon, t)}{t} dt \\ &= \int_{B_{\psi_\epsilon}(\epsilon^k, r^k)} T \wedge \alpha_{\psi_\epsilon}^p - \frac{k}{r^{kp}} \int_{\epsilon}^r s^{kp-1} \nu(dd^c T, \psi_\epsilon, s^k) ds + k \int_{\epsilon}^r \frac{\nu(dd^c T, \psi_\epsilon, s^k)}{s} ds \\ &= k^p \int_{B_{\varphi_\epsilon}(\epsilon, r)} T \wedge \alpha_{\varphi_\epsilon}^p - \frac{k^p}{r^{kp}} \int_{\epsilon}^r s^{kp-1} \nu(dd^c T, \varphi_\epsilon, s) ds + k^p \int_{\epsilon}^r \frac{\nu(dd^c T, \varphi_\epsilon, s)}{s} ds \\ &= k^p \left(\nu(T, \varphi_\epsilon, r) + \frac{1}{r^p} \int_{\epsilon}^r t^{p-1} \nu(dd^c T, \varphi_\epsilon, t) dt - \int_{\epsilon}^r \frac{\nu(dd^c T, \varphi_\epsilon, t)}{t} dt \right) \\ &\quad - \frac{k^p}{r^{kp}} \int_{\epsilon}^r s^{kp-1} \nu(dd^c T, \varphi_\epsilon, s) ds + k^p \int_{\epsilon}^r \frac{\nu(dd^c T, \varphi_\epsilon, s)}{s} ds \\ &= k^p \left(\nu(T, \varphi_\epsilon, r) + \int_{\epsilon}^r \frac{\nu(dd^c T, \varphi_\epsilon, t)}{t} \left(\frac{t^p}{r^p} - \frac{t^{kp}}{r^{kp}} \right) dt \right). \end{aligned} \quad (5)$$

Here we have used successively equality (4) with ψ_ϵ instead of φ_ϵ in the first equality, then the change of variable $t = s^k$, next the fact that $\nu(dd^c T, \psi_\epsilon, s^k) = k^{p-1} \nu(dd^c T, \varphi_\epsilon, s)$ (equality proved by Demailly in the case of closed positive currents), and finally equality (4).

When $\epsilon \rightarrow 0$, equality (5) implies equality (3).

Thanks to equality (3), for every $r \in]0, R(\varphi)[$, we have

$$\nu(T, \varphi, r) - \frac{1}{k^p} \nu(T, \varphi^k, r^k) = - \int_0^r \nu(dd^c T, \varphi, t) \left(\frac{t^{p-1}}{r^p} - \frac{t^{kp-1}}{r^{kp}} \right) dt.$$

The current $dd^c T$ is positive and closed, hence $\nu(dd^c T, \varphi, \cdot)$ is a non-negative increasing function and $\nu(dd^c T, \varphi) = 0$. We consider two disjoint cases:

- First case $k \geq 1$.

$$0 \leqslant \int_0^r \nu(dd^c T, \varphi, t) \left(\frac{t^{p-1}}{r^p} - \frac{t^{kp-1}}{r^{kp}} \right) dt \leqslant \nu(dd^c T, \varphi, r) \int_0^r \left(\frac{t^{p-1}}{r^p} - \frac{t^{kp-1}}{r^{kp}} \right) dt = \nu(dd^c T, \varphi, r) \frac{k-1}{kp}.$$

Then

$$-\frac{k-1}{kp} \nu(dd^c T, \varphi, r) \leqslant \nu(T, \varphi, r) - \frac{1}{k^p} \nu(T, \varphi^k, r^k) \leqslant 0.$$

- Second case $k < 1$. A similar calculation shows that we have

$$0 \leqslant \nu(T, \varphi, r) - \frac{1}{k^p} \nu(T, \varphi^k, r^k) \leqslant -\frac{k-1}{kp} \nu(dd^c T, \varphi, r).$$

In the two cases, both terms (of the right hand and the left hand) have the same limit 0 when $r \rightarrow 0^+$. This completes the proof of the proposition. \square

3. Applications

The following corollaries are immediate consequences of Theorem 1 and/or Proposition 1:

Corollary 1. Let T be a negative plurisubharmonic current of bidimension (p, p) on Ω and $\varphi \in \text{PSH}(T, \Omega)$. If there exists $k \in \mathcal{I}_T(\varphi)$ such that the function $t \mapsto \frac{\nu(dd^c T, \varphi, t^{1/k})}{t}$ is integrable on a neighborhood of 0, then the Lelong–Demainly number $\nu(T, \varphi)$ of T relatively to φ exists.

Proof. We can prove this corollary in two different ways:

- First way. Thanks to Proposition 1, to prove that $\nu(T, \varphi)$ exists, it suffices to show that $\nu(T, \varphi^k)$ exists; for this we observe that

$$\frac{\nu(dd^c T, \varphi^k, t)}{t} = \frac{\nu(dd^c T, \varphi^k, (t^{1/k})^k)}{t} = k^{p-1} \frac{\nu(dd^c T, \varphi, t^{1/k})}{t}$$

is integrable on a neighborhood of 0 (hypothesis). Hence, thanks to Theorem 1, $\nu(T, \varphi^k)$ exists.

- Second way. We remark that the integrability of $\frac{\nu(dd^c T, \varphi, t^{1/k})}{t}$ is equivalent to the integrability of $\frac{\nu(dd^c T, \varphi, t)}{t}$. Hence, thanks to Theorem 1, $\nu(T, \varphi)$ exists. In fact if we take $t = s^k$ we obtain

$$\int_0^{r_0} \frac{\nu(dd^c T, \varphi, t^{1/k})}{t} dt = k \int_0^{r_0^k} \frac{\nu(dd^c T, \varphi, s)}{s} ds. \quad \square$$

Corollary 2. Let T be a negative plurisubharmonic current of bidimension (p, p) on Ω and $\varphi \in \text{PSH}(T, \Omega)$ such that $\mathcal{I}_T(\varphi)$ is non-bounded (this is always the case if φ is C^2). Then for every $r \in]0, R(\varphi)[$ one has $\nu(T, \varphi, r) \leq - \int_0^r \nu(dd^c T, \varphi, t) \frac{t^{p-1}}{r^p} dt$. In particular, $\nu(T, \varphi, r) \leq -\frac{1-s^{-p}}{p} \nu(dd^c T, s\varphi, r) \forall s \geq 1$.

Proof. Thanks to Proposition 1, for every $k \in \mathcal{I}_T(\varphi)$, we have

$$u(k) := \frac{1}{k^p} \nu(T, \varphi^k, r^k) = \nu(T, \varphi, r) + \int_0^r \frac{\nu(dd^c T, \varphi, t)}{t} \left(\frac{t^p}{r^p} - \frac{t^{kp}}{r^{kp}} \right) dt.$$

The function u is non-positive and increasing on $\mathcal{I}_T(\varphi)$, so

$$\lim_{k \rightarrow +\infty, k \in \mathcal{I}_T(\varphi)} u(k) = \nu(T, \varphi, r) + \int_0^r \nu(dd^c T, \varphi, t) \frac{t^{p-1}}{r^p} dt \leq 0.$$

If $s = 1$ the inequality is clear. For $s > 1$ we have

$$\nu(T, \varphi, r) \leq - \int_0^r \nu(dd^c T, \varphi, t) \frac{t^{p-1}}{r^p} dt \leq - \int_{r/s}^r \nu(dd^c T, \varphi, t) \frac{t^{p-1}}{r^p} dt \leq -\frac{1-s^{-p}}{p} \nu(dd^c T, \varphi, \frac{r}{s}). \quad \square$$

Theorem 2. Let T be a positive or negative plurisubharmonic current of bidimension (p, p) on Ω and $\varphi, \psi \in \text{PSH}(T, \Omega)$ such that

$$\liminf_{\varphi(z) \rightarrow 0} \frac{\log \psi(z)}{\log \varphi(z)} \geq \ell$$

where $\mathcal{I}_T(\varphi)$ contains a neighborhood $\mathcal{V}(\ell)$ of ℓ . One has

- $\nu(T, \psi) \geq \ell^p \nu(T, \varphi)$ if T is positive.
- $\nu(T, \psi) \leq \ell^p \nu(T, \varphi)$ if T is negative (with the assumption that ψ satisfies condition (C)).

In particular, if $\log \psi(z) \sim \ell \log \varphi(z)$ for $\varphi(z) \in \mathcal{V}(0)$ then $\nu(T, \psi) = \ell^p \nu(T, \varphi)$.

This type of theorem is called a “comparison theorem”; such a statement has been proved by Demainly in the case of positive closed currents. Furthermore, one can deduce the invariance of the Lelong–Demainly numbers under changes of coordinate systems.

Proof. One can replace ψ by ψ^k with $1 < k \in \mathcal{I}_T(\psi)$, in order to assume that

$$\liminf_{\varphi(z) \rightarrow 0} \frac{\log \psi(z)}{\log \varphi(z)} > \ell.$$

As a consequence

$$\lim_{\varphi(z) \rightarrow 0} \frac{\psi(z)}{\varphi(z)^\ell} = 0.$$

For $\epsilon > 0$ small sufficiently, we set $\Psi_\epsilon := \psi + \epsilon \varphi^\ell \sim_{\varphi(z) \in \mathcal{V}(0)} \epsilon \varphi^\ell$. Thanks to the dominated convergence theorem,

$$\lim_{\epsilon \rightarrow 0} \frac{1}{r^p} \int_{\Psi_\epsilon < r} T \wedge \beta_{\Psi_\epsilon}^p = \frac{1}{r^p} \int_{\psi < r} T \wedge \beta_\psi^p.$$

- First assume $T \geq 0$. As the function $v(T, \Psi_\epsilon, \cdot)$ is increasing then

$$\begin{aligned} v(T, \Psi_\epsilon, r) &= \frac{1}{r^p} \int_{\Psi_\epsilon < r} T \wedge \beta_{\Psi_\epsilon}^p \geq \lim_{\rho \rightarrow 0} \frac{1}{\rho^p} \int_{\Psi_\epsilon < \rho} T \wedge \beta_{\Psi_\epsilon}^p = \lim_{\rho \rightarrow 0} \frac{1}{\rho^p} \int_{\epsilon \varphi^\ell < \rho} T \wedge \beta_{\Psi_\epsilon}^p \\ &\geq \lim_{\rho \rightarrow 0} \frac{1}{\rho^p} \int_{\epsilon \varphi^\ell < \rho} T \wedge \beta_{\epsilon \varphi^\ell}^p = v(T, \varphi^\ell) \end{aligned}$$

because $\Psi_\epsilon \sim \epsilon \varphi^\ell$ and $T \wedge (dd^c(\psi + \epsilon \varphi^\ell))^p \geq T \wedge (\epsilon dd^c(\varphi^\ell))^p$. If $\epsilon \rightarrow 0$, we obtain $v(T, \psi, r) \geq v(T, \varphi^\ell)$. Hence, if $r \rightarrow 0$, $v(T, \psi) \geq v(T, \varphi^\ell)$.

- Now, assume $T \leq 0$. Thanks to Theorem 1, equality (2) gives

$$v(T, \Psi_\epsilon, r) + \int_0^r \frac{v(dd^c T, \Psi_\epsilon, t)}{t} \left(\frac{t^p}{r^p} - 1 \right) dt \leq v(T, \Psi_\epsilon) = \lim_{\rho \rightarrow 0} \frac{1}{\rho^p} \int_{\Psi_\epsilon < \rho} T \wedge \beta_{\Psi_\epsilon}^p$$

and a similar calculation proves that

$$v(T, \Psi_\epsilon, r) + \int_0^r \frac{v(dd^c T, \Psi_\epsilon, t)}{t} \left(\frac{t^p}{r^p} - 1 \right) dt \leq v(T, \varphi^\ell).$$

So if $\epsilon \rightarrow 0$ we obtain

$$v(T, \psi, r) + \int_0^r \frac{v(dd^c T, \psi, t)}{t} \left(\frac{t^p}{r^p} - 1 \right) dt \leq v(T, \varphi^\ell).$$

When $r \rightarrow 0$, the last inequality gives $v(T, \psi) \leq v(T, \varphi^\ell)$.

In the particular case $\log \psi(z) \sim \ell \log \varphi(z)$, for every $\epsilon > 0$, there exists $\eta > 0$ so that if $|\varphi(z)| < \eta$, then $\varphi(z)^{\ell(1+\epsilon)} \leq \psi(z) \leq \varphi(z)^{\ell(1-\epsilon)}$ and for ϵ small enough, $\ell(1+\epsilon), \ell(1-\epsilon) \in \mathcal{V}(\ell) \subset \mathcal{I}_T(\varphi)$. Therefore we can apply the previous inequalities to obtain $\ell^p(1+\epsilon)^p v(T, \varphi) \leq v(T, \psi) \leq \ell^p(1-\epsilon)^p v(T, \varphi)$ (resp. $\ell^p(1-\epsilon)^p v(T, \varphi) \leq v(T, \psi) \leq \ell^p(1+\epsilon)^p v(T, \varphi)$) in the case $T \geq 0$ (resp. $T \leq 0$). Hence, when $\epsilon \rightarrow 0$, we obtain $v(T, \psi) = \ell^p v(T, \varphi)$. \square

Acknowledgements

I thank Professor Jean-Pierre Demailly for several suggestions which enabled me to clarify and improve the original form of this Note. I thank also Professor Khalifa Dabbek for useful discussions concerning this work.

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