



Properties of Iron Primary Cosmic Rays: Results from the Alpha Magnetic Spectrometer

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We report the observation of new properties of primary iron (Fe) cosmic rays in the rigidity range 2.65 GV to 3.0 TV with 0.62 million iron nuclei collected by the Alpha Magnetic Spectrometer experiment on the International Space Station. Above 80.5 GV the rigidity dependence of the cosmic ray Fe flux is identical to the rigidity dependence of the primary cosmic ray He, C, and O fluxes, with the Fe/O flux ratio being constant at 0.155±0.006. This shows that unexpectedly Fe and He, C, and O belong to the same class of primary cosmic rays which is different from the primary cosmic rays Ne, Mg, and Si class.

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1. Introduction

Primary iron cosmic rays are the most abundant heavy nuclei beyond silicon. They are thought to be mostly produced and accelerated in astrophysical sources. Precise knowledge of the iron spectrum in the GV–TV rigidity region provides important information on the origin, acceleration, and propagation processes of cosmic rays in the Galaxy [1].

We present the precise measurement of the Fe flux in the rigidity range from 2.65 GV to 3.0 TV based on 0.62×10^6 iron nuclei collected by the Alpha Magnetic Spectrometer (AMS) during its first 8.5 years (from May 19, 2011 to October 30, 2019) of operation on the International Space Station (ISS).

2. AMS-02 Detector and analysis

The AMS detector is a large-acceptance magnetic spectrometer, the full description of the detector is presented in [2] and references therein. The key elements used in this measurement are the permanent magnet, the nine layers of silicon tracker, L1-L9, and the four planes of time of flight (TOF) scintillation counters. Together with the permanent magnet, the silicon tracker measures the rigidity of charged cosmic rays.

In the first 8.5 years of operation AMS has collected 1.50×10^{11} cosmic ray events. Iron events are required to be downward going and to have a reconstructed track in the seven tracker layers placed on the top and inside the magnet. The track is also required to pass through tracker L1, and for the highest rigidity region R \ge 1.2 TV, through tracker L9. Charge measurements on tracker L1, the inner tracker (L2-L8), the upper TOF, and, for R \ge 1.2 TV, the lower TOF, and tracker L9 are required to be compatible with charge Z=26. Details of the Iron flux analysis procedure and particularly the studies of the systematic errors can be found in [3].

The material traversed by nuclei from the top of AMS to tracker L9 is composed primarily of carbon and aluminum. The survival probabilities of Fe nuclei due to interactions in the materials were measured using cosmic ray data collected by AMS as described in [4]. The simulation of nuclear interactions has been validated with data using all AMS measured nuclear charge changing cross sections (Fe \rightarrow He...Mn+X). Fig. 1 shows examples for Fe \rightarrow Cr+X, Fe \rightarrow Si+X, and Fe \rightarrow O+X.



Figure 1: Comparison of the simulated and measured $Fe \rightarrow Cr+X$, $Fe \rightarrow Si+X$, and $Fe \rightarrow O+X$ break-up probabilities between tracker L1 and tracker L2.

3. Properties of Iron Cosmic Ray nuclei

Fig. 2 shows the measured Fe flux as a function of rigidity \tilde{R} at the top of the AMS detector with the total errors, the sum in quadrature of statistical and systematic errors. In the figure the points are placed along the abscissa at \tilde{R} calculated for a flux $\propto R^{-2.7}$ [5]. For comparison, Fig. 2 also shows the latest AMS result on the oxygen flux from [2, 6].



Figure 2: The AMS iron flux (red dots) and oxygen flux (violet triangles) multiplied by $\tilde{R}^{2.7}$ with total errors as a function of rigidity.

To examine the rigidity dependence of the Fe flux, the variation of the flux spectral index with rigidity was obtained in a model independent way as $\gamma = d[\log(\Phi)]/d[\log(\Phi)]$ computed over nonoverlapping rigidity intervals bounded by 7.09, 12.0, 16.6, 28.8, 45.1, 80.5, 175.0, and 3000.0 GV. The results are presented in Fig. 3 together with the spectral index of the oxygen flux from [2]. As seen, above 80.5 GV the rigidity dependence of the iron flux and its spectral index follow the rigidity dependence of the oxygen flux and its spectral index.



Figure 3: The AMS iron flux spectral index (red dots) and oxygen flux spectral index (violet triangles) as functions of rigidity.

Fig. 4 shows the AMS iron flux as a function of kinetic energy per nucleon E_K together with measurements from other experiments [7–17].



Figure 4: The AMS iron flux as a function of kinetic energy per nucleon E_K multiplied by $E_K^{2,7}$ together with measurements from other experiments [7–17]. For the AMS measurement $E_K = (\sqrt{Z^2 \tilde{R}^2 + M^2} - M)/A$ where Z, M, and A are the ⁵⁶₂₆Fe nuclear charge, mass, and atomic mass numbers, respectively. Data from CALET are taken from [17] and data from other experiments are taken from [18].

To compare the rigidity dependence of the Fe flux with that of He, C, and O primary cosmic ray fluxes, which have identical rigidity dependence above 60 GV [2, 6], the ratio of the iron flux to the oxygen flux [2], Fe/O, was computed. Fig. 5 shows the AMS Fe/O ratio as a function of rigidity with total errors together with a constant value fit above 80.5 GV. The fit yields Fe/O=0.155±0.006 with χ^2 /d.o.f. = 8/11. This, together with Fig. 2, shows that Fe belongs to the same class of primary cosmic rays as He, C, and O.



Figure 5: The AMS Fe/O ratio as a function of rigidity with total errors and with red band (1σ) indicating the constant value fit (Fe/O=0.155 ± 0.006 with χ^2 /d.o.f. = 8/11) above 80.5 GV.

To compare the AMS result with previous measurements, the Fe/O ratio was converted from rigidity to kinetic energy per nucleon using the procedure described in [19]. Fig. 6 shows the

AMS Fe/O flux ratio as a function of kinetic energy per nucleon together with earlier measurements [7, 8, 11, 13, 15, 16]. As seen, the AMS result provides an accurate measurement of the energy dependence of the Fe/O flux ratio.



Figure 6: The AMS Fe/O flux ratio as a function of kinetic energy per nucleon E_K together with earlier measurements. To convert the AMS measurements from rigidity to kinetic energy per nucleon we used ⁵⁶Fe and ¹⁶O.

In summary, the rigidity dependence of Fe above 80.5 GV follows the rigidity dependence of O, see Fig. 2 and Fig. 5. Therefore as shown in Fig. 7, unexpectedly Fe belongs to the class of light primary cosmic rays He, C and O [2, 6], which is different from the rigidity dependence of the heavy primary cosmic rays Ne, Mg, and Si [20].



Figure 7: The rigidity dependence of the Fe spectrum compared with the rigidity dependence of the He, C, and O spectra and the Ne, Mg, and Si spectra above 80.5 GV. For clarity, the He, O, Ne, and Si data points above 400 GV are displaced horizontally. For display purposes only, the He, C, O, Ne, Si, and Fe spectra were rescaled as indicated. The shaded areas show the fit result of Eq. (5) of [20] to He, C, and O spectra (magenta) and Ne, Mg, and Si spectra (green).

4. Conclusion

We have presented the precision measurement of the cosmic ray Fe spectrum as a function of rigidity from 2.65 GV to 3.0 TV. Above 80.5 GV the rigidity dependence of the cosmic-ray Fe spectrum is identical to the rigidity dependence of the light primary cosmic-ray spectra He, C, and O, which is different from the rigidity dependence of the heavy primary cosmic-ray Ne, Mg, and Si spectra. In particular, above 80.5 GV the Fe/O ratio is well described by a constant value of 0.155±0.006. These are new and unexpected properties of primary iron cosmic rays.

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