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Undersampling Raster Scans in Spectromicroscopy for reduced dose and faster measurements

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Abstract: Combinations of spectroscopic analysis and microscopic techniques are used across many disciplines of scientific research, including material science, chemistry and biology. X-ray spectromicroscopy, in particular, is a powerful tool used for studying chemical state distributions at the micro and nano scales. With the beam fixed, a specimen is typically rastered through the probe with continuous motion and a range of multimodal data is collected at fixed time intervals.

The application of this technique is limited in some areas due to: long scanning times to collect the data, either because of the area/volume under study or the compositional properties of the specimen; and material degradation due to the dose absorbed during the measurement. In this work, we propose a novel approach for reducing the dose and scanning times by undersampling the raster data. This is achieved by skipping rows within scans and reconstructing the x-ray spectromicroscopic measurements using low-rank matrix completion. The new method is robust and allows for 5 to 6-fold reduction in sampling. Experimental results obtained on real data are illustrated.

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1. Introduction

X-ray spectromicroscopy combines x-ray spectroscopy and x-ray microscopy to map changes in chemical state across a specimen on the micro and nano scales. These techniques have been broadly applied to problems across life and physical sciences, such as chemical engineering [1], material science [2], and biology [3]. A spectromicroscopy experiment involves measuring an $n_1 \times n_2$ grid at n_E energy levels across the absorption edge of an element of interest. The resulting $n_1 \times n_2$ spectra represents a big data challenge; to extract meaningful information they are typically analyzed using PCA and cluster analysis to reduce to a mapping of representative spectra, or to a low-rank representation of the data.

While the technique has been generally successful, the application of spectromicroscopy to in-situ studies and in areas of soft matter or biological materials is limited by two main factors. First, the total experiment time required to collect the data to a given statistical significance; second, the total radiation dose over the collection and any resulting damage to the object, or changes to the chemical state, that may occur as a result. The issue of damage due to dose and long collection times occurs across both x-ray and electron optical systems.

To alleviate this issue, a variety of approaches to reduce the number of samples has been proposed. In electron tomography, *compressed sensing* schemes have been investigated to solve the missing wedge problem, or to reduce the number of angles used [4–6]. Random sampling or jittered row sampling has also been used with in-painting to reduce dose and experiment time [7, 8]. In the x-ray regime, sparse studies are limited, but a low-rank matrix decomposition approach using PCA analysis of spectrotomography datasets has also recently been demonstrated,

45 which merges the angular and energy measurements to reduce overall measurement time [9].

46 *Low-rank matrix completion* is a well-known inverse problem that was widely used in sensor
47 networks [10], computer vision [11] and medical imaging [12, 13]. See e.g. [14] for a review of
48 those and further inverse problems. The classical formulation of matrix completion [15] aims to
49 recover the missing elements of a low-rank matrix from an incomplete set of known entries, each
50 of which is sampled *independently* at random (for example, uniformly).

51 However, there are important differences between x-ray and electron systems that make inde-
52 pendent sampling unattractive for the reduction of the acquisition time in x-ray spectromicroscopy.
53 An electron beam is rapidly moved across the specimen and can be blanked electrostatically at
54 high rate to control dose, and recreate random patterns. In contrast, an x-ray beam is fixed, and
55 the specimen is moved mechanically instead. Moreover, a mechanical chopping of the x-ray
56 beam is needed to recreate sampling patterns. The mechanical operations in the x-ray regime
57 limit the translation of some of these schemes from equivalent electron experiments.

58 The use of non-uniform sampling patterns in matrix completion has been widely studied, with
59 the main goal of reducing the approximation error. Non-uniform sampling methods include
60 adaptive cross interpolation [16], maximizing the volume of sampled submatrices [17], adaptive
61 importance sampling using leverage scores [18] or application-oriented distributions [19], and
62 supervised learning using a training dataset [13]. However, all of these methods still assume
63 some arbitrary control over the sampling pattern rather than conventional acquisition patterns
64 such as a raster scan.

65 In this paper we propose a novel approach for undersampling and reconstructing low-rank
66 x-ray spectromicroscopy data that uses the *raster sampling pattern*. Our motivation with this
67 work is to deliver a solution that can be deployed routinely at spectromicroscopy facilities by
68 non-experts. This requires an approach where no intervention or tuning is needed to produce
69 results. The method should also work with the standard raster acquisition approaches at these
70 facilities, in particular our initial experimental focus was on optimizing fast low-dose in-situ
71 experiments to study the evolution of battery materials over time. Scanning along a line can
72 be carried out with a faster mechanical movement, hence the time per pixel can be reduced
73 compared to independent or adaptive sampling. In addition, we can lower the x-ray dose on
74 the specimen without compromising the recovery of missing entries: we develop a procedure
75 to generate a robust raster pattern such that each row and column of the matrix has at least one
76 sampled element.

77 Besides the sampling pattern, the performance of matrix completion depends on the cost
78 function. In addition to the squared misfit of the sampled elements, the cost function may
79 include regularization terms, promoting sparsity [13] or smoothness [20]. Alternatively, the
80 entire completion problem can be turned into a Bayesian inference problem by introducing a
81 prior distribution on the low-rank matrix factors, treated as random matrices, and a likelihood of
82 the observed data samples [21, 22]. As a by-product, sparsity-promoting priors may provide an
83 automatic selection of the rank as the number of nonzero posterior components [23]. However,
84 mathematical study of spectromicroscopy is still in some infancy, and lacks well-recognised
85 priors. The only regularization assumption that is generally valid is the low-rankness of the true
86 absorption distribution. Therefore, we start with a simple Alternating Steepest Descent algorithm
87 (ASD) [24] with the squared misfit cost only. For a reliable selection of the rank and number of
88 samples, we propose an algorithm, LoopedASD, which successively increases the rank from 1 to
89 some generous value (e.g. 20), taking a lower-rank result as the initial guess in each step. This
90 provides a smoother convergence which allows the KNEEDLE [25] algorithm (already used in
91 spectromicroscopy for the PCA analysis) to determine the final rank accurately.

92 Five datasets were used in this study. The first three (labelled *DS1*, *DS2*, *DS3*) are full datasets
93 that can be undersampled numerically after the experiment; the final two (labelled *DS4*, *DS5*) are
94 measured using both full and sparse raster patterns to verify the experimental implementation.

95 The sparse experimental measurements were conducted at a range of undersampling ratios
 96 (details on the different sampling approaches can be found in Section 3). The specimen were
 97 produced by mixing Fe_2O_3 , Fe_3O_4 and FeO powders. The powders were ground in a ball mill
 98 and then drop cast onto a silicon nitride membrane. The dimensions of each data set can be seen
 99 in Table 1. The pairs (DS1, DS2) and (DS4, DS5) are data derived from the same specimen but
 100 at different spatial dimensions, i.e. DS2 is DS1 focused on a smaller area, as is DS4 of DS5. We
 101 aim to illustrate results using all 5 datasets, but where space is restrictive we only show the 3
 102 independent datasets: DS1, DS3, DS5.

Table 1. Spectromicroscopy Meta Data

Dataset	Data Acquisition	n_E	n_1	n_2
DS1	Full	149	101	101
DS2	Full	150	92	79
DS3	Full	152	55	54
DS4	Full & Sparse	152	40	40
DS5	Full & Sparse	152	80	80

103 The paper is organised as follows. We begin in Section 2 with a description of the x-ray
 104 spectromicroscopy model, and derive its low rank nature. Next, we discuss general matrix
 105 completion and raster-aware sampling, which leads into our proposed matrix completion algorithm
 106 in Sections 3 & 4. Finally, in Sections 5 & 6, we discuss the results comparing the reconstructed
 107 sparse data against full data.

108 2. Low rank model

109 When the energy of a photon increases beyond the binding energy of a core electron, we see a
 110 sharp rise in a material's absorption - an *absorption edge*. The absorption coefficient, μ , will vary
 111 or be modulated by the local chemical environment. Measuring around the edge, x-ray near edge
 112 absorption spectroscopy (XANES) can be used as a fingerprinting tool to identify known standards
 113 or materials; investigation of the energy past the absorption edge, the Extended X-ray Absorption
 114 Fine Structure (EXAFS), can be used to extract information of nearest neighbour bonding and
 115 coordination. Further details regarding the different experimental setup and derivation of the
 116 following formula can be found in [26].

117 The variation of μ with energy depends on how an element is chemically bonded and when
 118 there are more than one chemical states present the absorption coefficients sum linearly. For a
 119 mixture of S materials, the measured x-ray absorption (or optical density), $D(E)$, can be modeled
 120 as,

$$D(E) = \sum_{s=1}^S \mu_s(E) t_s, \quad (1)$$

121 where $\mu_s(E)$ and t_s , $s = 1, \dots, S$ are the distinct absorption coefficients (cm^{-1}), and thicknesses
 122 (cm) of the S materials, respectively.

123 With a focused x-ray beam, this experiment can be performed over $n_1 \times n_2$ positions, and at
 124 n_E distinct energies. By stacking each spatial scans, the distribution of x-ray absorption spectra
 125 can now be represented as a 3D tensor $D \in \mathbb{R}^{n_E \times n_1 \times n_2}$,

$$D_{i j_1 j_2} = \sum_{s=1}^S \mu_s(E_i) (t_s)_{j_1 j_2}, \quad (2)$$

126 where $(t_s)_{j_1 j_2}$ is the thickness of material s at pixel (j_1, j_2) , and $\mu_s(E_i)$ is the absorption
 127 coefficient of material s at the i -th energy level. This dataset can be flattened by considering a
 128 matrix $A \in \mathbb{R}^{n_E \times N}$, such that

$$A_{ij} = \sum_{s=1}^S \mu_s(E_i) t_{sj} + \eta_{ij}, \quad (3)$$

129 where $j = 1, \dots, N$, $N = n_1 n_2$ indexes over all pixels, and $\eta \in \mathcal{N}(0, \delta^2 I)$ represents Gaussian
 130 noise with standard deviation $\delta \in \mathbb{R}$. With a slight abuse of notation, consider the matrices
 131 $\mu \in \mathbb{R}^{n_E \times S}$, $\mu_{is} = \mu_s(E_i)$ and $t \in \mathbb{R}^{S \times N}$; the columns of μ represent the absorption coefficients
 132 of each material within the specimen, and the rows of t represent the thickness/presence of each
 133 material within each pixel. This allows us to write Eq. (3) as

$$A = \mu t + \eta. \quad (4)$$

134 This is significant, as it illustrates that for the majority of experiments, where the specimen is made
 135 up of only a few materials or chemical states (S is low), the corresponding spectromicroscopy
 136 data is inherently *low rank*, as $\text{rank}(\mu t) = S$. With the addition of noise, A is approximately low
 137 rank.

138 To analyse the spatial distribution of the absorption spectra, standard techniques are used to
 139 filter and decompose A back to a smaller set of representative spectra; see, for instance, [27].
 140 First, PCA is applied to reduce the noise in the data by producing a rank- L approximation of the
 141 most significant components. We compute

$$A' = C' R', \quad (5)$$

142 with $A' \in \mathbb{R}^{n_E \times N}$, $C' \in \mathbb{R}^{n_E \times L}$, $R' \in \mathbb{R}^{L \times N}$. The variable L is chosen to capture as much
 143 variation in the data as possible with the smallest rank, and should approximate S : it is typically
 144 set to be the elbow point (point of maximum curvature) of the singular values of A , and can be
 145 selected automatically using algorithms like KNEEDLE [25].

146 The decomposition in Eq. (5) is similar to the model in Eq. (3), however the columns of C' are
 147 abstract spectra - linear combinations of the true spectra with no physical interpretation. Similarly,
 148 the rows of R' are the corresponding abstract thickness maps. Pixels are now clustered together
 149 based on the similarity of the normalised mixing factors of the PCA components. This is achieved
 150 by clustering the columns of R' using standard clustering algorithms such as kmeans [28] and
 151 lvq [29]. Taking the mean spectrum from the columns of A for each cluster increases the signal to
 152 noise ratio when compared to the measurements from an individual pixel, and produces accurate
 153 x-ray absorption spectra for the dominant material in each cluster.

154 In [27], it is noted that the first principal component from the PCA (the component with
 155 the greatest variation in the data) often describes the average x-ray absorption data across the
 156 whole specimen. In some cases the first principal component is discarded, or scaled down, and
 157 cluster analysis is only applied to the remaining components. This is done to emphasise the
 158 more subtle features and variations in the data and ensure the clustering results are determined
 159 by differences in materials not the thickness of the specimen. In this paper, we will refer to the
 160 process of discarding the first principal component as Reduction of Thickness Effect (RTE),
 161 which is achieved by simply removing the first row from R' in Eq. (5) before applying kmeans
 162 to its columns. Generally, the comparison between reconstructed and full data are worse when RTE
 163 is used, and they have been applied to several tests to provide the worst-case results. Any cluster
 164 results that have used RTE will be noted clearly. Further details on PCA, cluster analysis and
 165 RTE can be found in the supplemental document.

166 Figure 1 shows a schematic of the sparse spectromicroscopy process. In our proposed scheme,
 167 we measure only a small proportion of the data, and recover the missing entries later using low rank

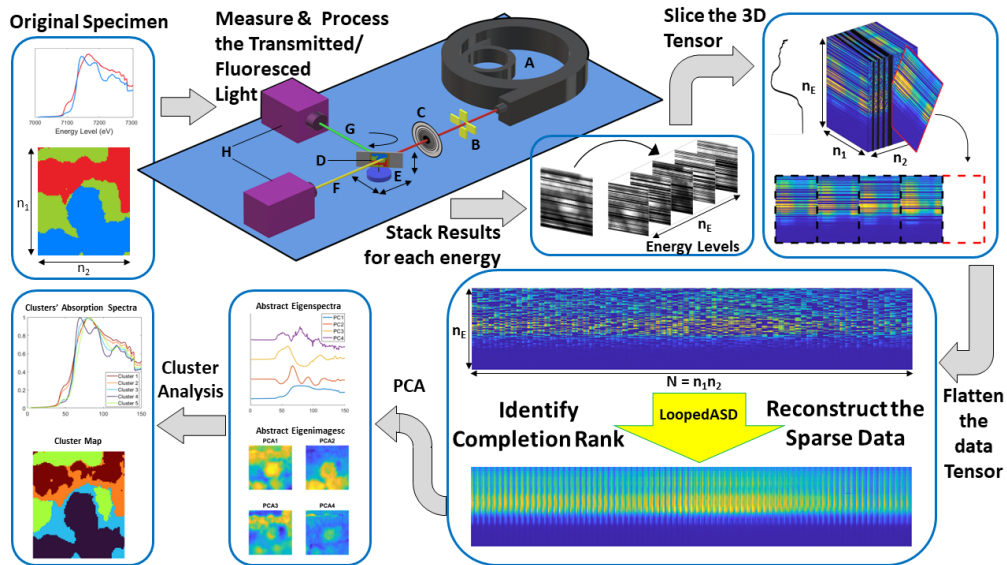


Fig. 1. Schematics of Sparse Spectromicroscopy using DS2. The original specimen contains a mixture of two materials: FeO in the blue region and Fe_2O_3 in the red region over the background (green region). The experimental setup is as follows: (A) a third/fourth generation synchrotron light source produces a beamline of energy E_i . (B) the intensity of the incident beam is measured. (C) the light is focused using an optic such as a zone plate or Kirkpatrick-Baez mirror to produce a micro or nanoscale beam. (D) The specimen is mounted on the stage, (E) which is able to move the specimen across in the XY plane to measure at different positions; the specimen can also be moved in the Z direction to bring the sample to the focus of the beam. Depending on the material and it's thickness, we measure *either* the transmitted flux, (F), *or* the fluoresced flux, (G), by counting photons using different detectors (H). The specimen is moved in a raster pattern; for sparse experiments we only sample pn_1 rows for each energy. The processed results for each energy are stacked as seen in the diagram. The tensor is flattened by concatenating slices of the sparse data. The missing entries are recovered using LoopedASD - our low rank matrix completion algorithm. PCA is applied to the completed data to produce the most significant components in the data - the eigenspectra and (reshaped) eigenimages seen in the schematic. The results of the cluster analysis with 5 clusters show similar results to the original specimen, with the locations of the two materials identified, and accurate absorption coefficients for each.

168 matrix completion. The remaining steps are consistent with a standard x-ray spectromicroscopy
 169 scheme. Precise details on the sampling are discussed in Section 3, and the completion methods
 170 are described in Section 4. The image illustrates the data acquisition, the formatting and flattening
 171 of the data tensor, the reconstruction of the sparse entries, and the PCA and cluster analysis.

172 3. Raster-aware sampling patterns

173 To formulate the sampling and reconstruction of x-ray spectromicroscopy data, we set out the
 174 following notation. Let $\Omega \subset \{1, \dots, n_E\} \times \{1, \dots, N\}$ be the set of known measurements, called
 175 the **sampling pattern**. For a matrix $X \in \mathbb{R}^{n_E \times N}$, we define the **sampling operator** \mathcal{P}_Ω as

$$\mathcal{P}_\Omega(X) = \begin{cases} X_{i,j}, & \text{if } (i, j) \in \Omega, \\ 0, & \text{if } (i, j) \notin \Omega. \end{cases} \quad (6)$$

176 Alternatively, the sampling pattern can be thought of as a binary matrix $\Omega \in \{0, 1\}^{n_E \times N}$ with
 177 1s in the locations of the known entries, and 0s everywhere else. It is then easy to compute
 178 $\mathcal{P}_\Omega(X) = \Omega \circ X$, where \circ is the Hadamard product. The key parameter for matrix completion is
 179 the **undersampling ratio**, p , the proportion of known entries:

$$p = \frac{|\Omega|}{n_E N}. \quad (7)$$

180 The standard matrix completion problem involves computing a matrix of minimal rank such
 181 that the known entries, indexed by Ω , are equal to the given values. However, this problem is
 182 notoriously difficult, and many algorithms will instead seek to solve easier, but provably related,
 183 problems; see [30–33].

184 For this application, the presence of noise means the spectromicroscopy datasets are only
 185 approximately low rank - i.e. there is no exact low rank matrix that would perfectly match the
 186 subset of known entries $\mathcal{P}_\Omega(A)$. A more useful and efficient approach to reconstruction is to
 187 fix the rank, r , then solve an optimisation problem to find the best rank- r approximation to the
 188 known entries. Thus, to reconstruct the sparse set of measurements, $\mathcal{P}_\Omega(A)$, we seek to solve,

$$\min_{Z \in \mathbb{R}^{n_E \times N}} \frac{1}{2} \|\mathcal{P}_\Omega(A) - \mathcal{P}_\Omega(Z)\|_F^2 \quad \text{subject to} \quad \text{rank}(Z) = r. \quad (8)$$

189 Here, we use the frobenius norm $\|\cdot\|_F$, defined for $X \in \mathbb{R}^{n_E \times N}$,

$$\|X\|_F^2 = \sum_{i,j} X_{ij}^2. \quad (9)$$

190 This approach is generally more effective (and far more computationally efficient) than attempting
 191 to solve the standard problem; we must, however, correctly input the completion rank r , since
 192 completion algorithms work best when r is close to the approximate rank of the full dataset.
 193 More details on accurately setting r can be found in the supplementary material.

194 3.1. Setting the sampling pattern

195 In many applications of low rank matrix completion, it is impossible to predict which entries
 196 will be known. To model this, the known entries of the sampling pattern are typically set at
 197 random. Usually, *Bernoulli sampling* is used, where each entry is sampled i.i.d. (independent
 198 and identically distributed) with probability p .

199 Unlike other settings, in x-ray spectromicroscopy we have complete control over the data
 200 acquisition and can set the scanner to implement specific patterns. However, when formulating
 201 the sampling selection, we must consider the physical restrictions of the experiment. For x-ray
 202 spectromicroscopy, the specimen is on an XY stage and moves at constant velocity through the
 203 beam in a *raster pattern* - scanning across each row in turn before switching to the next energy
 204 level and repeating the spatial scan. To maximise the efficiency of the modelled sparse scans, we
 205 must ensure the known entries of the sampling pattern are collected together into spatial rows to
 206 be scanned. The scanner can then move between known entries, quickly passing over the empty
 207 rows.

208 To model the raster aligned patterns, we introduce *Raster sampling*, where spatial rows of
 209 the specimen are sampled i.i.d with probability p . Illustrations of both Bernoulli and Raster
 210 sampling patterns on flattened data sets can be found in Figure 2, in which the known entries are
 211 highlighted yellow. Notice that the raster sampling pattern groups the known entries into the
 212 spatial rows of the specimen, which appear as short segments after the data is flattened.

213 Data that has been sampled with a Bernoulli pattern generally allows completion at lower un-
 214 dersampling ratios, because the known entries are spread more evenly. In practice, undersampling

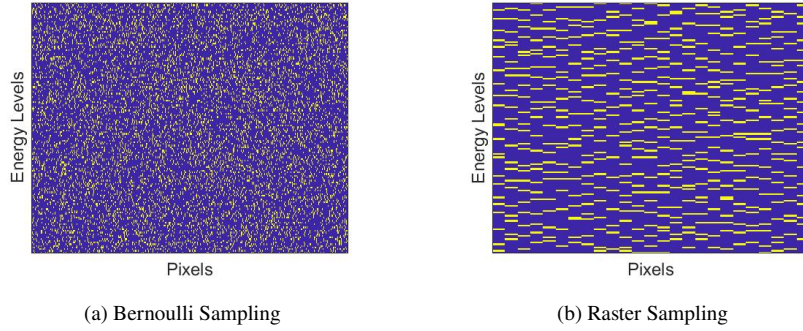


Fig. 2. Illustrations of Bernoulli and Raster sampling patterns. Patterns are shown on flattened datasets with $n_E = 150$, $n_1 = 25$, $n_2 = 25$, $p = 0.2$. Yellow indicates the point has been sampled, while purple indicates the entry is missing.

215 using Bernoulli sampling is possible, however it is only preferable under certain circumstances.
 216 Sampling patterns like these can be implemented for x-ray spectromicroscopy in two ways: either
 217 the specimen is moved through every position while rapidly blanking the beam, or the specimen
 218 is only positioned at the location of each known entry using stop-start motion. The former
 219 is difficult to implement and doesn't reduce the experiment time, however may be useful for
 220 dose reduction. Conversely, despite the potential use of lower undersampling ratios, stop-start
 221 motion is generally less efficient than the continuous motion used for raster sampling, resulting in
 222 longer experiments. Indeed, for typical dwell times of around $0.01s$, the time spent accelerating,
 223 decelerating and stabilising the scanner outweighs the time saved by scanning fewer entries.
 224 Despite this, Bernoulli sampling may be applicable for specimen that require longer dwell times
 225 (around $1s$), since reducing the number of known entries has a greater impact on the experiment.
 226 Developing methods for such specimen is left for further research.

227 One issue with Raster sampling is that the grouping of known entries increases the probability
 228 that certain rows are not scanned at all, especially for low undersampling ratios. Clearly, it
 229 is impossible to recover a row or column of the data set A (distinct from the spatial rows and
 230 columns of the specimen) that contains no known entries, so we must ensure that every spatial-row
 231 is measured at least once. By testing different sampling patterns, we have also seen that low rank
 232 completion is more reliable when the known entries are more evenly spread across the data.

233 To promote the spread of data, and ensure every row and column of A contains known entries,
 234 we have developed *Robust Raster sampling*. This is a variation of Raster sampling that ensures
 235 every row of the specimen is sampled exactly once before any row can be sampled a second time.
 236 Further details on Robust Raster Sampling can be found in the supplementary document.

237 4. Completion algorithm

238 To reconstruct undersampled data, we have developed a low rank matrix completion algorithm
 239 **LoopedASD**. This is a rank-incremental algorithm based on Alternating Steepest Descent
 240 (ASD) [24]; it uses the output of one ASD-completion as the input for the next higher-
 241 rank completion. After testing several different options, [31–33], it was found that this
 242 algorithm provided more accurate and reliable results on both simulated data and raster-sampled
 243 spectromicroscopy data.

244 ASD is an iterative method that fixes the rank, r , of its iterates by imposing the following
 245 decomposition:

$$Z = XY, \quad \text{for } Z \in \mathbb{R}^{n_E \times N}, X \in \mathbb{R}^{n_E \times r}, Y \in \mathbb{R}^{r \times N}. \quad (10)$$

246 Thus, the optimisation of the problem in Eq. (8) now becomes,

$$\min_{X,Y} f(X,Y), \quad \text{where} \quad f(X,Y) = \frac{1}{2} \|\mathcal{P}_\Omega(A) - \mathcal{P}_\Omega(XY)\|_F^2, \quad X \in \mathbb{R}^{n_E \times r}, \quad Y \in \mathbb{R}^{r \times N}. \quad (11)$$

247 We minimise the function f by alternately optimising the components X and Y using steepest
 248 descent with exact step sizes. Indeed, by fixing one component the gradient of f can be computed
 249 easily. It is then possible to analytically compute the exact step size required to minimise f along
 250 the gradient direction. Once the factor has been updated, we alternate the fixed component and
 251 repeat the process, iterating until suitable stopping conditions are satisfied.

252 Let $f(X,Y)$ be written as $f_Y(X)$ for fixed Y and $f_X(Y)$ for fixed X ; the corresponding gradients
 253 are written $\nabla f_Y(X)$ and $\nabla f_X(Y)$, and the exact step sizes are written η_X , η_Y . Beginning with
 254 random matrices $X_0 \in \mathbb{R}^{n_E \times r}$, $Y_0 \in \mathbb{R}^{r \times N}$, we implement ASD as follows,

$$\begin{cases} \text{Fix } Y_i, \text{ compute } \nabla f_{Y_i}(X_i) \text{ \& } \eta_{X_i} \\ X_{i+1} = X_i - \eta_{X_i} \nabla f_{Y_i}(X_i) \\ \text{Fix } X_{i+1}, \text{ compute } \nabla f_{X_{i+1}}(Y_i) \text{ \& } \eta_{Y_i} \\ Y_{i+1} = Y_i - \eta_{Y_i} \nabla f_{X_{i+1}}(Y_i). \end{cases} \quad (12)$$

255 The gradients and step sizes are given below. Full derivations can be found in the supplemental
 256 documents:

$$\nabla f_Y(X) = -(\mathcal{P}_\Omega(A) - \mathcal{P}_\Omega(XY^T))Y, \quad \nabla f_X(Y) = -X^T(\mathcal{P}_\Omega(A) - \mathcal{P}_\Omega(XY^T)). \quad (13)$$

$$\eta_X = \frac{\|\nabla f_Y(X)\|_F^2}{\|\mathcal{P}_\Omega(\nabla f_Y(X)Y^T)\|_F^2}, \quad \eta_Y = \frac{\|\nabla f_X(Y)\|_F^2}{\|\mathcal{P}_\Omega(X[\nabla f_X(Y)]^T)\|_F^2} \quad (14)$$

257 One advantage of ASD is that the residual $(\mathcal{P}_\Omega(A) - \mathcal{P}_\Omega(XY))$ can be easily updated between
 258 iterations, removing the need to compute the matrix product XY for each iteration. Thus, the per
 259 iteration cost of ASD has leading order $8|\Omega|r$ (see [24]).

260 A maximum number of iterations was set as a stopping condition, as well as a tolerance on the
 261 relative norm of the residual at each iteration,

$$\frac{\|\mathcal{P}_\Omega(A) - \mathcal{P}_\Omega(X_i Y_i)\|_F^2}{\|\mathcal{P}_\Omega(A)\|_F^2}. \quad (15)$$

262 Once either stopping condition is satisfied, the two factors are output by the algorithm and are
 263 denoted X^* and Y^* .

264 To evaluate the success of the completion algorithms, we compute the *completion error*, often
 265 denoted e_c . This is simply the relative norm of the difference between the true low rank matrix
 266 A and the output of the ASD algorithm $A^* = X^*Y^*$, defined as,

$$e_c = \frac{\|A - A^*\|_F^2}{\|A\|_F^2} \quad (16)$$

267 A more detailed description of the algorithm can be found in the supplemental material.

268 4.1. Improving results with LoopedASD

269 Following rigorous testing of ASD for raster sampling, it was found that there exists an approximate
 270 optimal undersample ratio, p^* , that depends linearly on, and is correlated with, the data's rank r .
 271 For $p > p^*$, the probability of a successful completion was almost certain, and the completion

272 errors were consistently low. For $p < p^*$, the completion errors increased and the completion
 273 rate (the proportion of tests that successfully recover data to a certain accuracy - usually 10^{-4})
 274 decreased rapidly to zero. Intuitively, the dependence of p^* on r makes sense: to reconstruct
 275 more complex datasets (with higher ranks), more known entries are required to successfully
 276 recover the missing ones.

277 *LoopedASD* was developed to take advantage of this heuristic: beginning with $r = 1$, we use
 278 the outputs of the $r = j$ completion as initial guesses for the $r = j + 1$ completion, iterating up to
 279 the completion rank set by the user.

280 Let $(X_0)^{(j)} \in \mathbb{R}^{n_E \times j}$, $(Y_0)^{(j)} \in \mathbb{R}^{j \times N}$ be the initial matrices for the j^{th} rank step, and
 281 $(X^*)^{(j)} \in \mathbb{R}^{n_E \times j}$, $(Y^*)^{(j)} \in \mathbb{R}^{j \times N}$ be the outputs of the j^{th} iteration of ASD with completion
 282 rank $r = j$. In order ensure the dimensions of the factors increase with each rank increment, we
 283 simply concatenate the j^{th} outputs $(X^*)^{(j)}$, $(Y^*)^{(j)}$ with a random column and row respectively
 284 to produce $(X_0)^{(j+1)}$, $(Y_0)^{(j+1)}$.

285 The aim of this variation is to allow the iterates to converge quickly to a rank-1 approximation
 286 of the sparse data, where fewer known entries are required. Once a rank- j approximation is
 287 known, the distance to the rank- $(j + 1)$ solution should be relatively low, and again fewer known
 288 entries are required to converge quickly to the next one. Ensuring iterates remain close to the
 289 minimum should also avoid convergence to a spurious local minima of f (which is non-convex).

290 The completion results of *LoopedASD* confirm the validity of this heuristic, as it yields good
 291 reconstructions more reliably at low undersampling ratios.

292 In addition to the usual stopping conditions, an early stopping procedure was required due
 293 to slow convergence rates when the iterates approached optimal solutions. In the case that
 294 the average change in the norm of the residual (Eq. (15)) over 50 iterations drops below a
 295 second tolerance (typically 10^{-5}), then the early stopping condition is satisfied and will halt the
 296 algorithm.

297 It has been noted that the completion rank must be set before implementing ASD and
 298 *LoopedASD*. This, however, is not an obvious choice since we cannot evaluate the accuracy of
 299 the different rank- k completions without prior knowledge of the results. To overcome this, we
 300 have developed a method that implements short completions over several ranks and evaluates the
 301 completion errors using cross validation so that we can efficiently identify the optimal completion
 302 rank to use. A full description and justification of the method can be found in the supplemental
 303 material.

304 5. Validation of the algorithm with numerical undersampling

305 We test the method in two distinct ways. First, we numerically sample full datasets. To produce
 306 these datasets, ‘unknown’ entries are set to zero according to a randomly generated robust raster
 307 sampling pattern. Then, in Section 6, we implement sparse scans on the beamline to determine if
 308 any additional issues arise and to test in a new setting outside of our test set.

309 To properly evaluate the performance of the methods described in this paper, several measures
 310 of success must be considered. We can compare the original and reconstructed datasets directly
 311 by computing the average completion error. Perhaps more significantly, we evaluate the difference
 312 between the cluster maps and the absorption spectra themselves. Finally, we can plot and
 313 visually compare the clusters and absorption spectra. Due to specimen drifts (described in
 314 Section 6), the first two ‘computational’ approaches can only be achieved using numerically
 315 sampled data; visual comparisons must be used when evaluating reconstructions using the sparse
 316 scan data. During these tests, we use the optimal completion rank r^* determined by the rank
 317 selection algorithm described in the supplemental document to produce the completed matrix
 318 $(A^*)^{(r^*)} = (X^*)^{(r^*)}(Y^*)^{(r^*)}$. A^* is then used instead of the full data set A in the analytic processes
 319 (PCA and clustering).

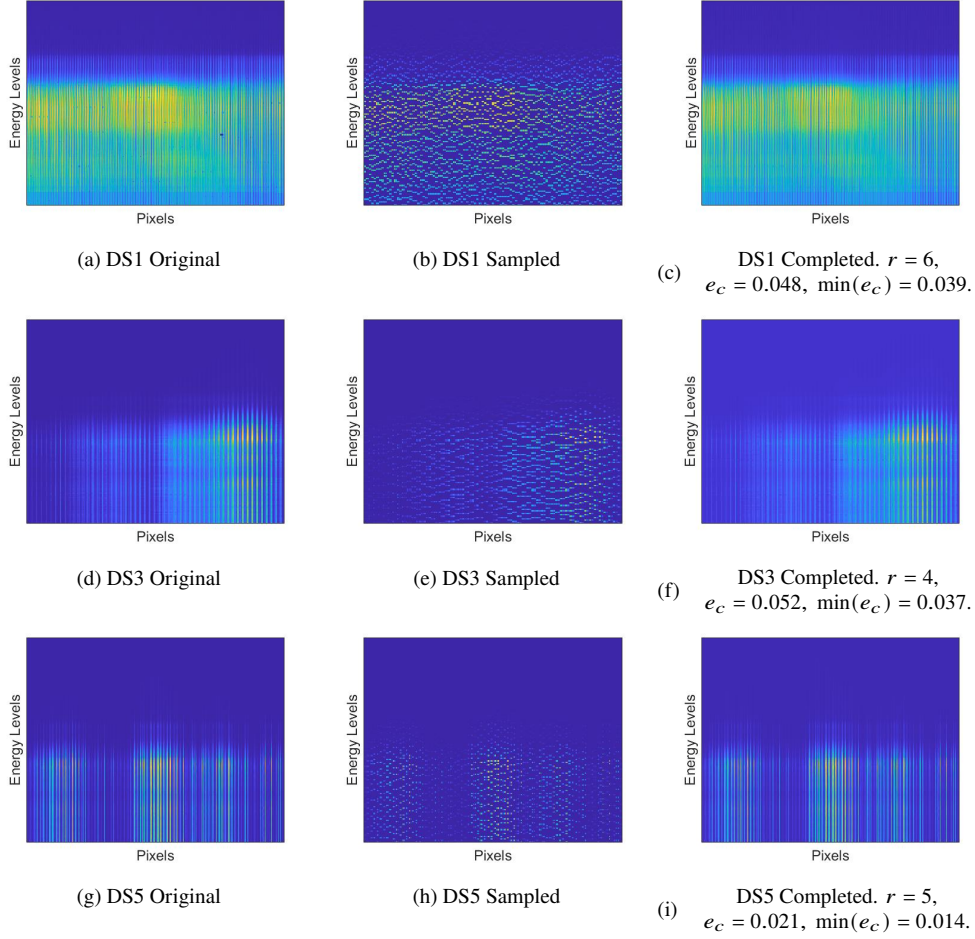


Fig. 3. Reconstructions of the three independent datasets, using LoopedASD as the completion algorithm. Each figure is a scaled colour image of the flattened dataset, with energy levels across the vertical axis, and all pixels across the horizontal axis. The left column shows the original datasets. The middle column shows the sampled data at $p = 0.20$. The right column shows the reconstructed data. Below each reconstruction is the completion rank, r , the completion error e_c (see Eq. (16)), and the minimum approximation error (see Eq. (17)) for that completion rank.

320 5.1. Completion errors of LoopedASD

321 We first examine the norm of the difference between data sets. When conducting experiments
 322 such as this, it is important to consider the constraints on the completion error, e_c (Eq. (16)), for
 323 a rank- r completion.

324 Consider a rank- r approximation of a matrix A , denoted $A^{(r)}$. We can compute the minimum
 325 approximation error of $A^{(r)}$ in the Frobenius norm using the Eckart-Young-Mirsky (EYM)
 326 Theorem [34]. We shall refer to this value as the minimal rank- r approximation error, and it is
 327 given by:

$$\min_{\text{rank}(A^{(r)})=r} \|A - A^{(r)}\|_F = \sqrt{\sigma_{r+1}^2 + \sigma_{r+2}^2 + \dots + \sigma_n^2}, \quad (17)$$

328 where $\sigma_1, \dots, \sigma_n$ are the singular values (SVs) of A . Since any completion result with a

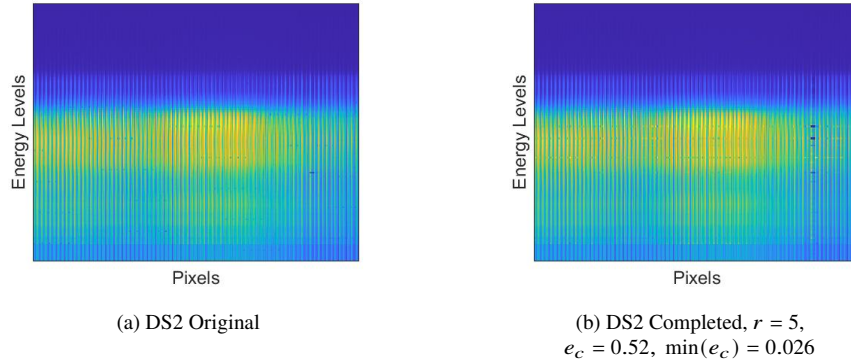


Fig. 4. Completion result following intentional sampling of the artifact in DS2 (dark spot towards the right). Note that the corrupted points are contained to the column of the initial artefact.

329 completion rank of r is just an example of a rank- r approximation, the corresponding minimum
 330 approximation errors are lower bounds for the corresponding completion errors. Thus the success
 331 of any completion should be evaluated by comparing back to the corresponding minimum error.

332 LoopedASD produces successful completions for all datasets, and for undersampling ratios as
 333 low as $p = 0.15$. By examining the mean completion errors, e_c , one can identify the optimal
 334 undersampling ratio, p^* . As before, for $p < p^*$ average completion errors rise quickly, and
 335 for $p > p^*$ average completion errors decrease slowly, reducing the benefit of taking further
 336 measurements. It was found that LoopedASD produces lower completion errors more reliably
 337 than ASD, in particular for undersampling ratios around p^* .

338 To help with the visualization of the completion, we plot the flattened data sets themselves as
 339 colourmap images. Purple points indicate a value of zero, while bright yellow points indicate
 340 higher measured values. In Figure 3 we provide a visual representation of the sampling and
 341 reconstruction process at the optimal rank, r^* , and with a consistent undersampling ratio of
 342 $p = 0.20$. Note that completions often appear smoother than the original data - since it is low
 343 rank, it will have filtered out much of the noise.

344 One advantage of ASD and LoopedASD is that the impact of any artifacts that have been
 345 sampled remain contained in their rows/columns in the dataset. In Figure 4 we illustrate this
 346 by intentionally sampling a corrupted entry in DS2. We can see that the majority of the image
 347 remains accurate and it is just the row and column that contained the original artefact that are
 348 affected. Indeed, it is very easy to identify the corrupted regions, which can be cut from the
 349 dataset. The example in Figure 4 initially had a completion error of 0.052, compared to the mean
 350 0.33 (the corresponding minimum approximation error is 0.026). Once the affected entries are
 351 removed, e_c is computed as 0.047 providing a much stronger result. In practice this property is
 352 very useful, since it implies the process is robust against artefacts found in the data.

353 5.2. Impact of completion on cluster analysis

354 We now compare the results of the cluster analysis from the full data and the sparse data. In some
 355 sense, this comparison is more relevant than simply comparing the completion errors because the
 356 success of the new methods will be determined by the quality of the cluster maps produced by
 357 the reconstructions, not necessarily the abstract difference between the two datasets.

358 We can evaluate the similarity of clusters using the Adjusted Rand Index (ARI) [35]. This is
 359 a symmetric score, with values between -1 and 1 . An ARI score of 1 indicates clusters are a
 360 perfect match, while a score of 0 implies classifications have effectively been assigned at random.

361 Generally we use the inverted score of $(1 - \text{ARI})$ to evaluate cluster quality so that 0 indicates a
 362 perfect match, and smaller scores are generally better.

363 On the other hand, we can compare individual absorption coefficients using the Euclidean
 364 2-norm, but we must ensure that we are taking the difference between coefficients for equivalent
 365 materials/clusters. Suppose we have performed cluster analysis on both the full data set and the
 366 completed data with N_{cluster} many cluster centres. Let $\mu_f \in \mathbb{R}^{n_E \times N_{\text{cluster}}}$ denote the absorption
 367 coefficients of the full data and $\mu_c \in \mathbb{R}^{n_E \times N_{\text{cluster}}}$ denote those from the completed data. In
 368 both cases, the i^{th} columns, $(\mu_f)_i$, $(\mu_c)_i$ respectively, represent the absorption coefficient
 369 corresponding to the i^{th} cluster. Note that the i^{th} clusters from the full data may not correspond to
 370 the i^{th} cluster from the completed data. We now align the clusters by taking each full-data cluster
 371 and finding the completed cluster that best fits it (minimises $(1 - \text{ARI})$); using this permutation
 372 index, we rearrange the columns of μ_c to get the aligned absorption coefficients, $\mu_a \in \mathbb{R}^{n_E \times N_{\text{cluster}}}$.
 373 We can now be sure the absorption coefficients represent the same area and the same material.
 374 We now wish to combine the relative differences of each cluster such that each cluster is weighted
 375 equally (not all absorption spectra are of the same magnitude) and is normalised for the number
 376 of clusters used. Thus, we compute the *spectral difference*, d_{spec} , as:

$$d_{\text{spec}} = \frac{1}{\sqrt{N_{\text{clusters}}}} \sqrt{\sum_{i=1}^{N_{\text{clusters}}} \left(\frac{\|(\mu_f)_i - (\mu_a)_i\|_2}{\|(\mu_a)_i\|_2} \right)^2}. \quad (18)$$

377 One complication is that cluster analysis is not deterministic, since there is some randomness in
 378 the starting vectors of the cluster centres. Because of this, we often find several common cluster
 379 maps resulting from the same dataset. This is true for both the full datasets and the reconstructed
 380 datasets, which can produce false ARI scores and spectral differences when opposing outputs are
 381 compared. To overcome this, we apply a clustering process to the clusters themselves - grouping
 382 similar maps together into ‘Super Clusters’. Taking the mode within each super cluster produces a
 383 series of representative cluster maps that are produced by the original data. We can now compare
 384 the reconstructed cluster against each representative in turn, recording the highest score as the
 385 most appropriate comparison.

386 The spectromicroscopy datasets were sampled numerically, their optimal ranks were computed,
 387 and the data reconstructed using LoopedASD with early stopping. The completion ranks were
 388 set to be the optimal rank, r^* , using the method described in the supplementary document.
 389 PCA is used to decompose the data into its most significant components (we set the number
 390 of components used $L = r^*$, so the rank is consistent). Finally, we perform the cluster analysis
 391 using kmeans with 5 clusters and compute the associated absorption spectra. For DS1, DS2,
 392 DS3 we used RTE (dropping the first principal component before applying kmeans) to provide
 393 the worst-case results. Note that if RTE had not been used, both the ARI scores and the spectral
 394 difference would improve. For DS4 and DS5, we avoided using RTE due to the poorer quality
 395 of the clustering results for both the full and reconstructed data. When using RTE on DS4 and
 396 DS5, the clusters are not smooth, the images and spectra were noisy, and the outcomes were not
 397 representative of typical results.

398 Averaging over many iterations, we plot the clustering results for each dataset and for various
 399 undersampling ratios in Figure 5. In Figure 5a we plot $\log_{10}(1 - \text{ARI})$ against the undersampling
 400 ratio, and in Figure 5b we plot the log of the spectral difference, $\log_{10}(d_{\text{spec}})$, against the
 401 undersampling ratio. The purpose of the plots is to illustrate the qualitative behaviour of the
 402 results, so have translated the individual plots vertically to improve clarity. The quantitative
 403 results can be found in Table 2. Here we see that, as we increase the undersampling ratio, the
 404 $(1 - \text{ARI})$ scores and spectral differences behave similarly to the completion error, e_c (Eq. (16)).
 405 Indeed, by plotting the logs of these scores we see the similarly shaped curves - a faster decline
 406 for lower p that flattens for higher p . Once again, we identify the optimal undersampling ratio,

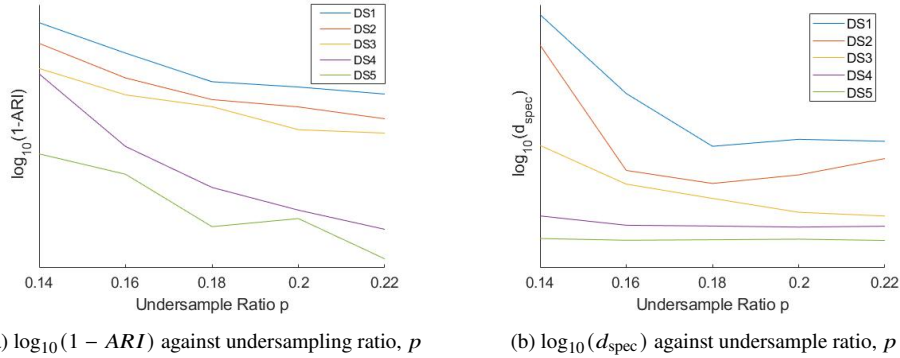


Fig. 5. Comparing \log_{10} of the cluster results against undersampling ratio using the optimal completion rank. In each figure, *lower scores are better*. Note that the individual plots have been translated vertically to improve the clarity of the image.

407 p^* , at the elbow of the plots - the points of maximum curvature. It is clear from the plot where
 408 the majority of the optimal undersampling ratios lie, however for DS2 and DS4 the elbows occur
 409 at different points across the two plots ($p = 0.16$ for Figure 5a and $p = 0.18$ for Figure 5b). In
 410 these cases, we take the higher value to ensure enough entries are known for better reliability.
 411 For DS4 and DS5, the pattern in Figure 5b is less clear and appears more as a constant value.
 412 This is simply because the variation in spectral difference over this range of undersample ratios is
 413 much smaller than for the other datasets, and by appropriately scaling these plots one can see the
 414 same curve as before.

415 In Table 2, we record the optimal completion rank, the corresponding optimal undersampling
 416 ratio (found at the elbow of the results in Figure 5), the completion errors, e_c , and the cluster
 417 results produced by these parameters.

Table 2. Optimal Clustering Results for numerically sampled Spectromicroscopy datasets. For the metrics used (e_c , $1 - \text{ARI}$, d_{spec}), *lower scores are better*.

dataset	Optimal Completion Rank, r^*	Optimal undersampling ratio, p^*	Mean Completion Error (16)	RTE used	Mean (1-ARI)	Mean Spectral Difference (18)
DS1	6	0.18	0.0483	Yes	0.106	0.082
DS2	5	0.18	0.0326	Yes	0.108	0.115
DS3	4	0.20	0.0515	Yes	0.207	0.509
DS4	7	0.18	0.025	No	0.009	0.096
DS5	5	0.18	0.027	No	0.006	0.152

418 In Figures 6 & 7, we visually compare the cluster maps and absorption spectra of the full
 419 data and the reconstructed data for DS1, DS2 and DS3 (DS4 and DS5 are used in Section 6 to
 420 illustrate the sparse scans). The data sets are sampled numerically with $p = 0.20$, the optimal
 421 completion rank was used, and each of these results uses RTE. These images were produced by
 422 Mantis X-ray [36], an open source software package used for analysing x-ray spectromicroscopy
 423 data.

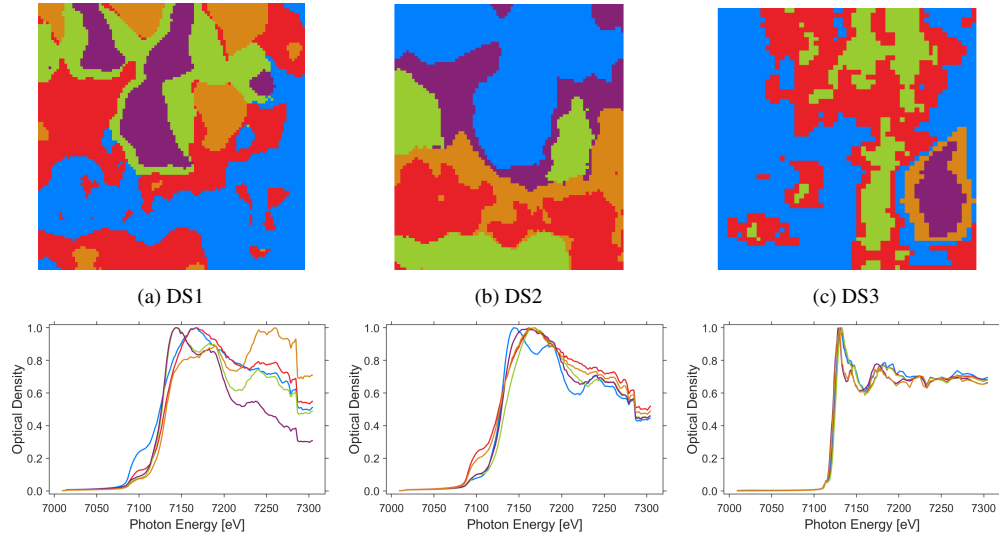


Fig. 6. Cluster and spectral results from full datasets using RTE. Results computed using Mantis-xray [36].

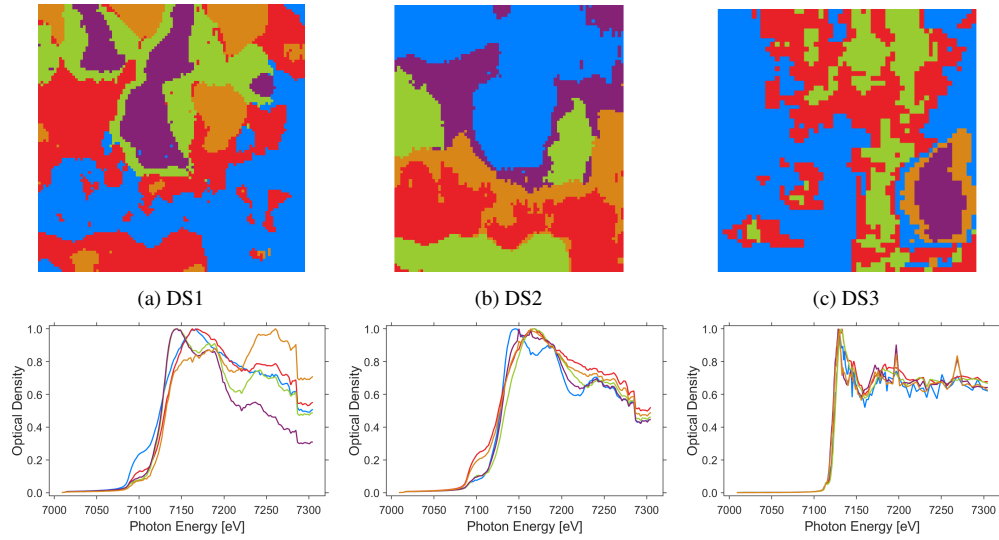


Fig. 7. Cluster and spectral results from reconstructions of sampled data, with $p = 0.20$, using RTE. Results computed using Mantis-xray [36]. Please note that for DS1 and DS2 there is a sharp downward going discontinuity at around 7270eV in both the full and completed results. This feature is present in the original data, and can be seen in the absorption spectra of many of the pixels with stronger signals. It is likely due to an error in the calibration of the energy stepping, which requires coordinated motion of an insertion device and monochromator, resulting in a strong variation in the intensity of the beam which was not normalised correctly.

424 We can see that both the cluster maps and absorption spectra are very similar between full and
 425 reconstructed data, with only some small (mostly insignificant) differences in each. The general
 426 outline of the reconstructions' clusters are very clear and are consistent with the full data results.
 427 The boundaries of the completion clusters lose some of their smoothness, but only a few isolated

Table 3. Time and efficiency (ratio of sparse and full scan times) of implementing sparse scans during experiments.

p	T (hrs : min : s) - DS4	T (hrs : min : s) - DS5	Efficiency - DS4	Efficiency - DS5
1.00	1:22:05.5	4:10:36.7	1.00	1.00
0.35	0:38:57.6	1:39:31.6	0.47	0.40
0.25	0:33:29.8	1:21:33.8	0.41	0.33
0.15	0:26:02.3	0:50:37.9	0.32	0.20

428 pixels have been misclassified with minimal impact to the overall image. Similarly the most
 429 important features of the absorption spectra have been preserved in the completions for DS1 and
 430 DS2, including the pre-edges, the peaks, and in particular the shifts in the absorption edges for
 431 each material. There are a few examples of noise affecting the data: DS3 is particularly noisy,
 432 where for some clusters we see the sharp spikes indicative of noise. Despite this, the general
 433 outline is still consistent.

434 Overall, the similarity of the clusters ensures the sparse data would be interpreted in the same
 435 way as the original data, and the absorption spectra are sufficiently clear to be able to identify the
 436 materials within each cluster.

437 6. Practical sparse scanning experiments

438 We now bring all of the above material together to implement sparse scanning in practice. Instead
 439 of numerically sampling a full dataset, we physically implemented a sparse robust raster sampling
 440 pattern. This was done on the same specimen, with the same pixel sizes and dwell times for
 441 $p \in \{0.15, 0.25, 0.35, 1\}$.

442 Unfortunately, it is not possible to numerically compare the results of the data due to
 443 experimental variations between the known entries in each dataset. Small changes in the ambient
 444 temperature can cause the specimen to drift very small distances between each spatial scan
 445 during the experiment. Since measurements are being taken on the micro/nano-scopic scale, this
 446 effect can create pixel level change from start to end. This is usually compensated for during the
 447 stacking process, but sparse scans have a much shorter experimental time, and registration to
 448 correct drifts cannot be performed in the same way as full datasets. The intensity of the incident
 449 beam will also vary with time and, although this can be normalized, some variation can still occur.
 450 Because of these potential spatial discrepancies and changes in intensity and background noise,
 451 the same known entries across different scans will show different values. Thus, it is impossible
 452 to know whether the difference between the full and reconstructed data is because of differences
 453 in the measurements or due to the limitations of the completion algorithm.

454 Despite this, we can still *visually* compare the results. For each of our sparse scans (DS4 and
 455 DS5) and at each undersampling ratio, we use LoopedASD to reconstruct the data. We equip
 456 the algorithm with the optimal completion rank for the data set, described in the supplementary
 457 documents. We use Mantis X-ray to perform the cluster analysis, setting $L = r^*$ and using 5
 458 clusters. Recall that RTE was **not** used for these results. In Figures 8 and 9, we plot these results.

459 Once again, despite a slight loss of smoothness around cluster boundaries and some individual
 460 miss-classified pixels, the overall structure of the cluster maps is near identical and clearly shows
 461 the same variations across the scans. Similarly, there are a few sharp discontinuities in the
 462 reconstructed absorption spectra, especially for lower undersample ratios, but each materials'
 463 XANES are still clearly identifiable.

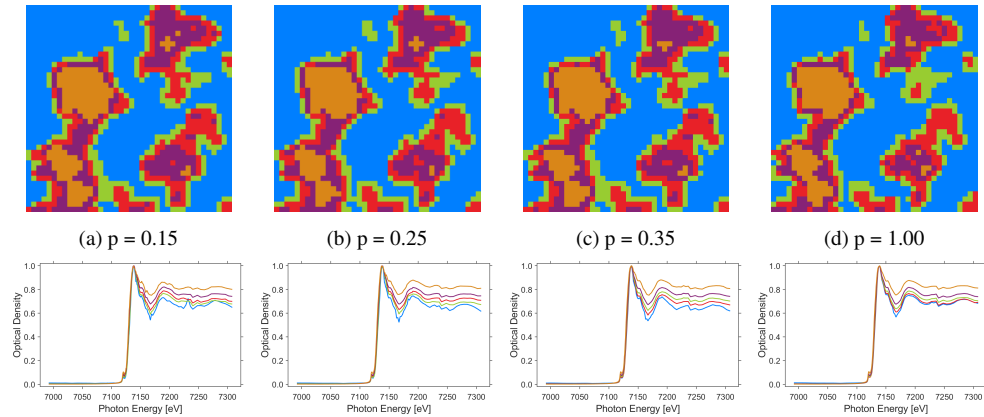


Fig. 8. Cluster and spectral results of sparse scanning for DS4. Measurements were taken at 15%, 25%, 35%, and 100% respectively. Cluster results computed using Mantis-xray [36]

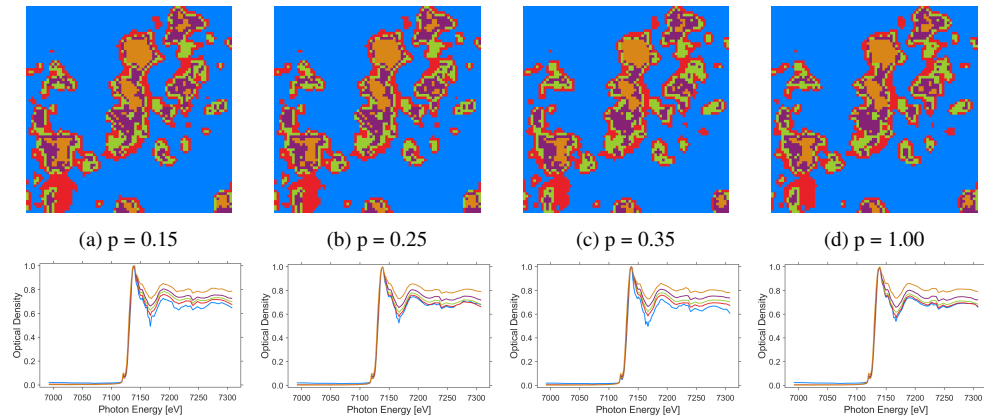


Fig. 9. Cluster and spectral results of sparse scanning for DS5. Measurements were taken at 15%, 25%, 35%, and 100% respectively. Cluster results computed by Mantis-xray [36].

464 After extracting the timestamps from the files' meta data, we can compare the run time for
 465 each sparse experiment, which have been summarised in Table 3.

466 The true time efficiency will never exactly match the undersampling ratio used due to 'dead
 467 time' as the scanner resets and for critical machine processes. Despite this, we see huge gains for
 468 the lowest undersample ratios, and a significant increase in experimental efficiency. In particular,
 469 we see greater improvements for the larger region of DS5, due to the higher measurement to
 470 dead time ratio. Another benefit for scanning larger regions is that the completion algorithms
 471 become more efficient. Because it is approximately low rank, the total number of entries grows
 472 much faster than the number of degrees of freedom - i.e. larger datasets are recoverable at lower
 473 undersampling ratios. Thus, it is in the interest of researchers using this approach to scan larger
 474 areas to improve both the scanning efficiency, and the quality of reconstruction.

475 7. Conclusion

476 We have demonstrated a new undersampling approach for spectromicroscopy data acquisition. By
477 taking advantage of the inherent low rank structure of the datasets, we have produced algorithms
478 that can accurately recover unknown entries from as little as 15% of the measurements when
479 raster sampling. We have illustrated the robustness of these algorithms to machine artefacts, and
480 how to derive parameters like the completion rank from the sparse data itself.

481 Finally, we showed the minimal impact reconstructing sparse data has on the cluster analysis
482 that is currently used to interpret the spectromicroscopy data. By implementing these methods, we
483 can conduct experiments at a much faster rate, over larger areas, and with lower dose on the sample:
484 the method consistently produces near-identical results using 20% of the measurements, with
485 potential improvements to 15% for larger samples. Alternate sampling schemes, compatible with
486 continuous motion scans, may be able to produce further reductions. These improvements will
487 help in the development of in-situ spectro-microscopy measurements and future developments
488 will reduce times in areas such as XANES nano-tomography and nano-EXAFS, which are
489 currently limited in their application by long acquisition times.

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506 **Supplemental document.** See Supplement 1 for supporting content.

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