

The effect of wind and plume height reconstruction methods on the accuracy of simple plume models- a second look at the 2010 Eyjafjallajökull eruption

Tobias Dürig¹, Magnús T. Gudmundsson¹, Thorbjörg Ágústsdóttir², Thórdís Högnadóttir¹, Louise S. Schmidt³

1. Institute of Earth Sciences, University of Iceland, Sturlugata 7, 101 Reykjavík, Iceland

2. Iceland GeoSurvey (ÍSOR)

3. Department of Geosciences, University of Oslo, Sem Saelands Vei 1, 0371 Oslo, Norway

* Corresponding author, tobi@hi.is

Abstract

Real-time monitoring of volcanic ash plumes with the aim to estimate the mass eruption rate is crucial for predicting atmospheric ash concentration. Mass eruption rates are usually assessed by 0D and 1D plume models, which are fast and require only a few observational input parameters, often only the plume height. A model's output, however, depends also on the plume height data handling strategy (sampling rate, gap reconstruction methods and statistical treatment), especially in long-term eruptions with incomplete plume height records. To represent such an eruption, we used Eyjafjallajökull 2010 to test the sensitivity of six simple and two explicitly wind-affected plume models against 22 data handling strategies. Based on photogrammetric measurements, the wind deflection of the plume was determined and used to re-calibrate radar-derived height data. The resulting data was then subjected to different data handling strategies, before being used as input for the plume models. The model results were compared to the erupted mass measured on the ground, allowing us to assess the prediction accuracy of each combination of data handling strategy and model. Combinations that provide highest prediction accuracies vary, depending on data coverage, eruption intensity, and fragmentation mechanism. However, for this type of moderate-to-weak eruption (VEI 3 in terms of maximum intensity), the most important factor was found to be the prevailing wind speed. When wind speeds exceed 20 m/s, most combinations of

33 strategies and models provide predictions that underestimate the erupted mass by more than
34 40%. Under such conditions, the optimal choice of data handling strategy and plume model
35 is of particular importance.

36

37 Introduction

38 Volcanic ash injected into the atmosphere during explosive eruptions can pose a serious threat for
39 aviation and air-travel infrastructure (Kienle et al. 1980; Grindle and Burcham 2002). Next to direct
40 observations of the ash cloud (e.g. satellite imagery and lidar retrievals) predictions of movement of
41 volcanic ash clouds are based on atmospheric ash dispersion models (Dacre et al. 2011; Kristiansen et
42 al. 2012; Dioguardi et al. 2016, 2020). Inaccurate predictions can on one hand lead to severe damage
43 to and even loss of aircraft (Guffanti et al. 2010), or the other hand to airport closures framed or
44 perceived as over-cautious (Harris et al. 2012; Harris 2015) and flight diversions or cancellations,
45 which involve large preventable costs (e.g., Brannigan 2011; Macrae 2011; Ragona et al. 2011). This
46 underlines the need for accurate, reliable and confident model forecasts, the lack of which is the
47 current bottle neck in our forecasting capability. Key to achieving such forecasts is thus to be able to
48 provide an ash dispersion model with accurate near-term eruption source parameters, i.e. physical
49 quantities, which characterize the eruptive source. In this regard the most important eruption source
50 parameter is the mass eruption rate (MER) (e.g., Degruyter and Bonadonna 2012; Mastin 2014;
51 Bonadonna et al. 2016), that is the mass flux of the eruptive gas-ash mixture, expressed in kg/s
52 (Wilson and Walker 1987).

53 MER cannot be directly measured. However, methods have been developed that infer MER using
54 observable properties of the plume. These include using video analyses of ash plumes and ejecta
55 (e.g., Wilson and Self 1980; Valade et al. 2014; Dürig et al. 2015b, a; Pioli and Harris 2019; Tournigand
56 et al. 2019), emitted infra-sound waves (Johnson and Ripepe 2011; Ripepe et al. 2013), thermal
57 infrared signatures (e.g., Harris 2013; Ripepe et al. 2013; Harris et al. 2013; Cerminara et al. 2015),
58 electrostatic field (Büttner et al. 2000; Calvari et al. 2012), interpretation of microwave radar signals
59 (Montopoli 2016; Marzano et al. 2020) or satellite-based estimates (e.g., Pouget et al. 2013;
60 Pavolonis et al. 2018; Gouhier et al. 2019; Bear-Crozier et al. 2020). For real-time MER assessment,
61 however many of these approaches are still in an experimental state or struggle with high
62 uncertainties as they often depend on data that are hard to obtain *in situ*, e.g. the vent geometry
63 (Dürig et al. 2015a). These methods will thus not be considered here.

64 The most straightforward and reliable observable in an explosive eruption is usually the height of the
65 eruptive column H . A number of plume models linking H with the mass eruption rate at the vent have

66 been developed (for in-depth overview see Costa et al. 2016). These range from “simple” theoretical
67 (Wilson and Walker 1987; Woods 1988) and empirical OD models (Sparks et al. 1997; Mastin et al.
68 2009), through explicitly wind-affected steady 1D models (Bursik 2001; Degruyter and Bonadonna
69 2012; Devenish 2013; Woodhouse et al. 2013; Mastin 2014; de’Michieli Vitturi et al. 2015; Folch et al.
70 2016; Aubry et al. 2017) to elaborate time-dependent multi-phase models in 2D (Neri et al. 1998) or
71 3D space (Esposti Ongaro et al. 2007; Suzuki and Koyaguchi 2012; Cerminara et al. 2016).

72 At present, real-time MER assessments must rely on simple OD and explicitly wind-affected 1D
73 models, which have the advantage of providing fast results (e.g., Sparks et al. 1997; Bursik 2001;
74 Woodhouse et al. 2013). A monitoring system that uses a suite of such models to assess the mass
75 eruption rate in near-real time is the software REFIR (Real-time Eruption source parameters
76 FutureVolc Information and Reconnaissance system) (Dürig et al. 2018; Dioguardi et al. 2020). Simple
77 OD models are limited by initial assumptions, accuracy of measured parameters (plume height and
78 mass) and simplifications on which they are based. To assess the effect of these shortcomings, this
79 study explores three aspects that might limit the models’ accuracies:

- 80 1. the effect of plume height uncertainties resulting from a stepwise horizontal ground-
81 based radar scanning technique on MER estimates.
- 82 2. the sensitivity of data handling strategies on model predictions when dealing with
83 incomplete sets of plume height recordings.
- 84 3. the impact of wind on MER prediction by non-explicitly wind-affected plume models
85 in comparison to the impact of using different data handling strategies.

86 The term “data handling strategies” refers to the statistical treatment of plume height data before it
87 is used as a model input. Note that with “plume height” we refer to the maximum elevation of the
88 plume above vent, unless stated otherwise. Since plume height H and mass eruption rate Q are
89 related in a highly non-linear way, the time-averaged mass eruption rate $\overline{Q(t)}$ is in most cases *not*
90 identical to Q resulting from using the time-averaged plume height $\overline{H(t)}$ as input. Thus:

$$91 \quad \overline{Q(t)} \neq Q(\overline{H(t)}) \quad (1)$$

92 Although rarely acknowledged, this fact implies that not only the plume height H itself, but also the
93 time interval over which it is measured, as well as the statistical strategy for its assessment, are
94 expected to affect a model’s outcome.

95 In our model sensitivity study, we focus on the 2010 Eyjafjallajökull (Iceland) eruption. This event
96 featured a wind-affected plume (Gudmundsson et al. 2012), which was monitored by a stepwise
97 horizontally scanning radar. The resulting plume height records were incomplete (Arason et al. 2011).

98 The 2010 Eyjafjallajökull eruption therefore represents an ideal test case for the exploration of the
99 three aspects mentioned above.

100 The Eyjafjallajökull 2010 eruption

101 The 2010 eruption of Eyjafjallajökull had four phases of activity.

- 102 • Phase I: lasting from 14 April 2010 until noon of 18 April, was an initial explosive phase with
103 phreatomagmatic activity (Gudmundsson et al. 2012; Dellino et al. 2012).
- 104 • Phase II: (second half of 18 April - 4 May) was a phase of low discharge that was effusive with
105 relatively weak but sustained explosive activity (Gudmundsson et al. 2012).
- 106 • Phase III: (5 - 17 May) saw a significant increase of explosive activity, coinciding with a
107 change in melt composition from benmoreite to trachyte (Gudmundsson et al. 2012). This
108 phase was characterized by the emission of distinct explosive ash pulses (Dürig et al. 2015b,
109 a) and is often referred to as “second explosive phase” (Gudmundsson et al. 2012; Dellino et
110 al. 2012).
- 111 • Phase IV: (18 - 22 May) was characterized by a decline in eruption activity and plume height.
112 Although the end of eruption is set on 22 May, minor isolated explosive events occurred on
113 4-8 June 2010 (Gudmundsson et al. 2012). No ash plumes were detected on radar after
114 10:20 UTC on 21 May, when the plume fell below the detection limit of 2.5 km (Arason et al.
115 2011).

116 The eruption’s fallout was measured in considerable detail at about 400 locations in Iceland
117 (Gudmundsson et al. 2012), allowing us to compare model predictions on the erupted mass with the
118 actual amount of mass deposited as tephra. For the part of the tephra that fell into the ocean south
119 and southeast of Iceland a considerable uncertainty exists, but the magnitude of this component can
120 nevertheless be estimated using well-established thinning behaviour of tephra layers with distance
121 (e.g., Thorarinsson 1954; Pyle 1989; Fierstein and Nathenson 1992). On 14-16 April, during phase I,
122 westerly upper-tropospheric winds prevailed, changing to northerly winds by 17 April, directing the
123 plume to the south (Gudmundsson et al. 2012). As a result, the erupted mass could be assessed
124 separately as phases *Ia* and *Ib* (see Table 1).

125 The monitoring conditions of this eruption are classified as relatively challenging since it is a
126 moderate-to-weak eruption that took place under adverse weather conditions (frequent clouds at
127 low and medium altitudes and strong wind) (Arason et al. 2011; Gudmundsson et al. 2012; Björnsson
128 et al. 2013), which meant direct observations of the volcanic plume were often not possible.

129 Data

130 The data used in this study are:

- 131 i) Plume height obtained by the Icelandic Meteorological Office (IMO) from the C-band
132 radar in Keflavík (Arason et al. 2011).
- 133 ii) Plume height as measured from photographs taken from inspection aircraft (Figure 1).
- 134 iii) Plume height as measured on photographs taken from the ground from areas to the
135 west of the volcano.
- 136 iv) Wind speed and direction at plume-relevant altitudes retrieved either from ERA5 or
137 ICRA. ERA5 is a global reanalysis product by the European Centre for Medium-Range
138 Weather Forecasts with a horizontal resolution of 30 km (Hersbach et al. 2020). ICRA is a
139 local reanalysis product by the Icelandic Met Office over Iceland, using the non-
140 hydrostatic numeric weather prediction model HARMONIE-AROME with a horizontal
141 resolution of 2.5 km (Nawri et al. 2017; Schmidt et al. 2018). Both models have a
142 temporal resolution of 1h. Since ICRA data were only available up to an altitude of 6 km
143 a.s.l., we used composite data sets for the ICRA labelled wind speed profiles: ICRA
144 reanalysis data below 6 km, and ERA5 reanalysis data above.

145 Radar data

146 Located at 155 km distance from the vent, the C-band radar at Keflavík airport, was the only weather
147 radar operating in Iceland at the time of the eruption. The radar's sampling strategy was to scan at
148 vertical angles of 0.5°, 0.9°, 1.3°, 2.4°, 3.5°, 4.5° and 6° (Arason et al. 2011). The width of the beam
149 was 0.9° providing some overlap between scans. For Eyjafjallajökull, the vertical distance between
150 the beams of the lowest scanning angles was about 1.1 km and about 3 km between 1.3° and 2.4°.
151 For a target above the volcano, this leads to a stepping in the plume height records, with preferential
152 sampling at about 2.5, 3.9, 5.0 and ~8 km height above sea level (for details see Arason et al., 2011).

153 Photos

154 The monitoring of vent activity and visual observations of plume behaviour took place on inspection
155 flights (see Fig. 1; for complete list of surveillance flights, see Appendix. 6.3 in Thorkelsson et al.
156 (2012)). For flights where the aircraft location was recorded with GPS, it was possible to determine
157 the position of the camera at the time of a photo with an accuracy higher than 1 km (Gudmundsson
158 et al. 2015). This is the maximum error in the distance between aircraft and the volcanic vents, or any
159 other reasonably well-defined target.

160 In most cases the plume was bent over by the wind, and its bearing was usually obtained from the
161 position of the aircraft when it was upwind and in line with the plume. Occasionally this was also
162 done by flying along the plume and determining the location relative to known landmarks. Errors in
163 distance to the top of the plume were in this way contained within 1 km in most cases, and less than
164 2 km in all cases. In the vast majority of cases this amounted to 5-10% of the distance.

165 The height of plume was either assessed directly by using the altimeter of the plane (Fig. 1d) or
166 calculated using the Python software Pixelcalc (Magnússon 2012). The plume top was referenced by
167 triangulation, using photos of the plume taken in short time intervals from different angles. Given a
168 known camera-target distance Pixelcalc converts distances between two points on the photo to
169 actual 3D distances at the location of the target by using the lens and camera specifications, which
170 are derived from the photo's metadata. The software corrects for Earth's curvature at distances
171 where it becomes significant. In most cases the plume height is obtained by measuring the distance
172 on the photo between the plume top and the surface, taking into account plume deflection by wind
173 and surface elevation under the plume top, to obtain height above sea level. In some cases, where
174 data on cloud top height are available, height over clouds (i.e., the distance between cloud top and
175 plume top) is determined. Such cloud heights are in some cases obtained by direct observation using
176 aircraft altimeter or through independent cloud height measurement using Pixelcalc. The
177 uncertainties of the heights and horizontal deflection distances vary from one photo to another,
178 depending on the uncertainties in distances, pixel size and image geometry. Assuming maximum
179 errors for these parameters, we found the uncertainties of plume height and deflection to be less
180 than 0.5 km.

181 In addition to the airborne imagery, photographs taken on the ground were used. For these the
182 location of the camera is accurately known which allowed us to calculate the distance to the plume
183 top. Independent information of plume azimuth was used to correct for plume top deflection relative
184 to vents.

185 The total number of photos taken on inspection flights are counted in thousands. Not all flights
186 provided photos that allow reliable determination of plume height, e.g., due to cloud cover. Only the
187 photos best suited for analysis were used, yielding in total 165 observations of plume heights,
188 covering 17 of the 38 eruption days examined. From these photos, 148 were taken in the air, and 17
189 on the ground. In 85 cases these photos could also be used to assess the plume deflection (i.e., the
190 horizontal shift of the plume top towards the vent location).

191 Methods

192 Models

193 For estimating the MER, we used six *simple* (non-explicitly wind-affected) plume models and two
194 wind-affected models, which are for simplicity referred to by the name of their first authors:

- 195 • *Wilson* (Wilson and Walker, 1987) - a theoretical model based on the theory of buoyant
196 plumes by Morton et al. (1956) which estimates the mass eruption rate Q by:

$$197 \quad Q_{Wilson} = (H/c)^4 \quad (2)$$

198 where H denotes the plume height (in m) and c is a constant which is $236 \text{ m (s/kg)}^{1/4}$.

199 To the group of non-explicitly wind-affected numerical models tested in this study we include three
200 derivatives of the Wilson model, which use Eq. (2), but with different constants c based on empirical
201 data from plumes with basaltic or andesitic tephra. They are referred to as:

- 202 • *Wehrmann* (Wehrmann et al. 2006), which uses for c a value of $295 \text{ m(s/kg)}^{1/4}$.
- 203 • *Scollo* (Scollo et al. 2007), using $247 \text{ m(s/kg)}^{1/4}$ for c .
- 204 • *Andronico* (Andronico et al. 2008), for which $244 \text{ m(s/kg)}^{1/4}$ is selected for c .

205 Two empirical models were used, which are based on correlations between plume height during
206 eruption and the quantity of tephra found in deposits:

- 207 • *Sparks* (Sparks et al., 1997) - an empirical model which approximates Q by:

$$208 \quad Q_{Sparks} = \rho \cdot (H/c)^{3.86} \quad (3)$$

209 where ρ is the dense-rock equivalent (DRE) density of the tephra within the plume. Following
210 Gudmundsson et al. (2012), a DRE density of 2600 kg/m^3 was used for all computations and the
211 constant c is given by Sparks et al. (1997) as $1670 \text{ m(s/kg)}^{1/3.86}$.

- 212 • *Mastin* (Mastin et al., 2009) – an empirical model with ρ defined as in (3) which estimates the
213 mass eruption rate by:

$$214 \quad Q_{Mastin} = \rho \cdot (H/c)^{4.15} \quad (4)$$

215 where c is given by Mastin et al. (2009) as $2000 \text{ m(s/kg)}^{1/4.15}$.

216 Since these models are empirically constrained they reflect the influence of the wind on plume-
217 height to some degree (Mastin 2014) and can therefore be seen as implicitly wind-affected.

218 To juxtapose these simple plume models with explicitly wind-affected models we included two
219 examples of the latter category:

- *Woodhouse* (Woodhouse et al. 2013) – an empirical relationship of a numerical 1D model, which estimates the MER by:

$$Q_{Woodhouse} = \left(\frac{h}{0.318} \cdot \frac{1 + 4.266\overline{W}_s + 0.3527\overline{W}_s^2}{1 + 1.373\overline{W}_s} \right)^{3.953} \quad (5)$$

220 where h denotes the centreline plume height (in km) and \overline{W}_s quantifies the strength of the wind
221 shear from the ground to a reference height H_1 , according to:

$$\overline{W}_s = 1.44 \frac{V_1}{\overline{N}H_1} \quad (6)$$

222 with V_1 being the wind speed at H_1 and \bar{N} being the average buoyancy frequency. The latter
 223 parameter is determined by:

$$\bar{N}^2 = \frac{1}{h} \int_0^h N^2(z) dz = \frac{1}{h} \frac{g}{C_{a0} T_{a0}} \int_0^h \left(1 + \frac{C_{a0}}{g} \frac{dT_a}{dz} \right) dz \quad (7)$$

224 with g , C , and T being the gravitational acceleration, heat capacity and temperature. Parameter z
 225 denotes the vertical coordinate above the source, subscript a refers to the atmosphere, and 0 to the
 226 volcanic source vent height. Note that a plume's top height H coincides with the elevation of its
 227 centreline h only for strong eruptions with vertical ash columns (Mastin 2014), which introduces an
 228 error when using H as an input for a weak or moderate eruption. Since for Eyjafjallajökull 2010 the
 229 exact conversion from H to h is unknown, for simplicity, in this study we assumed that the centreline
 230 height h was always 10% lower than H .

- 231 • *Aubry* (Aubry et al. 2017): A scaling relationship which uses multiple linear regression to
 232 relate the logarithm of the MER to the logarithm of the plume height H (in km), the logarithm
 233 of the average buoyancy frequency \bar{N} (eq. (7)), and a windscaling parameter \mathcal{W}_0^* :

$$\log(Q_{Aubry}) = C + C_H \log(H) + C_N \log(\bar{N}) + C_W \mathcal{W}_0^* \quad (8)$$

234 where the constants $C = 7.4$, $C_H = 3.4$, $C_N = 2.3$ and $C_W = 1.1$ are empirically determined. The
 235 wind scaling parameter \mathcal{W}_0^* is described by:

$$\mathcal{W}_0^* = \frac{W}{U_0} = \frac{W}{1.85 \sqrt{R n_0 T_0}} \quad (9)$$

236 where W is the horizontal wind speed averaged over the plume height, U_0 is the eruption exit
 237 velocity, $R=461.5$ J/K/kg is the gas constant of volcanic gas at vent, n_0 is the magma volatile content,
 238 and T_0 is the exit temperature. Here we assume that $n_0 = 3$ wt% (Woodhouse et al. 2013) and
 239 $T_0 = 1323$ K for Eyjafjallajökull (Keiding and Sigmarsson 2012).

240 For the goals of this study, we put our focus on the simple plume models. These have the advantage
 241 that they use plume height as the only independent variable. In their entirety, the six simple models
 242 represent the empirical variation of the predictions by the theoretical *Wilson* model. In contrast to
 243 simple models, the more sophisticated explicitly wind-affected models depend on additional
 244 variables and parameters, (e.g., atmospheric parameters, wind entrainment coefficient, volatile
 245 content, magmatic temperature, top vs centreline height) which are further sources of uncertainties
 246 that have to be taken into account when analysing the model results.

247 Plume height data handling strategies

248 The radar data set is based on sampling intervals of 5 minutes, but the data set is incomplete (Arason
249 et al. 2011). Reasons for missing data were (Arason et al. 2011):

- 250 i) The plume was below the visible horizon from Keflavík, which limited the lowest
251 observable plume to 2.5 km.
- 252 ii) The plume was obscured by precipitating clouds.
- 253 iii) The radar was malfunctioning.
- 254 iv) The plume was absent, and the data gap was “real”.

255 Here we explore 22 strategies to construct a data set with an improved coverage of the complete
256 eruption (see Table 2). In the example of Figure 2a, a data set was segmented into 3h and 6h static
257 blocks, composed of 36 and 72 bundled five-minute slots, respectively, illustrating how much of the
258 eruption was covered according to the blocks. For example, a data coverage of 50% means for a 6h
259 block that 36 of its 72 entries are missing. In this study, we distinguish between data voids and gaps.
260 *Voids* are defined as empty or missing 5-minute slots in a block that otherwise contains data. In
261 contrast, a completely empty block of data is called a *gap*. Figure 2b and c summarize the 3h- and 6h
262 gaps in the radar records of the Eyjafjallajökull 2010 eruption.

263 The tested strategies can be divided into five main categories (see Table 2). Aspects included in the
264 tested strategies are: sampling method, void reconstruction, gap reconstruction, timebase, time
265 steps and a factor introduced to correct for data reduction bias (Y-correction):

266 Static plume height reconstruction (*SH*)

267 Static plume height reconstruction (*SH*) strategies bundle data by splitting the data set into static
268 time intervals (blocks) and using plume height averages. The vast majority of studies using plume
269 heights for MER modelling apply *SH*-strategies. Studies on Eyjafjallajökull used the data by Arason et
270 al. (2011) with blocks typically ranging from 1 h (Woodhouse et al. 2013) to 6 h (Kaminski et al. 2011;
271 Bursik et al. 2012; Mastin 2014). Others report plume heights averaged over eruptive episodes that
272 can include multiple eruptive phases (Aubry et al. 2021).

273 **Sampling method:** *SH* strategies applied static sampling whereby the whole data set was simply split
274 into fixed time intervals (static blocks) of a given uniform duration (see Fig. 2f).

275 **Void reconstruction:** Voids, i.e., missing plume height values within a block, were filled with the
276 mean value of the interval (see Fig. 2e).

277 **Gap reconstruction:** When applied, the heights of the missing blocks were reconstructed. The
278 missing blocks were replaced by linearly interpolating between the previous and following block (see
279 Fig. 2f). Where only void reconstruction but no gap reconstruction was used, strategies are labelled

280 with *gap*. In contrast, static sampling strategies which applied both void and gap reconstruction are
281 termed *interpolated (int)*.

282 **Timebase:** describes the width (i.e., duration) of a block. *SH*-strategies were tested for four
283 timebases, covering the range from highest to lowest resolution (see time steps).

284 **Time steps:** describes which time interval is used for data sampling. For static sampling strategies,
285 the time step is identical to the timebase.

286 No data bundling is required when using five-minute time steps. In this case the integrated mass
287 M_{5min} over the tested eruption period of duration D (in seconds) was calculated by:

$$288 \quad M_{5min} = \sum_{j=1}^f Q_{j,model}(H_{5min,j}) \cdot \Delta t \quad (10)$$

289 with the time step $\Delta t=300s$, $f = D/\Delta t$, and $Q_{j,model}(H_{5min,j})$ being the mass eruption rate Q_j provided by a
290 specific plume model with $H_{5min,j}$ as input parameter.

291 By increasing the time steps, the data coverage is increased (see Fig. 2) at the cost of temporal
292 resolution. For example, when using static blocks with a time step of 3h, the integrated mass M_{3h} is
293 given by:

$$294 \quad M_{3h} = \sum_{j=1}^f Q_{j,model}(H_{3h,j}) \cdot \Delta t \quad (11)$$

295 where $H_{3h,j}$ is the 3 h - averaged plume height and the time step $\Delta t=300s \times 36 = 10,800s$.

296 Correspondingly, Δt is doubled when using a 6h time step, and a 6h plume height average is used as
297 model input. As time step of lowest temporal resolution, we used the complete eruptive phase for a
298 static block (see Table 1):

$$299 \quad M_{phase} = Q_{model}(H_{phase}) \cdot \Delta t \quad (12)$$

300 with Δt being the phase's duration and H_{phase} being the phase-averaged plume height. Equation (12)
301 was also applied for testing the complete eruption with $\Delta t = E$, where E is the whole recorded
302 eruption period. The corresponding diagrams are labelled as *all* (for all phases). In total, we used four
303 different time steps for *SH* strategies (*5min*, *3h*, *6h*, *phase*). For example, an *SH*-strategy with a
304 timebase of 6 h and applied gap reconstruction is denoted *SH_6h_int*.

305 **Y-correction:** To adjust for the bias induced by data reduction, four of the tested *SH* strategies
306 multiplied the resulting masses with a correction factor $1/Y$ (see below). Strategies with Y-correction
307 use the additional letter "Y" as identifier. For example, if Y-correction was applied to the strategy
308 *SH_6h_int*, it is denoted *SH_Y6h_int*.

309 Static Sampling of five-minute intervals (*S_5min_gap*)

310 Being a special case of the *SH* strategy, the *S_5min_gap* strategy uses equation (10) for the
311 computation of the total erupted mass and does not apply gap reconstruction. Since it uses the
312 highest possible temporal resolution, voids cease to exist, and every missing data is a gap, which is
313 treated as a “real measurement” by the *S_5min_gap* strategy.

314 Static MER reconstruction (*SM*)

315 When using average plume heights for reconstructing voids and gaps, we would expect to produce a
316 systematic arithmetic error in the computation of erupted mass by the equations (10)-(12), since,
317 according to equation (1), the average plume heights do not yield representative MERs. The static
318 MER reconstruction (*SM*) strategy follows therefore a different path. Instead of using average
319 heights, void and gap reconstruction is applied by reconstructing the “missing” MERs for each time
320 step.

321 **Sampling method:** *SM* strategies apply static sampling.

322 **Void reconstruction:** Voids were filled with the average MER value of the interval. The erupted mass
323 *M* of a block *B* was therefore computed by:

$$324 \quad M_{3h,B} = \left(\sum_{j=1}^{36} Q_{j,model}(H_{5min,j}) \cdot \Delta t \right) \cdot \frac{36}{(36-\lambda)} \quad (13)$$

$$325 \quad M_{6h,B} = \left(\sum_{j=1}^{72} Q_{j,model}(H_{5min,j}) \cdot \Delta t \right) \cdot \frac{72}{(72-\lambda)} \quad (14)$$

326 where λ is the number of voids in block *B* and $\Delta t = 300$ s.

327 **Gap reconstruction:** When applied, the MERs of the missing blocks were reconstructed. The missing
328 intervals were replaced by linearly interpolating between the masses of the previous and following
329 block.

330 **Timebase:** Two different timebases were tested (3 h and 6 h) with the *SM* strategies. In addition, we
331 tested a strategy with a five minute timebase and MER-reconstructed gap (*SM_5min_int*).

332 **Time steps:** see timebase.

333

334 Dynamic plume height and MER reconstruction (*DHM*)

335 The *DHM* strategies combine plume height and MER reconstruction, and apply dynamic sampling,
336 which is particularly useful for real-time monitoring purposes (Dürig et al. 2018; Dioguardi et al.
337 2020).

338 **Sampling method:** Dynamic sampling, using blocks of a uniform duration with moving start and end
339 times. The width of the block is defined by the timebase, while the temporal increment with which

340 the block is moved is described by the time step. For example, a dynamic sampling strategy using a
341 3h timebase means that at any given time the data of the last 3 hours are considered.

342 **Void reconstruction:** Voids were filled with the interval-averaged plume heights, as described for *SH*
343 strategies.

344 **Gap reconstruction:** The missing blocks were replaced by linearly interpolating between the masses
345 of the previous and following block, as described for *SM* strategies. For the presented DHM
346 strategies, gap reconstruction was always applied.

347 **Timebase:** Two different timebases were tested (3 h and 6 h) with the *DHM* strategies.

348 **Time steps:** All *DHM* strategies used five-minute time steps.

349

350 The REFIR strategies (*REFIR*)

351 The strategies discussed so far produce best MER estimates independently from plume height
352 uncertainties. To consider these, the *DHM* strategies were modified by computing the MER for a
353 dynamic block *B* with:

354
$$Q_{model,B} = (Q_{model}(H_B - s_B) + Q_{model}(H_B) + Q_{model}(H_B + s_B)) / 3 \quad (15)$$

355 where H_B is the average height and s_B the plume height uncertainty of the block at time t .

356 These strategies were computed with the software REFIR (Dürig et al. 2018; Dioguardi et al. 2020),
357 which uses equation (15) to obtain a best MER estimate. We note that while the software is designed
358 to apply a combination of models linked by weight factors, here we applied it exclusively to individual
359 models.

360 **Sampling method:** Dynamic sampling, see DHM.

361 **Void reconstruction:** see DHM.

362 **Gap reconstruction:** see DHM.

363 **Timebase:** Four different timebases were tested (15 min, 30 min, 1 h and 3 h) with *REFIR* strategies.

364 **Time steps:** All *REFIR* strategies used five-minute time steps.

365

366 Erupted tephra mass

367 The erupted mass M was used as validation parameter to test the capability of the assessed models
368 and data handling strategies in reproducing mass eruption rates. The individual model-derived
369 results for the four main stages of the eruption were compared to the total masses obtained for each

370 of the analyzed phases (Ia, Ib, II, III, IV and all, see Table 1), based on the mapped fallout and
 371 exponential dispersal models for the offshore part (Gudmundsson et al. 2012). We denote the
 372 measured range of uncertainty ME .

373 Computation of Y-ratios

374 To study the impact of data reduction on prediction quality, a subset of 49 selected six-hour blocks
 375 was used, for which the radar data provides high data coverage (containing at least 60 of 72 data
 376 points measured at five-minute intervals). The voids were filled with 6h averages, as described for
 377 the SH strategies. For each of the 6h blocks with the best data coverage the predicted mass erupted
 378 M_i was computed, based on:

379 I. Five minute radar height data H_{5min} . The integrated mass $M_{i_{5min}}$ was calculated by using eq.
 380 (8) with $f = 72$.

381 II. Mean, median and maximum heights (H_{3h_avg} , H_{3h_med} , H_{3h_max}) within two 3h intervals. The
 382 integrated mass $M_{i_{3h}}$ over a 6h interval was quantified by:

$$383 M_{i_{3h}} = \sum_{j=1}^2 Q_{j,model}(H_{3h,j}) \cdot \Delta t \quad (16)$$

384 with $\Delta t=10,800s$ ($= 3h$) and $H_{3h,j}$ being H_{3h_avg} , H_{3h_med} , H_{3h_max} .

385 III. Mean, median and maximum heights (H_{6h_avg} , H_{6h_med} , H_{6h_max}) within the 6h interval, using

$$386 M_{i_{6h}} = Q_{model}(H_{6h}) \cdot \Delta t \quad (17)$$

387 with $\Delta t=21600s$ ($= 6h$) and H_{6h} being H_{6h_avg} , H_{6h_med} or H_{6h_max} , depending on the tested procedure.

388 Varying plume model, time interval (3h, 6h) and statistical treatment (mean, median, maximum)
 389 results in 24 estimates for M_i for each of the 49 analyzed 6h blocks. These mass estimates were
 390 compared to the respective value $M_{i_{5min}}$ based on the 5 minutes data sets, by generating the ratios
 391 Y :

$$392 Y_i = M_i / M_{i_{5min}} \quad (18)$$

393 Strategies which apply Y-correction (see Table 2) multiply the estimated mass M with the according
 394 time interval-specific correction factor $1/Y$.

395 Plume height uncertainties

396 We take two types of plume height uncertainties into account for the MER calculations: i) the
 397 observation uncertainty of the corrected radar plume heights, and ii) the uncertainties introduced by
 398 using the average plume height of a block to calculate the MER.

399 The uncertainty of the radar, ΔH , is assumed to be 0.5 km based on the maximum uncertainty of the
 400 photo heights and the uncertainty estimates given in Aubry et al. (2021). The uncertainty arising from
 401 the block averaging of plume heights is assumed to be equal to two standard deviations σ within

402 each block, corresponding to a 95% significance level. The total plume height uncertainty of a block
403 s_B is then calculated as:

$$404 \quad s_B = \sqrt{\Delta H^2 + (2 \cdot \sigma(H_{block}))^2} \quad (19)$$

405

406 For strategies where no height averaging is used, σ is equal to zero, and therefore only radar
407 uncertainties are considered.

408 Erupted mass uncertainty

409 The results obtained depend critically on the error in erupted mass. The isopach map reported in
410 Gudmundsson et al. (2012) is based on about 400 data points obtained during and immediately after
411 the eruption. Post-depositional erosion and compaction that often limit the accuracy of tephra layer
412 thickness and size (Biass et al. 2013; Engwell et al. 2013) are therefore considered to be minor for the
413 Eyjafjallajökull data. Good spatial coverage in the near field where thicknesses are several meters on
414 a relatively regular glacier surface result in relative errors in thickness of only a few percent. In the
415 far field the relative errors rise, reaching up to 40% for thicknesses <0.2 cm. The resulting error
416 reported for volume on land is 15% (Gudmundsson et al., 2012). For the offshore part the error was
417 estimated as 40%, based on minimum and maximum extrapolations that are consistent with
418 observations in the Faroe Islands and mainland Europe (for more detail see methods in
419 Gudmundsson et al. 2012). The resulting best estimate of maximum total error was 25%. The
420 contemporary measurements of density of the tephra yielded $1400 \pm 40 \text{ kg m}^{-3}$ (Gudmundsson et al.,
421 2012) a relative error of 3%. Treating these two errors as independent, the effect of the error in
422 density is negligible, resulting in a relative error in mass of 25%.

423 Results

424 Plume heights and ERA5 wind speeds

425 In Figure 3, the radar-derived plume heights are compared with plume heights obtained by photos.
426 The comparison of the 165 data pairs suggests that on average, photo-based plume heights show an
427 offset from the radar data of 0.5 km (see Figures 3b and 3c). As a first step for data quality
428 improvement, this systematic shift was therefore corrected by adding 0.5 km to all radar-derived
429 plume heights (see Figure 3d).

430 Figure 4a shows the wind speed profiles over the vent during the Eyjafjallajökull 2010 eruption,
431 based on the ERA5 reanalysis data. The highest wind speeds, with up to 64 m/s, occurred during
432 phase I in regions above 6000 m at a time when the recorded plume heights reached similar altitudes

433 (for more detailed, phase-specific plots, see Online Resource 1). In Figure 4b, we show the
434 differences in wind speeds between ICRA and ERA5.

435 Evaluation of plume height data handling strategies and models

436 Mean values and standard deviations of the resulting Y -ratios based on the 49 selected 6h blocks are
437 presented in Table 3. Three statistical treatments were tested for model input, using average,
438 median and maximum plume height, respectively. The larger the deviation of Y from 1, the less
439 reliable the tested method of approximation (assuming $M_{i,5min}$ to be the true value). The Y -ratios
440 closest to 1 are found using median plume heights. However, using the median also results in
441 significantly larger standard deviations compared to average plume heights. When using 3h- or 6h-
442 blocks, all models must be corrected with $1/Y$ to be comparable to the summation of masses with
443 5min time step. For Y -corrections we focused on averaged heights, for which Y is almost constant for
444 all models. Y ranges between 0.89 and 0.91 for 6h-blocks and 0.91-0.92 for 3h-blocks. Therefore, Y -
445 correction strategies for 6h-blocks (denoted $Y6h$, see Table 2) were tested by using average heights
446 and a model-independent correction factor of $1/0.9 = 1.1$. Accordingly, a correction factor of $1/0.915$
447 $= 1.09$ was used for $Y3h$ strategies.

448 Figure 5 demonstrates how the temporal evolution of the erupted mass varies, depending on the
449 applied plume height strategy and plume model. In both cases, the whole eruption was used as the
450 reference period, with the measured range of the erupted mass (reported in Gudmundsson et al.
451 2012) indicated by the red error bar. The curves of Figure 5a show the modelled erupted mass using
452 strategy SH_3h_gap . This strategy uses 3h-averaged plume heights in combination with applied void-
453 reconstruction but without gap reconstruction (see Table 2). While for this strategy both wind-
454 affected models provide estimates that lie within the measured range, this applies only for one out
455 of the six simple plume models tested (*Wilson*). Figures 5b-d show the outcome of *Wilson* when
456 applying different strategies. Comparison of model-estimates to the measured mass for individual
457 eruptive phases is illustrated in Figure 6. When using strategy SH_3h_gap for phase I, most of the
458 models underestimate the erupted mass, while for phases II to IV *Wilson*, *Andronico*, *Scollo* and
459 *Sparks* provide estimates that coincide with the measured range. The same applies to *Woodhouse* for
460 the phases I, II and III, while this wind-affected model provides too large estimates for phase IV. The
461 other wind-affected model tested with SH_3h_gap , *Aubry*, provides estimated masses that are above
462 the measured range suggested by Gudmundsson et al. (2012), but provides estimates that are on the
463 lower margin for phase I. This combination results in a coinciding estimate, when considering the
464 whole eruptive period (see Fig. 5a). Systematically on the lower end of the estimated range of
465 modelled masses are *Wehrmann* and *Mastin*, suggesting, that SH_3h_gap is not the optimal choice
466 for these models. This overview is an example for just one of the 22 tested strategies with reference

467 to the whole eruption. The phase-specific results for *Wilson* model outcomes according to the
468 strategies in Fig. 5b can be found in Online Resource 2.

469 Figure 7 compares the measured masses of all reference periods with the estimates of four models
470 using six different strategies with a timebase of 3 hours. The ranges of measured mass for each
471 eruption phase are marked in green. For example, for phase III the *Mastin* model (see Fig. 7b)
472 achieves best fitting estimates when applying the *REFIR_3h* strategy. Figure 7 demonstrates the
473 impact that the choice of a strategy has on the modelled MER prediction. Phase II is the phase with
474 the most abundant gaps in the radar records (see Fig. 2). Thus, it is not surprising that this is also the
475 phase where the selected strategy has the biggest influence on the predicted MER. For example, for
476 *Mastin* the prediction with sampling strategy *REFIR_3h* is 246 % larger than with *SH_3h_gap*.

477 As a general finding, applying sampling strategies with gap reconstruction (*int*) always results in
478 larger masses than applying no gap reconstruction methods (*gap*). The *SM_3h_int* strategies provide
479 systematically larger results than *DHM_3h* or *SH_3h_int* strategies, due to the non-linear relation
480 between H and Q. For the same reason, the mass eruption rates $Q_{model}(H_B + s_B)$ and $Q_{model}(H_B - s_B)$ in
481 equation (15) are not symmetrically distributed towards $Q_{model}(H_B)$. Consequently, the *REFIR_3h*
482 strategy provides larger values as best MER estimate compared to other strategies.

483 When using the whole eruption as the reference period, the effect of the selected timebase for the
484 *REFIR* strategy on the model outcome (i.e., the erupted mass estimates) is at maximum 6.8% and
485 therefore small in comparison to the modelled mass uncertainties (see Online Resource 3). The
486 largest effect of the selected timebase on the outcome is found for phase IV, where the model
487 results vary with 12.0% at maximum (see Online Resource 4). Presented results for *REFIR* methods in
488 the following are limited to the two endmembers (*REFIR_3h* and *REFIR_15min*).

489 For the static sampling strategies *SH* and *SM*, the timebase can have some considerable impact on
490 the model outcome, especially for phases with abundant gaps. For example, for phase II, the mass
491 predictions increase by 39 to 42% (depending on the model), if strategy *SH_6h_gap* is applied instead
492 of *SH_3h_gap*.

493 [Discongruence indices and congruency maps](#)

494 A result table containing the modelled masses for all phases, data handling strategies and models can
495 be found in Online Resource 5. In order to systematically assess how well the model-estimates
496 coincide with the measured data, we defined discongruence indices (*DIs*) in Table 4, based on the
497 error *ME* of measured mass data (Table 1). A model prediction that lies within the range of the
498 measured mass is characterized by a discongruence index of 0. The more the model output deviates
499 from the ground truth, the larger the *DI*. To better visualize the results, we produced congruency

500 maps, where each index is assigned a colour. Examples for such congruency maps are shown in
501 Figure 8. There, the dark blue fields mark methods and strategies that lead to estimates that coincide
502 with the range of mass from ground-based measurements. Combinations of strategies and plume
503 models, which provide estimates that are just outside of the range of the measured values are
504 marked in light blue. Predictions that considerably over- or underestimate the erupted mass (*DI* of 5)
505 are displayed in orange. The resulting congruency map for the complete phase is presented as Online
506 Resource 6, and all phase-specific results are provided in Online Resource 7.

507 Based on the resulting *DIs*, it was possible to construct tables with their averaged values that allow us
508 to analyse the prediction qualities of models and data handling strategies for individual eruptive
509 phases, as well as for the complete eruption (Tables 5 and 6).

510 Phase-specific comparison of plume height data handling strategies

511 Table 5 lists for each phase the strategies' success, represented by discongruence indices averaged
512 over all tested plume models (separated between simple and explicitly wind-affected models). When
513 using the entire eruption (*all*) as the reference period, the overall average *DI* for all strategies is 0.7.
514 According to the results, the most successful strategy for the simple models is *REFIR_3h* with a *DI* of
515 0.3. For wind-affected models a number of strategies resulted in predicted mass values congruent to
516 the measured mass, including *SH_6h_gap*, *SH_6h_int*, *SM_3h_gap*, *DHM_6h* and *DHM_3h*.

517 According to the model-specific *DIs* averaged over all strategies (Table 6), the models by *Wilson* and
518 *Sparks* were found to have the highest prediction success rate, with *DIs* of 0.05 and 0.015,
519 respectively, when considering the entire eruption. When studying individual phases, *Andronico* and
520 *Sparks* turn out to be most successful, with an average *DI* of 0.05 for phase III.

521 Below we summarize the results for each of the phases:

522 **Phase Ia:** the overall average *DI* is 2.16, significantly larger than for any other phase. According to
523 Table 6, the most successful strategy for both simple and explicitly wind-affected models is
524 *REFIR_15min* with *DIs* of 1.7 and 0.5, respectively.

525 **Phase Ib:** here strategies provide higher success rates. Optimal strategies for simple models are the
526 *REFIR* strategies, as well as *SM_6h_gap* and *SM_6h_int*, featuring *DIs* of 0.5. For wind-affected
527 models, however, other strategies (e.g., *SH_6h_int* or *DHM_6h*) provide better results, while *REFIR*
528 strategies result in overestimates of predicted masses.

529 **Phase II:** due to relatively weak activity, phase II has the lowest data coverage per block (see Fig. 2d).
530 Many strategies and models tend to overestimate the erupted mass of this phase. This applies
531 especially to strategies where gaps were filled by interpolated data (*DHM*, *REFIR* and static sampling

532 methods with gap reconstruction). For simple models, the best results with a *DI* of 0.3 are found for
533 *SM_6h_gap*, *SH_6h_gap* and *SH_phase*, which are strategies with large timebases and no gap
534 reconstruction. For wind-affected models, the best results were achieved with *SM_3h_gap*,
535 *SH_3h_gap* and *SH_Y3h_gap*.

536 **Phase III:** featured increased explosive eruption at relatively low wind speeds. There, strategies
537 showed the highest success rates of all phases with an average *DI* of 0.45 (0.5 for simple models, 0.2
538 for explicitly wind-affected models). The most successful strategies for simple models were the *REFIR*
539 strategies with a *DI* of 0.3. For *Woodhouse* and *Aubry* most of the tested strategies result in
540 predictions that coincide with the measured mass. Exceptions are the *REFIR* strategies (*DI* = 1.0) , and
541 strategies that apply Y-correction (*DI* = 0.5).

542 **Phase IV:** For the weaker final eruptive phase, the optimal strategies for simple and explicitly wind-
543 affected models were found to be *REFIR_15min* (*DI* = 0.2) and *S_5min_gap* (*DI* = 0), respectively.

544 [Effect of using ERA5 vs ICRA weather data](#)

545 So far, the results presented for the explicitly wind-affected models (*Woodhouse*, *Aubry*) are based
546 on ERA5 wind speed data. How much do the results of these models (i.e., estimated erupted mass)
547 change, when the ICRA data of (considerably) higher spatial resolution are used? In most cases
548 (except for phase IV), the ICRA data leads to slightly larger mass estimates, but these changes are
549 almost insignificant, considering the uncertainties (see Online Resource 8 for *Woodhouse*). For
550 example, when applying the *REFIR 3h* strategy to *Woodhouse*, the differences between model-results
551 based on ERA5 versus ICRA wind speeds is only approximately 2% for the whole eruption (and 1.3%
552 of the uncertainty). The absence of sizeable changes can be explained by the fact that both sets of
553 weather data feature similar wind speeds over most of the eruption period (see Fig. 4b). The only
554 major difference (of more than 10 m/s) occurred during phase II, on 29 April below 2000 m a.s.l. For
555 this phase, the modelled mass differed by 4.7%, which is small considering the model uncertainties.
556 We thus conclude that for reanalysis of the Eyjafjallajökull 2010 eruption, using weather data of
557 higher spatial resolution only effects the overall outcome in a minor way. However, this result
558 should not be generalized. A reanalysis of other eruptions, especially short-termed ones, will
559 probably benefit from a higher spatial resolution of weather data.

560 [Wind effect on plume deflection and plume model predictions](#)

561 The horizontal deflection of the plume top plotted over the contemporary wind speed at plume top
562 level, based on ICRA reanalysis weather data (Figure 9a) shows an approximately linear correlation
563 between these parameters. This is also indicated by a Pearson correlation coefficient (*R*) of ~0.68. For
564 Figure 9a, wind speeds at variable elevation were used. To find out for which fixed altitude the wind
565 speeds correlate best with deflection, the according correlation coefficients were plotted, using wind

566 speeds of both ERA5 and ICRA data (Fig. 9b). Next to the raw data (*ERA5 1h* and *ICRA 1h*), wind
567 speeds averaged over 2h and 3h were also used. The best correlation at fixed elevation is observed
568 for 2h-averaged ERA5 wind speeds at 2900 m a.s.l. ($R = 0.65$). At that altitude ICRA wind speeds
569 correlate slightly less with deflection ($R = 0.63$), but below 2500 m a.s.l. it is the ICRA data that shows
570 a better correlation with wind speeds compared to ERA5.

571 From the models tested, not only *Woodhouse* and *Aubry* incorporate a wind-effect. There is an
572 implicit dependency on wind for the other (simple) models as well, as revealed by Figure 10. There,
573 the ratio of modelled vs measured mass was plotted over phase-averaged wind speeds. For the
574 modelled mass, three non-explicitly wind-affected models were investigated in combination with
575 four strategies: *Wilson* (*Wil*, as representative for the simple numerical OD models), *Sparks* (*Spa*) and
576 *Mastin* (*Mas*). Here we selected strategies that do not apply gap reconstruction, to avoid the bias
577 caused by the fact that there is a large variance in data gaps between the different phases. Two
578 general trends (indicated by trendlines) can be observed:

- 579 1) At low wind speeds the different strategies and models lead to predictions with a large span.
580 This range decreases for larger wind speeds.
- 581 2) With higher wind speeds all shown strategies and non-explicitly wind-affected models tend
582 to increasingly underestimate the actual erupted mass. While for a wind speed of 14.5 m/s
583 $M_{modelled} / M_{measured}$ averages at around 103%, this ratio drops for 16.2, 19.2 and 36.4 m/s to
584 79%, 67% and 46%, respectively. Thus, for wind speeds over 20 m/s the model predictions
585 underestimate the erupted mass on average by more than 40%.

586 Both trends are also confirmed in Figure 8, which compares the congruency maps of high-wind
587 eruptive phase Ia (Fig. 8a) with the phase of lowest wind speeds (phase II, Fig. 8b), thus contrasting
588 the two end members of Fig. 10. Compared to phase Ia, phase II was of much longer duration, and it
589 was characterized by abundant data gaps, which should result in large differences between gap-
590 reconstructing strategies and those which ignore them. Yet, for phase II, 54 combinations of
591 strategies and models provide predictions that lie within the uncertainties of the measurements *ME*
592 (tiles marked with dark blue color), as opposed to only 14 combinations found for phase Ia. While for
593 most combinations, the explicitly wind-affected models *Woodhouse* and *Aubry* appear to be robust
594 against this wind effect, for the simple models only very few combinations give predictions close to
595 the measurements.

596 Discussion

597 The comparison between radar and photographic data suggests that the former systematically
598 underestimated the plume height by 0.5 km. This finding is in contrast to the assessment of webcam

599 images analyzed by Arason et al. (2011), which suggested that the radar over-estimated the plume
600 heights. These webcams were, however, located to the north and west-northwest of the volcano,
601 while the plume was deflected away from the cameras, towards the south to southeast. Using the
602 distance to the vents instead of the plume top would therefore lead to an underestimation in
603 webcam-derived plume heights which might explain the differences with our results (see Fig. 3d). For
604 example, a plume top of 5 km height that is shifted towards the south with a deflection d of 6 km
605 would be perceived as being only 3 km high on images produced by the webcam at Þórólfsvell,
606 located 9 km to the north from the vents. The difference between photos and radar measurements is
607 mainly a consequence of the radar's horizontally stepping scanning procedure and the eruptive and
608 atmospheric circumstances during the 2010 eruption. This implies that such an offset could change in
609 a future eruption at Eyjafjallajökull, even if the same radar sensor is applied.

610 The data sets on which the empirical models are based include observations for various situations
611 regarding wind speed, but it is known that scenarios of weak eruptions under strong wind-conditions
612 are underrepresented (Mastin 2014). Furthermore, under high wind speeds plume height is reduced,
613 due to the facts that the plume is bent-over and that there is more efficient turbulent entrainment of
614 air into the column (Bursik 2001). For the phases of weak activity and relatively low wind speeds,
615 however, the empirical models by *Mastin* (phase II with 0.35) and *Sparks* (phase IV with 0.10) still
616 have shown to be among the ones with highest prediction success.

617 As a general finding, the explicitly wind-affected models tend to have larger ranges of uncertainty
618 and over rather than underestimate the measured mass (see, e.g., Fig. 11). Results from the simple
619 models fall at the other end of the modelled range, with *Wehrmann* being the most extreme
620 example, which tend to systematically underestimate the mass eruption rates, especially under
621 windy conditions. This discrepancy between estimates from simple and explicitly wind-affected
622 models under windy conditions is consistent with previous findings (Bursik 2001; Mastin 2014).

623 When analyzing the entire eruptive period by comparing the strategy-averaged *DIs* for simple plume
624 models (see Table 5), the *REFIR_15min* approach turned out to be the optimal strategy. In contrast,
625 using *REFIR* strategies with explicitly wind-affected models lead often to overestimated mass
626 predictions. These models are affected by larger uncertainties than the simple plume models (see,
627 e.g., Fig 11), which leads to larger MER predictions, due to the definition of the best estimate by
628 equation (15).

629 Of particular interest is also the comparison between phase Ib and III, which were similarly well
630 covered (9.3 vs 9.2 plume height measurements per hour), had only slightly dissimilar prevailing wind
631 speeds (19.2 vs 16.2 m/s), but differed in eruptive styles: phase Ib was phreatomagmatic, whereas

632 phase III was magmatic. Despite different outcomes for most of the other strategies, the *REFIR*
633 strategies appear to be the most robust for both types of eruptive styles. This robustness is
634 noteworthy, considering that phreatomagmatism is based on thermohydraulic fuel-coolant
635 interaction processes (Wohletz 1986; Büttner and Zimanowski 1998; Dürig et al. 2020), which
636 produce more fine-grained ash particles, release larger kinetic energy and generate steam - all of
637 which are factors that are expected to affect the plume rise behaviour (Koyaguchi and Woods 1996;
638 Sparks et al. 1997).

639 An important choice when selecting the plume height data handling strategy is the decision of
640 whether gap reconstruction is applied or not. In a real-time monitoring scenario (as well as in re-
641 analyses), where radar is the only source for plume heights, it is difficult to distinguish “real” gaps
642 from data missing for other reasons. For example, it was suggested that the data gap between 15
643 and 16 April 2010 (within phase Ia) was real, and a consequence of low plume heights (Mastin 2014).
644 According to our results, however, the optimal strategies for that phase are those that apply gap
645 reconstruction (e.g., *REFIR_15min*), indicating that the gap was not reflecting the status of the actual
646 plume. This is corroborated by photographs taken on that day, that prove the existence of a plume in
647 that period. In contrast, our findings for phase II indicate that the gaps during that phase were real.
648 In such scenarios of very weak eruptive activities, strategies with large timebases (such as *SH_phase*)
649 or applied gap reconstruction (such as *SM_6h_int*) result in overestimated mass predictions,
650 especially for explicitly wind-affected models. For modelling long-lasting eruptions with diverse
651 episodes, comparable to Eyjafjallajökull 2010, we would therefore recommend a mixture of
652 strategies, tailored to phase and model-type.

653 Similar to applying gap reconstruction, using a too large timebase result in real data gaps being
654 missed. Using a too short timebase, on the other hand, might lead to errors when monitoring an
655 evolving plume that has not reached stable buoyancy yet. Since this would violate the steady state
656 assumption of the models, the model predictions would not be reliable. We consider the latter
657 effect, however, as insignificant, when modelling a long-lasting eruption like Eyjafjallajökull 2010.

658 A key finding of this study is the dominating influence of the wind speed on the prediction quality of
659 plume height data for mass eruption rates, at least for moderate eruptions with bent-over plumes.
660 This is highlighted by the results of the phase-specific comparison of plume height data handling
661 strategies (Table 5) and illustrated in Figures 8 and 10. Phase Ia is characterized by the highest wind
662 speeds, and despite having a good data coverage (8.1 measurements per hour, see Table 1) the
663 strategies are on average considerably less successful than for phases Ib and III. All our results
664 indicate that for a moderate-to-weak eruption like Eyjafjallajökull 2010, it is increasingly challenging
665 for wind speeds exceeding ~20 m/s to find combinations of plume models and data handling

666 strategies with which the mass eruption rate can be accurately predicted in real-time. This inference
667 highlights how important it is to monitor such eruptions in real-time and use the ‘optimal’
668 combination of plume height data handling strategy and plume models.

669 The strategies presented were tested by focusing only on the best MER estimates, without
670 considering the ranges of uncertainty (with the *REFIR* strategies being notable exceptions). These can
671 become substantial and exceed the span of the measured mass by several orders of magnitude
672 (Figure 11, for uncertainty ranges in modelled MERs, see also online resource 5). The uncertainties
673 for $M_{modelled}$, as defined by equation (19), depend on model, timebase, strategy and the relative data
674 coverage of the studied period. For example, when considering the entire eruption period, the
675 uncertainties for *SH_3h_int* increase by 9 - 12% (depending on the model) compared to *SH_3h_gap*
676 (15 - 17% for *SM_3h_int* vs *SH_3h_gap*). For phase II alone, the uncertainty ranges between the two
677 strategies increase by up to 41% (up to 43% for *SM_3h_int* vs *SH_3h_gap*).

678 Selecting a short timebase has a significant effect on the modelled mass uncertainties. For example,
679 compared to the 3h-timebase, the *REFIR* uncertainties for a timebase of 15 minutes is reduced by up
680 to 24% when considering the whole eruption period. For phase IV, this decrease reaches 32.5%. The
681 largest uncertainties are found for the largest possible timebase, with strategy *SH_phase* (see Fig.
682 11b).

683 It is important to note that our error estimates did not consider model-inherent uncertainties (an
684 example for the uncertainties of the *Mastin* model is shown in Fig. 11a). For some strategies, the
685 model uncertainties could be larger than the ranges defined by equation (19). A model-strategy
686 combination with a large *DI* does not necessarily provide a “wrong” prediction, since its uncertainty
687 interval could (at least partly) overlap with the measured range. On the other hand, combinations
688 that result in predictions affected by a large uncertainty range are unfavorable for real-time
689 monitoring purposes. Our approach is optimized for finding the model-strategy pairs which provide
690 the best estimates that are closest to the measurement, but we note that this is not the only method
691 to assess the quality of strategies and models, and more advanced comparison methods might also
692 consider the range and overlap of modelled mass uncertainties.

693 The radar data by Arason et al. (2011) has been used by a plethora of studies on the 2010 ash plume
694 of Eyjafjallajökull (Kaminski et al. 2011; Bursik et al. 2012; Gudmundsson et al. 2012; Degruyter and
695 Bonadonna 2012; Devenish 2013; Ripepe et al. 2013; Woodhouse et al. 2013; Mastin 2014; Dürig et
696 al. 2015b; Dioguardi et al. 2020). Our suggested height correction still lies within the reported error
697 margins, but when used as input for plume models it will increase the MER predictions. Due to our
698 definition of uncertainties by equation (19), for strategies that do not apply height reconstruction,

699 the plume height uncertainties are only dependent on the radar height uncertainty itself. In Figure
700 11b, the effect of the height correction on the modelled mass is represented by the distance
701 between the lower error bar and the best estimate for these strategies (e.g., all *SM*-strategies).

702 In their study on Eyjafjallajökull, Woodhouse et al. (2013) applied an *SH* strategy with a timebase of
703 1h, but using the maximum plume height within each block instead of plume height averages. We
704 note that the 10% plume height reduction applied to convert *H* into *h* for Woodhouse is in the same
705 range as the height correction applied to the radar data. We therefore used comparable plume
706 heights as input as Woodhouse et al. (2013). Using the maximum plume heights per block instead of
707 averages resulted in MER predictions for phase Ia that exceeded our estimates by an order of
708 magnitude. We suggest that, when using Woodhouse to re-model the Eyjafjallajökull 2010 eruption,
709 using data handling strategies like *SH_3h_gap* with block-averaged plume heights might be the more
710 suitable strategy.

711 The *SH* strategies are the most commonly used for plume modelling (Kaminski et al. 2011; Bursik et
712 al. 2012; Degruyter and Bonadonna 2012; Woodhouse et al. 2013; Mastin 2014; Dürig et al. 2015b),
713 and our findings indicate that these strategies (especially *SH_3h_gap*) are often a good choice when
714 using explicitly wind-affected models, but they are far less successful with simple models. From the
715 dynamic sampling strategies, which are typically used for real-time monitoring, *DHM* strategies prove
716 to provide best results for explicitly wind-affected models. When used with simple models, however,
717 the optimal choice is the *REFIR* strategy.

718

719 Conclusions

720 Focusing on the different eruptive phases of the Eyjafjallajökull 2010 eruption, we used aerial and
721 ground-based photographs of the plume together with ground-based measurements of the erupted
722 mass to study and compare the predictive qualities of eight plume models in combination with 22
723 data handling strategies. Although the latter has a significant influence on the model outcome, their
724 influence has, to our knowledge, not yet been studied to this extent before.

725 The best re-analysis results for simple (non-explicitly wind-affected) plume models are in most cases
726 obtained by using strategies implemented in the software REFIR (Real-time Eruption source
727 parameters FutureVolc Information and Reconnaissance system), which apply dynamic sampling in
728 the form of a moving average and consider the plume height uncertainties. The only exception is
729 found for phase II, where the eruptive activity was weak, and the plume height was under the
730 detection limit of the radar for a considerable amount of time, causing data gaps in the records.

731 There, interpolating strategies generate ghost-data points, which lead to an overestimation of the
732 erupted mass. Under such conditions, strategies without gap reconstruction and with long timebases
733 (such as *SH_phase* or *SM_6h_gap*) are more reliable. For explicitly wind-affected models, the best
734 strategies were found to be static sampling strategies with short timebases and without gap
735 reconstruction (such as *SH_3h_gap*).

736 Our findings suggest that no single model has always the best answers, and different models and
737 data handling strategies work best under different plume and data conditions. Which model and
738 strategy to choose for optimal prediction results depends on data coverage, eruption intensity and,
739 according to our findings from Eyjafjallajökull, to a lesser extent on fragmentation mechanism (e.g.,
740 magmatic vs phreatomagmatic). The dominant factor on the prediction quality, however, was found
741 to be the wind speed. The higher the wind speed, the more the non-explicitly wind-affected models
742 tend to underestimate the MER. Conversely, the explicitly wind-affected *Woodhouse* model
743 overestimates it, when applying the *REFIR* strategies. Namely for wind speeds over 20 m/s, only a
744 few combinations of data handling strategies and plume models provide accurate predictions when
745 monitoring an Eyjafjallajökull 2010 – type of eruption.

746 A possible solution for obtaining accurate estimates in real time for such long-term eruptions of
747 variable and moderate intensity under varying wind speeds would therefore be to not rely on a single
748 combination of data handling strategy and plume model, but to be able to use a range of such
749 combinations, thus providing the possibility to adapt to the scenario monitored.

750

751

752 **Tables**

753 **Table 1:**

754 Modelled eruptive phases and ground truth. Modelled phases of the 2010 Eyjafjallajökull eruption
 755 are listed, together with the estimated ranges of erupted airborne tephra mass, as reported in
 756 Gudmundsson et al. (2012). The errors of erupted mass are denoted *ME*. In our study these values
 757 were used as ground truth for testing the model predictions with different plume height data
 758 handling strategies. Additional columns indicate duration, number of recorded radar measurements,
 759 data points (i.e., plume height measurements) per hour and time-averaged wind speed for each
 760 phase. Computation of wind speeds were based on ERA5 reanalysis data for the (radar-detected)
 761 plume heights.

phase	from (dd/mm)	until (dd/mm/yy)	duration (min)	erupted mass ± <i>ME</i> (×10 ¹⁰ kg)	number of radar measurements	data points per hour	wind speed m/s
Ia	14/04	16/04/10	3785	9.80 ± 2.10	511	8.1	36.4
Ib	17/04	18/04/10	2165	3.50 ± 0.98	335	9.3	19.2
I (total)	14/04	18/04/10	5945	13.30 ± 3.08	846	8.5	30.1
II	18/04	4/05/10	23765	4.20 ± 1.40	1005	2.5	11.8
III	5/05	17/05/10	18725	18.90 ± 4.90	2858	9.2	16.2
IV	18/05	21/05/10	4950	1.40 ± 0.42	430	5.2	14.5
all	14/04	21/05/10	53385	37.8 ± 9.8	5139	5.8	15.6

762

763 Table 2:

764 Tested data handling strategies. The columns list the aspects that were varied: sampling method,
 765 void reconstruction, gap reconstruction, timebase (i.e., duration of a block), time stepping and Y-
 766 correction.

Strategy	sampling	void reconstructed variable	gap reconstructed variable	timebase	time step	Y-correction
S_5min_gap	static	- n. a. -	- none -	5 min	5 min	- none -
SH_phase	static	height	- n. a. -	phase	phase	- none -
SH_6h_gap	static	height	- none -	6 h	6 h	- none -
SH_Y6h_gap	static	height	- none -	6 h	6 h	Yes
SH_3h_gap	static	height	- none -	3 h	3 h	- none -
SH_Y3h_gap	static	height	- none -	3 h	3 h	Yes
SH_6h_int	static	height	height	6 h	6 h	- none -
SH_Y6h_int	static	height	height	6 h	6 h	yes
SH_3h_int	static	height	height	3 h	3 h	- none -
SH_Y3h_int	static	height	height	3 h	3 h	yes
SH_5min_int	static	- n. a. -	height	5 min	5 min	- none -
SM_6h_gap	static	MER	- none -	6 h	6 h	- none -
SM_3h_gap	static	MER	- none -	3 h	3 h	- none -
SM_6h_int	static	MER	MER	6 h	6 h	- none -
SM_3h_int	static	MER	MER	3 h	3 h	- none -
SM_5min_int	static	- n. a. -	MER	5 min	5 min	- none -
DHM_3h	dynamic	height	MER	3 h	5 min	- none -
DHM_1h	dynamic	height	MER	1 h	5 min	- none -
REFIR_3h	dynamic	height	MER	3 h	5 min	- none -
REFIR_1h	dynamic	height	MER	1 h	5 min	- none -
REFIR_30min	dynamic	height	MER	30 min	5 min	- none -
REFIR_15min	dynamic	height	MER	15 min	5 min	- none -

767

768 Table 3:

769 Results for Y-values, based on 49 static 6h-blocks with a data coverage of 60/72 or higher. Numerical
 770 OD models stand for *Wilson, Wehrmann, Scollo* and *Andronico* models. For *Woodhouse*, ERA5
 771 weather data was used. Y-correction strategies multiplied the predicted mass with 1/Y.

		6h		3h	
		mean	stdev	mean	stdev
numerical OD models	average	0.90	0.10	0.91	0.09
	median	0.95	0.25	0.95	0.19
	maximum	2.36	2.15	2.06	1.45
Sparks	average	0.91	0.09	0.92	0.08
	median	0.96	0.24	0.95	0.19
	maximum	2.29	2.00	2.01	1.36
Mastin	average	0.89	0.11	0.91	0.10
	median	0.95	0.26	0.95	0.20
	maximum	2.44	2.31	2.12	1.55
Woodhouse	average	0.91	0.09	0.92	0.09
	median	0.97	0.26	0.95	0.19
	maximum	2.37	2.13	2.08	1.46

772

773

774 Table 4:

775 List of discongruence indices (*DIs*). A *DI* of 0 refers to a model-estimate, which lies within the range of
 776 measured mass according to Gudmundsson et al. (2012). This range is illustrated by red error bars in
 777 Fig. 5 and Fig. 6 and specified by the uncertainty *ME*. A *DI* of 0 is therefore congruent with the
 778 “ground truth”. The larger the discongruence index, the more the model-output deviates from the
 779 ground truth.

Discongruence Index (<i>DI</i>)	minimum deviation from measured mass	maximum deviation from measured mass
0	0	<i>ME</i>
1	<i>ME</i>	1.5 <i>ME</i>
2	1.5 <i>ME</i>	2 <i>ME</i>
3	2 <i>ME</i>	2.5 <i>ME</i>
4	2.5 <i>ME</i>	3 <i>ME</i>
5	3 <i>ME</i>	> 3 <i>ME</i>

780

781

782 Table 5:

783 Average discongruence indices listed by data handling strategy and reference period. The first value
 784 is the resulting *DI* for the six simple models tested (*Wilson, Andronico, Scollo, Wehrmann, Sparks,*
 785 *Mastin*), the second column presents the averages for the two wind-affected models (*Woodhouse,*
 786 *Aubry*). The smaller the value, the more often fits the modelled result with the measured mass.
 787 Values representing the most successful strategies for a phase are printed in bold.

Row	Ia		Ib		II		III		IV		all		average	
<i>S_5min_gap</i>	3.0	0.5	1.5	0	2.8	1.0	0.8	0	2.3	0	1.7	0	2.0	0.3
<i>SH_phase</i>	3.0	0.5	2.0	0	0.3	2.0	1.3	0	0.5	1.0	1.0	1.0	1.4	0.8
<i>SH_6h_gap</i>	2.8	0.5	1.5	0	0.3	1.5	0.5	0	0.5	1.0	1.3	0	1.2	0.5
<i>SH_Y6h_gap</i>	2.5	0.5	1.3	0.5	0.5	1.5	0.5	0.5	0.5	1.0	0.7	1.0	1.0	0.8
<i>SH_3h_gap</i>	3.5	0.5	1.7	0	0.5	0.5	0.5	0	0.5	1.0	1.3	0	1.3	0.3
<i>SH_Y3h_gap</i>	2.8	1.0	1.3	0	0.5	0.5	0.5	0	0.5	1.0	1.2	0	1.1	0.4
<i>SH_6h_int</i>	2.8	0.5	1.5	0	0.7	2.0	0.5	0	0.5	1.0	1.0	0	1.2	0.6
<i>SH_Y6h_int</i>	2.5	0.5	1.3	0.5	0.7	2.5	0.5	0.5	0.3	1.5	0.5	1.0	1.0	1.1
<i>SH_3h_int</i>	3.0	1.0	1.7	0	0.8	2.0	0.5	0	0.5	1.0	0.8	0	1.2	0.7
<i>SH_Y3h_int</i>	2.8	0.5	1.3	0	0.7	2.0	0.5	0.5	0.3	1.5	0.5	1.0	1.0	0.9
<i>SH_5min_int</i>	2.7	1.0	1.3	0	0.7	2.0	0.5	0	0.5	1.0	0.5	1.0	1.0	0.8
<i>SM_6h_gap</i>	2.2	0.5	0.5	0.5	0.3	1.5	0.5	0	0.5	1.0	0.5	0.5	0.8	0.7
<i>SM_3h_gap</i>	2.8	1.0	0.7	0.5	0.5	0.5	0.5	0	0.5	1.0	0.8	0	1.0	0.5
<i>SM_6h_int</i>	2.2	0.5	0.5	0.5	0.7	2.5	0.5	0	0.3	1.5	0.5	1.0	0.8	1.0
<i>SM_3h_int</i>	2.5	0.5	0.7	0.5	0.8	2.0	0.5	0	0.3	2.0	0.5	1.0	0.9	1.0
<i>SM_5min_int</i>	2.5	0.5	1.3	0	0.7	2.0	0.5	0	0.5	1.0	0.5	1.0	1.0	0.8
<i>DHM_6h</i>	3.0	1.0	1.7	0	0.7	2.0	0.5	0	0.5	1.0	0.8	0	1.2	0.7
<i>DHM_3h</i>	3.0	1.0	1.7	0	0.7	2.0	0.5	0	0.5	1.0	0.7	0	1.2	0.7
<i>REFIR_3h</i>	1.7	1.0	0.5	1.0	1.0	3.0	0.3	1.0	0.5	2.0	0.3	1.0	0.7	1.5
<i>REFIR_15min</i>	1.7	0.5	0.5	1.0	0.8	2.5	0.3	1.0	0.2	2.0	0.5	1.0	0.7	1.3
average	2.7	0.7	1.2	0.3	0.7	1.8	0.5	0.2	0.5	1.2	0.8	0.5	1.1	0.8

788

789

790

791 Table 6:

792 Average discongruence indices over all strategies, listed by reference period and plume model.

Row	Wilson	Wehrmann	Scollo	Andronico	Sparks	Mastin	Woodhouse	Aubry	average
Ia	1.55	4.90	2.20	1.85	1.85	3.55	0.95	0.40	2.16
Ib	0.40	3.10	0.70	0.70	0.70	1.75	0.10	0.40	0.98
II	1.05	0.80	0.70	0.70	0.80	0.35	1.10	2.45	0.99
III	0.10	2.00	0.10	0.05	0.05	0.95	0.10	0.25	0.45
IV	0.10	1.85	0.10	0.10	0.10	1.00	1.10	1.25	0.70
all	0.05	2.55	0.45	0.20	0.15	1.30	0.55	0.50	0.72
average	0.54	2.53	0.71	0.60	0.61	1.48	0.65	0.88	

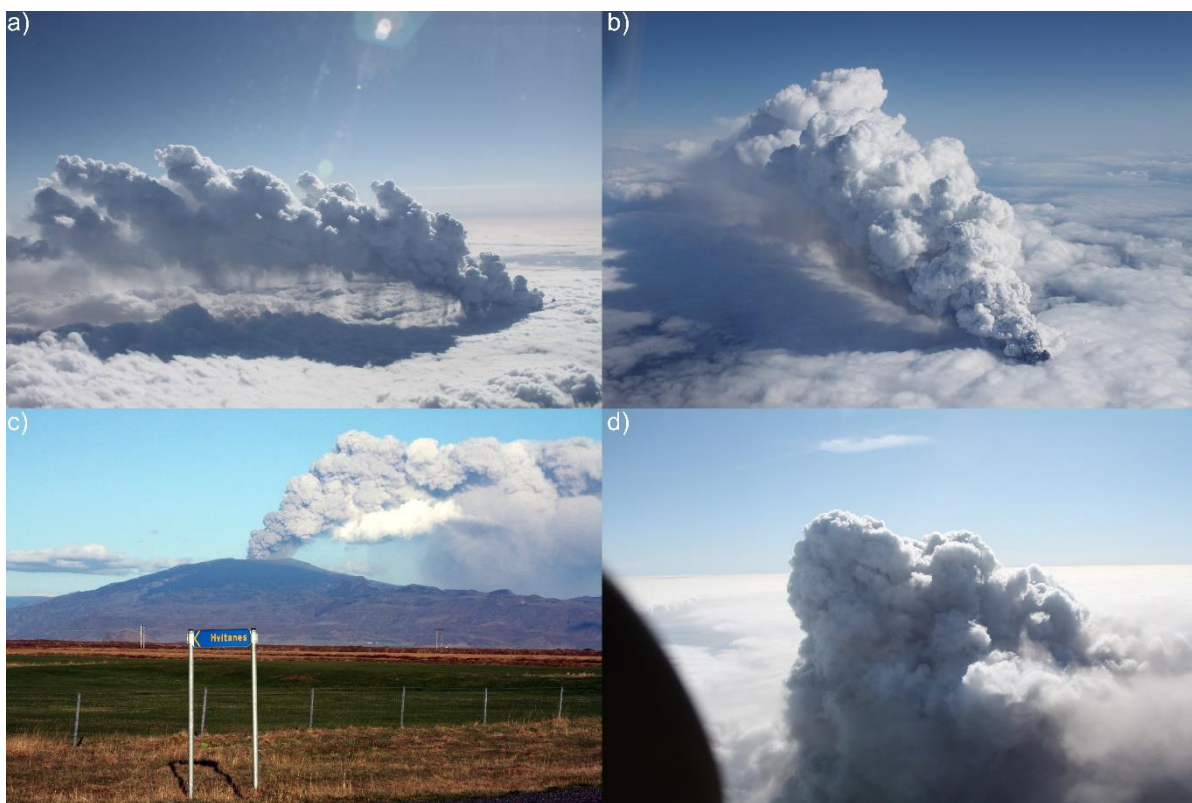
793

794

795 **Figures**

796 **Figure 1**

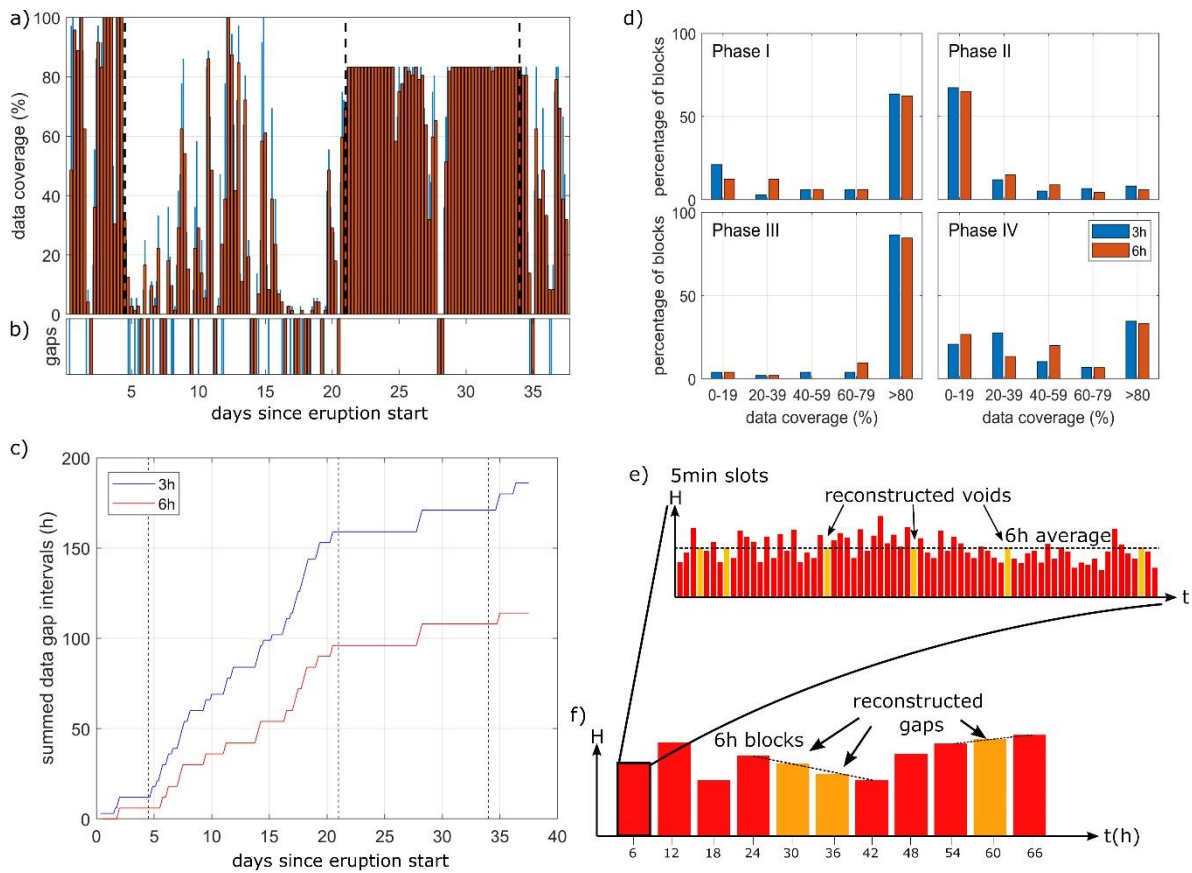
797 Examples of photos that were used to validate the plume height records from radar. a) Aerial picture
798 of the Eyjafjallajökull plume, taken on 14 April 2010 at 10:20 UTC. b) This photo was taken 5 minutes
799 later from a different angle. Since the GPS coordinates of the plane were logged, the horizontal
800 location of the plume top could be constrained by triangulation. At that time, the plume top reached
801 an altitude of (6.0 ± 0.3) km above vent. c) Example of a photo taken on the ground, from a distance
802 of 29 km from the vent. This picture dates from 10 May (19:50 UTC), when the plume was deflected
803 by 1.7 km towards 193° (SSW), with its top reaching an elevation of (2.2 ± 0.2) km above vent. d) This
804 photo was taken on the last day of the eruption (21 May) at 18:00 UTC. In this case the elevation of
805 the plume was directly assessed by flying at the same elevation as the plume top and using the on-
806 board altimeter. The aircraft flew at 12,000ft a.s.l., which translates into a plume height of (2.3 ± 0.1)
807 km. Photo credits: Th. Högnadóttir (a, b); M.T. Gudmundsson (c, d).



808

809 Figure 2

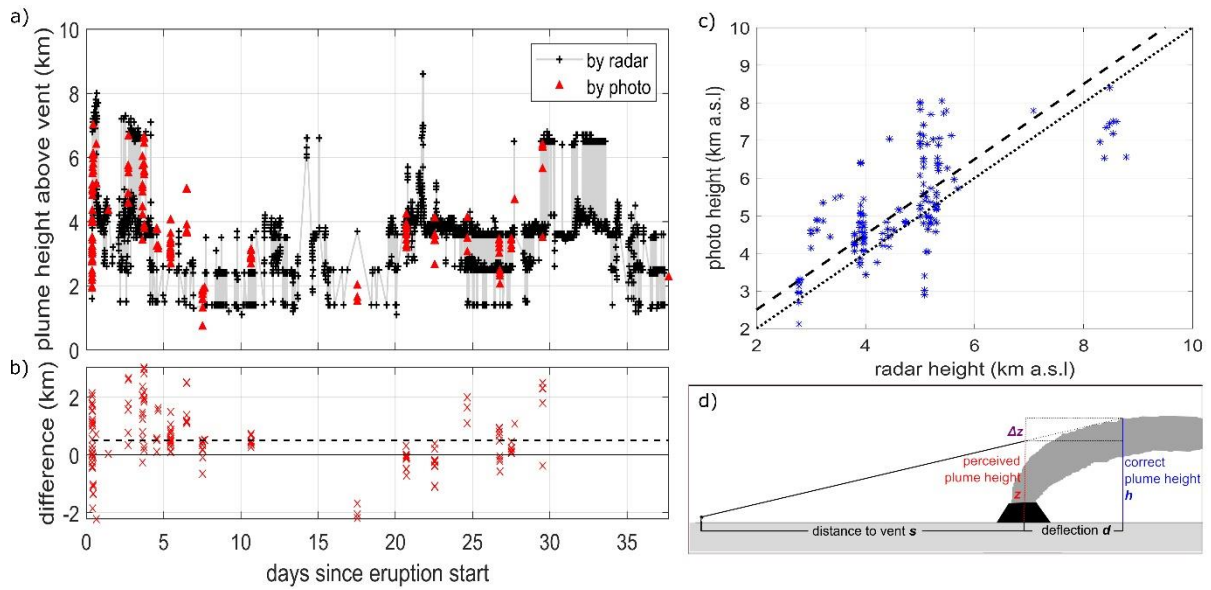
810 The radar record of the 2010 Eyjafjallajökull eruption. Data from Arason et al. (2011). a) Data
 811 coverage of 3h- (blue) and 6h -(red) blocks. For example, a 6h-block with 50% data coverage
 812 represents a set of 72 5-minute slots, of which 36 contain data. b) Timing of data gaps (i.e. blocks
 813 that are completely empty). c) Increasing the time step (i.e., block width), leads to a reduction of
 814 total overall gap time. d) Histograms show that data coverage of the blocks differed between the
 815 eruptive phases. e) Voids in the plume height records are reconstructed by filling them with the
 816 block's mean value (orange bars). f) Reconstructed gaps (orange bars) are obtained by interpolation.



817

818 Figure 3

819 Plume height validation by means of photographic plume analysis. a) Radar-based plume height
820 records of the Eyjafjallajökull 2010 eruption were compared with elevation data retrieved from
821 photos. b) Photo plume heights are systematically shifted upward by 0.5 km relative to radar plume
822 heights. c) This means that, statistically, the radar signal underestimated the actual plume height by
823 0.5 km. d) Due to windy conditions, the ash plume showed at times a considerable deflection d ,
824 which would lead to an underestimation of the plume if the distance to the plume top is assumed to
825 equal the distance to the vent s .

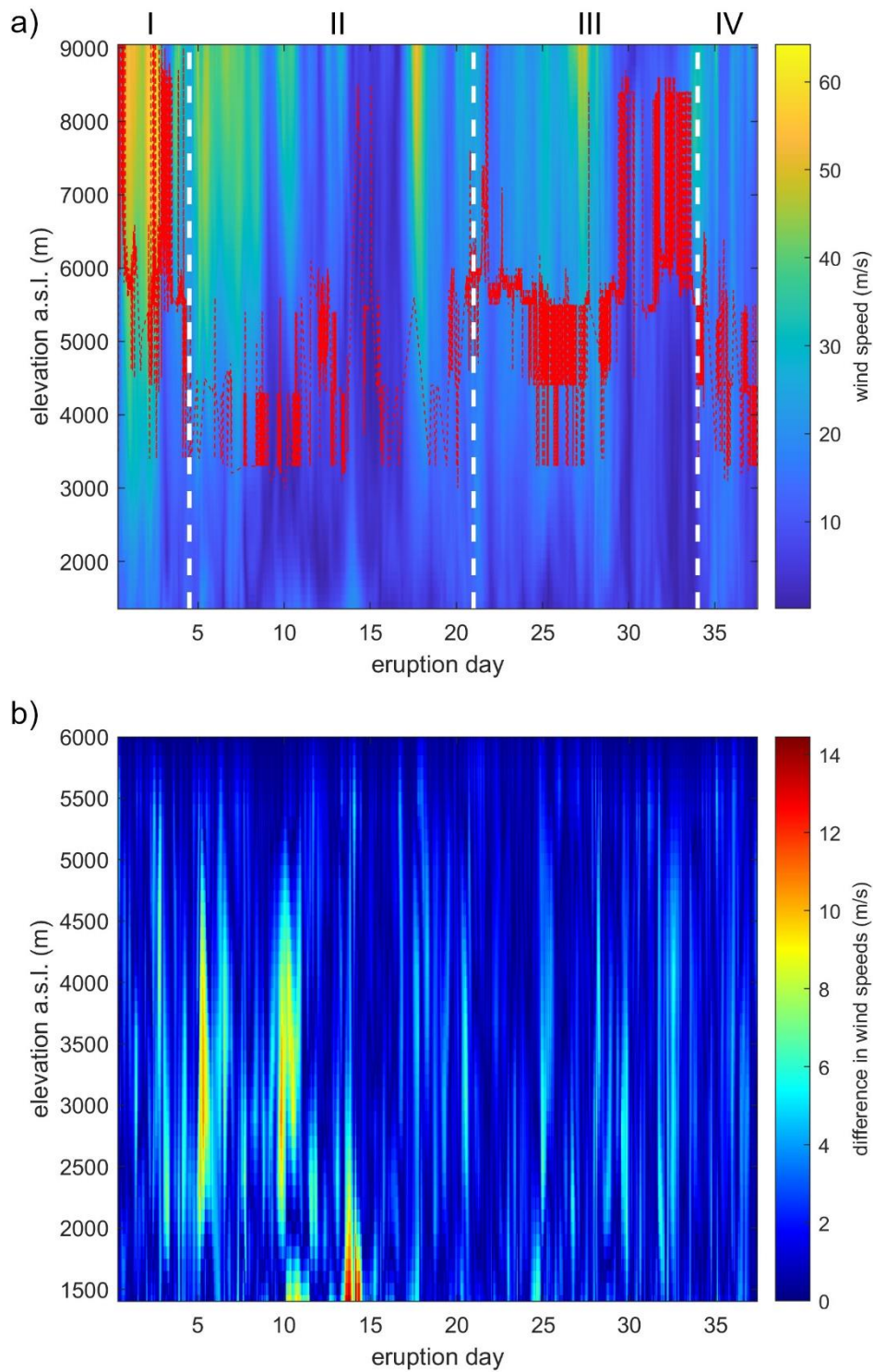


826

827

828 Figure 4

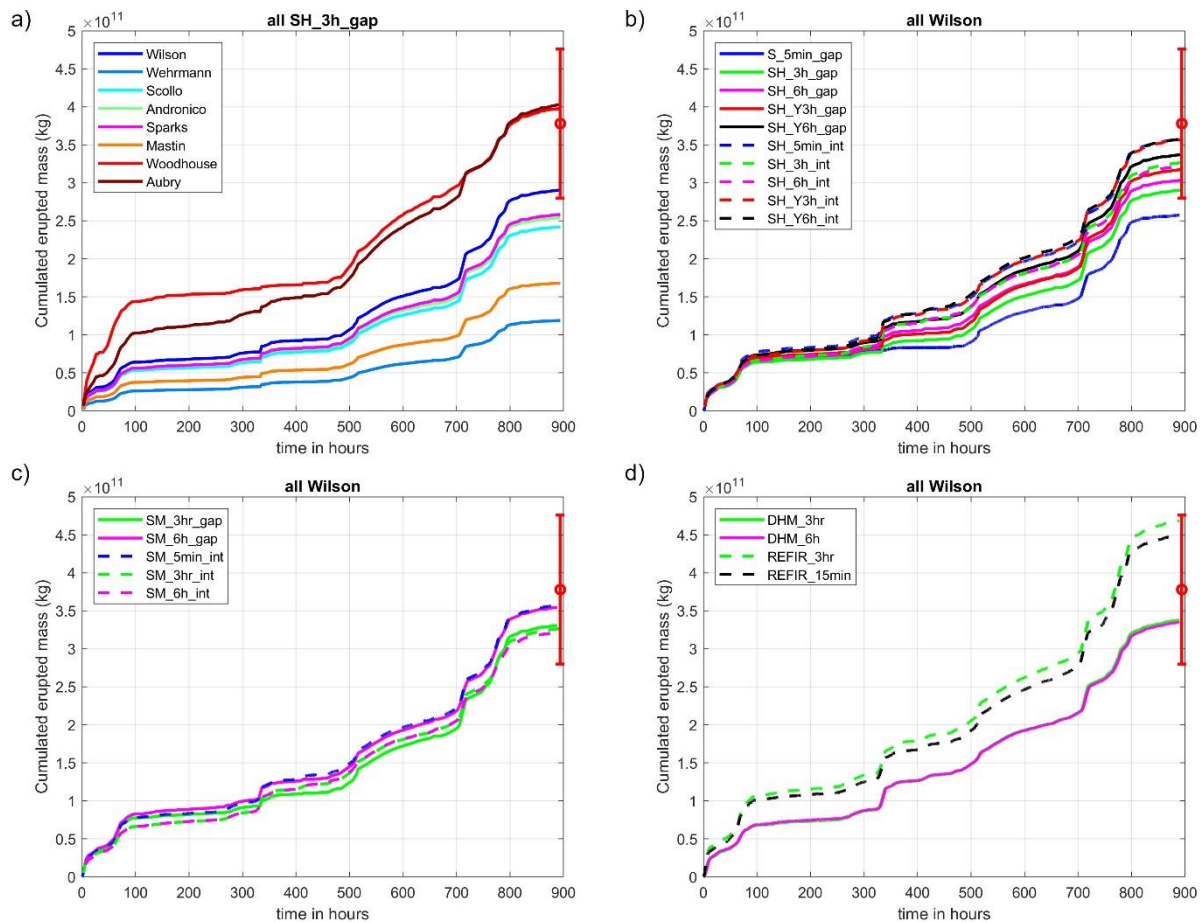
829 Wind speeds above Eyjafjallajökull during the 2010 eruption. a) Plume height (red curve) is plotted
830 together with altitude-specific wind speeds above the vent, based on ERA5 reanalysis data. Altitude
831 refers to sea level. Vertical dashed white lines mark margins of eruptive phases. b) Overall, the
832 absolute differences in wind speeds between ICRA and ERA5 were relatively small, with only sporadic
833 disparities, mainly at ground level on day 14.



834

835 Figure 5

836 Modelled erupted mass. Erupted mass is computed by multiplying the modelled mass eruption rates
837 by the block duration and summing the results over the reference period (here: the entire eruption).
838 The range of measured erupted mass as reported by Gudmundsson et al. (2012) (see also Table 1) is
839 displayed as red error bar with width ME . a) For each of the 8 tested models, 6h-averaged plume
840 heights were used as input with void but no gap reconstruction, thus applying strategy SH_3h_gap
841 (see Table 2). Wind speeds from ERA5 were used for *Woodhouse* and *Aubry*. b) Results for
842 S_5min_gap and SH plume height strategies applied to the *Wilson* model. c) Results for *Wilson* with
843 SM strategies. c) Predicted mass by *Wilson* with DHM and two $REFIR$ strategies.



844

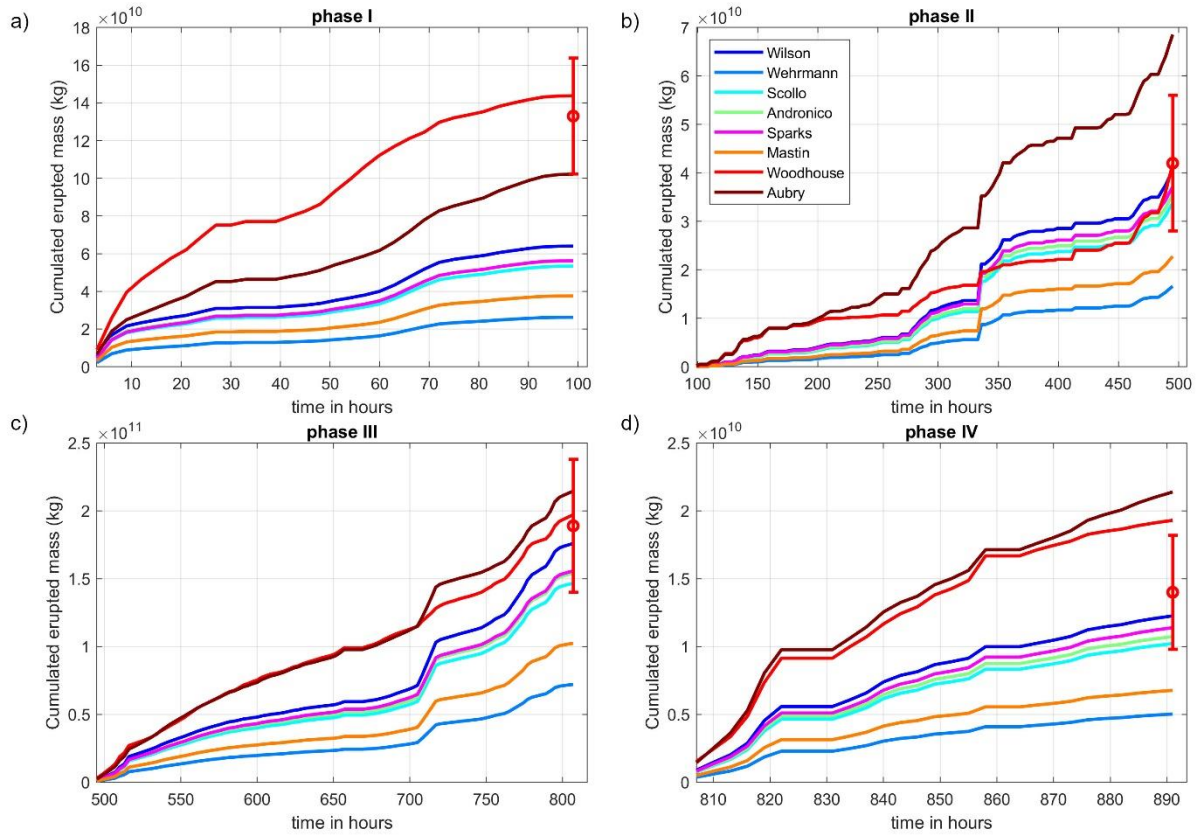
845

846

847 Figure 6

848 Modelled curves of erupted mass for each eruptive phase, applying strategy *SH_3h_gap* for eight
849 models. The range of measured erupted mass (ground truth) is indicated with a red error bar.

850 Reference periods are: a) phase I, b) phase II, c) phase III, d) phase IV.

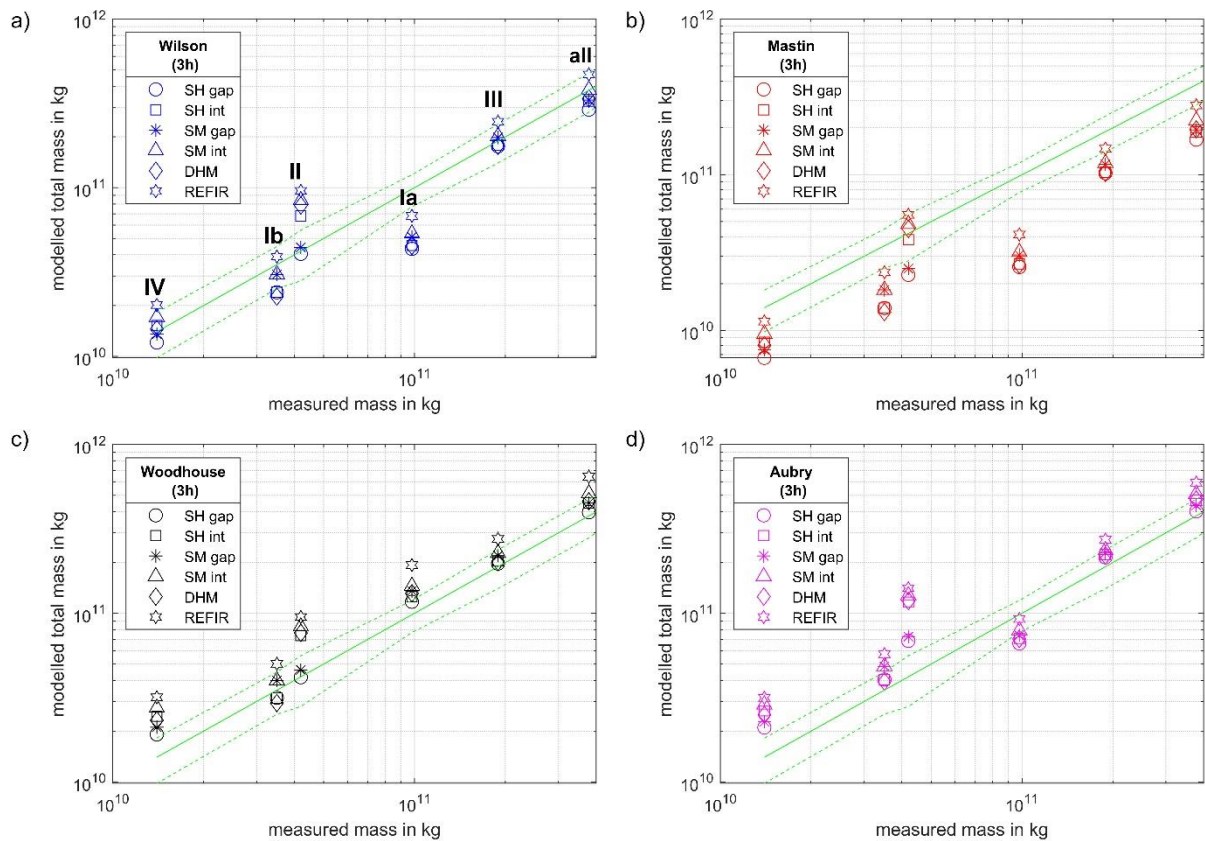


851

852

853 Figure 7

854 Phase-specific comparisons of modelled versus measured mass. Results for six strategies (marked by
855 different symbols) are presented for the whole eruption (*all*) and individual phases. The plots show
856 the model-specific outcomes for: a) *Wilson*, b) *Mastin*, c) *Woodhouse* d) *Aubry* (with using wind
857 speeds from ERA5 for both wind-affected models). The green lines indicate the range of erupted
858 mass, interpolated according to ground measurements. A symbol lying within this corridor indicates
859 that the respective strategy and plume model provides a 'successful' prediction under the tested
860 eruptive conditions.

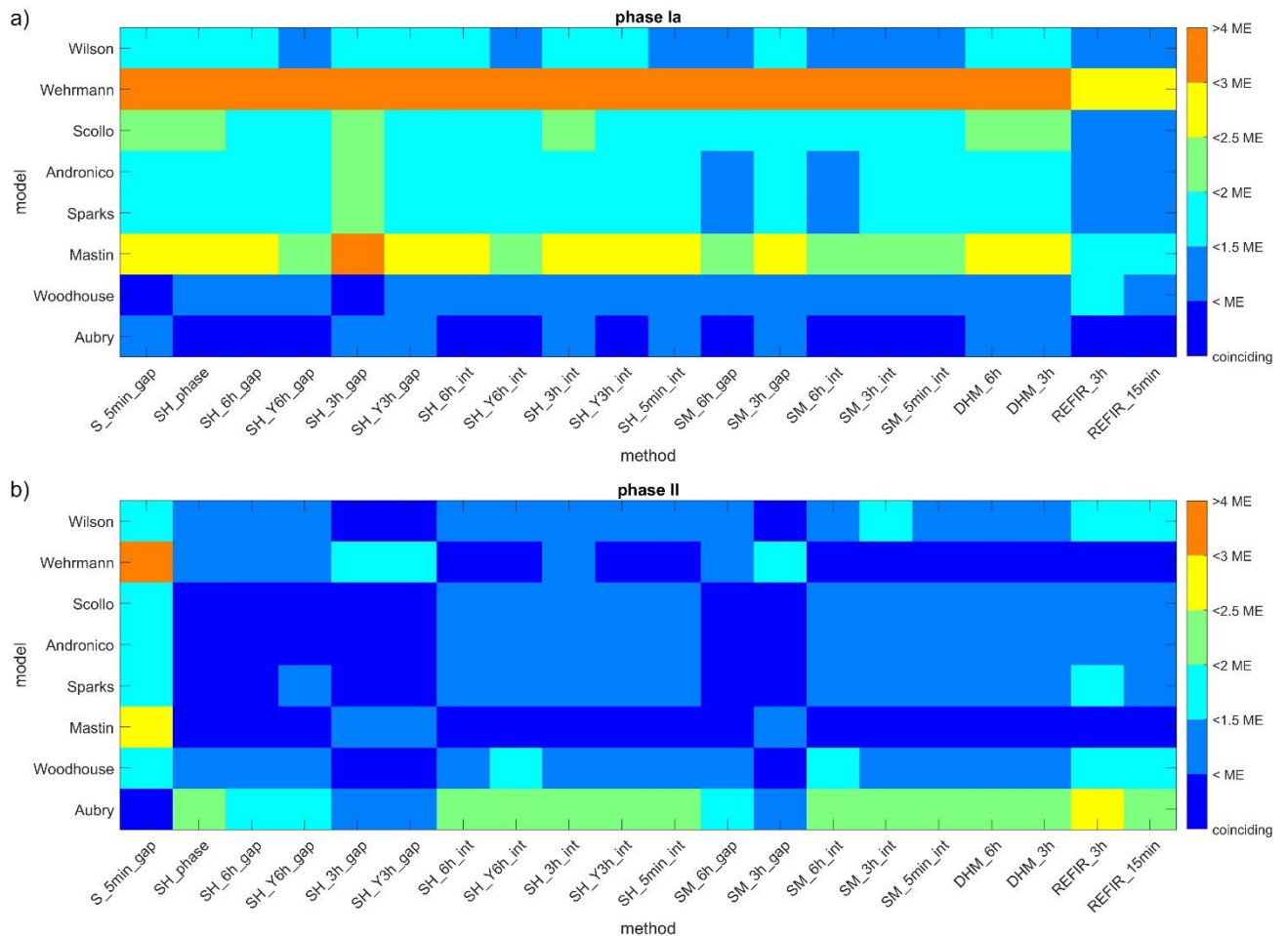


861

862

863 Figure 8

864 Congruency maps for combinations of strategies and models. ERA5 wind speeds were used for
865 *Woodhouse* and *Aubry*. Each cell's colour indicates how well the estimate of a strategy and model
866 pair fits with the measured data. Dark blue tiles indicate predictions that coincide with
867 measurements. a) Congruency map of all strategies and models for phase Ia, which was
868 characterized by the highest wind speeds. b) Congruency map for phase II, the phase with the most
869 gaps in the radar records.



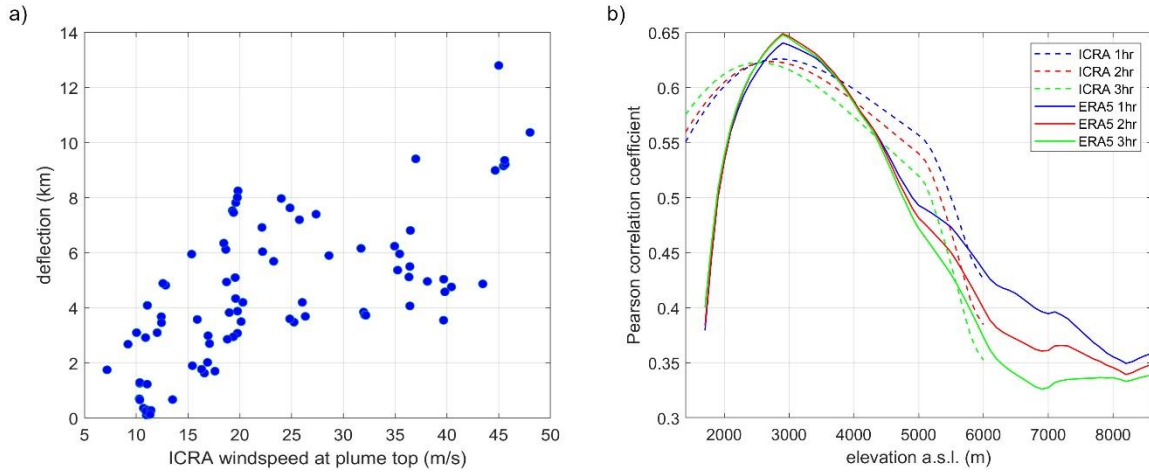
870

871

872 Figure 9

873 Wind-induced deflection of plume top. a) Photo-derived deflection of plume top is plotted versus
874 wind speeds at maximum plume height, according to ICRA reanalysis data. The data plots with a
875 Pearson correlation coefficient of 0.68. b) Correlation between deflection of the plume top and wind
876 speeds at different altitude levels (elevation above sea level).

877

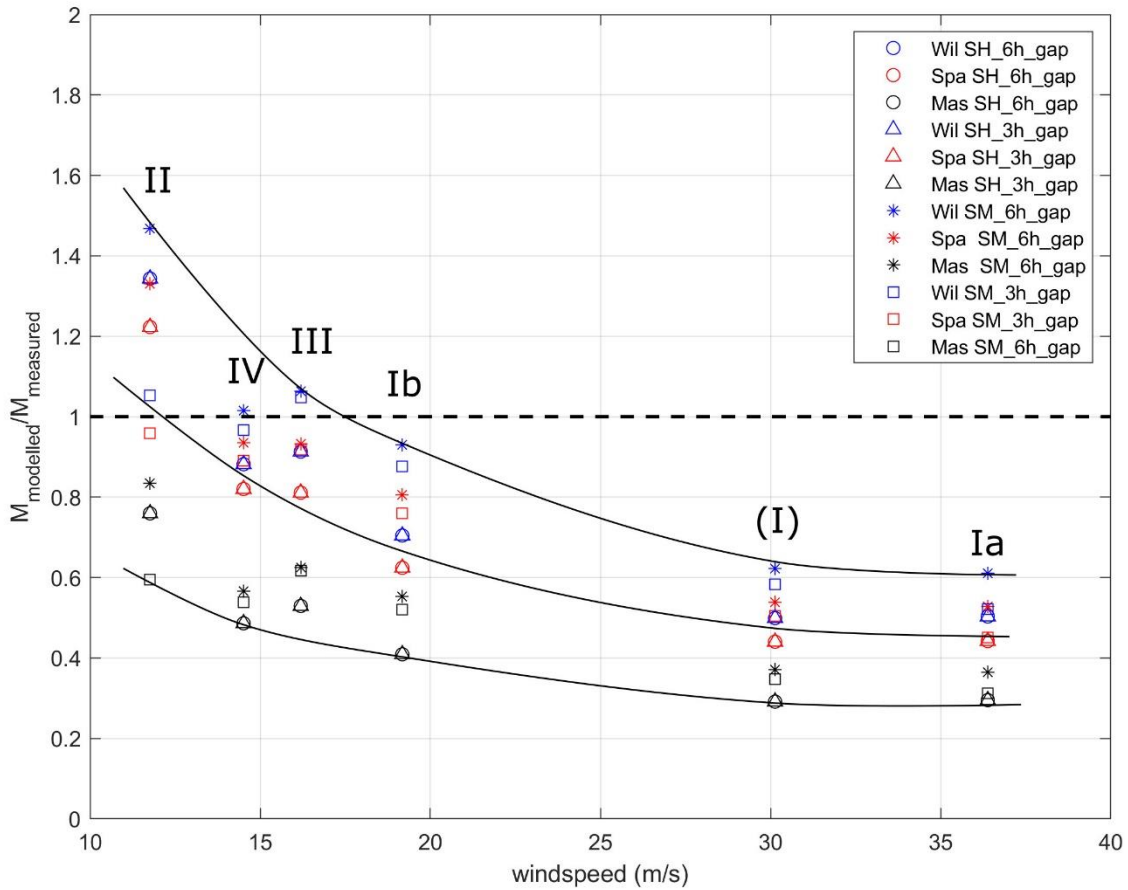


878

879

880 Figure 10

881 Wind-effect on simple (i.e., non-explicitly wind-affected) model predictions. Using four data handling
882 strategies for three models (*Wil*: *Wilson*, *Spa*: *Sparks*, *Mas*: *Mastin*), the ratio of modelled vs
883 measured mass is plotted over phase-averaged ERA5 wind speeds. The black curves serve for
884 visualization of what appears to be a general trend: the stronger the wind, the more the simple
885 plume models tend to underestimate the mass eruption rate, regardless of the applied strategy.



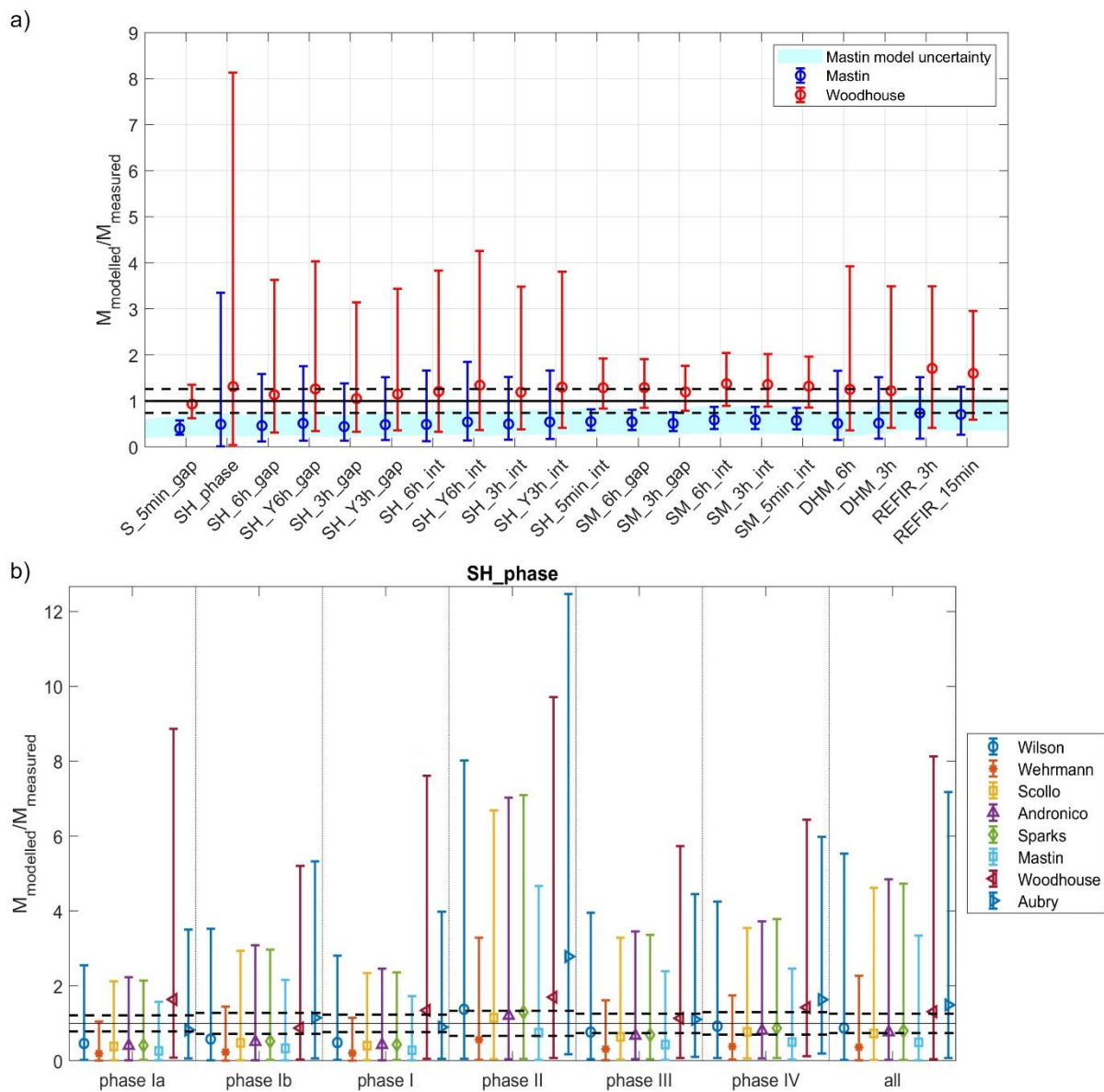
886

887

888 Figure 11

889 Ranges of uncertainty for modelled masses. a) Strategy-specific uncertainties for *Mastin* and
890 *Woodhouse*, when modelling the complete eruption period (*all*). For most strategies, the effect of
891 plume height errors according to equation (19) is larger than the model uncertainties (blue-shaded
892 background for *Mastin*). b) Ranges of predicted mass uncertainties for strategy *SH_phase* for all
893 studied eruptive phases and the complete eruption (*all*). Although depending on model and phase,
894 for this strategy the error bars are always considerably larger than the range of measured mass, due
895 to the large timebase.

896



897

898 Acknowledgements

899 We would like to thank Larry Mastin and two anonymous referees as well as handling editor
900 Julia Eychenne and executive editor Andrew Harris for their constructive comments which
901 helped to improve our manuscript. Rögnvaldur L. Magnússon did the geo-referencing of the
902 photos, which were mostly taken on board the aircraft of the Icelandic Coast Guard or those
903 of the private operator Eagle Air. Bolli Palmason from Icelandic Met Office is gratefully
904 acknowledged for providing ICRA data. The geo-referencing and photo analysis was
905 conducted under the EU Framework 7 FutureVolc project (2012-2016). This work
906 contributes to project MAXI-Plume, supported by the Icelandic Research Fund (Rannís),
907 grant Nr. 206527-051.

908 Online Resources

909 Online Resource 1 (.tif):

910 Corrected radar heights and ERA-5 wind speeds for different eruptive phases. a) phase I, b) phase II,
911 c) phase III, d) phase IV.

912 Online Resource 2 (.tif):

913 Phase-specific results for five plume height strategies applied to the Wilson model. The range of
914 measured erupted mass ("ground truth") is indicated with a red error bar. a) phase I, b) phase II, c)
915 phase III, d) phase IV.

916 Online Resource 3 (.tif):

917 Results for all models, using the REFIR strategies (*REFIR_15min*, *REFIR_30min*, *REFIR_1h* and
918 *REFIR_3h*), and the whole eruption (i.e, all phases) as reference period. *Woodhouse* and *Aubry* were
919 used with wind speeds from ERA5. Dashed horizontal lines indicate the range of measured erupted
920 mass.

921 Online Resource 4 (.tif):

922 Phase-specific results for all models, using the REFIR strategies (*REFIR_15min*, *REFIR_30min*,
923 *REFIR_1h* and *REFIR_3h*). Dashed horizontal lines indicate the range of measured erupted mass. a)
924 phase I, b) phase II, c) phase III, d) phase IV.

925 Online Resource 5 (.xlsx):

926 Modelled masses (in kg). Sheets are named after the plume height strategy applied. Columns show
927 eruptive phases. Rows show model results. Next to the best estimates, the range of uncertainty is
928 specified by providing the minimum (min) and maximum values (max).

929 Online Resource 6 (.jpg):

930 Congruency map for complete eruption (all).

931 Online Resource 7 (.jpg):

932 Phase-specific congruency maps for all strategies tested.

933 Online Resource 8 (.tif):

934 Phase-specific MER predictions for *Woodhouse* with *REFIR_3h*, using wind speeds of ERA 5 and ICRA.
935 In addition to 1h-timed wind speeds, 3h-averaged wind speeds were tested.

936

937

938 **Declarations**

939 **Funding**

940 TD was supported by the IRF (Rannís) PostDoctoral project grant 206527-051

941 **Conflicts of interest/Competing interests**

942 There are no competing interests

943 **Availability of data and material**

944 All data will be provided upon request to the lead author

945 **Code availability**

946 Not applicable

947 **References**

948 Andronico D, Scollo S, Caruso S, Cristaldi A (2008) The 2002–03 Etna explosive
949 activity: Tephra dispersal and features of the deposits. *J Geophys Res*
950 113:B04209. <https://doi.org/10.1029/2007JB005126>

951 Arason P, Petersen GN, Bjornsson H (2011) Observations of the altitude of the
952 volcanic plume during the eruption of Eyjafjallajökull, April–May 2010. *Earth Syst*
953 *Sci Data* 3:9–17. <https://doi.org/10.5194/essd-3-9-2011>

954 Aubry TJ, Engwell S, Bonadonna C, et al (2021) The Independent Volcanic Eruption
955 Source Parameter Archive (IVESPA, version 1.0): A new observational database
956 to support explosive eruptive column model validation and development. *J*
957 *Volcanol Geotherm Res* 417:107295.
958 <https://doi.org/10.1016/j.jvolgeores.2021.107295>

959 Aubry TJ, Jellinek AM, Carazzo G, et al (2017) A new analytical scaling for turbulent
960 wind-bent plumes: Comparison of scaling laws with analog experiments and a
961 new database of eruptive conditions for predicting the height of volcanic plumes.
962 *J Volcanol Geotherm Res* 343:233–251.
963 <https://doi.org/10.1016/j.jvolgeores.2017.07.006>

964 Bear-Crozier A, Pouget S, Bursik M, et al (2020) Automated detection and
965 measurement of volcanic cloud growth: towards a robust estimate of mass flux,

966 mass loading and eruption duration. *Nat Hazards* 101:1–38.
967 <https://doi.org/10.1007/s11069-019-03847-2>

968 Biass S, Bagheri G, Aeberhard W, Bonadonna C (2013) TError: towards a better
969 quantification of the uncertainty propagated during the characterization of tephra
970 deposits. *Stat Volcanol* 1:1–27. <https://doi.org/10.5038/2163-338X.1.2>

971 Björnsson H, Magnusson S, Arason P, Petersen GN (2013) Velocities in the plume of
972 the 2010 Eyjafjallajökull eruption. *J Geophys Res Atmos* 118:11,611–698,711.
973 <https://doi.org/10.1002/jgrd.50876>

974 Bonadonna C, Cioni R, Costa A, et al (2016) MeMoVolc report on classification and
975 dynamics of volcanic explosive eruptions. *Bull Volcanol* 78:84.
976 <https://doi.org/10.1007/s00445-016-1071-y>

977 Brannigan V (2011) Paradigms Lost: Emergency Safety Regulation under Scientific
978 and Technical Uncertainty. In: Alemanno A (ed) *Governing Disasters: The*
979 *Challenges of Emergency Risk Regulation*. Edward Elgar Publishing,
980 Cheltenham, pp 101–114

981 Bursik M (2001) Effect of wind on the rise height of volcanic plumes. *Geophys Res*
982 *Lett* 28:3621–3624. <https://doi.org/10.1029/2001GL013393>

983 Bursik M, Jones M, Carn S, et al (2012) Estimation and propagation of volcanic
984 source parameter uncertainty in an ash transport and dispersal model:
985 application to the Eyjafjallajökull plume of 14–16 April 2010. *Bull Volcanol*
986 74:2321–2338. <https://doi.org/10.1007/s00445-012-0665-2>

987 Büttner R, Zimanowski B (1998) Physics of thermohydraulic explosions. *Phys Rev E*
988 57:5726–5729. <https://doi.org/10.1103/PhysRevE.57.5726>

989 Büttner R, Zimanowski B, Röder H (2000) Short-time electrical effects during volcanic
990 eruption: Experiments and field measurements. *J Geophys Res Solid Earth*
991 105:2819–2827. <https://doi.org/10.1029/1999JB900370>

992 Calvari S, Büttner R, Cristaldi A, et al (2012) The 7 September 2008 Vulcanian
993 explosion at Stromboli volcano: Multiparametric characterization of the event and
994 quantification of the ejecta. *J Geophys Res Solid Earth* 117:B05201.

- 995 <https://doi.org/10.1029/2011JB009048>
- 996 Cerminara M, Esposti Ongaro T, Berselli LC (2016) ASHEE-1.0: a compressible,
997 equilibrium–Eulerian model for volcanic ash plumes. *Geosci Model*
998 *Dev* 9:697–730. <https://doi.org/10.5194/gmd-9-697-2016>
- 999 Cerminara M, Esposti Ongaro T, Valade S, Harris AJL (2015) Volcanic plume vent
1000 conditions retrieved from infrared images: A forward and inverse modeling
1001 approach. *J Volcanol Geotherm Res* 300:129–147.
1002 <https://doi.org/10.1016/j.jvolgeores.2014.12.015>
- 1003 Costa A, Suzuki YJ, Cerminara M, et al (2016) Results of the eruptive column model
1004 inter-comparison study. *J Volcanol Geotherm Res* 326:2–25.
1005 <https://doi.org/10.1016/j.jvolgeores.2016.01.017>
- 1006 Dacre HF, Grant ALM, Hogan RJ, et al (2011) Evaluating the structure and
1007 magnitude of the ash plume during the initial phase of the 2010 Eyjafjallajökull
1008 eruption using lidar observations and NAME simulations. *J Geophys Res*
1009 116:D00U03. <https://doi.org/10.1029/2011JD015608>
- 1010 de’Michieli Vitturi M, Neri A, Barsotti S (2015) PLUME-MoM 1.0: A new integral
1011 model of volcanic plumes based on the method of moments. *Geosci Model Dev*
1012 8:2447–2463
- 1013 Degruyter W, Bonadonna C (2012) Improving on mass flow rate estimates of volcanic
1014 eruptions. *Geophys Res Lett* 39:1–6. <https://doi.org/10.1029/2012GL052566>
- 1015 Dellino P, Gudmundsson MT, Larsen G, et al (2012) Ash from the Eyjafjallajökull
1016 eruption (Iceland): Fragmentation processes and aerodynamic behavior. *J*
1017 *Geophys Res Solid Earth* 117:B00C04. <https://doi.org/10.1029/2011JB008726>
- 1018 Devenish BJ (2013) Using simple plume models to refine the source mass flux of
1019 volcanic eruptions according to atmospheric conditions. *J Volcanol Geotherm*
1020 *Res* 256:118–127. <https://doi.org/10.1016/j.jvolgeores.2013.02.015>
- 1021 Dioguardi F, Beckett F, Dürig T, Stevenson JA (2020) The Impact of Eruption Source
1022 Parameter Uncertainties on Ash Dispersion Forecasts During Explosive Volcanic
1023 Eruptions. *J Geophys Res Atmos* 125:e2020JD032717.

- 1024 <https://doi.org/10.1029/2020JD032717>
- 1025 Dioguardi F, Dürig T, Engwell SL, et al (2016) Investigating Source Conditions and
1026 Controlling Parameters of Explosive Eruptions: Some Experimental-
1027 Observational- Modelling Case Studies. In: Németh K (ed) Updates in
1028 Volcanology - From Volcano Modelling to Volcano Geology. InTech, Rijeka
- 1029 Dürig T, Gudmundsson MT, Dellino P (2015a) Reconstruction of the geometry of
1030 volcanic vents by trajectory tracking of fast ejecta - the case of the Eyjafjallajökull
1031 2010 eruption (Iceland). *Earth, Planets Sp* 67:64.
1032 <https://doi.org/10.1186/s40623-015-0243-x>
- 1033 Dürig T, Gudmundsson MT, Dioguardi F, et al (2018) REFIR- A multi-parameter
1034 system for near real-time estimates of plume-height and mass eruption rate
1035 during explosive eruptions. *J Volcanol Geotherm Res* 360:61–83.
1036 <https://doi.org/10.1016/j.jvolgeores.2018.07.003>
- 1037 Dürig T, Gudmundsson MT, Karmann S, et al (2015b) Mass eruption rates in
1038 pulsating eruptions estimated from video analysis of the gas thrust-buoyancy
1039 transition—a case study of the 2010 eruption of Eyjafjallajökull, Iceland. *Earth,*
1040 *Planets Sp* 67:180. <https://doi.org/10.1186/s40623-015-0351-7>
- 1041 Dürig T, White JDL, Murch AP, et al (2020) Deep-sea eruptions boosted by induced
1042 fuel–coolant explosions. *Nat Geosci* 13:498–503.
1043 <https://doi.org/10.1038/s41561-020-0603-4>
- 1044 Engwell S, Sparks R, Aspinall W (2013) Quantifying uncertainties in the
1045 measurement of tephra fall thickness. *J Appl Volcanol* 2:5.
1046 <https://doi.org/10.1186/2191-5040-2-5>
- 1047 Esposti Ongaro T, Cavazzoni C, Erbacci G, et al (2007) A parallel multiphase flow
1048 code for the 3D simulation of explosive volcanic eruptions. *Parallel Comput*
1049 *33:541–560*. <https://doi.org/10.1016/j.parco.2007.04.003>
- 1050 Fierstein J, Nathenson M (1992) Another look at the calculation of fallout tephra
1051 volumes. *Bull Volcanol* 54:156–167. <https://doi.org/10.1007/BF00278005>
- 1052 Folch A, Costa A, Macedonio G (2016) FPLUME-1.0: An integral volcanic plume

- 1053 model accounting for ash aggregation. *Geosci Model Dev* 9:431–450.
1054 <https://doi.org/10.5194/gmd-9-431-2016>
- 1055 Gouhier M, Eychenne J, Azzaoui N, et al (2019) Low efficiency of large volcanic
1056 eruptions in transporting very fine ash into the atmosphere. *Sci Rep* 9:1449.
1057 <https://doi.org/10.1038/s41598-019-38595-7>
- 1058 Grindle TJ, Burcham FW (2002) Even minor volcanic ash encounters can cause
1059 major damage to aircraft. *ICAO J* 57:12–14
- 1060 Gudmundsson MT, Högnadóttir T, Dürig T, et al (2015) Field laboratory, aircraft
1061 observations and radar - Futurevolc Report D7.2. Reykjavík
- 1062 Gudmundsson MT, Thordarson T, Hoskuldsson A, et al (2012) Ash generation and
1063 distribution from the April-May 2010 eruption of Eyjafjallajökull, Iceland. *Sci Rep*
1064 2:1–12. <https://doi.org/10.1038/srep00572>
- 1065 Guffanti M, Casadevall TJ, Budding KE (2010) Encounters of aircraft with volcanic
1066 ash clouds: A compilation of known incidents, 1953-2009
- 1067 Harris A (2013) *Thermal Remote Sensing of Active Volcanoes: A User's Manual*.
1068 Cambridge University Press, New York
- 1069 Harris AJL (2015) Forecast communication through the newspaper Part 2:
1070 perceptions of uncertainty. *Bull Volcanol* 77:30. [https://doi.org/10.1007/s00445-](https://doi.org/10.1007/s00445-015-0902-6)
1071 [015-0902-6](https://doi.org/10.1007/s00445-015-0902-6)
- 1072 Harris AJL, Delle Donne D, Dehn J, et al (2013) Volcanic plume and bomb field
1073 masses from thermal infrared camera imagery. *Earth Planet Sci Lett* 365:77–85.
1074 <https://doi.org/10.1016/j.epsl.2013.01.004>
- 1075 Harris AJL, Gurioli L, Hughes EE, Lagreulet S (2012) Impact of the Eyjafjallajökull
1076 ash cloud: A newspaper perspective. *J Geophys Res Solid Earth* 117:.
1077 <https://doi.org/10.1029/2011JB008735>
- 1078 Hersbach H, Bell B, Berrisford P, et al (2020) The ERA5 global reanalysis. *Q J R*
1079 *Meteorol Soc* 146:1999–2049. <https://doi.org/10.1002/qj.3803>
- 1080 Johnson JB, Ripepe M (2011) Volcano infrasound: A review. *J Volcanol Geotherm*

- 1081 Res 206:61–69. <https://doi.org/10.1016/j.jvolgeores.2011.06.006>
- 1082 Kaminski E, Tait S, Ferrucci F, et al (2011) Estimation of ash injection in the
1083 atmosphere by basaltic volcanic plumes: The case of the Eyjafjallajökull 2010
1084 eruption. *J Geophys Res* 116:. <https://doi.org/10.1029/2011JB008297>
- 1085 Keiding JK, Sigmarsson O (2012) Geothermobarometry of the 2010 Eyjafjallajökull
1086 eruption: New constraints on Icelandic magma plumbing systems. *J Geophys
1087 Res Solid Earth*. <https://doi.org/10.1029/2011JB008829>
- 1088 Kienle J, Kyle PR, Self S, et al (1980) Ukinrek Maars, Alaska, I. April 1977 eruption
1089 sequence, petrology and tectonic setting. *J Volcanol Geotherm Res* 7:11–37.
1090 [https://doi.org/10.1016/0377-0273\(80\)90018-9](https://doi.org/10.1016/0377-0273(80)90018-9)
- 1091 Koyaguchi T, Woods AW (1996) On the formation of eruption columns following
1092 explosive mixing of magma and surface-water. *J Geophys Res B Solid Earth*.
1093 <https://doi.org/10.1029/95jb01687>
- 1094 Kristiansen NI, Stohl A, Prata AJ, et al (2012) Performance assessment of a volcanic
1095 ash transport model mini-ensemble used for inverse modeling of the 2010
1096 Eyjafjallajökull eruption. *J Geophys Res Atmos* 117:D00U11.
1097 <https://doi.org/10.1029/2011JD016844>
- 1098 Macrae D (2011) Which Risk and Who Decides When There Are So Many Players?
1099 In: Alemanno A (ed) *Governing Disasters: The Challenges of Emergency Risk
1100 Regulation*. Edward Elgar Publishing, Cheltenham, pp 13–24
- 1101 Magnússon RL (2012) Pixelcalc: Forrit til mælinga á stærð gosmakka út frá
1102 stafrænum myndum (Pixelcalc: Software to measure height of eruption plumes
1103 from digital photos). University of Iceland
- 1104 Marzano FS, Mereu L, Scollo S, et al (2020) Tephra Mass Eruption Rate From
1105 Ground-Based X-Band and L-Band Microwave Radars During the November 23,
1106 2013, Etna Paroxysm. *IEEE Trans Geosci Remote Sens* 58:3314–3327.
1107 <https://doi.org/10.1109/TGRS.2019.2953167>
- 1108 Mastin LG (2014) Testing the accuracy of a 1-D volcanic plume model in estimating
1109 mass eruption rate. *J Geophys Res Atmos* 119:2474–2495.

- 1110 <https://doi.org/10.1002/2013JD020604>
- 1111 Mastin LG, Guffanti M, Servranckx R, et al (2009) A multidisciplinary effort to assign
1112 realistic source parameters to models of volcanic ash-cloud transport and
1113 dispersion during eruptions. *J Volcanol Geotherm Res* 186:10–21.
1114 <https://doi.org/10.1016/j.jvolgeores.2009.01.008>
- 1115 Montopoli M (2016) Velocity profiles inside volcanic clouds from three-dimensional
1116 scanning microwave dual-polarization Doppler radars. *J Geophys Res Atmos*
1117 121:7881–7900. <https://doi.org/10.1002/2015JD023464>
- 1118 Morton BR, Taylor G, Turner JS (1956) Turbulent gravitational convection from
1119 maintained and instantaneous sources. *Proc R Soc London Ser A Math Phys Sci*
1120 234:1–23. <https://doi.org/10.1098/rspa.1956.0011>
- 1121 Nawri N, Pálmason B, Petersen NG, et al (2017) The ICRA atmospheric reanalysis
1122 project for Iceland, Report VÍ 2017-005, Icelandic Met Office, pp 37. Reykjavík
- 1123 Neri A, Papale P, Macedonio G, et al (1998) The role of magma composition and
1124 water content in explosive eruptions: 2. Pyroclastic dispersion dynamics. *J*
1125 *Volcanol Geotherm Res* 87:95–115. [https://doi.org/10.1016/S0377-](https://doi.org/10.1016/S0377-0273(98)00102-4)
1126 [0273\(98\)00102-4](https://doi.org/10.1016/S0377-0273(98)00102-4)
- 1127 Pavolonis MJ, Sieglaff J, Cintineo J (2018) Automated Detection of Explosive
1128 Volcanic Eruptions Using Satellite-Derived Cloud Vertical Growth Rates. *Earth*
1129 *Sp Sci*. <https://doi.org/10.1029/2018EA000410>
- 1130 Pioli L, Harris AJL (2019) Real-time geophysical monitoring of particle size
1131 distribution during volcanic explosions at stromboli volcano (Italy). *Front Earth*
1132 *Sci* 7:1–13. <https://doi.org/10.3389/feart.2019.00052>
- 1133 Pouget S, Bursik M, Webley P, et al (2013) Estimation of eruption source parameters
1134 from umbrella cloud or downwind plume growth rate. *J Volcanol Geotherm Res*
1135 258:100–112. <https://doi.org/10.1016/j.jvolgeores.2013.04.002>
- 1136 Pyle DM (1989) The thickness, volume and grainsize of tephra fall deposits. *Bull*
1137 *Volcanol* 51:1–15. <https://doi.org/10.1007/BF01086757>

- 1138 Ragona M, Hannstein F, Mazzocchi M (2011) The financial impact of the volcanic
1139 ash crisis on the European airline industry. In: Alemanno A (ed) *Governing*
1140 *Disasters: The Challenges of Emergency Risk Regulation*, Edward Elg.
1141 Cheltenham, pp 27–50
- 1142 Ripepe M, Bonadonna C, Folch A, et al (2013) Ash-plume dynamics and eruption
1143 source parameters by infrasound and thermal imagery: The 2010 Eyjafjallajökull
1144 eruption. *Earth Planet Sci Lett* 366:112–121.
1145 <https://doi.org/10.1016/j.epsl.2013.02.005>
- 1146 Schmidt L, Langen P, Aðalgeirsdóttir G, et al (2018) Sensitivity of Glacier Runoff to
1147 Winter Snow Thickness Investigated for Vatnajökull Ice Cap, Iceland, Using
1148 Numerical Models and Observations. *Atmosphere (Basel)* 9:450.
1149 <https://doi.org/10.3390/atmos9110450>
- 1150 Scollo S, Del Carlo P, Coltelli M (2007) Tephra fallout of 2001 Etna flank eruption:
1151 Analysis of the deposit and plume dispersion. *J Volcanol Geotherm Res*
1152 160:147–164. <https://doi.org/10.1016/j.jvolgeores.2006.09.007>
- 1153 Sparks RSJ, Bursik MI, Carey SN, et al (1997) *Volcanic plumes*. John Wiley & Sons,
1154 Chichester
- 1155 Suzuki YJ, Koyaguchi T (2012) 3-D numerical simulations of eruption column
1156 collapse: Effects of vent size on pressure-balanced jet/plumes. *J Volcanol*
1157 *Geotherm Res* 221–222:1–13. <https://doi.org/10.1016/j.jvolgeores.2012.01.013>
- 1158 Thorarinsson S (1954) The tephra-fall from Hekla on March 29th 1947. In: Einarsson
1159 T, Kjartansson G, Thorarinsson S (eds) *The Eruption of Hekla 1947–1948*.
1160 *Societas Scientarium Islandica*, Reykjavík, pp 1–68
- 1161 Thorkelsson B, Karlsdóttir S, Gylfason ÁG, et al (2012) The 2010 Eyjafjallajökull
1162 Eruption -Iceland Report to ICAO-June, pp 206. Reykjavík
- 1163 Tournigand P-Y, Fernández JJP, Taddeucci J, et al (2019) Time evolution of
1164 transient volcanic plumes: Insights from fractal analysis. *J Volcanol Geotherm*
1165 *Res* 371:59–71. <https://doi.org/10.1016/j.jvolgeores.2018.12.007>
- 1166 Valade SA, Harris AJL, Cerminara M (2014) *Plume Ascent Tracker: Interactive*

1167 Matlab software for analysis of ascending plumes in image data. Comput
1168 Geosci. <https://doi.org/10.1016/j.cageo.2013.12.015>

1169 Wehrmann H, Bonadonna C, Freundt A, et al (2006) Fontana Tephra: A basaltic
1170 Plinian eruption in Nicaragua. In: Rose WI, Bluth GJS, Carr MJ, et al. (eds)
1171 Volcanic Hazards in Central America. Geological Society of America, Boulder,
1172 Colorado, pp 209–223

1173 Wilson L, Self S (1980) Volcanic explosion clouds: Density, temperature, and particle
1174 content estimates from cloud motion. *J Geophys Res* 85:2567.
1175 <https://doi.org/10.1029/JB085iB05p02567>

1176 Wilson L, Walker GPL (1987) Explosive volcanic eruptions - VI. Ejecta dispersal in
1177 plinian eruptions: the control of eruption conditions and atmospheric properties.
1178 *Geophys J Int* 89:657–679. <https://doi.org/10.1111/j.1365-246X.1987.tb05186.x>

1179 Wohletz KH (1986) Explosive magma-water interactions: Thermodynamics, explosion
1180 mechanisms, and field studies. *Bull Volcanol* 48:245–264.
1181 <https://doi.org/10.1007/BF01081754>

1182 Woodhouse MJ, Hogg AJ, Phillips JC, Sparks RSJ (2013) Interaction between
1183 volcanic plumes and wind during the 2010 Eyjafjallajökull eruption, Iceland. *J*
1184 *Geophys Res Solid Earth* 118:92–109. <https://doi.org/10.1029/2012JB009592>

1185 Woods AW (1988) The fluid dynamics and thermodynamics of eruption columns. *Bull*
1186 *Volcanol* 50:169–193. <https://doi.org/10.1007/BF01079681>

1187