Regular Article

Mixing Bandt-Pompe and Lempel-Ziv approaches: another way to analyze the complexity of continuous-state sequences

S. Zozor^{1,2,a}, D. Mateos³, and P.W. Lamberti³

- ¹ Laboratoire Grenoblois d'Image, Parole, Signal et Automatique (GIPSA-Lab), CNRS, et Université de Grenoble, 961 rue de la Houille Blanche, 38402 Saint Martin d'Hères, France
- ² Instituto de Física de La Plata (IFLP), CONICET, and Departamento de Física, Facultad de Ciencias Exactas, Universidad Nacional de La Plata, C.C. 67, 1900 La Plata, Argentina
- ³ Facultad de Matemática, Astronomía y Física (FaMAF), CONICET, and Universidad Nacional de Córdoba, Avenidad Medina Allende, Ciudad Universitaria, X5000HUA, Córdoba, Argentina

Received 18 November 2013 / Received in final form 13 March 2014 Published online (Inserted Later) – © EDP Sciences, Società Italiana di Fisica, Springer-Verlag 2014

Abstract. In this paper, we propose to mix the approach underlying Bandt-Pompe permutation entropy with Lempel-Ziv complexity, to design what we call Lempel-Ziv permutation complexity. The principle consists of two steps: (i) transformation of a continuous-state series that is intrinsically multivariate or arises from embedding into a sequence of permutation vectors, where the components are the positions of the components of the initial vector when re-arranged; (ii) performing the Lempel-Ziv complexity for this series of 'symbols', as part of a discrete finite-size alphabet. On the one hand, the permutation entropy of Bandt-Pompe aims at the study of the entropy of such a sequence; i.e., the entropy of patterns in a sequence (e.g., local increases or decreases). On the other hand, the Lempel-Ziv complexity of a discrete-state sequence aims at the study of the temporal organization of the symbols (i.e., the rate of compressibility of the sequence). Thus, the Lempel-Ziv permutation complexity aims to take advantage of both of these methods. The potential from such a combined approach – of a permutation procedure and a complexity analysis – is evaluated through the illustration of some simulated data and some real data. In both cases, we compare the individual approaches and the combined approach.

1 1 Introduction

Many real signals result from very complex dynamics 2 and/or from coupled dynamics of many dimensional sys-3 tems. Various examples can be found in biology, such as 4 the reaction-diffusion process in cardiac electrical propa-5 gation that provides electrocardiograms, and the collective 6 actions of genes for the production of proteins in specific 7 quantities [1-3]. In finance, there is the example of the 8 variation in the price of an asset, which results from the 9 collective actions of the buyers and sellers [4], while statis-10 tical physics and social sciences also have huge numbers 11 of situations where 'complexity' emerges [5]. One of the 12 challenges is to describe these complex signals in a simple 13 way, to allow meaningful and relevant information to be 14 15 extracted [6-10].

The complex origin of such signals has led researchers to analyze these signals through tools that come either from the 'probability world', or conversely, from 'nonlinear dynamics'. The purpose is to characterize the degree of information or the complexity of the signals under analysis as well as possible. The first approach is statistical, and

the goal is to measure the spread of the distribution under-22 lying the data, or to detect any changes in the statistics. 23 The common tools that are used here come from infor-24 mation theory [8,9,11-13], or are correlation measures [3], 25 or come from spectral analysis [14]. The second approach 26 is devoted to signals that are produced by deterministic 27 (generally nonlinear) mechanisms, even if the sequence un-28 der analysis can appear to be somewhat 'random'. The 29 tools generally used for the description of such complex 30 signals come often from the chaos world, like Lyapunov 31 exponents, fractal dimensions, and others [6], or from the 32 concept of complexity in the sense of Kolmogorov (e.g., 33 Lempel-Ziv complexity) [7,10,15-17]. 34

The measures from information theory are very pow-35 erful, in a sense that they allow the quantification of a 36 degree of uncertainty (the rate) of a random sequence, 37 or of a sequence considered as randomly generated. How-38 ever, tools such as entropies can have some drawbacks 39 when used in practice. One of these occurs when dealing 40 with continuous-state data. In this case, the estimation 41 of a differential entropy from the data is not always an 42 easy task [18–20]. Some nonparametric estimators make 43 use of nearest neighbors, or of graph lengths, although 44 their properties are difficult to study [18–22]. More simple 45

^a e-mail: steeve.zozor@gipsa-lab.grenoble-inp.fr

estimators are based on 'plug-in' approaches [18]; namely, 1 the density is estimated using a Parzen-Rosenblatt ap-2 proach [23,24], and the estimation is plugged into the 3 mathematical expression of the entropy. The most sim-4 ple density estimator is based on a histogram, which is 5 equivalent to quantization of the data. The estimation 6 performance depends on this quantization (e.g., number 7 of thresholds, quantization intervals). To overcome this 8 potential difficulty, Bandt and Pompe proposed (i) to con-9 struct the multivariate trajectories from the scalar series, 10 i.e., an embedding; and (ii) to work with the so-called 11 vectors of permutation, i.e., for each point of the trajec-12 tory, its components are sorted, and each component of 13 the point is replaced by its position (rank) in the rear-14 ranged components [25]. Bandt and Pompe proposed then 15 to estimate the discrete entropy of the permutation vector 16 sequence, which led to the so-called *permutation entropy*, 17 and later on, to some variations of this measure [26-28]. 18 However, when dealing with sequences generated by a de-19 terministic process, such statistical measures can be inap-20 propriate because they measure an ensemble, or average, 21 behavior. 22

23 Conversely, for deterministic sequences generated by dynamical systems, there are a huge number of analysis 24 tools, like Lyapunov exponents, and fractal dimensions, 25 among others [29–31]. In general, the quantities under 26 study are relatively difficult to evaluate, and they require 27 long times of computation. As an example, there can be 28 29 the need to reconstruct a phase-space trajectory using sev-30 eral estimations to determine the embedding dimension 31 and the optimal delay, and then, in a second step, to es-32 timate some quantities from the reconstructed trajectory, 33 such as the whole Lyapunov spectrum, or just some exponents (e.g., positive, max), or dimensions [31,32]. More-34 35 over, these tools are generally designed specifically for the 36 study of chaotic series. A more natural concept of 'uncer-37 tainty' of a time series, whether chaotic or not, is that of its complexity in the sense of Kolmogorov. Roughly 38 speaking, this measures the minimal size of a binary pro-39 gram that can generate the sequence (i.e., the algorithmic 40 complexity) [33,34]. Among these, there is the Lempel-41 Ziv complexity, which is based on simple recursive copy-42 paste operations, as will be seen later [35,36]. This kind 43 of measure naturally finds applications in the compres-44 sion domain [33,36,37], and it is also used for signal anal-45 ysis [10,13,15,16]. A strength of this complexity is that 46 as it deals with a random discrete-state and ergodic se-47 quence, and when it is correctly normalized, it converges 48 to the entropy rate of the sequence [33,35]. In a sense, the 49 Lempel-Ziv complexity contains the concept of complex-50 ity both in the deterministic sense (Kolmogorov) and in 51 the statistical sense (Shannon). This property led to the 52 use of the Lempel-Ziv complexity for entropy estimation 53 purposes [21,38]. A possible drawback of the Lempel-Ziv 54 complexity is that it is defined for sequences that take 55 their values on a discrete (finite sized) alphabet. If it can 56 find natural applications that deal with discrete-state se-57 quences, such as DNA sequences or sequences generated 58 by logical circuits, while 'real-life' signals are generally 59

continuous states¹. Thus, to use the Lempel-Ziv complex-60 ity for signal characterization purposes, there is first the 61 need to quantize the data, which introduces some param-62 eters into the tuning. These parameters can influence the 63 behavior of the complexity of the quantized signal, as can 64 be seen, e.g., in reference [39], where for a logistic map, 65 some bifurcations are not (completely) captured by the 66 Lempel-Ziv complexity. 67

As can be imagined, there are many ways to overcome the drawbacks of purely statistical methods or purely deterministic approaches. Here, we concentrate on the 70 Lempel-Ziv complexity, using first the idea that underlies the Bandt-Pompe entropy, to 'quantize' a sequence to 72 analyze.

68

69

71

73

90

91

This report is organized as follows. In Section 2, we 74 first define the notation we use in the following sections. 75 Then we provide some basics on Bandt-Pompe entropy 76 (or permutation entropy). In the same section, we also 77 provide some basics on Lempel-Ziv complexity, proposing 78 then to 'mix' both of these approaches in Section 3, to 79 give what we call the Lempel-Ziv permutation complex-80 ity. In this same section, we provide some properties of 81 the Lempel-Ziv permutation complexity, including in an 82 Appendix the technical details and the description of a 83 practical way to calculate this complexity when dealing 84 with scalar sequences. We then illustrate in Section 4 how 85 the Lempel-Ziv permutation complexity can be used for 86 data analysis of both simulated sequences and biological 87 signals, and we finish the paper by drawing up our con-88 cluding remarks. 89

2 Notation and recall

2.1 Bandt-Pompe permutation entropy

The starting point of the Bandt-Pompe approach [25] ap-92 pears to take its origin from a study of chaos, and more 93 specifically, through the famous Takens' delay embedding 94 theorem [31,32]. The principle of this theorem is the re-95 construction of the state trajectory of a dynamical system 96 from the observation of one of its states. To fix the ideas, 97 consider a real-valued discrete-time series $\{X_t\}_{t\geq 0}$ that is 98 assumed to be a state of a multidimensional trajectory. 99 Consider two integers $d \geq 2$ and $\tau \geq 1$, and from the se-100 ries, let us then define a trajectory in the d-dimensional 101 space as: 102

$$\boldsymbol{Y}_{t}^{(d,\tau)} = \begin{bmatrix} X_{t-(d-1)\tau} & \dots & X_{t-\tau}X_{t} \end{bmatrix}^{t}, \quad t \ge (d-1)\tau$$
(1)

where the dimension d is known as the *embedding dimen*-103 sion, and where τ is called the *delay*. Takens' theorem 104

¹ When performing the acquisition of a signal in a computer, for example, the (discrete time) series is intrinsically a discretestate series due to the finite precision of the computer. However, this precision is generally high, so that the series can be assumed to be a continuous-state series. In particular, in general, the number of possible states is much higher that the number of samples to be analyzed.

79

1 gives conditions on d and τ such that $\boldsymbol{Y}_{t}^{(d,\tau)}$ preserves 2 the dynamical properties of the full dynamic system (e.g., 3 reconstruction of strange attractors) [31,32]. Many stud-4 ies have dealt with 'optimal' reconstruction of this phase 5 space; i.e., the choice of the correct embedding dimension, 6 and more particularly, the 'optimal' delay.

In reference [25], Bandt and Pompe did not focus es-7 pecially on chaotic signals, even if these signals serve as 8 illustrations. Thus, they did not focus on the phase-space 9 reconstruction problem. More precisely, they did not pro-10 vide discussion on the parameters d and τ . The only in-11 gredient they wished to conserve was the idea of taking 12 into account the dynamics of the system underlying an 13 observed signal. These questions of optimal reconstruc-14 tion also go beyond the scope of our paper, so we do not 15 discuss the choice of the embedding dimension and of the 16 delay in the sequel anymore. 17

Starting with the phase-space trajectory $\boldsymbol{Y}_{t}^{(d,\tau)}$, in-18 stead of focusing on the real-valued vectors, Bandt and 19 Pompe were interested in the order of the components of 20 the vectors. The principle consists first of the sorting (in 21 ascending order) of the components of $\mathbf{Y}_{t}^{(d,\tau)}$, and then the replacement of each component $X_{t-k\tau}$ by its rank/ position in the same large T22 23 tion in the sorted vector. This so-called *permutation vector* 24 is denoted as $\boldsymbol{\Pi}_{\boldsymbol{Y}^{(d,\tau)}}$ in the following. As an example, for 25 a vector $\boldsymbol{Y} = [Y_0^t \quad Y_1 \quad Y_2]^t$ such that $Y_2 \leq Y_0 \leq Y_1$, the permutation vector is $\boldsymbol{\Pi}_{\boldsymbol{Y}} = [1 \quad 2 \quad 0]^t$. Dealing with 26 27 28 random processes, it is then possible to define the *permu*-29 tation entropy as the Shannon entropy H of the (random) 30 permutation vector

$$H^{\pi}_{d,\tau}(X_t) \equiv H\left(\Pi_{\boldsymbol{Y}^{(d,\tau)}_t}\right). \tag{2}$$

For a stationary process, provided the size of the sequence 31 is large enough in terms of d!, the entropy can be estimated 32 via the frequencies of occurrence of any of the d! possi-33 ble permutation vectors in the sequence $\boldsymbol{Y}_{t}^{d,\tau}$. In their 34 paper, Bandt and Pompe defined the permutation en-35 tropy as the Shannon entropy of the frequencies of the 36 permutation vectors², which gives asymptotically the en-37 tropy $H_{d,\tau}^{\pi}(X_t)$ of equation (2) when dealing with a long-38 time (infinite) stationary and ergodic process, as indicated 39 in reference [25]. Starting from a sequence of length T, 40 $X_0 \ldots X_{T-1}$, in the sequel we write $H_{d,\tau}^{\pi}(X_{0:T-1})$ for the 41 entropy of the frequencies, to distinguish this from the 42 entropy of the random process. Several quantifiers of in-43 formation based on $\hat{H}^{\pi}_{d,\tau}(X_{0:T-1})$ were proposed in refer-44 ence [25], although such extensions go beyond the purpose 45 of the present paper. Thus, we do not present these here. 46

The idea behind permutation entropy is that the *d*! possible permutation vectors, also called *patterns*, might not have the same probability of occurrence, and thus,

this probability might unveil knowledge about the under-50 lying system. For a sequence of independent and iden-51 tically distributed (iid) variables, whatever the distribu-52 tion of the random variable, all of the patterns have the 53 same probability $\frac{1}{d!}$ of occurring (whatever the delay τ), 54 so that the permutation entropy is maximum and equal 55 to $\log(d!)$ [25]. Conversely, an important situation is rep-56 resented by the so-called *forbidden patterns*, which are 57 patterns that do not appear at all in the analyzed time 58 series [40-42]. As an example, it was shown in the logis-59 tic map $X_{t+1} = 4X_t(1 - X_t)$ that whatever the initial-60 ization X_0 , for d = 3 and $\tau = 1$, the permutation vec-61 tor $\begin{bmatrix} 2 & 1 & 0 \end{bmatrix}^t$ never appears. Such behavior shows how 62 the use of permutation vectors allows the distinguishing 63 between purely random sequences and deterministic se-64 quences (e.g., when the last one is chaotic, and thus ap-65 pears random): some authors have said that the presence 66 of forbidden patterns is an indicator of deterministic dy-67 namics [40-42]. This question remains, however, contro-68 versial, as it is possible to construct random series with 69 forbidden patterns [43], and conversely, a chaotic series 70 does not always show forbidden patterns [44]. 71

Note that if we work on a multidimensional sequence 72 $\{X_t\}_{t\geq 0}$, the permutation procedure can be performed 73 on each vector X_t , so that there are no embedding pro-74 cedures. To distinguish this situation from that of Bandt 75 and Pompe, we denote the permutation entropy and its 76 estimate as H^{π} and \hat{H}^{π} , respectively, without mention of 77 any delay and embedding dimension. 78

2.2 Lempel-Ziv complexity

Consider a finite-size sequence $S_{0:T-1} = S_0 \dots S_{T-1}$ of 80 symbols that take their values in an alphabet \mathcal{A} of fi-81 nite size $\alpha = |\mathcal{A}|$. In 1965, Kolmogorov introduced the 82 concept of the complexity of such a sequence as the size 83 of the smallest binary program that can produce the se-84 quence [33]. In an algorithmic sense, the Kolmogorov com-85 plexity measures the minimal 'information' contained in 86 the sequence, or the minimal information needed to gener-87 ate the sequence. Several years later, the seminal work of 88 Lempel and Ziv appeared [35], which dealt with the com-89 plexity of the Kolmogorov type of a sequence, restricting 90 this concept to the 'programs' based only on two opera-91 tions: recursive copy and paste operations. Their definition 92 lies in the two fundamental concepts of reproduction and 93 production: 94

- Reproduction: this consists of extending a se-95 quence $S_{0:T-1}$ of length T, adding a sequence $Q_{0:N-1}$ 96 via recursive copy-paste operations, which leads to 97 $S_{0:T+N-1}$, i.e., the first letter Q_0 is in $S_{0:T-1}$, let us 98 say $Q_0 = S_i$, the second one is the following one in the 99 extended sequence of size T + 1, i.e., $Q_1 = S_{i+1}$, etc.: 100 $Q_{0:N-1}$ is a subsequence of $S_{0:T+N-2}$. In a sense, all of 101 the 'information' of the extended sequence $S_{0:T+N-1}$ 102 is in $S_{0:T-1}$. 103
- **Production:** the extended sequence $S_{0:T+N-1}$ is now 104 such that $S_{0:T+N-2}$ can be reproduced by $S_{0:T-1}$. 105

² More precisely, in their paper, the permutation vector is defined as the time position of the component in the sorted vector, instead of the vector of the rank of the vector components. As there is a one-to-one mapping between the two ways of making, the entropy of the two vectors is the same.

4

1 The last symbol of the extension can also follow the 2 recursive copy-paste operation, so that the production 3 is a reproduction, but can be 'new'. Note thus that a

reproduction is a production, but the converse is false.

Any sequence can be viewed as constructed through a suc-5 cession of productions, called a history. As an example, a 6 sequence can be 'produced' symbol by symbol. However, a 7 given sequence does not have a unique history; several pro-8 cesses of productions can lead to the same sequence. In the 9 spirit of the Kolmogorov complexity, Lempel and Ziv were 10 interested in the optimal history; i.e., the minimal produc-11 tions needed to generate a sequence: the so-called Lempel-12 Ziv complexity, denoted as $C(S_{0:T-1})$ in the following, is 13 this minimal number of production steps needed for the 14 generation of $S_{0:T-1}$. In a sense, C describes the 'minimal' 15 information needed to generate the sequence by recursive 16 copy-paste operations. Thus, the approach of Lempel and 17 Ziv, and of several variations [36,37], naturally gave rise to 18 various algorithms of compression (including the famous 19 'gzip'). It can intuitively be understood that in a mini-20 mal sequence of production, all of the productions are not 21 reproductions, otherwise it would be possible to reduce 22 the number of steps [35]. This allowed the development 23 of simple algorithms for the evaluation of the Lempel-Ziv 24 complexity of a sequence [39]. 25

Surprisingly, although analyzing a sequence from a 26 completely deterministic point of view, it appears that 27 $C(S_{0:T-1})$ sometimes also contains the concept of infor-28 mation in a statistical sense. Indeed, it was shown in ref-29 erences [33,35] that for a random stationary and ergodic 30 process, when correctly normalized, the Lempel-Ziv com-31 plexity of the sequence tends to the entropy rate of the 32 33 process; i.e.,

$$\lim_{T \to +\infty} C(S_{0:T-1}) \frac{\log(T)}{T} = \lim_{T \to +\infty} \frac{H(S_{0:T-1})}{T}$$
(3)

where $H(S_{0:T-1})$ is the joint entropy of the *T* symbols, and the righthand side is the entropy rate (entropy per symbol) of the process. Such a property gave rise to the use of the Lempel-Ziv complexity for entropy estimation purposes [21,38].

Note that using the Lempel-Ziv complexity for analysis 39 purposes might not be envisaged if the size of the sequence 40 is not large enough in terms of the size of the alphabet. 41 Indeed, for small sequences compared to the size of the 42 alphabet, except for very elementary situations (e.g., con-43 stant signals, periodic signals), the complexity of the se-44 quence has a great probability of being close to the size of 45 the sequence. 46

47 **3** The Lempel-Ziv permutation complexity

As we have just seen, in a sense, the Lempel-Ziv complexity aims to capture a level of redundancy, or of regularity, in a sequence. Thus, this tool is interesting for the
analysis of signals that appear to be random, but that
hide some regularities, such as in chaotic sequences [39].

Conversely, viewing this complexity as an estimator of the 53 Shannon entropy when dealing with random sequences, its 54 use is also relevant in such a context. In some sense, it pro-55 vides a bridge between the two above-mentioned contexts. 56 However, a disadvantage of the Lempel-Ziv complexity is 57 that it is defined only for sequences of symbols taken in a 58 discrete (finite size) alphabet. Dealing with 'real-life' se-59 quences, a quantization has to be performed before its 60 use, as has been done in many of the studies dealing with 61 data analysis via this complexity [10,15,16]. Quantizing a 62 signal might have some consequences in the evaluation of 63 the complexity, and the effects of the parameters of the 64 quantizers appear difficult to evaluate. 65

Conversely, the permutation entropy also has some 66 drawbacks due to its statistical aspects. To illustrate why 67 sometimes it cannot capture the dynamics of a sequence, 68 consider the example of an iid scalar noise, versus a peri-69 odic scalar sequence of period T = 2. For an embedding 70 dimension d = 2 and a delay $\tau = 1$, in both cases the per-71 mutation vectors $\begin{bmatrix} 0 & 1 \end{bmatrix}^t$ and $\begin{bmatrix} 1 & 0 \end{bmatrix}^t$ appear with the same 72 frequency $\frac{1}{2}$ (assuming the length of the sequence is large 73 enough). Thus, the permutation entropy is equal in both 74 cases, and in this example it is thus not sensitive enough 75 to discriminate between the random iid sequence and the 76 periodic sequence³. Several variants to avoid such a draw-77 back can be imagined; e.g., taking into account the am-78 plitudes when constructing the permutation vectors. The 79 weighted-permutation entropy proposed in reference [28] 80 shows its efficiency for the detection of abrupt changes in 81 a sequence, but in the example given above, it will not 82 be able to discriminate between the two situations. More-83 over, when dealing with an intrinsic multidimensional se-84 quence, the permutation vectors do not clearly reflect any 85 dynamics. 86

To avoid the possible disadvantages of both methods, 87 we propose here to mix the Bandt-Pompe and Lempel-Ziv 88 approaches; i.e., to analyze the sequence of permutation 89 vectors via the Lempel-Ziv complexity. In this way, it is 90 expected that we can take advantage of both methods, 91 and thus reduce their respective drawbacks. In the fol-92 lowing, the so called *Lempel-Ziv* permutation complexity 93 of a finite length scalar sequence $X_{0:T-1}$ or a finite length 94 multivariate sequence $X_{0:T-1}$ are respectively denoted as: 95

$$C_{d,\tau}^{\pi}\left(X_{0:T-1}\right) \equiv C\left(\boldsymbol{\Pi}_{\boldsymbol{Y}_{(d-1)\tau}^{(d,\tau)}} \dots \boldsymbol{\Pi}_{\boldsymbol{Y}_{T-1}^{(d,\tau)}}\right) \qquad (4)$$

96

where $\boldsymbol{Y}_{t}^{(d,\tau)} = [X_{t-(d-1)\tau} \dots X_{t-\tau} X_{t}]^{t}$ and $\boldsymbol{\Pi}_{\boldsymbol{Y}_{T-1}^{(d,\tau)}}$ 97 is its permutation vector, and 98

$$C^{\pi}\left(\boldsymbol{X}_{0:T-1}\right) \equiv C\left(\boldsymbol{\Pi}_{\boldsymbol{X}_{0}}\dots\boldsymbol{\Pi}_{\boldsymbol{X}_{T-1}}\right).$$
 (5)

This way provides an answer to the necessity of working 99 with data taking the values on a finite size alphabet (here, 100

³ More rigorously, it is known that using the permutation entropy for data analysis, several embedding dimensions have to be tested. For d = 3 in this example, the permutation entropy makes the distinction between the iid noise and the periodic sequence.

59

76

77

78

84

94

95

96

the alphabet is $\mathcal{A} \equiv \{ [\pi(0) \quad \dots \quad \pi(d-1)]^t : \pi \in \Pi^{(d)} \}$ 1 of size $\alpha = d!$, where $\Pi^{(d)}$ is the ensemble of the d! pos-2 sible permutations on $\{0, \ldots, d-1\}$). Moreover, viewing 3 4 a permutation vector as quantization of the data, it is in-5 teresting to draw a parallel with dynamical quantization; 6 namely, of the sigma-delta type [45]. Indeed, dealing with scalar real-state sequences, in the case where $\tau = 1$ and 7 d = 2, for instance, the permutation vector is $\begin{bmatrix} 0 & 1 \end{bmatrix}^t$ if the 8 signal increases locally, and is $\begin{bmatrix} 1 & 0 \end{bmatrix}^t$ otherwise. In other 9 words, the two possible permutation vectors quantize the 10 variations of the signal in one bit. Roughly speaking, a 11 sigma-delta quantizer acts in a similar way⁴. For d > 2, 12 the same parallel should be made in some sense with the 13 so-called multi-stage sigma-delta quantizers [46]. This par-14 allel is another motivation to use permutation vectors as a 15 way to quantize a signal. Moreover, dealing with intrinsi-16 cally multivariate sequences, the permutation vectors can 17 be viewed as (vector) quantization of the real-valued vec-18 tors; this scheme does not need tuning parameters, con-19 trary to standard vector quantization schemes [45]. 20

Working on the permutation vectors maintains the 21 idea of studying the occurrences of patterns in a sequence. 22 By analyzing the permutation vectors via the Lempel-Ziv 23 complexity, a step is added because how the patterns are 24 temporarily organized is analyzed, rather than the fre-25 quency of occurrences. To stress this, let us come back 26 to the example of the permutation vector sequences of 27 an iid noise versus a periodic sequence of period T = 2. 28 As previously explained, the patterns $\begin{bmatrix} 0 & 1 \end{bmatrix}^t$ and $\begin{bmatrix} 1 & 0 \end{bmatrix}^t$ 29 appear with the same frequency in both cases. However, 30 the difference between the permutation vector sequences 31 in the two cases is that in the first case, the two patterns 32 appear in a random sequence, while in the second case, 33 they appear periodically: in the first case, the Lempel-Ziv 34 complexity is then high, while it is low (and equal to 3) 35 in the second case. With this very elementary example, 36 it can be seen why the Lempel-Ziv permutation complex-37 ity of a sequence can provide more information on the 38 dynamics; i.e., by analyzing how the patterns are orga-39 nized temporarily, not only in terms of the frequency of 40 occurrence. 41

Moreover, dealing with intrinsically multivariate se-42 quences, the argument of capturing the dynamics that un-43 derlie the sequence fails, as there is no embedding prior 44 to the quantization that is made with the construction of 45 the permutation vector. At least this question is not clear. 46 In essence, the Lempel-Ziv complexity will in a way cap-47 ture the dynamics of such a multivariate sequence, which 48 strengthens the interest for mixing both the Bandt-Pompe 49 and Lempel-Ziv approaches in this context. 50

The Lempel-Ziv permutation complexity has some 51 properties that have been inherited from the standard 52 Lempel-Ziv complexity. The first is the link with the 53 permutation entropy. Indeed, for a stationary ergodic 54 process that is scalar or multivariate, the sequence of 55

permutations remains stationary and ergodic, so that 56 equation (3) applies to the Lempel-Ziv complexity and the 57 entropy rate of this sequence, which can be written as: 58

$$\lim_{T \to \infty} C_{d,\tau}^{\pi} \left(X_{0:T-1} \right) \frac{\log T}{T} = \lim_{T \to \infty} \frac{H_{d,\tau}^{\pi} \left(X_{0:T-1} \right)}{T} \quad (6)$$

and

$$\lim_{T \to \infty} C^{\pi} \left(\boldsymbol{X}_{0:T-1} \right) \frac{\log T}{T} = \lim_{T \to \infty} \frac{H^{\pi} \left(\boldsymbol{X}_{0:T-1} \right)}{T}.$$
 (7)

The second property is the invariance of the Lempel-Ziv 60 permutation complexity to a given permutation applied 61 to the components of the vector of the initial series; i.e., 62 for any permutation matrix P, 63

> $C^{\pi}\left(\boldsymbol{P}\boldsymbol{X}_{0}\ldots\boldsymbol{P}\boldsymbol{X}_{T-1}\right)=C^{\pi}\left(\boldsymbol{X}_{0}\ldots\boldsymbol{X}_{T-1}\right).$ (8)

In other words, if a sequence of vectors X_t is constructed 64 from d scalar sequences, the choice of the order of the 65 components does not modify the value of the complex-66 ity of the 'joint' sequence. This property arises because 67 $\Pi_{PX_t} = P\Pi_{X_t}$ (permuting the components of a vector 68 results in permuting the components of its permutation 69 vector), together with the invariance of the Lempel-Ziv 70 complexity to a one-to-one transformation [17] 71

As shown by reference [17] for the Lempel-Ziv com-72 plexity, it is possible to build measures associated with the 73 Lempel-Ziv permutation complexity, although such possi-74 ble extensions go beyond the scope of the present paper. 75

Before moving on to put the Lempel-Ziv permutation complexity into action, let us just note the following additional choices:

- To take into account a finite resolution in data acqui-79 sition or to counteract possible low noise in the data, 80 we can introduce a radius of confidence δ ; i.e., if the 81 absolute value of the difference of two components is 82 strictly lower than δ , then they are considered to be 83 equal.
- Performing the permutation procedure, when two com-85 ponents of a vector are equal, we chose the 'smallest' 86 one as that with the lowest index (the oldest one in 87 the case of embedding). 88

(see Appendix for more details and further justification). 89 Once again, note that using the Lempel-Ziv permuta-90 tion complexity for analysis purposes might not be feasible 91 if the size of the sequence is not large enough in terms of 92 the size of the alphabet d!. 93

4 Illustrations based on synthetic and real data

4.1 Characterizing the logistic map

To illustrate how the Lempel-Ziv permutation complexity 97 can capture regularities in a signal, we consider here the 98 example of the famous logistic map 99

$$X_{t+1} = k X_t (1 - X_t), \qquad t \ge 0, \qquad k \in (0; 4].$$
(9)

More rigorously, it quantizes the difference between a sample and a prediction of this sample (the 'delta' part) in one bit. The prediction is made from all of the past samples, in general performing an integration or a summation (the 'sigma' part).

Page 6 of 12

1 We initialize X_0 randomly in [0; 1] so that the sequence 2 has a real value in the interval [0; 1]. This map has already 3 been taken as an illustration by both Bandt and Pompe in 4 reference [25], and Kaspar and Schuster in reference [39].

The logistic map has been studied for a long time, and 5 its behavior is well known and can be found in any text-6 book on chaos; e.g., [47,48]. Let us just recall that when k 7 increases, it shows more and more complex regimes: there 8 is an increasing sequence of values $k_{-1} = 0 < k_0 < \ldots < 0$ 9 $k_{\infty} \approx 3.56995$ such that, if $k \in (k_{n-1}; k_n]$, the out-10 put asymptotically oscillates between 2^n values, a phe-11 nomenon that is well known as *bifurcations*. For $k \ge k_{\infty}$, 12 the system is in a chaotic (unpredictable) regime. Roughly 13 speaking, it appears to behave randomly, although it is 14 produced by an elementary deterministic system. How-15 ever, in this zone, there remain some intervals, known as 16 islands of stability, in which the behavior is nonchaotic. 17 This briefly described behavior is summarized in the bi-18 furcation diagram plotted in Figure 1A. 19

Let us now study the regimes of the logistic map 20 versus k through the Lempel-Ziv permutation complex-21 ity proposed here. To this end, a sequence of size T =22 1000 is drawn and only the second half of the sequence, 23 which is assumed to be in the permanent regime, is ana-24 lyzed. The behavior of $C^{\pi}_{(d,\tau)}$ versus k is depicted in Fig-25 ure 1G, and this is compared to the permutation entropy 26 (Figs. 1D–1F), to the Lempel-Ziv complexity performed 27 on a static 2-level quantization of the signal $\mathbb{1}_{(.5;1]}(X_t)$, 28 where 1 is the indicator function (Fig. 1C), and to the 29 Lyapunov exponents (Fig. 1B). Roughly speaking, the 30 Lyapunov exponent⁵ measures the exponential conver-31 gence or divergence of two trajectories for two close initial 32 conditions: a positive Lyapunov exponent is a signature of 33 chaos [47, 48]. 34

The behavior of each descriptor can be interpreted as follows:

The Lyapunov exponent: this exponent clearly de-37 scribes the chaotic character of the logistic sequence 38 (when it is positive) versus its non-chaotic character 39 (when it is negative). However, as already mentioned 40 in the literature, this is not precise enough to distin-41 guish different types of behavior in nonchaotic regimes. 42 The Lempel-Ziv complexity $c(\{\mathbb{1}_{(.5;1]}(X_t)\})$: as claimed 43 by Kaspar and Schuster, this measure is more pre-44 cise than the Lyapunov exponent. In particular, the 45 complexity is very high in chaotic regimes, while it 46 is low in nonchaotic regimes. However, it can be seen 47 that the bifurcations are not detected very well. This 48 is clearly due to the quantization threshold. Indeed, 49 for k < 3.237, the system asymptotically oscillates be-50 51 tween two values > 0.5, the threshold that was chosen 52 by Kaspar and Schuster, which explains why the com-53 plexity fails to detect the bifurcations. The same phe-54 nomenon appears for the following bifurcations. Note



Fig. 1. Characterization of the logistic map versus k. (A) The bifurcation diagram; i.e., the values taken by the series in the permanent regime for each value of k. (B) The Lyapunov exponent λ . (C) The Lempel-Ziv complexity of the quantized signal $\mathbb{1}_{(.5,1]}(X_t)$ as in reference [39]. (D–F) The permutation entropy $\hat{H}^{\pi}_{(d,\tau)}$ with a delay $\tau = 1$ when $(d, \delta) = (3, 0)$ (D), $(d, \delta) = (5, 0)$ (E), and $(d, \delta) = (3, 10^{-3})$ (F). (G) The Lempel-Ziv permutation complexity $C^{\pi}_{(d,\tau)}$ for $(d, \tau, \delta) = (3, 1, 10^{-3})$.

that choosing a threshold of 2/3 for this system leads 55 to the detection of the first bifurcation, but the other 56 bifurcations remain undetected. 57

- The permutation entropies $\hat{H}^{\pi}_{(d,1)}$: in both cases of 58 d = 3 and d = 5, the permutation entropy precisely 59 characterizes the different regimes of the logistic map. 60 In particular, it is high in chaotic regions, while it is low 61 in nonchaotic regions; e.g., as is the case in the islands 62

⁵ For a discrete map of the type $X_{t+1} = f(X_t)$, this coefficient is given by $\lambda = \lim_{T \to \infty} \frac{1}{T} \sum_{t=1}^{T} \log f'(X_t)$ [47,48]. Practically speaking, this is calculated for a large T.

of stability. This is particularly true for the 'high' embedding dimension; e.g., d = 5. Note that for $\delta = 0$, the first bifurcation is not detected here. This is due to the small oscillations that remain around the limit value when $k \in (2; k_0]$. The consequence is that the permutation entropy fails to detect the first bifurcation, as the damped oscillatory behavior of the system for $k \in (2; k_0]$ is seen in the same manner as the sustained oscillations of the system when $k \in (k_0; k_1]$. Obviously, the permutation entropy ($\delta = 0$) detects this oscillatory behavior which is inherent to the system. However, if we are not really interested in the signal itself, but in its asymptotic regime, the small fluctuations can be viewed as perturbations. Choosing $\delta > 0$ allows the 'filtering' of these perturbations. In this case, even if the permutation entropy does not characterize the logistic sequence itself, it very precisely characterizes the

asymptotic regimes of the sequence, as can be seen in Figure 1. Indeed, in this case, the bifurcations are very well detected, even for 'low' embedding dimensions. – The Lempel-Ziv permutation complexity $C^{\pi}_{(d,1)}$: at a first glance, this measure behaves like the permutation

entropy. In particular, the same effects of detection 23 or not of the bifurcation occur if $\delta = 0$ (not plotted 24 in Fig. 1) or $\delta > 0$. Note, however, that even in the 25 low embedding dimension, the complexity appears to 26 better characterize the constant, oscillatory or chaotic 27 28 regimes. Indeed while $H^{\pi}_{(3,1)}$ is roughly constant when the chaos appears (for k slightly > k_{∞}), the complexity 29 greatly increases. 30

4.2 Detecting of a sudden change in a three-dimensional signal

1

2

3

4

5

6

7

8

9

10

11 12

13

14

15

16

17

18

19

20

21

22

To illustrate how the proposed measure can outperform 33 the permutation entropy in assessing the degree of com-34 plexity of some signals, let us consider a multidimen-35 sional series X_t composed first of N_c points issued from 36 37 a d-dimensional logistic series, followed by N_n points of 38 both spatially and temporally iid noise. The *d*-dimensional 39 logistic map we have chosen here for our purpose is de-40 scribed by the following equation:

$$\boldsymbol{X}_{t+1} = k \left(\boldsymbol{K} \boldsymbol{X}_t + \boldsymbol{1} \right) \odot \boldsymbol{X}_t \odot \left(\boldsymbol{1} - \boldsymbol{X}_t \right) \qquad (10)$$

where X_t is a *d*-dimensional vector, $\mathbf{1} = \begin{bmatrix} 1 & \dots & 1 \end{bmatrix}^t$, K41 is a $d \times d$ coupling matrix, and \odot is the component-wise 42 product $(t \ge 0)$. When **K** is zero, the *d* logistics are de-43 coupled. For the opposite, when K = 3P with P as the 44 cyclic permutation matrix of one place to the left, or when 45 $K = 11^{t}$, the map corresponds to the models proposed by 46 Lopez-Ruiz and Fournier-Prunaret in the 2-dimensional 47 and 3-dimensional contexts to model symbiotic interac-48 tions between species, where parameter k represented the 49 growth rate of the species [49,50]. In both the cases of 50 d=2 and d=3, according to the value of k, these maps 51 show regular orbits or chaotic orbits. We do not describe 52 53 here the richness of these maps, but instead direct the 54 reader to [49,50].



Fig. 2. Detection of a sudden change in a 3-dimensional sequence composed of $N_c = 2500$ points of a coupled 3-dimensional logistic map given by equation (10) followed by $N_n = 2500$ points of pure random noise (uniform). (A) A snapshot of the first component of such a sequence, with 2000 sequences then analyzed through a sliding window of $N_w = 500$ points, moving sample by sample. (B–D) Ten snapshots of the permutation entropies (B), the Lempel-Ziv complexities of a quantized version of the vectors (C), and Lempel-Ziv permutation complexity (D) are shown. Right: the corresponding histograms of the values taken by the measure, showing the windows in the chaotic part (solid line) and in the noise part (dashed line). The chaotic map is here strongly coupled, with K = 3P and k = 1.01.



Fig. 3. Same as for Figure 2 for a weakly coupled chaotic map, with K = .01P and k = 3.96.

For our purposes, we have chosen to study what hap-55 pens when the N_c first points of the sequence are gener-56 ated by the 3-dimensional map (d = 3) showing chaotic 57 behavior. We considered two cases: in the first, the cou-58 pling is K = 3P and k = 1.01; and in the second, 59 K = .01P and k = 3.96. In the first case, the compo-60 nents are strongly coupled, while they are weakly coupled 61 in the second case. A snapshot of these logistic map se-62 quences followed by pure noise is shown in Figures 2A 63 and 3A. Visually, it is relatively difficult to detect the 64 instant where the nature of the signal changes. Let us 65 then analyze the signal through sliding windows of size 66 N_w , moving point by point. In each window of the anal-67 ysis $(\mathbf{X}_{t-N_w+1},\ldots,\mathbf{X}_t), t = N_w - 1,\ldots$, we evaluate 68 the permutation entropy, the Lempel-Ziv complexity of 69 a quantized version of the components (by $\mathbb{1}_{[.5;+\infty)}$), and 70 the Lempel-Ziv permutation complexity. The results ver-71 sus t are plotted in Figures 2B-2D and 3B-3D, where 72 10 realizations are shown. On the right of Figures 2B–2D 73 1 and 3B-3D, the corresponding histograms are shown⁶ for 2 the values taken by each measure using 4×10^6 snapshots of 3 the chaotic map (solid lines) and the noise (dashed lines). 4 In these examples, the interpretations are the 5 following:

The permutation entropy: this index cannot detect the 6 change in the nature of the signal, as can be seen in 7 the snapshots for both the strong and weak coupling 8 (Figs. 2B and 3B). This is because, in these exam-9 ples, the patterns obtained in the permutation vectors 10 performed on the components appear with similar fre-11 quencies to the chaotic regime and in the noise regime. 12 By statistically analyzing these patterns, the dynamics 13 underlying the data are lost. The difficulty in the dis-14 crimination between chaos and noise is also illustrated 15 by the probabilities taken by the values of H^{π} : roughly 16 speaking, the probability of error in a discrimination 17 task is a function of the surface shared by the two 18 distributions. 19

- The Lempel-Ziv complexity: when looking at the case 20 of the strong coupling between the components of the 21 logistic, the Lempel-Ziv complexity performed on the 22 basic quantized version of the vector clearly discrim-23 inates between chaos and noise. However, when the 24 components are weakly coupled, this is no more the 25 case. This is clearly seen in the histograms that over-26 lap in the weak coupling case (Fig. 3C) while they are 27 28 separated in the strong coupling situation (Fig. 2C). Our interpretation of this effect is that, in a sense, the 29 Lempel-Ziv analyzes the components almost individu-30 ally: in the weak coupling case, it does not 'see' that 31 the components follow exactly the same dynamics and 32 are, in a sense, linked by these common dynamics. 33
- 34 The Lempel-Ziv permutation complexity: in both types 35 of coupling, this measure unambiguously detects the change in the nature. This can be viewed both in 36 37 the snapshots and in the probability distributions of the values taken by this measure (Figs. 2D and 3D). 38 39 Clearly, there is no overlap between the two his-40 tograms, which confirms that there is no probability 41 of error in the discrimination between the chaos and 42 noise in this illustration. From the curves, it would ap-43 pear that for both cases, the Lempel-Ziv permutation entropy shows a weaker dispersion around its mean 44 value than does the standard Lempel-Ziv complexity. 45
- These illustrations show that in spite of the power of the 46 permutation entropy to discriminate between chaos and 47 randomness, for instance, there are situations in which 48 this tool fails in this task. Using the Lempel-Ziv complex-49 ity of a basic quantized version of the sequence can be 50 an alternative, but this remains dependent on the quan-51 tification. Moreover, in this example, when there is no 52 coupling or there is weak coupling between the compo-53 nents, the permutation vector takes into account that the 54



Fig. 4. Electroencephalogram records of the analysis of a secondary generalized tonic-clonic epileptic seizure. The analysis was performed with a sliding window of 10 s (1024 points) moving sample by sample. (A) The original EEG. (B, C) The Lempel-Ziv analysis was performed on a 2-level quantization (B) and a 16-level quantization (C), and the quantizers were uniform over the dynamics of the analyzed window. (D, E) For both the permutation entropy (D) and the Lempel-Ziv permutation complexity (E), the permutation vectors were evaluated from a reconstructed phase-space trajectory with an embedded dimension d = 4 and a delay $\tau = 1$. The confidence radius was chosen as zero. The vertical dotted lines denote the characteristic times of T_1, T_2, T_3 and T_4 .

components follow exactly the same dynamics, which is55what the standard Lempel-Ziv complexity appears not to56do. For these interpretations, basically, we believe that57dealing with an intrinsic multidimensional sequence, the58Lempel-Ziv permutation complexity should be preferred59to the permutation entropy and the standard Lempel-Ziv60complexity.61

62

4.3 Epileptic electroencephalogram analysis

The electroencephalogram (EEG) signal analyzed in this 63 illustration corresponds to a scalp EEG record of a sec-64 ondary generalized tonic-clonic epileptic seizure, recorded 65 from a central right location (C4) of the scalp. This EEG 66 record is one of the EEGs studied by Rosso et al. in refer-67 ences [51-53]. It was obtained from a 39-year-old female 68 patient with a diagnosis of pharmaco-resistant epilepsy 69 (temporal lobe epilepsy), and no other accompanying dis-70 orders. The EEG signal is shown in Figure 4A. The epile-71 ptic seizure started at $T_1 = 80$ s, with a *discharge* of slow 72 waves that are superposed by fast waves with a lower 73 amplitude. This discharge lasts beyond $\Delta T = 8$ s, and 74 has a mean amplitude of 100 μ V. During the tonic-clonic 75 epileptic seizure, there are very high amplitudes that con-76 taminate the seizure recording, and the patient had to 77 be treated with an inhibitor of muscle responses. After 78

⁶ In the case of the Lempel-Ziv complexities, as these values can only take on discrete values between 2 and 500, their probability distributions are discrete. By misuse of representation, we have plotted them as continuous distributions to make their reading easier.

a short period, a desynchronization phase, known as the epileptic recruiting rhythm, appears in a frequency band centered at about 10 Hz, and it rapidly increases in ampli-

tude. After approximately 10 s, a progressive increase of 4 the lower frequencies (0.5-3.5 Hz) was observed [54]. For 5 the EEG studied here, this phase appears at $T_2 = 90$ s. 6 It is also possible to establish the beginning of the clonic 7 phase, at around $T_3 = 125$ s, and the end of the seizure at 8 $T_4 = 155$ s, where there is an abrupt decay of the signal 9

1

2

3

10

amplitude. The recorded signal has a duration of 180 s, and the 11 sampling frequency was 102.4 Hz (1024 samples/10 s) so 12 that we dispose of 18432 samples. To analyze the sig-13 nal, we again consider the methodology proposed in this 14 paper; namely, the evaluation of the Lempel-Ziv permu-15 tation complexity. This result is compared to that given 16 by the standard Lempel-Ziv performed on a static quan-17 tized version of the signal, and with the permutation en-18 tropy. The analysis was performed with sliding windows 19 of size $N_w = 1024$ points (10 s), which moved sample 20 by sample. Here, two quantized version are considered: 21 a 2-level Q_2 and a 16-level Q_{16} , both of which are uni-22 form over the range of the signal in the window of analy-23 sis. For the permutation measures, the permutation vec-24 tors were constructed with an embedding dimension and 25 a delay, of d = 4 and $\tau = 1$, respectively. We chose here 26 a radius of confidence of zero. The results are shown in 27 Figures 4B-4E. 28

The interpretations of these analyses are the following: 29

30 - The Lempel-Ziv complexity: for both the 2-level and 31 16-level quantization, this measure cannot detect any 32 change in the analyzed series. Although not plotted 33 here, we also tested 4-level and 8-level uniform quan-34 tizers, which leads to the same conclusion.

35 The permutation entropy: in this signal, the permutation entropy detects the appearance of the epileptic 36 seizure at $T_1 = 80$ s, which is visible in the signal. The 37 increase in the entropy measures a change in the na-38 ture of the signal; it is not just a change in amplitude, 39 otherwise the nature of the sequence of the permuta-40 tion vectors would not have been changed, and nor 41 would its entropy. Similarly, the characteristic times 42 $T_2 = 90$ s (not very visible in the signal), $T_3 = 125$ s 43 (the clonic phase) and $T_4 = 155$ s (end of seizure) that 44 are visible in the signal are also detected (as decreases 45 and an increase in the permutation entropy, respec-46 tively). However, the characteristic time T_3 is not well 47 detected by the permutation entropy. 48

The Lempel-Ziv permutation complexity: it can be seen 49 that the characteristic times detected by the per-50 mutation entropy are also clearly detected by the 51 Lempel-Ziv permutation complexity. The shape of this 52 complexity is very similar to that of the permutation 53 entropy. In particular, the Lempel-Ziv permutation 54 complexity detects a modification of the signal after 55 the time $T_2 = 90$ s, a change that is not particularly 56 detectable visually: at the peak, the analyzed window 57 is completely inside the 'complex part' of the crisis, but 58 the decrease indicates that the signal becomes more 59

and more organized. Finally, the Lempel-Ziv permuta-60 tion complexity better detects the modification of the 61 signal after the time $T_3 = 125$ s than the permutation 62 entropy. 63

Note that both the permutation entropy and the Lempel-64 Ziv permutation complexity appear to indicate the ap-65 pearance of an event at time 110 s, as seen by their in-66 creases. We have no interpretation yet as to this possible 67 event. Finally, the abrupt change that was detected by the 68 standard Lempel-Ziv complexity at time 165 s is only a 69 consequence of the abrupt change in the dynamics. 70

We can see in this example that the measure of com-71 plexity introduced in this paper increases steeply and very 72 precisely in time when the patient starts the seizure, and 73 even more, it can detect the different states of the tonic-74 clonic epileptic seizure. Note also the high level of the 75 complexity at the end of the signal compared to that at 76 the beginning. This level indicates that the signal remains 77 'disorganized'. A possible interpretation of such high com-78 plexity is that even if the epileptic sequence is apparently 79 ended, complex activity remains consequent to the cri-80 sis. A longer post-epilepsy sequence would be needed to 81 verify whether the complexity decreases to the low value 82 observed before the crisis. 83

As this signal serves essentially as an illustration, and 84 as our goal here is not to carry out deep EEG analyses, we 85 will not go further with this analysis. We also do not com-86 pare our result here to those obtain in references [51-53], 87 which merits a study in itself.

5 Discussion

Data analysis has a long history and still gives rise to a 90 huge amount of research. Among the challenges, especially 91 for the analysis of natural signals such as biomedical sig-92 nals, there is the need to characterize the degree of organi-93 zation or the degree of complexity of signal sequences, the 94 problem of detecting sudden sequence changes that are not 95 detectable visually, and the problem of characterization of 96 the nature of specific changes in a sequence. The literature 97 on information theory on the one hand, and on dynami-98 cal systems analysis on the other, provides an important 99 number of tools and methods to solve these challenges. 100

In this paper, we propose a tool that mixes two very 101 well known approaches: the permutation entropy and the 102 Lempel-Ziv complexity. The idea is to try to take the ad-103 vantage of both of these approaches, the first of which is 104 statistical, and the second of which is deterministic. 105

The Lempel-Ziv complexity has long been known and 106 was initially introduced in the compression domain. How-107 ever, it has been shown to be powerful for data analysis. 108 On the other hand, the permutation entropy allows a part 109 of the dynamics of a signal underlying data to be cap-110 tured when it is performed on reconstructed phase-space 111 signals. Moreover, in some sense, it is based on a kind 112 of quantization of the data, by considering only the ten-113 dencies rather than the values of the sequence. From this 114 last, it appears natural to quantize data, as has been done 115

89

88

Page 10 of 12

via the permutation vectors of a vector sequence (natu-1 ral or reconstructed) followed by the evaluation of the the 2 complexity of such a quantized sequence. The association 3 of these two approaches has here 'given birth' to what 4 we have named the Lempel-Ziv permutation complexity, 5 which is at the heart of our proposal. 6

In this paper, in particular, we have shown how the 7 Lempel-Ziv permutation complexity of a sequence can 8 precisely capture the degree of organization of such se-9 ries. When dealing with scalar sequences, the Lempel-Ziv 10 permutation complexity appears to give similar results to 11 those of the permutation entropy, even if one measure is 12 statistical while the other is purely deterministic. How-13 ever, when dealing with intrinsic multidimensional sig-14 nals, without procedures of phase-space reconstruction, 15 the entropy performed on the permutation vectors built 16 from the vector sequences cannot capture the dynamics 17 that underlie the data. Indeed, the calculation of the fre-18 quency of occurrence of such permutation vectors is then a 19 point-by-point analysis, and the links between successive 20 points are lost. Conversely, as the Lempel-Ziv complex-21 ity aims to detect regularities in a sequence by analyzing 22 how the symbols (numerical scalar samples, vectors, or 23 any kind of symbol) can be predicted algorithmically from 24 the past symbols, it captures the dynamics of the signal. 25 Doing this analysis for the permutation vector sequences 26 allows the natural solving of the question of quantization 27 of the data, as by definition, the Lempel-Ziv complexity 28 works with sequences of symbols lying on a discrete fi-29 nite size alphabet. As shown in our illustration, we can 30 imagine many situations for which the Lempel-Ziv per-31 mutation complexity can capture a degree of organiza-32 tion, while the permutation entropy fails, especially when 33 dealing with multidimensional signals; i.e., without phase-34 space (re)construction. 35

S. Zozor is grateful to the Région Rhône-Alpes (France) for the 36 grant that enabled this work. D. Mateos is a Fellowship holder 37 of SeCyT, UNC. 38

Appendix: Technical details 39

Before detailing a possible practical implementation, we 40 should point out that when two components of a vector 41 are equal, an ambiguity remains when performing the per-42 mutation procedure. Such a situation appears with a prob-43 ability of zero for continuous state iid random sequences, 44 but it can appear in constant or periodic sequences, for 45 instance. To avoid such an ambiguity, Bandt and Pompe 46 proposed to add a small perturbation to the values, which 47 is equivalent to choosing randomly the 'smallest' value be-48 tween two equal values. For instance, in the example of a 49 constant sequence, in this way, the permutation vectors 50 reflect only the behavior of the perturbation, and thus 51 both the permutation entropy and the Lempel-Ziv per-52 mutation complexity are of the noise and not of the signal 53 under analysis. To overcome such a difficulty, we chose 54 here to consider that the 'smallest' of two equal values as 55

the 'oldest' one, as has also been done in the literature. 56 In the example of a constant signal, the sequence of per-57 mutation vectors will be constant, which can then capture the low complexity of the sequence. 59

Conversely, an observed sequence can be corrupted by 60 a low noise. This corrupting noise can hide the complexity 61 of the sequence when the permutation vectors are evalu-62 ated. The example of a constant signal again illustrates 63 such an impact of the noise. To counteract perturbations, 64 a way to denoise or filter the observed sequence can consist 65 of choosing a value $\delta \geq 0$ so that for two components Y(i)66 and Y(j) of a (phase-space) vector, if $|Y(i) - Y(j)| \leq \delta$ 67 then Y(i) and Y(j) are interpreted as equal. In a sense, δ 68 is a radius of confidence in the measured data. If $\delta = 0$, 69 this means that we have perfect confidence in the mea-70 sured data, while for $\delta > 0$ we take into account possible 71 perturbations in the measures. In other words, δ can be 72 chosen to be equal to the resolution of the acquisition. 73

Practically, to evaluated $C^{\pi}_{d,\tau}(X_t)$, and to avoid two 74 passes through the sequence, this can be done recursively, 75 by alternating the calculation of the permutation vectors 76 and the up-dating of the complexity: 77

- Step 0. Construction of the first d-dimensional vector 78 $Y = Y_t^{(d,\tau)}$ and evaluation of the first permu-79 tation vector $\boldsymbol{\Pi}_t = \boldsymbol{\Pi}_{\boldsymbol{Y}}, t = 0$; storage of this 80 permutation vector in a stack, and initialization 81 of the Lempel-Ziv algorithm (implicitly, the first 82 production step). 83
- Step 1. $t \leftarrow t+1$: replacement of **Y** by the new vector of the trajectory, evaluation of the new permutation vector Π_t to be stored in the stack.
- Step 2. Up-dating of the Lempel-Ziv complexity using 87 this permutation vector, and go to step 1. 88

In the case where $\tau = 1$, the evaluation of the permutation 89 vector $\boldsymbol{\Pi}_t$ at time t can be simplified by using $\boldsymbol{\Pi}_{t-1}$. 90 Indeed, in the constructed trajectory vector Y, the first 91 point $X_{\text{out}} = Y(0)$ disappears, the other d-1 components 92 are shifted, and the next point of the scalar sequence X_t 93 appears as the last component of Y. The permutation of 94 component *i* (previously $i+1, i = 1, \ldots, d-1$) changes only 95 if either $X_t \ge Y(i)$ and $X_{out} \le Y(i)$ (the rank decreases) 96 or $X_t < Y(i)$ and $X_{out} > Y(i)$ (the rank increases). This 97 up-dating of the rank can thus be made with d doublet of 98 comparisons (seeking also the rank of the new point X_t). 99

For the Lempel-Ziv complexity, when beginning a new 100 production step, the algorithm of [39] consists of testing all 101 of the letters of the already constructed history as possible 102 pointers of a production step, and retaining the letter that 103 gives the greatest production step: this pointer gives what 104 is then called an *exhaustive* production step. 105

The global recursive algorithm is described in detail 106 by the diagram flow shown in Figure A.1; in this simple 107 case, $\tau = 1$. For $\tau > 1$, the same scheme holds, except 108 that we have to first store the τ permutation vectors, then 109 store the τ vectors \boldsymbol{Y} , let us say $\boldsymbol{Y}_0, \ldots, \boldsymbol{Y}_{\tau-1}$, and use 110 both $\boldsymbol{Y}_{t \mod \tau}$ and $\boldsymbol{R}_{t-\tau}$ to recursively evaluate $\boldsymbol{\Pi}_t$. For a 111 non-zero radius of confidence, in the algorithm described 112 in Figure A.1, x > y (and respectively, $x \ge y$) is then 113

58

84

85

86



Lempel–Ziv complexity

Bandt–Pompe permutation

Fig. A.1. Diagram flow of the algorithm evaluating the Lempel-Ziv permutation complexity $C_{d,\tau}^{\pi}$ for a scalar sequence. In this diagram, $\tau = 1$ (see text for the extension to any τ), and the size of the sequence is denoted as T. w_e marks when a word is exhaustive or not, l is the beginning of an exhaustive word, j is the tested pointer, and k_m is the size of the current exhaustive word [35,39].

1 replaced by $x > y + \delta$ (respectively, $x \ge y + \delta$) and x < y2 (respectively, $x \le y$) by $x < y - \delta$ (respectively, $x \le y - \delta$).

Note that there are various fast algorithms that rank a vector [55,56]. In general, these work by recursively partitioning the points to be ranked in a partially ordered manner (through a tree), performing a brute-force sorting in the last partitions, and coming back to the overall ensemble. In general, the computational cost is in $O(d \log d)$, instead of $O(d^2)$ for a totally brute force method. Such 9 approaches should be used in our algorithm, using the 10 partitions at step t - 1 to determine that at step t, ex-11 pecting a computational cost in $O(\log d)$ instead of d. 12 However, in practice, the Bandt-Pompe entropy (and here 13 the Lempel-Ziv permutation complexity) is studied in low 14 dimensions, so that the computational cost of a brute force 15 approach is relatively close to that of fast approaches. 16

Page 12 of 12

Thus, we will not go deeper into such possible improve-1 ments of the proposed algorithm. 2

- Finally, note that contrary to the permutation entropy, 3
- the Lempel-Ziv complexities can be evaluated online, i.e., 4
- up-dated acquisition by acquisition. 5

6 References

- 1. S. Zozor, O. Blanc, V. Jacquemet, N. Virag, J.M. Vesin, 7 E. Pruvot, L. Kappenberger, C. Henriquez, IEEE Trans. 8 9 Biomed. Eng. 50, 412 (2003)
- 10 2.S.A. Kauffman, The Origins of Order: Self-Organization 11 and Selection in Evolution, 1st edn. (Oxford University Press, Oxford, 1993) 12
- 3. G.W. Botteron, J.M. Smith, IEEE Trans. Biomed. Eng. 13 **42**, 548 (1995) 14
- 4. F.F. Ferreira, G. Francisco, B.S. Machado, 15 Ρ. 16 Muruganandam, Physica A 321, 619 (2003)
- 17 5. W.B. Arthur, Am. Econ. Rev. 84, 406 (1994)
- 6. M. Rajković, Physica A 287, 383 (2000) 18
- 19 7. M. Rajković, Z. Mihailović, Physica A **325**, 40 (2003)
- C. Cysarz, H. Bettermann, P.V. Leeuwen, Am. J. Physiol. 20 8. 21 Heart Circ. Physiol. 278, H2163 (2000)
- R.Q. Quiroga, J. Arnhold, K. Lehnertz, P. Grassberger, 22 9. Phys. Rev. E 62, 8380 (2000) 23
- 24 10. X.S. Zhang, Y.S. Zhu, N.V. Thakor, Z.Z. Wang, IEEE Trans. Biomed. Eng. 46, 548 (1999) 25
- 11. G.A. Darbellay, D. Wuertz, Physica A 287, 429 (2000) 26
- 12. W. Ebeling, L. Molgedey, J. Kurths, U. Schwarz, Entropy, 27 complexity, predictability and data analysis of time series 28 and letter sequences, in *Theory of Disaster*, edited by A. 29 30 Bundle, H.-J. Schellnhuber (Springer Verlag, Berlin, 2000)
- 31 13. M.E. Torres, L.G. Gamero, Physica A **286**, 457 (2000)
- T.H. Evrett, J.R. Moorman, L.C. Kok, J.G. Akar, D.E. 32 14. Haines, Circulation 103, 2857 (2001) 33
- 15. N. Radhakrishnan, IEEE Trans. Biomed. Eng. 49, 1371 34 (2002)35
- 36 16. J. Szczepański, J.M. Amigó, E. Wajnryb, M.V. Sanchez-Vives, Network: Computation in Neural Systems 14, 335 37 (2003)38
- S. Zozor, P. Ravier, O. Buttelli, Physica A 345, 285 (2005) 39 17.
- J. Beirlant, E.J. Dudewicz, L. Györfi, E.C. van der Meulen, 40 18. 41 Int. J. Math. Stat. Sci. 6, 17 (1997)
- 19. N. Leonenko, L. Pronzato, V. Savani, Ann. Statist. 36, 42 2153 (2008) 43
- 20. P.O. Amblard, S. Zozor, O. Michel, A.M. Cuculescu, On 44 the estimation of the entropy using kth nearest neigh-45 46 bours, in International Conference on Mathematics in Signal Processing (IMA'08) (Circenter, United Kingdom, 47 2008)48
- T. Schürmann, P. Grassberger, Chaos 6, 414 (1996) 49 21.
- T. Schürmann, J. Phys. A 36, L295 (2004) 50 22.
- 51 23.M. Rosenblatt, Ann. Math. Statist. 27, 832 (1956)
- 24. E. Parzen, Ann. Math. Statist. 33, 1065 (1962) 52
- 25. C. Bandt, B. Pompe, Phys. Rev. Lett. 88, 174102 (2002) 53
- 26. K. Keller, M. Sinn, Physica A 356, 114 (2005) 54
- 55 27.C. Bian, C. Qin, Q.D.Y. Ma, Q. Shen, Phys. Rev. E 85, 56 021606(2012)
- 28.B. Fadlallah, B. Chen, A. Keil, J. Príncipe, Phys. Rev. E 57 58 87, 022911 (2013)

- 29. P. Grassberger, I. Procaccia, Physica D 9, 189 (1983)
- 30. A. Lasota, M.C. Mackey, Chaos, Fractals, and Noise; Stochastic Aspects of Dynamics, Applied Mathematical Sciences 97, 2nd edn. (Springer Verlag, New York, 1994)

59

60

61

62

63

64

65

66

67

68

69

70

71

72

73

74

75

76

77

78

79

80

81

82

83

84

85

86

87

88

89

90

91

92

93

94

95

96

97

98

99

109

120

- 31. J.C. Robinson, Dimensions, Embeddings, and Attractors (Cambridge University Press, Cambdrige, UK, 2011)
- 32. F. Takens, Detecting strange attractors in turbulence, in Dynamical Systems and Turbulence of Lecture Notes in Mathematics, edited by D. Rand, L.S. Young (Springer Verlag, Warwick, 1981), Vol. 898, pp. 366-383
- 33. T.M. Cover, J.A. Thomas, Elements of Information Theory, 2nd edn. (John Wiley & Sons, Hoboken, New Jersey, 2006)
- 34. P. Gács, J.T. Tromp, P.M.B. Vitányi, IEEE Trans. Inf. Theory 47, 2443 (2001)
- 35.A. Lempel, J. Ziv, IEEE Trans. Inf. Theory 22, 75 (1976)
- 36. J. Ziv, A. Lempel, IEEE Trans. Inf. Theory 23, 337 (1977)
- 37. T.A. Welch, IEEE Trans. Comput. 17, 8 (1984)
- 38. G. Hansel, Lecture Notes in Computer Science (Electronic Dictionaries and Automata in Computational Linguistics) **377**. 51 (1989)
- 39. F. Kaspar, H.G. Schuster, Phys. Rev. A 36, 842 (1987)
- 40. J.M. Amigó, L. Kocarev, J. Szczepanski, Phys. Lett. A 355, 27 (2006)
- 41. J.M. Amigó, S. Zambrano, M.A.F. Sanjuán, Europhys. Lett. 79, 50001 (2007)
- 42. J.M. Amigó, Permutation Complexity in Dynamical Systems (Springer Verlag, Heidelberg, 2010)
- 43. O.A. Rosso, L.C. Carpi, P.M. Saco, M. Gómez Ravetti, Physica A **391**, 42 (2012)
- 44. O.A. Rosso, F. Olivares, L. Zunino, L.D. Micco, A.L.L. Aquino, A. Plastino, H.A. Larrondo, Eur. Phys. J. B 86, 116(2013)
- 45. A. Gersho, R.M. Gray, Vector quantization and signal compression (Kluwer, Boston, 1992)
- W. Chou, P.W. Wong, R.M. Gray, IEEE Trans. Inf. 46. Theory **35**, 784 (1989)
- 47. K.T. Alligood, T.D. Sauer, J.A. Yorke, Chaos: An Introduction to Dynamical Systems (Springer Verlag, New-York, 1996)
- 48. S.H. Strogatz, Nonlinear Dynamics and Chaos (Westview Press (Perseus Books Publishing), Cambridge, MA, USA, 100 1994)101
- 49. D. Fournier-Prunaret, R. López-Ruiz, A. Taha, Route 102 to chaos in three-dimensional maps of logistic type, in 103 Proceedings of the European Conference on Iteration 104 (ECIT '04) (Batschuns, Austria, Theory Grazer 105 Mathematische Berichte, Institut für Matematik Karl-106 Franzens Universität Graz, 2004), Vol. 350, pp. 82–95 107
- 50. R. López-Ruiz, D. Fournier-Prunaret, Math. Biosci. Eng. 108 1, 307 (2004)
- 51. O.A. Rosso, S. Blanco, A. Rabinowicz, Signal Process. 83, 110 1275 (2003) 111
- 52. O.A. Rosso, M.T. Martín, A. Figliola, K. Keller, A. 112 Plastino, J. Neurosci. Methods 153, 163 (2006) 113
- 53. M. Pereyra, P.W. Lamberti, O.A. Rosso, Physica A 379, 114 122(2007)115
- 54. H. Gaustaut, R.J. Broughton, Epileptic Seizure: Clinical 116 and Electrographic Feature, Diagnosis and Treatment 117 (Charles C. Thomas Publisher, Springfield, IL, 1972) 118 119
- 55. C.A.R. Hoare, Comput. J. 5, 10 (1962)
- 56. R. Sedgewick, Programming Techniques 21, 847 (1978)