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ARTICLE

Optimisation of a Venturi shaped structure around a vertical axis tidal turbine

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ABSTRACT

Tidal energy is one of the worlds most predicable renewable energy sources and therefore holds great potential to be a valuable building block for the decarbonisation of the electricity production. This paper focuses on a vertical axis tidal turbine utilising a Venturi shaped outside structure (shroud) to accelerate the flow speed at the turbine. This concept is known as Davidson Hill Venturi (DHV) turbine. By constructing the nozzle and diffuser using hydrofoils, initial demonstrations indicate increased system efficiency. However, due to the potential number of hydrofoil geometric and structural variations, only a general description of the location of the hydrofoils is provided in order to facilitate modelling while allowing for future geometric variations to be devised. The conducted investigations focus on the influence of the nozzle and diffuser sections as the main geometry variations, identifying the length component in the orthogonal direction as the dominant parameter. A combined variation indicates higher velocity values are connected with larger forces, which must be supported by the devices structure. Slight improvements from the provided reference geometry were found as variations of hydrofoil placement and spacing were simulated. Thus, the main conclusion is that the reference geometry needs only small adaptations, which will be investigated further in a 3D-simulation study, including turbine interaction and rotation.

KEYWORDS

Tidal turbine; vertical axis; optimisation; Davidson Hill Venturi (DHV); ANSYS-CFX; SpaceClaim; Python;

1. Introduction

The kinetic energy of tidal currents are a renewable energy source, which is highly predicable and can provide a key component for the future fully de-carbonised electrical energy production (Bahaj, 2013; Khojasteh et al., 2022). A wide range of turbine concepts are currently investigated (Roshanmanesh, Hayati, & Papaelias, 2020). For example, they can be classified by the orientation of their rotation axis in vertical and horizontal axis turbines (Behrouzi, Nakisa, Maimun, & Ahmed, 2016; Khan, Bhuyan, Iqbal, & Quaiocoe, 2009) or also in floating and bottom mounted devices. A wide range of different blades are used including open centre horizontal axis turbines (Belloni, Willden, & Houlby, 2017; Borg, Xiao, Allsop, Incecik, & Peyrard, 2020, 2021),

33 twisted blades (Mosbahi, Elgasri, Lajnef, Mosbahi, & Driss, 2021) or deflector blades
34 (Patel & Patel, 2022). Blades can be either fixed, provide a variable pitch (Gu, Lin,
35 Xu, Liu, & Li, 2018; H. Liu, Li, Lin, Li, & Gu, 2020) and/or yaw angle (Modali,
36 Vinod, & Banerjee, 2021; qi Wang, Xu, qing Zhu, & Wang, 2018) or even change the
37 shape (Pisetta, Le Mestre, & Viola, 2022) to further improve the efficiency. Further
38 classifications can be made using the generator, nacelle types, supporting structure
39 and generators (Roshanmanesh et al., 2020).

40 In addition to the turbine structure itself, various methods are available to increase
41 the flow speed at the turbine and therewith improve the overall performance of the
42 turbine. For example, Z. Liu, Wang, Shi, and Qu (2019) investigates four additional
43 hydrofoils placed as a guide-vane diffuser using ANSYS-Fluent. Hua-Ming, Xiao-Kun,
44 Lin, Lu-Qiong, and Qiao-Rui (2020) used the software STAR-CCM+ to investigate
45 different ducts in a current flow limited by close side banks. This investigation was
46 conducted as a 3D simulation but without a turbine present. A three bladed vertical
47 axis turbine was simulated by El-Sawy, Shehata, Elbatran, and Tawfiq (2022) under
48 river conditions. An integration in a floating configuration is also possible (Hardisty,
49 2008). More commonly used for free stream turbines is a diffusor similar to the ones
50 for wind turbines (Arumugam, Ramalingam, & Bhaganagar, 2021; Noorollahi, Ghan-
51 bari, & Tahani, 2020; Nunes, Brasil Junior, & Oliveira, 2020). Those structures can
52 be optimised for on single flow direction or allow a bi-directional usage (Fleming &
53 Willden, 2016), which enables a fixed installation. Such ducted geometries can be very
54 simple structures but also include a variety of complex shapes and combined struc-
55 tures, as investigated by Huang et al. (2022). This paper investigates a Venturi shaped
56 structure, which assembles multiple hydrofoils in order to accelerate the available flow
57 through the vertical axis turbine (Kirke, 2011). A summary of shapes and the influence
58 of diffusors and horizontal axis turbines can be found in Nunes et al. (2020). Walker
59 and Thies (2021) indicates the percentage of failed tidal turbines, are higher hence the
60 loads due to the increased velocity are higher. Nevertheless, they indicated that with
61 improved materials and optimised concepts the potential of ducted turbines can be
62 significant.

63 Ongoing research challenges can be found in the simulation of the turbine itself,
64 which can be simplified by various approximations (Baratchi, Jeans, & Gerber, 2020).
65 Ke, Wen-Quan, and Yan (2020) compares bares a horizontal axis turbine with a range
66 of diffusors under a distributed inlet velocity as well as the interaction of multiple
67 turbines in an array. Multiple vertical axis turbine arrays are investigated by Sun,
68 Ji, Zhang, Li, and Wang (2021). Ahmed, Apsley, Afgan, Stallard, and Stansby (2017)
69 studied the influence of the velocity distribution and compared it with field data and
70 Badshah, Badshah, and Kadir (2018) used a Fluid-Structure-Analysis to quantify the
71 influence of a realistic velocity distribution in relation to a homogeneous approach.

72 The overall design of tidal turbines is strongly reliant on numerical simulations
73 (Nachtane, Tarfaoui, Goda, & Rouway, 2020). Additional, experimental investigations
74 allow specific measurement data to be compared leading to improvements within the
75 numerical models. A good example can be found in the paper by Badoe et al. (2022),
76 who simulated up to three tidal turbines under a comparable complex flow conditions
77 representing the unique FloWave Ocean Energy Research Facility (part of the Uni-
78 versity of Edinburgh). This facility provides a raisable floor for the dry installation of
79 not only tidal turbine model and ensures a highly repeatable flow condition of up to
80 1.6 m/s rotatable by 360° due to the circular arrangement of the flow drives. W. Liu et
81 al. (2022) validated their numerical model of a diffuser-augmented tidal turbine with
82 a towing tank experiment. Feng et al. (2022) used a flume to compare the wake inter-

83 action between two ducted horizontal axis turbines. Obviously, the next step for the
84 validation is the deployment of the optimised structure combined with a competitive
85 measurement system to reduce the uncertainties in the validation.

86 The presented research work uses a similar approach to Maduka and Li (2021), who
87 investigated a ducted vertical turbine using a reduced 2D-approach and simplified the
88 turbine with an actuator disc approach in the numerical software OpenFOAM. In
89 contrast to this, the current paper focuses on an application of a ducted vertical axis
90 tidal turbine and introduces a basic description of the Venturi shaped shroud of the
91 Davidson Hill Venturi (DHV) turbine. The supporting structure of the turbine device
92 is in part assembled using hydrofoils arranged either side of the inlet, parallel to the
93 axis of the turbine, improving the acceleration of the flow at the turbine. All numerical
94 simulations are conducted with a commercial code ANSYS-CFX. The geometry gen-
95 eration is provided in the SpaceClaim specific IronPython as well as the open Python
96 code (Gabl, Burchell, Hill, & Ingram, 2022). This ensures that the conducted varia-
97 tions can be reproduced and expanded with any solver. The methodology is described
98 Section 2, which includes a reference geometry (Sec. 2.4) as well as the verification of
99 the 2D-numerical simulation (Sec. 2.6). A key output of the paper is the variation of
100 the main geometry parameters as well as the embedded hydrofoils (Sec. 3). Section 4
101 provides a discussion of the chosen methodological approach. The investigation is con-
102 ducted with a reduced complexity of numerical simulations and allow the identification
103 of the optimum shape, which will be further numerically tested in a fully 3D-setup
104 to bring this concept closer to a commercial deployment producing fully predicable
105 renewable energy from river, canal and tidal flows.

106 **2. Materials and methods**

107 ***2.1. Overview***

108 A standard tidal turbine requires the blades, generator and support structure. One of
109 the unique features of the Davidson Hill Venturi (DHV) turbine is the outer struc-
110 ture, which contains multiple hydrofoils forming a Venturi channel (DHV Turbines
111 Ltd, 2022; Kirke, 2011). This concept was invented by Aaron Davidson and Craig Hill
112 and is commercialised by DHV Turbines Ltd. The smallest cross sectional area occurs
113 at the turbine position, which augments the incoming flow resulting in increased
114 fluid velocities acting on the turbine. As part of the ongoing research work, a gener-
115 alised geometry description is suggested, providing the current conducted parameter
116 variation as well as future additional work. The conducted numerical simulations are
117 limited to a 2D-approach and focuses on the guide structure only, neglecting interac-
118 tion with the turbine. Section 2.2 introduces the solver and basic numerical settings.
119 A local coordinate system is introduced in Section 2.3 for the geometry description,
120 shown as the reference geometry in Section 2.4. The wide range of numerical results
121 are analysed, with primary focus on the turbine cross section (Sec 2.5). Section 2.6
122 summarises the key aspects of the verification process, namely the size of the fluid
123 domain and the mesh test. All these components are required for the variation of the
124 geometry presented in Section 3 and a discussion of the methodology can be found
125 in Section 4. The optimised geometry will be further numerically investigated and
126 refined, with future experimental investigation and deployment required in order to
127 validate the modelling data.

128 *2.2. Solver and numerical settings*

129 The presented numerical investigation uses the commercial Computational Fluid Dy-
130 namic (CFD) software ANSYS-CFX (version 2020 R2). This software was used for a
131 wide range of numerical studies in specific aspects of tidal turbines. Sun et al. (2022)
132 found that a four bladed vertical axis turbine brings significant advantages in the
133 starting performance based on simulations with CFX. Sun, Ma, Wang, and Li (2019)
134 focused on the fluctuation of the angular speed and also conducted specific validation
135 experiments in the laboratory. Rehman et al. (2021) investigated the wake behind a
136 single horizontal axis turbine utilising this software and Allmark et al. (2020) inte-
137 grated it into the design and blade optimisation of a scaled turbine. In combination
138 with other ANSYS products, further capabilities to couple physics and conducted
139 Fluid-Structure-Interaction (FSI) simulations are attainable, as shown by Badshah,
140 Badshah, and Jan (2020) and Ullah et al. (2019). While a key application area of the
141 software is the investigation and optimisation of turbines (Mulu, Cervantes, Devals,
142 Vu, & Guibault, 2015; Picone, Sinagra, Aricó, & Tucciarelli, 2021), a wide range of
143 further applications can be found (Lee, Seong, & Kang, 2018; Richter et al., 2021).
144 Recent validation experiments include the optimisation of an inlet-outlet structure
145 (Bermúdez et al., 2017) and spillway (Andersson, Andreasson, & Lundström, 2013).
146 The high degree of validation within the numerical solver allows this paper to focus
147 fully on the verification as presented in Section 2.6.

148 All simulations use the assumption of a 2D-approach, which is realised by a single
149 cell size in the vertical z-direction. The size of the fluid domain is chosen based on the
150 verification process presented in the Section 2.6. A constant homogeneous velocity is set
151 at the inlet face. Nominally the inlet speed is set to 1 m/s, however for completeness the
152 effect of speed variation is assessed and can be found in Section 3.9. On the opposing
153 side of the fluid domain, the outlet is set as a pressure outlet with a constant static
154 pressure of 1 bar to ensure that cavitation is not an issue. All other sides of the fluid
155 domain box are set to symmetry with the investigated structure allocated wall (no slip
156 wall option with a roughness setting of smooth wall). The simulated fluid is standard
157 water (density = 997.0 [kg m⁻³]), while isothermal approach is used at a constant
158 temperature of 10°C. No buoyancy is considered.

159 The standard steady state solver computes the Reynolds-Averaged Navier-Stokes
160 (RANS) equations, while the closure problem is solved based on the Shear Stress
161 Transport (SST) turbulence model (F. Menter, 1993, 1994; F. R. Menter, Kuntz, &
162 Langtry, 2003). This turbulence model blends a $k - \omega$ and $K - \epsilon$ turbulence model
163 to combined the advantage of each specific model. The SST model is widely used for
164 a broad range of solver (Cindori, Čajić, Džijan, Juretić, & Kozmar, 2022; He, Zhao,
165 Wan, & Wang, 2022; Sarkar & Savory, 2021; Yi, Wang, Sun, Huang, & Zheng, 2017).
166 Due to the fact that the geometry is varied in a significant way and introduces different
167 flow regimes, the auto time scale function of the solver is used with a time scale factor
168 of 1. For some exemplary cases a modification to 0.1 and 10 was investigated, which
169 resulted in a slower convergence behaviour of the numerical solution. The solver runs
170 in all cases using double precision to reduce the potential influence of rounding errors.
171 Three additional monitor points of velocities in the cross section of the turbine are
172 used to ensure that the solution is converged and can be stopped. Further verification
173 studies can be found in Section 2.6.

174 **2.3. Local coordinate system**

175 A local coordinate system is defined so that the z-axis is coincident to that of the
 176 vertical axis turbine. The positive x-axis is orientated along the main flow direction
 177 resulting in a positive velocity. Thus the x-z-plane is the symmetry plane for the
 178 geometry and the y-axis aligns according to the right handed coordinate system. The
 179 origin for the 2D-dimensional simulations was placed on the bottom part of the axis.

180 **2.4. Reference geometry and parameters**

181 The overall geometry of the DHV turbine includes a large number of variables, with
 182 each having significant impact on the overall efficiency, production cost and mainte-
 183 nance of the turbine. In order to reduce the number of parameters within the design,
 184 the first step focuses on the outer Venturi structure and assumes that the turbine itself
 185 can be optimised in a following assessment. Similar assumptions are made for the sup-
 186 port structure. The previously conducted experimental and numerical investigations
 187 ensured the principal validation of the turbine concept (Kirke, 2006, 2011) as well as a
 188 fair amount of parameter variations. A full investigation of all parameters and combi-
 189 nations would exceed the time frame of any research project. Hence, a limited amount
 190 of combinations are investigated starting from the later described reference geometry.
 191 Nevertheless, a key step is the full parametrisation of the geometry so that the later
 192 gained results can be easily documented, reproduced and expanded upon. The decision
 193 was made to provide this in the form of a Python code, which is available in Gabl et
 194 al. (2022). This also publishes a specific IronPython version, which is allows for the
 195 geometry definition in 3D CAD modelling software SpaceClaim. Therewith a full in-
 196 tegration in the ANSYS WORKBENCH is possible and the fluid solver ANSYS-CFX
 197 was used for this project.

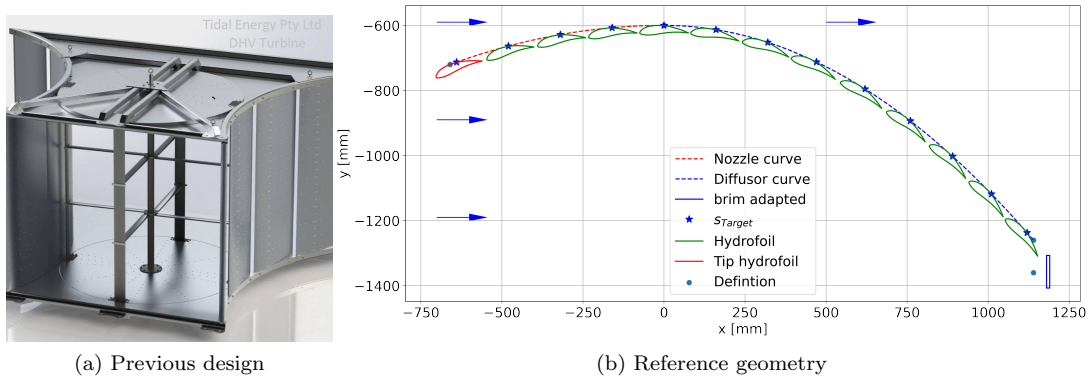


Figure 1. Overview of the previous design (a) provided by DHV Turbines Ltd and the half side of reference geometry (b) and the vertical axis turbine rotating in the origin of the coordinate system

198 For the following geometry investigation an initial reference geometry (Figure 1 (b))
 199 was utilised, based on a previous DHV demonstrator project. These values are pre-
 200 sented in combination with the specific geometry parameters as follows:

- 201 • One of the defining parameters of the housing geometry is the radius of the
 202 vertical axis turbine R_{Turb} . It is chosen to 0.6 m for the reference geometry and
 203 later used to standardise all length of the geometry.
- 204 • Both sides of the Ventri channel are symmetrical along the x-z-plane and split

205 at the y-z-plane (turbine location) in two parts: nozzle and diffusor. As shown
 206 in Figure 1 (b), this shape is defined by three points for the main curves. The
 207 centre one is defined by the turbine with R_{Turb} and the coordinates of the outer
 208 two points are defined by the R_{Front} and L_{Front} as well as R_{Back} and L_{Back} .
 209 In the case of the reference geometry, shown in Figure 1 (b), the combination
 210 of $[R_{Front}, L_{Front}, R_{Back}, L_{Back}]$ equal $[1.2, 1.1, 2.1, 1.9] \cdot R_{Turb}$ was chosen. The
 211 radius describes the distance/length in the y-direction according to the local
 212 coordinate system.

- A bell curve is defined between these points, chosen based on the Alkhabbaz, Yang, Tongphong, and Lee (2022), which includes the work of Khamlaj and Rumpfkeil (2018b, 2018a). Eq. (1) presents the nozzle (with a negative x-value indicated by x_{neg}) and diffusor (x_{pos}) part aligned along the flow direction on the right side of the geometry (Fig. 1 (b)).

$$\begin{aligned}
 y_{neg} &= R_{Turb} - 1 - x_{neg}^2 \cdot (R_{Front} - R_{Turb}) / L_{Front}^2 \\
 y_{pos} &= R_{Turb} - 1 - x_{pos}^2 \cdot (R_{Back} - R_{Turb}) / L_{Back}^2 \\
 \text{with } -1 \cdot L_{Front} &\leq x_{neg} \leq 0 \text{ and } 0 \leq x_{pos} \leq L_{Back}
 \end{aligned} \tag{1}$$

- The hydrofoil profiles are then located to these curved forms. For the reference geometry a GOE 222 (MVA H.33) AIRFOIL (Airfoil Tools, 2022) with a chord length c_{Foil} of 150 mm was chosen. The local reference point for the hydrofoil is placed at the mid point of the chord length at the maximum height (Gabl et al., 2022). This ensures that the hydrofoil does not interfere with the cross sectional area of the turbine. Downstream of the central position on each side, the bell curves split into segments with a constant chord length of s_{Target} equal 160 mm. The implementation is shown in Gabl et al. (2022) and results in an integer number of hydrofoils for each side. As shown in Figure 1 (b), the individual reference point for the last hydrofoil is still inside of the definition point but the hydrofoil reaches outside of the curve. Limiting the design to one single size hydrofoil with constant chord spacing might swallow smaller geometry changes but seems to be more practical than variations of the hydrofoil sizes. Variable hydrofoil size would result in significant higher cost of the structure. Furthermore, each hydrofoil is rotated based on the differential of the curve to ensure that the tangential direction is similar to the bell shape. Additional variables, which allows for a constant changes to the angle of attack of all hydrofoils as well as a specific modification of the first hydrofoil (red foil in Fig. 1 (b)) in the flow direction, are included in the Python code. The additional angle of the orientation of the first hydrofoil in the flow direction is investigated in Section 16.
- In addition to the hydrofoils, a brim is included in the reference geometry. The height (y-direction) is 100 mm and the thickness is 10 mm. These values were primarily chosen to be small but still clearly have an influence in the mesh generation. A specific investigation of the size of the additional brim is presented in Section 3.3.

238 The last two noteworthy aspects of the geometry definition, namely the hydrofoils as
 239 well as the brim, are clear differences to the previous tested design of the DHV turbine.
 240 In some aspects a smaller hydrofoil has advantages compared to the larger profiles
 241 previously deployed. Smaller shapes can be extruded (ideal for a mass production)
 242 instead of welding single pre-cut parts. In the original configuration of the DHV turbine
 243 a brim structure was absent from the design. A wide range of investigations show the

244 efficiency of such an additional blockage element for wind turbines (Alkhabbaz et al.,
 245 2022; Arumugam et al., 2021; Khamlaj & Rumpfkeil, 2018b, 2018a; Nunes et al., 2020).
 246 To test the influence in the case of a DHV turbine a small brim section is initially
 247 added on both sides at the end of the bell shape. Its upstream edge, closest to the
 248 symmetry plane, is similarly placed to the reference points of the hydrofoils but, with
 249 half the chord distance Δs_{Brim} equal to 80 mm. The variation of the geometry is
 250 discussed and the results are presented in Section 3.

251 **2.5. Post-Processing**

252 The post-processing of the numerical results focuses on the velocity distribution rep-
 253 resented by the variable vel . The actual flow direction is not considered and further
 254 detailed analysis can be used for the evaluation of the full geometry (Gabl, Achleitner,
 255 Neuner, & Aufleger, 2014). Three different control sections are defined. One full cut
 256 through the fluid domain as a y-z-plane (coordinate system is described in Sec 2.3)
 257 with an x-value of -3 m for the inflow conditions and another at x=6 m for the outflow.
 258 The main control cross section is set in the centre of the turbine axis orientated in the
 259 y-z-direction but limited to $R_{Turb}+10$ mm to each side. This allows specific analyse
 260 the flow between the two structures at the location of the turbine. Values referencing
 261 this cross section are indicated with the index c . The other two cross sections are used
 262 for the verification of the inflow and outflow and are not specifically reported in this
 263 paper. For the main analysis the following parameters are used:

- 264 • area-averaged velocity at the turbine cross section \overline{vel}_c using the ANSYS CFD-
 265 Post function $areaAve(vel)$
- 266 • total forces in x-direction F_x based on the calculation $sum(Force X)$
- 267 • kinetic energy flux coefficient α_c based on the Eq. (2)
- 268 • extreme values (minimum and maximum) of vel_c

269 The first two values, namely \overline{vel}_c and F_x , are the main parameters for the analysis
 270 presented in Section 3. The average of the full cross section was deliberately used for
 271 the velocity value, while the velocity distribution was evaluated with the help of the
 272 α_c -value. This coefficient, which is used to correct the standard Bernoulli's equation
 273 for the real kinetic energy $E_{kin,real}$, is typically larger than the homogeneous approach
 274 for the theoretical value $E_{kin,theo}$. Fully developed laminar flow in a pipe cross section
 275 results in a parabolic velocity distribution and a α -value of 2. Typical turbulent flows
 276 are in the range of 1.2 (Gabl et al., 2014; Gabl & Righetti, 2018; Ward-Smith, 1980).
 277 The α -value is calculated based on the following equation:

$$\alpha_c = \frac{E_{kin,real}}{E_{kin,theo}} = \frac{1}{A} \cdot \int_A \left(\frac{vel}{vel_m} \right)^3 dA \quad (2)$$

278 This equation can be included in ANSYS CFD-Post using the CFX Expression
 279 Language, or CEL, as an expression with the location $Plane2$ references the cross
 280 section at the turbine:

```
281 alphaC = areaInt(vel^3)@Plane2/((areaAve(vel)@Plane2)^3*area()@Plane2)
```

282 In addition to this global evaluation value, the minimum and maximum value of the
 283 cross section is analysed to find the range of the flow velocities at the turbine. Note
 284 that the α and extreme values are only shown if they are particular interesting for

285 the analysis. Section 3 includes example analysis of velocity contour plots. Those plots
 286 always show the x-y-plane and the influence of the z-direction can be neglected due
 287 to the a 2D nature of the simulation. If not otherwise stated, the colour bar for the
 288 velocity plots are limited from 0 to 3.5 m/s, which allows ease of comparison between
 289 the various geometry investigations provided.

290 2.6. Verification

291 The verification process includes an investigation of the fluid domain size based on a
 292 first pass approximation of the computation mesh, which was subsequently modified
 293 in a second step. A compromise between the distance of the boundary conditions to
 294 the main investigated part as well as the overall computational cost of the simulations
 295 is targeted, while ensuring that the resulting outcomes are independent of the chosen
 296 mesh input variable. For all presented simulations, the reference geometry was investi-
 297 gated as introduced in Section 2.4. The derived variables and analyses are presented
 298 in Section 2.5.

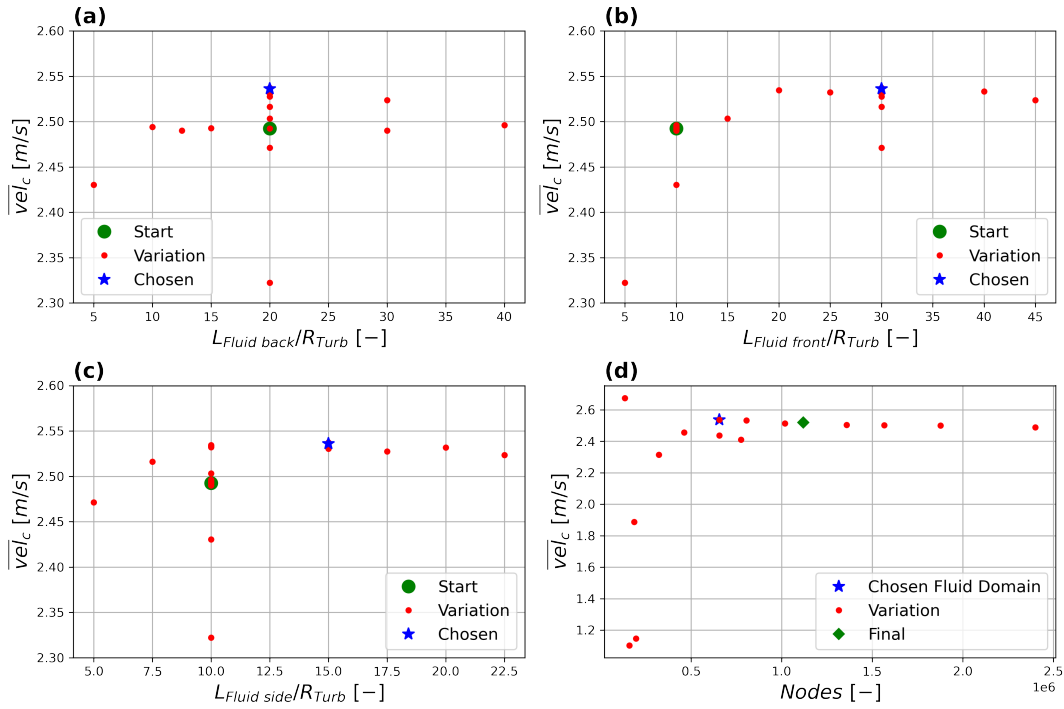


Figure 2. Summary of the verification process with a specific mark of the start as well as the final chosen value — (a-c) present the full set of fluid domain changes but as a result of to the individual investigated length — (d) comparison of the different investigated meshes.

299 An initial volume for the fluid domain was chosen based on a rectangular cuboid with
 300 three length (all referenced from the origin of the local coordinate system, Section 2.3)
 301 and a thickness. The near side length $L_{Fluid\ front}$ defines the dimension in the negative
 302 x-direction and was initial chosen to be equal $10 \cdot R_{Turb}$. While, both sides were
 303 extended out to $L_{Fluid\ side}$ of $10 \cdot R_{Turb}$ from the symmetry plane, resulting in a total
 304 width of $20 \cdot R_{Turb}$. The downstream part of the fluid domain was expanded with
 305 a length of $L_{Fluid\ back}=20 \cdot R_{Turb}$. This values is highlighted as a starting value in
 306 the graphs provided in Figure 2 (a-c). All three graphs show the same outputs but

307 mapped against the alternative investigated lengths. To start with, the downstream
 308 section of the fluid domain was varied. Figure 2 (a) shows that the range of 10 to
 309 $40 \cdot R_{Turb}$ results in very similar values for the chosen indicator \overline{vel}_c . The upstream
 310 section was chosen to be too short in the first assumption and hence expanded to
 311 $30 \cdot R_{Turb}$. Similarly, the side lengths were expanded to $15 \cdot R_{Turb}$. All dimensions
 312 were deliberately increased to accommodate the expansion of the bell shapes and any
 313 expected modifications of the geometry. The total thickness of the fluid domain was
 314 set to the depth of a single cell of 10 mm, thus resulting in a 2D-simulation. The
 315 thickness (fluid domain in the z-direction) was modified to 5 and 100 mm, which had
 316 no influence on the velocity distribution.

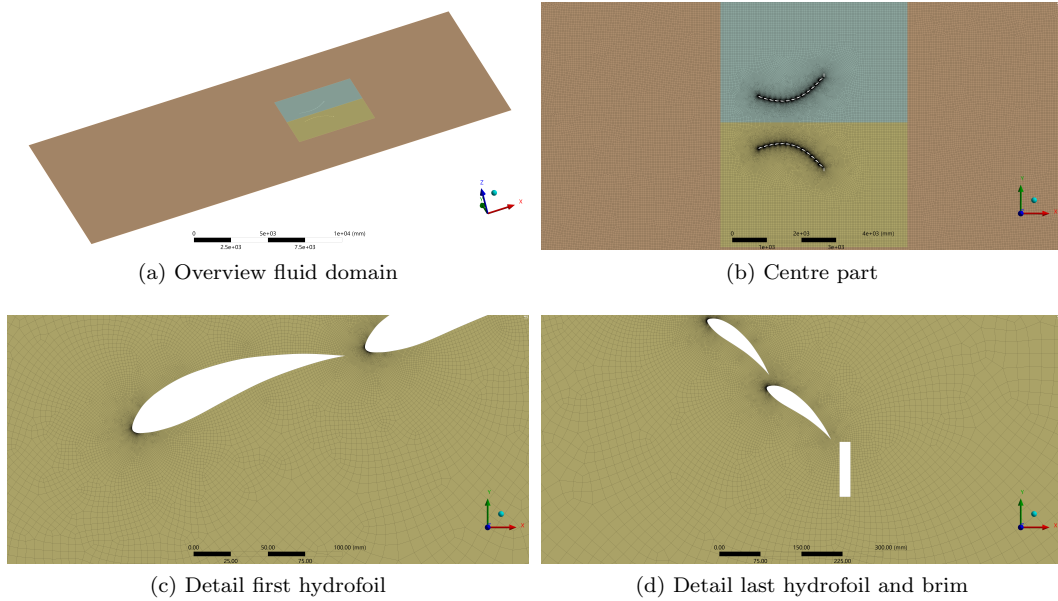


Figure 3. Final meshing used for the 2D-investigation using the reference geometry.

317 Following the variation of the fluid domain, the meshing strategy and resolution was
 318 modified. Figure 2 (d) presents the overview of the conducted mesh test. All meshes
 319 were 2D-dimensional with one cell height in the z-direction. Due to the automated
 320 geometry generation and high degree of variability, only global mesh parameters were
 321 set, which influence the complete fluid domain. The final mesh deploys the following
 322 settings with the chosen values: Element size (50 mm), growth rate (1.05 [-]), max
 323 size (100 mm), defeature size (0.05 mm), curvature min size (0.05 mm) and curvature
 324 normal angle (0.5 deg). The MultiZone method was used to fill the fluid domain with
 325 hexagonal shaped cells and in total the mesh included over 1.1. million nodes. Due to
 326 the fact that this meshing method is limited to a single core, mesh generation took a
 327 comparably long time and the decision was made to split the complete fluid domain
 328 in three sub domains. As shown in Figure 3, each side of the geometry is cut out with
 329 a total length in x-direction of $9 \cdot R_{Turb}$ (3 upstream and $6 \cdot R_{Turb}$ downstream of the
 330 turbine axis) and individual width of $6 \cdot R_{Turb}$. This results in an interface boundary
 331 but allows to distribute the mesh generation to more processors.

332 As a final assessment the selected mesh was run within the fluid domain, which was
 333 expanded by the factor 1.5 in the x-y-plane. For this expanded case the area-average
 334 value of the velocities at the turbine were smaller by 0.01277 m/s representing 0.504%
 335 of to the original average value.

336 **3. Results**

337 **3.1. Overview**

338 The presented results can be grouped into four different categories: (a) the shape of
 339 the bell curve described by the definition points, (b) the angle of the first hydrofoil,
 340 (c) the shape and size of the hydrofoils and (d) the inflow velocity, which is presented
 341 in the Section 3.9.

342 An overview of the major variation grid for the nozzle and diffuser part is shown
 343 in Figure 4. The reference value, presented in Section 2.4, is marked. This figure also
 344 provides the tangential angle of the hydrofoil to the outside of the bell curve. In
 345 addition to the colour bar, two exemplary values are marked with 10 and 55° as a
 346 reference. The latter presented results also include further refinement, however, this
 347 is not included in this overview. Both parts are varied separately, while the other
 348 side is kept at the reference geometry. Section 3.2 presents the variation of the nozzle
 349 geometry and the investigation of the diffuser is combined with the extension of the
 350 brim in Section 3.3. For all those variations the radius of the vertical turbine R_{Turb}
 351 is kept constant at 600 mm. Section 3.4 provides the investigation of this parameter in
 352 two ways and compliments the geometry variation. The narrowing of the structures
 353 geometry resulted in an interesting result, which is further investigated in Section 3.5
 354 by scaling both sides of the structure with a constant factor. In addition, the angle of
 355 the first hydrofoil in the flow direction is investigated in Section 3.6.

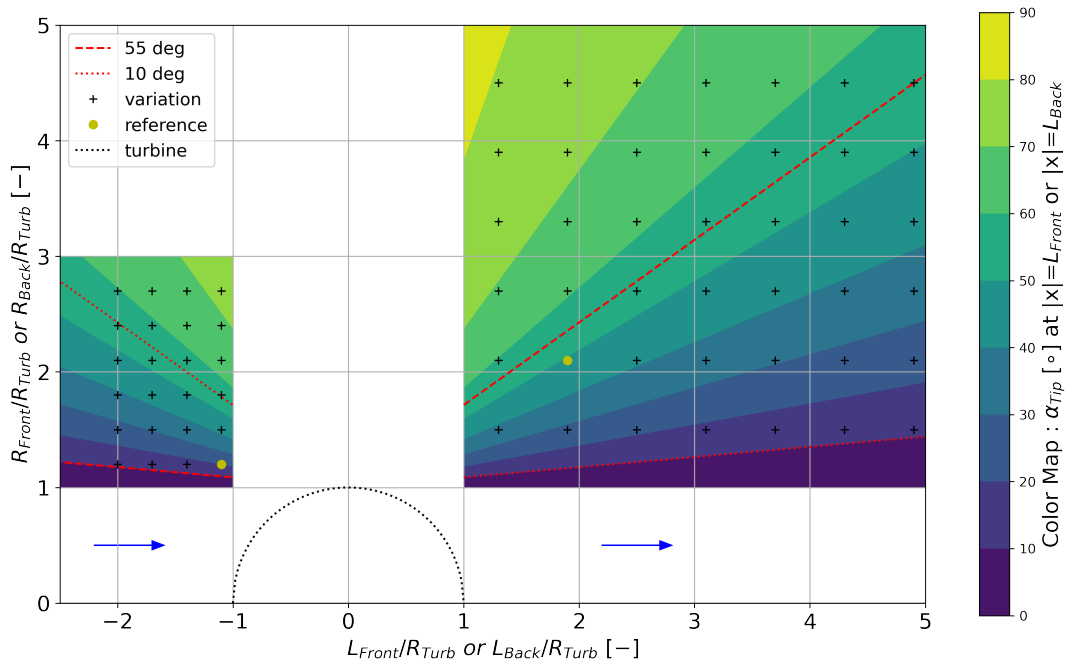


Figure 4. Overview of the main grid for the geometry variation normalised by the turbine radius R_{Turb} with a colour map of the angle α_{Tip} of the first (for the nozzle) or last hydrofoil using this definition point. A value of $\alpha_{Tip} = 0^\circ$ indicates orientation along the x-axis and 90° along the y-axis. The reference value is described in Section 2.4.

356 The second main group of the variation is the investigation of the hydrofoils itself.
 357 Section 3.7 scales the length of the individual hydrofoil as well as the distance between
 358 the standard foil. Up to this point, all simulations are conducted with a hydrofoil profile

359 based on the GOE 222 Airfoil (Airfoil Tools, 2022) and only within Section 3.8 is the
 360 hydrofoil geometry changed. Due to various factors, it was assumed that only one
 361 single type of hydrofoil is used at any time in the variation. A discussion of the used
 362 methodological approach is provided in Section 4.

363 3.2. Variation of the nozzle

364 The first variation of the geometry is to alter the the nozzle of the Venturi structure.
 365 For those cases, the downstream part is fixed with reference values, as described in
 366 Section 2.4. The results for this variation are provided in Figure 5, which presents the
 367 area-average velocity value \overline{vel}_c of the turbine cross section in the upper row and the
 368 total forces F_x in the x-direction. Note that the latter is normalised by the reference
 369 value. Each graph shows the same numerical results but, the left column presents it in
 370 relation to the standardised value of the inlet radius R_{Front} and the length L_{Front} in the
 371 right column of Figure 5. A variation grid of 180 mm ($0.3 \cdot R_{Turb}$) was investigated for
 372 both variables defining the inlet bell. In addition, specific addition points were run to
 373 further investigate specific combinations and enrich specific parts with a high gradient.

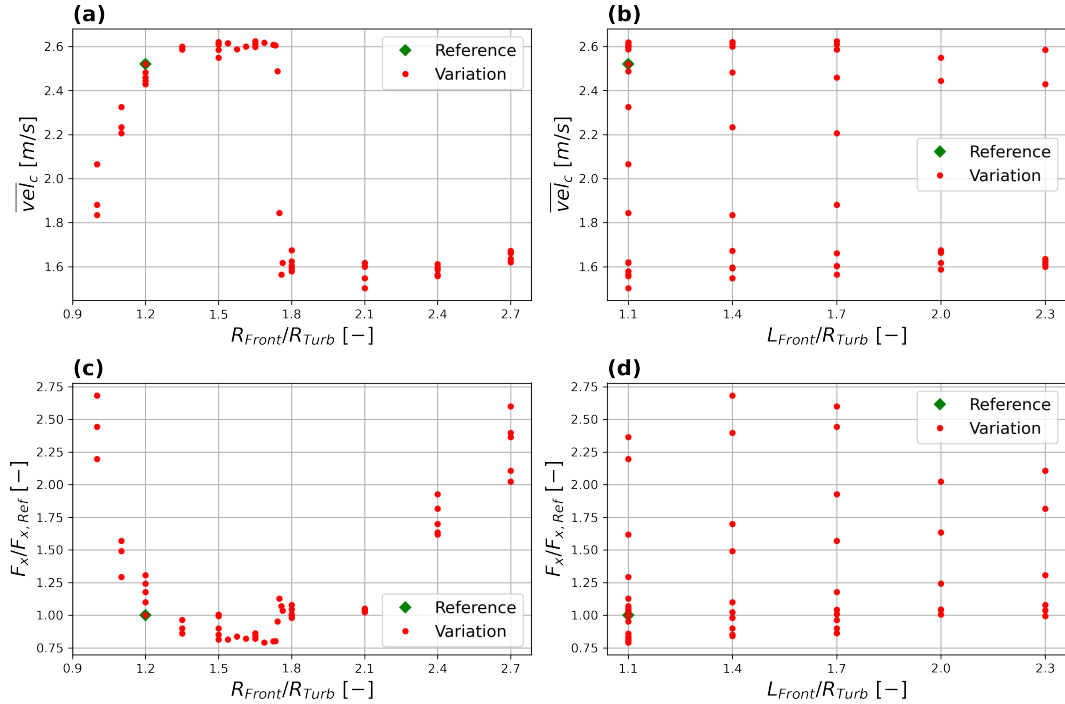


Figure 5. Variation of the nozzle with a fixed diffuser $[R_{Back}, L_{Back}] = [2.1, 1.9] \cdot R_{Turb}$ — area-averaged velocity at the turbine cross section \overline{vel}_c depending on the opening radius R_{Front} (a) and sorted by the length L_{Front} (b) — Forces in x-direction F_x standardised by the reference value (c-d) — all length standardised by the turbine radius R_{Turb} .

374 By varying R_{Front} and length L_{Front} and comparing the outcomes, the dimension
 375 in the orthogonal direction to the main flow direction, namely R_{Front} , is shown to
 376 have a significant influence. Figure 5 (a) indicates that an increasing opening width
 377 R_{Front} results in an increasing area-averaged velocity at the turbine cross section \overline{vel}_c
 378 up to a certain level, which is followed by a drop in the efficiency of the inlet structure.
 379 This specific section is refined down to a level of 5 mm steps for the R_{Front} . The jump

380 occurs for a R_{Front} value in the range of $1.74583 \cdot R_{Turb}$ (between 1045 and 1050 mm).
 381 In reality, such close points would not be advisable as construction tolerances could
 382 potentially have a significant influence of the turbines operation. Therefore, a higher
 383 level can be achieved starting with $1.5 \cdot R_{Turb}$ in order that the aforementioned jump
 384 can be avoided.

385 The total force in the x-direction F_x acting on the structure increases for both
 386 smaller and larger values with a minimum around $R_{Front}=1.5 \cdot R_{Turb}$. Consequently,
 387 both values would indicate that the R_{Front} could be slightly increased.

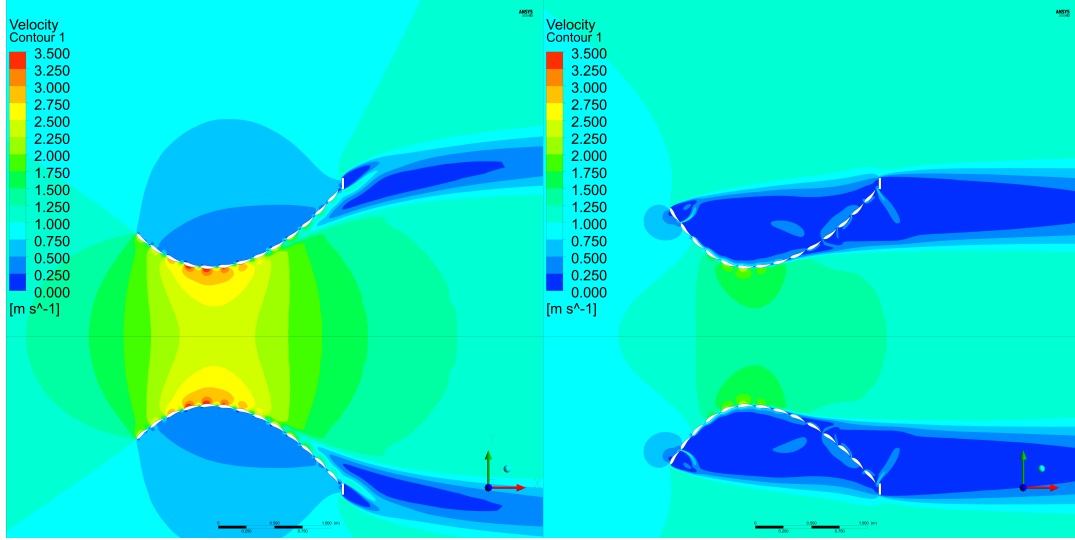


Figure 6. Velocity plot comparing $R_{Front}=1.5 \cdot R_{Turb}$ (left) and $R_{Front}=1.8 \cdot R_{Turb}$ (right) with a fixed $L_{Front}=1.1 \cdot R_{Turb}$.

388 As shown in Figure 6, a velocity drop occurs due to recirculation zone changes. A
 389 larger R_{Front} value results in the side structure acting like a blockage and reducing the
 390 flow speed inside to near zero. Consequently, this is not only influenced by the nozzle
 391 but also of the diffuser, including the brim as well as the flow speed. In the case of the
 392 larger R_{Front} (Fig. 6 right), a leakage of the flow behind the first hydrofoil is obvious,
 393 which indicates that this combination does not guide the flow as intended.

394 The main conclusion, which can be drawn based on this variation assessment, is that
 395 the inlet dimension of the reference geometry is in a satisfactory range. The parameter
 396 R_{Front} is critical but an expansion close to or even beyond the outlet R_{Back} results in
 397 no further improvements. Such a drop must be prevented and the final design should
 398 tolerate some uncertainty to stay in the correct flow regime.

399 **3.3. Variation of the diffuser including brim**

400 The variation of the nozzle part of the Venturi structure indicated that the diffuser
 401 section should reach further out in the orthogonal direction as the front part. Conse-
 402 quently, smaller R_{Back} values are excluded and larger steps for the variation of 360 mm
 403 ($0.6 \cdot R_{Turb}$) are investigated. Figure 7 provides an overview similar to Figure 5, which
 404 is described in Section 3.2. Similarly to the front, the parameter R_{Back} has a more
 405 significant influence than the length L_{Back} . An increase of R_{Back} not only directly
 406 enhances the average velocity at the turbine section vel_c but also the intensify the
 407 total forces acting on the structure. It has to be highlighted that the increase from

408 2.5 m/s to close to 3.25 m/s with the larger R_{Back} results in scaled forces, multiplied
 409 by a factor of about 8.

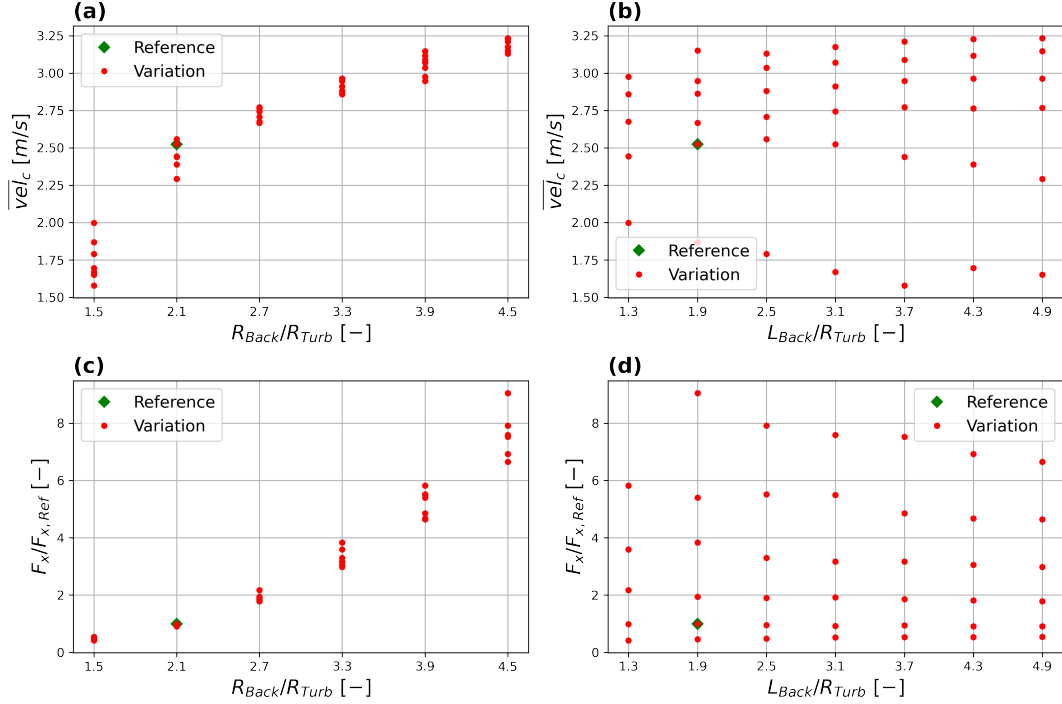


Figure 7. Variation of the diffuser with a fixed nozzle $[R_{Front}, L_{Front}] = [1.2, 1.1] \cdot R_{Turb}$ — area-averaged velocity at the turbine cross section \overline{vel}_c depending on the opening radius R_{Back} (a) and sorted by the length L_{Back} (b) — Forces in x-direction F_x standardised by the reference value (c-d) — all length standardised by the turbine radius R_{Turb} .

410 Hence large R_{Back} act similar to a brim, the result effects of including this additional
 411 blockage element are described within this section. Starting from the reference geom-
 412 etry, which is presented in Section 2.4, the length of the additional brim is changed,
 413 results of which are presented in Figure 8. With an initial R_{Back} of $2.1 \cdot R_{Turb}$, the
 414 brim reaches out further $2.4 \cdot R_{Turb}$ and consequently covers a blockage of $4.5 \cdot R_{Turb}$,
 415 which is similar to the maximum expansion of the geometry parameter R_{Back} of the
 416 diffuser (Fig. 7). An increasing h_{Brim} causes a direct increase of the indicator \overline{vel}_c
 417 at the turbine cross section but also a significant increase of the forces F_x acting on
 418 the overall structure. The conducted variation of the h_{Brim} indicates that the velocity
 419 converges in the range of 3.2 m/s, while the forces are increased by a factor of 8. This
 420 is very similar to the variation of the R_{Back} and explains why both parameters are
 421 detailed jointly within this section. Clearly, the optimisation of both parameters is not
 422 driven by the hydraulic characteristic but instead by construction costs.

423 Figure 9 presents the extreme cases for the variable R_{Back} with the shortest length
 424 L_{Back} . The geometry with the largest R_{Back} introduces a massive blockage and acts
 425 more like a brim with holes. Based on the literature (Nunes et al., 2020), an inclination
 426 of the brim to the front instead of to the back might be advantageous. The observed
 427 convergence behaviour of the simulation with more extreme values indicates that the
 428 flow is not longer steady state and a transient solver would be more appropriate. This
 429 is also obvious in the Figure 9 (right), which shows that the downstream jet is slightly
 430 orientated in the positive y-direction, indicating flow instability caused by this extreme

431 values of the diffusor part of the geometry.

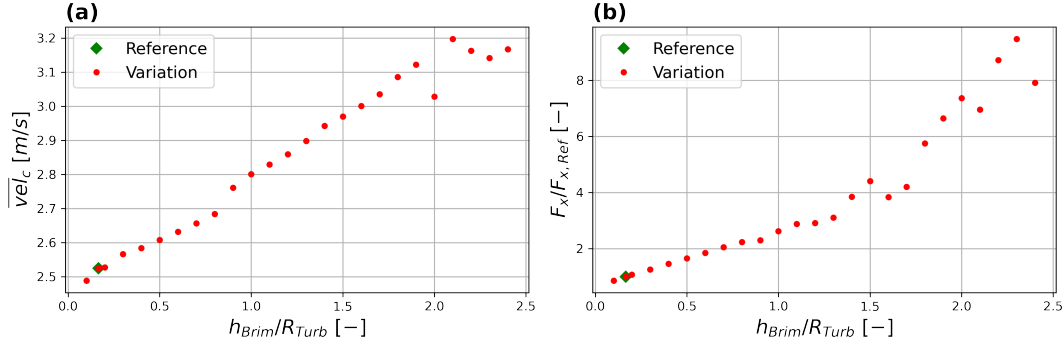


Figure 8. Variation of the height h_{Brim} of the additional brim — area-averaged velocity at the turbine cross section vel_c (a) and Forces in x-direction F_x standardised by the reference value (b) — all length standardised by the turbine radius R_{Turb} .

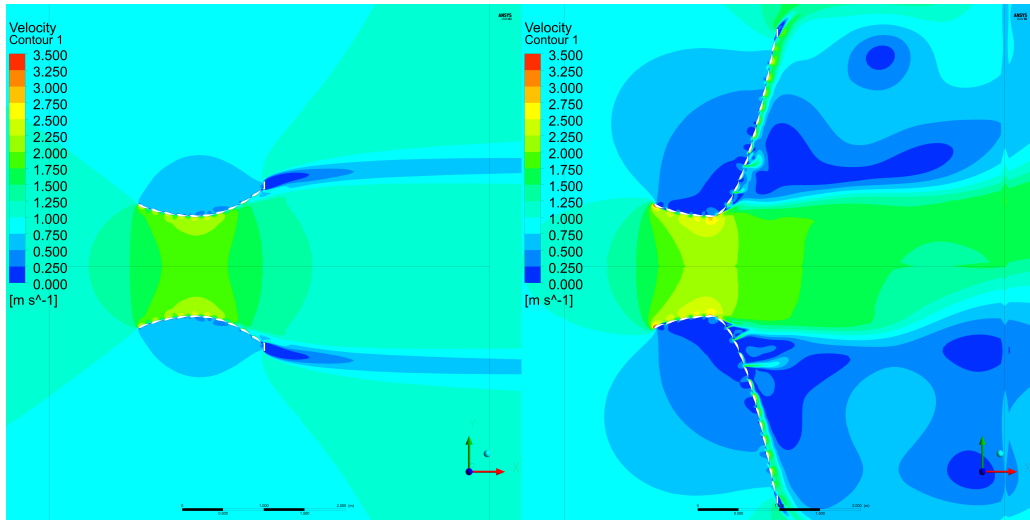


Figure 9. Velocity plot comparing $R_{Back}=1.5 \cdot R_{Turb}$ (left) and $R_{Back}=4.5 \cdot R_{Turb}$ (right) with a fixed $L_{Back}=1.3 \cdot R_{Turb}$.

432 An expansion of the R_{Back} as well as the h_{Brim} results in not only higher velocities
 433 but also increased forces. The extreme expansion of R_{Back} can be combined or even
 434 replaced with a larger brim, which might be more cost effective than the hydrofoils.
 435 But for a real application the additional loading will be prohibitive if too large values
 436 are chosen. An addition of a second turbine instead of a single large turbine should be
 437 considered in this case.

438 3.4. Variation of the turbine radius

439 For all previously reported variations of the geometry the turbine radius R_{Turb} is set
 440 to 600 mm, which is used as the reference value $R_{Turb,Ref}$ for the following variation
 441 of this parameter. Two different concepts of the variation are investigated: (a) the
 442 dimension R_{Turb} is varied independently of the other parameters (Figs. 10 and 11)
 443 and (b) the change of the R_{Turb} is similarly applied for the other two length in the

444 y-direction of the geometry. The latter results in a parallel movement of the reference
 445 geometry reducing or widening the middle cross section (Figs. 12 and 13).

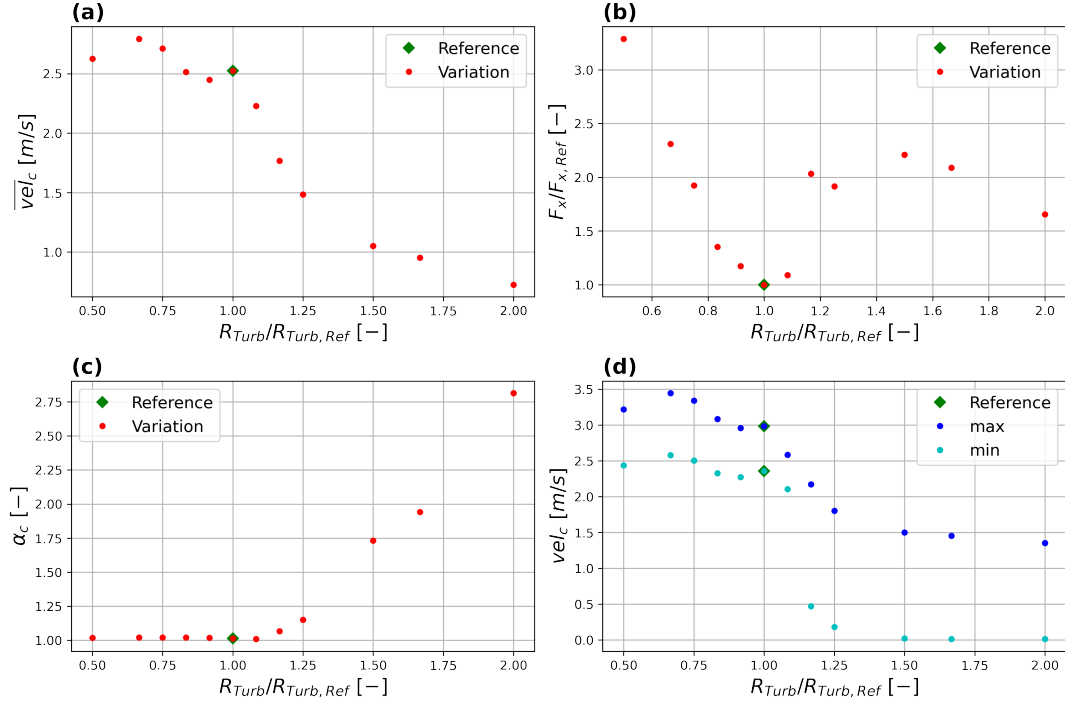


Figure 10. Variation of the turbine radius R_{Turb} , while all other geometry parameters are kept to the reference values — area-averaged velocity at the turbine cross section \overline{vel}_c (a), forces in x-direction F_x standardised by the reference value (b), kinetic energy flux coefficient α_c for the turbine cross section (c) and the extreme values of the velocities in the turbine cross section (d) — all length standardised by the turbine radius of the reference geometry $R_{Turb,Ref} = 600$ mm.

446 The independent variation of the turbine radius R_{Turb} is added more for complete-
 447 ness and not for their actual usability within the optimisation. Since the possibility to
 448 install the turbine is not specifically checked, these cases are not realistic. For those
 449 simulations, the inlet dimension of the nozzle were fixed with $R_{Front} = 1.2 \cdot R_{Turb,Ref}$
 450 and consequently all R_{Turb} values, which exceed this length, the complete flow regime
 451 changes. Figure 10 shows that for larger R_{Turb} values the \overline{vel}_c is significantly reduced.
 452 This results in flow speeds under 1 m/s (boundary conditions at the inlet), which
 453 would be worse than having no supporting structure. The analysis of the α_c indicates
 454 an extreme increase of inhomogeneous flow and the minimum value of the velocity
 455 falls down to 0 m/s. Such an expansion is clearly not a good choice. Figure 11 clearly
 456 show that the recirculation zone moves from the outside to the inner area, which is
 457 far from a favourable result.

458 More realistic than the previously single variation of the R_{Turb} is a parallel change
 459 of the structure. Where all parameters in the y-direction, namely R_{Turb} , R_{Front} and
 460 R_{Back} , are change with the identical value. This results in a constant offset in the y-
 461 direction keeping the geometry symmetrical to the x-axis. The results of this variation
 462 are presented in Figure 12 covering a range of a quarter of the initial reference radius
 463 to the doubled value. Reducing the space between the side Venturi structures increases
 464 the average velocity at the turbine cross section \overline{vel}_c but also results in a higher total
 465 force F_x on the structure. Reducing R_{Turb} to $0.25 \cdot R_{Turb,Ref}$ increases flow speed in
 466 the range of additional 1 m/s while also multiplying the resulting forces by a factor of

467 2.5 in relation to the reference values with a $R_{Turb,Ref} = 600$ mm.

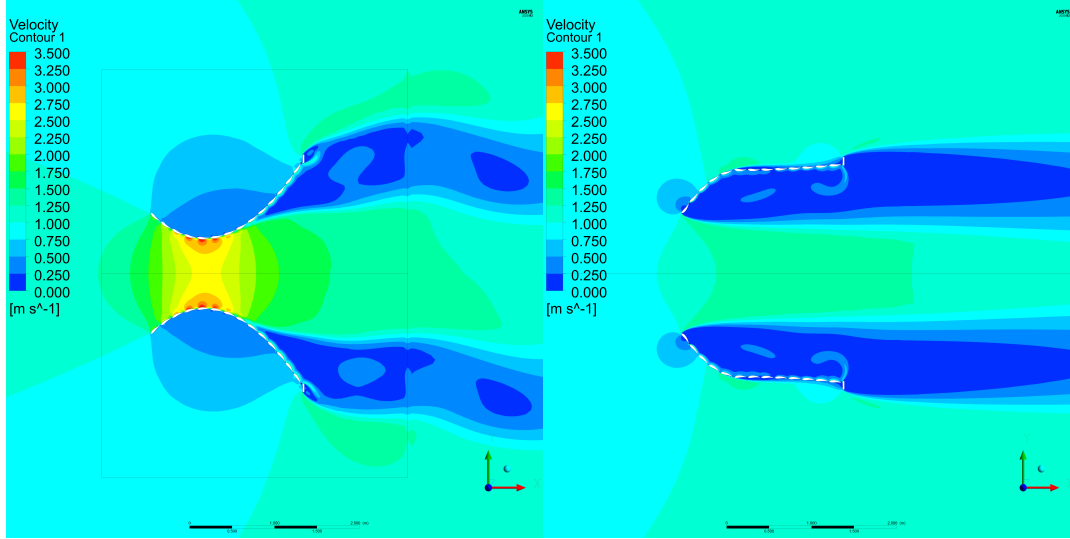


Figure 11. Velocity plot comparing the reference geometry with a turbine radius R_{Turb} of 400 mm, which is equal to $0.66 \cdot R_{Turb,Ref}$ (left,) in comparison with $2 \cdot R_{Turb,Ref}$ (right).

468 The velocity distribution becomes more homogeneous with the a smaller R_{Turb} indi-
 469 cating by the α_c -value and a comparable small velocity difference of 0.5 m/s (maximum
 470 subtracted by the minimum velocity). In this case the cross section at the turbine is
 471 reduced to a quarter but the discharge, which flows through the cross section reaches
 472 approximate 34% of the discharge through the reference geometry.

473 The smallest investigated distance between both sides is compared with the largest
 474 configuration in the Figure 13. It shows the comparable high velocity in the turbine
 475 cross section, which results in a jet at the outlet. The obvious asymmetry of the velocity
 476 distribution in the wake of the turbine indicates that an unstable flow is generated. It
 477 seems that each side has a range of influence and if the distance is reduced, both sides
 478 support each other. To test this assumption, a specific variation of the full structure
 479 is presented in Section 3.5. In any case, this investigation has to be repeated in a
 480 full 3D-model and with a full turbine model in order to assess if higher velocities can
 481 generate a positive impact on energy production.

482 This section includes a theoretical investigation of the single variation of the turbine
 483 radius R_{Turb} , which resulted in some completely different flow regimes. The second
 484 variation proposes a parallel offsetting of both side geometries with the scaling of all
 485 three defining lengths in the y-direction. Narrowing the channel between both sides
 486 resulted in increased velocities and forces. This indicates that a combined variation of
 487 nozzle and diffuser have the potential to improve the velocity at the turbine further.
 488 How far this is beneficial for the actual energy production at the turbine will need to
 489 be investigated separately.

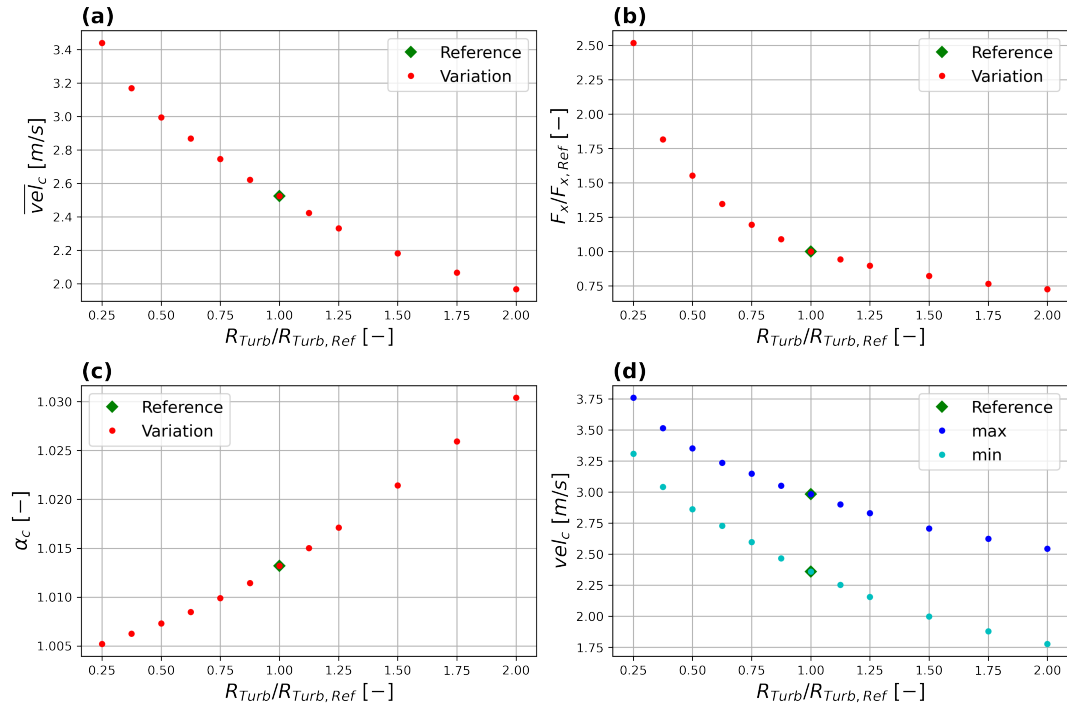


Figure 12. Variation of the turbine radius R_{Turb} including R_{Front} and R_{Back} resulting in parallel move of the sides — area-averaged velocity at the turbine cross section \overline{vel}_c (a), forces in x-direction F_x standardised by the reference value (b), kinetic energy flux coefficient α_c for the turbine cross section (c) and the extreme values of the velocities in the turbine cross section (d) — all length standardised by the turbine radius of the reference geometry $R_{Turb,Ref} = 600$ mm.

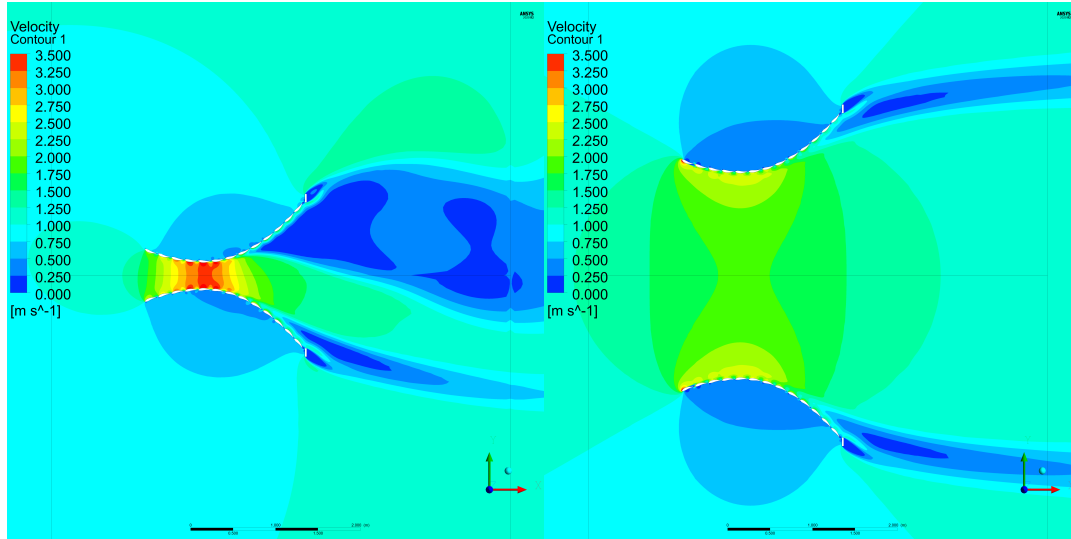


Figure 13. Velocity plot comparing two cases with a parallel move of the structure with the turbine radius R_{Turb} : $0.25 \cdot R_{Turb,Ref}$ (left) and $2 \cdot R_{Turb,Ref}$ (right).

490 **3.5. Combination – scaling the structure**

491 The second part of the previous Section 3.4 shows that a parallel offsetting of the full
 492 structure to the centre results in increased speed at the centre of the turbine (Fig 12).
 493 A preliminary hypotheses is the aforementioned range of influence one side has to the
 494 other, which overlaps if the sides are close enough. To test this, a combined change
 495 of the nozzle and diffuser part is investigated in this section. The factor fac is used
 496 to scale the geometry while keeping the turbine radius R_{Turb} constant. This factor is
 497 applied directly to the length L_{Front} and L_{Back} as well as to the difference value of
 498 the R_{Front} and R_{Back} as presented in Eq. (3).

$$fac = \frac{L_{Front}}{L_{Front,Ref}} = \frac{L_{Back}}{L_{Back,Ref}} = \frac{R_{Front} - R_{Turb}}{R_{Front,Ref} - R_{Turb}} = \frac{R_{Back} - R_{Turb}}{R_{Back,Ref} - R_{Turb}} \quad (3)$$

499 This is obviously not perfectly scaled since the hydrofoils remained one type, main-
 500 taining specific lengths and distances along the channel, but it does provided an ap-
 501 proximation allowing the simulation of comparable geometry changes as conducted
 502 with the variation of the R_{Turb} (Fig. 12). The results of this scaling are shown in
 503 Figure 14. A smaller factor fac down to 0.5 is comparable to a larger R_{Turb} (larger
 504 distance between the two sides). Both result in a reduction of the velocity at the
 505 turbine cross section \overline{vel}_c and the total forces F_x . A variation in the other direction
 506 shows similarity with the previously reported $R_{Turb}/R_{Turb,Ref} = 0.5$ (Fig. 12) a \overline{vel}_c
 507 of approximately 3 m/s, which can also be found for a scaling factor $fac = 2$ in Fig-
 508 ure 14. This indicates that both geometries are comparable, however, the forces show
 509 that $R_{Turb}/R_{Turb,Ref}$ equates to 1.5 times (Fig. 12) of the reference value, while it
 510 is 2.5 times (Fig. 14) the loading for the scaled version of the geometry. This scaled
 511 geometry has a $R_{Back} = 3.2 \cdot R_{Turb,Ref}$ and $L_{Back} = 3.8 \cdot R_{Turb,Ref}$. Although com-
 512 parable in geometry, with the front section fixed to the reference geometry, the results
 513 show \overline{vel}_c of around 3 m/s with higher ratios of $F_x/F_{x,Ref}$ greater than 4 (Fig. 7).
 514 The analysis in Figure 14 also includes a further adaptation for the extreme values
 515 while only the R_{Front} and R_{Back} is modified by the factor fac . For those two specific
 516 cases, the lengths L_{Front} and L_{Back} are kept at the reference length (Sec 2.4). This
 517 was completed in Section 3.2 and 3.3 which identify the R -dimensions as dominate
 518 parameters. The direct comparison to the values for the full scaling with a factor fac
 519 of 0.5 and 2 show, that those specifically marked parameters result in a lower \overline{vel}_c ,
 520 higher F_x , comparable α_c -values and the range of the extreme values is wider and
 521 moved. Consequently, full scaling shows better results than utilising only R -scaling.

522 Figure 15 shows the velocity plots of the four mentioned extreme values. The two
 523 top row pictures present the results of the full scaling according to the Eq. (3). For the
 524 geometries in the bottom row, lengths L_{Front} and L_{Back} are not scaled and only R_{Front}
 525 and R_{Back} are. These results indicate that the R -length in the geometry are dominate
 526 but that scaling in length along the flow direction should also be considered. How
 527 far such an adaptation results in a real performance improvement must be checked
 528 separately.

529 **3.6. Angle of the first hydrofoil**

530 This specific section focuses on the orientation of the first hydrofoil in the flow di-
 531 rection, describing a key function which provides the addition of an initial direction

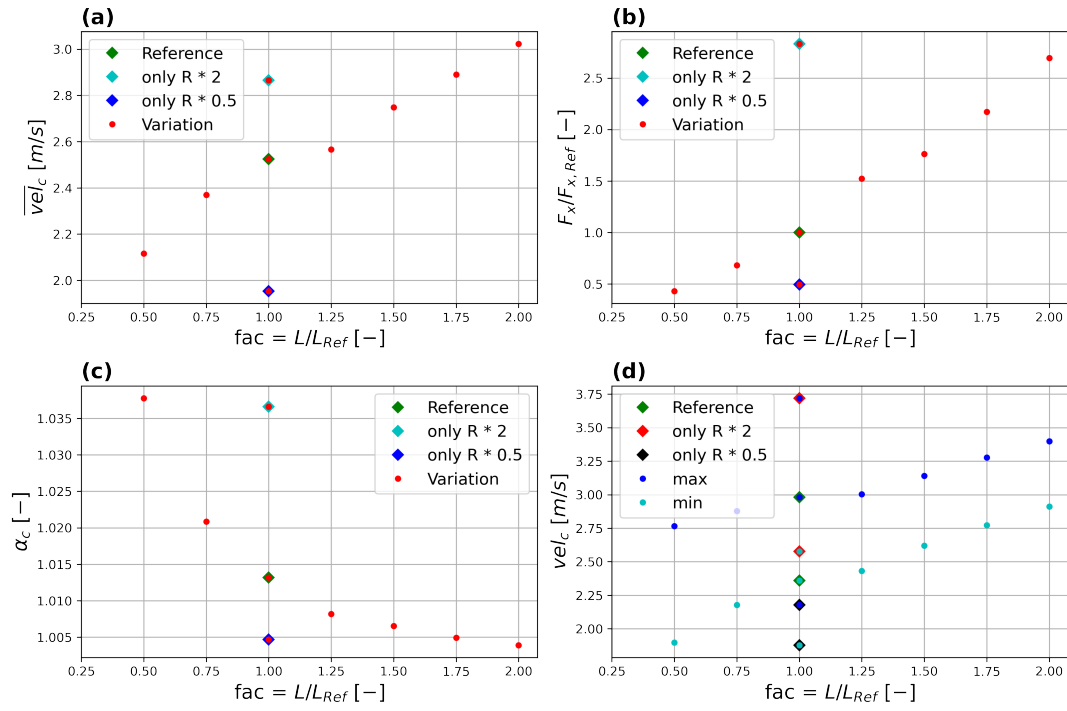


Figure 14. Scaling by different factors f of the structure with a fixed turbine radius R_{Turb} — area-averaged velocity at the turbine cross section \overline{vel}_c (a), forces in x-direction F_x standardised by the reference value (b), kinetic energy flux coefficient α_c for the turbine cross section (c) and the extreme values of the velocities in the turbine cross section (d).

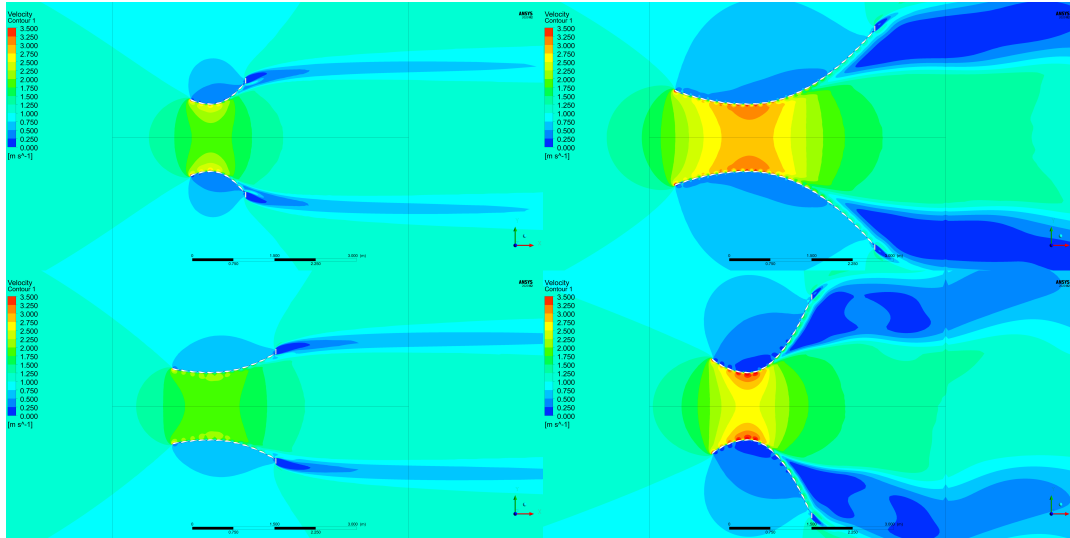


Figure 15. Velocity plot comparing four different scaled geometries: factor $fac = 0.5$ (left, top) and 2 (right, top) — the scaling is only applied for the R -values with a $fac = 0.5$ (left, bottom) and 2 (right, bottom).

532 change included to benefit of the Venturi structure. The parameter $\Delta\alpha_{FirstFoil}$ is
 533 introduced to describe this additional angle of this specific foil. For all other cases,
 534 the hydrofoil is rotated according to the tangential direction of the bell shape at the
 535 definition point of the individual hydrofoil.

536 The conducted variation of this additional angle covers a range of $\pm 90^\circ$ from the
 537 original tangential orientation with a $\Delta\alpha_{FirstFoil} = 0^\circ$. Figure 16 presents the re-
 538 sults, showing that incremental negative angles improve the area-averaged velocity at
 539 the turbine cross section \overline{vel}_c occurring at velocities slightly over 2.6 m/s around a
 540 $\Delta\alpha_{FirstFoil}$ of -20° . This is also connected with a reduction of the total forces in the
 541 x-direction. Consequently, such a modification of the first hydrofoil is advantageous.

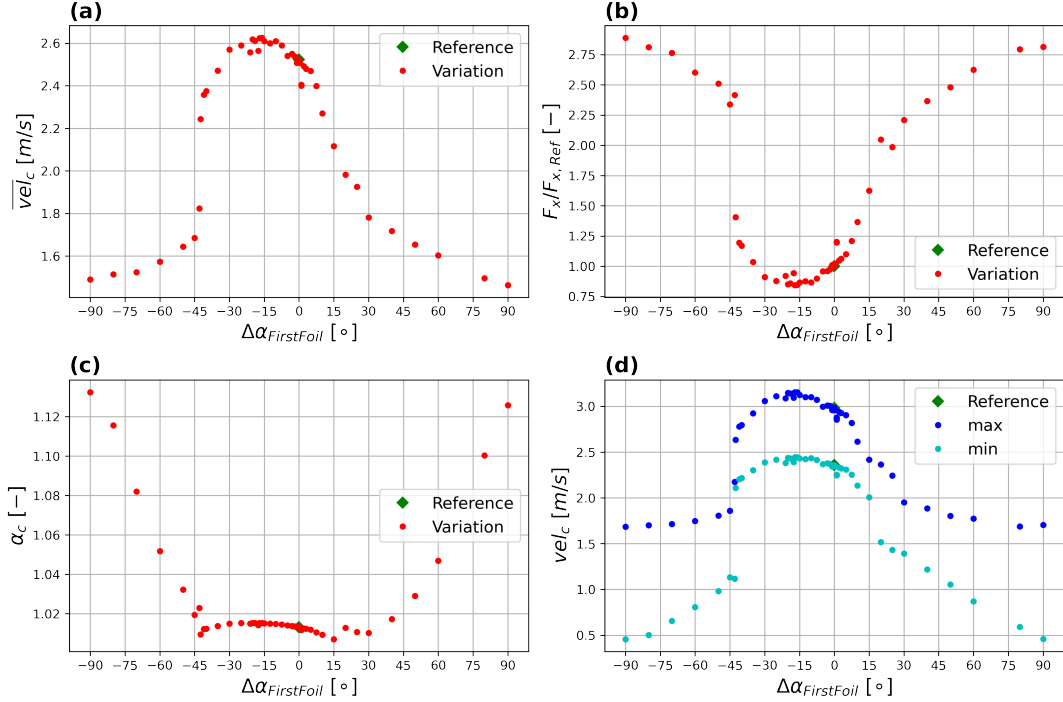


Figure 16. Variation of the additional angle of the first hydrofoil $\Delta\alpha_{FirstFoil}$ — area-averaged velocity at the turbine cross section \overline{vel}_c (a), forces in x-direction F_x standardised by the reference value (b), kinetic energy flux coefficient α_c for the turbine cross section (c) and the extreme values of the velocities in the turbine cross section (d).

542 The conducted variation of the parameter is deliberately expanded to investigate
 543 the range of influence in more detail. Extreme values of $\pm 90^\circ$, for example, result in a
 544 significant reduction of the velocity value \overline{vel}_c , while the kinetic energy flux coefficient
 545 α_c is significantly increased. This indicates that the first hydrofoil introduces a dis-
 546 turbance, which has a downstream influence. Obviously, this can also be seen in the
 547 evaluation of the extreme values in the turbine cross section shown in Figure 16 (d).
 548 Minimum values are reduced to 0.5 m/s , equal to half of the inlet velocity of 1 m/s .

549 Figure 16 indicates that an optimum value can be reached by including an additional
 550 angle of around -20° . The benefits are presented in the top row of Figure 17. A value of
 551 -20° for the parameter $\Delta\alpha_{FirstFoil}$ results in the end of the first hydrofoil being placed
 552 slightly closer to the symmetry plane (x-z-plane) leading to an overall flow velocity
 553 increase. A rotation in the opposite direction generates a step and the flow detaches.
 554 The ramification of which is an overall reduced flow speed at the turbine cross section

555 as well as higher forces on the structure. The extreme values of $\pm 90^\circ$ are also shown in
 556 the Figure 17 representing the boundaries of this variations. Both cases are obviously
 557 not ideal.

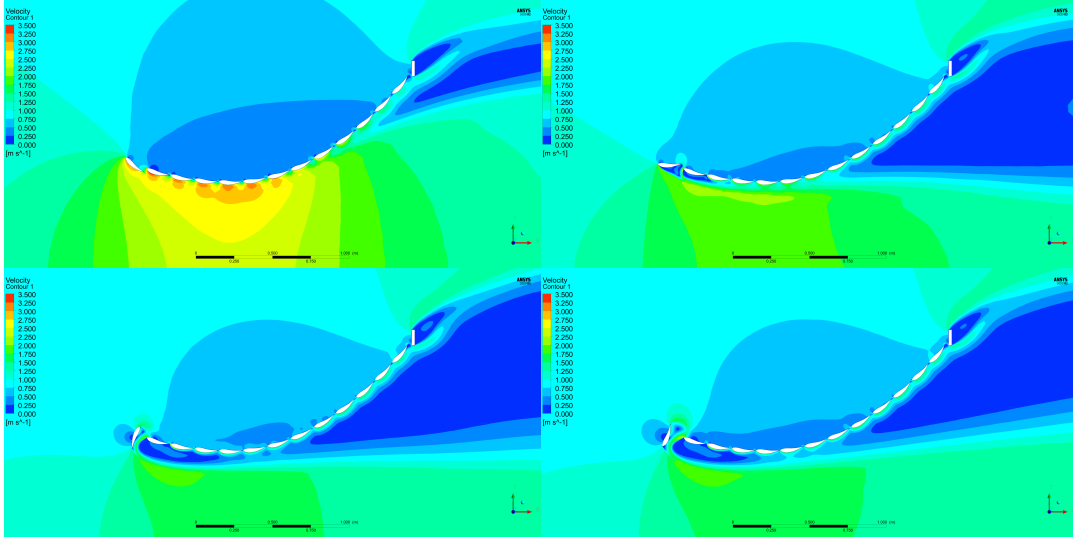


Figure 17. Velocity plot comparing four different additional angles applied on the first hydrofoil in the flow direction: $\Delta\alpha_{FirstFoil} = -20^\circ$ (left, top), 20° (right, top), -90° (left, bottom) and 90° (right, bottom).

558 The orientation of the first hydrofoil in the flow direction has a significant influence,
 559 while the precise optimum angle is dependant on the size and type of hydrofoil profile
 560 implemented. Consequently, this parameter is one of the last criteria to optimise after
 561 key geometry decisions are made.

562 3.7. Hydrofoil length and distance

563 The previous variation focused on the main component of the geometry utilising a
 564 single hydrofoil GOE 222 (MVA H.33) AIRFOIL (Airfoil Tools, 2022) with a chord
 565 length c_{Foil} of 150 mm and a linear chord distance s_{Target} of 160 mm between the
 566 definition point of each hydrofoil on the bell curves. The type of the hydrofoil is varied
 567 in Section 3.8 and this section focuses on the length as well as the distance between
 568 the foils.

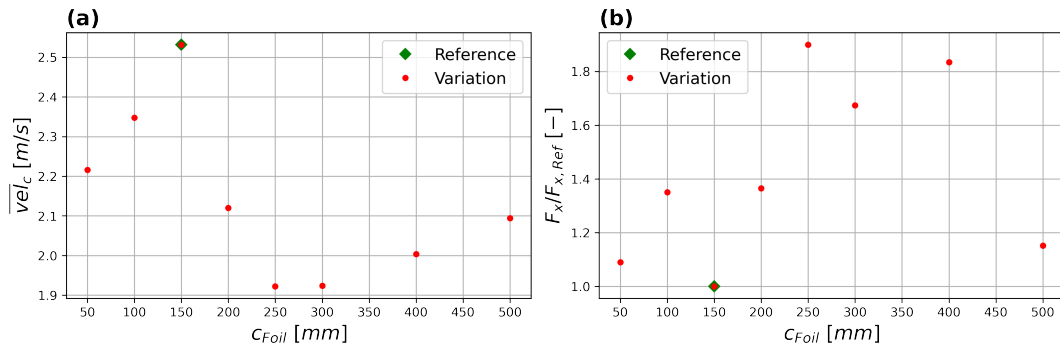


Figure 18. Variation of the chord length c_{Foil} of the hydrofoil — area-averaged velocity at the turbine cross section \overline{vel}_c (a) and Forces in x-direction F_x standardised by the reference value (b).

569 Figure 18 shows the results for different chord length of the same standard hydrofoil.
 570 A range between 50 mm and 500 mm is presented. The smaller blades use the same
 571 concept for the calculation with $s_{Target} = c_{Foil} + 10$ mm, while additional values were
 572 increased to 20 mm for the 400 mm foil and 30 mm for the 500 mm foil. This ensured
 573 that the individual elements are separated. Looking only at this analysis, the average
 574 velocity at the turbine section \overline{vel}_c is maximised and the forces F_x are at a low level
 575 (Fig. 18). The four examples shown in Figure 19 indicated that the choice of the blade
 576 size has a significant influence on the structure.

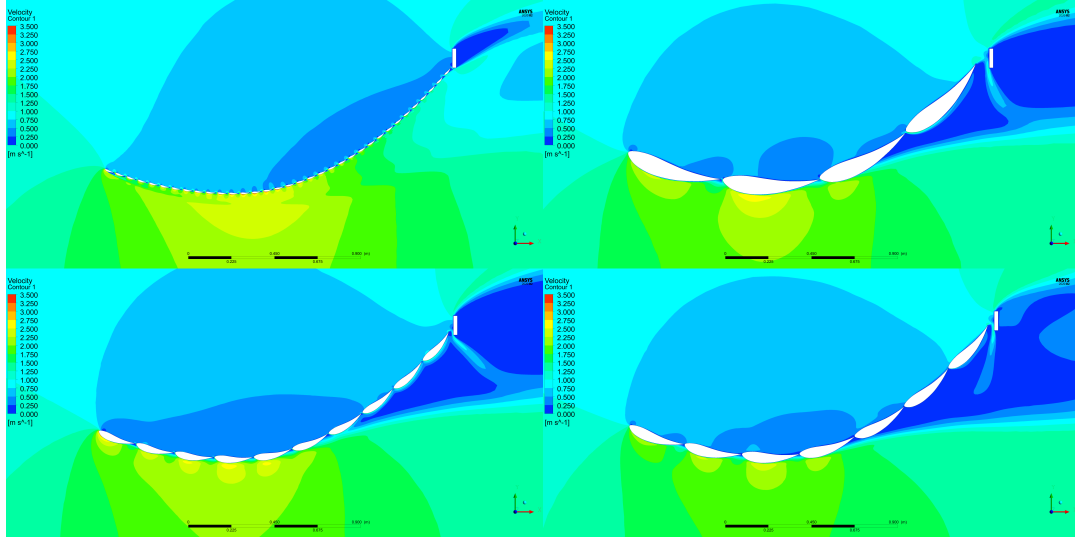


Figure 19. Velocity plot comparing four different length of the hydrofoil: 50 mm(left,top), 200 mm (left, bottom), 300 mm (right, bottom) and 500 mm (right, top).

577 The previous sections show that specific parameters can have a significant influence
 578 on the overall performance of the turbine. It also must be noted that the definition
 579 point of all cases is identical but the bell shape can only be filled with the same
 580 hydrofoils based on the constant linear chord distance along the curves. Smaller profiles
 581 can better reach a precise endpoint and the larger profiles would have to be spaced
 582 perfectly to reach a better comparability. Keeping the limitation of this comparison
 583 in mind, all investigated length performed in a very comparable way and might need
 584 more detail optimisation to reach an optimum.

585 In a second step, the linear chord distance s_{Target} is varied to investigate the initial
 586 choice of 160 mm. This is conducted for the hydrofoil GOE 222 (MVA H.33) AIRFOIL
 587 (Airfoil Tools, 2022) only, with a chord length c_{Foil} set to 150 mm. s_{Target} is varied
 588 between 155 mm and 300 mm, which is double the chord length of the hydrofoil. For
 589 this variation the step size Δs is changed from 10 mm to 1 mm to better represent
 590 the small alterations of the parameters within each simulation. This value is used to
 591 find the required x-distances for the following definition point, indicating that these
 592 changes result in a reduction of the \overline{vel}_c for the reference geometry. Consequently,
 593 for this case, two variation points are reported in Figure 20 for the s_{Target} length
 594 of 160 mm. A slight expansion results in an overall improvement of the \overline{vel}_c and
 595 reduction of the total forces F_x can be observed, with an optimum point at 166 mm.
 596 It is not surprising, that larger values introduce large gaps in the structure and hence
 597 the efficiency is reduced. Figure 21 shows the comparison of the a s_{Target} of 166 mm
 598 and the largest value of 300 mm. With a larger gap between the individual hydrofoils

599 the behaviour of the structure changes from a joint guidance of the flow into separate
 600 obstacles in the flow.

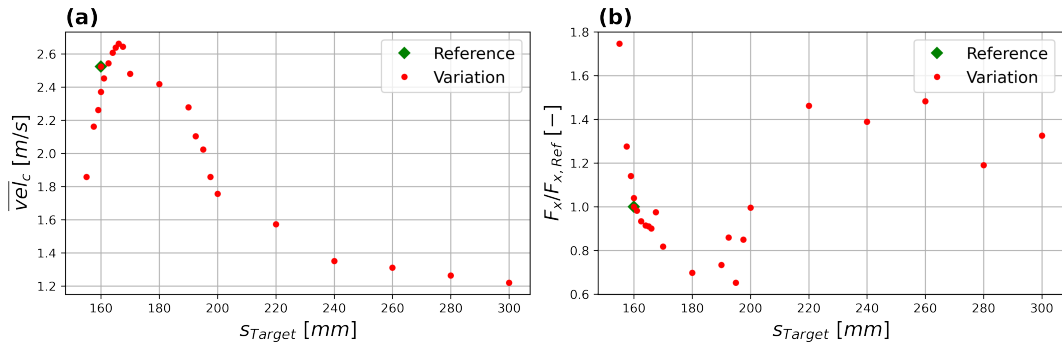


Figure 20. Variation of the length s_{Target} , which defines the distance between the reference points on the bell shape — area-averaged velocity at the turbine cross section \overline{vel}_c (a) and Forces in x-direction F_x standardised by the reference value (b).

601 This variation of hydrofoil chord length and the definition point indicates that
 602 the initial chosen values are in a good range. Nevertheless, the distance between the
 603 definition points of the hydrofoils indicated some optimisation potential. However,
 604 extremes in hydrofoil sizing do not provide better results. The decisive point for this
 605 decision is more the likely the availability and costs of the production of the hydrofoils.

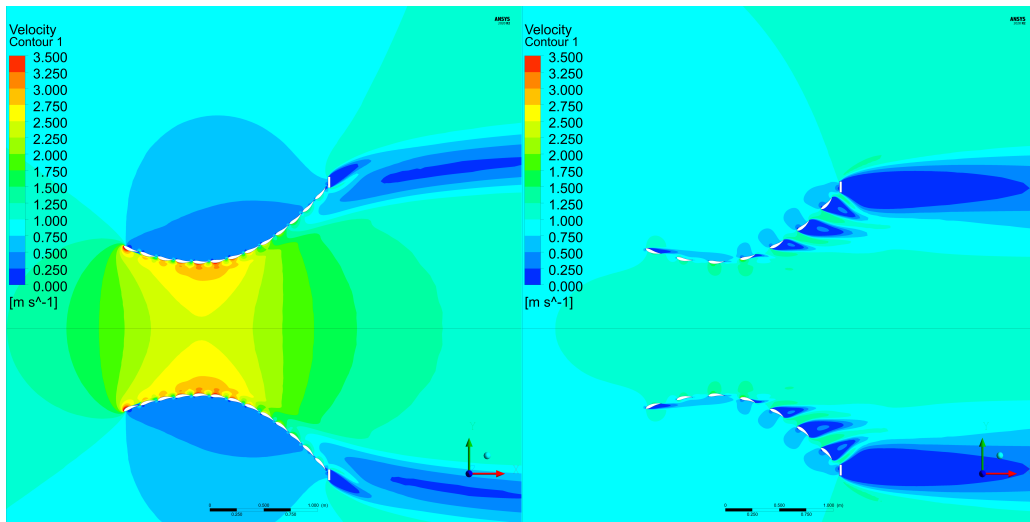


Figure 21. Velocity plot comparing two different chord length s_{Target} between the definition points of the hydrofoils: 166 mm(left) and 300 mm (right).

606 3.8. Hydrofoil type

607 In this section the type of hydrofoil is altered, while keeping a constant chord length
 608 $c_{Foil} = 150$ mm and distance between the definition points for the individual foils
 609 $s_{Target} = 160$ mm. Similar to the GOE 222, the geometry data is sourced from the
 610 online database Airfoil Tools (2022), which provides over 1600 different airfoil types.
 611 The selection of additional foils were conducted based on a random choice, while
 612 trying to cover a wider range of different shapes. In addition, a circular cross section
 613 is also tested, for which the s_{Target} has to be increased to 165 mm (Note that for this
 614 simulation the definition points are located on the inside of the turbine and not in the
 615 centre line of the hydrofoils). Figure 22 provides an overview of the investigated types.

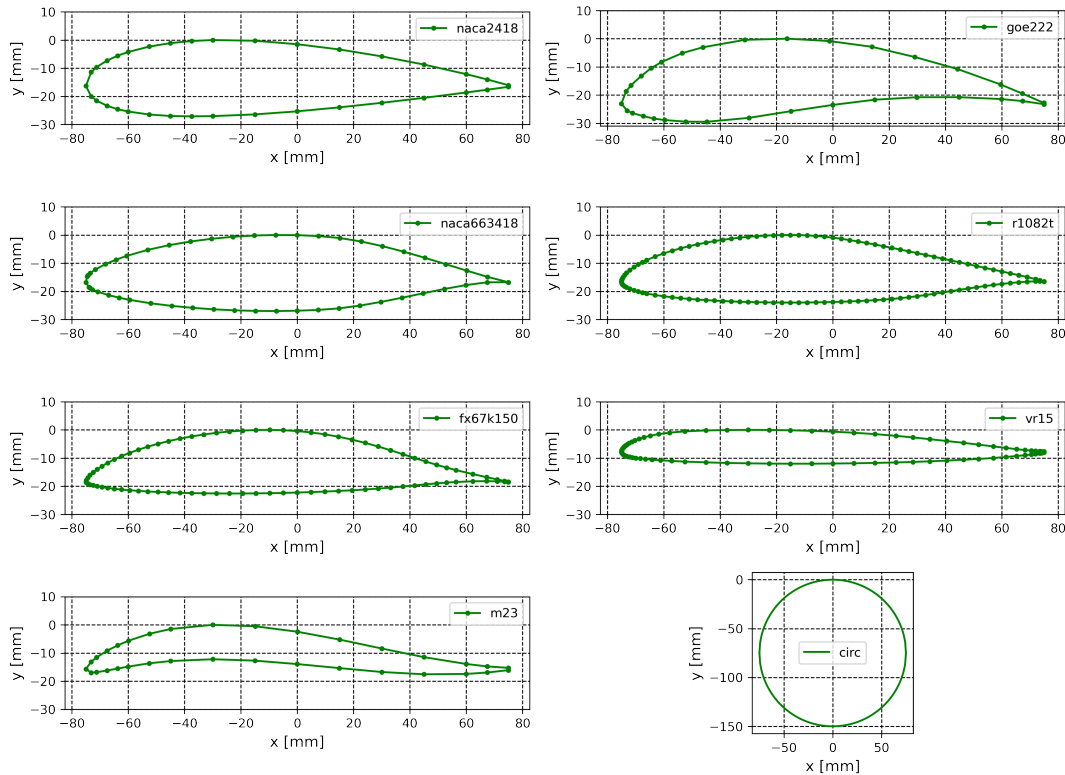


Figure 22. Overview of the investigated hydrofoils.

616 The results of the variation of the hydrofoil types is presented in Figure 23, sorted
 617 depending on the area-averaged velocity at the turbine cross section \overline{vel}_c . A slight
 618 improvement in velocity as well as the total forces in the x-direction is noted with the
 619 NACA 2418 hydrofoil. All other foils resulted in worse conditions with a massive reduc-
 620 tion for the circular cross section. The latter requires a much more massive structure
 621 while only producing a slightly higher total force. The kinetic energy flux coefficient
 622 α_c is very similar for all investigated hydrofoils with the largest inhomogeneity shown
 623 for the circular profile, which provides the most extreme values for the velocities in
 624 the cross section.

625 Figure 24 shows the comparison of four of the eight investigated options. The top
 626 row shows the best and worst options side by side and it is obvious that the circular
 627 cross section is not a good option for this approach. The comparable large radius of
 628 the circle in the lateral direction result in a comparable massive structure and it works

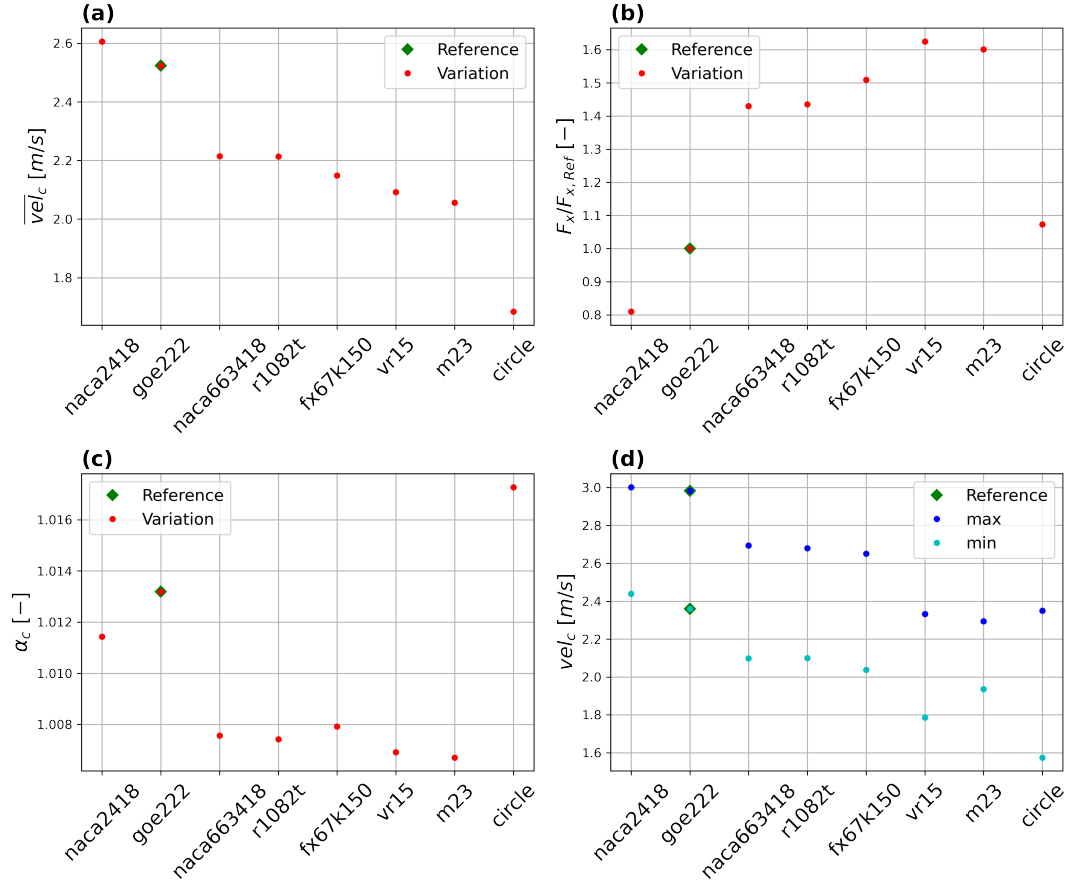


Figure 23. Variation of the hydrofoil types — area-averaged velocity at the turbine cross section \overline{vel}_c (a), forces in x-direction F_x standardised by the reference value (b), kinetic energy flux coefficient α_c for the turbine cross section (c) and the extreme values of the velocities in the turbine cross section (d).

629 more like a solid Venturi channel with additional wall roughness. Interestingly the tip
 630 of the foil of the very thin M23 hydrofoil. The flow detaches very early and results
 631 in a small recirculation zone behind the first foil. On the other hand, the two foils
 632 on the left cause a velocity peak at this location. This indicates that the fixed angle
 633 along the bell curve with an orientation along the tangential direction is not ideal for
 634 all foils. A correction of the first foil in the incoming flow direction could result in an
 635 improvement of the results for this specific foil.

636 This section analyses seven additional hydrofoils for comparison with the originally
 637 used GOE 222 foil. Slight improvement can be found changing to a NACA 2418,
 638 which is more symmetric and has the potential to be less expensive for manufacture.
 639 The distance between the definition points of the hydrofoils may allow for further
 640 improvements, such as that shown in Section 3.7. Importantly, it was found that the
 641 orientation of the first foil in the flow direction is critical, and hence a correction
 642 angle for this specific profile was added in the Python code for the generation of the
 643 geometry in Gabl et al. (2022) and a variation is presented in Section 3.6.

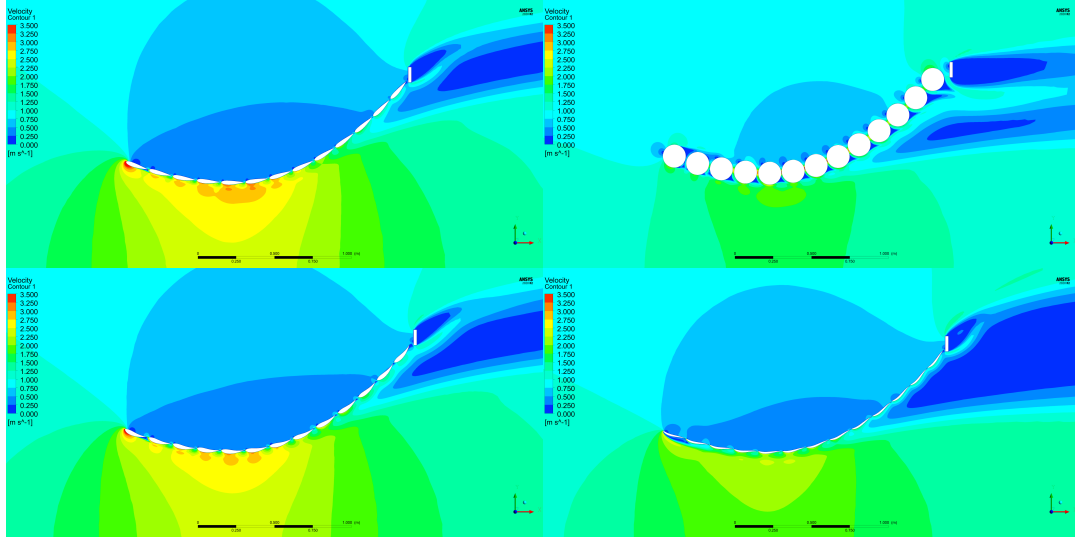


Figure 24. Velocity plot comparing the reference geometry with three other types of hydrofoils: NACA 2418 (left,top), GOE 222 (left, bottom), M23 (right, bottom) and circle (right, top).

644 3.9. Inlet velocity

645 All previous simulation assumed that the inlet velocity is constant 1 m/s, which is
 646 introduced homogeneously at the inlet boundary conditions. In this section, the value
 647 is varied over a range of 0.5 to 2 m/s, while keeping all the other assumptions the
 648 same, including the constant reference geometry (Sec 2.4). Figure 25 presents the
 649 results of this variation. Contrary to the previous analysis, the velocity at the turbine
 650 cross section \overline{vel}_c is normalised by the inlet velocity v_{Inlet} . Hence the standard v_{Inlet}
 651 is 1 m/s, the values can be directly compared with the previous results, while being non-
 652 dimensionalised. This value changes in relation with v_{Inlet} in a range of approximately
 653 ± 0.2 [-] from the reference value. A doubling of the v_{Inlet} causes an increase of the
 654 forces by a factor of approximate 3.5 [-]. Tidal flows are generally very predictable but
 655 nevertheless extreme flow speeds can occur resulting in extreme loads on the structure.

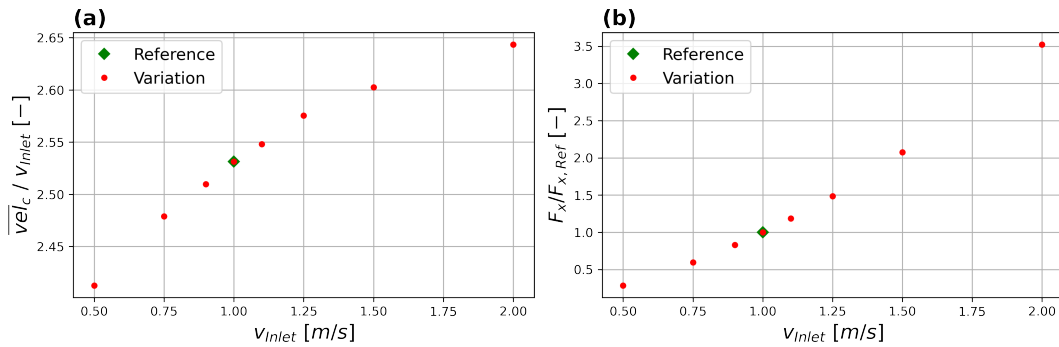


Figure 25. Variation of the inlet velocity v_{Inlet} to investigate the standardised velocity at the turbine cross section \overline{vel}_c (a) and the total forces F_x on the structure in x-direction (b).

656 Figure 26 provides the slowest and fastest investigated flow speeds of 0.5 and 2 m/s
 657 respectively. A separate colour bar is used for each result, with maximum values scaled
 658 provided according to the inlet velocity. The downstream section with close to 0 m/s

659 is smaller for the increased flow speed. A higher gradient in the cross section can
 660 be observed for the higher speed, which can also be seen in an increased normalised
 661 $\overline{vel}_c/v_{Inlet}$ value.

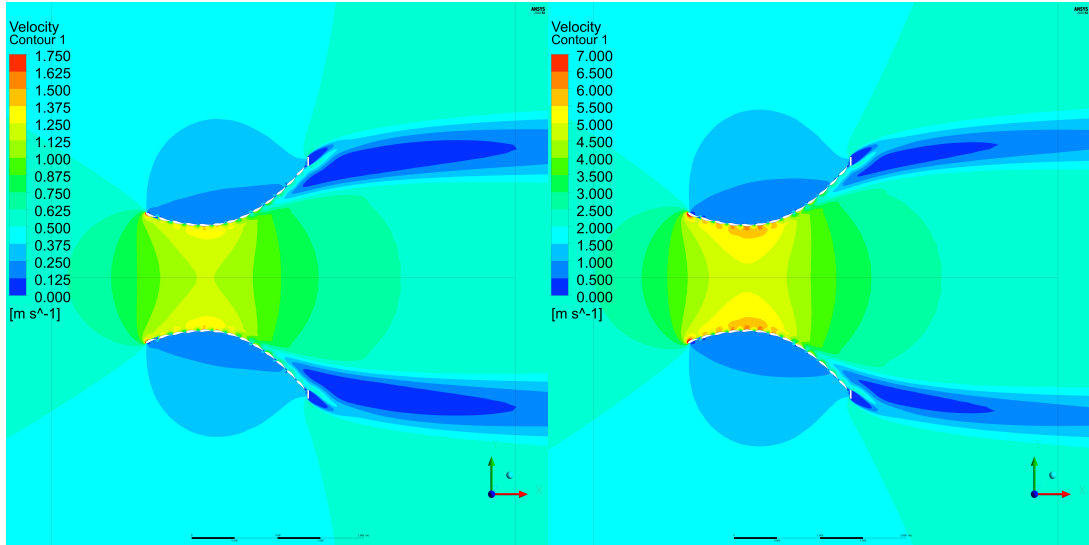


Figure 26. Variation of the inlet velocity v_{Inlet} with 0.5 m/s (left) and 2 m/s (right) — colour bar scaled with the same factor.

662 In summary, the change of inlet velocity has an influence on the velocity at the
 663 turbine cross section \overline{vel}_c making it necessary to optimise the specific design for each
 664 deployment site. Extreme flow speeds also have to be considered in the design of the
 665 support structure to ensure the survivable of the device.

666 4. Discussion

667 The obvious limitation of this presented work is that the numerical simulations are
 668 limited to a 2D-approach. All the used computation meshes have only one cell in the
 669 vertical direction and hence assuming that there is no change in the z-direction. This
 670 neglects potential velocity differences, which are significant closer to the ground Ahmed
 671 et al. (2017); Badshah et al. (2018); Ke et al. (2020). Furthermore, the structure has a
 672 clear top an bottom part, which is in the previous tests a plane cross section (Fig. 1 (a)).
 673 This causes additional interactions and has to be included in the overall optimisation.

674 It is assumed that the current direction is perfectly aligned, homogeneously dis-
 675 tributed and constant. Obviously this is an idealisation and even in deep current
 676 streams, the interaction of waves as well as large turbulence in the flow causes vari-
 677 ations in the velocity distribution or changes in the flow directions. In all cases, only
 678 the steady state solver is used and some results indicated that even with a fully perfect
 679 inflow condition the resulting wake can be an unstable flow. The following detailed
 680 optimisation has to consider the usage of a fully transient solver.

681 The computational grid is a standard point, which can be always improved. All cur-
 682 rent simulations use a comparable fine mesh around the profiles but not a designated
 683 inflation layer. Consequently, the y^+ - values are comparable high, but similar to com-
 684 parable investigations (Maduka & Li, 2021) for a further optimisation a designated

685 refinement close to the wall should be considered. The conducted verification process
686 is presented in Section 2.6 but limited to the reference geometry. It is assumed that the
687 mesh independence of the results is also given for changes of the geometry. Section 3
688 shows that some variations of the geometry cause changes in the location of the recir-
689 culating zones and the concentration of the wake in a jet. For those cases, an additional
690 check of the fluid domain and the mesh resolution would be advisable. Nevertheless,
691 those geometries are not likely to be chosen in the future as their performance is not
692 an improvement.

693 Further limitations can be found in the geometry definition. All cases are built
694 using a single hydrofoil and with identical chord distance set on the bell curves. Hence
695 only full hydrofoils are added, not all changes in the geometry definition have an
696 immediate impact in the actual geometry. A certain threshold has to be exceeded
697 to trigger the addition or removal of a profile. Nevertheless, the conducted geometry
698 generation is ideal for such an exploitative approach but can be refined in an additional
699 detailed optimisation. Furthermore, the hydrofoils are always orientated according to
700 the tangential direction. All of this aspects can be individually modified and varied
701 for each hydrofoil. Obviously, the number of potential options for the investigations
702 increases dramatically.

703 The presented research work focuses on the Venturi structure and neglects the
704 influence of the vertical axis turbine. In a following step, the found improvements
705 will be checked, if the changes are really beneficial for the overall energy production.
706 Currently, only fully symmetrical geometries are investigated but there may be benefits
707 in modifying one side, allowing for optimisation of the geometry where the rotating
708 turbine moves against the flow direction. Further detailed investigations are needed to
709 further improve the overall efficiency of the Venturi shaped structure.

710 5. Conclusions

711 The presented research work focuses on the hydraulic performance of the surround-
712 ing structure of the Davidson Hill Venturi (DHV) tidal turbine. Multiple hydrofoils
713 are placed in a Venturi shape to increase the flow speed at the vertical axis turbine.
714 The first step was to provide a generalised description of the structure (Gabl et al.,
715 2022), which was done in Python allowing future expansions, including using a differ-
716 ent numerical solver. In the second step, this geometry description was adapted for
717 the ANSYS-Workbench with the commercial solver ANSYS-CFX, the files for which
718 are available in the connected datashare (Gabl et al., 2022). A wide range of geom-
719 etry parameters were investigated in a 2D-approach without the turbine based on a
720 reference geometry (Sec 2.4). This variation included the nozzle part as well as the
721 diffuser, including the brim and an additional angle of the first hydrofoil in the flow
722 direction. Larger structures resulted in an improved velocity at the turbine cross sec-
723 tion, however, also increased the forces on the structure. The variation of the turbine
724 radius on both sides simultaneously showed potential for further improvements with-
725 out resulting in extreme operational forces. It can be concluded that the hydrofoil
726 chord length does not significantly influence operational criteria, but that fabrication
727 costs would need to be considered as part of an overall device optimisation. By in-
728 creasing the gap between the hydrofoils as well as the exchanging the GOE 222 with
729 at NACA 2418 hydrofoil, some improvements in output were shown. In addition to the
730 variations of the geometry, the inlet velocity was varied showing that better results
731 can be achieved with higher flow speeds but no significant decrease for lower speed.

732 Overall, these variations show that the chosen reference geometry results provide good
733 system performance, with only small improvements being achieved by tweaking de-
734 sign constraints. The next step for the research work is to expand the results from
735 the 2D-approach and integrate the full supporting structure as well as the turbine
736 optimisation.

737 Disclosure statement

738 M.H. is the owner of the company, which commercialises the tidal turbine concept
739 presented in this work. The authors declare no further conflict of interest.

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