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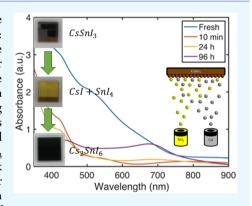
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## Optical Characterization of Lead-Free Cs<sub>2</sub>Snl<sub>6</sub> Double Perovskite Fabricated from Degraded and Reconstructed CsSnl<sub>3</sub> Films

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Supporting Information

ABSTRACT: Halide perovskites have experienced a huge development in the past years, but they still have two major challenges for their massive implantation: the long-term stability and the use of lead. One of the most obvious lead-free candidates to replace these perovskites is CsSnI<sub>3</sub>, but due to its poor environmental stability, it has been discarded for the fabrication of stable devices. Nevertheless, ambient degradation of CsSnI<sub>3</sub> and ulterior reconstruction produce a relatively stable lead-free Cs<sub>2</sub>SnI<sub>6</sub> double perovskite with interesting optical properties that have not been deeply characterized previously. In this work, the potential use for the optical properties of Cs<sub>2</sub>SnI<sub>6</sub> is studied and compared with that of the most common halide perovskite, CH3NH3PbI3 (MAPbI<sub>3</sub>). The Cs<sub>2</sub>SnI<sub>6</sub> films stayed in a standard atmosphere for a week without showing any signs of degradation. They also demonstrated better reflective behavior than MAPbI<sub>3</sub> and higher absorption in the 650 and 730 nm spectral range, making this material interesting for the development of



photodetectors in this region. This study demonstrates that Cs<sub>2</sub>SnI<sub>6</sub> is a promising material for photodevices, as it highlights its main characteristics and optical parameters, giving an original view on the use of the double perovskite, but at the same time emphasizing the need to improve the electrical properties for the development of efficient optoelectronic devices.

KEYWORDS: optical characterization, lead-free, perovskite, coevaporation deposition, tin halide, full inorganic

#### INTRODUCTION

Lead halide perovskites (APbX<sub>3</sub>, with A being an organic/ inorganic cation and X an halide anion) have experienced a huge development in past years, due to their suitability for the fabrication of solar cells, going from an initial 3.81% power conversion efficiency<sup>1</sup> to the recently achieved 25.2%.<sup>2</sup> These perovskites have demonstrated their great absorption capabilities, high photoluminescence quantum yield, and cheap fabrication processes. All these facts make them great candidates not only for the solar cell world but also for other optoelectronic applications such as solid light emitters<sup>3,4</sup> or as photodetectors.<sup>5-7</sup> However, these systems present a significant drawback facing their commercialization due to the toxicity of the lead. Consequently, a lot of studies focus on the development of lead-free perovskites with characteristics similar to those of their lead counterparts.8-11

There are a lot of possible lead-free candidates. 12 However, tin halide perovskites have drawn the attention of the scientific community, due to their similarities with the lead halide perovskites. 13,14 In particular, the fully inorganic CsSnI<sub>3</sub> achieved 0.88% efficiency in its first use as a Schottky solar cell. 15 Chung et al. studied the electrical properties of CsSnI<sub>3</sub>, discussing its metallic behavior instead of a semiconducting one. 16 They argued that this behavior is due to the tin vacancies produced in the structure during its formation. 17-19 Following that, Marshall et al. included a 10% excess of  $\rm SnI_2$  in the formation of the film, or 20% of  $\rm SnF_2^{~20}$  to cover the tin vacancies, achieving an efficiency around 2-2.76%. 21,2

The main drawback of this perovskite is the poor stability in environmental conditions, because the Sn<sup>2+</sup> is easily oxidized to Sn<sup>4+</sup>. Lee et al. tried to avoid this oxidation using SnF<sub>2</sub> as a reducing agent in the formation of the formamidinium tin iodide (FASnI<sub>3</sub>) perovskite, achieving a power conversion efficiency of 4.8%. 23 Recently, Chen et al. mixed the CsSnI<sub>3</sub> with germanium (Ge), making a "double perovskite" (Cs<sub>2</sub>SnGeI<sub>6</sub>) with 7.11% efficiency.<sup>24</sup> However, a majority of these studies were carried under a nitrogen atmosphere. Under these conditions, Sn<sup>4+</sup> arises, transforming CsSnI<sub>3</sub> perovskite into Cs<sub>2</sub>SnI<sub>6</sub> double perovskite. Nevertheless, this material itself constitutes a more air stable candidate than its nonoxidized counterpart.<sup>25</sup> It has been previously used as a hole transport layer (HTL) in dyesensitized solar cells, 26-28 and also as nanowires 29 or quantum rods.<sup>30</sup> Several synthesis methods have been employed for its preparation. 25,31,32 However, the optical properties of this material have not been broadly studied in the literature. In this work, the optical properties of Cs<sub>2</sub>SnI<sub>6</sub>, obtained as the final

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degradation product of CsSnI<sub>3</sub> in air, are studied in depth. In order to compare the properties of the thin film material and of the bulk, the powder is also synthesized with a novel and easy precipitation method. Moreover, with a comparative analysis of Cs<sub>2</sub>SnI<sub>6</sub> and of the widely used perovskite, the methylammonium lead iodide (MAPbI<sub>3</sub>) properties show the potential for Cs<sub>2</sub>SnI<sub>6</sub> not only in terms of ambient stability but also from the optical point of view. With this study, we intend to pave the way for further electric studies of this promising lead-free material, targeting the fabrication of efficient optoelectronic devices.

#### EXPERIMENTAL SECTION

Thin Film Fabrication. The CsSnI<sub>3</sub> thin film was deposited using a coevaporation technique, as can be schematically seen in Figure 1. We

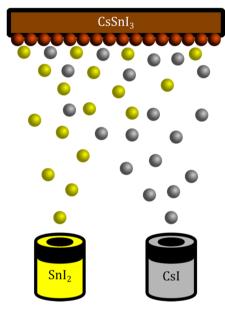


Figure 1. Scheme of the thermal coevaporation fabrication process for the preparation of CsSnI<sub>3</sub>.

used an MBraun thermal evaporator, filling two crucibles with precursors: one with CsI (99.9%, Sigma-Aldrich) and the other one with SnI<sub>2</sub> (99.9%, Sigma-Aldrich). The glass substrates were previously cleaned with a soapy water ultrasonic bath, rinsed with MiliQ water, ultrasonicated with ethanol and isopropanol (1:1), and then ultrasonicated again with ethanol and acetone (1:1).

The coevaporation of the binary salts leads to the self-assembly of the two precursors and to the formation of the lead-free perovskite (eq 1):

$$1CsI + 1SnI_2 \rightarrow CsSnI_3 \tag{1}$$

The coevaporation process involved heating the two crucibles simultaneously under vacuum, and then sublimation until the desired deposition rate was reached for both precursors.

In order to achieve the correct stoichiometry (1:1) in the final film, it is convenient to calculate the deposition rate ratio. It is, in fact, reasonable that the material volume ratio is proportional to the film thickness, and in turn to the evaporation rate. Thus, the volume ratio, depending on the molecular weight and densities (eq 2), is equal to the rate ratio.

$$\frac{1 \text{ mol CsI}}{1 \text{ mol SnI}_2} \rightarrow \frac{259.81 \text{ g CsI}}{372.52 \text{ g SnI}_2} \rightarrow \frac{57.61 \text{ mL CsI}}{49 \text{ mL SnI}_2}$$

$$= \frac{\text{deposition rate}_{\text{CsI}}}{\text{deposition rate}_{\text{SnI}_2}} = 1.176$$
(2)

This means that, to achieve the correct stoichiometry, the deposition rate of CsI must be 1.176 times bigger than that of SnI<sub>2</sub>. With this simple calculation, we can customize the ratio between the two precursors and make films with CsI or SnI<sub>2</sub> excess, with different properties. <sup>22,33</sup> The optimized temperature values for each crucible are ~270 °C for SnI<sub>2</sub> and ~510 °C for CsI. All these coevaporations took place with the evaporator at  $3 \times 10^{-6}$  Torr pressure, inside a nitrogen filled glovebox.

Reconstruction Conditions. The conversions to Cs<sub>2</sub>SnI<sub>6</sub> were carried out in the dark, with a temperature between 21 and 23 °C and relative humidity of 28-35%. The samples were stored in aluminum wrapped boxes in laboratory conditions, in air.

Cs<sub>2</sub>SnI<sub>6</sub> Powder Synthesis. Pure Cs<sub>2</sub>SnI<sub>6</sub> powder was synthesized with the following procedure. A 1 mmol portion of CsI (259.81 mg) was mixed with 1 mmol of SnI<sub>2</sub> (372.52 mg) in a vial. A 1 mL portion of acetone was added to the mixture to form a black suspension, which gently stirred at 750 rpm for a couple of hours, in order to completely disperse the precursors. After the evaporation of the acetone in ambient conditions, it was possible to collect the precipitated black crystals of

Materials Characterization. X-ray diffraction (XRD) measurements of the powders and the thin films were taken by an X-ray diffractometer (D8 Advance, Bruker-AXS) (Cu K $\alpha$ , wavelength  $\lambda$  = 1.5406 Å). A UV-vis spectrophotometer (Cary 300 UV-vis, Agilent) was used to measure the absorbance, reflectance, and transmittance of the films. Photoluminescence (PL) measurements were done using a spectrofluorometer (Fluorolog, Horiba). The photoluminescence quantum yield (PLQY) measurement was taken using an integrating sphere (Hamamatsu). Thicknesses of the films were obtained using a profilometer (Dektak 6M, Veeco). All of the microscopy was done using a scanning electron microscope (SEM) (JSM-7001F, JEOL) which was also equipped with energy dispersive X-ray spectroscopy (EDS).

#### ■ RESULTS AND DISCUSSION

CsSnI<sub>3</sub> samples prepared on glass substrates by coevaporation (see Experimental Section) inside the glovebox were characterized under ambient conditions. After 10 min of ambient air exposure, samples experienced a fast degradation to their precursors (Figure 2a) with a decrease of the absorbance (Figure 2b). Accordingly, their color changed from dark brown (CsSnI<sub>3</sub>, left) to yellow (CsI + SnI<sub>2</sub>, center). The poor stability of the CsSnI<sub>3</sub> in air is mainly due to the decomposition (eq 3) and oxidation (eq 4) mechanisms.

$$CsSnI_3$$
 + environmental conditions  $\rightarrow CsI + SnI_2$  (3)

$$2SnI_2 + O_2 \rightarrow SnI_4 + SnO_2 \tag{4}$$

The oxidation from Sn<sup>2+</sup> to Sn<sup>4+</sup> leads to the phase conversion to Cs<sub>2</sub>SnI<sub>6</sub>, in which half of the tin atoms are missed, because they form tin oxide. The complete reaction is shown in eq 5.

$$2CsSnI_3 + O_2 \rightarrow Cs_2SnI_6 + SnO_2$$
 (5)

This CsSnI<sub>3</sub> degradation and Cs<sub>2</sub>SnI<sub>6</sub> reconstruction has been demonstrated by Qiu et al.<sup>34</sup> but with some differences. In our case, the degradation and reconstruction take between 3 and 5 days to get the film completely dark (Figure 2a, right) as compared to the 24 h needed in that study. This enhanced stability of initial CsSnI3 is attributed to better quality of the deposition method employed here, as they reported a two-step sequential deposition by alternating thermal evaporation depositions, followed by a thermal annealing to provoke a solid-state reaction between the layers. This method increases the chance of getting pinholes and an inhomogeneous morphology, compared to the film obtained by thermal coevaporation.<sup>34</sup> The absorbance spectra reported in Figure

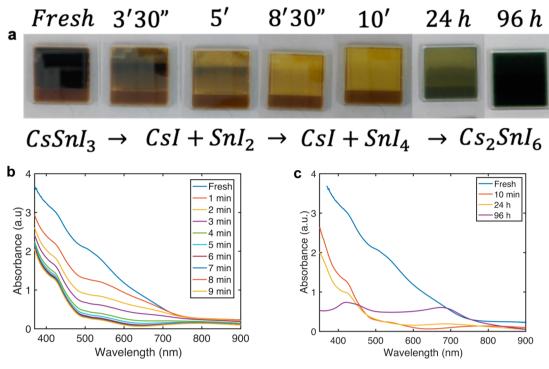


Figure 2. (a) Degradation, oxidation of CsSnI<sub>3</sub>, and reconstruction of the Cs<sub>2</sub>SnI<sub>6</sub> thin film. (b) Absorbance decay due to CsSnI<sub>3</sub> degradation. (c) Absorbance recovery due to formation of Cs<sub>2</sub>SnI<sub>6</sub>.

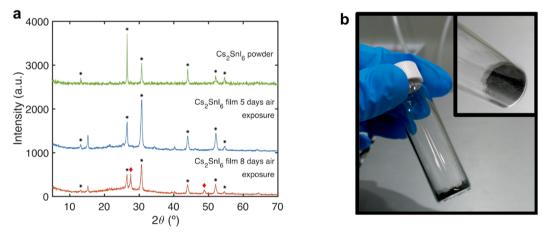


Figure 3. (a) XRD patterns of the pure powder and the thin films of  $CsSnI_3$  after 5 days and 8 days in air. After the 5th day the  $CsSnI_3$  film is totally converted into  $Cs_2SnI_6$ . Black stars show  $Cs_2SnI_6$  peaks, and red diamonds show the CsI peaks. (b) Picture of the  $Cs_2SnI_6$  pure powder.

2c show that the film recovers its capacity to absorb the wavelengths between 700 and 800 nm completely, but it loses the absorption capacity in shorter wavelengths and has a planar response with two humps, instead of the original ramp shape. This characteristic behavior, resulting in an average visible transmittance (AVT) of 42.1%, doubling the value of the AVT reported for 150 nm MAPI films, 35,36 is attractive for a semitransparent application.<sup>37</sup> In this case the thickness, measured using a profilometer and by SEM cross section, giving a result of 230 nm (Figure S1a,b), was not sacrificed to obtain a high transmittance value, avoiding the problems related to the contact between the deposited layer above and below the perovskite, leading to a short circuit in optoelectronic devices. Note that slower degradation of CsSnI<sub>3</sub> is observed in the inferior side of the analyzed sample, Figure 2a, where FTO (fluorine doped SnO<sub>2</sub>) is exposed in the substrate. This fact points to a possible effect of fluorine in FTO as a reducing  ${\rm agent}^{23}$  in the enhanced stability of  ${\rm CsSnI_3}$  in this region. However, deep analysis of this observation is beyond the scope of this work.

To compare the properties in bulk, we synthesized the  $Cs_2SnI_6$  powder following the procedure explained in the Experimental Section. The X-ray diffraction (XRD) patterns of the powder and of the thin film are shown in Figure 3a. Here, we considered the powder (Figure 3b) as a reference pure material, with characteristic peaks at  $13.16^\circ$ ,  $26.52^\circ$ ,  $30.72^\circ$ ,  $44^\circ$ ,  $52.04^\circ$ , and  $54.68^\circ$ , all of them representing lattice planes of  $Cs_2SnI_6$ .  $^{26,32,33}$  After comparing these peaks with a film exposed to ambient air for 5 days, it is possible to recognize similar patterns in the degraded and reconstructed samples. However, 3 days after this, another two peaks have arisen, at  $27.6^\circ$  and  $48.76^\circ$ , both belonging to CsI, one of the precursors. This fact implies that the films maintain their air stability for 1 week, approximately. Then, they start decomposing into their

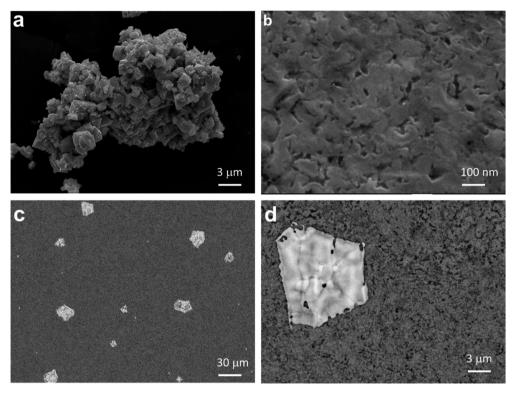


Figure 4. Scanning electron microscopy (SEM) images of (a) pure powder sample and (b)  $Cs_2SnI_6$  samples. (c) After 11 days on the SEM, the top view of the  $Cs_2SnI_6$  CsI crystals became visible. (d) Zoom-in of the CsI crystal detail.

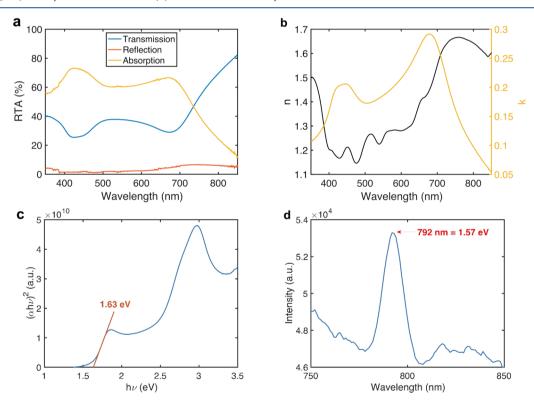
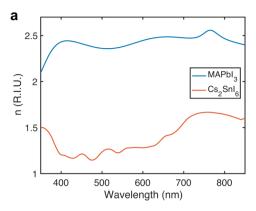


Figure 5. Optical characterization of  $Cs_2SnI_6$  thin films: (a) reflection, transmission, and absorption measurements; (b) refractive index and extinction coefficient values; (c) Tauc plot determining the band gap; and (d) photoluminescence peak matching with the band gap.

precursors again. This highlights that, despite the  $\mathrm{Cs_2SnI_6}$  double perovskite presenting an enhanced stability compared with most of its perovskite counterparts, it also requires encapsulation for long-term stability.

To double-check the presence of the CsI, we measured the morphology by scanning electron microsopy (SEM) (Figure 4). The pure powder presents crystals with tiny bright dots in the secondary electron image but not in the backscattered one. In



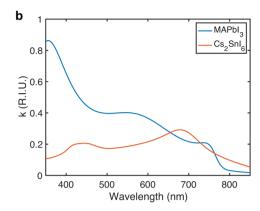


Figure 6. (a) Refractive index and (b) extinction coefficient comparison between classical MAPbI<sub>3</sub> and lead-free (Cs<sub>2</sub>SnI<sub>6</sub>) perovskite.

fact, the crystal and the bright dots were analyzed using energy-dispersive X-ray spectroscopy (EDS) measurements, with both giving the same result,  $Cs \approx 22\%$ ,  $Sn \approx 11\%$ , and  $I \approx 66\%$ , matching the  $Cs_2SnI_6$  structure (Figure S2a). Figure 4b shows that the surface of the thin film before the final decomposition is rough. Although it was not possible to measure the  $CsSnI_3$  film, because of the fast degradation, we are confident that the morphology is homogeneous, as the coevaporation gives a perfect mirror-like film (Figure S3). Consequently, the roughness observed in  $Cs_2SnI_6$  is likely generated during the oxidation and reconstruction process, as briefly mentioned above.

In Figure 4c,d, we can see the surface of the Cs<sub>2</sub>SnI<sub>6</sub> thin film starting the degradation process, because crystals of CsI start appearing. The composition of the crystals was confirmed by EDS measurements (Figure S2b).

Optical properties of the Cs<sub>2</sub>SnI<sub>6</sub> thin films prepared by degradation and reconstruction were systematically analyzed. We measured and calculated the absorption coefficient, band gap, refractive index, and extinction coefficient.

The absorption, reflection, and transmission of the films were measured (Figure 5a), and the corresponding values were used to obtain the absorption coefficient,  $\alpha$ , using the relation below:<sup>38</sup>

$$\alpha = \frac{1}{t} \ln \left[ \frac{(1-R)^2}{T} \right] \approx 2.303 \frac{A}{t}$$
 (6)

Here, *t* is the thickness of the film, and *R*, *T*, and *A* are reflectance, transmittance, and absorbance of the film, respectively. The approximation showed in eq 6 is accepted if the reflectance is close to zero. In this study, the approximation matches the real case, due to the low reflection of the film (Figure 5a).

The values of refractive index and extinction coefficient as a function of wavelength, see Figure 5b, were obtained following eqs 7 and  $8:^{38}$ 

$$\kappa = \frac{\alpha \lambda}{4\pi} \tag{7}$$

$$n = \frac{-2(R+1) \pm \sqrt{(2(R+1))^2 - 4(R-1)^2(1+\kappa^2)}}{2(R-1)}$$

In these equations,  $\lambda$  is the wavelength, R the measured reflectance, n the refractive index, and  $\kappa$  the extinction coefficient.

Cs<sub>2</sub>SnI<sub>6</sub> thin films presented a band gap of 1.63 eV, see Figure 5c, as it was calculated using the linear interpolation of the Tauc plot. This plot compares  $(\alpha h \nu)^n$  (n=2, considering Cs<sub>2</sub>SnI<sub>6</sub> as a direct band gap material) against the energy of the photons  $(h \nu)$ , where h is Planck's constant and  $\nu$  the photon's frequency. Previous literature studies provide a wide range of band gap values for Cs<sub>2</sub>SnI<sub>6</sub>, between 1.2 and 1.6 eV.<sup>25–28,32,34</sup> Our work is in good agreement with studies pointing to a high band gap in this range. <sup>25,27,32,34</sup> Regarding the optical band gap, the photoluminescence (PL) of the thin films results in the emission peak (Figure 5d) at 792 nm. Translated to energy, this value is 1.57 eV, which matches with the result obtained from the Tauc plot, with a relatively small Stokes shift. No significant PL quantum yield was measured for this samples, pointing to a high nonradiative recombination of the material.

In order to put this material in context with other materials of the perovskite family, we compared it with the most studied perovskite, MAPbI<sub>3</sub>. Figure 6a,b shows the comparison of refractive index and extinction coefficient between both materials. Two main aspects are remarkable here: the first one, regarding n, is the huge difference between the MAPbI<sub>3</sub> (around 2.5) and the Cs<sub>2</sub>SnI<sub>6</sub> (around 1.5) values. This fact can make Cs<sub>2</sub>SnI<sub>6</sub> better than MAPbI<sub>3</sub> in terms of reflection and interface index matching, depending on the surrounding layers. The second important fact is regarding κ, as MAPbI<sub>3</sub> can be considered a better absorber than Cs2SnI6; nevertheless, between 650 and 730 nm, Cs<sub>2</sub>SnI<sub>6</sub> has a higher extinction coefficient than MAPbI3, which implies that it is a better absorber in this spectral interval. This fact points to a good potential for Cs<sub>2</sub>SnI<sub>6</sub> for the development of photodetectors in this range.

With these results it is also possible to calculate the ideal short-circuit current density that a solar cell made with  $Cs_2SnI_6$  could have. We have done this using the generation rate and by supposing an ideal case with no electrical losses.<sup>40</sup>

$$J_{SC} = q \int_{0}^{T} G(x) dx$$

$$J_{SC} = q \int_{0}^{T} \int_{\lambda_{\min}}^{\lambda_{\max}} \frac{\lambda}{hc} Irr(\lambda) [1 - r(\lambda)] \alpha(\lambda) e^{-\alpha(\lambda)x} d\lambda dx$$

$$J_{SC} = \frac{q}{hc} \int_{\lambda_{\min}}^{\lambda_{\max}} \lambda Irr(\lambda) [1 - r(\lambda)] (1 - e^{-\alpha(\lambda)T}) d\lambda$$

In these equations, q is the electron charge, G(x) the generation rate, h Planck's constant, c the speed of light,  $\lambda$  the incident wavelength in vacuum,  $\alpha$  the absorption coefficient, T the

thickness of the film, r the reflectance, and Irr the irradiance value of the incident light.

Using this equations, under the sun spectrum (AM1.5G), we obtain a  $J_{\rm SC}$  value of 18.86 mA/cm² in a device with an active layer thickness of 230 nm. The maximum value is obtained by supposing ideal conditions (zero reflectance and an optimum optical thickness), and it equals 33.68 mA/cm². So, making thicker devices will give us higher current density values (e.g., 26.01 mA/cm² at 500 nm thickness). In thicker cases, several problems start appearing, due to the increase in the distance that the carriers have to travel to reach the electrodes. Nevertheless, these electrical issues are far away from our actual study.

#### CONCLUSIONS

In summary, we have prepared a Cs<sub>2</sub>SnI<sub>6</sub> thin film from the degradation and reconstruction of CsSnI3 first prepared by coevaporation. The structural properties of the Cs<sub>2</sub>SnI<sub>6</sub> samples were compared with Cs<sub>2</sub>SnI<sub>6</sub> powders, prepared by a novel method, obtaining a good agreement. Cs<sub>2</sub>SnI<sub>6</sub> thin films exposed to ambient standard conditions of oxygen and moisture do not exhibit degradation for 1 week. We have successfully characterized the powder and films to obtain their crystalline structure, band gap, morphology, and optical properties (PL, absorption coefficient, band gap, refractive index and extinction coefficient). These values have been compared with MAPbI<sub>3</sub> perovskite, revealing that Cs<sub>2</sub>SnI<sub>6</sub> presents better reflective properties than MAPbI<sub>3</sub>. In addition, despite MAPbI<sub>3</sub> presenting higher absorption than Cs<sub>2</sub>SnI<sub>6</sub> in most of the visible wavelength range, this is not the case in the spectral range between 650 and 730 nm, where Cs<sub>2</sub>SnI<sub>6</sub> presents higher light absorption, making this material interesting for the development of photodetectors in this specific range or for semitransparent applications. Further research will be needed to improve the electrical properties of this material limited by high nonradiative recombination as we confirm by PL characterization. All these facts make this lead-free double perovskite, Cs2SnI6, a great candidate as an absorber for photodevices from an optical point of view and stress the need for further work on the improvement of the electrical properties.

#### ASSOCIATED CONTENT

#### **S** Supporting Information

The Supporting Information is available free of charge at https://pubs.acs.org/doi/10.1021/acsaem.9b01827.

Scanning electron microscopy cross section, energy-dispersive X-ray spectroscopy results, and pictures of the evaporated films (PDF)

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Notes

The authors declare no competing financial interest.

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