

Vertical coupling and transport in high-latitude ionosphere

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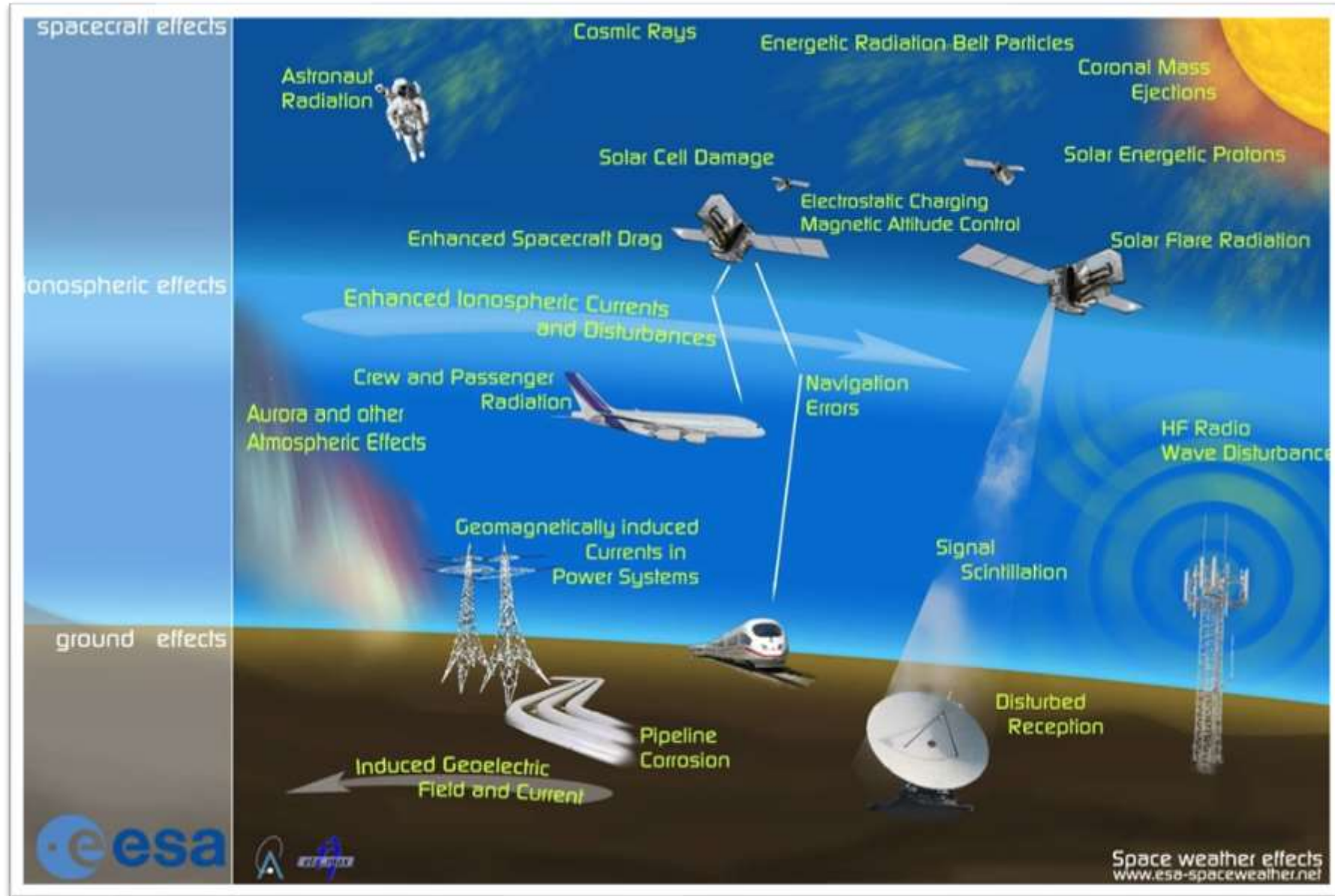
Seminar at the GFZ
Potsdam, 6-Sep-2022



Knowledge for Tomorrow



Intro: Space weather



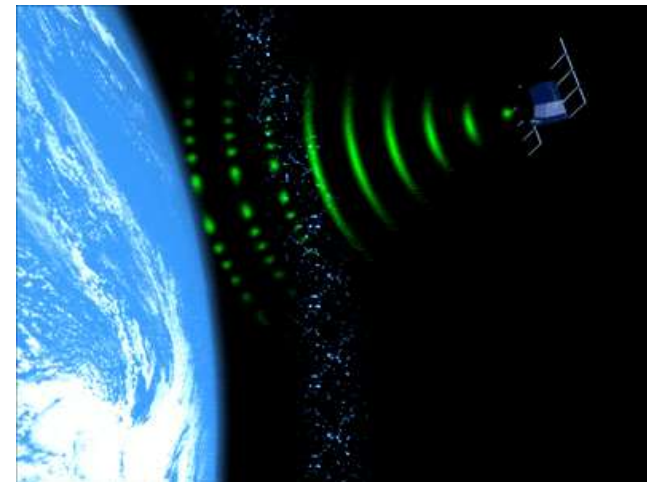
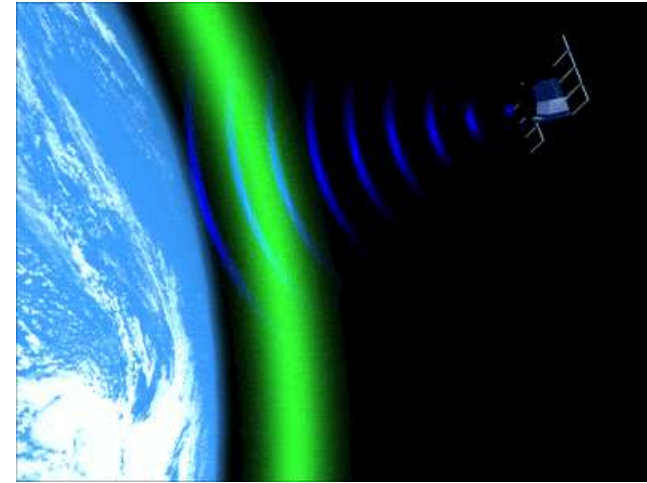
Ionospheric effects on GNSS

The ionosphere causes major problems for GNSS navigation:

Group delay to the signal propagation time that is proportional to the total electron content. This can change the apparent position by tens of metres.

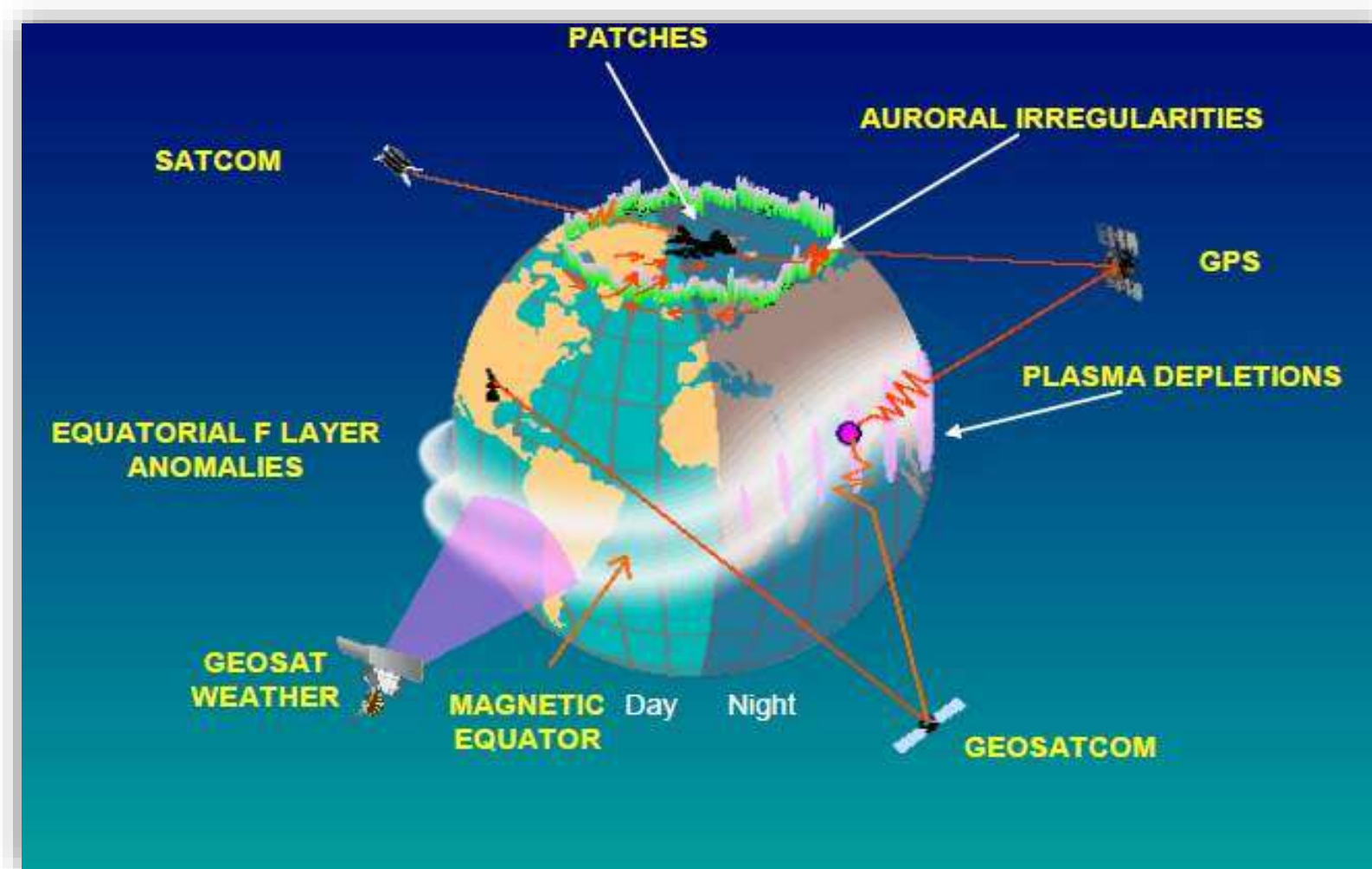
Scintillations of the signal are related to small-scale irregularities in plasma density. These can cause temporary loss of the signals.

While these effects may not be an issue for most GNSS users but they are important for safety-critical applications.



Credit: Univ. of Bath

Global SatCom Outage Regions



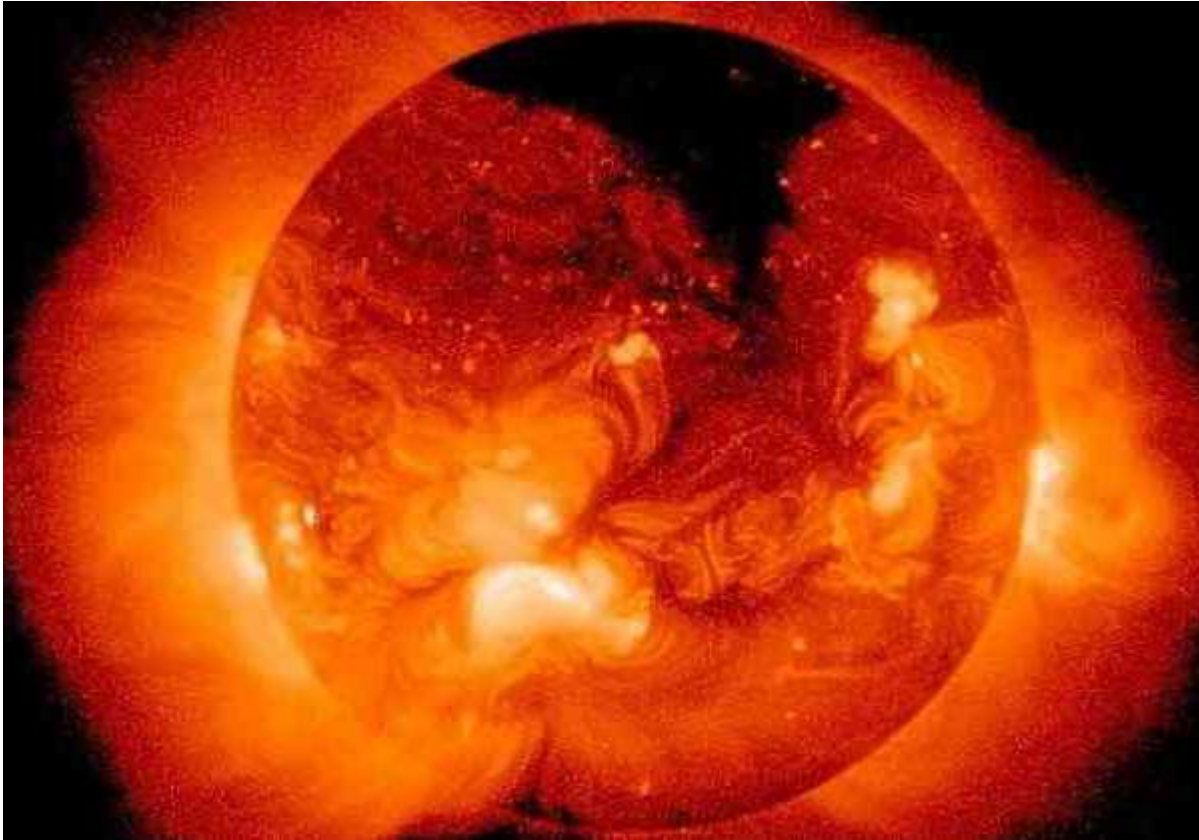
Courtesy of H.C. Carlson



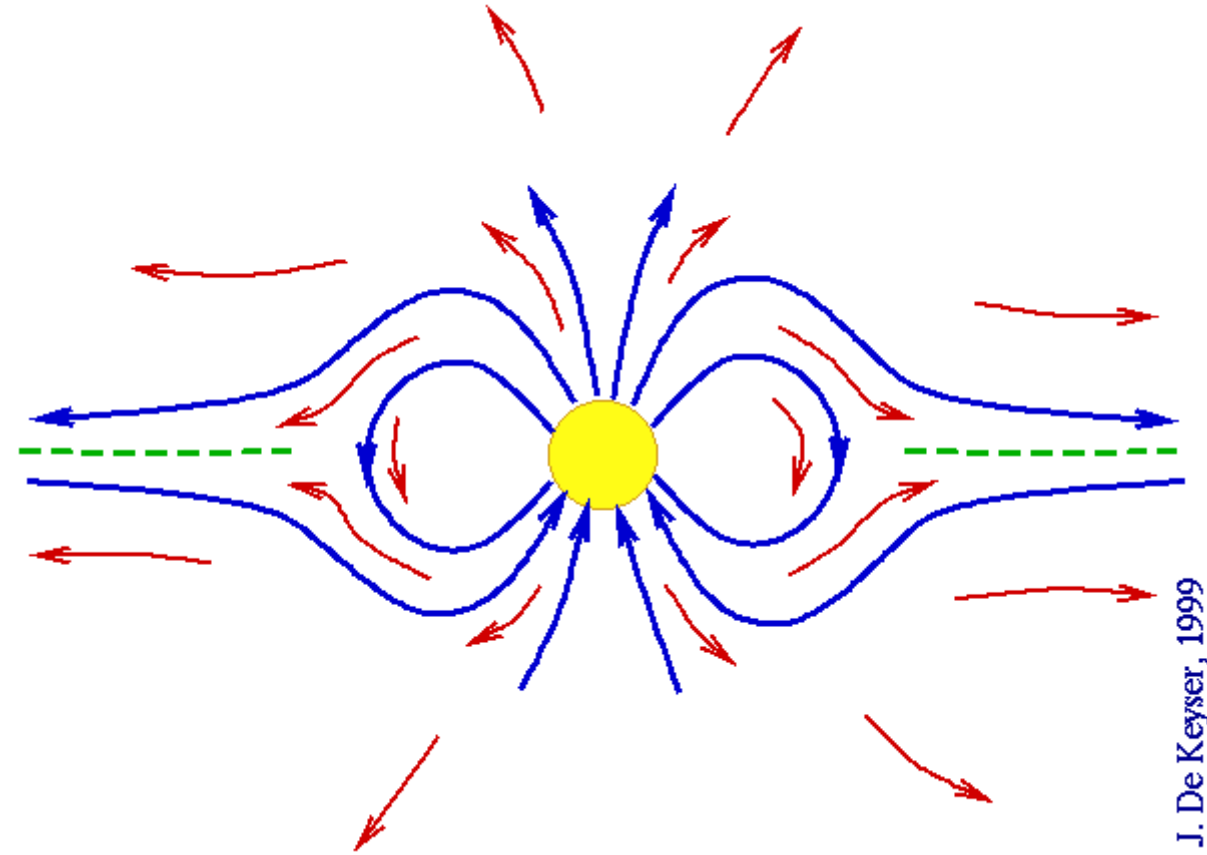
Origins of solar wind and interplanetary magnetic field (IMF)

Two distinct regions of the solar corona:

- Equatorial latitudes.
- Polar latitudes: coronal holes.



SOLAR-A probe image, courtesy of NASA



J. De Keyser, 1999

Diffusion and frozen flux

Assumptions:

- magnetised plasma;
- collisionless plasma (collision frequencies are much less than gyrofrequency, but not fully negligible);
- “cold” or “warm” plasma (particle energies substantially below relativistic, ~10s - 100s of eV).

From Faraday’s law and generalised Ohm’s law, eliminating \mathbf{E} field:

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{v} \times \mathbf{B} - \mathbf{j}/\sigma_0)$$

Using $\nabla \cdot \mathbf{B} = 0$ and neglecting displacement currents (as usual in plasma physics):

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{v} \times \mathbf{B}) + \frac{1}{\mu_0 \sigma_0} \nabla^2 \mathbf{B}$$

Advection
term

Diffusion
term



Magnetic diffusion

Assumptions:

- plasma at rest (or moving with constant velocity);
- conductivity is finite.

$$\frac{\partial \mathbf{B}}{\partial t} = D_m \nabla^2 \mathbf{B} \quad \leftarrow \text{only diffusion term left}$$

with the magnetic diffusion coefficient: $D_m = (\mu_0 \sigma_0)^{-1}$

Solution is given by: $B = B_0 \exp(\pm t/\tau_d)$

with the magnetic diffusion time: $\tau_d = \mu_0 \sigma_0 L_B^2$

Consider typical solar wind:

Density $\sim 5 \text{ cm}^{-3}$

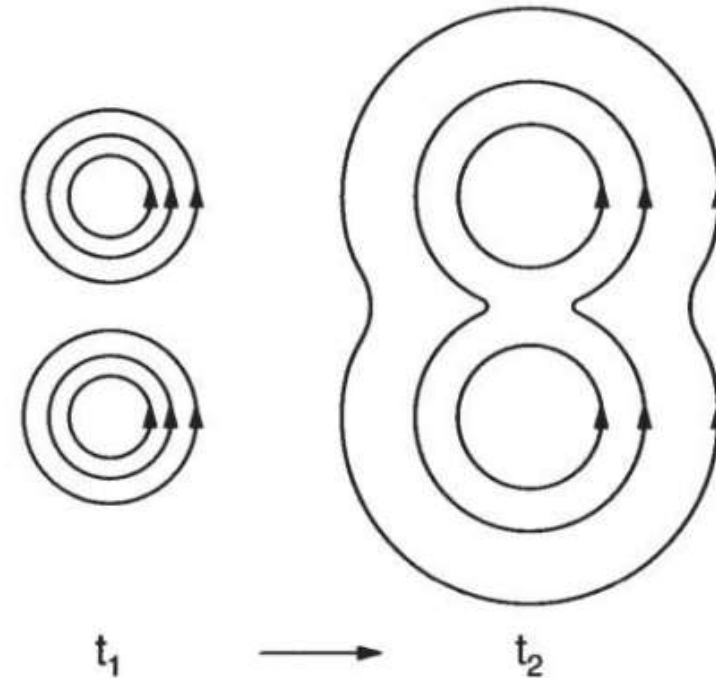
Temperature $\sim 50 \text{ eV}$

$$\sigma_0 = \frac{n_e e^2}{m_e \nu_c}$$

For the travel time from Sun to Earth ($\sim 3 \text{ days}$)

$$L_B \sim 10^3 \text{ m}$$

In solar wind the magnetic diffusion is negligible!



Frozen-in condition

Assumption:

- collisionless plasma with infinite conductivity ($\sigma \rightarrow \infty$)

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{v} \times \mathbf{B})$$

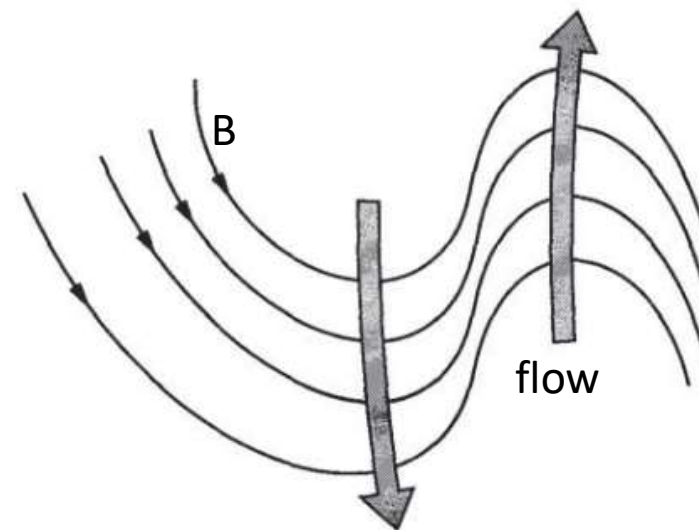
or equivalent

$$\mathbf{E} + \mathbf{v} \times \mathbf{B} = 0$$

Hydromagnetic theorem:

In infinitely conductive plasma, the total magnetic induction encircled by a closed loop remains unchanged

-> magnetic field lines are “frozen” into the plasma flow



Frozen magnetic field lines



Magnetic merging and reconnection

Back to the case of finite conductivity (both advection and diffusion terms present)

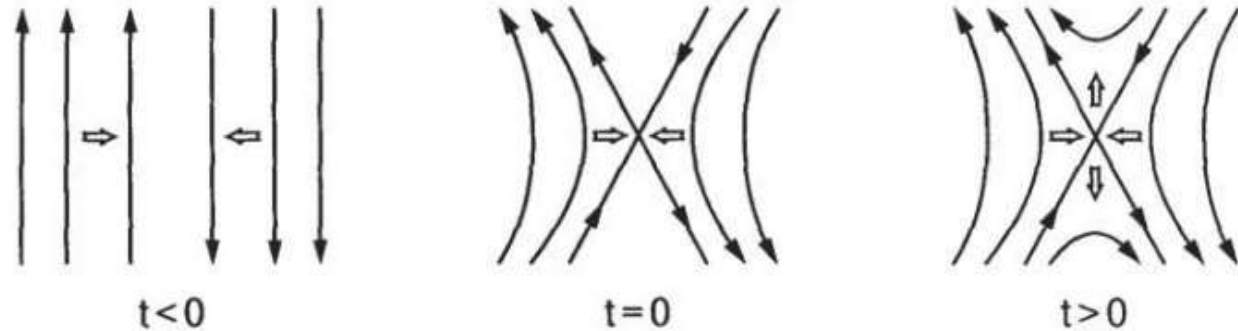
$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{v} \times \mathbf{B}) + \frac{1}{\mu_0 \sigma_0} \nabla^2 \mathbf{B}$$

or in dimensionless form

$$\frac{B}{\tau} = \frac{V B}{L_B} + \frac{B}{\tau_d}$$

Magnetic Reynolds number

$$R_m = \mu_0 \sigma_0 L_B V$$



Magnetic field line merging

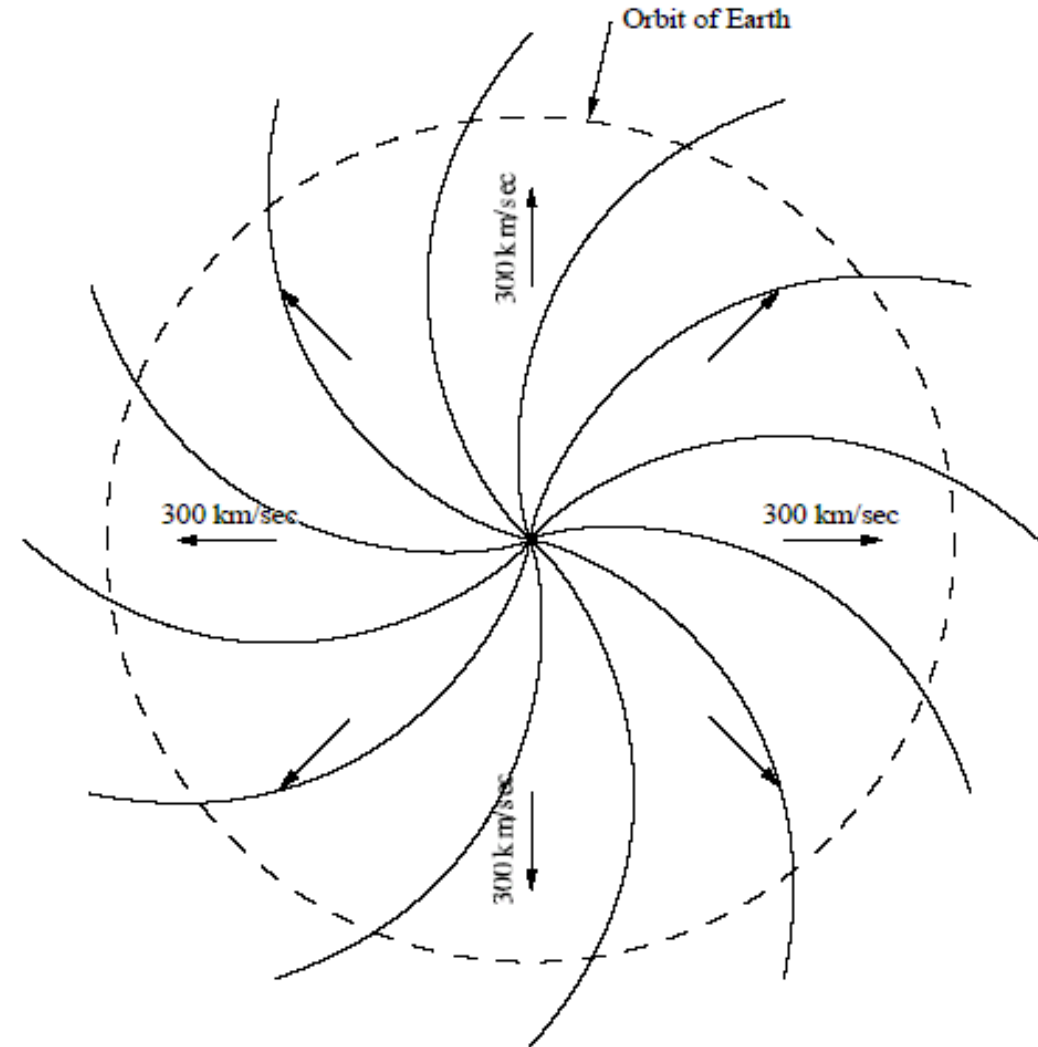
For $R_m \gg 1$: diffusion is negligible (in solar wind $R_m \approx 10^{17}$)

For $R_m \sim 1$: diffusion is substantial and can dominate -> **merging/reconnection**



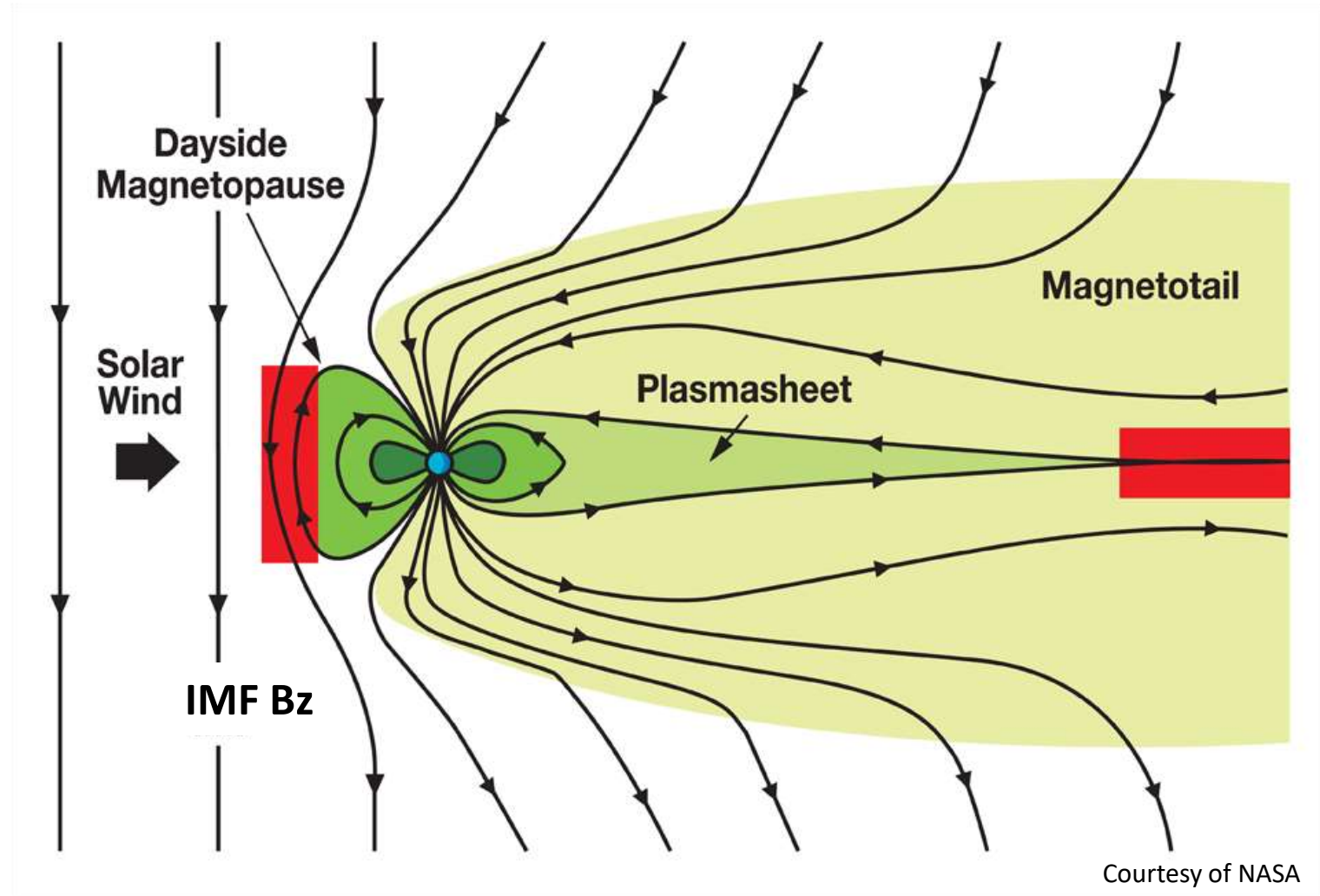
Interplanetary magnetic field (IMF): Parker spiral

- Solar rotation drags out frozen solar wind magnetic field forming **Parker spiral** (after Gene Parker)
- Winding angle depends on wind speed, but on average
- ~ 45 deg at Earth orbit
- ~ 90 deg 10 AU



Earth's Magnetosphere

- Magnetosphere is the region of space where the Earth's own magnetic field dominates.
- Under southward IMF conditions, the merging of magnetic field lines is possible at the nose of magnetosphere (dayside reconnection).
- Closed magnetic field lines have both ends linked to the Earth; open field lines have one end linked to the solar wind.

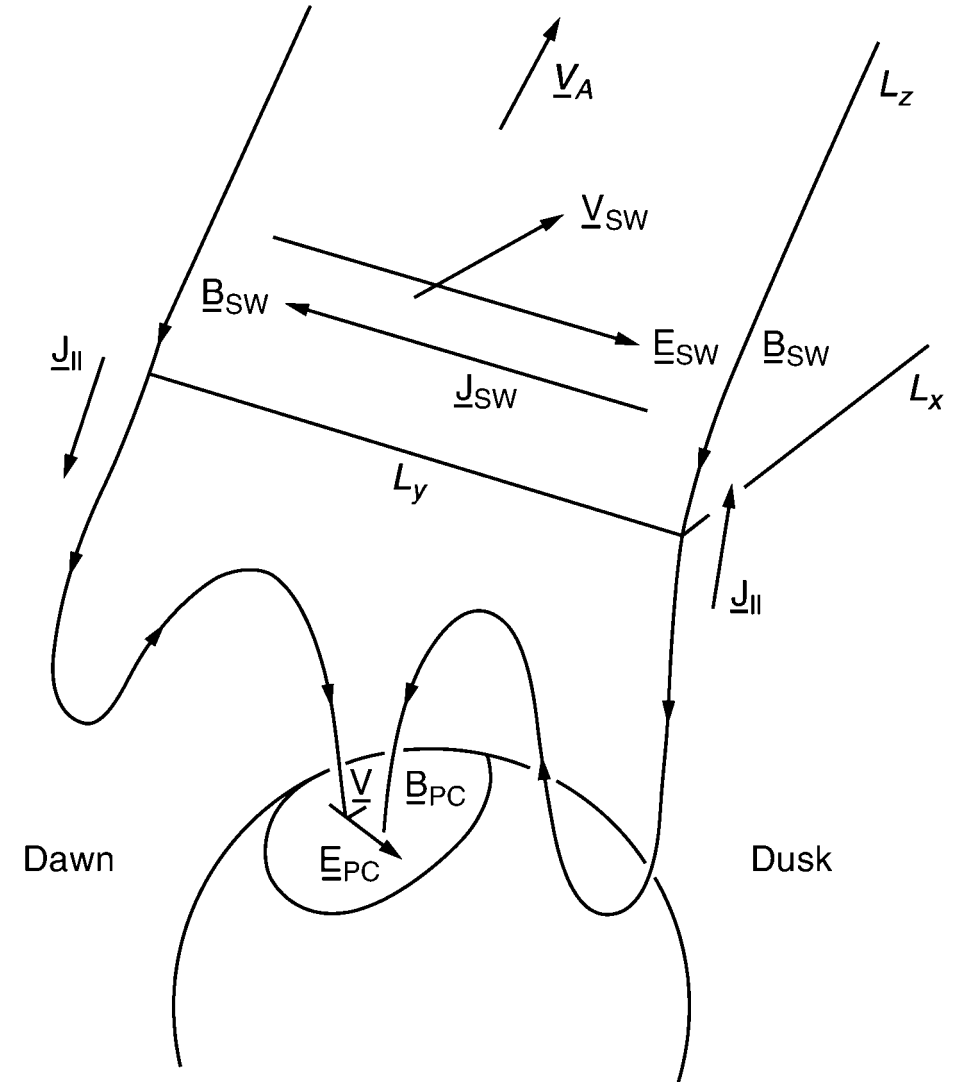


Coupling under southward IMF

Solar wind electric field (frozen-in plasma)

$$\mathbf{E}_{\text{SW}} = -\mathbf{V}_{\text{SW}} \times \mathbf{B}_{\text{SW}}$$

Under a southward IMF, a dawn-to-dusk solar wind electric field \mathbf{E}_{SW} maps to the polar cap ionosphere.



Satellite navigation systems

GPS system

~ 30 navigational GPS spacecraft

55° inclination orbits

Broadcast two frequencies in L1 and
L2 bands (1.57 and 1.23 GHz).

Other similar GNSS systems:

GLONASS, Galileo, Beidou

Ionosphere is a dispersive media

Total electron content (TEC) along the signal path
from GPS spacecraft to the receiver can be inferred
by measuring the phase advance and/or the group
delay between different frequencies.



GNSS tomography example-1

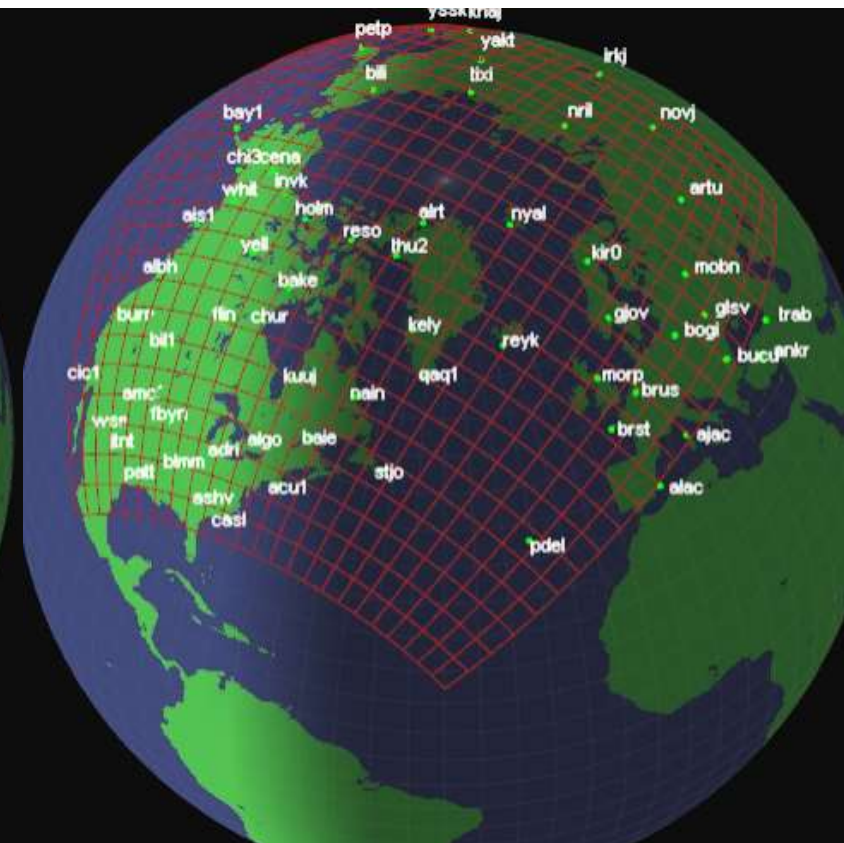
Network of ground dual-frequency GNSS receivers at high latitudes.

Tomographic inversion of the GNSS data should reveal plasma dynamics.

Network of GNSS receivers



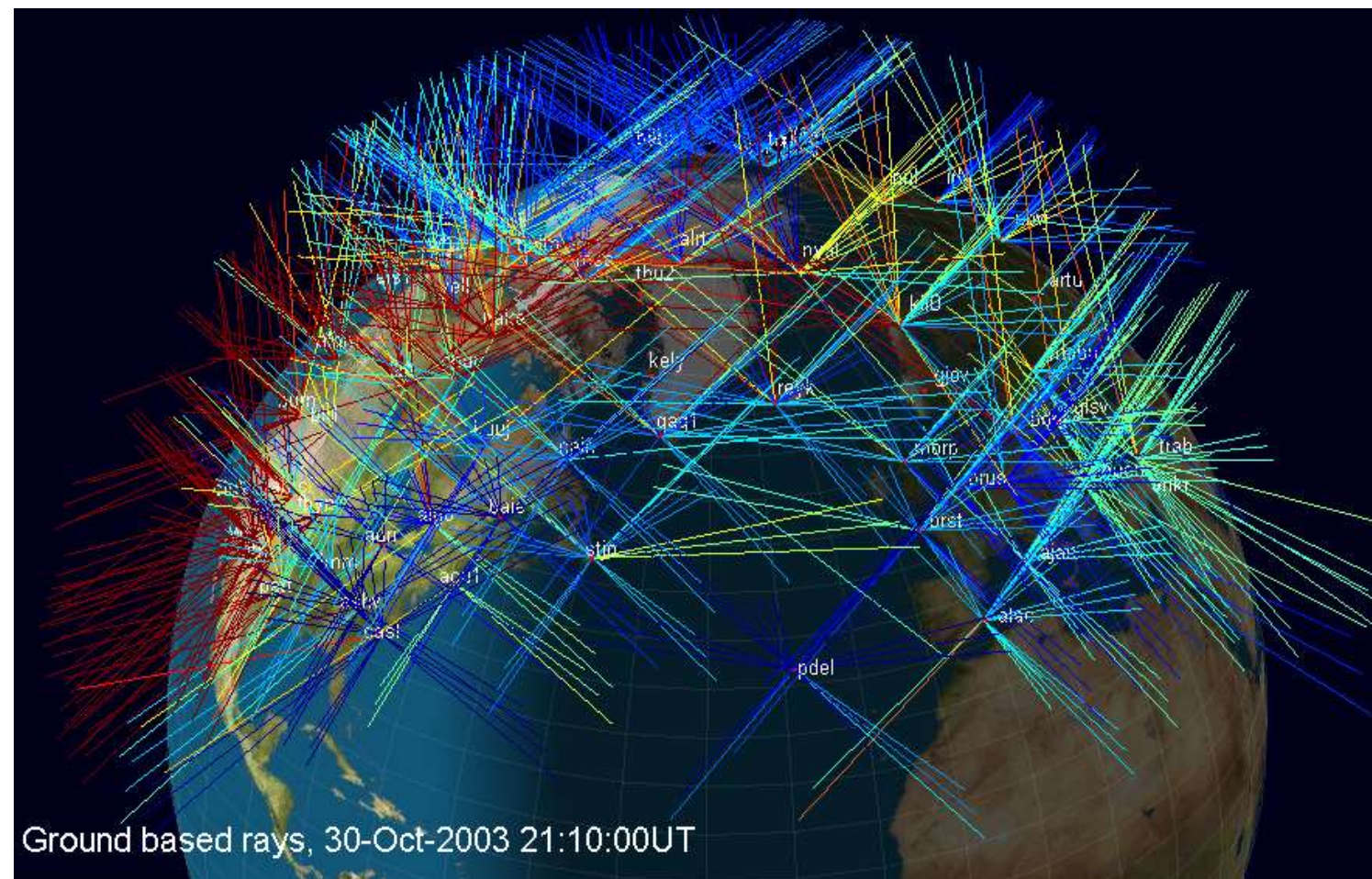
Projection of the tomographic grid



GNSS tomography example-2

Rays of GPS signals received during the major 30-Oct-2003 geomagnetic storm.

Colour shows plasma content along the ray.

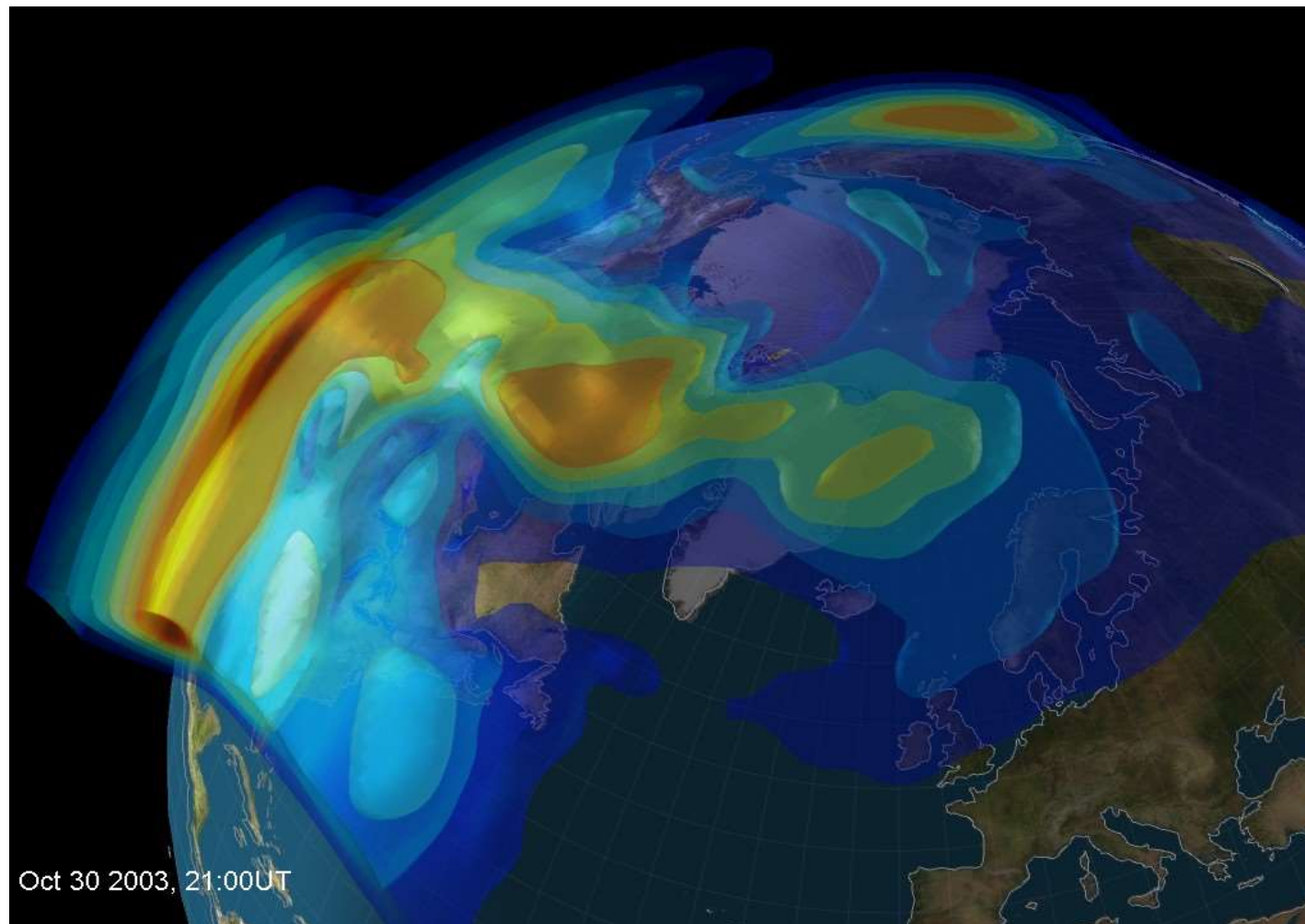


GNSS tomography example-3

Results of the tomographic reconstruction: plasma density

The global distribution of ionospheric plasma density can be deduced from characteristics of GPS signals acquired by ground-based network of GPS receivers.

Plasma follows general anti-sunward cross-polar convection.

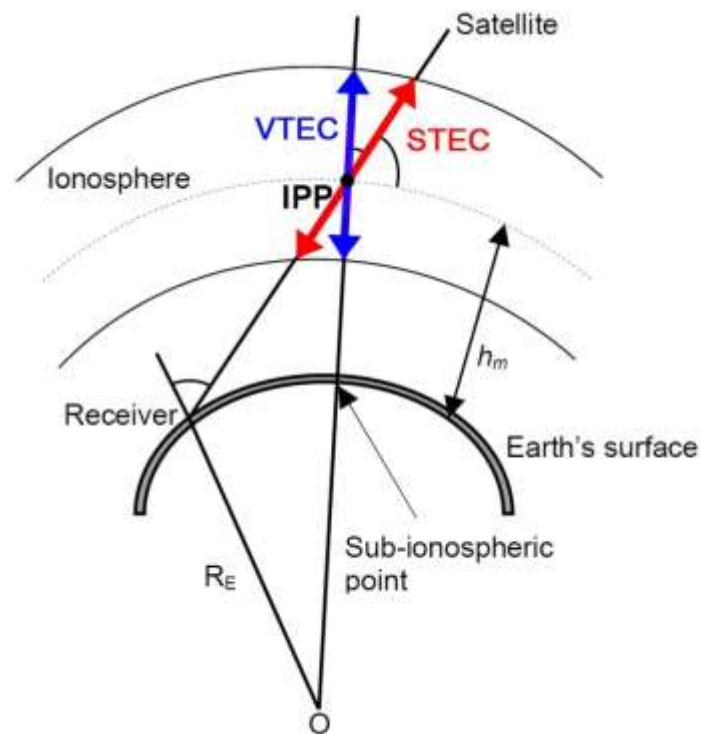


Mitchell et al, AGU Mono. 2008, doi:10.1029/181GM09

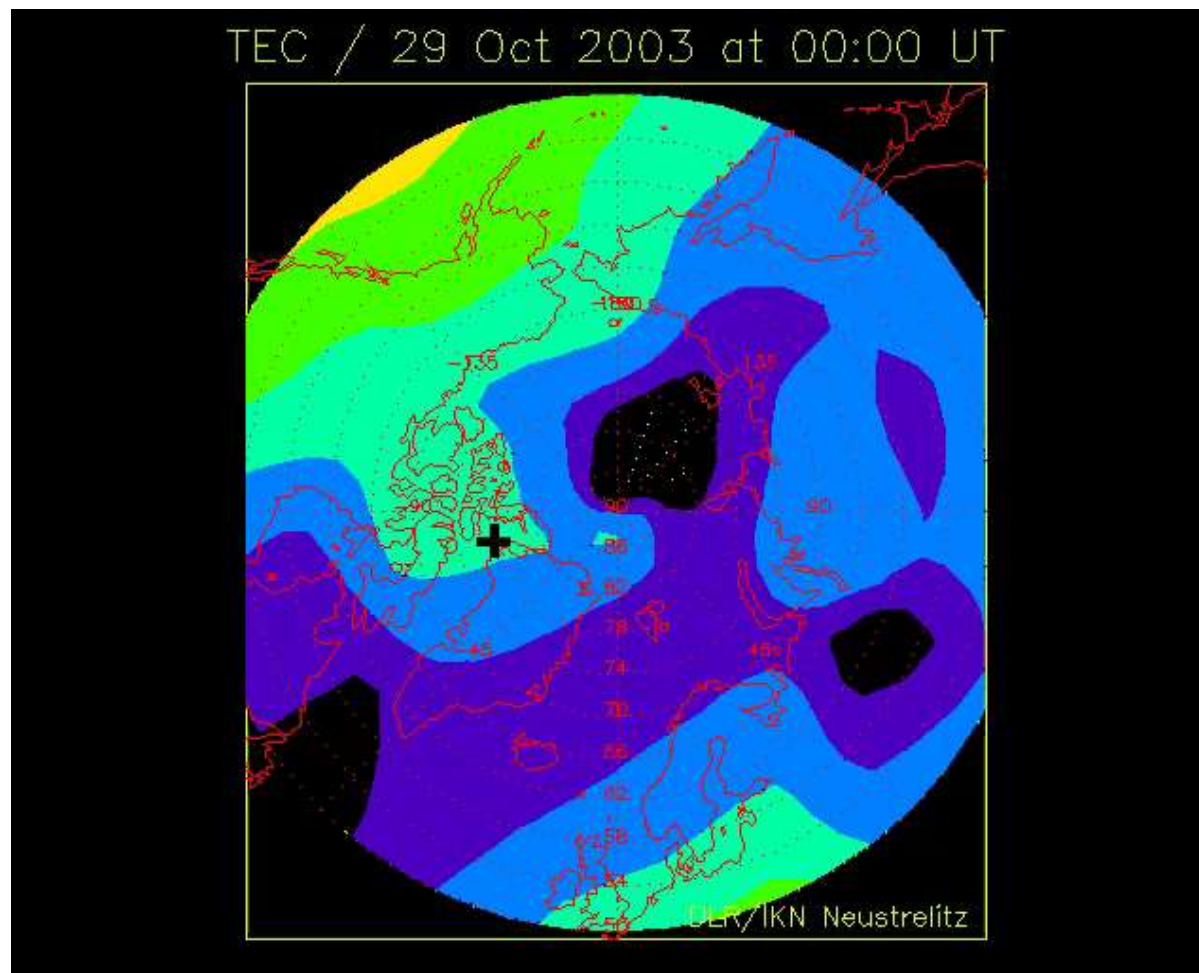
Direct reconstruction of TEC

Simpler direct reconstructions are also possible.

Only assumptions about ionospheric equivalent height (centre of mass) are needed.



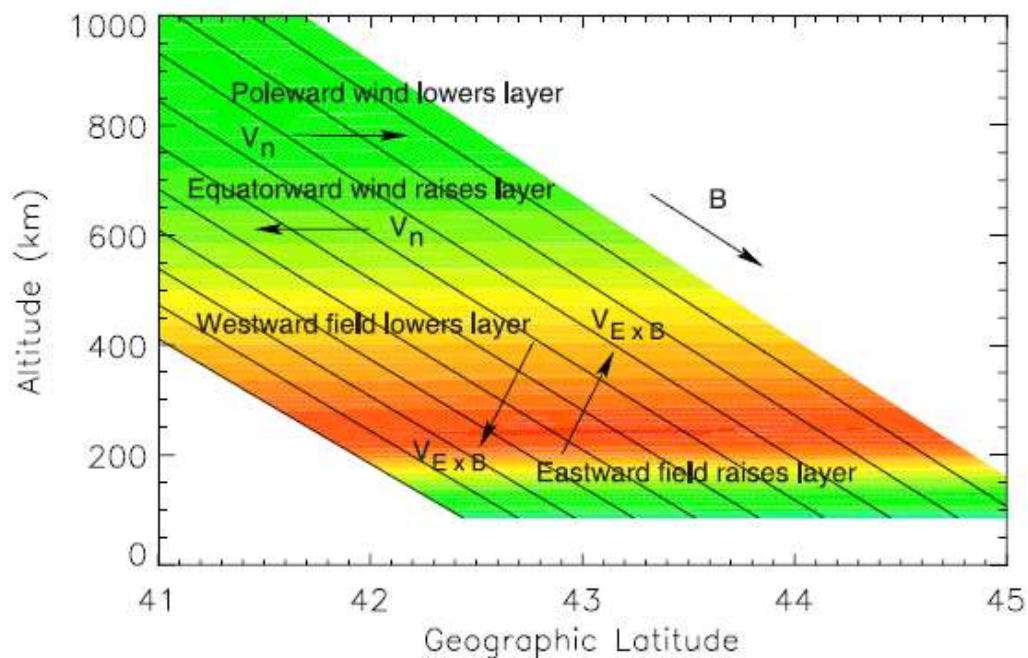
From Jaacob et al., 2008



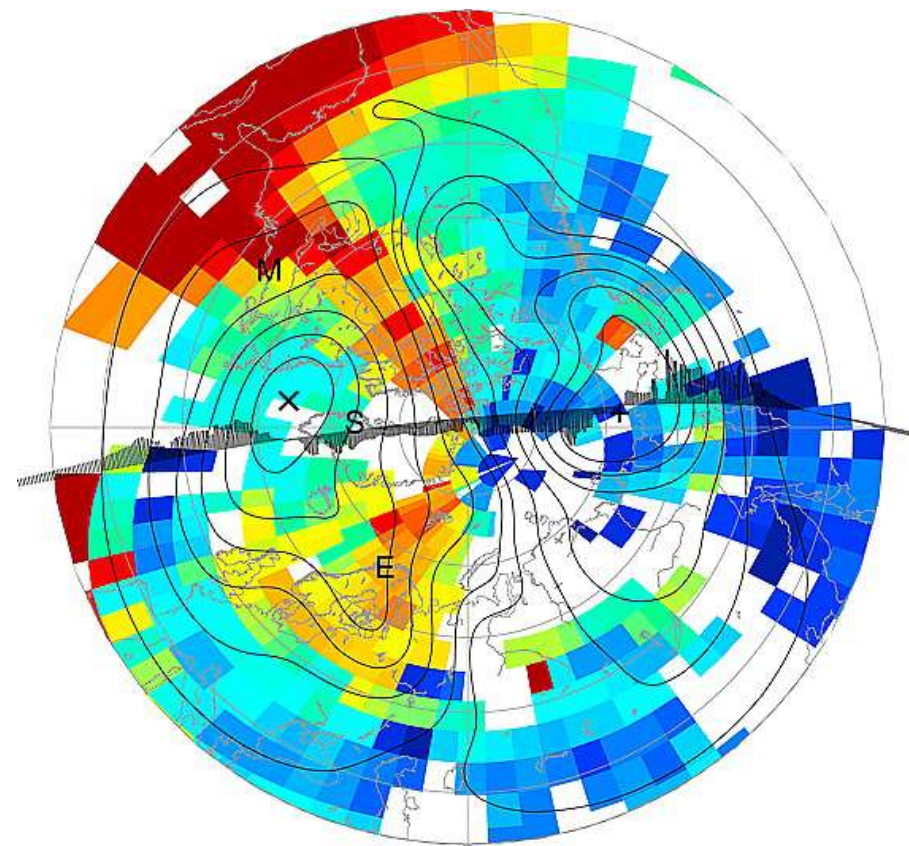
Jakowski et al., JASTP, 2008, doi:10.1016/j.jastp.2005.02.023

Formation of polar tongue of ionisation during storms

- Neutral winds can push plasma along (up) the field lines, to where the plasma recombination rate is slower.
- $E \times B$ drift may have a vertical component at mid-latitudes.
- Other mechanisms (e.g., chemical or compositional changes) could be more important during storm recovery phases.



Swisdak et al., 2006, doi:10.1029/2005GL024973

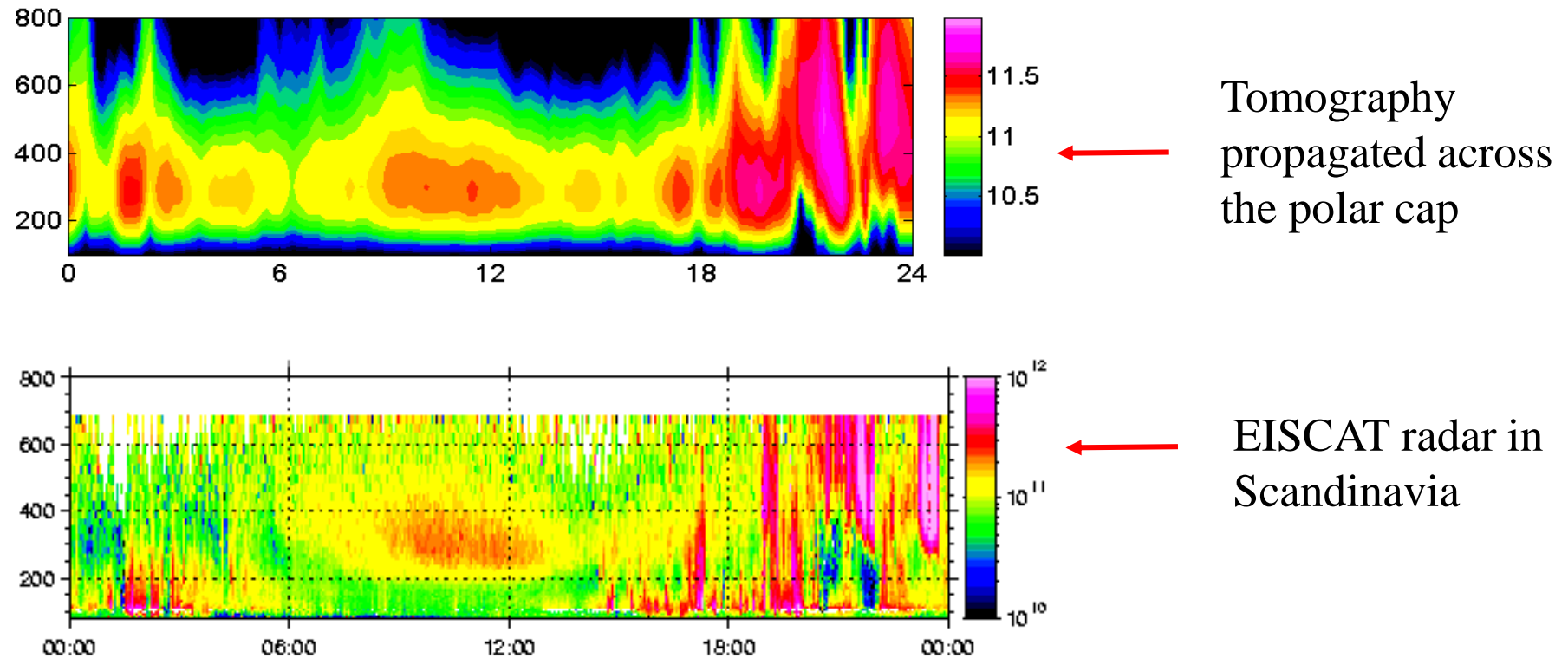


Foster et al., 2005, doi:10.1029/2004JA010928



Comparison with ground radar observations

Electron density during 30-Oct-2003 storm as a function of height and UT

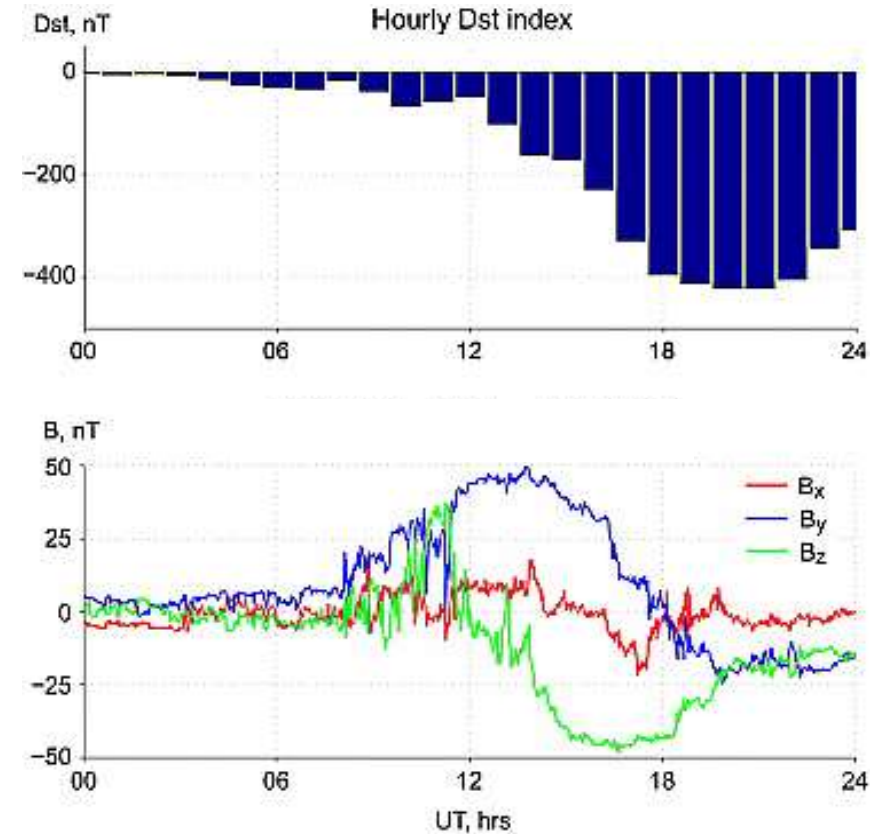


Mitchell et al., 2008, doi:10.1029/181GM09



Example: 20-Nov-2003 geomagnetic superstorm

- The storm is isolated with a quiet pre-storm day (19-Nov).
- Formation of the high-latitude anomaly is expected during the main phase (12 - 21UT).
- The tongue would be forming in North American sector (dayside during the main phase) spreading anti-sunward.
- More on the 20-Nov-2003 superstorm:
 - Foster et al., 2005, doi:10.1029/2004JA010928
 - Pokhotelov et al., 2008, doi:10.1029/2008JA013109
 - Pokhotelov et al., 2021, doi:10.5194/angeo-39-833-2021

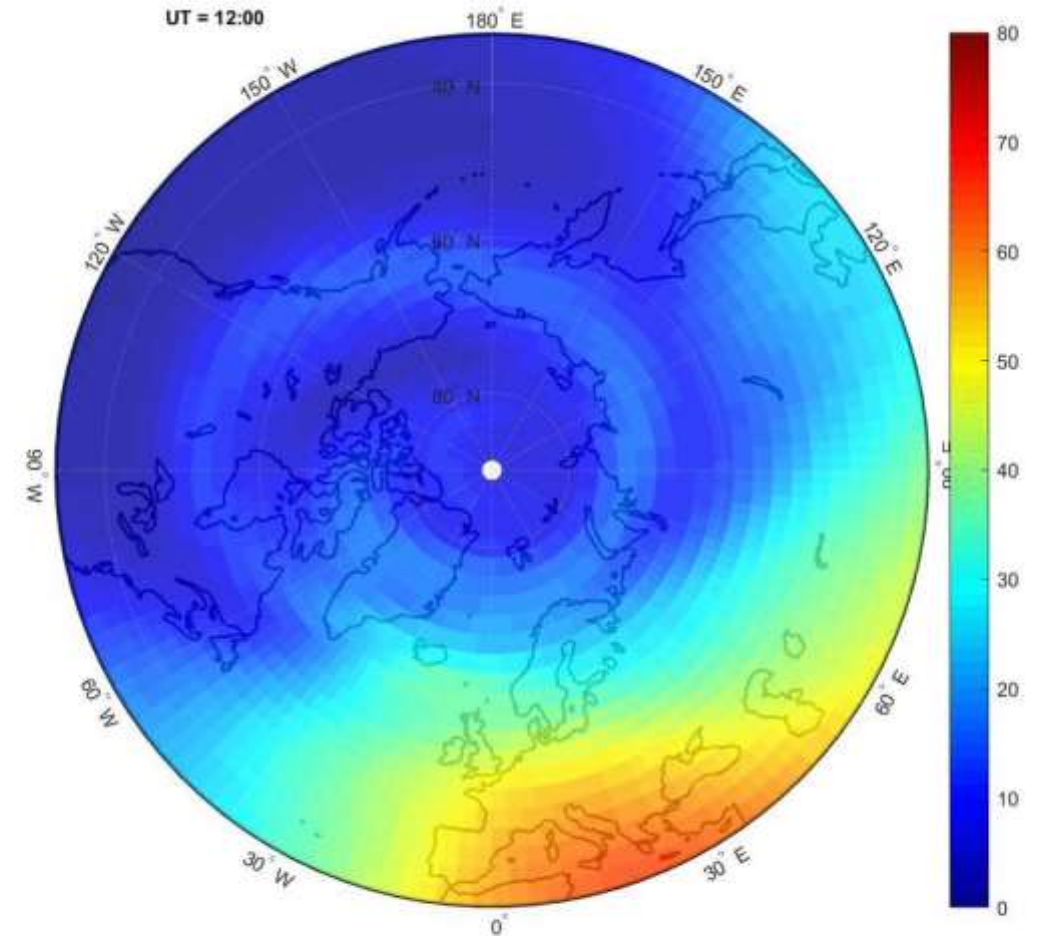


Pokhotelov et al., 2008
doi:10.1029/2008JA013109



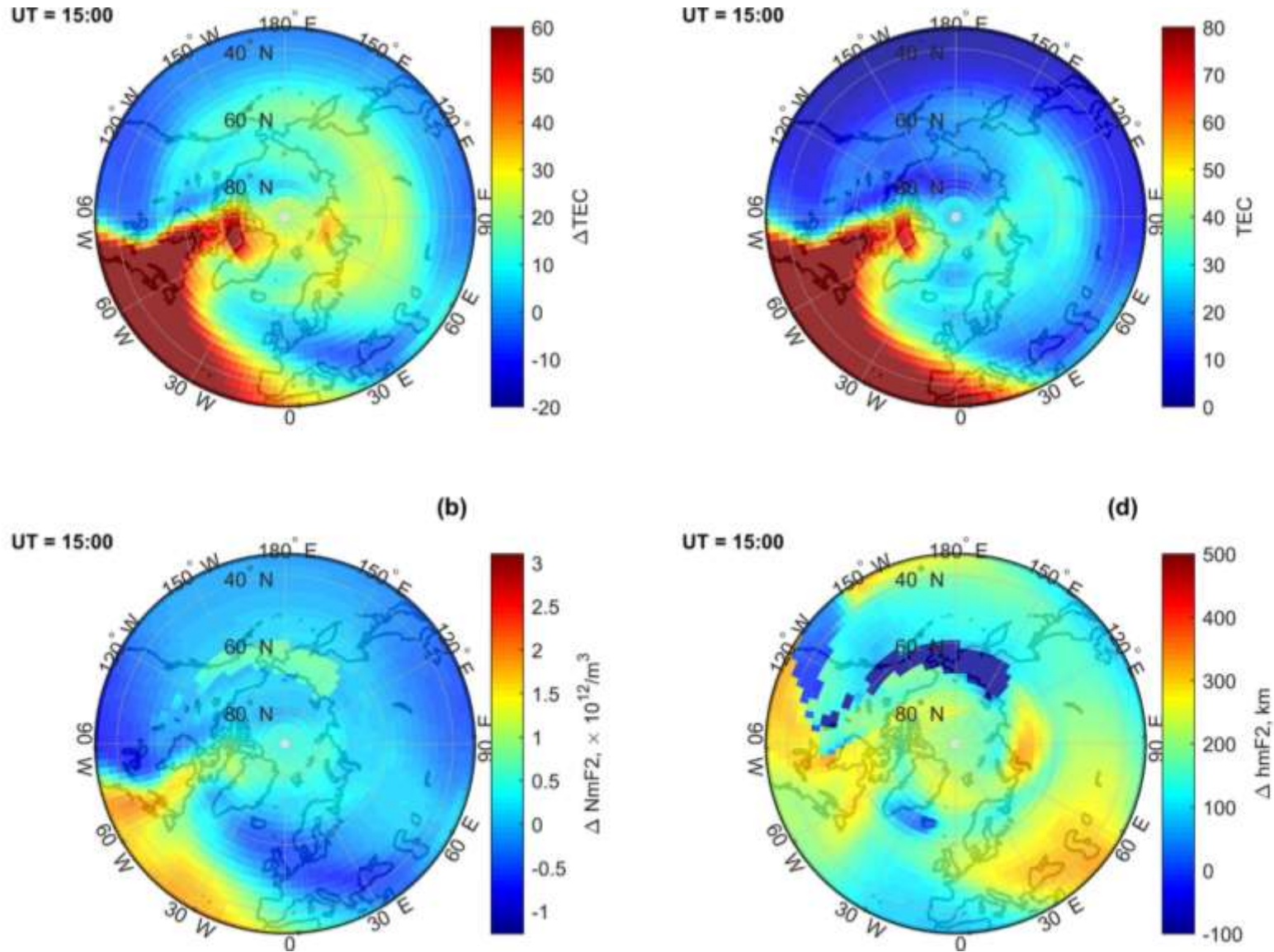
TIE-GCM Simulations: polar cap view

- TIE-GCM simulation of the TOI with Weimer model.
- Animated polar projections (above 30°N).
- TOI maximises over North American Atlantic sector during the main phase (16-18 UT).
- At 18-20 UT large amounts of transported plasma reach over the polar cap into Scandinavian sector.



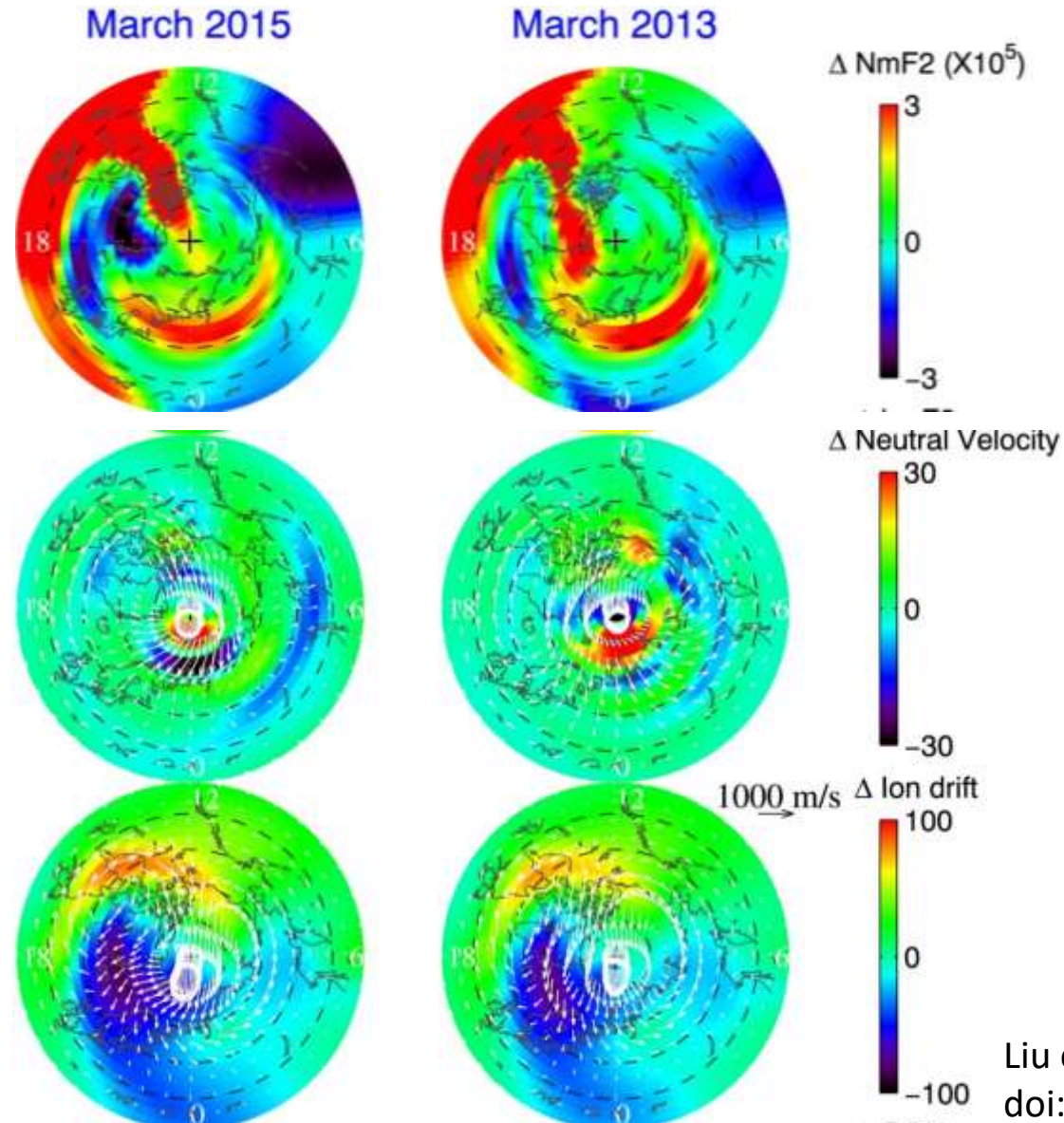
Pokhotelov et al., 2021, doi:10.5194/angeo-39-833-2021

TEC and simulated ionospheric heights from TIEGCM



Pokhotelov et al., 2021
doi:10.5194/angeo-39-833-2021

Comparison to the simulations of different geomagnetic storms



- TIE-GCM simulation of the TOI with Weimer ExB transport model for Mar 2015 (left) and Mar 2013 (right) storms.

- Geomagnetic storms of different origins:

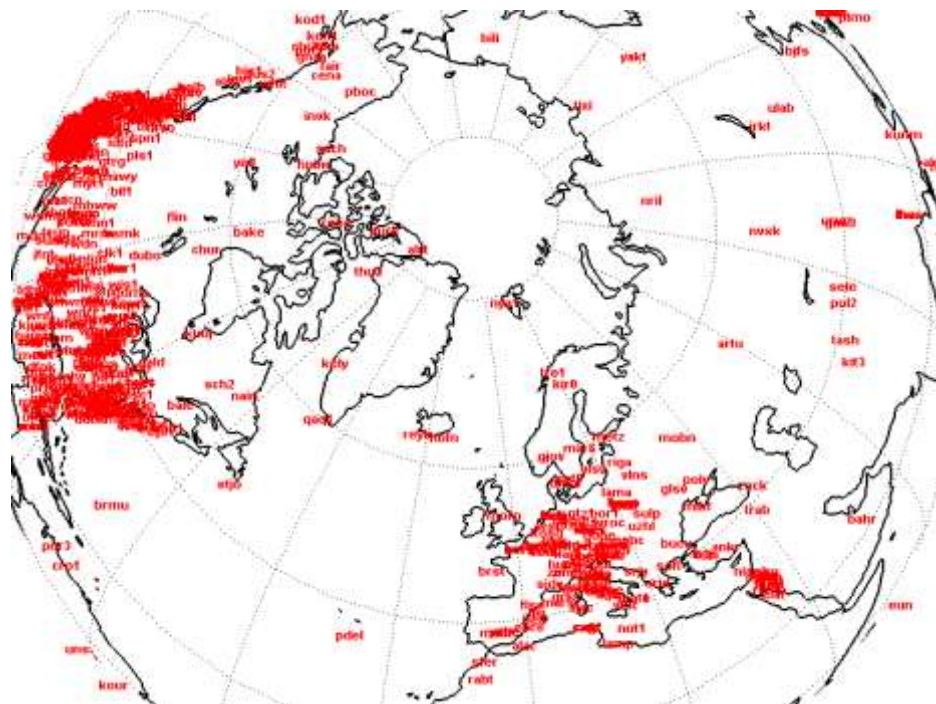
- CIR- and HSS-driven storms: Pokhotelov et al., 2009 doi:10.1029/2009JA014216; 2010 doi:10.1098/rspa.2010.0080

Liu et al., JGR, 2016
doi:10.1002/2016JA022882

Ground networks of GNSS receivers

Global scale

- IGS network distributed over the globe.
- Coverage varies greatly with location.



Mesoscale

Dense networks of GNSS receivers in Fennoscandia

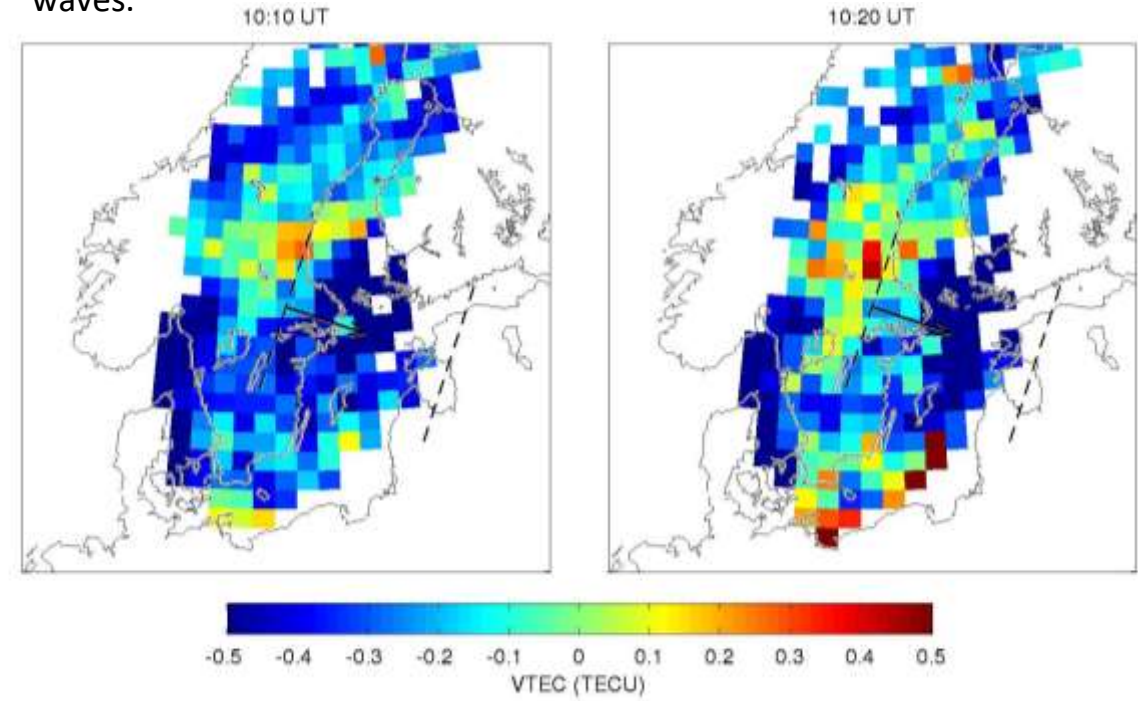
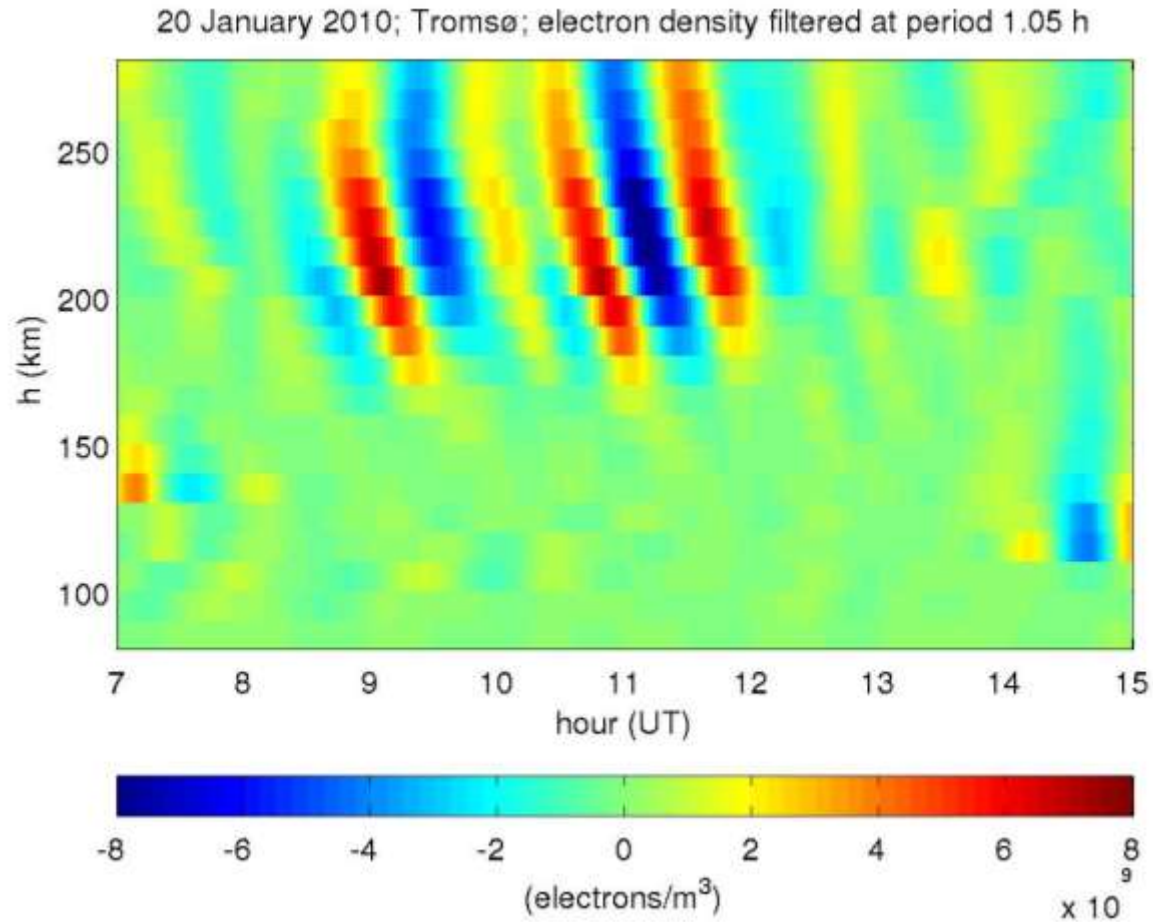
- 86 stations in Finland;
- over 200 stations in Sweden;
- more stations in Norway.



Ionospheric waves in GNSS and ISR data

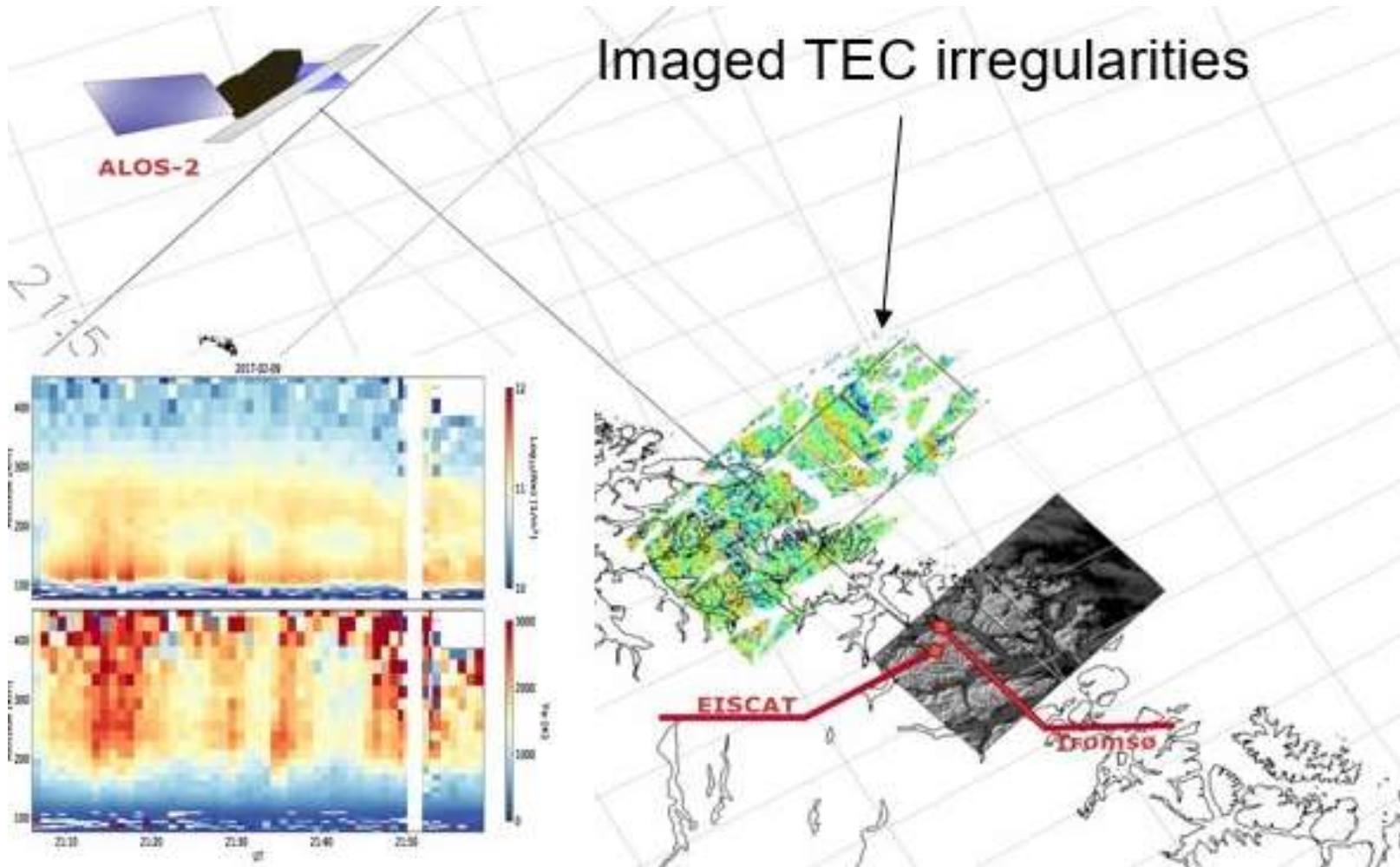
Travelling ionospheric disturbances (TIDs) can be sensed simultaneously by EISCAT radar (vertical structure) and by GNSS receiver network (horizontal structure).

The analysis of GNSS data (~200 receivers in Sweden) was inspired by Japanese research (e.g., Tsugawa et al. 2004) on detecting ionospheric effects from extra-long “pre-tsunami” waves.



van de Kamp et al, 2014 doi:10.5194/angeo-32-1511-2014

SAR imaging of auroral ionosphere



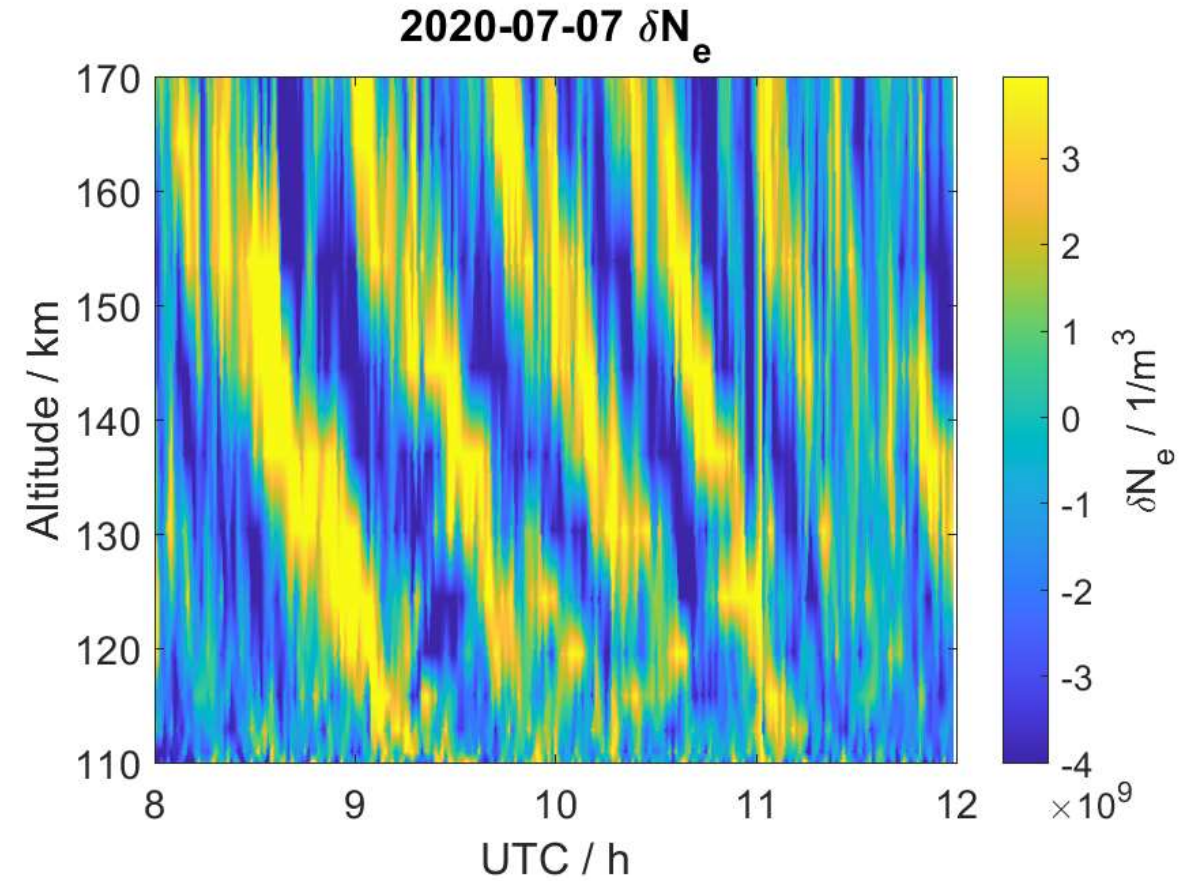
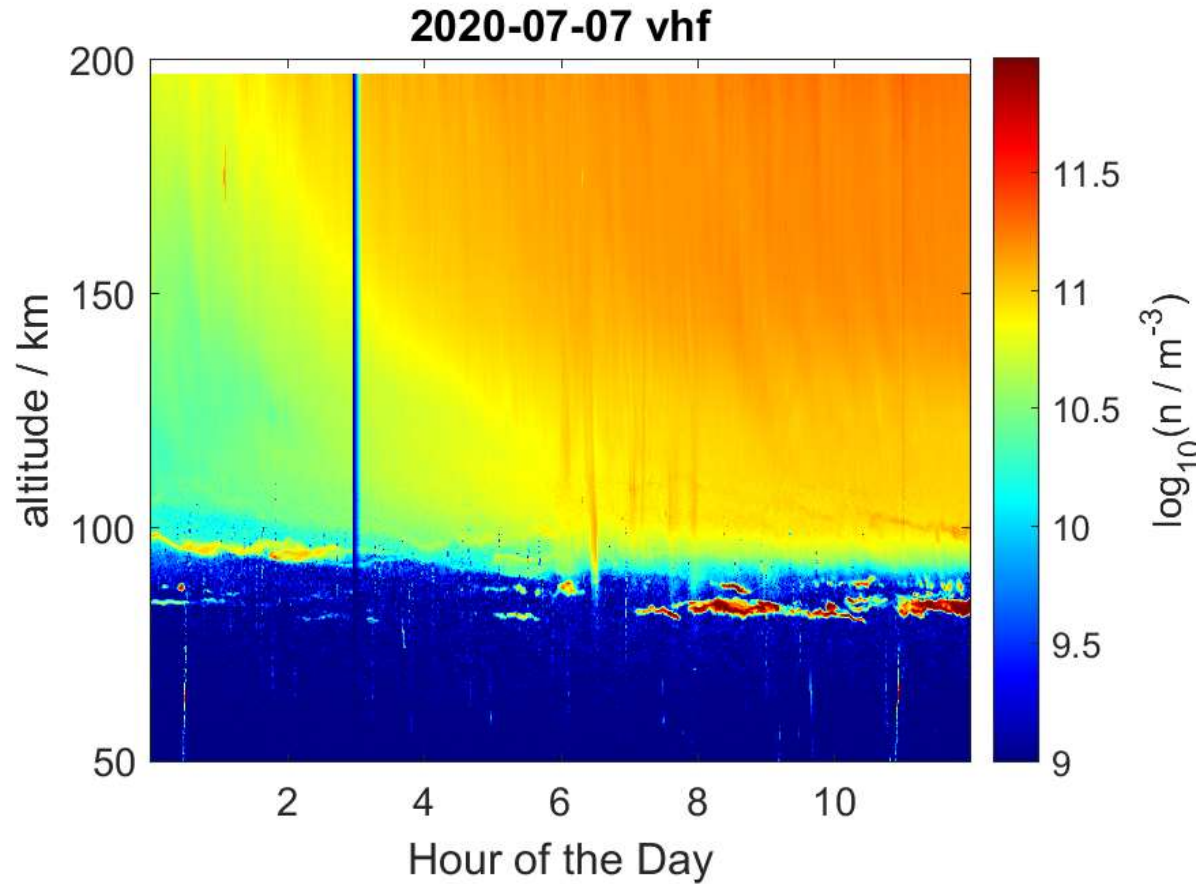
Synthetic Aperture Radar (SAR)

Imaging high-latitude plasma
Imaging high-latitude plasma density irregularities resulting from particle precipitation: spaceborne L-band SAR and EISCAT observations.

Sato et al., 2018, doi:10.1186/s40623-018-0934-1



Neutral winds and TIDs from EISCAT



**Bistatic EISCAT experiment from 2020
(jointly with Univ. of Bern and Univ of Tromsøe)**

Günzkofer et al., in preparation



Frictional heating theory, models and measurements

$$T_i = T_n + \frac{m_n}{3k} (u_i - u_n)^2 = T_n + \frac{25}{3} \cdot 10^{-4} \cdot (u_i - u_n)^2$$

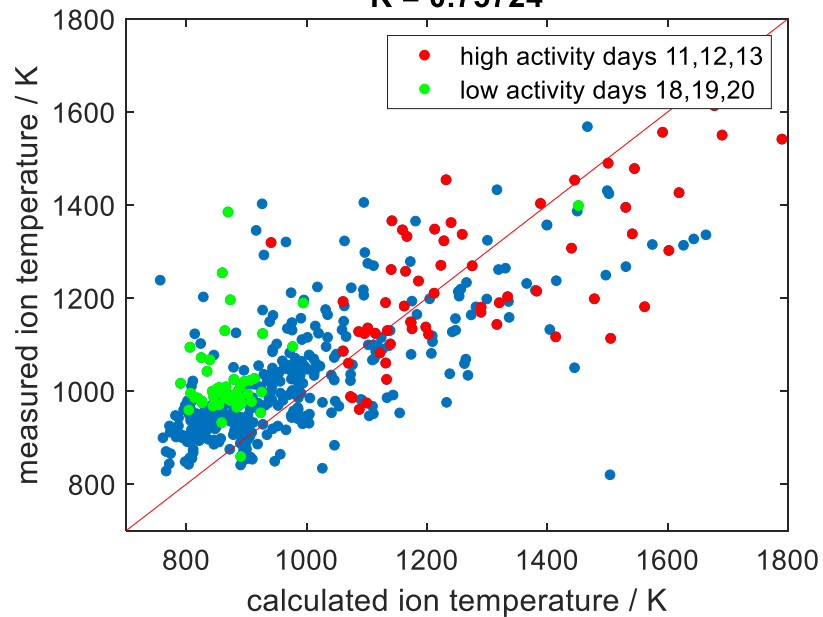
$$\frac{m_n}{3kB^2} = 0.33 \cdot 10^6$$

$T_n; u_n$: model neutral atmosphere

at 300 km

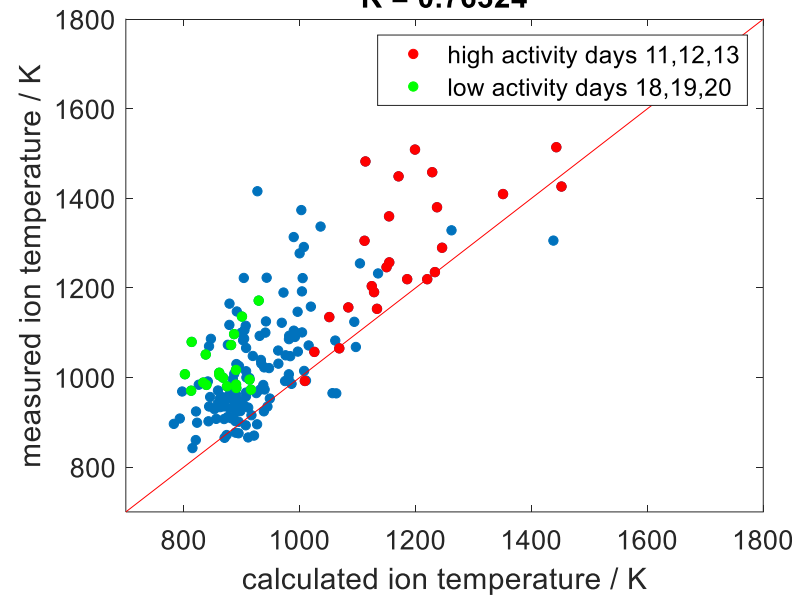
u_i : EISCAT

$R = 0.75724$



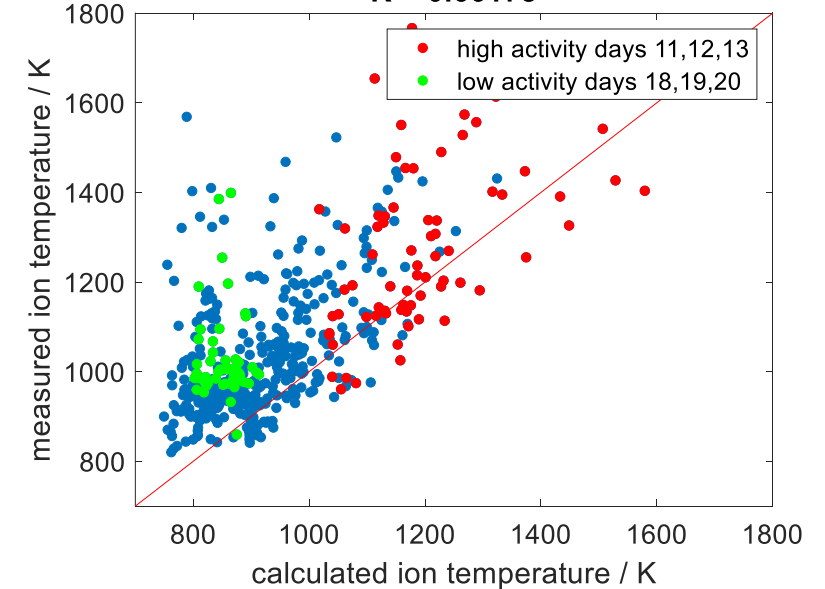
u_i : WACCM-X SD

$R = 0.76324$



u_i : TIE-GCM

$R = 0.66178$



Günzkofer et al., in preparation



Summary: neutral dynamics issues

- There seem to be a mismatch between polar cap neutral winds, produced by different circulation models (WACCM-X, GAIA, etc.).
- Number of earlier works (e.g., Burns et al., 2004; Liu et al., 2016) suggested that neutral winds during geomagnetic storms blow generally anti-sunward across the polar cap, following much faster plasma drifts.
- The anti-sunward neutral wind can be due to enhanced ion drag imposed by the ion drifts.
- An efficiency of such drag mechanism needs to be validated, especially at lower altitudes of the transition region (~120km).



Summary: plasma dynamics issues

- There is a mismatch between polar cap ExB plasma convection, produced by different convection models (Heelis, Weimer, SuperDARN, AMIE, etc.), especially in the extent of how far the ExB convection expands during geomagnetic storms.
- Electrodynamical transport plays the dominant role in the storm main phase, in the recovery phase a complex interplay between different mechanisms is involved.
- Relative roles of co-rotational and convective forces are poorly understood.
- Mechanisms of plasma uplifts at middle (sub-auroral) latitudes need to be understood, but difficult to measure. Mid-latitude SuperDARN and/or incoherent scatter radars should be used.



References

- Jakowski et al, JASTP, 2007 <https://doi.org/10.1016/j.jastp.2005.02.023>
- Mitchell, Yin, Spencer, Pokhotelov, AGU Monograph 181, 2008, <https://doi.org/10.1029/181GM09>
- Pokhotelov et al., JGR, 2008 <https://doi.org/10.1029/2008JA013109>
- van der Kamp, Pokhotelov, Kauristie, AG, 2014 <https://doi.org/10.5194/angeo-32-1511-2014>
- Sato, Kim, Jakowski, Häggström, EPS, 2018 <https://doi.org/10.1186/s40623-018-0934-1>
- Pokhotelov, Fernandez-Gomez, Borries, AG, 2021 <https://doi.org/10.5194/angeo-39-833-2021>
- Günzkofer, Pokhotelov, et al., JGR, 2022 <https://doi.org/10.1002/essoar.10511964.1> (preprint)



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<https://www.researchgate.net/profile/Dimitry-Pokhotelov>

