The Correlation of the N_A Measurements by Counting ^{28}Si Atoms

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An additional value of the Avogadro constant was obtained by counting the atoms in isotopically enriched Si spheres. With respect to the previous determination, the spheres were etched and repolished to eliminate metal contaminations and to improve the roundness. In addition, all the input quantities—molar mass, lattice parameter, mass, and volume—were remeasured aiming at a smaller uncertainty. In order to make the values given in Andreas *et al.* [Metrologia **48**, S1 (2011)] and Azuma *et al.* [Metrologia **52**, 360 (2015)] usable for a least squares adjustment, we report about the estimate of their correlation. © 2015 AIP Publishing LLC. [http://dx.doi.org/10.1063/1.4921240]

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1. Introduction

In 2011, the International Avogadro Coordination published a comprehensive survey of the measurements contributing to the determination of the Avogadro constant, N_A, by counting the atoms in two ^{28}Si-enriched single crystals—named AVO28-S5 and AVO28-S8—shaped as nearly perfect spheres.[3] This count is a candidate to the realization of the kilogram definition based on fixing the value of the Planck constant.

A limit of the 2011 determination was a metallic contamination—occurring as metal silicides—of the sphere surfaces. This contamination was selectively removed by Freckle™ etch;[4] subsequently, both spheres were repolished to improve their roundness. Eventually, measurements were repeated by using the repolished spheres—now, named AVO28-S5c and AVO28-S8c—and an additional N_A value was published in 2015.[2]

To make both the values usable for the least squares adjustment of the Planck constant value and the kilogram redefinition,[5,6] we report about the estimate of their correlation. Section 2 outlines the measurement procedure; the model used to calculate the correlation of the elemental contributions to the error budgets is given in Sec. 3. In Sec. 4, we examine in detail the contributions to the N_A uncertainty and iterate the correlation calculation up to the N_A values. The result is given in Sec. 5.

2. Measurement Equation

The value of the Avogadro constant,

$$N_A = \frac{8MV}{a^3 m},\qquad(1)$$

was obtained from measurements of the molar volume, VM/m, and lattice parameter, a, of two chemically pure ^{28}Si single-crystals shaped as nearly perfect spheres having about 93 mm diameter. In (1), m and V are the mass and volume, M is the mean molar mass, $a^3/8$ is the atom volume, and 8 is the number of atoms in the cubic unit cell.

From (1), it follows that the N_A determination requires the measurement of (i) the lattice parameter—by combined x-ray and optical interferometry,[7,8] (ii) the amount of substance fraction of the Si isotopes and, then, of the molar mass—by absolute mass-spectrometry,[9–12] and (iii) the sphere mass and volume.[13–15]

Silicon crystals contain impurities, interstitial atoms, and vacancies, which implies that the measured mass value does not correspond to that of an ideal Si crystal and that the crystal lattice may be distorted. This means that the spheres must be characterized both structurally and chemically so that the appropriate corrections are applied.[16–19] The mass, thickness, and chemical composition of the surface layers must be taken into account; they are measured by optical and x-ray spectroscopy and reflectometry.[20]

3. Correlation Model

The calculation of the correlation of the 2011 and 2015 results requires the correlation of each pair of input data—molar mass, lattice parameter, volume, and mass. In turn, it is necessary to estimate the correlations of all the elemental contributions to the uncertainties of the 2011 and 2015 input data.

To exemplify the estimation of these correlations, let us consider an elemental contribution to the total uncertainty of a given input datum. We indicate by x_{11} and x_{15} the 2011 and 2015 estimates—which might be zero—of the relevant corrections y and z. Hence,

$$x_{11} = y + u_{11},\qquad(2a)$$

$$x_{15} = z + u_{15},\qquad(2b)$$

where u_{11} and u_{15} are zero-mean errors having variances σ_{u11}^2 and σ_{u15}^2. Each error is split in two parts, that is,

$$u_{11} = \kappa_{11}\epsilon_0 + \epsilon_{11},\qquad(3a)$$

$$u_{15} = \kappa_{15}\epsilon_0 + \epsilon_{15},\qquad(3b)$$

where ϵ_0 is a zero-mean systematic error affecting identically both the estimates, κ_{11} and κ_{15} are scale factors, and ϵ_{11} and ϵ_{15} are zero-mean random errors. All these errors are assumed to be uncorrelated. In matrix notation, $\mathbf{u} = \mathbb{W}\boldsymbol{\epsilon}$, where $\mathbf{u} = [u_{11}, u_{15}]^T$, $\boldsymbol{\epsilon} = [\epsilon_0, \epsilon_{11}, \epsilon_{15}]^T$, and the design matrix,

$$\mathbb{W} = \begin{pmatrix} \kappa_{11} & 1 & 0 \\ \kappa_{15} & 0 & 1 \end{pmatrix},\qquad(4)$$

maps the $\boldsymbol{\epsilon}$ into $\mathbf{u}$. Since the $\boldsymbol{\epsilon}$ contributions to the total errors $\mathbf{u}$ are uncorrelated, their covariance matrix is

$$C_{\epsilon\epsilon} = \begin{pmatrix} \sigma_0^2 & 0 & 0 \\ 0 & \sigma_{11}^2 & 0 \\ 0 & 0 & \sigma_{15}^2 \end{pmatrix},\qquad(5)$$

where σ_0^2 is the variance of the systematic error and σ_{11}^2 and σ_{15}^2 are the variances of the random errors. Eventually, the covariance matrix of $\mathbf{u}$,

$$C_{uu} = \mathbb{W}C_{\epsilon\epsilon}\mathbb{W}^T = \begin{pmatrix} \kappa_{11}^2\sigma_0^2 + \sigma_{11}^2 & \kappa_{11}\kappa_{15}\sigma_0^2 \\ \kappa_{11}\kappa_{15}\sigma_0^2 & \kappa_{15}^2\sigma_0^2 + \sigma_{15}^2 \end{pmatrix},\qquad(6)$$

is given by the law of propagation of the uncertainty. It is worth noting that the covariance matrices of $\{u_{11}, u_{15}\}$ and $\{x_{11}, x_{15}\}$ are the same. Therefore, we can refer either to the correlation of the applied corrections or of the correction errors.

The correlation of u_{11} and u_{15} is the ratio between $\kappa_{11}\kappa_{15}\sigma_0^2$ and the geometric mean $\widetilde{\sigma}_u^2 = \sigma_{u11}\sigma_{u15}$ of the variances $\sigma_{u11}^2 = \kappa_{11}^2\sigma_0^2 + \sigma_{11}^2$ and $\sigma_{u15}^2 = \kappa_{15}^2\sigma_0^2 + \sigma_{15}^2$ of u_{11} and u_{15}, that is, $\rho = \kappa_{11}\kappa_{15}\sigma_0^2/\widetilde{\sigma}_u^2$. The $\kappa_{11}\kappa_{15}$ sign determines if the elemental uncertainty-contributions are correlated, $\kappa_{11}\kappa_{15} > 0$, or anti-correlated, $\kappa_{11}\kappa_{15} < 0$.

As a guideline, it can be noted that the correlation,

$$\rho = \frac{\kappa_{11}\sigma_0}{\sigma_{11}}\frac{\kappa_{15}\sigma_0}{\sigma_{15}}, \tag{7}$$

is the product of the systematic fractions of the total uncertainties, $\kappa_{11}\sigma_0/\sigma_{11}$ and $\kappa_{15}\sigma_0/\sigma_{15}$. Two limit cases are worth consideration. The first is when the same systematic error affects both the correction estimates, that is, when $\kappa_{11} = \kappa_{15} = 1$. In this case, $\rho = \sigma_0^2/\widetilde{\sigma}_u^2$. The second is when the scales of the systematic error are proportional to the standard deviations of the total errors, that is, when $\kappa_{11} \to \kappa_{11}\sigma_{u11}/\sigma_0$ and $\kappa_{15} \to \kappa_{15}\sigma_{u15}/\sigma_0$. In this case, $\rho = \kappa_{11}\kappa_{15}$.

4. Correlation of the Input Quantities

In the following, we examine the uncertainties and correlations of the 2011 and 2015 values of the quantities—molar mass, lattice parameter, volume, and mass—measured to determine N_A. These quantities were separately measured for each sphere and the two N_A values thus obtained were subsequently averaged. In addition, at least two laboratories measured the input quantities, but the lattice parameter, and the results were averaged. To simplify the analysis, we consider only the N_A values obtained by using AVO28-S5 and AVO28-S5c. It is worth noting that since these values do not benefit of the averaging, the associated uncertainties are a bit larger than those given in Refs. 1 and 2.

To estimate the correlation of the 2011 and 2015 values of the input quantities, we split each elemental contribution to their total uncertainties in systematic and random parts. Next, we calculate each correlation as the product of the 2011 and 2015 systematic fractions.

4.1. Molar mass

The molar mass,

$$M = M_{28} + (M_{29} - M_{28})x_{29} + (M_{30} - M_{28})x_{30}, \tag{8}$$

is given in terms of the molar masses M_i and amount-of-substance fractions x_i of the Si isotopes. For the 2015 determination, the amount-of-substance fractions were remeasured independently and *ab initio* by the PTB, NMIJ, and National Institute of Standards and Technologies, using isotope dilution[21] and multicollector inductively coupled plasma mass spectrometers,[10–12] and by the INRIM, via neutron activation.[22] Contrary to the 2011 measurement, tetramethylammonium hydroxide was used as solvent and diluent to reduce the baseline level of the ion currents in the mass spectrometry. The values of the molar mass of the Si isotopes are extremely well known and contributed negligibly to the total uncertainties. Therefore, the 2011 and 2015 values of the molar values are uncorrelated.

4.2. Lattice parameter

We determined the lattice parameter midway the spheres AVO28-S5 and AVO28-S8 by measuring—via a combined x-ray and optical interferometer—the same crystal displacement

TABLE 1. Relative uncertainties (expressed in parts per 10^9) and correlations of the elemental contributions to the uncertainties of the determinations of the AVO28-S5 and AVO28-S5c lattice parameters. The data are from tables 1 and I of Refs. 7 and 8. The extrapolation uncertainty is from Appendix C of Ref. 1

Contribution	AVO28-S5 uncertainty	AVO28-S5c uncertainty	Correlation
Data averaging	0.36	0.722	0.00
Wavelength	0.03	0.033	0.81
Laser beam diffraction	0.65	0.597	0.25
Laser beam alignment	0.77	0.480	0.00
Beam walks	1.20	0.577	0.00
Abbe's errors	1.50	0.611	0.00
Movement direction	0.65	0.214	0.22
Temperature	2.55	0.497	0.05
Thermal strain	–	0.641	–
Self-weigh	0.30	0.377	0.00
Aberrations[a]	0.64	0.642	1.00
Extrapolation	0.70	0.697	1.00
Total	3.57	1.89	0.15

[a]The 2011 uncertainty has been updated to account for the surface stress.

in terms of the periods of the traveling x-ray and optical fringes.[7,8] Since the point defects (i.e., impurity atoms and vacancies) strain the crystal lattice, the AVO28-S5 and AVO28-S5c lattice-parameters,

$$a = a_0\left[1 + \sum_i \beta_i(N_i - N_{0,i})\right], \tag{9}$$

were obtained—in 2011 and 2015, respectively—by extrapolating the value a_0 measured by x-ray/optical interferometry to the spheres' position.[17] In (9), N_i and $N_{0,i}$ are the concentrations of the i-th defect in the x-ray interferometer and spheres and β_i is the relevant sensitivity coefficient. Table 1 summarizes the uncertainties and correlations of the elemental contributions to the uncertainties of the 2011 and 2015 determinations.

The x-ray/optical interferometer used in 2011 was disassembled, upgraded, and reassembled. Therefore, the data averages are assumed uncorrelated.

In 2015, a 532 nm frequency-doubled Nd:YAG laser substituted for the 633 nm diode laser. The wavelength is affected by the pressure of the residual gas in the vacuum chamber. Since none of the chamber, pumps, and pressure gauge were changed, the systematic fractions of the wavelength errors are both estimated as 0.9, to which a 0.81 correlation will correspond.

In 2015, the optical interferometer, the chamber windows, and the delivery, collimation, modulation, pointing, and detection systems of the laser beam were rebuilt or replaced. Therefore, the errors due to beam wanderings synchronous with the displacements of the x-ray interferometer are uncorrelated.

The models and ancillary measurements required to cope with diffraction did not yet pass an experimental verification. Since we cannot exclude imperfections in the modeling and measurements, we assign to the systematic fractions of the 2011 and 2015 total uncertainties, $\kappa_{11}\sigma_0/\sigma_{11}$ and $\kappa_{15}\sigma_0/\sigma_{15}$

in (7), a uniform probability in the $[0, 1]$ interval. Therefore, $\rho \sim -\mathrm{If}(0 < z < 1)\ln(z)$ and the expected correlation is 0.25.

Since in 2015 the laser beam was aligned by using a different procedure and the estimate of the relevant correction was refined, the alignment errors are not correlated.

The offsets between the centroids of the x-ray and laser-beam spots were independently nullified. Consequently, the correlation of the Abbe errors is zero.

The error due to the projections of the crystal displacements on the normals to the front mirror and diffracting planes is the baseline misalignment $\alpha \approx 10(2)$ μrad—which is common to both the measurements—times the angular distance $\beta \approx 100(10)$ μrad (2011) or $\beta \approx 70(10)$ μrad (2015) of the movement direction from the normals' bisecting-plane. Consequently, the systematic fractions and correlation of the correction errors are equal to 1/3, 2/3, and 0.22, respectively.

As regards the temperature measurements, we used the same thermometer and fixed-point cells in both the 2011 and 2015 measurements, but the thermometer was calibrated anew and read by means of a new resistance bridge. Correlation arises because of the cell offsets and deviations of the thermometer reading from the interpolation function, which together contribute to the uncertainty by about 0.11 mK. Since the total measurement uncertainties are about 1 mK and 0.2 mK, the systematic fractions and correlation of the measurement errors are 1/10, 1/2, and 0.05.

The self-weigh strain of the x-ray interferometer depends on the distance of the support points; since we do not envisage any correlation between these distances and between the strain estimates, the strain corrections are not correlated.

Since we always used the same crystal, the interferometer aberrations affected the measurements in the same way and are fully correlated. In 2015, the associated uncertainty was better evaluated and an additional contribution due to surface stress was included;[23] therefore, the 2011 contribution in Table 1 has been updated.

The lattice parameters of AVO28-S5 and AVO28-S5c were extrapolated by taking the gradients of the point-defect concentrations between the spheres and the interferometer into account. Since, apart the tiniest effect of the newly determined gradient of the nitrogen concentration, the same extrapolations were carried out in 2011 and 2015, they are fully correlated.

By using a matrix notation, the total errors, $\mathbf{u}_a = [u_{a,11}, u_{a,15}]^\mathrm{T}$, affecting the 2011 and 2015 measurements of the lattice parameter are

$$\mathbf{u}_a = \mathbb{A}\mathbf{u}, \tag{10}$$

where $\mathbf{u} = [u_{1,11}, \ldots, u_{11,11}, u_{1,15}, \ldots, u_{12,15}]^\mathrm{T}$ are the elemental terms listed in Table 1 and the 2×23 design matrix is

$$\mathbb{A} = \begin{pmatrix} 1 & \ldots & 1 & 0 & \ldots & 0 \\ 0 & \ldots & 0 & 1 & \ldots & 1 \end{pmatrix}. \tag{11}$$

After writing the 23×23 covariance matrix C_{uu} of $\mathbf{u}$ by using the uncertainties and correlations listed in Table 1, the matrix of the relative variances and covariances of the lattice parameter values is

$$C_{aa} = \mathbb{A}C_{uu}\mathbb{A}^T = \begin{pmatrix} 12.8 & 1.03 \\ 1.03 & 3.56 \end{pmatrix} \times 10^{-18}, \tag{12}$$

yielding a correlation of 15%. Actually, the equation $\mathbb{A}C_{uu}\mathbb{A}^T$ holds only if C_{uu} contains the absolute variances and covariances. In (12), the use of relative quantities is possible because the two measured values of N_A are extremely close.

4.3. Volume

The volume of the Si sphere-core, $V = \pi D^3/6$, is given in terms of the mean diameter

$$D = D_0 + 2(\Delta R - t_{\mathrm{SL}}), \tag{13}$$

where D_0 is the mean optical diameter, the ΔR correction takes the phase shift of the laser-beam reflection into account, $D_0 + 2\Delta R$ is the mechanical diameter, and t_{SL} is the thickness of the surface layers.

Optical interferometers were used to measure the diameters. The measurement takes advantage of a differential approach: first, it measured the spacing of an optical cavity; second, the sphere is placed into the cavity and the gaps are measured. Diameters, measured in many different directions, are obtained by difference. Tables 2 and 3 give the uncertainties and correlations of the elemental contributions to the uncertainties of the 2011 and 2015 measurements.

The optical values D_0 of the sphere diameters are affected by a number of errors—e.g., in the analysis of the interference pattern, laser frequency measurement, phase-step setting, and optical alignments. Since the measurements were repeated anew, they are not correlated.

The temperature measurements were carried out by using the same thermometers and fixed-point cells, but the thermometers were calibrated anew. Correlation arises because of the cell offsets and deviations of the thermometer from the interpolation function, which together contribute to the uncertainty by about 0.28 mK. By using the values in Table 2 as an example, since total measurement uncertainties are about 0.62 mK, we estimate that the systematic fractions and correlation of the measurement errors are 0.45, 0.45, and 0.21.

TABLE 2. Relative uncertainties (expressed in parts per 10^9) and correlations of the elemental contributions to the uncertainties of the determinations of the volumes of the Si-cores of AVO28-S5 and AVO28-S5c at the NMIJ. The data are from tables 4 and 5 of Refs. 2 and 14

	AVO28-S5 uncertainty	AVO28-S5c uncertainty	Correlation
Interferogram analysis	36.3	10.3	0.00
Temperature	4.8	4.8	0.21
Diffraction	16.0	16.0	0.25
Surface[a]	2.5	0.8	0.52
Volume estimation	29.7	3.3	0.00
Total	49.9	19.9	0.07

[a]Optical effect $2(\Delta R - t_{\mathrm{SL}})$ of the surface.

TABLE 3. Relative uncertainties (expressed in parts per 10^9) and correlations of the elemental contributions to the uncertainties of the determinations of the volumes of the Si-cores of AVO28-S5 and AVO28-S5c at the PTB. The data are from tables 1 and 6 of Refs. 2 and 15

	AVO28-S5 uncertainty	AVO28-S5c uncertainty	Correlation
Interferogram analysis	4.0	4.0	0.00
Temperature	6.0	6.0	0.13
Diffraction	35.0	25.0	0.25
Surface[a]	4.8	0.5	0.52
Volume estimation	6.0	6.0	0.00
Total	36.6	26.7	0.23

[a]Optical effect $2(\Delta R - t_{SL})$ of the surface.

Wavefront distortions and parasitic interferences affect the optical diameters. The former relates to effects as retrace errors due to imperfections of both the sphere and the reference etalons, to misalignments, and to diffraction. The latter relates to interferences between the stray light reflected by the surfaces of the optical system. This problem affected the 2011 measurement of the NMIJ, but it was eliminated by redesigning the interferometer etalons. Also because of the poor understanding and modeling of the relevant phenomena, it is difficult to quantify the variance of the relevant errors. Though the different interferometers and setups used by the NMIJ and PTB and the newly shaped spheres might suggest a relatively low correlation, we cautiously assume again that the systematic fractions of the total uncertainties, $\kappa_{11}\sigma_0/\sigma_{11}$ and $\kappa_{15}\sigma_0/\sigma_{15}$ in (7), have a uniform probability in the [0, 1] interval. Therefore, as in Sec. 4.2, the expected correlation is 0.25.

The optical value of the mean diameter is corrected for the optical effect $2(\Delta R - t_{SL})$ of the surface—modeled as a layered structure of SiO$_2$, metal contaminants, chemisorbed water, and hydrocarbons. This correction combines the defective optical thickness ΔR—due to the reflection phase-shift—and the total geometrical thickness of the surface layers t_{SL}. These corrective terms are strongly correlated—because both depend on the surface model and the thicknesses of each model layer—and cancel nearly exactly, as well as their estimate uncertainties. The largest sources of uncertainty are the chemisorbed water and carbonaceous and (only in 2011) metallic contaminations. In turn, the uncertainties of these layers affect also the estimate of the SiO$_2$ thickness. The 2011 and 2015 errors of the thicknesses of the chemisorbed water and carbonaceous contaminations are strongly correlated because the former thicknesses were estimated from the same literature data and the latter were measured by using same reference. Since, before the 2015 measurements, the metallic contamination was eliminated and the sphere surface re-oxidized, we assume that the 2011 and 2015 model errors did not correlate the estimates of the SiO$_2$ thicknesses. Additionally, we assume that the systematic fractions of the $2(\Delta R - t_{SL})$ errors are the same 0.61 and 0.85 systematic-fractions of the t_{SL} errors calculated from tables 10 and 4a in Refs. 2 and 20. Accordingly, the correlation of the $2(\Delta R - t_{SL})$ errors is 0.52.

The coverage of the diameter measurements, positioning errors, and goodness of fit affect the volume (the subjects

are the coverage of the diameter measurements, positioning errors, etc.) calculation. Since, owing to the repolishing, the surface geometries differed substantially, the last entries of Tables 2 and 3 are uncorrelated.

As an example, the matrix of the relative variances and covariances of the NMIJ's volume values is

$$C_{VV} = \begin{pmatrix} 2485.1 & 69.9 \\ 69.9 & 396.7 \end{pmatrix} \times 10^{-18}, \qquad (14)$$

yielding a correlation of 7%.

The matrix of the relative variances and covariances of the 2011 and 2015 weighted averages of the NMIJ and PTB values is

$$C_{VV} = \mathbb{A}\Phi\mathbb{A}^T = \begin{pmatrix} 995..3 & 73.7 \\ 73.7 & 277.5 \end{pmatrix} \times 10^{-18}, \qquad (15)$$

where

$$\mathbb{A} = \begin{pmatrix} 1 & 1 & 0 & 0 \\ 0 & 0 & 1 & 1 \end{pmatrix} \qquad (16)$$

is the design matrix of the weighted average of the $[V_{11}^{NMIJ}, V_{11}^{PTB}, V_{15}^{NMIJ}, V_{15}^{PTB}]^T$ list of the input values and, by neglecting the correlation of the NMIJ and PTB estimates of the optical effect of the surface,

$$\Phi = \begin{pmatrix} 2485.1 & 0 & 69.9 & 0 \\ 0 & 1336.0 & 0 & 224.7 \\ 69.9 & 0 & 396.7 & 0 \\ 0 & 224.7 & 0 & 713.3 \end{pmatrix} \times 10^{-18} \qquad (17)$$

is the joint matrix of the relative variances and covariances of the 2011 and 2015 NMIJ and PTB volume-values. Eventually, the uncertainty and correlation of the 2011 and 2015 weighted averages of the NMIJ and PTB values are given in Table 4.

4.4. Mass

The sphere masses were determined by comparisons against PtIr or stainless-steel standards both in low pressure argon (by NMIJ in 2011) and *in vacuo*. Tables 5–7 give the uncertainties and correlations of the elemental contributions to the uncertainties of the 2011 and 2015 weighings at the NMIJ, PTB, and BIPM.

As an example, the model equations of the mass measurements at the NMIJ are

$$m_{Si} = m_0 + \Delta m, \qquad (18a)$$
$$m_0 = m_0^{air} + m_{H_2O}, \qquad (18b)$$
$$\Delta m = \Delta l / S + \rho(V_{Si} - V_0) + m_0 \gamma \Delta h / g, \qquad (18c)$$

TABLE 4. Uncertainty—expressed in parts per 10^9—and correlation of the weighted averages of the NMIJ and PTB determinations of the volumes of the Si-cores of AVO28-S5 and AVO28-S5c

AVO28-S5 uncertainty	AVO28-S5c uncertainty	Correlation
30.9	16.7	0.14

TABLE 5. Relative uncertainties (expressed in parts per 10^9) and correlations of the elemental contributions to the uncertainties of the determinations of the masses of the Si-cores of AVO28-S5 and AVO28-S5c at the NMIJ

	AVO28-S5 uncertainty	AVO28-S5c uncertainty	Correlation
Balance indication	3.7	5.4	0.65
Washing reproducibility	4.0	4.0	1.00
Air buoyancy	0.03	0.03	0.00
Reference-mass sorption	2.1	2.1	1.00
Center of mass	0.3	0.3	1.00
BIPM mass scale	3.0	3.0	1.00
Surface layer mass	14.5	10.0	0.27
Point defect mass	2.4	3.8	0.63
Total	16.1	13.2	0.41

TABLE 7. Relative uncertainties (expressed in parts per 10^9) and correlations of the elemental contributions to the uncertainties of the determinations of the masses of the Si-cores of AVO28-S5 and AVO28-S5c at the BIPM

	AVO28-S5 uncertainty	AVO28-S5c uncertainty	Correlation
Weighing A0[a]	1.7	1.6	0.54
Weighing[b]	0.9	2.7	0.07
Sorption of A0	0.6	0.8	0.80
BIPM mass scale	3.0	3.0	1.00
Surface layer mass	14.5	10.0	0.27
Point defect mass	2.4	3.8	0.63
Total	15.1	11.6	0.32

[a] A0 was compared against the reference mass in air.
[b] AVO28-S5 and AVO28-S5c were compared against the A0 *in vacuo*.

where m_{Si} and m_0 are the AVO28-S5c and standard masses in vacuum, Δm is the observed mass difference, m_0^{air} is the standard mass in air, m_{H_2O} is the mass of the desorbed water, Δl is the balance indication, S is the balance sensitivity, ρ is the Ar density (or zero, when weighing *in vacuo*), V_{Si} and V_0 are the AVO28-S5c and standard volumes, γ is the gravity vertical-gradient, g is the gravitational acceleration, and Δh is the height difference between the centers of mass of AVO28-S5c and the standard.

In 2014, the BIPM carried out a calibration campaign with respect to the international prototype, in anticipation of the planned redefinition of the kilogram.[24] This campaign brought to light the existence of an offset in the BIPM as-maintained mass unit, which was traceable to the international prototype in 1992. The corrections to apply depend on the date of the BIPM calibrations used for traceability. In the case of NMIJ, both the corrections were $-30.1(3.0)$ μg; in the case of PTB, the corrections were $-30.4(3.0)$ μg and $-35.4(3.0)$ μg. Therefore, the 2011 and 2015 mass values, expressed in the mass unit maintained formerly at the BIPM, were similarly reduced. For the purpose of this analysis, we assume that these corrections are fully correlated.

Such a correlation had always existed: formerly, the BIPM estimated its calibration uncertainty as 7 μg. This uncertainty was dominated by the lack of knowledge of the mass of the working standards with respect to the international prototype, which was for the last time available in 1992 and since

TABLE 6. Relative uncertainties (expressed in parts per 10^9) and correlations of the elemental contributions to the uncertainties of the determinations of the masses of the Si-cores of AVO28-S5 and AVO28-S5c at the PTB

	AVO28-S5 uncertainty	AVO28-S5c uncertainty	Correlation
Mass of standards[a]	6.7	6.6	0.23
Weighing difference	0.5	0.5	0.00
Balance linearity	0.7	0.4	0.57
Balance sensitivity	1.7	0.2	0.12
Mass position	0.5	0.5	0.50
Reference-mass sorption	9.1	1.0	0.00
Center of mass	0.3	0.3	1.00
Surface layer mass	14.5	10.0	0.27
Point defect mass	2.4	3.8	0.63
Total	18.6	12.6	0.24

[a] With auxiliary weights and correction for the BIPM mass scale.

then was stored in a safe. This error correlated all the mass calibrations made by the BIPM. Therefore, the correlation due to the traceability of the mass values to the BIPM is not new and it is not related to the upgrade of the BIPM mass scale.

To obtain the mass of the Si core, the sphere mass is corrected by subtracting the total mass of the surface layers. The SiO_2 mass is determined as density times volume; since the same 2.2(1) g/cm^3 density was used in the 2011 and 2015 determinations, the density error contributed systematically. The same 0.028(8) μg/cm^2 sorption coefficient of the chemisorbed water was used in 2011 as well in 2015;[25] therefore, the estimated errors of the chemisorbed-water masses are identical. The masses of the carbonaceous contaminations were determined by comparing the C-K to Si-L peak-ratio against the same reference carbon layer; furthermore, the same $CH_{1.5}$ stoichiometry was assumed. Consequently, the estimated masses of the carbonaceous contaminations are taken fully correlated. Since, in the 2015 measurements, the metallic contamination was removed, it does not contribute to the correlation. Eventually, we extracted the contribution of the SiO_2 density, which is systematic, from the total uncertainty of the SiO_2 mass and left out the contribution of the SiO_2 thickness, which is random. The systematic fractions and correlation of the total-mass errors—0.41, 0.66, and 0.27, respectively—have been obtained from tables 11 and 4b in Refs. [2] and [20].

The point-defect mass was taken into account by correcting the mass values to obtain the mass of a crystal having Si atoms at all lattice sites and no interstitial atoms. The mass deficits—8.1(2.4) μg and 3.8(3.8) μg for AVO28-S5[1] and AVO28-S5c,[2] respectively—were calculated on the basis of the same contaminant and vacancy concentrations, but the 2015 measurement took an additional $-0.3(2)$ μg correction because of the nitrogen contamination.[2] Furthermore, after the 2011 determination, the sphere AVO28-S5c was thermally oxidized and $-4(3)$ μg were further subtracted because of a possible metal diffusion into the crystal.[2] Consequently, the systematic fractions and correlation of the 2011 and 2015 estimate errors of the point-defect mass are 1, 0.63, and 0.63, respectively.

As an example, the matrix of the relative variances and covariances of the NMIJ's mass values is

TABLE 8. Uncertainty budgets—expressed in parts per 10^9—of the 2011 and 2015 N_A determinations using AVO28-S5 and AVO28-S5c. The uncertainties are from tables 8 and 12 of Refs. 1 and 2

	AVO28-S5 uncertainty	AVO28-S5c uncertainty	Correlation
Molar mass	7.9	5.4	0.00
Unit cell volume	10.5	5.5	0.15
Sphere volume	30.2	16.0	0.14
Sphere mass	15.0	11.0	0.32
Total	36.2	20.9	0.17

$$C_{mm} = \begin{pmatrix} 259.2 & 87.4 \\ 87.4 & 173.1 \end{pmatrix} \times 10^{-18}, \tag{19}$$

yielding a correlation of 41%.

The 2011 and 2015 mass values of the NMIJ, PTB, and BIPM were averaged by taking the correlations between the individual results into account; details are given in Refs. 2 and 13. Next, the same corrections were applied for the masses of the surface layers and point defects. For the sake of simplicity, in Table 8, we use the 32% correlation of the BIPM values as that of the averaged Si-core masses of AVO28-S5 and AVO28-S5c.

5. Correlation of the N_A Measurements

Table 8 summarizes the uncertainty and correlation of the contributions to the uncertainty of the 2011 and 2015 determinations of N_A using AVO28-S5 and AVO28-S5c. The correlations are the exemplar values given in Tables 1, 4, and 7, but the uncertainty contributions are taken from tables 8 and 12 of Refs. 1 and 2. The small differences between the uncertainties given in Table 8 and in Secs. 4.2 and 4.3 are due to the difficulty to reconstruct the rounding in the calculations. The joint matrix of the relative variances and covariances is

$$C_{uu} = \begin{pmatrix} 62 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 110 & 0 & 0 & 0 & 9 & 0 & 0 \\ 0 & 0 & 912 & 0 & 0 & 0 & 68 & 0 \\ 0 & 0 & 0 & 225 & 0 & 0 & 0 & 53 \\ 0 & 0 & 0 & 0 & 29 & 0 & 0 & 0 \\ 0 & 9 & 0 & 0 & 0 & 30 & 0 & 0 \\ 0 & 0 & 68 & 0 & 0 & 0 & 256 & 0 \\ 0 & 0 & 0 & 53 & 0 & 0 & 0 & 121 \end{pmatrix} \times 10^{-18}. \tag{20}$$

The matrix mapping the contributions in Table 8 into the 2011 and 2015 total uncertainties is

$$\mathbb{A} = \begin{pmatrix} 1 & 1 & 1 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 1 & 1 & 1 \end{pmatrix}. \tag{21}$$

Eventually, the matrix of the relative variance and covariances of the two N_A values is

$$C_{N_A N_A} = \mathbb{A} C_{uu} \mathbb{A}^T = \begin{pmatrix} 1310 & 129 \\ 129 & 436 \end{pmatrix} \times 10^{-18}, \tag{22}$$

to which a 17% correlation will correspond.

6. Conclusions

The forthcoming definition of the kilogram based on a fixed value of the Planck constant is prompting the development of extremely refined technologies and experiments to measure both the Planck and Avogadro constants with the utmost accuracy. In order to gain the maximum benefit from the measurement results, to extract from them all the available information, and to evidence discrepancies, a similarly refined analysis of the measurement uncertainty is desirable. This paper investigated the correlation of the N_A values reported by the International Avogadro Coordination. Despite these values being obtained by counting twice the atoms in the same ^{28}Si-enriched spheres, owing to the sphere-surface repolishing and the upgrades of many of the measurement technologies, their correlation is only 17%.

In Ref. 2, the value $N_A = 6.022\,140\,82(18) \times 10^{23}$ mol^{-1} given in Ref. 1 has been updated to $N_A = 6.022\,140\,99(18) \times 10^{23}$ mol^{-1}, by taking the recalibration of the mass standards into account. The value derived from the new measurements repeated by using the repolished spheres is $N_A = 6.022\,140\,76(12) \times 10^{23}$ mol^{-1}.[2] Taking the correlation into account, the average is

$$N_A = 6.022\,140\,82(11) \times 10^{23} \text{ mol}^{-1}, \tag{23}$$

with a relative uncertainty of 1.8×10^{-8}.

To assess the consistency of the N_A values given in Refs. 1 and 2, we observe that their difference is $38(33) \times 10^{-9} N_A$, where we took the 17% correlation into account in the uncertainty calculation.

Acknowledgments

7. References

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