

# Comparative study of DR Nanoantenna Array For Optical Frequencies

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**Abstract**—The frequency dependent complex permittivity of plasmonic materials is one of the most decisive parameters in their resonant characteristics. This work shows how the optical responses of gold (Au), and silver (Ag) optical nanoantennas are affected by their size, shape and their frequency dependent optical functions. The optical functions of these metals are described by the Drude-Lorentz model in which both the free electrons contributions and harmonic oscillator (SPRs) activities are considered. The change of material in the nanostrip and the ground plane results in the change in different parameters like return loss, VSWR, etc.

**Keyword**- optical antennas, nanoplasmonics, plasmonic metals, optical energy.

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## I. INTRODUCTION

Optical antennas stand for a class of novel optical detectors that have the potential to revolutionize optical interconnections, image sensing and related fields by adapting radio wave techniques to the optical regime. Optical antennas couple up with electromagnetic radiation in the visible and infrared wavelengths in the same way that radio electric antennas do at the corresponding wavelengths [1]. The optical antenna size is in the range of the detected wavelength from a few hundred nanometers to a few microns. Optical antennas are advantageous in the detection of light showing polarization dependence, tunability, and a potential rapid time response. They also can be assumed as point detectors and directionally sensitive.

Nanoplasmonics is a branch of optical material science devoted on the phenomena of nanoscale nanostructured metal systems. A unique property of such systems is their ability to keep the optical energy concentrated on the nanoscale because of modes called surface plasmons (SPs). The existence of SPs depends entirely on the fact that dielectric function  $\epsilon_m$  has a negative real part,  $\text{Re } \epsilon_m < 0$  [2]. The SPs are well pronounced as resonances when the losses are quite small, i.e.,  $\text{Im } \epsilon_m \ll -\text{Re } \epsilon_m$ . This is a known property of a good plasmonic metal, valid, e.g., for silver in the most of the visible region. A substance is a good plasmonic metal if these two properties are satisfied simultaneously.

$$\text{Re } \epsilon_m < 0, \text{Im } \epsilon_m \ll -\text{Re } \epsilon_m \dots \dots (1)$$

There is an extent to which an electromagnetic wave can be concentrated. The nanoplasmonics is about concentration of electromechanical energy at optical frequencies on the nanoscale. The scale of the concentration of electromagnetic energy can be determined by the wavelength that can be understood from Fig. 1 (a).

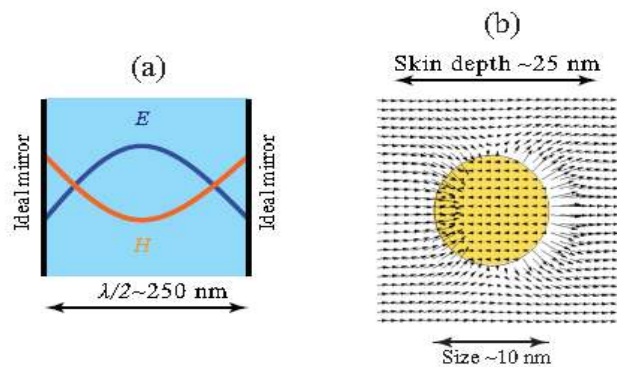


Fig. 1(a) localisation of optical fields by ideal mirror (b) by a gold nano particle.

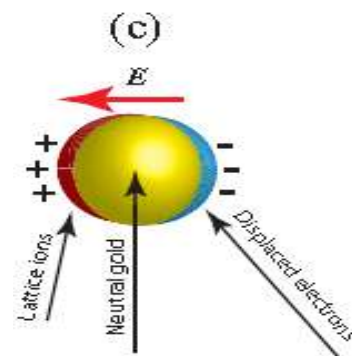


Fig.1(c) schematic of charge separation.

It is difficult to achieve a localization of the optical energy to smaller volume than  $\lambda^3/8$  apart from the ideal mirrors gives the best confinement of electromagnetic waves. There are two implied assumptions: (i) The optical energy is electromagnetic energy, (ii) The best confinement is brought through ideal mirrors. Both these assumptions must be removed to get nanolocalization of optical energy. Consider a nanoplasmonic system whose size is less than or comparable to the skin depth.

$$l_s = \lambda \left[ \operatorname{Re} \left( \frac{-\epsilon_m^2}{\epsilon_m + \epsilon_d} \right)^{\frac{1}{2}} \right]^{-1} \dots\dots\dots (2)$$

Where  $\lambda = \omega / c$  is the reduced vacuum wavelength. For single-valence plasmonic metals (silver, gold, copper, alkaline metals)  $l_s \approx 25$  nm covering the entire optical region. In this paper we are analysing the effect of the two metals i.e. gold and silver respectively on the optical DR nanoantenna array.[3]

For such a plasmonic nanosystem with  $R \leq l_s$ , the optical electric field penetrates the entire system and drives oscillations of the metal electrons. The total energy of the system in this case is a sum of the potential energy of the electrons in the electric field and their mechanical kinetic energy. While the magnetic field is present, non-relativistic electrons' interaction with it is weak proportional to a small parameter  $\frac{v_F}{c} \sim \alpha \sim 10^{-2}$ , where  $v_F$  is the electron speed at the Fermi surface,  $c$  is speed of light, and  $\alpha = e^2/hc$  is the fine structure constant. Thus in this limit, which is conventionally called quasistatic, the effects of the magnetic component of the total energy is relatively small. Hence, this total energy is mostly *electromechanical* (and not *electromagnetic*) energy.

At optical frequencies, electrons in metals have considerable inertia and cannot respond instantaneously to the driving fields. Typically, the skin depth is on the order of tens of nanometers, comparable to the dimensions of the antenna. Traditional design rules that prescribe antenna parameters only in terms of an external wavelength are thus no longer valid. The metal must be rigorously treated as a strongly coupled plasma, which leads to the antenna "seeing" a reduced effective wavelength. This effective wavelength,  $\lambda_{eff}$ , is related to the external (incident) wavelength,  $\lambda$ , by a surprisingly simple relation

$$\lambda_{eff} = n_1 + n_2 \left[ \frac{\lambda}{\lambda_p} \right] \dots\dots\dots (3)$$

where  $\lambda_p$  is the plasma wavelength of the metal, and  $n_1$  and  $n_2$  are constants that depend on the geometry and dielectric parameters of the antenna.  $\lambda_{eff}$  is shorter, by a factor of 2 to 6, than the free space,  $\lambda$ , for typical metals (e.g., gold, silver) and realistic antenna thicknesses. The shortening of wavelength from  $\lambda$  to  $\lambda_{eff}$  has interesting implications. For example, it implies that the radiation resistance of an optical half-wave antenna is on the order of just a few Ohms.

## II. DESIGN

The proposed antenna design is classified with the help of two materials of the nanostrip as well as the ground plane. Initially the material assigned for the nanostrip and the ground plane was silver. The optical dielectric nanoantenna array has four radiating elements of cylindrical shape and is composed of silicon with the height ( $h$ ) = 325 nm and diameter ( $d$ ) = 510 nm.

A DRA can be analysed in terms of its shape and modes of excitation. The electromagnetic near fields inside the resonating body are the prime sources of interpreting the far field radiation patterns and thus help to understand the overall characteristics of the antenna. Though many design shapes are available, however the best one is that which is either readily available or that is simple to cut, grind and mould. It is for this reason that a circular cylindrical shape has been chosen for the proposed designs. The geometry shape of the DRA is analysed in terms of its resonant modes, near-field distribution inside the resonator, far-field radiation into the space, resonant frequency and impedance bandwidth. The overall performance of an array can be evaluated in terms of the following parameters;

- i. Geometry and dimensions of DRA elements.
- ii. The spacing among elements.
- iii. The number of elements used.
- iv. The feed mechanism.[9]

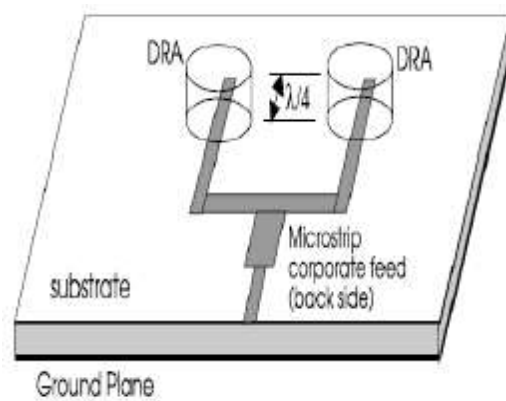


Fig. 2: feed structure to construct arrays

The resonators are above the nanostrip and positioned to obtain a maximum wideband. The antenna substrate is composed by  $\text{SiO}_2$  and the ground plane is composed by silver. The thickness of the  $\text{SiO}_2$  layer between the nanostrip and the ground plane is  $h_1 = 145$  nm. There is a layer of thickness  $h_3 = 10$  nm made of  $\text{SiO}_2$  between the nanostrip and the DR. The nanostrip has  $w = 340$  nm of width and  $h_2 = 20$  nm of height. The metallic regions of the nanostrip are composed by silver. Whose dispersive properties were described by the Drude's model assuming  $\epsilon_{inf} = 5$ ,  $f_p = 2175$  THz, and  $\gamma = 4.35$  THz [4][5]

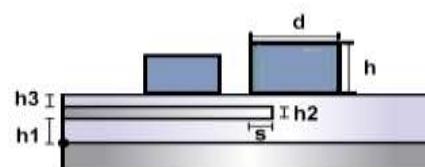


Fig 3: side view of the DR antenna array.

The DRA was investigated in nanoscale and operating at optical frequencies. The resonators are composed of silicon

( $\epsilon_r = 11.56$ ). The resonance frequency of DR's  $HE_{11\delta}$  mode in free space, can be estimated from [6]

$$f = \frac{6.324c}{(2\pi a \sqrt{\epsilon_r + 2})} \left[ 0.27 + 0.36 \left( \frac{d}{2h} \right) + 0.002 \left( \frac{d}{2h} \right)^2 \right] \dots\dots\dots(4)$$

According to this equation the mode can resonate at 193.5 THz when  $d$  and  $h$  are assumed proportionally. The length of nanostrip which is assumed to be under the resonator is taken as "s", which plays a very crucial role in the variation of the results. [7]

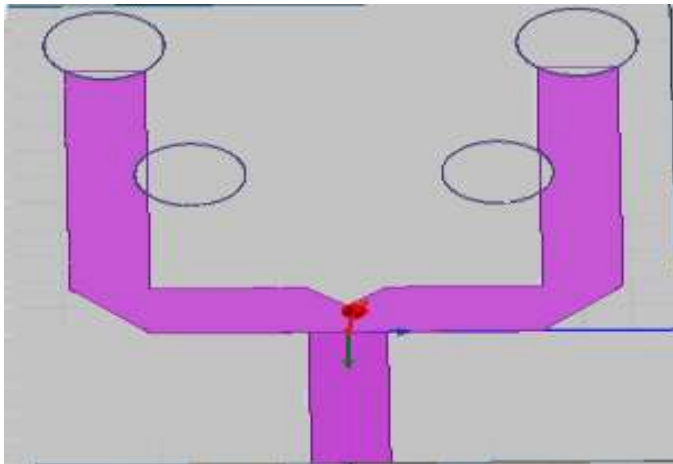


Fig 4: Top view of the nanostrip and ground plane is composed of silver

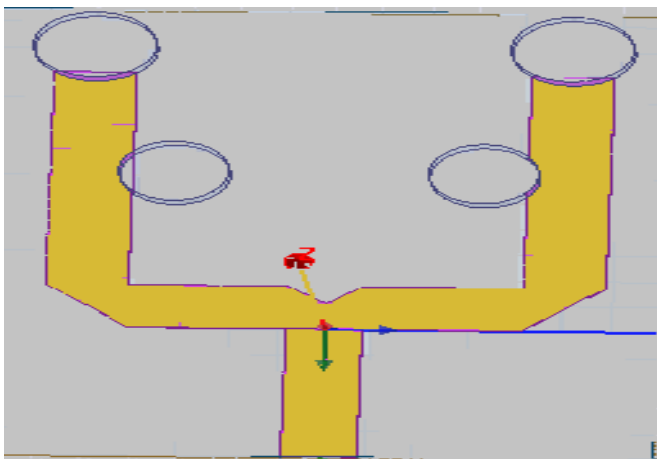


Fig 5: Top view of the nanostrip and the ground plane is composed of gold

The antenna design considered, has four cylindrical resonator elements, each of equal height and diameter, one pair is kept near the ends of nanostrip and another pair is kept at the midpoint of the length of nanostrip.

### III. RESULTS

The results are simulated on HFSS14.0 software licensed version which is a FEM (Finite Element Method) based simulator. It is used to calculate various Electromagnetic behavior of the structure such as basic EM field quantities, S-parameters, resonance of the structure, etc. The return loss

of the optical DR nanoantenna of material silver and gold at the optical frequency 193.5THz are -18.26db and -23.5db.

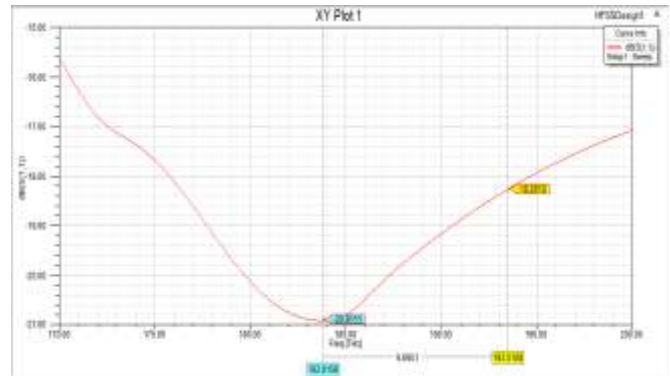


Fig 6(a): return loss for silver nanostrip and ground plane

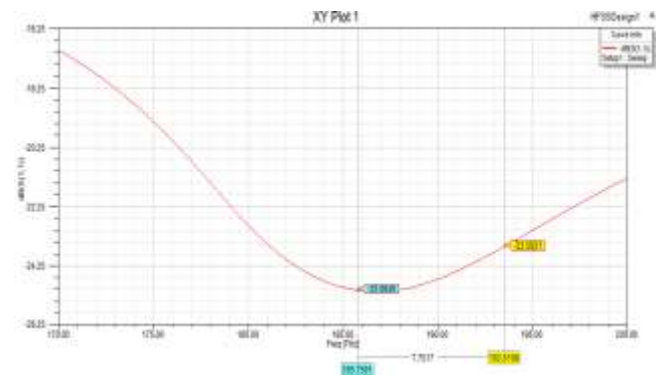


Fig. 6(b):return loss for gold nanostrip and ground plane

The voltage standing wave ratio (VSWR) is a measurement of the impedance mismatch between the transmitter and the antenna. Large value of VSWR corresponds to the high mismatch. Minimum value of VSWR corresponds to a perfect match that is unity. The VSWR v/s frequency plot for the proposed antenna is shown in Fig.7. The VSWR for silver nanostructure and ground plane is 1.19 and for gold nanostructure and ground plane is 1.11. It may be observed that the values of the VSWR are less than 2 for the whole bandwidth which is within the required limits .

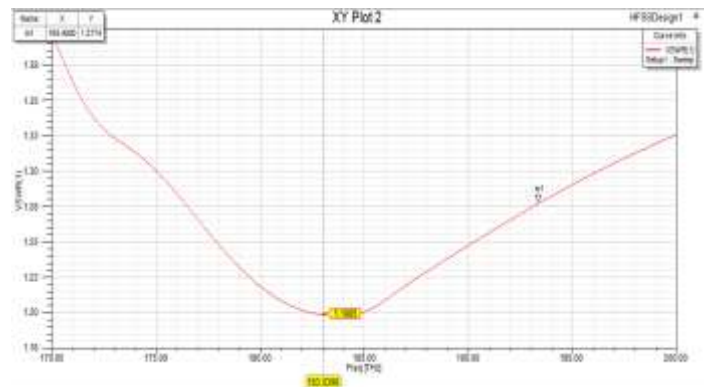


Fig 7(a) : VSWR for silver nanostrip and ground plane

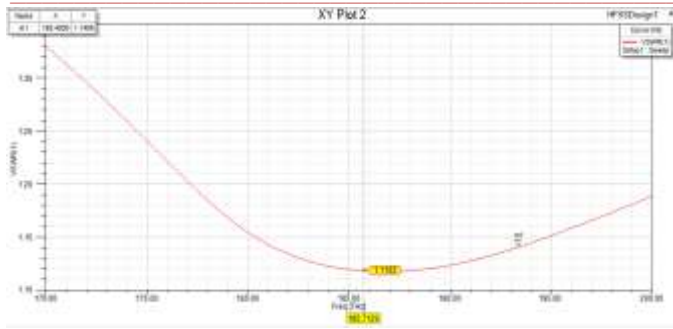


Fig. 7(b): VSWR for gold nanostrip and ground plane

#### IV. CONCLUSION

The optical DR nanoantenna array having four dielectric resonator elements has been compared separately by assigning materials like silver and gold on the nanostrip and the ground plane. The design having gold as nanostrip and the ground plane have better results than the silver one. This design can be used for many applications in the field of medicine, photovoltaic, spectroscopy and near field microscopy. [8] New ideas and developments are emerging at a rapid pace, and it is now clear that the optical antenna concept will provide new opportunities for optoelectronic architectures and devices. Today, the building blocks for optical antennas are plasmonic nanostructures that can be fabricated either from the bottom up by colloidal chemistry or from the top down with established nanofabrication techniques, such as electron-beam lithography and focused ion-beam milling. It is also conceivable that future optical antenna designs will draw inspiration from biological systems, such as light harvesting proteins in photosynthesis.

#### REFERENCES

- [1] Boreman G 2002 Divide and conquer *OE Mag.* **2** 47–8
- [2] D. J. Bergman and D. Stroud, “Properties of macroscopically inhomogeneous media,” in “Solid State Physics,” vol. 46, H. Ehrenreich and D. Turnbull, eds. (Academic Press, 1992), pp. 148–270.
- [3] L. D. Landau and E. M. Lifshitz, *Electrodynamics of Continuous Media* (Pergamon, 1984).
- [4] P. B. Johnson and R. W. Christy, “Optical Constants of the Noble Metals,” *Phys. Rev. B*, vol. 6, no 12, p. 4370–4379, dez. 1972
- [5] H. Iizuka, N. Engheta, H. Fujikawa, and K. Sato, “Arm-edge conditions in plasmonic folded dipole nanoantennas,” *Opt. Express* **19**(13), 12325–12335 (2011).
- [6] R. K. Mongia and P. Bhartia, “Dielectric resonator antennas—a review and general design relation frequency and bandwidth,” *Int. J. Microwave Millimeter-Wave Computer-Aided Engineering* **4**(3), 230–247(1994)
- [7] Gilliard N. Malheiros-Silveira, IEEE, Student Member and Hugo E. Hernández-Figueroa, IEEE, Senior Member, “Dielectric Resonator Nanoantenna Array for Optical Frequencies” ©2013 IEEE

- [8] A E Krasnok, I S Maksymov, A I Denisyuk, P A Belov, A E Miroshnichenko, C R Simovski, Yu S Kivshar, “Optical nanoantennas” #2013 *Uspekhi Fizicheskikh Nauk*, Russian Academy of Sciences
- [9] A. Petosa, R.K. Mongia, A. Ittipiboon and J.S. Wight, “Experimental investigation on feed structures for linear arrays of dielectric resonator antennas”, *Proceedings of IEEE AP-S Conference*, California, USA, pp. 1982–1985, 1995