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Xianfan Wang

Elfed Lewis

Pengfei Wang

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Investigation of the self-imaging position of a singlemode-multimode-singlemode optical fiber structure

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TABLE 2 Performance comparison of the proposed balun with other works

References	Design method	Output isolation	Size (mm ²)	Frequency ratio	Bandwidth ($ S_{11} > 10$ dB, GHz)	Flexible implementation
5	TDC ^a	No	16.8 × 40.5	N/A	0.08	×
6	CRLHTL ^b	No	29 × 30.5	N/A	1.25	×
8	BLS	No	40 × 55	N/A	N/A	×
9	DSPSL	Exist	51.9 × 105	1.62–3	>0.10 ^c	×
This work	TSPSCL	Exist	55 × 70	2.38–4.12	1.30/0.62	✓

^aTrans-directional coupler.

^bComposite right/left handed transmission line.

^c $|S_{11}| > 15$ dB.

design. Therefore, this novel planar balun will be extensively utilized in dual-band radio frequency circuits and systems.

ACKNOWLEDGMENTS

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Investigation of the self-imaging position of a singlemode-multimode-singlemode optical fiber structure

Xianfan Wang¹ | Elfed Lewis² | Pengfei Wang^{1,3}

¹Key Lab of In-fiber Integrated Optics, Ministry Education of China, Harbin Engineering University, Harbin 150000, China

²Optical Fibre Sensors Research Centre, Department of Electronic and Computer Engineering, University of Limerick, Limerick, Ireland

³Photonics Research Centre, Dublin Institute of Technology, Kevin Street, Dublin 8, Ireland

Correspondence

Pengfei Wang, Photonics Research Centre, Dublin Institute of Technology, Kevin Street, Dublin 8, Ireland.

Email: pengfei.wang@dit.ie

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Abstract

In this paper, a theoretical and experimental investigation of the multimode interference induced self-imaging phenomenon in a singlemode-multimode-singlemode (SMS) fiber structure is described. The simulated transmission spectra based on guided modal propagation analysis and beam propagation method are presented and compared with experimental obtained results. The temperature dependence of the self-imaging position of the SMS fiber structure is also investigated. These investigations provide excellent opportunities for multiple future applications including telecommunications and sensing.

KEYWORDS

multimode interference, self-imaging position, temperature dependence

1 | INTRODUCTION

The SMS fiber structure has been widely investigated and is a relatively simple structure being a multimode fiber spliced between two single-mode fibers. It can be used as the basis for fabricating fiber optic sensors with a range of advantages such as ease of fabrication, high sensitivity, compact size, and potential low cost,^{1–4} which compares favorably with other technologies in fiber optic sensors including the fiber Bragg Grating (FBG) and long period grating (LPG).^{5,6}

The SMS fiber structure is an excellent candidate for fiber optic sensing applications due to the existence of inherent multimode interference (MMI) in the multimode fiber (MMF), and this hybrid fiber structure can be developed as an edge filter,⁷ bandpass filter,⁸ as well as other optical fiber based photonic components by controlling the length and selecting the type of MMF sandwiched between two SMFs. Self-imaging, as a basic principle of the multimode interferometer (MMI) has been utilized to design a number of MMI based couplers,⁹ splitters,¹⁰ combiners,¹¹ optical switches,¹² and multi/demultiplexers.¹³ In addition to significant experimental progress, investigations of the theoretical aspects of MMI have recently been published.

In 1995, Soldano and Pennings demonstrated a theoretical analysis of the self-imaging properties in a planar multimode waveguide utilizing a guided modal propagation analysis (MPA), in which the formation of multiple re-imaging distance at certain lengths of MMI was successfully derived.¹⁴ Since then another analysis developed by Wang et al.¹⁵ using a numerical model based on the wide angle-beam propagation method (WA-BPM) through scanning the length of the MMF section has been demonstrated for theoretical prediction of the characteristics of MMI devices.

Over the past decade, the temperature dependence of the SMS fiber structure has also been widely investigated based on either wavelength encoding or intensity variation.^{2,16–20} The results of these investigations have shown that the thermo-optic coefficient (TOC) of MMF makes a more significant contribution to the temperature dependence of the SMS than that of the thermal expansion coefficient (TEC). To date no study of the temperature dependence of self-imaging properties within the MMF has been reported in the literature.

In this paper, a comparison between the theoretical models on the modeling of light propagation within the SMS fiber hybrid structure is presented, which has not been addressed in the literature. All analyses and predictions are verified numerically and experimentally. Additionally, a comprehensive investigation of the temperature dependence of self-imaging properties within the SMS fiber structure is presented. The temperature dependence of the SMS structure is investigated by monitoring the modal field distribution and the variation of the intensity of the output beam at the self-imaging position of the MMF.

2 | COMPARISON OF MPA AND BPM ON THE CALCULATED SELF-IMAGING POSITION

Figure 1 shows the basic structure of the SMS structure, which comprises a short length of MMF sandwiched between two standard SMFs. The significant parameters of the SMF are the radius of the core and cladding being 4.15 and 62.5 μm respectively, and the refractive index of the core and cladding being 1.4504 and 1.4447, respectively. In the case of the MMF, the radii of the core and cladding are 52.5 and 62.5 μm respectively, and the refractive index of core and cladding is 1.4446 and 1.4271, respectively.

The guided-modal propagation analysis (MPA) has been previously applied to this type of structure and widely used to analyze the light propagation characteristics and the spectral response within the SMS fiber structure. According to waveguide propagation theory within the SMS structure, the electric field can be expressed as²¹

$$E(r, z) = \sum_{m=1}^M c_m F_m(r) \exp(i\beta_m z) \quad (1)$$

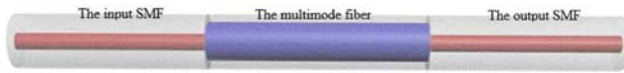


FIGURE 1 Schematic of SMS fiber structure [Color figure can be viewed at wileyonlinelibrary.com]

$$L_s(z) = 10 \log_{10} \left(\left| \sum_{m=1}^M c_m^2 \exp(i\beta_m z) \right|^2 \right) \quad (2)$$

where c_m is the excitation coefficient of the m th eigenmode, M is the excited mode number of the LP_{0m} multimode fiber, and β_m is the propagation constant of the m th eigenmode of the multimode fiber.

Using Equations 1 and 2, the spatial optical intensity distribution within the MMF can be obtained as shown in Figure 2. It is seen that the intensity changes as the light propagates along the MMF, and reaches a maximum within the axial distance range of [40,000, 45,000] (μm), which corresponds to the self-imaging position.

In order to obtain a comparison with the simulated results using MPA, the Beam PROP simulation engine (a part of RSoft Component Design Suite) based on the wide-angle beam propagation method (WA-BPM) was used to simulate the light transmission within the MMF. In order to achieve a more accurate simulated result, the Padé order (3, 3) and the boundary condition of perfectly matched layer (PML) were adopted. When an eigenmode of the SMF at the wavelength of 1550 nm is used as the input field, it is possible to calculate the field amplitude distribution of propagation within the MMF and this is shown in Figure 3. The maximum field intensity was also found to occur in the range of [40,000, 50,000]. Furthermore, the transmission of the SMS fiber structure was calculated and the approximate theoretical relationship between transmission loss and MMF length was obtained as¹⁵

$$L_s(l) = 10 \log_{10} \left[\frac{\left| \int_0^\infty E(l, r) F(r) r dr \right|^2}{\int_0^\infty |E(l, r)|^2 r dr \int_0^\infty |F(r)|^2 r dr} \right] \quad (3)$$

where l is MMF length, $E(l, r)$ is the calculated field at the interface between the MMF section and the output SMF, and $F(r)$ is the eigenmode of the output SMF.

Based on Equations 2 and 3, the transmission loss as a function of the length of the MMF was calculated and the result shown in Figure 4. From Figure 4, it is clear that a very small discrepancy exists between the calculated transmission spectra using MPA and BPM. Specifically, for the calculated results using the MPA method, the self-imaging position occurs at an axial position of 42,877 μm with a loss of -0.25 dB, while for that of the BPM, the self-imaging position is at 43,125 μm with a loss of -0.633 dB along the propagation direction of the MMF.

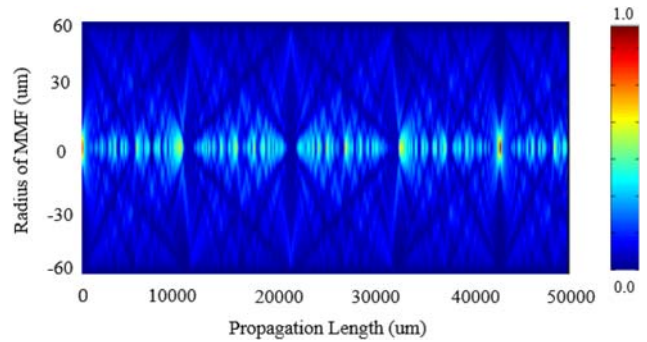


FIGURE 2 Optical intensity distribution within MMF using MPA [Color figure can be viewed at wileyonlinelibrary.com]

3 | EXPERIMENTAL DETERMINATION OF THE TRANSMISSION SPECTRAL OF THE SMS

Two identical SMS fiber samples were fabricated using a short piece of MMF about 43,000 μm spliced between two SMFs using a commercial fusion splicer (Fujikura 62S). The MMF lengths were initially approximately determined using a high precision cleaver (Fujikura CT-32), then accurately controlled by polishing using high grade fiber polishing papers. Finally, the length of the MMF of the first SMS sample was accurately set to 42877 μm , the self-imaging position as predicted by the MPA simulation. The MMF length of the second SMS sample was similarly set to 43125 μm , the self-imaging position predicted by the WA-BPM simulation. Figure 5 shows the experimental setup for measuring the transmission spectra of the SMS fiber sample. A Broadband Source (NKT) was used for launching the light signal into the SMS fiber samples and the optical spectrum analyzer (OSA) was used for recording the transmission spectra received. In Figure 6, the red dot line represents the transmission optical spectrum of the first SMS fiber sample. It is clear from Figure 6 that the central wavelength (the wavelength of minimum transmission loss) is located around 1549 nm. The black solid line represents the transmission optical spectrum of the second SMS fiber sample, and the central wavelength

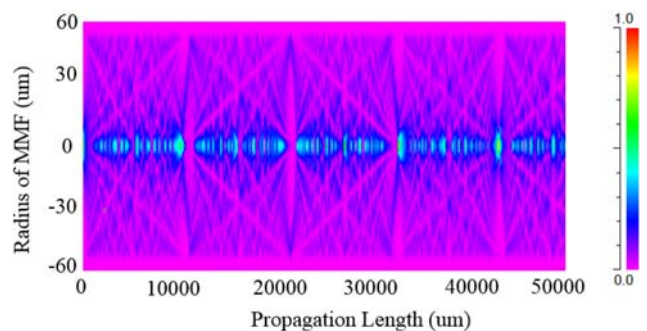


FIGURE 3 Optical intensity distribution within MMF using WA-BPM [Color figure can be viewed at wileyonlinelibrary.com]

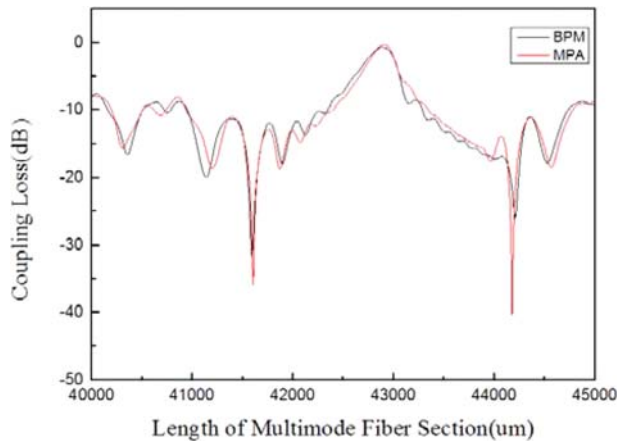


FIGURE 4 Comparison of the self-imaging phenomenon utilizing MPA (red line) and BPM (black line) [Color figure can be viewed at wileyonlinelibrary.com]

in this case is located close to 1545 nm. The results of Figure 6 indicate that the discrepancy between the experimental and the calculated self-imaging position using the MPA method is much smaller than that of the BPM. Therefore, the MPA method was selected for determining the self-imaging position of the SMS structure in the investigation described in this article.

4 | TEMPERATURE DEPENDENCE OF SELF-IMAGING PROPERTIES OF THE SMS FIBER STRUCTURE

The results of Figures 2–4 indicate that the self-imaging position plays a key role in determining the light propagation characteristics of the SMS fiber structure, and the MPA model provides the most accurate numerical vehicle for analysis and prediction of the SMS device performance.

The temperature dependence of the SMS fiber structure is well known and has been widely investigated over the past decade, but is exclusively based on wavelength shift or intensity variation monitored by an optical spectrum analyzer. The effect of temperature on the SMS fiber structure arises mainly from the temperature dependence of the refractive index and physical dimensional changes of the fiber core and cladding. These changes can be determined by thermo-optic coefficient (TOC) and the thermal expansion coefficient (TEC). The changes in core radius (R), MMF length (L), and refractive index (n) as a function of temperature variation (ΔT) can be expressed as follows¹⁶

$$R_{(\text{SMF,MMF})T} = R_{(\text{SMF,MMF})0} + \alpha \cdot R_{(\text{SMF,MMF})0} \cdot \Delta T \quad (4)$$

$$L_{(1,2)T} = L_{(1,2)0} + \alpha \cdot L_{(1,2)0} \cdot \Delta T \quad (5)$$

$$n_{(\text{core,clad})T} = n_{(\text{core,clad})0} + \xi \cdot n_{(\text{core,clad})0} \cdot \Delta T \quad (6)$$

where α and ξ are the TEC and TOC, respectively. In the case of the fiber types used in the structure of this investiga-

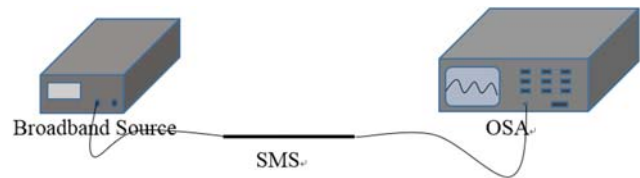


FIGURE 5 Experimental setup for measuring the transmission spectrum of SMS [Color figure can be viewed at wileyonlinelibrary.com]

tion, $\alpha = 5 \times 10^{-7} / ^\circ\text{C}$ and $\xi = 6.9 \times 10^{-6} / ^\circ\text{C}$ for both the SMF and MMF. It is evident that the TOC has a greater effect on the parameters of the SMS fibers than the TEC.

The temperature dependent transmission spectrum of the SMS Structure can be simulated using the MPA method taking the TOC and TEC into consideration. The relationship between transmission losses @1550 nm as a function of propagation distance in the axial region very close to the self-imaging position at the different temperatures is shown in Figure 7. Figure 7A shows the transmission loss in close proximity to the self-imaging position for a temperature range 20 to 100°C with an interval of 20°C. It is clear that there is a significant shift in the self-imaging position with changing temperature. It is also clear from Figure 7B, that the shift in self imaging position is a near ideal linear relationship with the temperature, having a linear regression coefficient value $R^2 = 1.0000$, and its slope is $0.3215 \mu\text{m}/^\circ\text{C}$. A knowledge of the self-imaging position can therefore be a useful tool to investigate the temperature dependent performance of the SMS fiber structure.

Using the MPA method, the output beam profile at the self-imaging position of the MMF fiber has been investigated, the length of the MMF in the simulation being chosen as 42,877 μm where self-imaging phenomenon is known to occur. The calculated 2D and 3D modal field distribution, calculated as normalized intensity referred to a maximum value of 3.29 dB at a temperature of 20°C and 100°C, are shown in Figure 8. From the simulated results in Figure 8A–D, it is clear

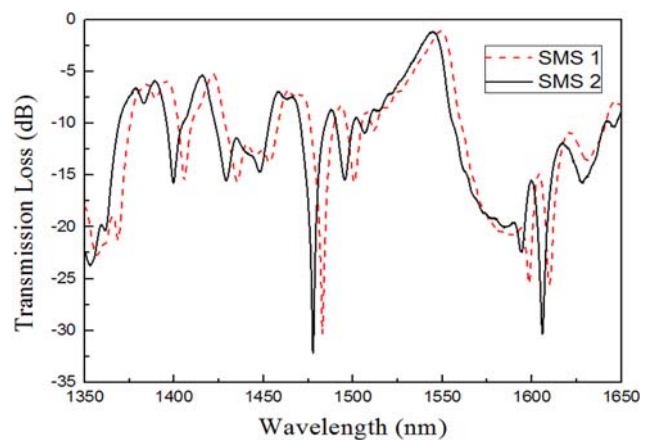


FIGURE 6 Measured transmission spectra of the studied SMS fibers [Color figure can be viewed at wileyonlinelibrary.com]

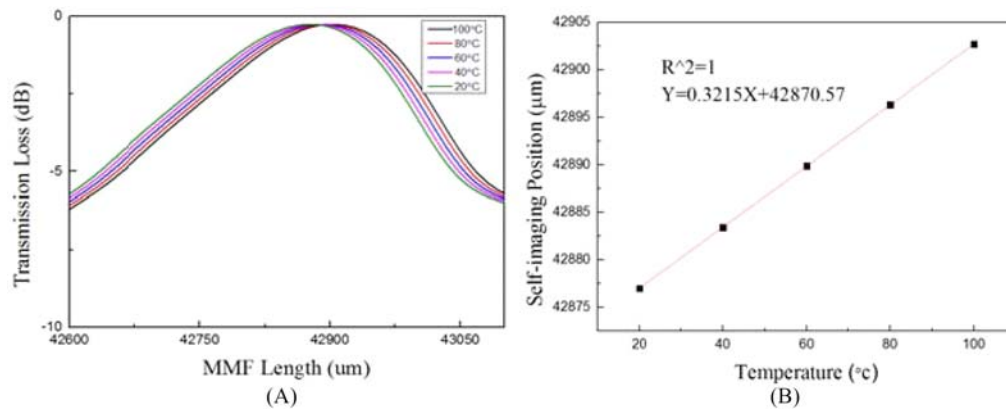


FIGURE 7 (A) The transmission spectra around self-imaging position at the different temperatures; (B) self-imaging position as a function of temperature and its linear fit [Color figure can be viewed at wileyonlinelibrary.com]

that it has a maximum intensity of -0.25 dB at the temperature of 20°C and as the temperature increases, the intensity decreases to -0.35 dB when the temperature increases to 100°C .

To verify the simulated results, the length of MMF was set at the self-imaging distance about $42,877\ \mu\text{m}$ and spliced with a single SMF, forming an SMF-MMF fiber structure. A tunable laser (Santec TSI-710) was used as the light source for launching the signal to the SMF-MMF structure at the source's wavelength of $1550\ \text{nm}$, the SMF-MMF fiber struc-

ture was placed on a thermoelectric Peltier cooler who's temperature was accurately controlled using a current controller (TTi, TSX3510P) as shown in Figure 9. A three-dimensional high precision translation stage with a displacement resolution of $50\ \text{nm}$ was used to adjust the position of the cooler along with the MMF fiber. The output light of the MMF was projected to a laser beam profiler (CCD, Newport), which was connected with a windows computer for monitoring the output modal field distribution, and the typical output is shown in Figure 10. In addition, an optical power meter was

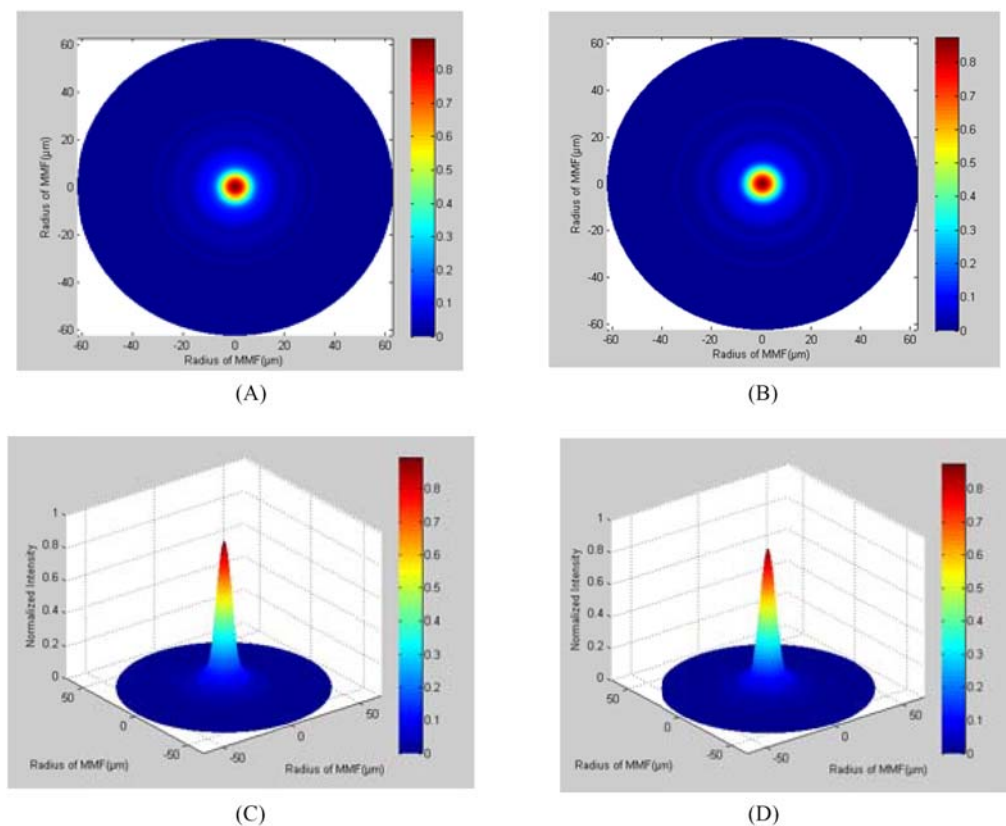


FIGURE 8 (A) and (C) 2D/3D modal field distribution at 20°C ; (B) and (D) 2D/3D modal field distribution at 100°C [Color figure can be viewed at wileyonlinelibrary.com]

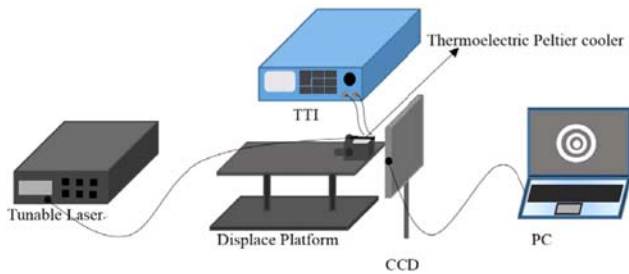


FIGURE 9 Schematic of experimental set-up measuring modal field distribution [Color figure can be viewed at wileyonlinelibrary.com]

placed at the end of the SMF-MMF fiber structure and was used to measure the optical intensity variations at the wavelength of 1550 nm at different temperatures. Figure 11 shows the output intensity as a function of temperature in the range 20–100°C. It indicates that the output intensity decreases when the surrounding temperature increases, and the measured results agree well with the simulated results (shown in Figure 8).

From Figure 10A,B and C,D, it is clear that at room temperature (20°C), the captured modal field has a maximum intensity of -7.849 dB, because it is closest to the modal field distribution maximum at the self-imaging position. When temperature increases to 100°C, the measured intensity drops to -8.733 dB because of the redshifts of the self-imaging position of the MMF. To better evaluate the difference between the theoretical model and the measured results presented in Figures 9 and 10, the optical intensities at a fixed wavelength as a function of different temperature was measured. Figure 11 shows both measured and simulated optical intensities at the wavelength of 1550 nm as a function

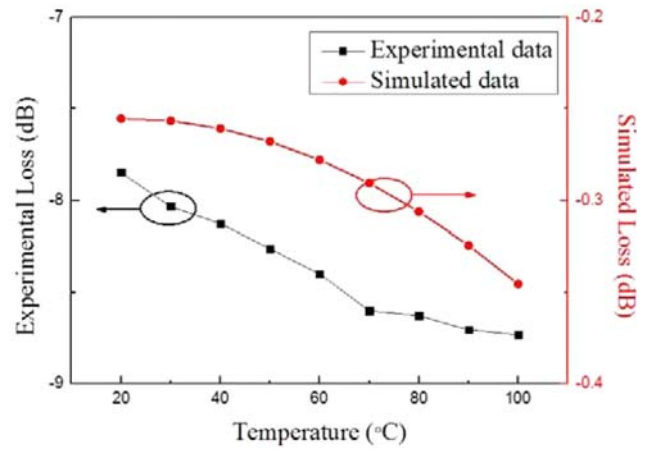


FIGURE 11 The optical intensities at a wavelength of 1550 nm as a function of temperatures (20–100°C) [Color figure can be viewed at wileyonlinelibrary.com]

of the ambient temperature ranging from 20 to 100°C with an interval of 10°C. From Figure 11, it is clear that there is general agreement, about 90% in difference between the measured results and simulated data. The discrepancy between the theoretical and experimental results are mainly caused by several impact factors, for example, airflow induced temperature vibrations and the resolution of the CCD etc.

5 | CONCLUSION

The multimode interference induced self-imaging phenomenon in the SMS fiber structure has been studied numerically and experimentally in this work. Detailed comparisons

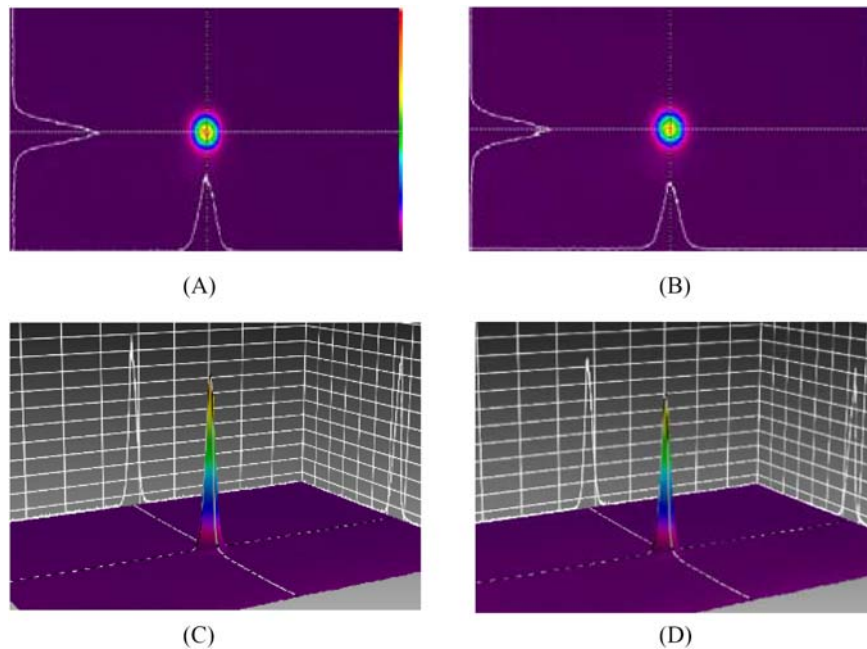


FIGURE 10 (A) and (C) 2D/3D modal field distribution at the room temperature (20°C); (B) and (D) 2D/3D modal field distribution at 100°C [Color figure can be viewed at wileyonlinelibrary.com]

between theoretical and experimental results have been carried out to investigate the accuracy of two different models based on the WA-BPM and MPA. The temperature dependence of the SMS fiber structure has also been investigated theoretically and experimentally. The good agreement between the experimental and simulated results indicates that the approximated guided-MPA method is superior when used to evaluate the optical performance of the SMS fiber structure or other multimode interference based photonic devices. The temperature dependent self-imaging behavior may offer excellent opportunities for multiple future applications including telecommunications, near-field micro-scanning/manipulation and compact sensing systems.

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Miniaturized high-band UWB array antenna

Fumiya Iwai¹ | Kuniaki Yoshitomi² | Haruichi Kanaya¹

¹Department of Electronics, Kyushu University, Fukuoka 819-0395, Japan

²E-JUST Center, Kyushu University, Fukuoka 819-0395, Japan

Correspondence

Haruichi Kanaya, Department of Electronics, Kyushu University, Fukuoka 819-0395, Japan.

Email: kanaya@ed.kyushu-u.ac.jp