One-neutron transfer study of ^{137}Xe and systematics of $13/2^+_1$ and $13/2^+_2$ levels in $N = 83$ nuclei

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Excited states in ^{137}Xe have been studied by using the near-barrier single-neutron transfer reactions ^{13}C(^{136}Xe,^{12}Cγ)^{137}Xe and ^{9}Be(^{136}Xe,^{8}Beγ)^{137}Xe in inverse kinematics. Particle-γ and particle-$\gamma\gamma$ coincidence measurements have been performed with the Phoswich Wall and Digital Gammasphere detector arrays. Evidence is found for a $13/2^+_2$ level ($E = 3137$ keV) and for additional high-lying $3/2^-$ and $5/2^-$ states. The results are discussed in the framework of realistic shell-model calculations. These calculations are also extended to the $13/2^+_1$ and $13/2^+_2$ levels in the $N = 83$ isotonic chain. They indicate that there is a need for a value of the neutron $0i_{13/2}$ single-particle energy ($E_{\mathrm{SPE}} = 2366$ keV) lower than the one proposed in the literature. It is also demonstrated that the population patterns of the $j = \ell \pm 1/2$ single-particle states in ^{137}Xe are different for the two targets used in these measurements and the implications of this effect are addressed.

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I. INTRODUCTION

New developments in instrumentation and experimental techniques for precision spectroscopy have recently enabled progress in the identification of single-neutron states in the ^{132}Sn ($Z = 50$, $N = 82$) region—see, e.g., Refs. [1–5]. As a result, for the $N = 83$ odd-mass nuclei from Sn to Sm ($Z = 62$), a nearly complete set of energy levels and spectroscopic factors related to the orbitals of interest is now available, i.e., for the $1f_{7/2}$, $2p_{3/2}$, $2p_{1/2}$, $1f_{5/2}$, $0h_{9/2}$, and the unique-parity $0i_{13/2}$ states. In a few cases, such as ^{133}Sn, excited-state lifetime information has been obtained as well [5]. This body of data has provided stringent tests of shell-model calculations [6,7], herewith improving their predictive power for properties of even more neutron-rich nuclei in this region.

While the systematics of the single-particle states in the $N = 82$–126 shell appears to be by now fairly complete, issues remain with the $0i_{13/2}$ orbital and the corresponding single-particle energy. Strictly speaking, this quantity should be derived from the energy spectrum of a nucleus with a single valence particle. However, the $13/2^+$ state in ^{133}Sn has not yet been observed. Due to this lack of information, researchers have resorted to (and reached meaningful conclusions from) the three available "sources" that shed light on the neutron $0i_{13/2}$ single-particle energy (SPE): (i) the 10^+ level in ^{134}Sb with a presumably dominant proton-neutron configuration $(\pi g_{7/2}\nu i_{13/2})_{10^+}$ [8,9]; (ii) the $27/2^-$ and $29/2^-$ states in ^{135}Sb likely associated with the configurations $[\pi g_{7/2}\nu(i_{13/2}f_{7/2})_{10^-}]_{27/2^-}$ and $[\pi g_{7/2}\nu(i_{13/2}h_{9/2})_{11^-}]_{29/2^-}$, respectively [10]; and (iii) the first $13/2^+$ level in ^{135}Te and heavier odd-mass $N = 83$ nuclei [4]. These types of information are complementary to each other, and all of them should be considered when dealing with the $\nu 0i_{13/2}$ orbital.

The subject of the present study is ^{137}Xe ($Z = 54$). For the description of the $13/2^+_1$ level in the $N = 83$ isotones, the interaction and mixing with a higher-lying second $13/2^+$ state should be taken into account [11]. This $13/2^+_2$ state is expected to be primarily composed of members of the $2^+ \otimes 0i_{13/2}$ and $3^- \otimes 1f_{7/2}$ multiplets [11]. These configurations represent quadrupole and octupole vibrational excitations of the $N = 82$ "core" coupled to the $13/2^+$ excited and $7/2^-$ ground states of the corresponding $N = 83$ isotope. For example, in ^{137}Xe, the expected $13/2^+_1$-$13/2^+_2$ admixture and the uncertainty of the single-particle energy of the neutron used in calculations are thought to be responsible for the comparatively poor agreement between the measured and the predicted $13/2^+_1$ level energy [6]. Note that the same work [6] indicates good agreement between theory and experiment for all the other states of interest. To address these issues, the location of the $13/2^+_2$ level in ^{137}Xe is helpful.

Previous experimental studies of ^{137}Xe have been carried out with an emphasis on the high-spin yrast sequence [12], on the low-spin structure [3,13], and on the precise location of the $13/2^+_1$ intruder state [4]; the resulting level scheme is proposed in Ref. [14]. The present study follows the experimental approach of Ref. [4], i.e., it uses one-neutron transfer reactions in inverse kinematics with a ^{136}Xe beam on two different targets, ^{13}C and ^{9}Be. The observed difference in the population of excited states in ^{137}Xe is analyzed in detail, while new levels are reported, including the $13/2^+_2$ level of interest.

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The experimental findings are complemented by shell-model calculations. A lower value of the $\nu 0i_{13/2}$ single-particle energy than that used in Ref. [8] is introduced to calculate both the $13/2^+_1$ and $13/2^+_2$ states for a series of odd-mass $N = 83$ nuclei, and the energy levels in ^{134,135}Sb associated with the $i_{13/2}$ orbital. The overall agreement between theory and experiment is found to be satisfactory. The comparison of the results obtained with the two $0i_{13/2}$ energy values used in the calculations allows one to propose a realistic range for the $\nu 0i_{13/2}$ single-particle energy.

The $13/2^+_2$ level is observed only with the ^{13}C target. Likewise, the data indicate that the population of the $13/2^+_1$ state is enhanced when compared to the data set from the ^{9}Be target. The present study was carried out, in part, with the intent to verify that the different population patterns follow the ℓ- and j-selection rules for nucleon transfer discussed in Ref. [15]. For example, the population of a $0i_{13/2}$, $j = \ell + 1/2$ ($j_>$) state is expected to be more likely with the odd nucleon (of either the target or the beam) residing in a $j = \ell - 1/2$ ($j_<$) rather than a $j_>$ orbital. Hence, by using a ^{13}C target with a $0p_{1/2}$ ($j_<$) valence neutron, a transfer to the $0i_{13/2}$ orbital is favored over that to the available $j_<$ orbitals. Conversely, using the same beam and a ^{9}Be target, with a $0p_{3/2}$ ($j_>$) valence neutron, should introduce a preference for feeding the $j_<$ orbitals, such as the $2p_{1/2}$ and $1f_{5/2}$ ones, in ^{137}Xe.

II. EXPERIMENTAL CONDITIONS AND ANALYSIS PROCEDURES

The experiment was performed at the ATLAS accelerator at Argonne National Laboratory. A 560-MeV ^{136}Xe beam impinged on two targets: a ^{13}C foil with 99% isotopic enrichment and 0.15-mg/cm^2 thickness and a monoisotopic 1.5-mg/cm^2 ^{9}Be foil. The beam intensities were 500 and 70 ppA, and the runs took 40 and 14 h with the ^{13}C and ^{9}Be targets, respectively. The detection setup comprised Digital Gammasphere [16], with 92 Compton-shielded HPGe detectors arranged in 16 angular rings around the target [17], and the Phoswich Wall, a 256-element fast-plastic + CsI(Tl) charged-particle detector array [18]. The latter array was located downstream from the target, with a laboratory-angle coverage of $9^\circ \leqslant \theta_{\text{lab}} \leqslant 72^\circ$, and enabled the correlation of a specific targetlike fragment (TLF) with coincident γ rays emitted by the corresponding projectilelike fragment (PLF). The event trigger required that a Phoswich Wall element and at least two HPGe detectors fired prior to suppressing Compton-scattering signals.

In the offline analysis, a gating procedure was applied to the so-called (A,C) particle map, a combination of the fast-plastic and "late" CsI(Tl) signals (see Ref. [18] for details), *and* the prompt peak of the time spectrum of the measured particle with respect to the accelerator radiofrequency. The map gating condition for the different runs required the presence of TLF carbon or $^8\text{Be} \rightarrow 2\alpha$ events. A typical particle map for the data obtained with the ^{9}Be target is presented in Fig. 1. The Doppler-shift correction applied to the γ-ray spectra relied on the PLF velocity vector reconstructed event by event for the binary reaction and took advantage of the high degree of pixelation of the Phoswich Wall.

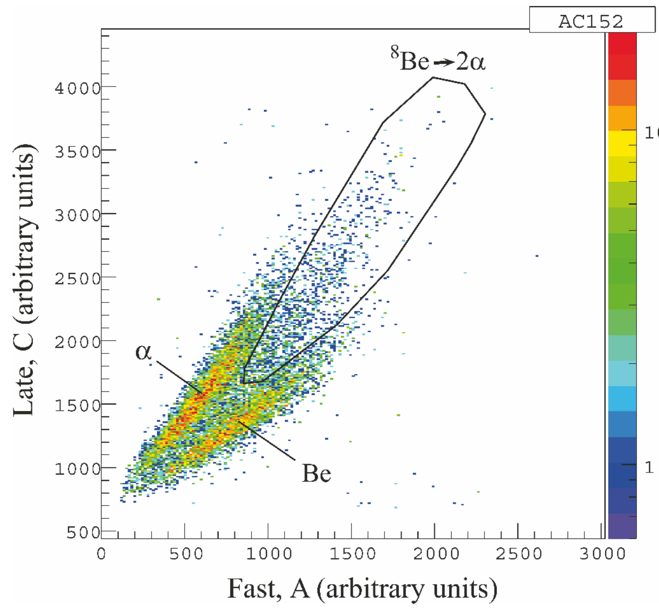

FIG. 1. An (A,C) particle map for the data taken with the ^{9}Be target. The two-dimensional gate selects $^8\text{Be} \rightarrow 2\alpha$ events. The other two particle groups represent beryllium and α-particle/^{5}He events associated with projectile Coulomb excitation and incomplete fusion, respectively. The map is for pixel 152 of the Phoswich Wall. The color scale shown on the right provides the range of z-axis values in this example.

For each run, two types of two-dimensional histograms of γ-ray energies were created: (i) an E_γ-E_γ matrix and (ii) a set of "angle-dependent" $E_\gamma(\chi)$-E_γ (any) matrices. The matrices in (ii) allowed measuring the γ-ray anisotropies relative to the spin direction of the fragment nucleus following the procedure of Ref. [19]. Here, χ represents the angle between the emitted γ ray and the spin direction (binned into 10° increments), whereas any stands for no angle requirement. All these histograms were analyzed using the RADWARE analysis package [20].

III. EXPERIMENTAL RESULTS

This section reports new information on the ^{137}Xe level scheme obtained in the present experiment. As different population patterns were observed for the excited states in ^{137}Xe as a function of the target used, the information is summarized in two separate diagrams.

Representative γ-ray spectra for the desired one-neutron transfer channel in both reactions are provided in Figs. 2 and 3, respectively, and are used to justify the locations of new transitions in the level schemes of Fig. 4. The total projection in panel (a) of Fig. 2 has been produced with a carbon gate, and the strongest γ rays correspond to one-neutron pickup and transfer (^{135,137}Xe) as well as to projectile Coulomb excitation (^{136}Xe). In the case of ^{135}Xe, some of the γ rays are only assigned tentatively and labeled as such by a filled symbol [21]. In addition to the xenon nuclei, the proton-transfer product ^{137}Cs and, in a lesser amount, the ^{139,140}Ba nuclei are present as well; the latter originate from $^{136}\text{Xe} + \alpha$ or ^{5}He incomplete

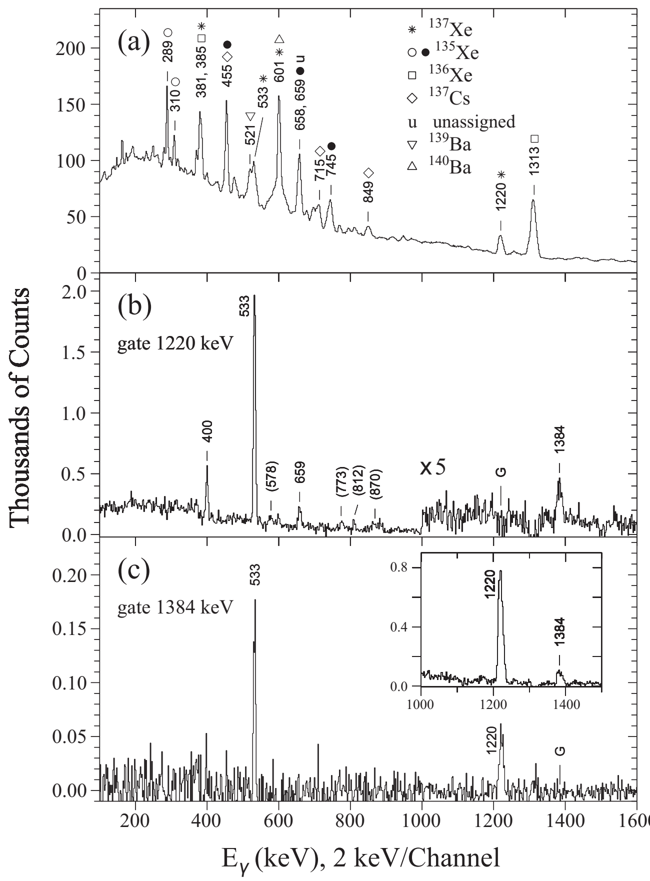

FIG. 2. Representative γ-ray spectra for the 560-MeV ^{136}Xe + ^{13}C reaction. (a) Total projection of a γ-γ coincidence matrix gated with carbon ions in the A-C map (see text). Transitions are labeled by their energies in keV. For ^{135}Xe, the open and filled symbols indicate firm and tentative assignments, respectively. (b) Coincidence spectrum for ^{137}Xe obtained by gating on the 1220-keV transition. Note the change in scale at a γ-ray energy of 1000 keV. Transitions with labels given in parentheses are associated with an interfering 1218-keV line (see text). (c) Similar to (b) but with gating transitions of 1384 (main panel) and 533 keV (inset). The position of the gate is indicated by the letter G.

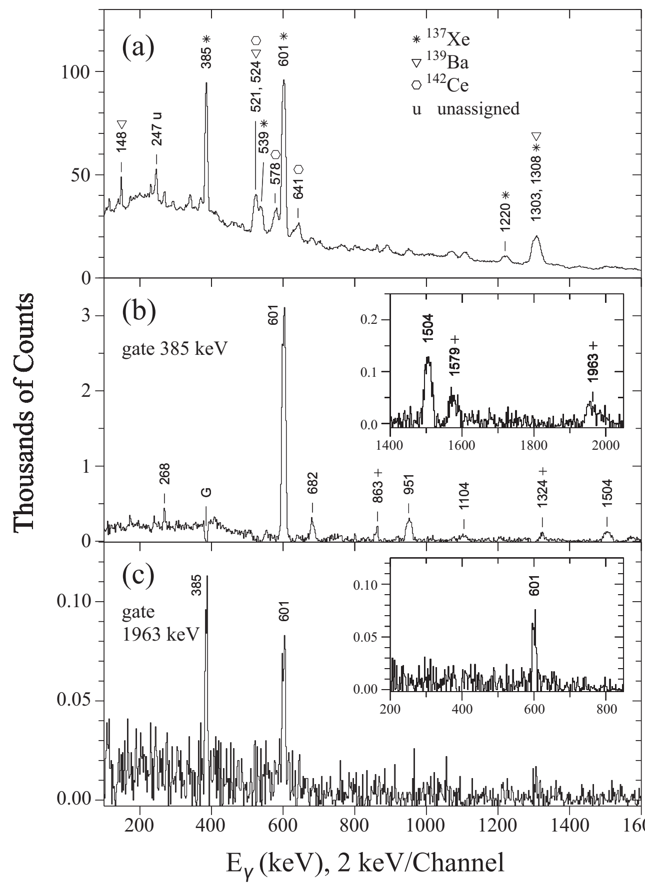

FIG. 3. Representative γ-ray spectra for the 560-MeV ^{136}Xe + ^{9}Be reaction. (a) Total projection of a γ-γ matrix gated on 2α events (see text). Note that the 1303-keV peak is composed of a ^{137}Xe ground-state transition and ^{139}Ba lines with 1306 keV $\leqslant E_\gamma \leqslant$ 1319 keV. (b) Coincidence spectrum for ^{137}Xe obtained by gating on the 385-keV transition. The inset provides the extension of the spectrum toward higher energies. The letter G indicates the position of the gate. The γ rays labeled by a plus sign are newly observed. (c) Similar to (b) but for the 1963- and 2349-keV gating transitions. The latter gated spectrum is shown in the inset.

fusion reactions. In these cases, ^{12}B (^{137}Cs) or one of the remaining ^{9}Be or ^{8}Be fragments (^{139,140}Ba) was detected and has leaked into the particle coincidence gate.

Panels (b) and (c) display carbon- and γ-gated spectra for ^{137}Xe. The spectrum gated on the 1220-keV ground-state transition [panel (b)] displays the expected lines [4,14] and a new 1384-keV γ ray. In turn, the 1384-keV gate of panel (c) shows the 533- and 1220-keV lines only. This observation requires the placement of the new γ ray on top of the 533-keV $13/2^+ \rightarrow 11/2^-$ transition as proposed in panel (a) of Fig. 4. The new state at 3137 keV is the most notable feature of this level scheme. Note that it is absent in Fig. 4(b).

Figure 3 presents a set of representative ^{137}Xe spectra from the ^{9}Be data. The total projection of panel (a) is dominated by the 601- and 385-keV transitions from the first and second excited states in ^{137}Xe. These states are crucial for determining the low-spin part of the level scheme. A comparison of the intensities of the 601- and 1220-keV ground-state transitions with those in panel (a) of Fig. 2 confirms that the choice of the two targets leads to the anticipated differences in the population pattern of the nucleus (see Table I and related discussion). In Fig. 2(a), binary reaction products other than ^{137}Xe are also substantially reduced. Note that projectile Coulomb excitation is excluded by the coincidence requirement of 2α events. Hence, the competing reaction channels are mainly those resulting from incomplete fusion. In this context, the Ce lines are attributed to ^{136}Xe reactions on the oxygen originating from target oxidation.

Panels (b) and (c) display 2α- and γ-gated coincidence spectra for ^{137}Xe. The coincidence spectrum gated by the 385-keV line leads to the observation of new ^{137}Xe transitions with respective energies of E_γ = 863, 1324, 1579, and 1963 keV. In addition, the spectrum gated by the 601-keV ground-state transition (not shown) suggests that a weak 2349-keV γ ray bypasses the 1963-keV transition. A spectrum gated on one of

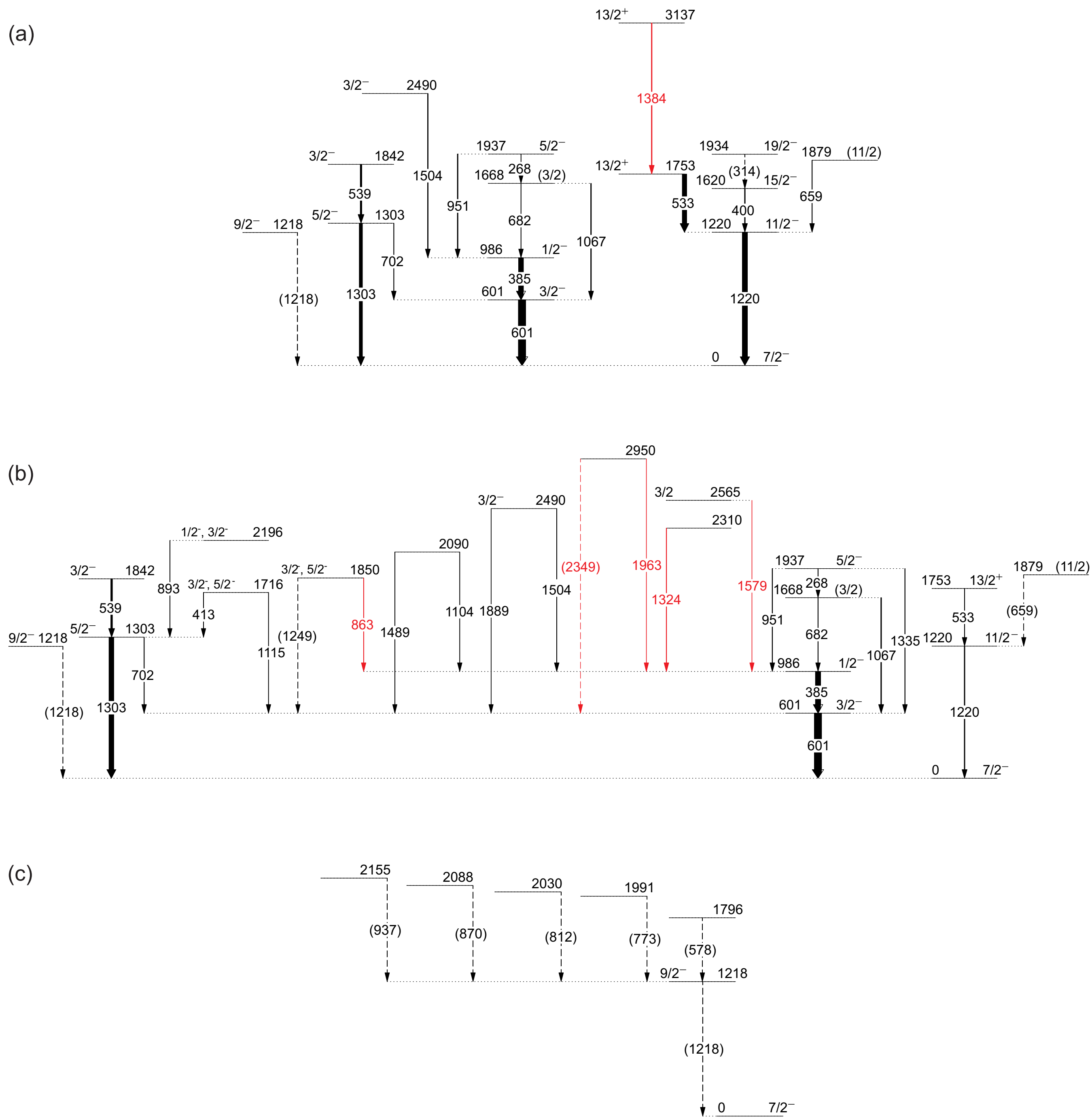

FIG. 4. The level scheme for ^{137}Xe obtained in the 560-MeV ^{136}Xe + ^{13}C (a) and ^{136}Xe + ^{9}Be (b) reactions with the feeding of the 9/2^{-} state shown separately (c). The widths of the arrows are proportional to the measured γ-ray intensities. The energies are in keV. The assignments given in parentheses are tentative. The transitions marked in red are new to this work.

these new transitions ($E_\gamma = 1963$ keV) is presented in panel (c): it shows the 385- and 601-keV lines only. The 1579-keV transition confirms the existence of a previously reported level ($E = 2565$ keV). Likewise, the 863-keV γ ray represents a newly observed decay branch of a known state, whereas the 1324-keV line and the 1963- and 2349-keV pair of transitions establish two new levels.

For most of the newly observed transitions, spin and parity assignments are proposed based on a γ-ray angular-distribution analysis. Figure 5 provides sample angular distributions for three γ rays measured in the carbon data: a known stretched electric dipole ($E1$) (a) and quadrupole ($E2$) transition (b), and the newly observed 1384-keV line (c). These have been fitted with a standard Legendre polynomial expression, and the fit results are included in the figure. The characteristic A_2/A_0 and A_4/A_0 coefficients derived from the fits are reported in Table I, together with other information on the transitions of interest. Values obtained for known transitions

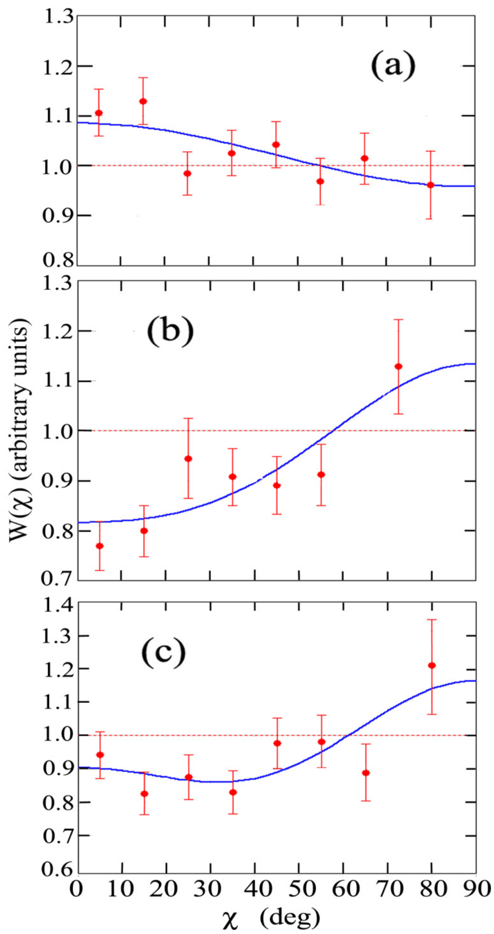

FIG. 5. Representative γ-ray angular distributions, with respect to the spin direction, for transitions in ^{137}Xe (^{136}Xe + ^{13}C data). Panels (a)–(c) are for the 533-, 400-, and 1384-keV transitions, respectively. The former two cases are established stretched dipole ($E1$) and quadrupole ($E2$) transitions, respectively.

[Figs. 5(a) and 5(b)] are in line with expectations for the present analysis, using the spin-alignment method, where stretched dipole and quadrupole transitions have positive and negative A_2/A_0 coefficients, respectively. The angular distribution of the 1384-keV transition is consistent with either a quadrupole ($E2$) or an unstretched-dipole (no spin change) assignment. However, the large value of the A_4/A_0 coefficient indicates a mixed multipole character. Since a transition of the $E2 + M3$ type is unlikely, the $M1 + E2$ assignment is preferred. Here, the $E2$ mixing fraction is estimated to be $18^{+7}_{-4}\%$. In addition, the following considerations apply: (i) the new level at 3137 keV is observed only with the ^{13}C target, herewith suggesting a $13/2^+$ rather than a $17/2^+$ assignment; (ii) the $13/2^+_2$ states in the $N = 83$ isotones ^{143}Nd and ^{145}Sm also decay to the respective $13/2^+_1$ levels [14]. These additional arguments support the conclusion reached from the angular distribution, and a $13/2^+_2$ assignment follows for the 3137-keV level.

In the present study, the assignments of important known transitions have been confirmed as well. For the 986-keV level, the A_2/A_0 coefficient of the 385-keV line is consistent with 0, indicating an isotropic transition from this state and supporting a $1/2^-$ assignment. Although the A_2/A_0 coefficient has a large error, this result removes the previous $1/2^-$,$3/2^-$ ambiguity [3]. The isotropy of the 601-keV $3/2^- \rightarrow 7/2^-$ γ ray (again with a large error for the A_2/A_0 coefficient) is attributed to the loss of alignment at the 986-keV $1/2^-$ state, which is by far the strongest feeder level of the above ground-state transition. The 1936-keV level is reassigned as $5/2^-$ since the rather strong 951-keV γ ray is of quadrupole ($E2$) character.

The decay of the 1879-keV level to the $11/2^-_1$ 1220-keV state is confirmed; the 659-keV transition has an intensity of $I_\gamma = 4.5(6)$ in the ^{13}C measurement ($I_\gamma \lesssim 3$ in the ^{9}Be run). This 1879-keV state is tentatively assigned 11/2 based on intensity considerations. An alternative 13/2 assignment seems to be ruled out since, in the ^{13}C measurement, this 1879-keV level is less populated than the $13/2^+_2$ off-yrast state. Similarly, a 9/2 assignment appears unlikely in view of (i) the weak population of the $9/2^-_1$ 1218-keV level, which is comparable to the intensity of the 659-keV transition (see below) and (ii) the nonobservation of the $(9/2^-_2)$ 1590-keV level in the present experiment. Hence, the 1879-keV level is viewed as a candidate for the $11/2^-_2$ state.

The lowest-lying states of single-particle character in ^{137}Xe are populated with both targets, but with markedly different strengths: in the ^{13}C measurement, the $13/2^+_1$ level ($0i_{13/2}$ candidate) and the $11/2^-_1$ state to which it decays are prominently present, whereas, in the ^{9}Be data, these states are weakly populated compared to the $5/2^-_1$ ($1f_{5/2}$), $1/2^-$ ($2p_{1/2}$), and $3/2^-_1$ ($2p_{3/2}$) levels. Despite this difference in the population patterns of the single-particle states, the ^{13}C and ^{9}Be measurements share the common feature that the known high-spin yrast levels, which feed the $11/2^-_1$ state [12], are suppressed. Specifically, the populations of the $15/2^-$ 1620-keV level and its feeder states are not competitive with that of the $13/2^+_1$ level. The present findings are depicted in Fig. 6 where the decay intensity for a given level is plotted as a function of the excitation energy. Here, corrections for internal conversion have been applied where possible. These are small compared to the uncertainties of the γ-ray intensities.

Note that the $3/2^-$, $5/2^-$ 1716-keV (decay-intensity $\leqslant$6), and (11/2) 1879-keV levels have been excluded from Fig. 6 for simplicity. Both the ^{13}C and the ^{9}Be measurements also provide evidence for the population of the $9/2^-_1$ 1218-keV level, which is partially fed by transitions with $E_\gamma = 578$, 773, 812, and 870 keV [cf. Fig. 2(b)] and directly decays to the ground state [14] . However, as was the case in Ref. [3], the population of the $9/2^-_1$ level ($0h_{9/2}$ candidate) is weak with respect to, e.g., the $5/2^-_1$ state as the aforementioned feeder transitions have a combined intensity of $\Sigma I_\gamma \lesssim 3$. In view of these low intensities, it is not possible to draw a conclusion about a potential difference in the population of the $9/2^-_1$ level between the two data sets.

IV. DISCUSSION

A. Differences in the two reactions

The differing population patterns of Fig. 6 follow the ℓ- and j-selection rules expected for one-nucleon transfer [15]. Clearly, the $0i_{13/2}$ ($j_>$) state is populated more strongly with the ^{13}C target, where the valence neutron occupies

TABLE I. Information for selected γ-ray transitions in ^{137}Xe from the present work. The table is organized according to the level schemes of Fig. 4.

E (keV)[a]	$I_i^\pi \rightarrow I_f^\pi$ [b]	E_γ (keV)	I_γ [c]	A_2/A_0	A_4/A_0
		(a) ^{13}C target			
1220.1(5)	$11/2_1^- \rightarrow 7/2^-$	1220.1(5)	71(4)	$-0.049^{+0.034}_{-0.039}$	$0.011^{+0.034}_{-0.039}$
1620.0(7)	$15/2^- \rightarrow 11/2_1^-$	399.9(3)	9.7(11)	$-0.232^{+0.027}_{-0.023}$	$0.048^{+0.029}_{-0.049}$
1752.8(6)	$13/2_1^+ \rightarrow 11/2_1^-$	532.7(3)	61(4)	$0.085^{+0.024}_{-0.032}$	$0.000^{+0.044}_{-0.044}$
3137(1)[d]	$13/2_2^+ \rightarrow 13/2_1^+$	1384(1)[e]	9.1(11)	$-0.229^{+0.044}_{-0.045}$	$0.132^{+0.046}_{-0.069}$
		(b) ^{9}Be target			
601.1(4)	$3/2^- \rightarrow 7/2^-$	601.1(4)	100(4)	~ 0[f]	~ 0[f]
986.4(5)	$1/2^- \rightarrow 3/2_1^-$	385.3(3)	71(3)	~ 0[f]	~ 0[f]
1220.1(6)	$11/2_1^- \rightarrow 7/2^-$	1220.1(5)	9.6(10)		
1753(1)	$13/2_1^+ \rightarrow 11/2_1^-$	532.7(3)	4.3(6)		
1850(1)	$3/2^-,5/2^- \rightarrow 1/2^-$	863(1)[e]	4.2(5)		
1937(1)	$5/2^- \rightarrow 1/2^-$	951(1)	10(1)	$-0.113^{+0.007}_{-0.011}$	$0.029^{+0.011}_{-0.011}$
2310(2)[d]	- $\rightarrow 1/2^-$	1324(2)[e]	4.3(5)		
2490(2)	$3/2^- \rightarrow 1/2^-$	1504(2)	6.2(6)	$0.042^{+0.027}_{-0.040}$	$-0.039^{+0.061}_{-0.029}$
2565(3)	$3/2 \rightarrow 1/2^-$	1579(3)[e]	4.7(5)	$0.140^{+0.090}_{-0.098}$	$0.178^{+0.148}_{-0.089}$
2950(3)[d]	- $\rightarrow 1/2^-$	1963(3)[e]	3.4(5)		
2950(3)[d]	- $\rightarrow 3/2^-$	2349(3)[e]	2.1(4)		

[a]Energy of the depopulated state.
[b]Spins and parities of the levels linked by the transition involved.
[c]Relative γ-ray intensity of the transition normalized to 100 for the 601-keV ground-state transition.
[d]Newly observed level.
[e]Newly observed γ ray.
[f]A_k/A_0 ($k = 2$ or 4) coefficient consistent with 0.

the $0p_{1/2}$ ($j_<$) orbital, whereas, for the ^{9}Be target with its odd $0p_{3/2}$ ($j_>$) neutron, the states in ^{137}Xe based on the $2p_{1/2}$ and $1f_{5/2}$ ($j_<$) orbitals are preferred. In the following considerations, the net intensity, obtained from the total intensities out of and into a level ($I_{\rm net} = \Sigma I_{\rm out} - \Sigma I_{\rm in}$), is taken as a measure of the direct population of the state of interest. This quantity is, after a common normalization, compared with the cross sections from distorted-wave Born approximation (DWBA) calculations. The latter were performed with the FRESCO code [22] where the midtarget energies of the reactions were used and the assumption was made that the states are based on pure single-neutron configurations. The pertinent details of this comparison are summarized in Table II. The $I_{\rm net}$ values commonly represent a considerable fraction ($\sim$1/2–3/4) of the corresponding decay intensities in Fig. 6. Hence, each part of the present comparison is affected by unobserved side feeding, but to a comparable degree. Uncertainties in the calculations on the spectroscopic factors impact the comparison as well. (The $I_{\rm net}$ values for the $3/2^-$ state, which are not part of the present comparison, are close to zero.) Given that the procedure has limited accuracy, the $I_{\rm net}^{\rm rel}$ and $\sigma_{\rm DWBA}^{\rm rel}$ values of Table II exhibit a reasonably close correspondence. Particularly, the very different relative yields for the $13/2_1^+$ level in the two reactions are accounted for.

The weak population of the yrast states with $I \geqslant 15/2$ can, perhaps, be explained in a similar fashion. Since the ground states of ^{137}Xe and the target correspond to orbital angular momenta $\ell = 3$ and 1, respectively, the angular momentum transfer should not exceed a total value of 4; i.e., a spin difference of $15/2 - 7/2$.

B. Comparisons with shell-model calculations

In this section, a comparison is carried out between, on one hand, the experimental $13/2_1^+$ and $13/2_2^+$ levels in ^{137}Xe and the neighboring isotones and, on the other hand, the results of realistic shell-model calculations. The new information on other levels in ^{137}Xe is addressed by the calculations as well.

The relevant $N = 83$ systematics are presented in Fig. 7. Besides ^{137}Xe, the $13/2_2^+$ states have been established in the neighboring odd-mass nuclei ^{143}Nd and ^{145}Sm ($Z = 60,62$), whereas candidate $13/2_2^+$ levels are reported for ^{139}Ba, ^{141}Ce, and ^{147}Gd ($Z = 56$, 58, and 64) [14]. For the latter levels, the decay is unknown, and/or the spin-parity assignment is uncertain. In Fig. 7, these $13/2_1^+$ and $13/2_2^+$ levels are also compared with the results of realistic shell-model calculations. The latter use the two-body effective Hamiltonian introduced in Ref. [6]. These matrix elements were derived from the CD Bonn nucleon-nucleon potential [23], by way of time-dependent perturbation theory [24], within a model space including the $0g_{7/2},1d,2s,0h_{11/2}$ and $1f,2p,0h_{9/2},0i_{13/2}$ orbitals for protons and neutrons, respectively. As in Ref. [6], the adopted single-proton and single-neutron energies were taken, where possible, from inspection of the ^{133}Sb and ^{133}Sn level schemes [14]. However, the $\pi s_{1/2}$ and $\nu i_{13/2}$ energies are not available yet from these two semimagic nuclei, and thus, additional information is needed. The position of the proton

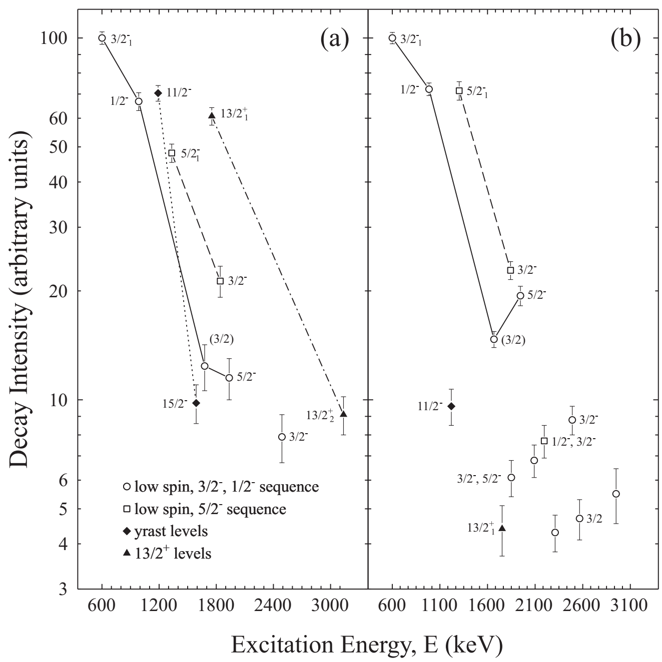

FIG. 6. Decay intensity versus energy for levels in ^{137}Xe obtained in the (a) carbon and (b) beryllium measurements. The sequences of levels based on the $3/2^-_1$ and $1/2^-$ states, the $5/2^-_1$ level, the $11/2^-_1$ yrast, and the $13/2^+_1$ unique-parity states are represented by circles, squares, diamonds, and triangles, respectively. The data points are labeled by the corresponding spin-parity assignments [except the levels in panel (b) with $E = 2090$, 2310, and 2949 keV]. The most strongly populated states in a sequence are connected by straight lines.

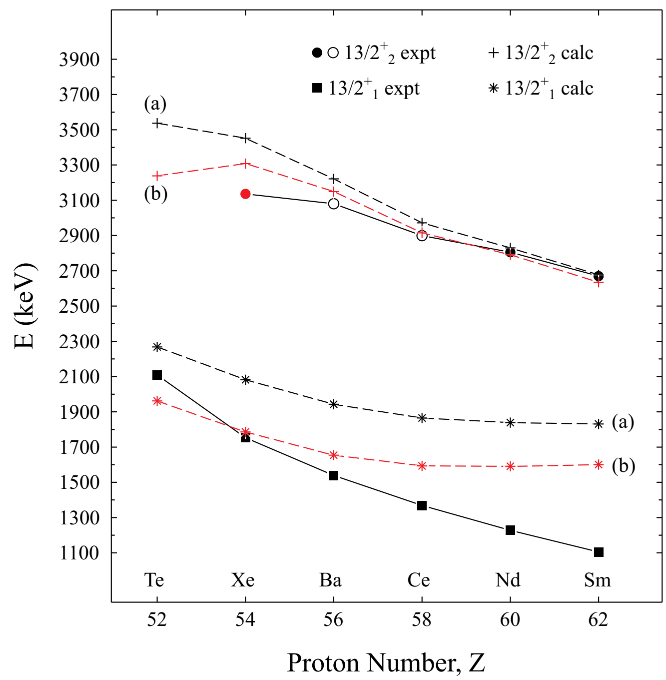

FIG. 7. The $13/2^+_1$ and $13/2^+_2$ systematics from Refs. [4,14] and the present data (red circle), and results from two sets of calculations described in Ref. [6] and in the text (red crosses and asterisks), and presented in the text as (a) calc1 and (b) calc2, respectively. The filled and open circles distinguish between firmly established and candidate $13/2^+_2$ levels, respectively. See the text for details.

$2s_{1/2}$ orbital was determined based on the $1/2^+$ 2150-keV level in ^{137}Cs [25]. For the position of the neutron $0i_{13/2}$ orbital, two alternative procedures have been used here, and the corresponding calculations are referred to hereafter as "calc1" and "calc2." For the calc1 computations, the $\nu 0i_{13/2}$ energy was estimated based on the 10^+ level in ^{134}Sb, which has been assigned a $\pi g_{7/2}\nu i_{13/2}$ configuration and is located at 2434 keV with respect to the yrast 7^- state [8,9]. This procedure leads to a value of 2694 keV and is the same as that adopted in Ref. [6]. For the calc2 calculations, the estimate has been modified such that the energy of the $13/2^+_1$ level in ^{137}Xe is reproduced. In this context, this $13/2^+_1$ level is essentially described as a $0i_{13/2}$ neutron coupled to the proton wave function, which represents the "core"; i.e., its amplitude squared is 92% in the total wave function, as discussed further below. The adopted value of the $\nu 0i_{13/2}$ SPE is 2366 keV in this approach.

TABLE II. Measured relative yields (normalized net total intensities) and results of DWBA calculations for prominent states in the one-neutron transfer study of ^{137}Xe.

E (keV)	I^π	^{13}C target			^{9}Be target		
		I_{net} [a]	I^{rel}_{net} [b]	σ^{rel}_{DWBA} [c]	I_{net} [a]	I^{rel}_{net} [b]	σ^{rel}_{DWBA} [c]
986	$1/2^-$	47(4)	100(9)	100	32(3)	100(11)	100
1303	$5/2^-_1$	27(4)	58(8)	80	39(5)	124(14)	120
1753	$13/2^+_1$	52(4)	111(8)	89	4.3(6)	14(2)	8

[a]Decay intensity of the level minus the observed feeding intensity.
[b]Net intensity normalized to 100 for the $1/2^-$ level.
[c]Relative cross section calculated for the level normalized to 100 for the $1/2^-$ state.

Focusing first on the $13/2^+_1$ state, the calculations with the two effective interactions display (Fig. 7) the same general trend with proton number with the calc2 results being, in general, closer to the data than the calc1 ones because of the choice of the $\nu 0i_{13/2}$ single-particle energy. The agreement between experiment and theory, although satisfactory for ^{135}Te and ^{137}Xe, becomes gradually poorer for $Z > 56$. This is likely, in part, due to the fact that the two-body effective interaction has been derived from data on systems with two valence nucleons, which may be less adequate for the description of nuclei with a larger number of such nucleons where many-body effects could play a role. In addition, as Z increases, the nature of the $13/2^+_1$ state itself may well be changing, as discussed further below. Regarding the $13/2^+_2$ level, a dependence on the $\nu 0i_{13/2}$ single-particle energy is visible only at low Z, e.g., for ^{135}Te and ^{137}Xe, with the calc2 results closer to the data than the calc1 ones. Thus, the trend exhibited by the experimental data for the $13/2^+_2$ levels as a function of Z is reproduced well, in general, with the largest discrepancies occurring closer to $Z = 50$, a behavior opposite to that noted for the $13/2^+_1$ states.

The applicability of the calc2 approach was checked further for the relevant states in ^{134}Sb and ^{135}Sb, and the results are summarized in Table III. For the 10^+ state in ^{134}Sb, the

TABLE III. A comparison of calculated and experimental energies of states in ^{134}Sb and ^{135}Sb associated with the $i_{13/2}$ neutron orbital.

Nucleus	I^π	E (keV)		
		calc1	calc2	Expt.
^{134}Sb	10^+	2824[a]	2515	2713[b]
^{135}Sb	$27/2^-$	3557	3468	3249[c]
^{135}Sb	$29/2^-$	4416	4119	3688[c]

[a]Reference [6].
[b]Adopted value based on the information in References [8,9,26].
[c]Reference [10].

calc1 energy is closer to the adopted experimental value of 2713 keV than the value computed for calc2. On the other hand, a better agreement is reached for the $27/2^-$ and $29/2^-$ experimental levels in ^{135}Sb with the calc2 calculations. This finding presumably illustrates the fact that an optimal value for the $\nu 0i_{13/2}$ single-particle energy has yet to be found. However, it also suggests that a value lower than the one proposed in the literature, $E_{\rm SPE} = 2694$ keV [6], is required. From the discussion above, it can be concluded that the recommended value for the $\nu 0i_{13/2}$ single-particle energy to be used in shell-model calculations for $N = 83$ and 84 nuclei should be in the 2360-keV $\leqslant E_{\rm SPE}(\nu 0i_{13/2}) \leqslant$ 2600-keV range.

As alluded to above, the character of the $13/2_1^+$ state may well be changing as Z increases. Figure 8(a) presents the calc2 values of the $\nu 0i_{13/2}$ effective single-particle energy (ESPE), with respect to the $1f_{7/2}$ ESPE of the ground state, in the isotonic chain together with the energy of the 3^- state in the corresponding $N = 82$ nucleus. It can be seen that the 3^- excitation decreases rapidly in energy with Z and approaches the $\nu 0i_{13/2}$ ESPE value. As a result, sizable admixtures of the $3^- \otimes 1f_{7/2}$ configuration into the $13/2_1^+$ wave function are to be expected. Panel (a) also displays a constructed energy that is obtained by adding the excitation energies of each of the yrast 2^+ and 4^+ levels in the $N = 82$ nucleus to the $\nu 0i_{13/2}$ ESPE value of the corresponding isotope. These $N = 82$ excitation energies are provided separately in panel (b) of the figure. The set of constructed energies in Fig. 8(a) represents the quadrupole vibrational excitations in the $N = 83$ system, which are lowest for $Z \leqslant 56$. Here, these energies cross with those of the octupole vibrational excitations (represented by the 3^- curve), but the constructed values increase at higher Z. Consequently, excitations involving the 2^+ and 4^+ states of the $N = 82$ nucleus can play a significant role, but mostly for $Z \leqslant 56$. Hence, admixtures of types $2^+ \otimes 0i_{13/2}$ and $4^+ \otimes 0i_{13/2}$ are to be expected in addition to the $3^- \otimes 1f_{7/2}$ one in the wave functions of the $13/2_1^+$ and $13/2_2^+$ states with relative contributions varying with Z. Shell-model calculations reflect these observations. In the case of ^{137}Xe, for example, the calculated wave function of the $13/2_2^+$ level contains amplitude squared values of 34, 31, and 10%, respectively, for the $4^+ \otimes 0i_{13/2}$, $2^+ \otimes 0i_{13/2}$, and $3^- \otimes 1f_{7/2}$ configurations, as compared to a value of 5% for the $0^+ \otimes 0i_{13/2}$ component. The wave function of the $13/2_1^+$ state, on the other hand, is dominated by the $i_{13/2}$ orbital, with the $0^+ \otimes 0i_{13/2}$ component representing an amplitude squared of 77%, whereas the other significant components, $2^+ \otimes 0i_{13/2}$ and $4^+ \otimes 0i_{13/2}$, contribute together a value of 15%.

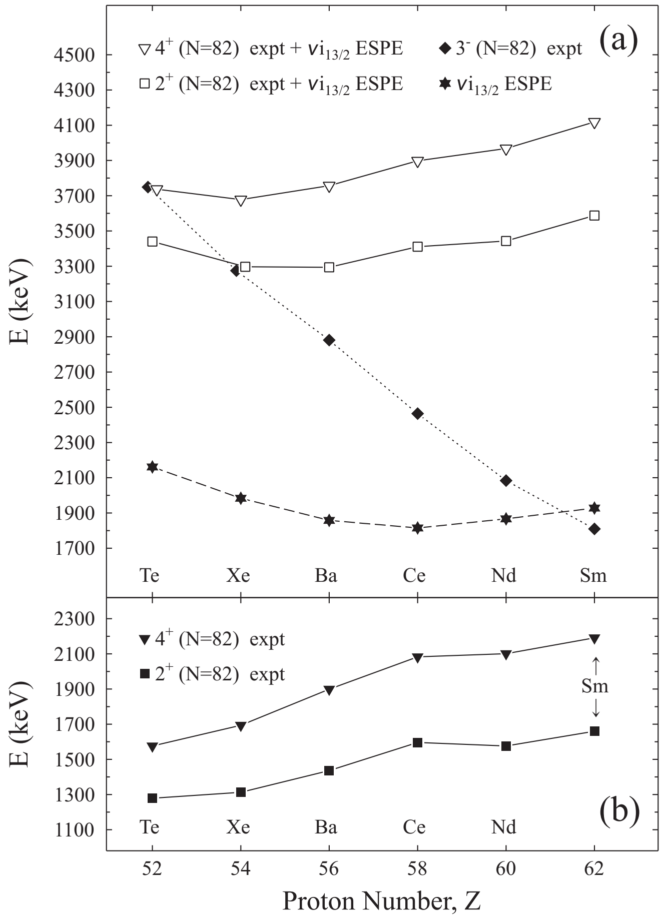

FIG. 8. A combined plot similar to Fig. 7. (a) The $\nu 0i_{13/2}$ effective single-particle energies, obtained from the calc2 calculations, are shown together with the 3^- energy levels in the corresponding $N = 82$ isotopes and a quantity reflecting excitations involving the coupling of an $i_{13/2}$ neutron with the yrast 2^+ and 4^+ states in the same even-mass nuclei (see text). (b) Shown are the yrast 2^+ and 4^+ energy levels as considered in the top panel. The experimental data are reported in Refs. [4,14].

Finally, the new information on negative-parity levels in ^{137}Xe is discussed. Here, the focus is on the higher-lying $3/2^-$ and $5/2^-$ levels, but other states are considered as well. The calculated excitation energies are compared with those obtained from the experiment in Table IV. Here, the calc2 values are reported, but there are no significant differences with calc1 results, except for the $13/2^+$ states discussed above. Irrespective of certain ambiguities in the experimental spin values, the observed and calculated $3/2^-$ and $5/2^-$ levels are reasonably close with discrepancies of 200 keV or less. The candidate $11/2_2^-$ level at 1879 keV can be associated with a 1760-keV state, increasing the confidence in the proposed assignment. This state has a $2^+ \otimes 1f_{7/2}$ "core excited" configuration as does the $11/2_1^-$ state.

TABLE IV. A comparison of calculated and experimental energies in ^{137}Xe. Note that the former ones are calc2 values (see text).

Calculation		Experiment	
I^π	E (keV)	I^π	E (keV)
$3/2_1^-$	729	$3/2_1^-$	601[a]
$3/2_2^-$	1709	(3/2)	1668[a]
$3/2_3^-$	1783	$3/2^-,5/2^-$	1716[a]
$3/2_4^-$	2063	$(3/2^-)$	1842[a]
$3/2_5^-$	2186	$3/2^-,5/2^-$	1850[a]
$3/2_6^-$	2233	$3/2^-$	2196[a]
$3/2_7^-$	2287	$3/2^-$	2490
$3/2_8^-$	2343	3/2	2565
$5/2_1^-$	1345	$5/2_1^-$	1303[a,b]
$5/2_2^-$	1664	$5/2^-,7/2^-$	1534[b]
$5/2_3^-$	1795	$3/2^-,5/2^-$	1716[a]
$5/2_4^-$	1890	$3/2^-,5/2^-$	1850[a]
$5/2_5^-$	2039	$5/2^-$[c]	1937
$7/2_2^-$	1589	$5/2^-,7/2^-$	1534[b]
$9/2_1^-$	1324	$9/2_1^-$	1218[a,b]
$9/2_2^-$	1584	$(9/2_2^-)$	1590[b]
$11/2_1^-$	1452	$11/2_1^-$	1220[a]
$11/2_2^-$	1760	(11/2)[d]	1879[a]
$13/2_1^+$	1786[e]	$13/2_1^+$	1753[a,f]
$13/2_2^+$	3308[e]	$13/2_2^+$	3137

[a]Reference [14].
[b]Reference [3].
[c]New spin-parity assignment, cf. Table I.
[d]Tentative spin assignment based on intensity considerations.
[e]Cf. Fig. 7.
[f]Reference [4].

V. CONCLUSIONS

The $13/2_2^+$ level and a couple of new $3/2^-$ and $5/2^-$ levels in ^{137}Xe have been observed by using a ^{136}Xe beam and ^{13}C and ^{9}Be targets and performing particle-γ coincidence measurements with the Phoswich Wall and Digital Gammasphere detector arrays. The observation of the $13/2_2^+$ level adds important information to the otherwise detailed knowledge of the ^{137}Xe level scheme. The shell-model calculations performed in the course of this work focused on the systematics of the $13/2_1^+$ and $13/2_2^+$ levels in ^{137}Xe and neighboring $N = 83$ nuclei. Their primary outcome is to provide a realistic range for the $\nu 0i_{13/2}$ single-particle energies (2360 keV $\leqslant E_{\rm SPE}(\nu 0i_{13/2}) \leqslant$ 2600 keV). Specifically, the lower-limit value of this range is based on the present calculations, which are guided by the $13/2_1^+$ level energy. The reasonable agreement of the calculations with the experimental $13/2_2^+$ level energies, including that of the newly observed one in ^{137}Xe, and additional comparisons with related levels in ^{134,135}Sb support this range of values further. The calculations also support the view that couplings of core excitations with $i_{13/2}$ and $f_{7/2}$ neutrons contribute to the wave functions of these two states and that the associated amplitudes change as a function of Z.

Furthermore, it is demonstrated that the population patterns of the $j_>$ and $j_<$ single-particle states in ^{137}Xe differ significantly depending on whether a ^{13}C or a ^{9}Be target (differing in the valence-neutron j value) is used. Hence, the "two-target" approach may well be instrumental in identifying the dominant single-particle character of a specific excitation. The technique clearly has potential for nuclear structure investigations using direct reactions with low-intensity rare-isotope beams.

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