Atmospheric
Measurement
Techniques

Testing the efficacy of atmospheric boundary layer height detection algorithms using uncrewed aircraft system data from MOSAiC

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Abstract. During the Multidisciplinary drifting Observatory for the Study of Arctic Climate (MOSAiC) expedition, meteorological conditions over the lowest 1 km of the atmosphere were sampled with the DataHawk2 (DH2) fixed-wing uncrewed aircraft system (UAS). These in situ observations of the central Arctic atmosphere are some of the most extensive to date and provide unique insight into the atmospheric boundary layer (ABL) structure. The ABL is an important component of the Arctic climate, as it can be closely coupled to cloud properties, surface fluxes, and the atmospheric radiation budget. The high temporal resolution of the UAS observations allows us to manually identify the ABL height (Z_{ABL}) for 65 out of the total 89 flights conducted over the central Arctic Ocean between 23 March and 26 July 2020 by visually analyzing profiles of virtual potential temperature, humidity, and bulk Richardson number. Comparing this subjective Z_{ABL} with Z_{ABL} identified by various previously published automated objective methods allows us to determine which objective methods are most successful at accurately identifying Z_{ABL} in the central Arctic environment and how the success of the methods differs based on stability regime. The objective methods we use are the Liu–Liang, Heffter, virtual potential temperature gradient maximum, and bulk Richardson number methods. In the process of testing these objective methods on the DH2 data, numerical thresholds were adapted to work best for the UAS-based sampling. To determine if conclusions are robust across different measurement platforms, the subjective and objective Z_{ABL} deter-

mination processes were repeated using the radiosonde profile closest in time to each DH2 flight. For both the DH2 and radiosonde data, it is determined that the bulk Richardson number method is the most successful at identifying Z_{ABL}, while the Liu–Liang method is least successful. The results of this study are expected to be beneficial for upcoming observational and modeling efforts regarding the central Arctic ABL.

1 Introduction

The transfer of energy between the Earth's surface and the overlying atmosphere, particularly at high latitudes, remains an area of substantial uncertainty in our understanding of the global climate system (de Boer et al., 2012; Tjernström et al., 2012; Karlsson and Svensson, 2013). The consequences of this uncertainty are significant, with global climate model projections of present-day sea ice demonstrated to fall short of simulating the observed rate of change (Stroeve et al., 2007, 2012). The thermodynamic structure of the lower atmosphere plays a central role in regulating cloud life cycle and radiative transfer, and their influence on atmospheric energy transport (Tjernström et al., 2004; Karlsson and Svensson, 2013; Brooks et al., 2017). Significant insight can be gained by measurements collected over the central Arctic Ocean pack ice, focused on the structure of the lower atmosphere, its spatial and temporal variabil-

ity, the intensity of turbulent energy fluxes, and its connection to surface features. To provide such measurements, uncrewed aircraft were deployed in the lower atmosphere during legs 3 (March through May 2020) and 4 (June through August 2020) of MOSAiC (Multidisciplinary drifting Observatory for the Study of Arctic Climate; Shupe et al. 2020), a year-long expedition that took place from October 2019 to September 2020 in which the icebreaker RV *Polarstern* (Alfred-Wegener-Institut Helmholtz-Zentrum für Polar- und Meeresforschung, 2017) was frozen into the central Arctic Ocean sea ice pack and allowed to passively drift across the central Arctic for an entire year (Fig. 1). Additional information on measurements taken of the atmosphere and sea ice during MOSAiC can be found at Shupe et al. (2022) and Nicolaus et al. (2022), respectively.

One important indicator of the extent to which energy may be transferred between the Earth's surface and overlying atmosphere is the atmospheric boundary layer (ABL) height. The ABL is the turbulent lowest part of the atmosphere that is directly influenced by the Earth's surface (Stull, 1988; Marsik et al., 1995). In the central Arctic, the ABL is impacted by interactions between the atmosphere and underlying surface, including both sea ice and open-water portions, which can cause either buoyantly or mechanically produced turbulence. The generation of buoyant turbulence can occur through surface energy fluxes emitted from open water regions such as leads (Lüpkes et al., 2008), cold air advection, especially over thin ice (Vihma et al., 2005), or turbulent mixing below cloud base due to cloud top radiative cooling (Tjernström et al., 2004). Mechanical generation, which is the dominant driver of turbulence in the central Arctic (Brooks et al., 2017), can occur due to the interaction between the atmosphere and surface roughness features such as ridges and ice edges (Andreas et al., 2010) or oceanic waves (Jenkins et al., 2012) or due to the presence of a low-level jet (Brooks et al., 2017; Banta et al., 2003). Solar heating of the Earth's surface and the subsequent formation of buoyant thermals, which is a dominant forcing of the ABL in most parts of the planet (Marsik et al., 1995), plays only a minor role in the central Arctic due to the relatively reflective surfaces found there.

The Arctic ABL is usually either stable or near neutral, while a convective ABL is rarely observed (Brooks et al., 2017; Esau and Sorokina, 2010). A stable boundary layer forms when there is a deficit of radiation at the surface or when warmer air is advected over a cooler surface and can range from being nearly well mixed with moderate turbulence to nearly laminar (Stull, 1988). A neutral boundary layer occurs when air at the surface is neutrally buoyant (Sivaraman et al., 2013); this is primarily due to mechanically generated turbulence that mixes air between the surface and the above atmosphere (Brooks et al., 2017). A convective boundary layer forms when convective thermals create positive buoyancy (Liu and Liang, 2010) and an air parcel at the surface rises adiabatically until becoming neutrally buoyant; when this phenomenon occurs in the central Arctic, it

is likely due to the presence of open water such as leads or polynyas (Lüpkes et al., 2008). While the various forms that the Arctic ABL may take are complex, most of the time the Arctic ABL is capped by a temperature inversion (which may extend to the surface for a stable ABL) and the local maximum in the potential temperature gradient, marking the entrainment zone, which is a stable layer that makes the transition from the ABL to the free atmosphere (Stull, 1988). One important difference between the Arctic ABL and that in the mid-latitudes is that there is usually no residual layer above a stable Arctic ABL due to the lack of a diurnal cycle. Additionally, the Arctic ABL is typically much shallower than that at mid-latitudes (Esau and Sorokina, 2010). These discrepancies cause certain ABL height detection methods to fail when applied to Arctic data.

Knowing the height of the Arctic ABL is important for many applications. First, it is a metric that represents the altitude up to which the atmosphere is directly impacted by surface processes. This can then inform the extent to which the surface interacts with atmospheric features such as clouds (and their influence on radiative transfer in the lower atmosphere), low-level jets (LLJs), and temperature inversion layers, which all have important implications for Arctic warming (Serreze and Barry, 2011). For example, a shallow, stable ABL is more likely to be observed with clear skies above (Brooks et al., 2017), which promotes longwave cooling of the surface and decoupling from the above atmosphere. In this instance, a surface-based temperature inversion is likely to constrain warming to the surface, which contributes to Arctic amplification (Lesins et al., 2012). ABL height (hereafter Z_{ABL}) plays an important role in many other applications, including transfer of air pollutants and weather forecasting (Garratt, 1994), and the proper parameterization of the ABL in numerical weather prediction models. Since any determination of Z_{ABL} is simply an approximation, the most value can be gained if this approximation is as accurate as possible. The goal of the current work is to determine which methods, based on thermodynamic and kinematic uncrewed aircraft system (UAS) profile data, can best accomplish this.

The depth of the ABL has been previously defined using a variety of approaches that involve visualizing the profiles of different thermodynamic and kinematic variables, which are listed in Table 1, along with some examples of associated literature that references use of that variable. Each of these profiles typically exhibits a distinct change in vertical structure at the top of the ABL. Additional methods may exist, such as analyzing the vertical gradient of aerosol content, but are not listed since the current study focuses on Z_{ABL} determination using thermodynamic and kinematic processes.

Due to the different atmospheric dynamics involved in each of the above approaches, the definition of Z_{ABL} is often debatable amongst experts. Depending on one's purpose for knowing Z_{ABL}, different approaches may be most applicable. Of these methods, some of the most widely used ones and the ones applied in the current analysis of a central Arctic
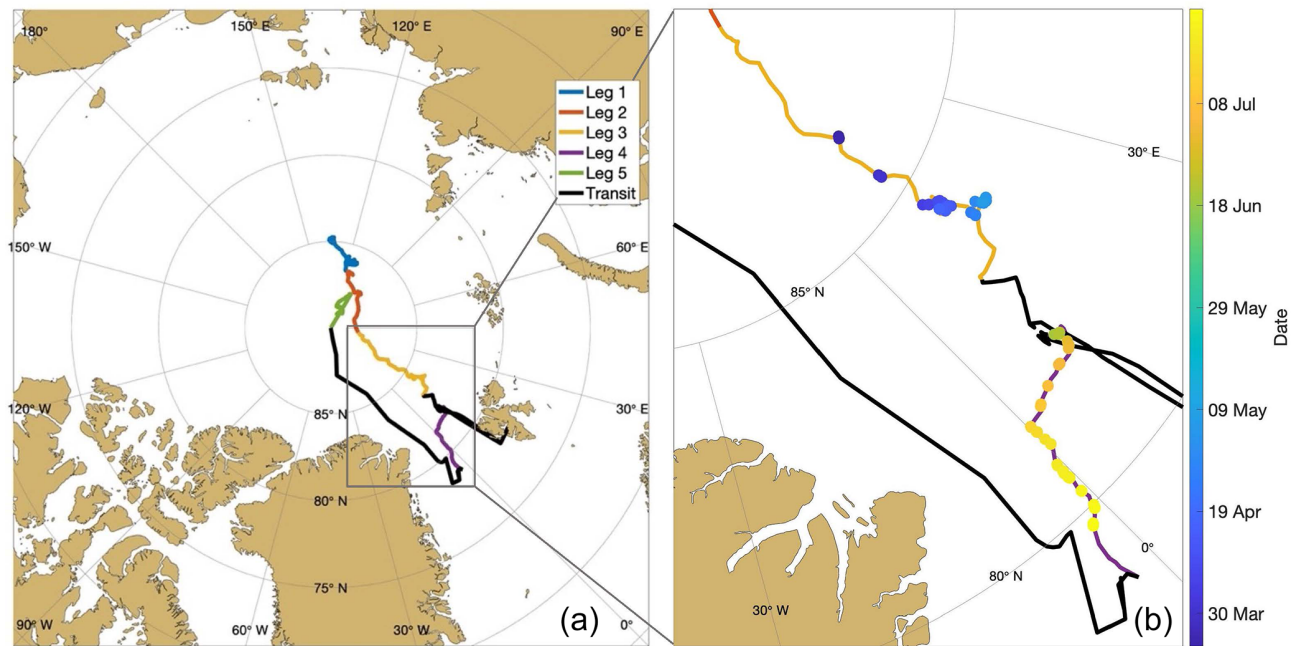
Figure 1. (a) The drift track of the *Polarstern*, separated by color into the five different legs. The black "transit" line indicates when the ship was traveling under its own power between legs 3 and 4 and between legs 4 and 5. **(b)** The zoomed-in portion of the *Polarstern* drift during which DH2 flights were conducted (legs 3 and 4). The locations of all of the DH2 flights are overlaid on the drift track and color coded by date: blue-tinted dots indicate flights conducted during leg 3, and yellow-tinted dots indicate flights conducted during leg 4.

dataset to determine Z_{ABL} are the ones that involve analysis of virtual potential temperature (θ_v), vertical gradient of virtual potential temperature ($d\theta_v/dz$), humidity (relative and absolute), bulk Richardson number (Ri_b), and wind speed profiles. The current focus is on these variables because the physical basis for each one as an indication of Z_{ABL} is relevant for the Arctic atmosphere. Specifically, θ_v helps identify the entrainment zone above the ABL, the vertical gradient of humidity either decreases or increases noticeably above the ABL (Dai et al., 2014), Ri_b helps identify where turbulence (usually caused by strong wind shear or surface roughness in the Arctic ABL; Grachev et al., 2005) ceases above the ABL, and wind speed helps identify the top of the ABL when it is capped by an LLJ, as the ABL top is often at or just below the LLJ core (Stull, 1988). Other methods, such as that using temperature inversion top to identify Z_{ABL} (Collaud Coen et al., 2014), do not perform well in the Arctic where a weak temperature inversion can extend well above the ABL. Though turbulent kinetic energy is recognized as perhaps the most valuable profile for Z_{ABL} identification (Stull, 1988; Seibert et al., 2000; Dai et al., 2014; Zhang et al., 2014), these data are not available to aid in the current study.

High-resolution data collected by the DataHawk2 UAS allow for determination of Z_{ABL} with high accuracy through manual visual analysis. However, visually determining Z_{ABL} case by case is time consuming for processing a large dataset. Therefore, the UAS-derived dataset is leveraged to com-

pare manually (or "subjectively") determined Z_{ABL} with that identified through previously published automated (or "objective") methods. While this subjective Z_{ABL} may not necessarily be the "true" ABL top, as the definition of this quantity can be debatable among experts and Z_{ABL} is not constant over time, it is the best estimate of Z_{ABL} given the available data. This evaluation is completed to identify objective methods that can accurately diagnose Z_{ABL} across a larger dataset of central Arctic atmospheric conditions.

To subjectively identify Z_{ABL} in each atmospheric profile from DH2 data, the stability regime of the ABL (stable, neutral, or convective) is categorized and Z_{ABL} is visually identified through combined evaluation of θ_v, humidity (both relative humidity, RH, and mixing ratio), and Ri_b profiles. Objective identification of Z_{ABL} is derived through the application of four previously published methods: the Liu–Liang method (Liu and Liang, 2010), the Heffter method (Heffter, 1980), the virtual potential temperature gradient maximum (TGRDM) method (Dai et al., 2014), and the Ri_b method (Sivaraman et al., 2013), all adapted to best suit the DH2 profiles examined (aside from the Heffter method, which was kept as standard). Following this, statistical comparisons between the objective and subjective Z_{ABL} are conducted. Next, the objective methods are applied in their adapted form to radiosonde profiles nearest in time to each DH2 flight to determine if these methods are robust across different measurement platforms for central Arctic conditions. Finally, dis-

Table 1. List of quantities previously used to identify Z_{ABL}, as well as some associated literature in which each variable is referenced.

Quantity used	Application of quantity	Previous literature
Virtual potential temperature (θ_v)	θ_v difference across θ_v inversion exceeds a threshold at the top of the ABL	Heffter (1980); Pesenson (2003); Sivaraman et al. (2013)
	θ_v at the top of an unstable ABL equals θ_v at the surface	Stull (1988); Liu and Liang (2010); Collaud Coen et al. (2014); Seibert et al. (2000)
Vertical gradient of virtual potential temperature ($d\theta_v/dz$)	Comparing $d\theta_v/dz$ to a threshold differentiates between ABL, entrainment zone, or free atmosphere	Heffter (1980); Stull (1988); Steeneveld et al. (2007); Liu and Liang (2010); Dai et al. (2011); Sivaraman et al. (2013); Dai et al. (2014); Zhang et al. (2014)
	Local maximum in $d\theta_v/dz$ at the top of the ABL	Dai et al. (2014)
Vertical gradient of temperature (dT/dz)	$dT/dz =$ zero at the top of a stable ABL	Stull (1988); Seibert et al. (2000); Dai et al. (2014); Collaud Coen et al. (2014)
	$dT/dz \leq$ the dry adiabatic lapse rate at the top of an unstable ABL	Collaud Coen et al. (2014)
Bulk Richardson number (Ri_b)	Ri_b exceeds critical value above the ABL	Stull (1988); Seibert et al. (2000); Zilitinkevich and Baklanov (2002); Steeneveld et al. (2007); Georgoulias et al. (2009); Dai et al. (2011, 2014); Sivaraman et al. (2013); Zhang et al. (2014); Collaud Coen et al. (2014)
Total wind speed	Low-level jet occurs at the top of a stable ABL	Stull (1988); Seibert et al. (2000); Steeneveld et al. (2007); Liu and Liang (2010); Sivaraman et al. (2013); Zhang et al. (2014)
Wind shear	Component-wise wind shear is greatly reduced above the ABL	Dai et al. (2011, 2014); Zhang et al. (2014)
Liquid water content and absolute humidity	Air moisture decreases drastically at the top of the ABL	Stull (1988); Seibert et al. (2000); Pesenson (2003); Dai et al. (2014)
Turbulent kinetic energy (TKE)	TKE ceases at the top of the ABL	Stull (1988); Seibert et al. (2000); Dai et al. (2014); Zhang et al. (2014)

cussion is included about the features that do or do not lend themselves to accurate identification of Z_{ABL} by the objective methods, and findings are summarized to support future studies seeking to identify Z_{ABL} quickly, objectively, and accurately across large atmospheric datasets collected in the central Arctic.

2 Data and methods

2.1 The DataHawk2

Data presented in this study were obtained between 23 March and 26 July 2020 using the University of Colorado DataHawk2 (DH2) UAS (de Boer et al., 2022). Flights were conducted from the sea ice alongside the *Polarstern*, known as the MOSAiC floe, ranging in location from 86.2° N, 15.8° E on 23 March, to 79.8° N, 1.9° W on 26 July 2020 (Fig. 1). Throughout this period, the MOSAiC floe evolved from snow-covered rigid ice situated in the high Arctic to being covered with melt ponds and leads close to the sea ice edge.

The surface atmospheric temperatures also transitioned from nearly $-35\,°C$ at the beginning of leg 3 to hovering near $0\,°C$ throughout the entirety of leg 4.

The DH2 (Hamilton et al., 2022) is a fixed-wing, battery-powered UAS (1.1 m wingspan, 1.8 kg weight, 40 min endurance) carrying various meteorological sensors, which measure the state of the atmosphere in Earth-relative coordinates. Instrumentation includes a fine wire array providing high-frequency (800 Hz) information on temperature and air speed, multiple sensors for temperature and relative humidity (Vaisala RSS421 measuring at 5 Hz and SHT-85 measuring at 100 Hz), and upward and downward looking thermopile sensors to provide infrared brightness temperatures of the sky and surface. Air pressure is measured at 5 Hz by the Vaisala RSS421 sensor. Altitude estimates are obtained using a global navigation satellite system (GNSS) receiver and barometer onboard. The altitude used in the current analysis is a high-resolution (800 Hz) barometric pressure altitude, which is corrected for drift using the GNSS altitude.

Measurements of attitude from the inertial measurement unit, airspeed from a Pitot static probe, and ground speed from the GPS receiver support the derivation of high-frequency (10 Hz) horizontal wind estimates. First, the "standard" approach, as laid out in van den Kroonenberg et al. (2008) and Hamilton et al. (2022), is applied, which derives wind estimates by combining GPS velocity measurements in the wind triangle using attitude estimates to rotate airframe-relative winds to Earth-relative coordinates. Additionally, a "hybrid" approach, as laid out in Lawrence and Balsley (2013) and Hamilton et al. (2022), is applied, which derives wind estimates by primarily using airspeed magnitude and GPS velocity, with secondary use of attitude estimates. For the purposes of this study, we use the DH2 winds derived from the "hybrid" approach. Please see Hamilton et al. (2022) and de Boer et al. (2022) for additional details on the wind derivation as specifically applied to the DH2. In addition, because take-offs and landings were flown manually by a remote pilot, the winds calculated during these times were found to be less reliable and accurate. As a result, we do not use DH2 winds calculated below 30 m altitude for this study. A brief description of the processing methods for the above variables are provided in the metadata for the DH2 dataset used for the current study (Jozef et al., 2021).

Combined, these sensors provide a comprehensive picture of atmospheric thermodynamic and kinematic state along with some context on the surface and sky condition under which these measurements were obtained. Table 2 lists the resolution, repeatability (standard deviation of difference between two successive repeated calibrations), and response time for the Vaisala RSS421 sensor. Uncertainty in the wind speed estimation is not provided, as determining this is still in progress.

Measurements collected by the DH2 are logged at different frequencies, requiring the implementation of a time alignment process to assure that the time index for each data point of each variable is consistent with all other measurements. Data collected by the DH2 during MOSAiC are available for public download through the National Science Foundation Arctic Data Center at https://doi.org/10.18739/A2KH0F08V (Jozef et al., 2021).

During MOSAiC, DH2 flights were conducted whenever flight weather criteria were met and when the team was able to access the ice alongside the *Polarstern*. The weather criteria include near-surface wind speeds with a sustained average below $10\,\mathrm{m\,s^{-1}}$ and gusts below $14\,\mathrm{m\,s^{-1}}$, as well as sufficient visibility to maintain visual contact with the aircraft at all times during flight. In addition, DH2 flights required coordination with other MOSAiC activities, especially those impacting air space over the MOSAiC floe, including manned helicopter flights and other UAS and tethersonde operations.

The most common flight pattern conducted with the DH2, and the flight pattern from which data for this analysis was acquired, was a profiling flight in which the plane flew a spiral ascent and descent pattern, with a radius of 75–100 m between the surface and 1 km altitude (or cloud base, if lower than 1 km), with the aircraft ascending and descending at a rate of $2\,\mathrm{m\,s^{-1}}$ and flying at an airspeed of 14–$18\,\mathrm{m\,s^{-1}}$. Each profiling flight lasted an average of 30 min, with some shorter flights when the air temperature was at its coldest ($\sim -35\,°\mathrm{C}$) near the beginning of leg 3 and some longer flights when the air temperature was much warmer ($\sim 0\,°\mathrm{C}$) during leg 4. Throughout the measurement period, 89 flights were conducted with the DH2. In the present study, 65 of these flights are found to have a clearly identifiable Z_{ABL} within the altitude range sampled. The remaining flights sampled only the lowest portion of the atmosphere due to cloud cover or other unfavorable environmental conditions and therefore did not observe the full depth of the ABL.

2.1.1 Preparing the DataHawk2 data for analysis

The primary profiles of interest for subjective and objective Z_{ABL} identification are θ_{v}, humidity (RH and mixing ratio), wind speed, $\mathrm{Ri_b}$, and $\mathrm{d}\theta_{\mathrm{v}}/\mathrm{d}z$. θ_{v} was calculated using RSS421 temperature, pressure, and RH. Differences in response times of the RSS421 temperature and RH sensors has a negligible impact on the calculation of θ_{v} because the moisture content in the Arctic atmosphere is so low that θ and θ_{v} values typically differ on the order of less than 1 K. Regardless, the addition of humidity does not change the structure and location of features in the θ_{v} profile, which is what is important for Z_{ABL} identification. To further eliminate the effects of differences in sensor response times during ascent and descent and for ease of visualization, we average the θ_{v}, humidity, and wind speed variables over 1 m altitude bins throughout the entire flight (e.g., values at 10.5 m are averaged from 10 to 11 m). This also mitigates the effect of changes in atmospheric conditions near the surface throughout the span of a flight, although the near-surface observations largely remained constant during a given flight. A value of 1 m is chosen as an averaging bin because using a greater bin value would eliminate much of the fine-scale detail in the θ_{v} and humidity profiles which the DH2 provides and which makes its data a valuable resource in honing Z_{ABL} detection methods. However, since fine-scale fluctuations in wind speeds evident at the 1 m scale are usually artifacts of the wind estimation routines applied to a circular flight pattern, we additionally apply a 60 m running mean, which eliminates small-scale wiggles while retaining the important large-scale features. Next, we exclude periods of manual flight during takeoff and landing (this is usually at altitudes below 5 m) since measurements during manual flight are prone to inaccuracies due to the irregular flight pattern. Lastly, we exclude the first 5 s of flight, as the initial measurements after take-off may be faulty due to hysteresis associated with the sensor sitting still at the surface before launch.

Using the 1 m averaged θ_{v} and wind speed component profiles, we calculate the $\mathrm{Ri_b}$ profile. $\mathrm{Ri_b}$ is calculated at altitude,

Table 2. Accuracy and reliability of the variables recorded by the Vaisala RSS421 sensors used in this study.

Variable	Resolution	Repeatability	Response time
Pressure	0.01 hPa	0.4 hPa	–
Temperature	0.01 °C	0.1 °C	0.5 s
Humidity	0.1 %RH	2 %RH	< 0.3 s (at 20 °C) to < 10 s (at −40 °C)

z, using the following equation from Stull (1988):

$$\mathrm{Ri_b}(z) = \frac{\left(\frac{g}{\overline{\theta_v}}\right)\Delta\theta_v \Delta z}{\Delta u^2 + \Delta v^2}, \tag{1}$$

where g is acceleration due to gravity, $\overline{\theta_v}$ is mean virtual potential temperature over the altitude range being considered, z is altitude, u is zonal wind, v is meridional wind, and Δ represents the difference over the altitude range used to calculate $\mathrm{Ri_b}$ throughout the profile. The only way that $\mathrm{Ri_b}$ can be negative is if the value for $\Delta\theta_v$ is negative, indicating a convective atmosphere with buoyancy-driven generation of the turbulence. $\mathrm{Ri_b}$ profiles are created by calculating $\mathrm{Ri_b}$ over a 30 m altitude range (Δz), at 5 m resolution (i.e., between 30 and 60 m, then between 35 and 65 m, and so on), rather than using the ground as the reference level, in order to isolate local likelihood of turbulence rather than that over the full depth from the surface (Stull, 1988; Georgoulias et al., 2009; Dai et al., 2014).

Since we do not use DH2 winds below 30 m, an intermediate $\mathrm{Ri_b}$ value between the surface and 30 m is calculated using an assumed zero wind at the surface. This results in $\mathrm{Ri_b}$ values at 15, 45, 50, 55 m, and so on. It is not crucial to consider the drift speed of the ice for the calculation of this initial $\mathrm{Ri_b}$ value since the ice drift speed during MOSAiC was on average less than 0.1 m s^{-1} (Krumpen et al., 2021), and the maximum drift speed during the DH2 flights was about 0.3 m s^{-1}, which is negligible compared to the speed of the observed winds. Nevertheless, any error in $\mathrm{Ri_b}$ that ensues due to the drift speed of the ice is limited to the first level where $\mathrm{Ri_b}$ is determined. Lastly, the $\mathrm{d}\theta_v/\mathrm{d}z$ profile is similarly created by calculating it over an altitude range of 30 m, at 5 m resolution.

The above profiles are used to determine stability regime, visually identify Z_{ABL} using criteria found in this paper, and objectively identify Z_{ABL} using the four published methods. For the remainder of this paper, Z_{ABL} determined from manual visual identification is referred to as the "subjective" Z_{ABL} and that determined by the published methods (which are automated algorithms performed by computers) are referred to as "objective" Z_{ABL}. These terms are used as a simplification to differentiate between manual and automated methods, though they both consider much of the same underlying physical processes that dictate ABL structure and height.

2.2 Determining the stability regime

Some of the methods for both subjectively and objectively identifying Z_{ABL} differ depending on the stability regime, and thus the sampled regime is first identified for each DH2 flight. The three possible stability regimes considered include a convective boundary layer (CBL), stable boundary layer (SBL), and neutral boundary layer (NBL; Liu and Liang, 2010). In a CBL, θ_v near the surface is greater than that of the overlying ABL (Stull, 1988). In an SBL, the vertical gradient of θ_v is positive (Stull, 1988). In an NBL, θ_v at the surface is approximately the same value as that of the overlying remainder of the ABL (Stull, 1988).

Therefore, stability regimes are identified by comparing θ_v between the lowest altitude sampled by the DH2 ("i" in the below equations; typically ~ 5 m since altitudes below this are usually sampled with manual flight) and 40 m above, using Eqs. (2)–(4) below, which are adapted from Liu and Liang (2010).

$$\theta_{v_{i+40\,m}} - \theta_{v_i} < -\delta_s = \mathrm{CBL} \tag{2}$$
$$\theta_{v_{i+40\,m}} - \theta_{v_i} > +\delta_s = \mathrm{SBL} \tag{3}$$
$$-\delta_s \leq \theta_{v_{i+40\,m}} - \theta_{v_i} \leq +\delta_s = \mathrm{NBL} \tag{4}$$

In these equations, δ_s is a stability threshold that represents the minimum positive or negative vertical difference of θ_v near the surface necessary for the ABL to qualify as an SBL or CBL, respectively. If this minimum is not either negatively (in the case of a CBL) or positively (in the case of an SBL) reached, the ABL is identified as an NBL (Liu and Liang, 2010). In an idealized case, δ_s would be zero. However, in practice it must be specified as a small positive number, and this number depends on the surface characteristics as well as inherent uncertainties or noise in the measurements. For profiles over ocean and ice, this threshold has been defined to be 0.2 K (Liu and Liang, 2010).

While Liu and Liang (2010) compare θ_v between pressure levels that equate to approximately 40 and 160 m in the conditions we sampled, this range was found to be inadequate for differentiating between an SBL, NBL or CBL in the Arctic, where the top of the ABL is often below 160 m, and sometimes even below 40 m. Therefore, considering the θ_v change below ~ 45 m more accurately reflects the stability regime of the Arctic ABL. Once the stability regime is identified, criteria based on the θ_v, humidity, and $\mathrm{Ri_b}$ profiles are applied to subjectively determine Z_{ABL}. For the current dataset, 31 SBL cases, 32 NBL cases, and 2 CBL cases were identified.

2.3 Subjective identification of atmospheric boundary layer height

There is no one best method for subjectively identifying Z_{ABL} that is agreed upon throughout the scientific community, evident by the many methods outlined in Table 1, and therefore a subjectively determined Z_{ABL} is prone to error. The best we can do to increase the confidence in a subjectively determined Z_{ABL} is to take into account several of the most commonly used methods and establish criteria which are applied consistently across all profiles. We describe these criteria below.

To subjectively identify Z_{ABL}, the θ_v profile is first analyzed, as the θ_v profile changes structure above the ABL (Stull, 1988). For a CBL and NBL, above the ABL, θ_v changes from decreasing or constant with height to increasing with height, marking the entrainment zone (Stull, 1988). The structure of an SBL, however, can vary a lot more (Mayer et al., 2012; Steeneveld et al., 2007; Zilitinkevich and Baklanov, 2002). In an ideal SBL case, the θ_v inversion is at its strongest (greatest vertical gradient of θ_v) near the surface and transitions to the free atmosphere (nearly constant or gradually increasing θ_v with altitude) above the SBL, with no entrainment zone (Stull, 1988). Z_{ABL} is then identified as the altitude of the shift from the surface-based θ_v inversion to the free atmosphere (Stull, 1988). In reality, the structure of an SBL is often not that simple, and the height of an SBL can be difficult to identify based on θ_v alone (Stull, 1988; Zhang et al., 2014). SBLs in the DH2 dataset often include a weaker surface-based θ_v inversion capped by a layer of enhanced stability (stronger θ_v inversion), reminiscent of an entrainment zone, likely because of surface-drag-induced turbulence close to the surface. ABLs with this structure form as the near-surface atmosphere fluctuates between weakly stable and near-neutral (Brooks et al., 2017). In more difficult cases such as these, the top of the SBL can be better determined by supplementing the θ_v profile with the RH and mixing ratio profiles, which usually have an obvious transition at the top of the ABL (Dai et al., 2014). This transition can manifest as either a shift from zero or positive to negative vertical gradient of humidity or as a humidity inversion. Use of the humidity profiles can also increase the confidence in identification of CBL and NBL height.

In addition, the Ri_b profile can aid in Z_{ABL} identification (Zhang et al., 2014). Ri_b is an approximation of the ratio between buoyantly produced (from thermals) or suppressed (from static stability) turbulence and mechanically produced turbulence (from wind shear; Sivaraman et al., 2013). Therefore, Ri_b can help to identify the top of the ABL under the assumption that turbulence ceases above the ABL (Stull, 1988). In the limit of layer thickness becoming small, Ri_b can be compared to a critical value of ~ 0.25 (Stull, 1988), with Ri_b below the critical value indicating an atmosphere that is likely to become or remain turbulent and Ri_b above the critical value indicating that an already laminar layer will not become turbulent, as static stability is strong enough to suppress mechanically generated turbulence. However, Ri_b does not always assume a small layer thickness, and thus a critical value is not well defined for Ri_b. Thus, for Ri_b near the critical value, there is uncertainty in the likelihood of turbulence (American Meteorological Society, 2012). However, since we calculate the profile of Ri_b over layers with a consistent thickness of 30 m, we can assume that the threshold for the likelihood of turbulence should at least be consistent throughout the profile. Additionally, since 30 m is a somewhat shallow thickness, there is less uncertainty in the likelihood of turbulence for Ri_b near the critical value of 0.25 than if we calculated Ri_b over an ever-increasing distance as we progress upward from the surface when always using the ground as a reference level.

Different studies have found the appropriate critical Richardson number to range from as low as 0.15 to as high as 7.2 in coarse-resolution models (Dai et al., 2014), but across the board, lower Ri_b is expected in the ABL, and higher Ri_b is expected above the ABL (Seibert et al., 2000). This increase in Ri_b above the ABL is in large part due to the decrease in wind shear. By examining Ri_b profiles for the DH2 flights, this transition from low values (near zero) to high values (with an increase of a few digits above the lower altitude values) can aid in identifying the top of the ABL.

Table 3 outlines the subjective criteria applied to determine Z_{ABL} depending on stability regime, which are separated depending on how many kinks there are in the θ_v profile that might indicate the entrainment zone. The term "kink" refers to a dramatic shift in slope (i.e., drastic change in vertical gradient). The primary methods applied to determine Z_{ABL} are those in which there are either one or two θ_v kinks, where we rely most heavily on the θ_v profile, and secondarily on the humidity and Ri_b profiles. For SBL cases, the humidity profiles often provide more insight than the Ri_b profile in identifying Z_{ABL}. In only a few especially difficult cases, we relied primarily on the Ri_b profiles. Figure 2 shows some examples of the subjective method applied.

When applying the above criteria, Z_{ABL} for the majority of cases (about 85 %) was clearly identifiable (i.e., relevant θ_v and humidity kinks were at the same altitude). For the other cases, Z_{ABL} was more ambiguous (e.g., Figs. S9, S12, S28, S34, S40, S42, S44, S48, S54, and S62 in the Supplement), meaning there were multiple features that one could argue marked the ABL top (i.e., the θ_v and humidity kinks that could both be interpreted as Z_{ABL} were at different altitudes). In these instances, depending on which feature is chosen, Z_{ABL} could differ by on average about 10–30 m, but preferential treatment is given to the kink that also corresponds to an increase in Ri_b. Additionally, if kinks in the RH and mixing ratio profiles occur at different altitudes, preferential treatment is given to the kink that occurs at the same altitude as that in the θ_v and/or Ri_b profiles. Following this, we determine the uncertainty in the subjective Z_{ABL} to be less than 30 m. Uncertainty in the height of a kink in an indi-

Table 3. Subjective criteria for identifying Z_{ABL}, depending on stability regime.

	One θ_v kink	Multiple θ_v kinks	No clear θ_v kinks
Convective boundary layer (CBL)	Z_{ABL} is the altitude at which the vertical gradient of θ_v is positive and may be the bottom of a layer of enhanced stability (greater vertical gradient of θ_v above), corresponding to a kink in the relative and/or absolute humidity profiles and an increase in Ri_b. Example: Fig. 2a		
Neutral boundary layer (NBL)	Z_{ABL} is the altitude of the singular θ_v kink marking the bottom of the lowest θ_v inversion. Example: Fig. 2b	Z_{ABL} is the altitude of the θ_v kink near the bottom of the lowest θ_v inversion that corresponds to a kink in the humidity profiles and an increase in Ri_b. Example: Fig. 2c	Z_{ABL} is the altitude of a faint θ_v slope shift that is identified via a corresponding kink in the humidity profiles and increase in Ri_b. Example: Fig. 2d
Stable boundary layer (SBL)	Z_{ABL} is the altitude of the θ_v kink marking the bottom of a layer of enhanced stability (greater vertical gradient of θ_v), corresponding to a kink in the humidity profiles and sometimes an increase in Ri_b. Example: Fig. 2e		Z_{ABL} is the altitude of a faint θ_v slope shift that is identified via a corresponding kink in the humidity profiles and sometimes an increase in Ri_b. Example: Fig. 2f

vidual profile is only subject to the vertical averaging procedure and sensor response time, and thus it is on the order of only ~ 1 m.

2.4 Objective identification of atmospheric boundary layer height

The strength of the subjective method described above is the knowledge of the expert, which cannot be automated (outside of possibly a machine learning algorithm, which would be costly and may still not be fully reliable). However, such expert knowledge and the time necessary to individually assess profiles is not always available. Thus, an automated method may often be preferred. Four such methods for objectively determining Z_{ABL} are applied and evaluated. Each of these methods relies on profiles of either $d\theta_v/dz$ or Ri_b, some in combination with the θ_v and/or wind speed profiles. Because the $d\theta_v/dz$ and Ri_b profiles are calculated over an altitude range of 30 m with 5 m resolution, objective Z_{ABL} detection methods which ultimately rely on these profiles can be determined with a resolution of 5 m. If they ultimately rely on the θ_v or wind speed profiles, Z_{ABL} can be determined with 1 m resolution. Figure 3 shows the application of all objective methods for an SBL and NBL case. A CBL case is not shown, as there were only two CBLs identified in the DH2 profiles, and they are rare in the central Arctic.

2.4.1 Liu–Liang method

The application of the Liu–Liang method depends on whether the profile includes a CBL, SBL, or NBL, which is determined using Eqs. (2)–(4). To implement the Liu–Liang method for a CBL profile, we first find the lowest altitude at which θ_v exceeds its the lowest DH2 value by

0.1 K. Then, Z_{ABL} is identified at the next lowest altitude in which $d\theta_v/dz$ exceeds $0.05 \, \text{K} \, (100 \, \text{m})^{-1}$ (Liu and Liang, 2010). For an NBL, Z_{ABL} is identified as the altitude at which $d\theta_v/dz$ first exceeds $2.5 \, \text{K} \, (100 \, \text{m})^{-1}$, which is adapted from a threshold of $0.05 \, \text{K} \, (100 \, \text{m})^{-1}$ used in Liu and Liang (2010), as this threshold was found to be inappropriate for the current dataset (Z_{ABL} found with the original threshold was always far too low). The basis of this method is to identify the entrainment zone at the top of the ABL through an increased value of $d\theta_v/dz$. The need for a greater threshold for NBL height identification in the current study is likely because the vertical resolution of sounding data used in the development of the Liu–Liang method was ~ 40–50 m (Liu and Liang, 2010), which would result in a much smoother $d\theta_v/dz$ profile than what is possible with the DH2 data. However, it would not make sense to interpolate the DH2 profiles to a resolution of 40–50 m before applying the Liu–Liang method, as this would eliminate the ability the identify key features in the often shallow Arctic ABL.

For an SBL, the Liu–Liang method searches for a potential Z_{ABL} associated with either minimal turbulence due to the lack of buoyancy within the ABL or greater turbulence in the ABL due to the presence of wind shear (Liu and Liang, 2010), with both being scenarios that may dictate Z_{ABL} for an SBL (Stull, 1988). Thus, SBL height is defined as either the top of the bulk stable (θ_v inversion) layer starting from the ground or the height of the LLJ maximum if present, whichever is lower (Liu and Liang, 2010). The top of the bulk stable layer is identified where the surface-based θ_v inversion has consistently diminished, and LLJ presence is identified by searching for wind speeds reaching a maximum that is at least $2 \, \text{m} \, \text{s}^{-1}$ stronger than the local minima above and below (Stull, 1988; Liu and Liang, 2010). For greater de-
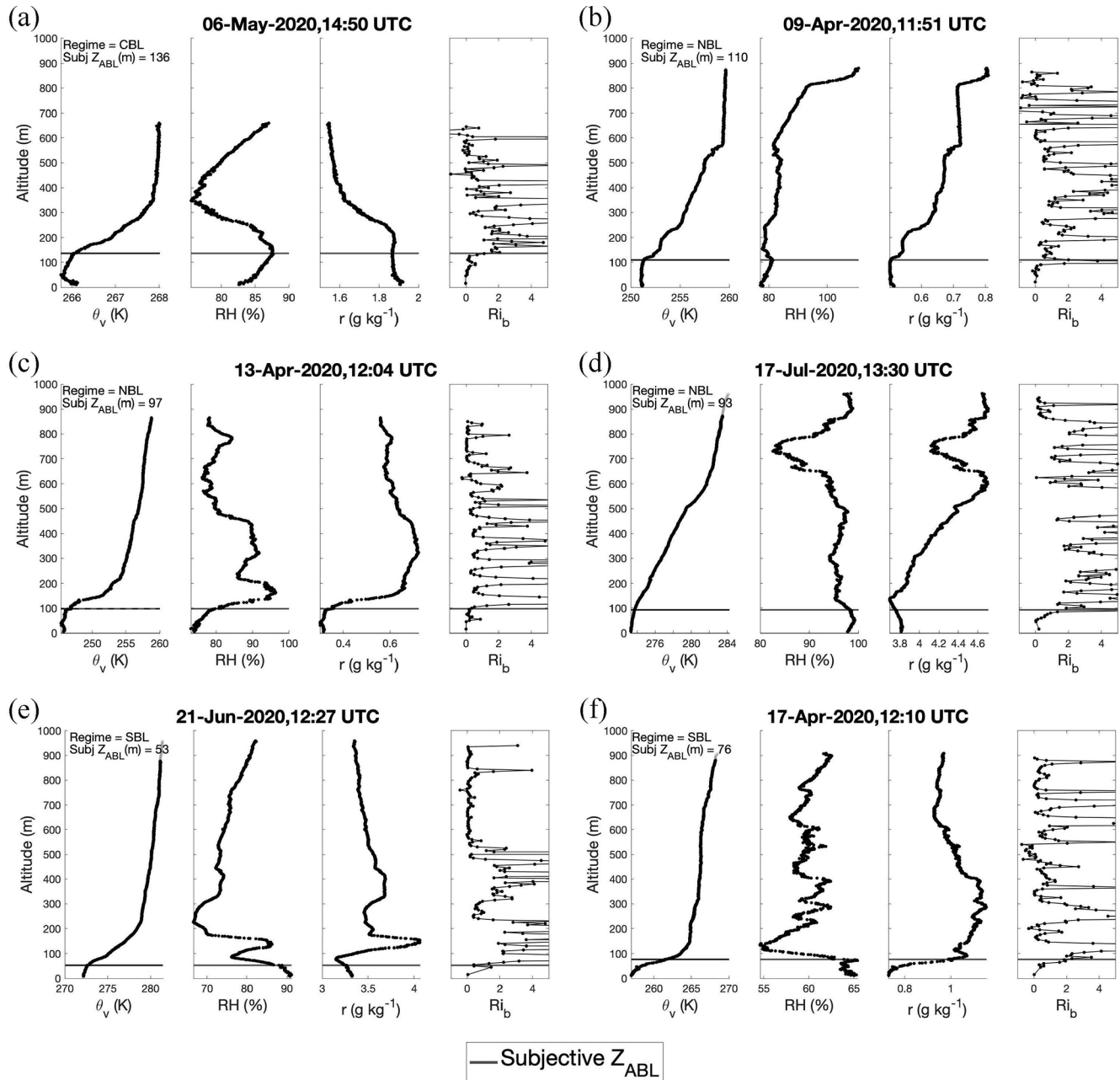
Figure 2. For each flight shown here, the θ_v profile is plotted in the left panel, the RH and mixing ratio profiles are plotted in the middle two panels, and the Ri_b profile is plotted in the right panel. Subjective Z_{ABL} is marked with a horizontal black line in each panel and is written, along with stability regime, on the left panel. (a) Example of a CBL case. (b–d) Examples of NBL cases. (e, f) Examples of SBL cases.

tail on these methods and the guiding equations, see Liu and Liang (2010). Figure S1 shows an example of the Liu–Liang method applied to a case for each stability regime.

2.4.2 Heffter method

The Heffter method uses θ_v difference across a θ_v inversion ($d\theta_v$) as an indication of Z_{ABL} (Sivaraman et al., 2013), by identifying the lowest θ_v inversion layer where $d\theta_v/dz$ is

greater than $0.5\,\mathrm{K}\,(100\,\mathrm{m})^{-1}$ throughout the θ_v inversion and $d\theta_v$ is at least $2\,\mathrm{K}$ (Heffter, 1980; Pesenson, 2003; Sivaraman et al., 2013). Within this θ_v inversion, the altitude at which θ_v first becomes more than $2\,\mathrm{K}$ greater than θ_v at the bottom of the θ_v inversion is labeled as Z_{ABL} (Marsik et al., 1995; Delle Monache et al., 2004; Snyder and Strawbridge, 2004; Sivaraman et al., 2013).

For a CBL or NBL, this method is meant to determine the altitude of the elevated θ_v inversion marking the entrainment
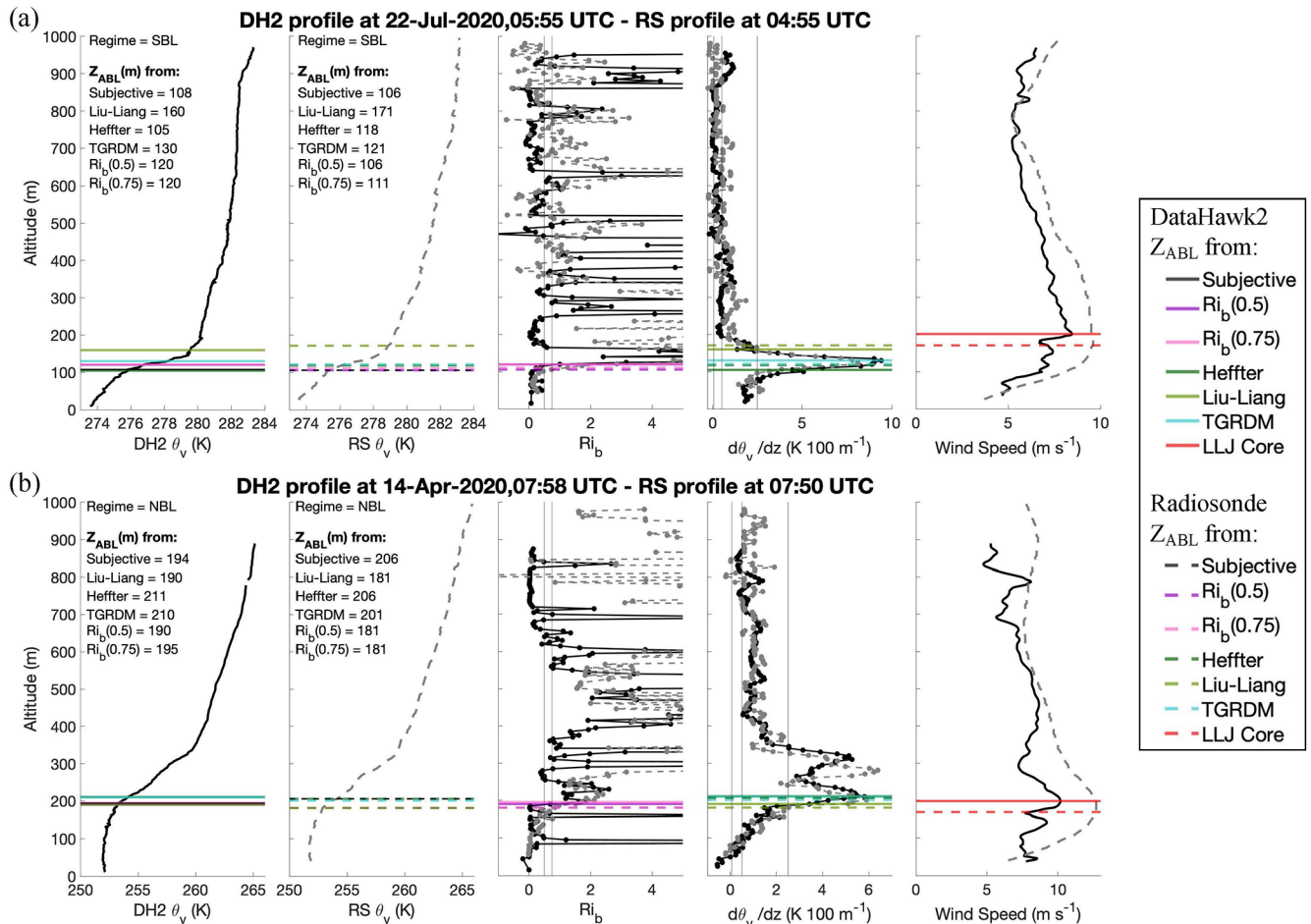
Figure 3. Demonstration of Z_{ABL} identification using all objective methods on both the DH2 (represented by solid lines) and corresponding radiosonde (represented by dashed lines) for an **(a)** SBL and **(b)** NBL case. From left to right the following profiles are shown: θ_v profile from the DH2, θ_v profile from the radiosonde, Ri_b profiles from the DH2 (solid black) and the radiosonde (dashed grey), $d\theta_v/dz$ profiles from the DH2 (solid black) and the radiosonde (dashed grey), and wind speed profiles from the DH2 (solid black) and radiosonde (dashed grey). The legend on the right indicates the Z_{ABL} detection method associated with each horizontal line in the figure. The LLJ core is not in itself a Z_{ABL} detection method but plays into the Liu–Liang method, and thus it is included. Each Z_{ABL} is written on the corresponding platform's θ_v profile.

zone between the well-mixed ABL and free atmosphere (Pesenson, 2003). For an SBL, this method determines where the change in strength of the surface θ_v inversion marks the transition from the ABL to residual layer (if one exists) or free atmosphere above (Stull, 1988). For greater detail on this method and the guiding equations, see Heffter (1980) or Sivaraman et al. (2013). Figure S2 in the Supplement shows an example of the Heffter method applied to a case for each stability regime.

2.4.3 Virtual potential temperature gradient maximum (TGRDM) method

The final $d\theta_v/dz$-based method used to find Z_{ABL} is the virtual potential temperature gradient maximum (TGRDM) method (Dai et al., 2014). Since the ABL is typically capped

by a well-defined θ_v inversion layer (Stull, 1988), even in a weakly stable case, we expect to see a local maximum in the $d\theta_v/dz$ profile at this point. By finding the maximum in the $d\theta_v/dz$ profile, the altitude at which the θ_v inversion is at its strongest and weakens above is identified. To apply this method, local maxima in the $d\theta_v/dz$ profile where $d\theta_v/dz$ is at least $1.75\,\mathrm{K}\,(100\,\mathrm{m})^{-1}$ greater than the local minimum $d\theta_v/dz$ above are identified. Z_{ABL} is set to the altitude of this lowest peak. Figure S3 in the Supplement shows an example of the TGRDM method applied to a case for each stability regime.

2.4.4 Bulk Richardson number method

Finally, a bulk Richardson number method for finding the ABL top is applied by determining the altitude at which Ri_b

exceeds a threshold value, which indicates where turbulence was likely no longer able to form in a laminar atmosphere. Previous literature suggests a wide range of critical values, with 0.25 (Stull, 1988) being the most widely accepted value, though a value of 0.5 is also often used (Sivaraman et al., 2013; Zhang et al., 2014). To determine a viable threshold value for the identifying Z_{ABL} in the DH2 data, a comparison between Z_{ABL} determined from a range of threshold values (we used 0.25, 0.5, 0.75, 1.0, 1.25, and 1.5) and the subjective Z_{ABL} was conducted. In identifying Z_{ABL} from these different threshold values, the level above which Ri_b was consistently greater than the threshold value was found. For this dataset, four consecutive data points (20 m) were required to be above the threshold value. We include this requirement due to the method of calculating Ri_b over a rolling 30 m range, rather than always with the ground as the reference layer, as it is possible for Ri_b to locally exceed the threshold but still be within the ABL. Thus, only when the Ri_b consistently exceeds the threshold, indicating that the bulk likelihood for turbulence has ceased, can we be confident that the top of the ABL has been reached.

Thus, the bottom of the lowest 20 m thick layer in which Ri_b exceeds the threshold value is identified as Z_{ABL}. The threshold values deemed to identify Z_{ABL} closest to that identified by the subjective method was 0.5 followed by 0.75. Therefore, further Z_{ABL} presented using the Ri_b method is calculated with threshold values of 0.5 (hereafter called $Ri_b(0.5)$) and 0.75 (hereafter called $Ri_b(0.75)$). Figure S4 in the Supplement shows an example of the Ri_b method applied to a case for each stability regime.

2.5 Applying the objective methods to radiosonde profiles

As discussed above, some of the objective methods used in this study were modified from their original descriptions to better work with the Arctic UAS data. Primarily, this includes changing the altitude range for determining stability regime, adjusting the threshold for calculating Liu–Liang NBL height, adding the $1.75\,\mathrm{K}\,(100\,\mathrm{m})^{-1}$ criterion to the TGRDM method, and choosing the best threshold values and specifying the necessary vertical distance for the Ri_b method. These adaptations are necessary in part because previous implementations involved analysis of radiosonde profiles, which have a lower vertical resolution than the DH2 profiles, and in mid-latitude locations, where the ABL structure is often quite different than that observed in the Arctic. Thus, profiles of θ_v, humidity, and wind speed from the balloon-borne radiosondes that were launched at least four times per day from the deck of the *Polarstern* (Maturilli et al., 2021) during MOSAiC are leveraged to determine if the objective methods used to identify Z_{ABL} from the UAS data are robust across platforms, despite differences in sampling methods.

To do this, radiosonde profiles with launch times closest to the DH2 flight times (within at most $\sim 3\,\mathrm{h}$) are used, repeating the same processes for subjective and objective Z_{ABL} identification and comparison. In eight instances, there were two DH2 flights in closest time proximity to the same radiosonde launch, and thus we use data from a total of 57 different radiosonde profiles. The specs for the Vaisala RS41-SGP sensor, which recorded the radiosonde variables, are the same as those listed in Table 2 for the DH2's RSS421 sensor, with the addition of pressure, temperature, and humidity uncertainty of 1.0 hPa, 0.3 °C, and 4 %, respectively, and a wind uncertainty and resolution of 0.15 and $0.1\,\mathrm{m\,s^{-1}}$, respectively, for velocity and of 2° and 0.1°, respectively, for direction. The radiosonde samples had a frequency of 1 Hz and an approximate climb rate of $5\,\mathrm{m\,s^{-1}}$, which results in data with a vertical resolution of $\sim 5\,\mathrm{m}$. Altitude measurements are calculated with the hydrostatic equation using the initial pressure at 10 m. Before proceeding with analysis, profiles of temperature, wind, and humidity from the radiosondes were visually compared to those from the corresponding DH2 flight to confirm that the measurements were similar to each other.

Prior to applying the objective methods, data below 23 m altitude were removed, as the lowest part of the radiosonde profiles were found to show inaccurately warm temperatures for several cases (Maturilli et al., 2021), due to the *Polarstern* acting as a "heat island." Additionally, in some cases the radiosonde data showed anomalously warm measurements some distance above 23 m, which is assumed to be the result of the balloon passing through the *Polarstern*'s exhaust plume. These measurements were adjusted by interpolating the temperature between the closest good measurements above and below where the radiosonde was presumably in the ship's plume. Applying these adjustments means that radiosonde data near the surface are not available for the determination of the stability regime. Therefore, we adapt the methods applied to the DH2 data in Eqs. (2)–(4) and instead calculate $d\theta_v$ between the lowest radiosonde measurement and 30 m above or the subjective Z_{ABL} if it is lower. We then compare this $d\theta_v$ to the appropriate threshold value, δ_s, that is equal to $(0.2\,\mathrm{K}/40\,\mathrm{m} = 0.005\,\mathrm{K\,m^{-1}})$ times the Δz used. For example, if the Δz of 30 m is used, the value of δ_s is 0.15 K. These adaptations in themselves do not result in the identification of a different stability regime than is found in the DH2 profiles; instead, differences in stability regime between the two platforms may result from the lack of near-surface observations from the radiosonde or a change in atmospheric structure between the two corresponding launches.

Figure 3 shows two examples (one SBL and one NBL) of all of the objective methods applied to both a DH2 flight and its corresponding radiosonde. These examples show that the subjective Z_{ABL} identified using the DH2 and radiosonde data are similar (differing by only 2 m for the SBL and 12 m for the NBL) and that the objective methods reveal a similar outcome when applied to the radiosonde data as they do

for the DH2 data for both cases. Similar figures for all DH2 and radiosonde profiles used in this study can be found in Figs. S5–S69 in the Supplement.

While the radiosonde and DH2 profiles generally exhibit a similar structure due to the close time and space proximity (the radiosondes were launched $< 100\,\mathrm{m}$ from the DH2 flights), the subjective Z_{ABL} identified in those profiles differ by 1–101 m. In general, the deviation between Z_{ABL} from the DH2 and the radiosonde increases with increasing time proximity. Figure 4 shows the absolute difference between DH2 and radiosonde subjective Z_{ABL} (Fig. 4a), as well as the absolute difference between the DH2 and radiosonde objective Z_{ABL} for each method (Fig. 4b) as a function of time difference in minutes between the DH2 and radiosonde launch. The best fit linear regression for each method shows that as time between the DH2 and radiosonde launch increases, the differences in Z_{ABL} increase as well (though minimally). However, the increase in absolute difference between subjective Z_{ABL} from the DH2 and radiosonde as time between the launches increases is not significant at the 5 % significance level (probability value of 0.74). Therefore, we are confident that Z_{ABL} does not significantly change for DH2 and radiosonde launches up to 3.16 h apart, which justifies the use of the radiosonde closest in time to each DH2 to test if there is similar efficacy of the different objective methods.

3 Results and discussion

3.1 Efficacy of objective Z_{ABL} identification methods

Whereas the objective methods all rely on information from one variable (or two, in the case of the Liu–Liang method for an SBL), the subjective method uses a combination of methods which can only be weighted properly by visual analysis. This is why the subjective method arguably results in a more accurate Z_{ABL} identification and provides a good basis for comparison with Z_{ABL} identified by the objective methods.

To determine how well the different objective methods worked, Z_{ABL} identified by each objective method is compared to the subjective Z_{ABL}. Figure 5 shows scatter plots comparing the objective to the subjective Z_{ABL} in each case, along with the associated best fit linear regression, coefficient of determination (R^2), slope, and probability value (p value) resulting from a paired two sample t test. For instances in which there were two DH2 flights in closest time proximity to the same radiosonde launch, the results from that radiosonde profile are plotted only once.

The R^2 value demonstrates how much of the variation in objective Z_{ABL} can be explained by the difference in subjective Z_{ABL}. Slope values (m) are also included to help evaluate the level of correspondence between the subjective and objective Z_{ABL} by comparison to an ideal value of $m = 1.00$. Additionally, looking at the intercept combined with the slope value tells us whether the objective method

tends to overestimate or underestimate Z_{ABL} compared to the subjective method. Lastly, the p value tells us whether the relationship between subjective and objective Z_{ABL} can be considered statistically significant at the 5 % significance level (a p value less than 0.05 indicates that there is a 95 % chance the relationship is due to true correlation).

Based on the DH2 data in these scatter plots, the method that gives the greatest R^2 is the $\mathrm{Ri_b}(0.5)$ method ($R^2 = 0.653$, Fig. 5d), followed by the $\mathrm{Ri_b}(0.75)$ method ($R^2 = 0.537$, Fig. 5e). These are followed closely by the Heffter method ($R^2 = 0.485$, Fig. 5b). The TGRDM method has the fourth-highest R^2 ($R^2 = 0.316$, Fig. 5c). The only objective method with a very low R^2 is the Liu–Liang method ($R^2 = 0.0907$, Fig. 5a). The slope values for all methods fall within $m = 1.00 \pm 0.30$, the closest to 1.00 being the $\mathrm{Ri_b}(0.75)$ method ($m = 1.02$), followed by the TGRDM method ($m = 1.10$) and Heffter method ($m = 1.18$). These slope values greater than 1.00 and positive intercept indicate that these methods generally overestimate Z_{ABL} when applied to the DH2 data, compared to the subjective Z_{ABL}. The results of the $\mathrm{Ri_b}(0.5)$ method and the Liu–Liang method, however, are more complex, as the slope values are both less than 1.00 ($m = 0.721$ and 0.708, respectively), but the intercepts are both positive. This indicates that these methods overestimate Z_{ABL} for a shallow ABL, but underestimate it for a deep ABL when applied to the DH2 data. Comparing the p values for all relationships to the 5 % significance level, the relationship between subjective and objective Z_{ABL} can be considered significant for every method (p value is less than 0.05). These p values follow the same order as the R^2 values, with the lowest p value found for the $\mathrm{Ri_b}(0.5)$ (indicating the highest significance) and the highest p value for the Liu–Liang method (indicating the lowest significance).

The radiosonde data (referred to as "RS" in Fig. 5) give a slightly different conclusion. Here, the method that gives the greatest R^2 is the Heffter method ($R^2 = 0.558$, Fig. 5b), followed by the $\mathrm{Ri_b}(0.5)$ method ($R^2 = 0.420$, Fig. 5d). The $\mathrm{Ri_b}(0.75)$ method and the TGRDM method have lower R^2 ($R^2 = 0.207$ and 0.225 in Fig. 5e and c, respectively). As was the case for the DH2 data, the only objective method with a very low R^2 is the Liu–Liang method ($R^2 = 0.00597$, Fig. 5a), which is also echoed by a slope value far from 1.00 ($m = 0.171$). The slope values for the rest of the methods are not as close to 1.00 as they are for the DH2 data, but they all fall within $m = 1.00 \pm 0.50$. The TGRDM has a slope value of $m = 1.00$, and the method with the next closest value to 1.00 is the Heffter method at $m = 1.13$. Both of these methods have a positive intercept, which indicates that these methods tend to overestimate Z_{ABL} when applied to the radiosonde data used in the current study. The rest of the methods have a slope of less than 1.00 and positive intercept, indicating that they tend to overestimate Z_{ABL} for a shallow ABL but underestimate it for a deep ABL when applied to the radiosonde data used in the current study. However, as R^2 for the Liu–Liang method is very low, this indicates that there is
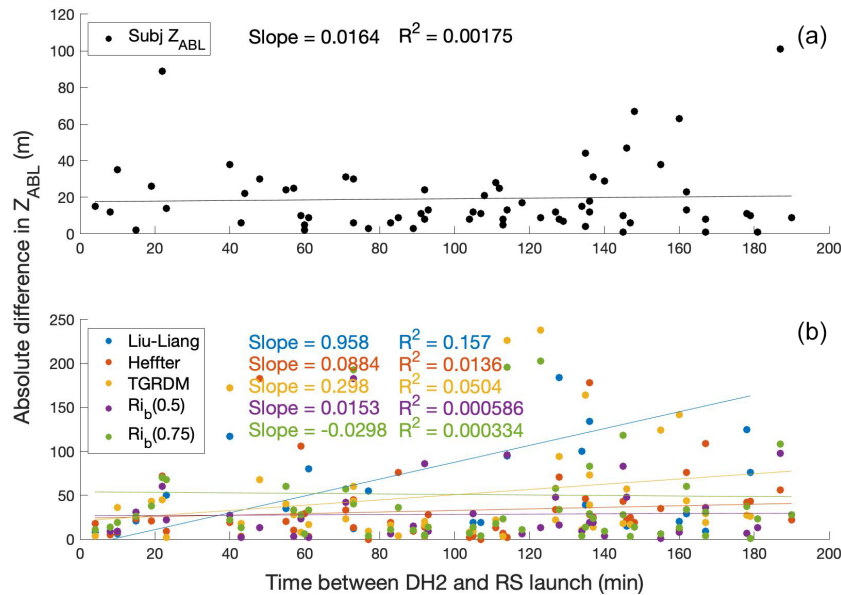
Figure 4. Absolute difference between subjective Z_{ABL} from the DH2 and subjective Z_{ABL} from the radiosonde closest in time to the DH2 launch (black dots, **a**) and absolute difference between objective Z_{ABL} from the DH2 and objective Z_{ABL} from the radiosonde closest in time to the DH2 launch (colored dots, **b**) versus absolute time difference in minutes between the DH2 and radiosonde launches. A few outlier points are not shown, as they lie outside the range of the y axis. Lines of best fit are included for the subjective Z_{ABL} and for each objective method, and the slope and R^2 value of each line is written next to the legend.

not much correlation between the objective and subjective Z_{ABL} for this method, and thus analysis of the slope does not provide reliable information. Lastly, the p values follow the same order as the R^2 values, with the lowest p value found for the Heffter method (indicating the highest significance) and the highest p value for the Liu–Liang method (indicating the lowest significance). Unlike the DH2 results, for the radiosonde the p values for all relationships compared to the 5 % significance level show that the relationship between subjective and objective Z_{ABL} can be considered significant for every method except the Liu–Liang method, in which the p value is greater than 0.05.

Lastly, Fig. 5f compares subjective Z_{ABL} from the radiosondes to subjective Z_{ABL} from the DH2. The high R^2 (0.752) indicates a rather strong correlation between subjective Z_{ABL} from both platforms, which demonstrates that Z_{ABL} usually did not change much between the DH2 and radiosonde launches in each case. Interestingly, there is enhanced deviation from the line of best fit for a shallower ABL and better agreement for a deeper ABL. However, this might simply be due to the greater number of samples with Z_{ABL} below ~ 200 m. The very low p value of 2.62×10^{-18} demonstrates the high significance in the relationship between Z_{ABL} from the DH2 and radiosondes.

Figure 6 shows the results presented in Fig. 5 but separated by stability regime, where the top half shows results for only SBLs, and the bottom half shows results for only NBLs. One primary takeaway from separating the results into stability regime is that for both platforms the TGRDM method

performs better for SBLs than it does for NBLs. Similarly, the Heffter method performs better for SBLs than NBLs for the DH2 data and performs similarly for the radiosonde data. This discrepancy is likely because these two methods search for a θ_v inversion to identify Z_{ABL}, which is often more defined for an SBL than NBL. Next, for the DH2 data, the Ri_b methods show less dependency on stability, with rather high R^2 for both regimes; however, the higher threshold performs better for NBL cases. Additionally, when splitting into stability regimes, the discrepancy between DH2 and radiosonde results increases for some methods. For example, the Ri_b method has more outliers for radiosonde NBL cases (Fig. 6i and j), causing R^2 to be rather low. For this category, the $Ri_b(0.5)$ method performs better, suggesting that the lower threshold value is more robust across platforms. Lastly, the Liu–Liang method, aside from a few outliers, performs rather well for NBL cases (Fig. 6f).

Additional analysis was completed to assess the cumulative frequency distribution for the difference in objective Z_{ABL} relative to the subjective Z_{ABL}. To do this, relative difference between the objective and subjective Z_{ABL} in each case and for each method was determined. These results are included in Fig. 7a for the DH2 profiles and in Fig. 7b for the radiosonde profiles. For example, about 26 % of the time, the Liu–Liang Z_{ABL} was within 10 % of the subjective Z_{ABL} for the DH2 data.

Figure 7a shows that for the DH2 profiles the $Ri_b(0.75)$ method results in the highest percent of cases within 10 % of the subjective Z_{ABL}, followed by the $Ri_b(0.5)$ method. In-
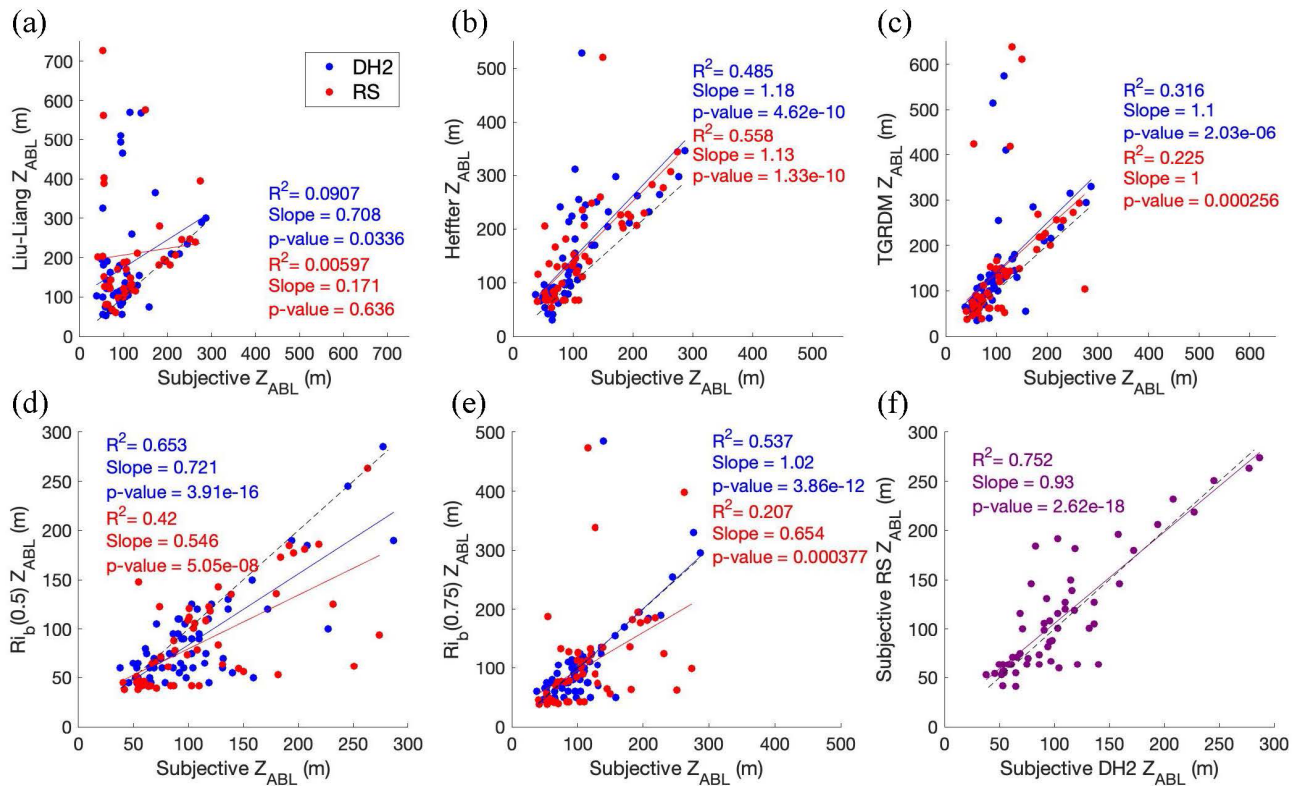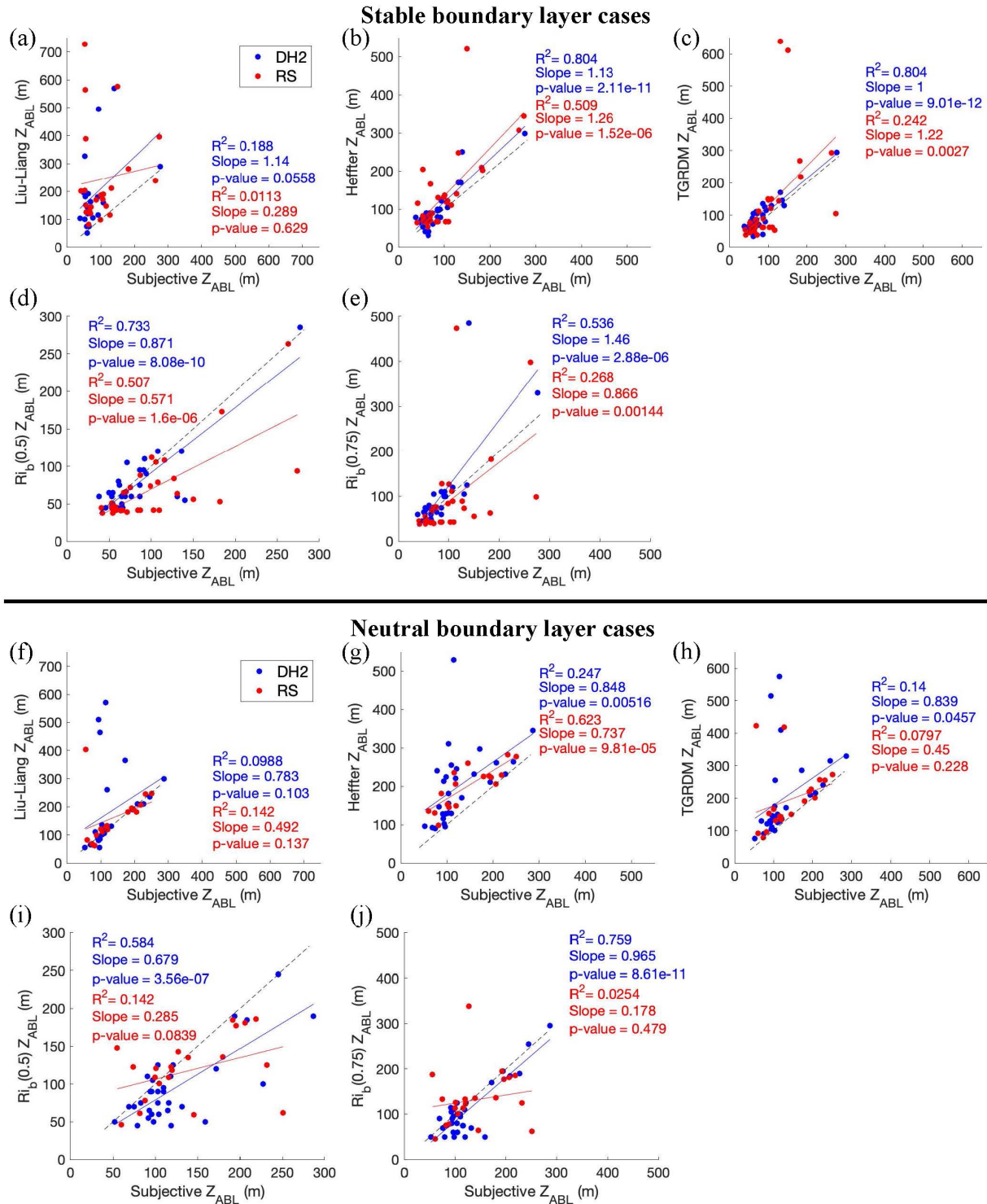
Figure 5. Relationships between subjective Z_{ABL} and objective Z_{ABL} from the **(a)** Liu–Liang method (50 DH2 samples and 40 RS samples), **(b)** Heffter method (61 DH2 samples and 53 RS samples), **(c)** TGRDM method (62 DH2 samples and 55 RS samples), and **(d, e)** Ri_b method (65 DH2 samples and 57 RS samples). Blue dots represent DH2 data and red dots represent radiosonde data. The solid blue line (solid red line) in each panel is the line of best fit for the DH2 (radiosonde) data. **(f)** Relationship between subjective Z_{ABL} from the radiosonde and subjective Z_{ABL} from the DH2 with line of best fit in purple (57 samples). Each panel is overlaid by the corresponding R^2, slope value, and p value. The dashed black line in each panel is a line with slope of 1.00 and y intercept of 0, for reference.

terestingly, the Liu–Liang method results in the third highest percent of cases within 10 % of the subjective Z_{ABL}. However, the Liu–Liang method falls behind other methods as the relative difference range is increased above 20 %. Additionally, the Liu–Liang method has the highest percent of cases in which no Z_{ABL} is found at all for the DH2 profiles, as well as about 20 % of cases that have greater than 100 % difference from the subjective Z_{ABL}. This trend indicates that while the Liu–Liang method sometimes works to find a Z_{ABL} close to the subjective Z_{ABL}, it also fails to find a Z_{ABL} close to the subjective Z_{ABL} or to find any Z_{ABL} in many cases. The primary reason for the failure of the Liu–Liang method, which is listed in Table 4 and discussed further in Sect. 3.2 below, is the high prevalence of a weak θ_v inversion that persists throughout the entire lower atmosphere in the Arctic. Another important finding is that the Ri_b method using either threshold value never fails to find a Z_{ABL}, and the number of cases within each relative difference range is greater for the Ri_b method than that for all other methods.

The information presented in the bar graph for the radiosonde profiles (Fig. 7b) leads to a similar conclusion. As for the DH2 profiles, the Ri_b method results in the high-

est percent of cases within 10 % of the subjective Z_{ABL} (but for this platform, the $Ri_b(0.5)$ method does best). Here, the Liu–Liang method results in the fourth highest percent of cases within 10 % of the subjective Z_{ABL} and performs more poorly as the relative difference range is increased. The Liu–Liang method also has the highest percent of cases in which no Z_{ABL} is found at all, followed by the Heffter and TGRDM methods, which was also true for the DH2 data. As for the DH2, there are no radiosonde cases in which the Ri_b method with either threshold value finds no Z_{ABL}. The main difference between Fig. 7b of the radiosonde data and Fig. 7a of the DH2 data is that while the $Ri_b(0.75)$ method applied to the DH2 data was always more successful than the $Ri_b(0.5)$ method for relative difference ranges below 70 %, for the radiosonde data the $Ri_b(0.5)$ method proves to always be more successful than the $Ri_b(0.75)$ method. We suspect that this results from the radiosonde data being more smoothed, which produces less sporadic Ri_b values as the atmosphere transitions from the ABL to the free atmosphere compared to the less smoothed DH2 data. This smoothing of the radiosonde data is applied by the Vaisala software to remove any effect of the chaotic pendulum swing directly af-

Figure 6. Relationships between subjective Z_{ABL} and objective Z_{ABL} for only stable cases (top) and only neutral cases (bottom) from the **(a, f)** Liu–Liang method (20 DH2 and 23 RS samples for SBL cases; 28 DH2 and 17 RS samples for NBL cases), **(b, g)** Heffter method (30 DH2 and 35 RS samples for SBL cases; 30 DH2 and 18 RS samples for NBL cases), **(c, h)** TGRDM method (31 DH2 and 35 RS samples for SBL cases; 29 DH2 and 20 RS samples for NBL cases), and **(d, e** and **i, j)** Ri_b method (31 DH2 and 35 RS samples for SBL cases; 32 DH2 and 22 RS samples for NBL cases). Blue dots represent DH2 data, and red dots represent radiosonde data. The solid blue line (solid red line) in each panel is the line of best fit for the DH2 (radiosonde) data. Each panel is overlaid by the corresponding R^2, slope value, and p value. The dashed black line in each panel is a line with slope of 1.00 and y intercept of 0, for reference.
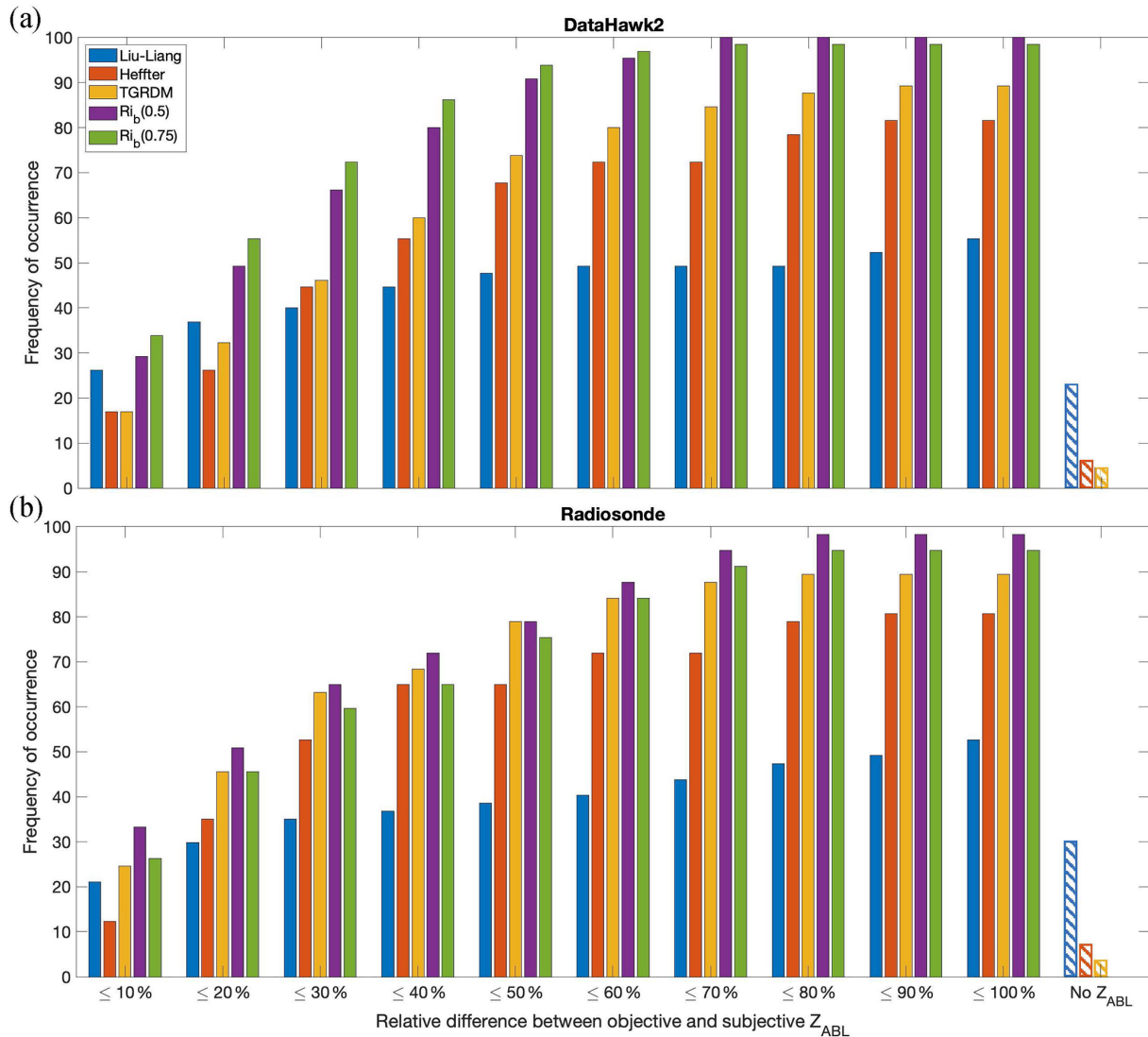
Figure 7. Bar plot showing what percent of **(a)** DH2 cases and **(b)** radiosonde cases give an objective Z_{ABL} within different relative difference ranges from the subjective Z_{ABL} using the different objective methods. Plot also shows the percent of cases for each method where no Z_{ABL} is found (labeled as "No Z_{ABL}").

ter launch while the wire unwinds. Thus, a lower threshold Ri_b value may be better applicable when more smoothing or filtering procedures are applied to a dataset.

Figures S70 and S71 in the Supplement show the results presented in Fig. 7 but separated by stability regime, where Fig. S70 shows results for only SBLs and Fig. S71 shows results for only NBLs. The primary takeaways from separating the results into stability regime is that for both the DH2 and radiosonde, the Ri_b method has the most cases and the Liu–Liang method has the least cases with objective Z_{ABL} within 10 % of the subjective Z_{ABL} for SBLs, though the Heffter and TGRDM methods also do well. For NBLs, the Liu–Liang method actually has the most cases with objective Z_{ABL} within 10 % of the subjective Z_{ABL}, followed by the Ri_b method for both platforms.

After comparing Z_{ABL} from the different objective methods to the subjective Z_{ABL} for both the DH2 and the radiosondes (Figs. 5 and 7), it is found that, with the exception of the Liu–Liang method, all other methods generally provide a reasonable estimate of Z_{ABL} for both datasets, with the Ri_b method being most favorable. This is in agreement with Seibert et al. (2000), Dai et al. (2014), and Zhang et al. (2014), who found an Ri_b-based method to be preferred when mechanically produced turbulence dominates, as is true in the central Arctic (Brooks et al., 2017). Additionally, the efficacy of each method is similar for the DH2 and the radiosonde data, as is indicated by similar patterns in the scatter plots (Fig. 5) and bar plots (Fig. 7), despite occasional differences in radiosonde versus DH2-based Z_{ABL} estimates, which likely result from the differences in sam-

Table 4. Summary of the features that lead to failure by each objective method, along with examples of DH2 cases that exemplify each failure, which can be found in the figures in the Supplement. The last column indicates the figures in the Supplement associated with cases in which the objective Z_{ABL} had a greater than 50 % difference from the subjective Z_{ABL} or no objective Z_{ABL} was found.

Objective method	Features which lead to failure	Resulting failure	Examples	Cases with > 50 % difference in Z_{ABL}
Liu–Liang	1. a weak θ_v inversion persists throughout the whole profile a. LLJ core altitude is well above the ABL top b. no LLJ 2. NBL capped by weak θ_v inversion	1a. overestimation of Z_{ABL} 1b. no Z_{ABL} found 2. overestimation of Z_{ABL}	1a. S6 on 24 March at 12:09 UTC 1b. S33 on 30 April at 14:07 UTC 2. S54 on 17 July at 13:30 UTC	S6, S9, S10, S11, S13, S14, S17, S18, S19, S24, S29, S30, S31, S32, S33, S34, S35, S39, S41, S46, S48, S49, S52, S54, S55, S57, S58, S59, S60, S62, S64, S65, S66, S68
Heffter	1. SBL height is not the altitude at which θ_v is 2 K warmer than θ_v at the surface a. SBL extends higher b. SBL does not extend as high 2. NBL capped by weak θ_v inversion 3. only shallow, weak θ_v inversion(s)	1a. underestimation of Z_{ABL} 1b. overestimation of Z_{ABL} 2. overestimation of Z_{ABL} 3. no Z_{ABL} found	1a. S5 on 23 March at 13:52 UTC 1b. S42 on 21 June at 13:13 UTC 2. S52 on 18 July at 13:10 UTC 3. S40 on 6 May at 14:50 UTC	S4, S15, S16, S17, S25, S29, S32, S33, S34, S40, S41, S45, S47, S51, S52, S54, S55, S56, S58, S59, S66
TGRDM	1. θ_v inversion is strongest at the surface 2. θ_v inversion is strongest within the entrainment zone 3. only shallow, weak θ_v inversion(s)	1. underestimation of Z_{ABL} 2. overestimation of Z_{ABL} 3. no Z_{ABL} found	1. S10 on 7 April (radiosonde profile) 2. S64 on 22 July at 07:37 UTC 3. S57 on 20 July at 11:28 UTC	S12, S13, S14, S24, S25, S29, S32, S45, S46, S52, S54, S57, S58, S59, S60, S64, S66
Ri_b	1. Ri_b is not capturing transition from turbulent to laminar atmosphere 2. threshold value is not accurate	1 and 2. overestimation or underestimation of Z_{ABL}	1 and 2. S8 on 29 March at 12:24 UTC and S45 on 30 June at 08:39 UTC	$Ri_b(0.5)$: S8, S17, S18, S52, S57, S66 $Ri_b(0.75)$: S17, S52, S57, S66

pling methods between the two platforms. Most specifically, the DH2 samples very close to the surface (~ 5 m) in most cases, and thus it observes important ABL features that support accurate stability and Z_{ABL} identification, whereas the radiosonde, which only samples down to 23 m at the lowest, may miss these features. Additionally, the DH2 samples with higher vertical resolution (due to higher time resolution of instrumentation and slower climb rate), again contributing to its ability to record complex fine-scale features that the radiosonde might miss. However, the similarity in efficacy of the objective methods between both platforms supports the fact that the objective Z_{ABL} identification methods that were adjusted using the high-resolution DH2 data are indeed robust across platforms with different sampling methods.

This is further explored by re-running the analysis with DH2 profiles averaged over 5, 10, and 20 m bins instead of 1 m bins to determine how sensitive the efficacy of the methods is to the vertical resolution of the data. When comparing objective Z_{ABL} found using the coarser data to the orig-

inal subjective Z_{ABL} for each method, the F test reveals that the R^2 values generally do not differ significantly from those found using 1 m binned data at the 5 % significance level. The only exceptions are the Liu–Liang method at all larger bin sizes and the Heffter method when using a 10 or 20 m bin size, which all manifest in lower R^2 value than those found using 1 m binned data. This reveals that the Liu–Liang method performs even more poorly at lower vertical resolution and that the Heffter method starts to perform more poorly at a vertical resolution of 10 m. On the other hand, the Ri_b and TGRDM methods remain just as successful when vertical resolution is reduced, and the preferred Ri_b threshold value does not appear to depend on vertical resolution. For vertical resolution of 30 m or coarser, the altitude range over which Ri_b is calculated would have to be increased, and at this point a lower threshold Ri_b value may be more applicable.

While we state an uncertainty in the subjective Z_{ABL} to be less than 30 m, this is only applicable to a handful of DH2

flights ($\sim 15\%$), whereas the majority have an uncertainty on the order of only ~ 1 m due to the vertical averaging procedure and sensor response time. Therefore, we do not expect this uncertainty to make any significant effect on the results.

3.2   When the objective methods fail

Table 4 lists the most common features that cause each objective method to fail (meaning the objective Z_{ABL} is much different than the subjective Z_{ABL}), along with the corresponding failure (either overestimation or underestimation or no Z_{ABL} found) and an example of such a situation shown in the figures in the Supplement. As shown in Sect. 3.1, while the Liu–Liang method sometimes works well, it is not reliable across a wide range of different profile structures. Option 1a causes failure because the $d\theta_v/dz$ criteria are not met anywhere in the profile, meaning that the method reverts to using the LLJ core height as Z_{ABL}. However, the LLJ core was observed to usually be above the subjective Z_{ABL} (supported by Stull, 1988; Jakobson et al., 2013; and Mahrt et al., 2014). This cause for failure agrees with Dai et al. (2014), which found that using LLJ core height to define SBL top produces results inconsistent with those from other methods. The Liu–Liang method likely performs better for NBL cases (as is evident in Figs. 6 and S71) than SBL cases because the Liu–Liang method for an NBL is not dependent on the sufficient diminishment of the θ_v inversion or the presence or altitude of a LLJ.

Any of the other objective methods would be a good choice for objectively determining Z_{ABL} for a dataset similar to the DH2 and radiosonde datasets (high-resolution profiles in the central Arctic environment). However, each method still struggles in some situations. The primary downfall of the Heffter method is that it identifies Z_{ABL} as the point where θ_v is 2 K warmer than θ_v at the bottom of the θ_v inversion. Failures noted in options 1–3 in Table 4 all occur when this criterion does not accurately identify the ABL top. The primary downfall of the TGRDM method, as noted in options 1–2 in Table 4, is that the strongest point of the θ_v inversion is not always at the ABL top. The TGRDM method also fails to find any Z_{ABL} if there is no θ_v inversion strong enough to exceed the threshold necessary for Z_{ABL} identification as laid out in Sect. 2.4.3. Lastly, the failure of the Ri_b method occurs due to the difficulty of defining an accurate threshold value that correctly captures the likelihood of turbulence for all cases.

The last column in Table 4 lists the cases in which the objective Z_{ABL} differs by more than 50% from the subjective Z_{ABL} for the DH2 data or no Z_{ABL} was found, which can be referenced in the figures in the Supplement for all examples of the profile structures that are not as conducive to the success of the different objective methods.

Aside from what is listed in Table 4, the objective methods may produce results different than those found by the subjective method due to the consideration of different variables. Primarily, none of the objective methods directly consider the profiles of RH or mixing ratio (of course, humidity is indirectly considered through the virtual potential temperature profiles). Additionally, the Liu–Liang method for a CBL or NBL, as well as the Heffter and TGRDM methods, do not consider wind shear in the identification of Z_{ABL}.

When applying these objective methods to a large dataset to automatically identify Z_{ABL}, it is recommended that some level of pre-screening is applied to flag cases that contain the features or structural patterns summarized in Table 4 that can objectively be identified and would make certain objective methods have difficulty identifying Z_{ABL} (for example, one can screen for whether the θ_v persists throughout the entire profile or where the θ_v maximum occurs) and choosing which objective method to use based on that. While not all features in Table 4 may be possible to pre-screen for, this list should at least help to identify some cases in which certain objective methods are likely to fail.

On the simplest level, one could choose which objective Z_{ABL} detection method to use based on stability regime. Given the results in Figs. 6, S70, and S71, the best choice to use for SBLs might be the Heffter method (highest R^2 and higher frequency of cases within 10% of the subjective Z_{ABL} when compared to NBL cases from both the DH2 and radiosonde data), and the best choice to use for NBLs might be the Ri_b method with either threshold value (highest R^2 values from the DH2 data and higher frequency of cases within 10% of the subjective Z_{ABL} when compared to SBL cases from both the DH2 and radiosonde data). However, when separating out the efficacy of the objective methods depending on stability regime, the Ri_b method has a combination of a high R^2 values and a high percentage of cases with objective Z_{ABL} within 10% of the subjective Z_{ABL} for both stability regimes, and thus this would be the best choice to apply to all profiles if one wanted to choose a single method, preferably with the threshold value of 0.5.

Overall, the objective methods are more likely to agree with each other as well as with the subjective Z_{ABL} for cases with more simplistic structures, such as those with strong θ_v inversions with a base at or just below the top of the ABL, those with LLJ core altitude at or just above the top of the ABL, and those with consistently and somewhat gradually increasing θ_v with altitude above the entrainment zone.

4   Summary and conclusions

By comparing subjective Z_{ABL} identified visually in θ_v, humidity (both RH and mixing ratio), and Ri_b profiles to objectively determined Z_{ABL}, the performance of several published methods (i.e., Liu–Liang, Heffter, TGRDM, and Ri_b) are evaluated across 65 DH2 UAS profiles. When comparing objective to subjective Z_{ABL} for each DH2 case, the method that is most successful (combination of high R^2 value, low p value, and slope close to 1.00) is the Ri_b method with either threshold value of 0.5 or 0.75 (Fig. 5). When calculat-

ing the percent of DH2 cases in which the objective Z_{ABL} is within certain relative difference ranges from the subjective Z_{ABL}, the Ri_b method is also most successful (Fig. 7). The Heffter and TGRDM methods also produce reasonable results according to Figs. 5 and 7. The only objective method that largely fails at accurately identifying Z_{ABL} is the Liu–Liang method.

In the process of applying these different objective methods to the DH2 data, some threshold values were modified to be better applicable to the UAS dataset. While these adjustments were made to best suit the 65 DH2 profiles analyzed in this study, which occurred between March and July of 2020, these adjustments should yield better results for identifying Z_{ABL} over sea ice during any season and location in the central Arctic. We hypothesize this because the ABL structures sampled by the DH2 in the current study were diverse and encompass the variety of ABL structures commonly observed in the central Arctic (which are typically shallow and either stable or neutral) throughout the entire year. Additionally, since the locations of the DH2 flights in this study range from deep in the Arctic pack ice to near the marginal ice zone, we are confident that the adjustments made will be applicable for identifying Z_{ABL} in either environment.

Testing these adjustments outside of the 65 DH2 flights, the modified techniques were also applied to the radiosonde profiles closest in time to each DH2 flight, to determine if the methods work similarly on data from another sensing platform with different sampling methods. Radiosonde profiles closest in time proximity to the DH2 flights were used under the assumption that the ABL structure would change minimally between the launch of the two platforms (supported by Fig. 4), and thus applying the methods of subjective and objective Z_{ABL} detection would lead to a similar conclusion. For the radiosonde data, the Heffter and Ri_b methods prove most successful in terms of having a high R^2 value, low p value, and slope closest to 1.00 when compared to the other objective methods (Fig. 5). Additionally, the Ri_b method also proves most successful when looking at the percent of cases in which the objective Z_{ABL} was within different relative difference ranges for the radiosondes, as it did for the DH2 (Fig. 7). Once again, the only method that consistently provided unfavorable results is the Liu–Liang method. These similar conclusions demonstrate that the adapted objective methods are indeed robust across platforms despite differences in sampling method, which suggest that one can take the methods and apply them to UAS, radiosonde, or other profile data alike without having to tweak them.

These findings show that no single method works well 100 % of the time. Given this, the best way to accurately identify Z_{ABL} across a variety of conditions in the Arctic atmosphere is to visually analyze the θ_v, humidity, and Ri_b profiles for each case individually. However, as subjective identification is time consuming and requires expert knowledge of the physical processes that dictate ABL structure, in the case of large datasets that require automated processing tech-

niques, the current study reveals that the Ri_b, Heffter, or TGRDM methods are most suitable for such a task, with the preferred method being the Ri_b method with threshold value of 0.5. For data with vertical resolution of 10 m or coarser, the Heffter method is no longer recommended. The Liu–Liang method does not provide consistent results in accurately identifying Arctic Z_{ABL} in many cases, especially for SBLs (Fig. S70). The most common occurrence of failure of the objective methods exists for NBLs capped by a weak θ_v inversion, and thus a clear θ_v slope change between the ABL and entrainment zone is difficult for automated methods to find. In such cases, the Ri_b method was found to be most reliable for identifying Z_{ABL}. A full list of features that cause each objective method to fail is provided in Table 4. The objective methods may also fail if the near-surface atmosphere is not well sampled, for example in the case of the radiosonde data; if ABL stability is defined by what is happening near the surface (e.g., a shallow convective layer), then this is missed by radiosonde profiles which only begin 23 m or higher, and stability regime could be incorrectly diagnosed. This highlights the value of platforms that can sample the near-surface atmosphere, such as the DH2. To accommodate the above problems, a semi-automatic approach may be beneficial in which one would apply all the recommended objective methods and visually inspect only the profiles for which the resulting Z_{ABL} diverges greatly.

The methods and results of this study for stability regime and Z_{ABL} identification are currently being applied to the entire year of radiosonde data collected during the MOSAiC expedition (October 2019 – September 2020) to create a data product containing year-long statistics on ABL characteristics in the central Arctic. Additional metrics, such as LLJ height and speed, and temperature inversion layer depth and strength will be included in this product for eventual publication. The value of the DH2 data and methods used in the current study comes from the uniqueness of the location and timing of the profiles collected. Therefore, these data provide a unique opportunity to evaluate any additional Z_{ABL} detection schemes that were not addressed in this study or that have yet to be developed and can be used to learn about the intricacies of additional structural components of the Arctic atmosphere such as the entrainment zone. Lastly, we are working to derive turbulence parameters from the DH2 fine-wire measurements, which will enhance the value of the DH2 data in ABL studies.

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