



HAL
open science

Second-order cone programming for rolling friction contact mechanics

Minh-Hoang Nguyen, Vincent Acary, Paul Armand

► **To cite this version:**

Minh-Hoang Nguyen, Vincent Acary, Paul Armand. Second-order cone programming for rolling friction contact mechanics. SMAI MODE 2022 - Journées du groupe MODE de la Société de Mathématiques Appliquées et Industrielles, May 2022, Limoges, France. pp.1-1. hal-03761228

HAL Id: hal-03761228

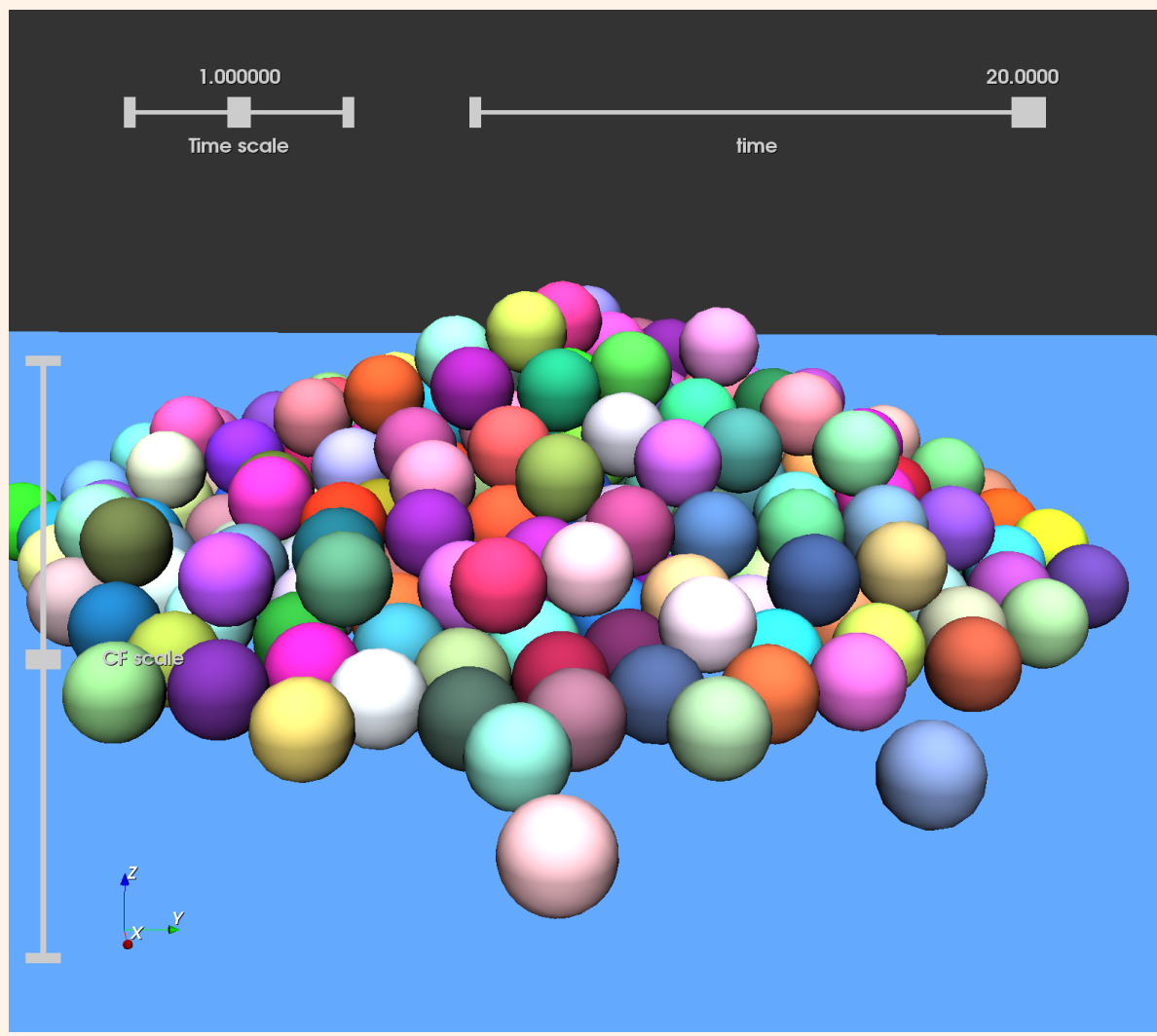
<https://hal.inria.fr/hal-03761228>

Submitted on 26 Aug 2022

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

GENERAL FRICTION CONTACT PROBLEM



$$\begin{aligned} Mv + f - H^T r &= 0 \\ Hv + w &= u \\ \mathcal{F}^* \ni u + \Phi(u) \perp r &\in \mathcal{F} \end{aligned}$$

where

- **Parameters:** n No. of contact points, m No. of degrees of freedom.
- **Given data:** $0 \prec M = M^T \in \mathbb{R}^{m \times m}$, $H \in \mathbb{R}^{5n \times m}$, $f \in \mathbb{R}^m$, $w \in \mathbb{R}^{5n}$.
- **Convex problem:** $\Phi(u)$ is fixed.
- **Variables:** global velocity $v \in \mathbb{R}^m$, velocity $u \in \mathbb{R}^{5n}$, reaction $r \in \mathbb{R}^{5n}$.

ORIGINAL ROLLING FRICTION CONE

$$\begin{aligned} \mathcal{F} &= \prod_{i=1}^n \mathcal{R}_i := \{r_i = (r_{i0}, \tilde{r}_i, \tilde{r}_i) \in \mathbb{R}^5 : \max\{\|\tilde{r}_i\|_2, \|\tilde{r}_i\|_2\} \leq r_{i0}\} \\ \mathcal{F}^* &= \prod_{i=1}^n \mathcal{R}_i^* := \{u_i = (u_{i0}, \tilde{u}_i, \tilde{u}_i) \in \mathbb{R}^5 : \|\tilde{u}_i\|_2 + \|\tilde{u}_i\|_2 \leq u_{i0}\} \end{aligned}$$

$$\begin{aligned} \min_{v,u} \frac{1}{2} v^T M v + f^T v & \quad \max_{v,r} -\frac{1}{2} v^T M v - w^T r \\ \text{(P)} \quad \text{s.t. } H v + w = u & \quad \text{(D)} \quad \text{s.t. } M v + f - H^T r = 0 \\ & \quad u \in \mathcal{F}^* & \quad r \in \mathcal{F} \end{aligned}$$

Since \mathcal{F} is not self-dual, the Jordan algebra cannot be used as it is.

NESTEROV & TODD (NT) SCALING

To get a symmetric and non-singular system, we need a change of variables at each iteration of the **Interior Point algorithm**. Define a vector p such that

$$\hat{z} := Q_p z = Q_{p-1} y =: \tilde{y}$$

where Q_p is quadratic representation of p .

$$\begin{aligned} \min_{v,\hat{z}} \frac{1}{2} v^T M v + f^T v & \quad \max_{v,\tilde{y}} -\frac{1}{2} v^T M v - w^T (K Q_p \tilde{y}) \\ \text{(P}_{NT}) \quad \text{s.t. } H v + w = J Q_{p-1} \hat{z} & \quad \text{(D}_{NT}) \quad \text{s.t. } M v + f - H^T K Q_p \tilde{y} = 0 \\ & \quad \hat{z} \in \mathcal{L}^{2n} & \quad \tilde{y} \in \mathcal{L}^{2n} \end{aligned}$$

REFORMULATION WITH PRODUCT OF LORENTZ CONES

To simplify, set $n = 1$. By letting $u_0 = t + t'$, we obtain

$$\begin{aligned} \min_{v,z} \frac{1}{2} v^T M v + f^T v & \quad \max_{v,y} -\frac{1}{2} v^T M v - w^T K y \\ \text{(P')} \quad \text{s.t. } H v + w = J z & \quad \text{(D')} \quad \text{s.t. } M v + f - H^T K y = 0 \\ & \quad z \in \mathcal{L}^2 & \quad y \in \mathcal{L}^2 \end{aligned}$$

where

$$J = \begin{pmatrix} 1 & 1 \\ I_2 & I_2 \end{pmatrix}, \quad K = \begin{pmatrix} 1 & 0 \\ I_2 & I_2 \end{pmatrix}, \quad z = \begin{pmatrix} t \\ \tilde{u} \\ t' \\ \tilde{u} \end{pmatrix} \quad \text{and} \quad y = \begin{pmatrix} r_0 \\ \tilde{r} \\ r_0 \\ \tilde{r} \end{pmatrix}.$$

Note that

- $Jz = u$, $Ky = r$ and $KJ^T = I$,
- \mathcal{L}^2 is self-dual, then Jordan algebra can be used.

SOLVER

The **perturbed KKT system** is

$$\begin{aligned} M v + f - H^T K Q_p \tilde{y} &= 0, \\ J Q_{p-1} \hat{z} - H v - w &= 0, \\ \hat{z} \circ \tilde{y} &= 2\mu e, \\ (\hat{z}, \tilde{y}) &\in \text{int}(\mathcal{L}^4). \end{aligned}$$

The **linear system** to solve at each iteration is of the form

$$\begin{bmatrix} M & -H^T & 0 \\ -H & 0 & J \\ 0 & J^T & Q_{p^2} \end{bmatrix} \begin{bmatrix} d^v \\ d^r \\ d^z \end{bmatrix} = - \begin{bmatrix} M v + f - H^T r \\ J z - H v - w \\ J^T r - 2\mu z^{-1} \end{bmatrix}.$$

One difficulty is that Q_{p^2} is badly conditioned near the optimal solution.

EQUAL STEPLENGTHS

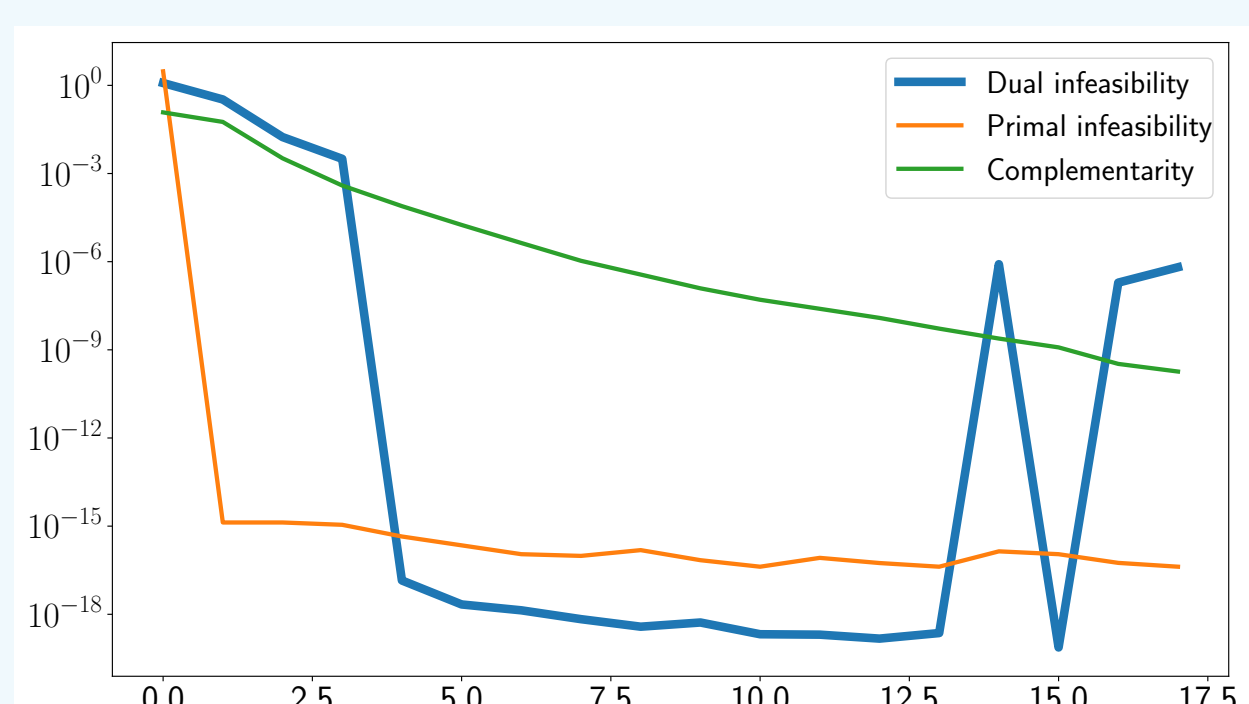


Figure 1. $n = 70$, steplengths $\alpha_p \neq \alpha_b$

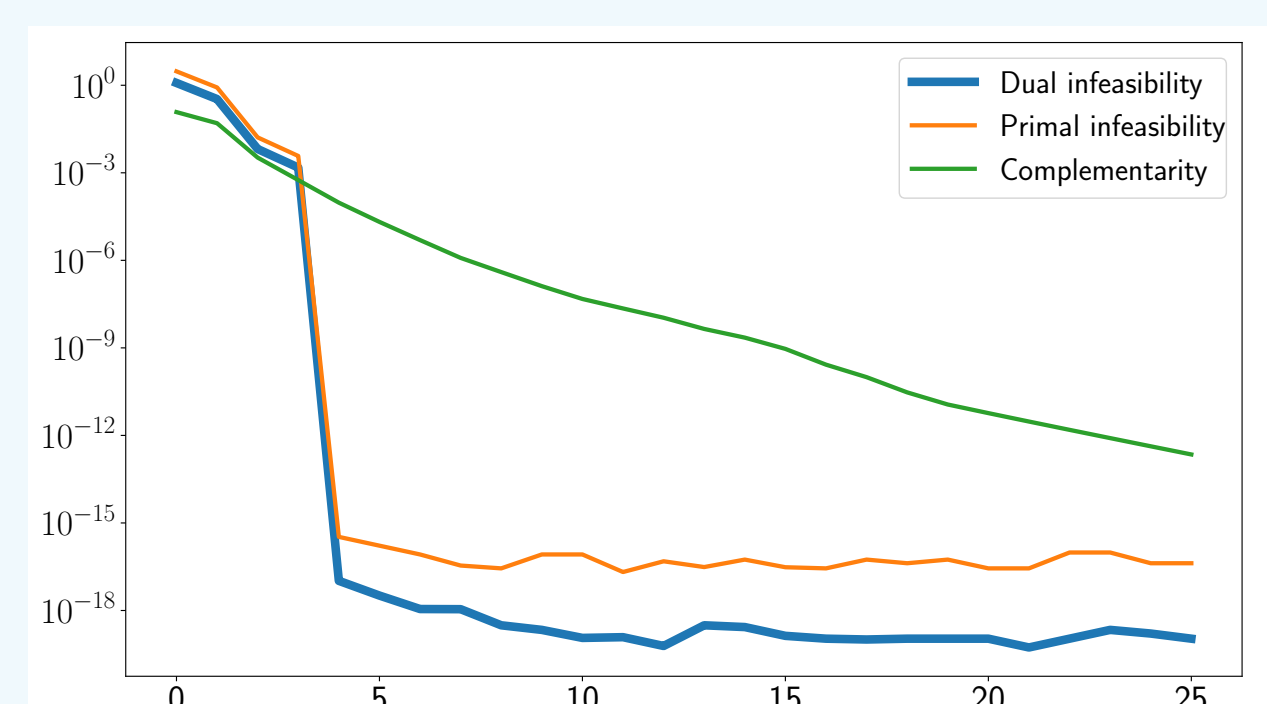
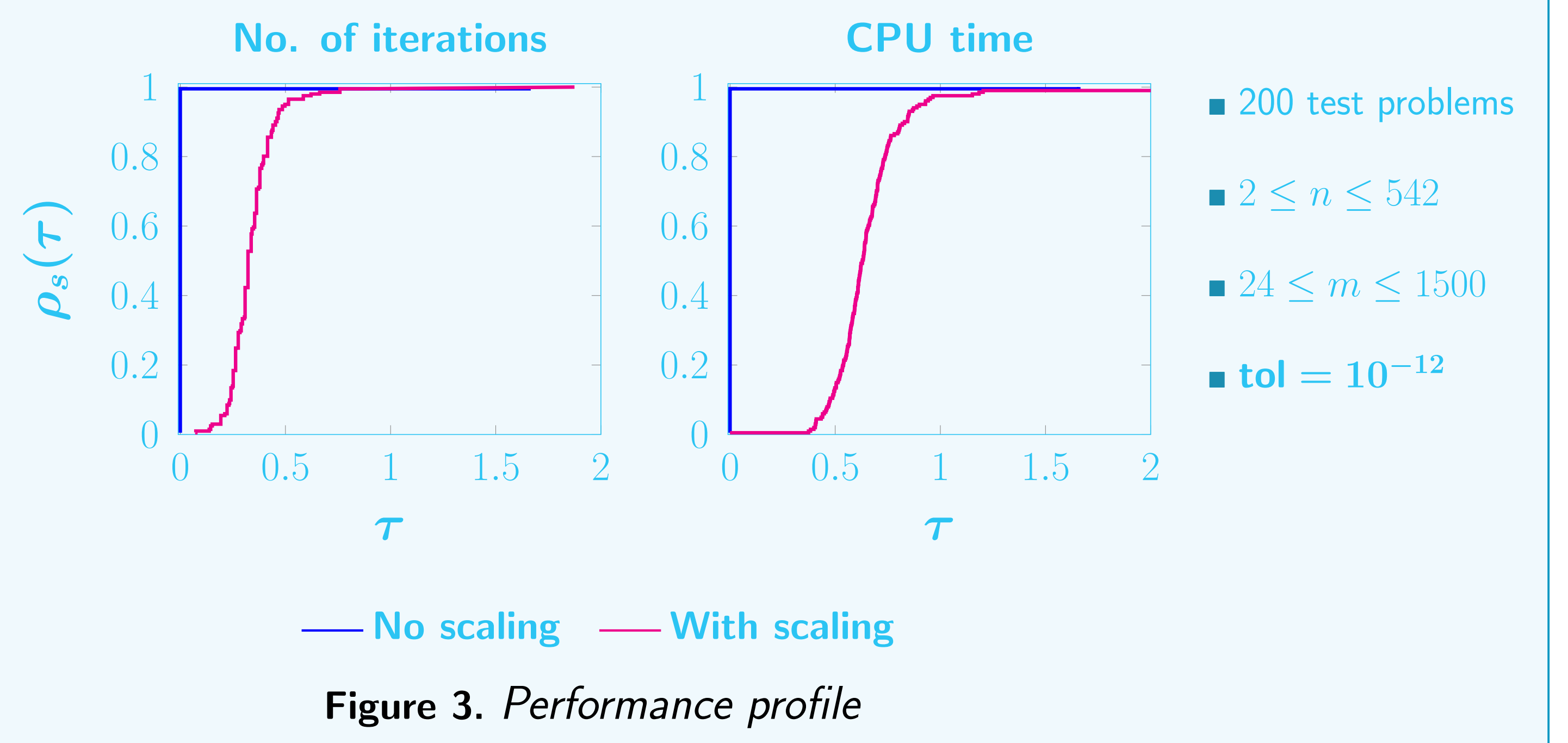


Figure 2. $n = 70$, steplengths $\alpha_p = \alpha_b$

Primal-dual infeasibilities

$$\begin{aligned} \text{pinfeas}^+ &= (1 - \alpha_p) \text{pinfeas} + \alpha_p \xi_p & \text{where} \\ \text{dinfeas}^+ &= (1 - \alpha_b) \text{dinfeas} + \alpha_b \xi_b + (\alpha_b - \alpha_p) M d^v & \xi_p, \xi_b \text{ are residual vectors.} \end{aligned}$$

EFFICIENCY & ROBUSTNESS



THEORETICAL RESULT FOR FRICTION CONTACT MECHANICS WITHOUT ROLLING FRICTION

- For the case without rolling friction, the original friction cone \mathcal{F} is the Lorentz cone \mathcal{L} . If (P) is feasible, then it has a unique optimal solution (v^*, u^*) . It is not the case for (D), because H is rank deficient.
- If the **Slater's conditions** and the **strict complementarity** hold, there are 3 possible cases:

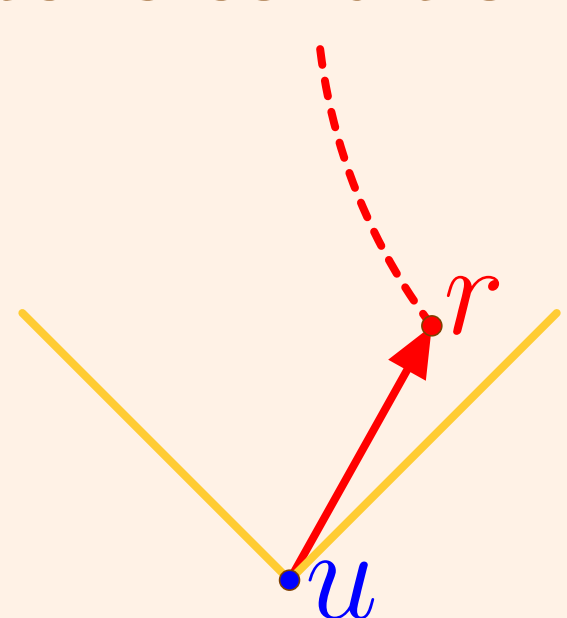


Figure 4. $u = 0$, $r \in \text{int}(\mathcal{L})$

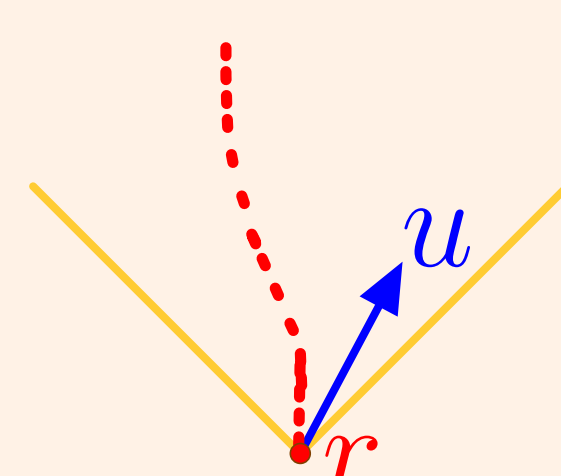


Figure 5. $u \in \text{int}(\mathcal{L})$, $r = 0$

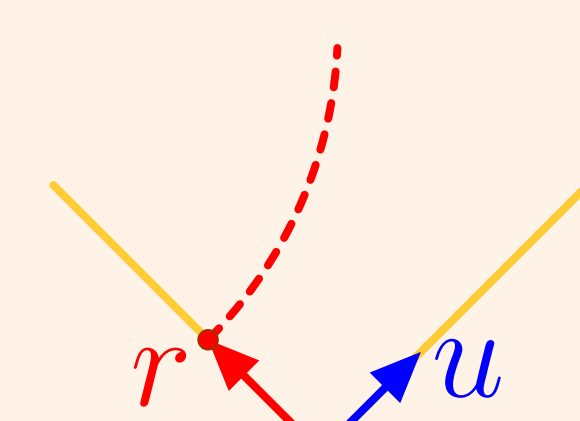


Figure 6. $u \in \text{bd}(\mathcal{L})$, $r \in \text{bd}(\mathcal{L})$

Then, the optimal solution r will converge to the **analytic center** of the optimal set \hat{R} which is the unique optimal solution of the problem

$$\min_{r \in \text{ri}(\hat{R})} \psi(r) := - \sum_{i \in B} \log \det r_i - \sum_{i \in R} \log r_{i,0},$$

where (B, N, R) is a partition of $\{1, \dots, n\}$.

- This result can be extended to the case of rolling friction.