

In Situ Study of Zinc Peroxide Decomposition to Zinc Oxide by X-Ray Absorption Spectroscopy and Reverse Monte Carlo Simulations

Alexei Kuzmin,^{*} Inga Pudza, and Konstantin Klementiev

The Zn K-edge X-ray absorption spectroscopy is used to investigate in situ the decomposition of zinc peroxide (ZnO_2) to zinc oxide (ZnO). Principal component and linear combination analyses of the extended X-ray absorption fine structure (EXAFS) spectra are employed to identify the phase composition of the oxide upon heating to 900 °C. Only the ZnO_2 phase is found up to 180 °C, whereas only the nanocrystalline ZnO phase is occurring above 250 °C. Detailed structural information on the temperature dependence of the local environment of zinc atoms is obtained using the reverse Monte Carlo (RMC) simulations. A strong increase of disorder is found upon approaching the decomposition temperature, evidenced by the broadening of Zn–O and Zn–Zn pair distribution functions (PDFs) and related mean-square relative displacements.

to track the phase composition during the decomposition of zinc peroxide to zinc oxide.

A more refined analysis of X-ray absorption spectra can be performed based on ab initio multiple-scattering (MS) theory^[16–20] to extract detailed structural information. If one is not limited to only the first coordination shell of the absorbing atom, then properly restricted MS expansion over a selected number of scattering paths is widely used.^[17,21,22] While such an approach is commonly adopted, it has a drawback when the analysis is extended to the outer coordination shells due to the rapidly increasing number of parameters.^[2] As a result, a correlation between

1. Introduction


X-ray absorption spectroscopy is a powerful technique for in situ and operando investigations of the local environment in materials with varying degrees of crystallinity and disorder.^[1–6] Good sensitivity of the method to the short-range order around a selected element makes it useful to track the phase transitions in complex compounds.^[7–12] Technically, this means that, first, it is necessary to determine the type of phase(s) and then extract the structural information using a proper method of analysis.^[6,13] Phase identification can be performed by the comparison of the experimental data with a set of preselected reference compounds using principal component analysis (PCA) and/or linear combination analysis (LCA) procedure.^[14,15] The PCA method allows one to determine the minimum number and type of probable phases present in the sample, whereas their amount can be further quantified by the LCA. This approach will be used here

the parameters often leads to their nonphysical values despite a good agreement between the calculated and experimental extended X-ray absorption fine structure (EXAFS) spectra. To overcome this problem, constraints can be used on the correlation between the parameters or the range of their valid values. A natural way to address this issue is to employ the reverse Monte Carlo (RMC) method choosing the starting model based on diffraction data, i.e., lattice parameters, and allowing only small atomic displacements during the simulation.^[23,24] Such an approach maintains the average structure of the material allowing to account for thermal and weak static disorder within ab initio MS formalism.^[25] Here, it will be demonstrated by analyzing the temperature-dependent Zn K-edge EXAFS spectra for zinc peroxide and zinc oxide as an example.

Zinc oxide exists at ambient conditions in two phases—wurtzite-type ZnO ^[26,27] and peroxide ZnO_2 .^[28,29] In both structures, oxygen atoms compose the first coordination shell of zinc. However, the coordination of zinc atoms is tetrahedral in ZnO , while it is octahedral in ZnO_2 . A peculiar crystallographic structure of ZnO_2 containing peroxide (O_2)^{2–} groups is responsible for its low decomposition temperature of about 200–230 °C,^[30,31] which can be further affected by reducing the size of ZnO_2 crystallites.^[28,32] Upon decomposition, ZnO_2 transforms to nanocrystalline ZnO (nano- ZnO). It was suggested in ref. [32] that a formation of an intermediate amorphous phase can occur before the growth of the nano- ZnO phase. The use of ZnO_2 as a precursor for the production of ZnO nanoparticles and thin films has been demonstrated in several studies.^[30,32–36] Among many applications,^[29] such nanoparticles may be of interest in printed optoelectronics as part of zinc oxide inkjet inks.^[37–39]

A. Kuzmin, I. Pudza
Institute of Solid State Physics
University of Latvia
Kengaraga street 8, LV-1063 Riga, Latvia
E-mail: a.kuzmin@cfi.lu.lv

K. Klementiev
MAX IV Laboratory
Lund University
PO BOX 118, SE-22100 Lund, Sweden

 The ORCID identification number(s) for the author(s) of this article can be found under <https://doi.org/10.1002/pssb.202200001>.

DOI: 10.1002/pssb.202200001

The process of ZnO₂ decomposition to ZnO has recently been studied in the temperature range from ambient to 500 °C by in situ combined X-ray diffraction (XRD) and X-ray absorption spectroscopy in ref. [36]. The results obtained were supplemented by thermogravimetric analysis (TGA). The weight loss observed in TGA and the change of the ZnO₂ (022) reflection in the XRD patterns indicated the decomposition of ZnO₂ and the formation of ZnO in the temperature range from 175 to 250 °C.^[36] It was also proposed in ref. [36] from the analysis of XRD patterns that a poorly crystalline (amorphous like) ZnO phase exists in the temperature range of about 250–350 °C, while above 350 °C it becomes more ordered. At the same time, an analysis of the Zn K-edge X-ray absorption near edge structure (XANES) and EXAFS suggested that a large number of defects are present in the ZnO obtained in this way.^[36]

In this study, we performed a high-temperature (20–900 °C) Zn K-edge X-ray absorption spectroscopy study of two zinc oxide phases (ZnO and ZnO₂) to probe their temperature-dependent local structure and lattice dynamics and to elucidate the mechanism of ZnO₂-to-ZnO decomposition. The PCA method was used to determine the number of phases that exist at each temperature, while the LCA method was used to determine the type of each phase. As a result, the temperature ranges of the existence and coexistence of the two phases were established. This information was used to analyze the local atomic structure of each phase in detail using RMC simulations.

2. Data Analysis

Selected in situ experimental Zn K-edge XANES and EXAFS spectra are shown in **Figure 1**. The EXAFS spectra were extracted following the conventional procedure.^[2]

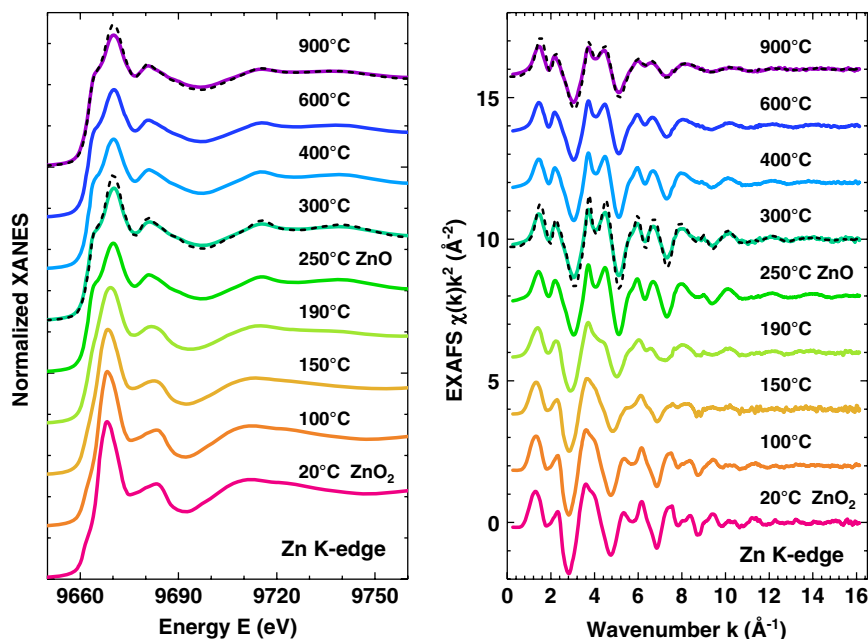


Figure 1. In situ experimental Zn K-edge X-ray absorption near edge structure (XANES) and extended X-ray absorption fine structure (EXAFS) spectra obtained in the temperature range of 20–900 °C during the decomposition of ZnO₂ to ZnO. The spectra for reference microcrystalline ZnO are shown by dashed lines at 300 and 900 °C for comparison.

The analysis of the extracted Zn K-edge EXAFS spectra was performed using the PCA method to evaluate the phase composition of the sample, which changes upon the ZnO₂ decomposition. Two main components corresponding to ZnO₂ and wurtzite-type ZnO phases were identified. The amount of the ZnO₂ fraction obtained by the LCA method as a function of temperature is shown in **Figure 2**. As one can see, two zinc oxide phases coexist in the temperature range of 190–240 °C, whereas only the ZnO₂ phase is present in the sample up to 180 °C, and only the ZnO phase exists above 250 °C. Therefore, the EXAFS spectra acquired in these two temperature ranges were analyzed further in detail based on the crystallographic structure of the respective zinc oxide phase using the RMC simulations.^[23,24] The EvAX code^[40] implementing the RMC method with the evolutionary algorithm approach was used. The method was previously applied to the analysis of the low-temperature EXAFS spectra of microcrystalline ZnO^[25,41] in and demonstrated good sensitivity to small variations of the local atomic structure upon the lattice thermal expansion.

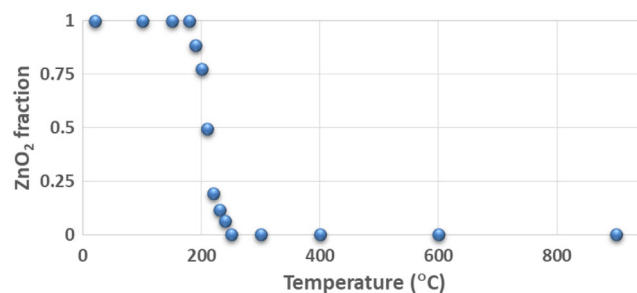


Figure 2. Decomposition process of ZnO₂ to ZnO according to the linear combination analysis of the Zn K-edge EXAFS data.

To perform RMC simulations, an initial model of the oxide structure (a simulation box) was constructed as a supercell with a size of $4a \times 4a \times 4a$ for ZnO_2 or $4a \times 4a \times 2c$ for ZnO based on

the crystallographic parameters of the unit cell obtained by X-ray ($a = 4.896 \text{ \AA}$ for ZnO_2)^[29] or neutron ($a = 3.250 \text{ \AA}$ and $c = 5.204 \text{ \AA}$ for ZnO)^[42] diffraction. Periodic boundary

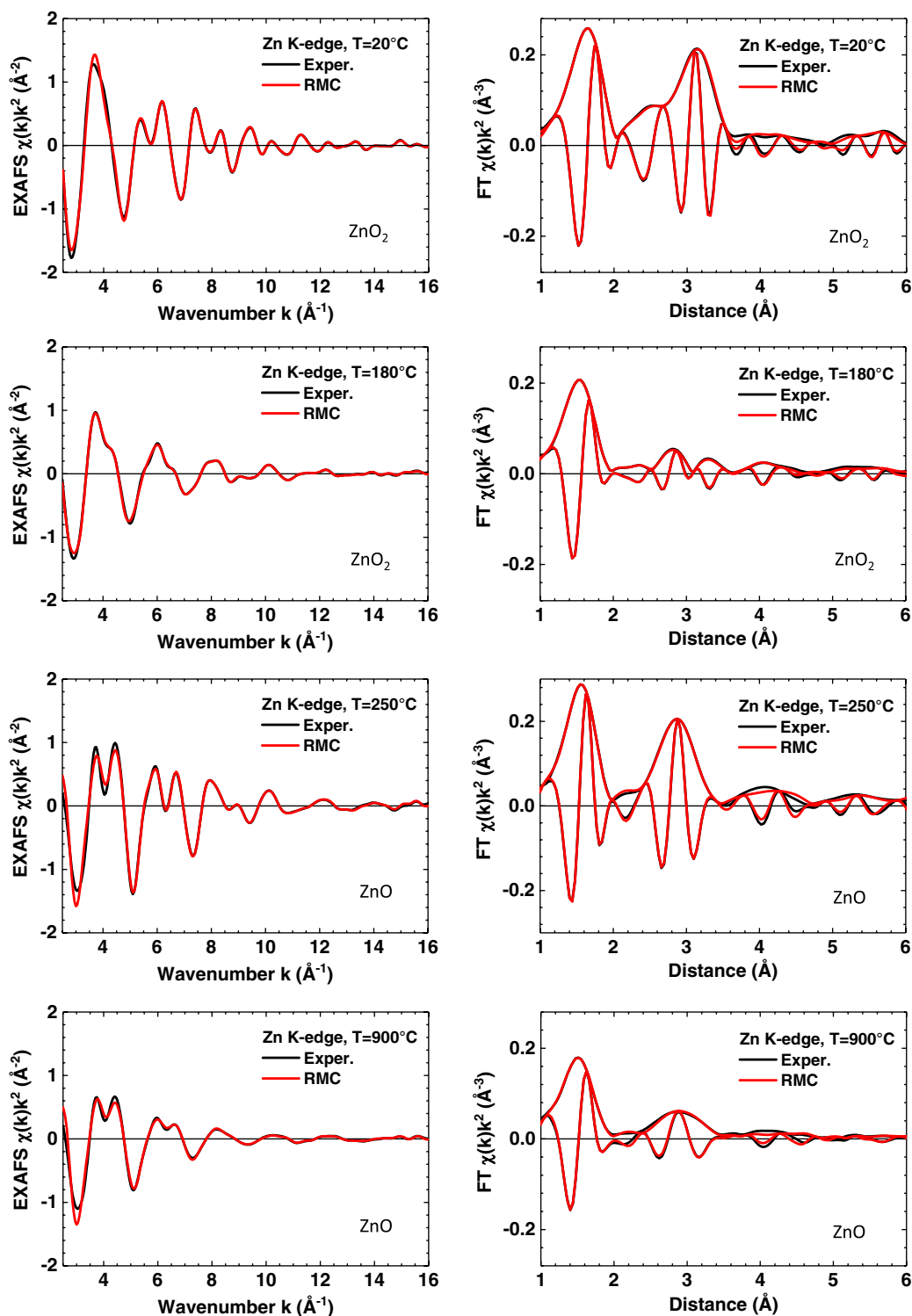


Figure 3. Comparison of the experimental and reverse Monte Carlo (RMC) calculated Zn K-edge EXAFS spectra $\chi(k)k^2$ (left panels) and their Fourier transforms (FTs) (right panels) for selected temperatures. The results for 20 and 180 °C correspond to the ZnO_2 phase, whereas for 250 and 900 °C to the ZnO phase obtained upon the decomposition of the peroxide.

conditions were employed to avoid problems with boundary effects. The size of the supercell was selected to be large enough to avoid artificial correlation effects and to collect reasonable statistics during acceptable calculation time. The number of atoms in the supercell was equal to the stoichiometric one for both oxide phases. Several RMC simulations with different pseudo-random number sequences were performed to accumulate even better statistics. The Zn K-edge EXAFS spectra were calculated for each zinc atom in the supercell and averaged to obtain the configuration-averaged (CA) EXAFS spectrum. At each step of the RMC simulation, all atoms in the supercell were randomly displaced with the maximum allowed displacement of 0.4 Å to account for thermal and/or static disorder. The difference between the Morlet wavelet transforms (WTs) of the experimental and CA-EXAFS $\chi(k)k^2$ spectra was used as a criterion for optimization.^[43] The use of the WT guarantees the agreement between experiment and theory simultaneously in k and R spaces.

Ab initio real-space multiple-scattering FEFF8.50L code^[17,44] was used to calculate theoretical CA-EXAFS spectra. The EXAFS amplitude damping due to the photoelectron inelastic losses was accounted for using the complex exchange-correlation Hedin–Lundqvist potential,^[45] and the EXAFS amplitude reduction factor S_0^2 was set to 1.0.^[2] Both single-scattering (SS) and multiple-scattering (MS) contributions were taken into account. MS contributions included all scatterings up to the 8th order. A large number of scattering paths was reduced by grouping similar ones and including into the analysis only important paths with the relative amplitude of corresponding

partial contribution to the total EXAFS greater than 0.1–1%.^[24] As a result of such path selection, only 60 and 43 scattering paths were involved in the cases of ZnO₂ and ZnO, respectively.

Comparison of the experimental and RMC calculated Zn K-edge EXAFS spectra $\chi(k)k^2$ for selected temperatures is shown in Figure 3 and 4 in the k and R spaces. Note that the peaks in the Fourier transforms (FTs) are located at distances that are slightly shorter than their crystallographic values because the FTs were not corrected for the phase shift present in the EXAFS equation.^[2] The coordinates of atoms in the final RMC simulation box were used to calculate pair distribution functions (PDFs) g_{Zn-O} and g_{Zn-Zn} and bond angle distribution functions (BADFs) for the O–Zn–O angles (Figure 5).

3. Results and Discussion

In situ Zn K-edge XANES and EXAFS spectra reported in Figure 1 indicate the change of the local zinc environment upon heating of ZnO₂ from 20 to 900 °C. The principal component analysis of a series of the Zn K-edge EXAFS spectra measured in situ during the ZnO₂-to-ZnO decomposition as a function of temperature suggests that there are only two main components, which are related to the two phases of zinc oxide. Therefore, the linear combination analysis of the EXAFS spectra using the reference data for ZnO₂ and wurtzite-type ZnO allowed us to determine the amount of the two phases upon heating (Figure 2). Only the peroxide phase exists up to 180 °C. The decomposition of ZnO₂ occurs in the temperature range of 190–240 °C, where both zinc oxide phases coexist. Finally, only

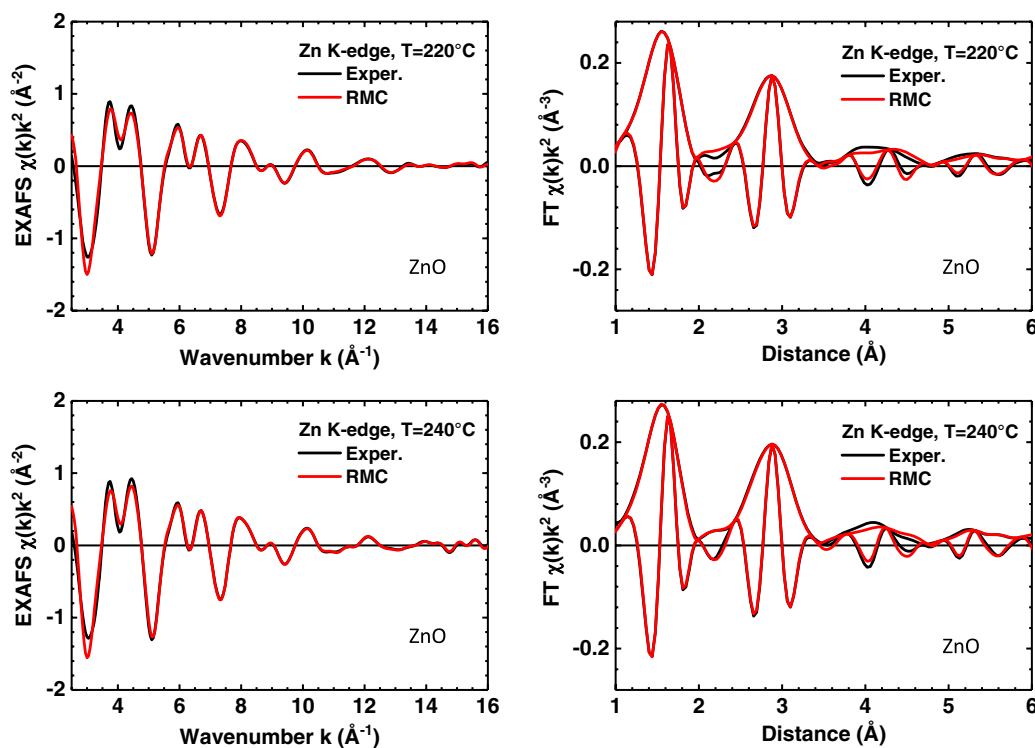


Figure 4. Comparison of the experimental and calculated (RMC) Zn K-edge EXAFS spectra $\chi(k)k^2$ (left panels) and their FTs (right panels) in the intermediate temperature range at 220 and 240 °C. The RMC simulations were performed using the wurtzite-type ZnO structure model.

the ZnO phase is present above 250 °C. However, a comparison with the reference spectra of microcrystalline ZnO at 300 and 900 °C (see dashed lines in Figure 1) indicates some expected differences due to the nanocrystalline nature of the decomposition product.^[36]

Based on the results of the PCA/LCA calculations (Figure 2), the RMC simulations of the Zn K-edge EXAFS spectra were performed considering two structural models—ZnO₂ for temperatures up to 180 °C and wurtzite-type ZnO for temperatures above 250 °C. Both models reproduce well the experimental data at respective temperatures in *k* and *R* spaces (Figure 3). Note that good agreement is observed in the outer coordination shells up to 6 Å with some small deviation for the peak around 4 Å in the Fourier transforms above 250 °C in the ZnO phase. The observed deviations could be related to the fact that our RMC model with the periodic boundary conditions corresponds to the infinite crystal and neglects structural disorder at the crystallite surfaces that affects stronger the outer coordination shells.

We have also found that the experimental EXAFS spectra in the intermediate temperature range from 220 to 240 °C can be well fitted using the wurtzite-type ZnO structure (Figure 4) since the expected contribution of the ZnO₂ phase from the PCA/LCA calculations is relatively small. Indeed, the observed agreement is the same as for higher temperatures where only the ZnO phase is present.

Comparison of the EXAFS spectra and their Fourier transforms in Figures 3 and 4 suggests their strong dependence on the oxide phase which is related to the difference in the zinc coordination (ZnO₆ vs ZnO₄) and the arrangement of the structural units in the crystallographic structures. While the thermal disorder contributes to the damping of EXAFS oscillations at large *k*-values and the reduction of peaks in the Fourier transforms upon increasing temperature, the variation of the crystallographic structure from ZnO₂ to ZnO produces a dominant effect well observed in the ranges of the first and, especially, second coordination shells (peaks at around 1.5 and 2.8 Å in FTs).

Temperature-dependent atomic PDFs *g*(*r*) for Zn–O and Zn–Zn atom pairs obtained from the results of the RMC simulations are shown in Figure 5. They reflect nicely the evolution of the local atomic environment around zinc atoms. Both PDFs in the ZnO₂ phase broaden significantly upon approaching the decomposition temperature when ZnO₂ transforms to nano-ZnO. From 220 °C, the local environment of the wurtzite-type ZnO phase starts to dominate. However, further increase in temperature promotes the growth of ZnO crystallites, which starts to compete with the increasing thermal disorder. The variation of the local structure in zinc oxides can be also described in terms of the BADFs for the O–Zn–O angles (Figure 5). The BADFs change as expected from 90° in ZnO₆ octahedra to about 109° in ZnO₄ tetrahedra and broaden upon increasing temperature.

Finally, we compared in Figure 6 low (24 °C) and high (900 °C) temperature PDFs *g*(*r*) for Zn–O and Zn–Zn atom pairs in microcrystalline wurtzite-type ZnO^[46] and the one produced from ZnO₂. Their good agreement suggests that the crystallinity of nano-ZnO produced by decomposition almost recovers upon heating to 900 °C. Nevertheless, some broadening of the outer shell peaks remains in both PDFs after cooling to 24 °C.

The PDFs reported in Figure 5 were used to extract information on the mean-square relative displacements (MSRDs) for the

Zn–O and Zn–Zn atom pairs corresponding to the first and second coordination shells of zinc, respectively (Figure 7). The values of MSRDs were evaluated as the second moment of the PDFs in the range of the corresponding peaks. Strong bonding between Zn and O atoms is responsible for the lower MSRD $\sigma^2(\text{Zn–O})$ values than in the case of the Zn–Zn atom pairs. Both MSRDs $\sigma^2(\text{Zn–O})$ and $\sigma^2(\text{Zn–Zn})$ increase strongly in ZnO₂ before the decomposition temperature, whereas their

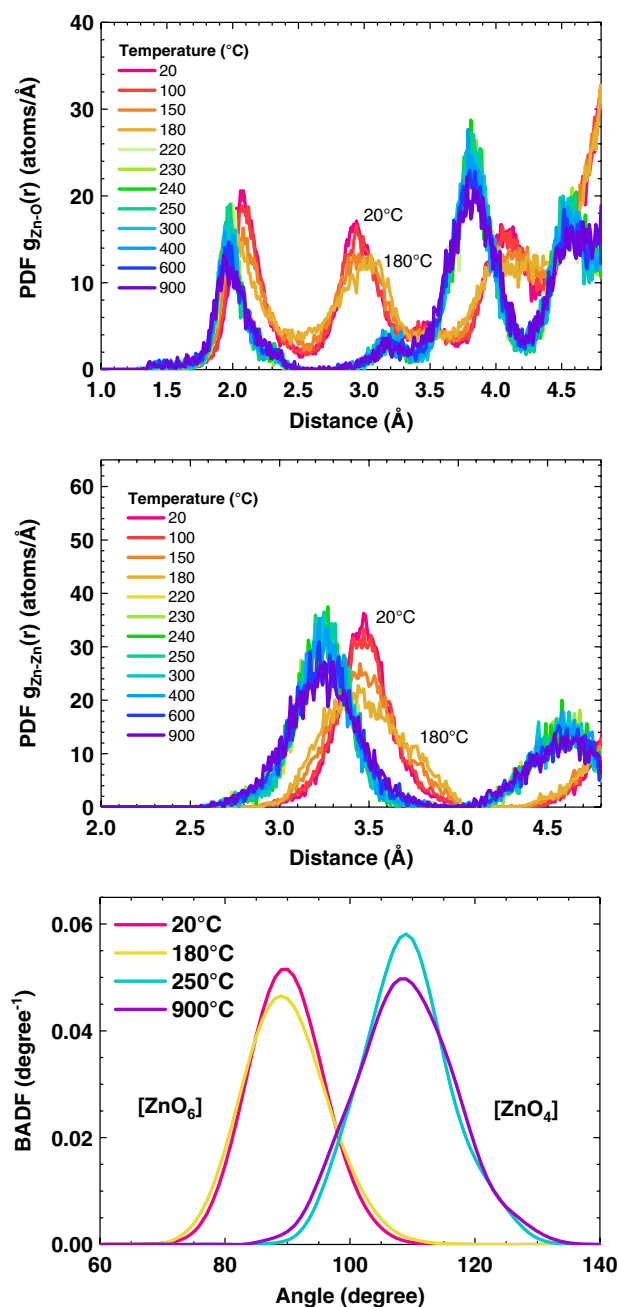


Figure 5. Temperature-dependent atomic pair distribution functions (PDFs) *g*(*r*) for Zn–O and Zn–Zn atom pairs and bond angle distribution functions (BADFs) for the O–Zn–O angles calculated from the results of the RMC simulations of the Zn K-edge EXAFS spectra.

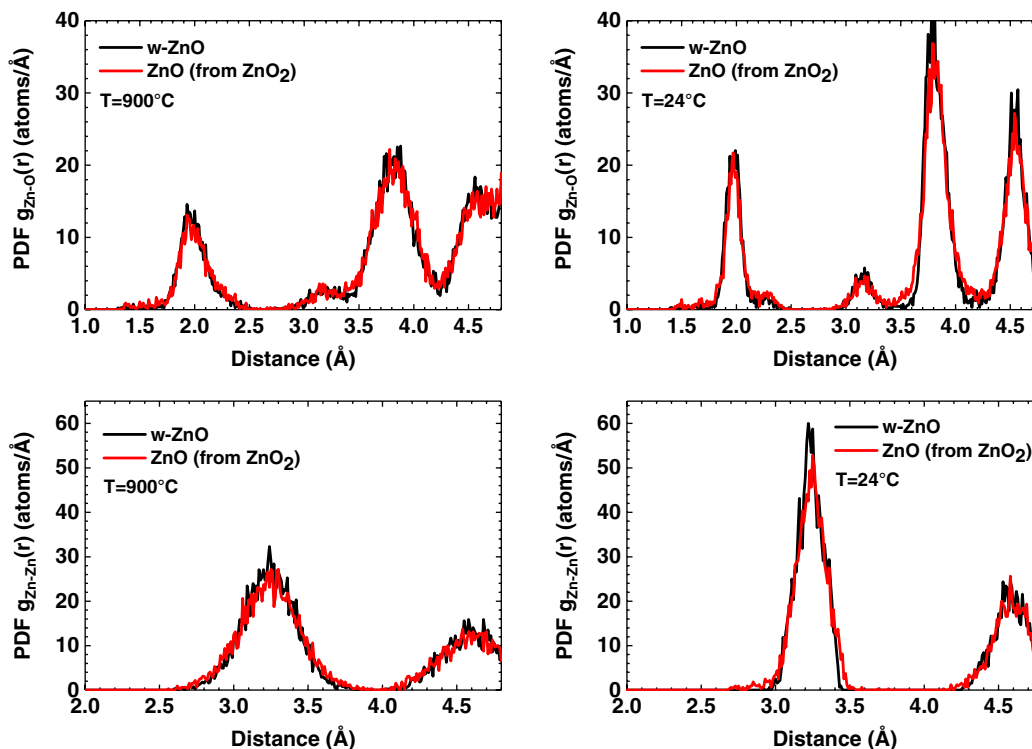


Figure 6. Comparison of the atomic PDFs $g(r)$ for Zn–O and Zn–Zn atom pairs at 900 °C and after cooling down to 24 °C in w-ZnO and ZnO obtained by decomposition of ZnO_2 . PDFs were calculated from the results of the RMC simulations of the Zn K-edge EXAFS spectra.

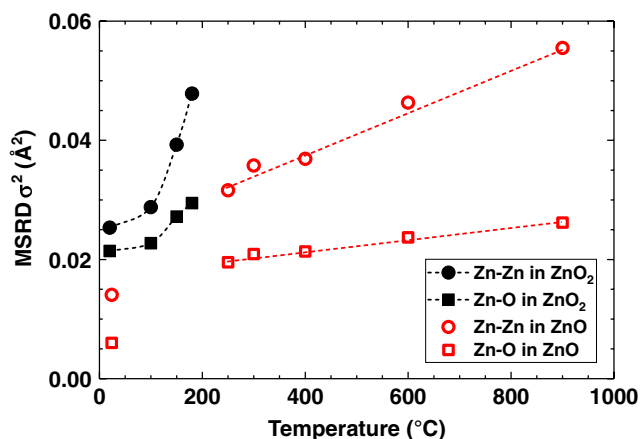


Figure 7. Temperature dependence of the mean-square relative displacements (MSRDs) σ^2 for Zn–O and Zn–Zn atom pairs in the first and second coordination shells, respectively. The dashed lines are guides for the eye. The MSRD values for ZnO at 24 °C are after cooling.

temperature dependence in the ZnO phase is almost linear. An increase in the size of the ZnO nanocrystals causes a slower than expected growth of the MSRDs with an increase in temperature above 250 °C. Such behavior of MSRDs is due to a decrease in static disorder partly compensates for the increase in thermal disorder. Note that the values of both MSRDs after cooling to 24 °C are close to that in bulk w-ZnO as supported by the similarity of their PDFs in Figure 6.

4. Conclusions

In situ high-temperature (20–900 °C) Zn K-edge X-ray absorption spectroscopy study has been performed to follow the changes of the local atomic structure during the decomposition of zinc peroxide to zinc oxide, occurring in the range of 190–240 °C. The use of the principal component and linear combination analyses of the EXAFS spectra allowed us to establish the temperature ranges of the existence and coexistence of the two oxide phases. We found that only the peroxide phase exists up to 180 °C, whereas only the nano-ZnO phase occurs above 250 °C. These conclusions were further supported by the results of the RMC simulations, which provided us with different distribution functions describing the local environment and its temperature dependence. The calculated atomic PDFs $g(r)$ for Zn–O and Zn–Zn atom pairs suggest an increase of disorder close to the decomposition temperature. This effect is also well observed in the temperature dependence of the related mean-square relative displacements. An increase of temperature to 900 °C promotes the growth of ZnO crystallites.

5. Experimental Section

Nanocrystalline ZnO_2 with the crystallite size of about 22 nm from ref. [29] and commercial polycrystalline wurtzite-type ZnO (99.99% purity, Alfa Aesar) were used in the powder form. For high-temperature experiments, each zinc oxide powder was mixed with BN powder and pressed into a pellet. The amount of zinc oxide was adjusted to give the value of the Zn K-edge (9659 eV) jump $\Delta\mu$ close to one.

Temperature-dependent Zn K-edge X-ray absorption experiments were performed in transmission mode on the BALDER beamline located at the 3.0 GeV storage ring of MAX IV Laboratory.^[47] X-ray beam from the in-vacuum wiggler source was monochromatized using a liquid-nitrogen-cooled double-crystal Si(111) monochromator and measured by two ionization chambers located before and after the sample and filled with N₂ and Ar gases. The uncoated Si collimating mirror and silica focusing mirror was used for harmonic reduction. The sample was heated in a nitrogen atmosphere using the Linkam furnace in the temperature range from 20 to 900 °C and then cooled down to 24 °C. In situ X-ray absorption spectra were collected at selected temperatures. Each EXAFS scan took ≈60 s, while the temperature was maintained fixed. To assure reproducibility, two repeats were taken for each temperature.

Acknowledgements

The authors wish to thank Dr. R. Kalendarev for the synthesis of ZnO₂ sample. A.K. would like to thank the financial support of the ERDF Project No. 1.1.1.1/20/A/060. The experiment at the MAX IV synchrotron was performed within the project 20190823. Institute of Solid State Physics, University of Latvia as the Center of Excellence has received funding from the European Union's Horizon 2020 Framework Programme H2020-WIDESPREAD-01-2016-2017-TeamingPhase2 under grant agreement No. 739508, project CAMART2.

Conflict of Interest

The authors declare no conflict of interest.

Data Availability Statement

The data that support the findings of this study are available from the corresponding author upon reasonable request.

Keywords

EXAFS, phase transition, reverse Monte Carlo method, X-ray absorption spectroscopy, ZnO, ZnO₂

Received: January 3, 2022

Revised: March 13, 2022

Published online:

- [1] S. Bordiga, E. Groppo, G. Agostini, J. A. van Bokhoven, C. Lamberti, *Chem. Rev.* **2013**, *113*, 1736.
- [2] A. Kuzmin, J. Chaboy, *IUCrj* **2014**, *1*, 571.
- [3] L. Mino, E. Borfecchia, J. Segura-Ruiz, C. Giannini, G. Martinez-Criado, C. Lamberti, *Rev. Mod. Phys.* **2018**, *90*, 025007.
- [4] W. Grünert, K. Klementiev, *Phys. Sci. Rev.* **2020**, *5*, 20170181.
- [5] P. Ghigna, E. Quartarone, *J. Phys. Energy* **2021**, *3*, 032006.
- [6] J. Timoshenko, B. Roldan Cuenya, *Chem. Rev.* **2021**, *121*, 882.
- [7] H. Arai, K. Sato, Y. Orikasa, H. Murayama, I. Takahashi, Y. Koyama, Y. Uchimoto, Z. Ogumi, *J. Mater. Chem. A* **2013**, *1*, 10442.
- [8] S. Boccato, R. Torchio, S. Anzellini, E. Boulard, F. Guyot, T. Irifune, M. Harmand, I. Kantor, F. Miozzi, P. Parisiades, A. D. Rosa, D. Antonangeli, G. Morard, *Sci. Rep.* **2020**, *10*, 11663.
- [9] I. Pudza, A. Anspoks, A. Cintins, A. Kalinko, E. Welter, A. Kuzmin, *Mater. Today Commun.* **2021**, *28*, 102607.
- [10] T. Miyanaga, *Symmetry* **2021**, *13*, 1315.
- [11] X. Hao, H. Zhu, Z. Guo, H. Li, Y. Gong, D. Chen, *J. Phys.: Condens. Matter* **2021**, *33*, 215402.
- [12] H. Dua, R. Shukla, R. S. Dhaka, *Phys. Rev. B* **2021**, *103*, 174107.
- [13] A. Kuzmin, J. Timoshenko, A. Kalinko, I. Jonane, A. Anspoks, *Radiat. Phys. Chem.* **2020**, *175*, 108112.
- [14] T. Ressler, J. Wong, J. Roos, I. L. Smith, *Environ. Sci. Technol.* **2000**, *34*, 950.
- [15] F. Panfili, A. Manceau, G. Sarret, L. Spadini, T. Kirpichtchikova, V. Bert, A. Laboudigue, M. A. Marcus, N. Ahamdach, M.-F. Libert, *Geochim. Cosmochim. Acta* **2005**, *69*, 2265.
- [16] A. Filipponi, A. Di Cicco, T. Tyson, C. Natoli, *Solid State Commun.* **1991**, *78*, 265.
- [17] J. J. Rehr, R. C. Albers, *Rev. Mod. Phys.* **2000**, *72*, 621.
- [18] A. Filipponi, A. Di Cicco, C. R. Natoli, *Phys. Rev. B* **1995**, *52*, 15122.
- [19] C. R. Natoli, M. Benfatto, S. Della Longa, K. Hatada, *J. Synchrotron Radiat.* **2003**, *10*, 26.
- [20] A. Di Cicco, *Radiat. Phys. Chem.* **2020**, *175*, 108077.
- [21] S. I. Zabinsky, J. J. Rehr, A. Ankudinov, R. C. Albers, M. J. Eller, *Phys. Rev. B* **1995**, *52*, 2995.
- [22] J. J. Rehr, J. J. Kas, M. P. Prange, A. P. Sorini, Y. Takimoto, F. Vila, C. R. Phys. **2009**, *10*, 548.
- [23] J. Timoshenko, A. Kuzmin, J. Purans, *Comput. Phys. Commun.* **2012**, *183*, 1237.
- [24] J. Timoshenko, A. Kuzmin, J. Purans, *J. Phys.: Condens. Matter* **2014**, *26*, 055401.
- [25] J. Timoshenko, A. Anspoks, A. Kalinko, A. Kuzmin, *Acta Mater.* **2014**, *79*, 194.
- [26] U. Özgür, Y. I. Alivov, C. Liu, A. Teke, M. A. Reshchikov, S. Doğan, V. Avrutin, S.-J. Cho, H. Morkoç, *J. Appl. Phys.* **2005**, *98*, 041301.
- [27] A. Janotti, C. G. V. de Walle, *Rep. Prog. Phys.* **2009**, *72*, 126501.
- [28] W. Chen, Y. Lu, M. Wang, L. Kroner, H. Paul, H. Fecht, J. Bednarcik, K. Stahl, Z. Zhang, U. Wiedwald, U. Kaiser, P. Ziemann, T. Kikegawa, C. D. Wu, J. Z. Jiang, *J. Phys. Chem. C* **2009**, *113*, 1320.
- [29] D. Bocharov, A. Chesnokov, G. Chikvaidze, J. Gabrusenoks, R. Ignatans, R. Kalendarev, M. Krack, K. Kundzins, A. Kuzmin, N. Mironova-Ulmane, I. Pudza, L. Puust, I. Sildos, E. Vasil'chenko, M. Zubkins, J. Purans, *J. Phys. Chem. Solids* **2022**, *160*, 110318.
- [30] N. Uekawa, N. Mochizuki, J. Kajiwara, F. Mori, Y. J. Wu, K. Kakegawa, *Phys. Chem. Chem. Phys.* **2003**, *5*, 929.
- [31] C. Bergs, P. Simon, Y. Prots, A. Pich, *RSC Adv.* **2016**, *6*, 84777.
- [32] T. Daley, E. Raj, S. Ramos, G. Cibin, A. Dent, T. I. Hyde, G. Sankar, *J. Phys.: Conf. Ser.* **2013**, *430*, 012080.
- [33] Y. Zhang, X. Wu, X. Y. Hu, R. Guo, *J. Cryst. Growth* **2005**, *280*, 250.
- [34] X. Han, R. Liu, W. Chen, Z. Xu, *Thin Solid Films* **2008**, *516*, 4025.
- [35] P. Rawat, Shalu, R. Nagarajan, *Mater. Lett.* **2018**, *212*, 178.
- [36] T. Daley, K. B. Opuni, E. Raj, A. J. Dent, G. Cibin, T. I. Hyde, G. Sankar, *J. Phys.: Condens. Matter* **2021**, *33*, 264002.
- [37] S. Sharma, S. S. Pande, P. Swaminathan, *RSC Adv.* **2017**, *7*, 39411.
- [38] G. Arrabito, Y. Aleeva, R. Pezzilli, V. Ferrara, P. G. Medaglia, B. Pignataro, G. Prestopino, *Crystals* **2020**, *10*, 449.
- [39] J. Morales-Mendoza, G. Herrera-Pérez, E. López-Martinez, D. Lardizabal-Gutierrez, F. Paraguay-Delgado, *J. Nanopart. Res.* **2021**, *23*, 34.
- [40] J. Timoshenko, EvAX code, **2017**, <http://www.dragon.lv/evax/> (accessed: March 2022).
- [41] J. Timoshenko, A. Anspoks, A. Kalinko, A. Kuzmin, *Phys. Status Solidi C* **2014**, *11*, 1472.
- [42] J. Albertsson, S. C. Abrahams, Å. Kvik, *Acta Crystallogr. B* **1989**, *45*, 34.
- [43] J. Timoshenko, A. Kuzmin, *Comput. Phys. Commun.* **2009**, *180*, 920.
- [44] A. Ankudinov, B. Ravel, J. Rehr, S. Conradson, *Phys. Rev. B* **1998**, *58*, 7565.
- [45] L. Hedin, S. Lundqvist, *J. Phys. C: Solid State Phys.* **1971**, *4*, 2064.
- [46] D. Bocharov, I. Pudza, K. Klementiev, M. Krack, A. Kuzmin, *Materials* **2021**, *14*, 5206.
- [47] K. Klementiev, K. Norén, S. Carlson, K. G. V. S. Clauss, I. Persson, *J. Phys.: Conf. Ser.* **2016**, *712*, 012023.