



## RESEARCH ARTICLE

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# Satellite-Based Detection of Secondary Droplet Activation in Convective Clouds

### Key Points:

- A new method to interpret the cloud microstructure of deep convective clouds from satellite remote sensing is presented in this study
- The microphysical process of secondary activation of droplets above cloud bases is demonstrated with a model and detected in observations
- The satellite detection of the secondary activation microphysical zone is introduced here for the first time

### Supporting Information:

Supporting Information may be found in the online version of this article.

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**Abstract** We present a new approach of analyzing and interpreting vertical profiles of cloud microstructure obtained by satellite remote sensing. The method is based on a spectral bin microphysics adiabatic parcel model and aims to elucidate the effects of aerosols on the evolution of convective clouds and related microphysical processes, including the activation of cloud condensation nuclei (CCN), the growth of cloud droplets, and the formation of precipitation. Characteristic features in the vertical profiles of effective radius ( $r_e$ ) and temperature ( $T$ ) reveal different microphysical zones in convective clouds related to the change increase of  $r_e$  with decreasing  $T$ . The classification of the different microphysical zones includes the (a) condensational growth of droplets, (b) growth by coalescence, (c) rainout, (d) secondary droplet activation zone (SAZ), (e) mixed-phase of ice particles and water droplets, and (f) glaciation of the cloud. The detection of the SAZ is introduced here for the first time. This method allows us to identify the activation of aerosol particles above cloud base and their role in the invigoration of deep convective clouds.

**Plain Language Summary** Using satellite remote sensing, we can obtain the vertical profiles of cloud microphysical processes in developed clouds. In this study, we present a new way of analyzing these profiles and understand the different processes that cloud droplets undergo during the development of the cloud. These processes include the turning of aerosols into cloud droplets, and the droplets' coalescence into rain drops. The change of droplets' size with height, as obtained by the satellite, reflect the different microphysical processes inside the cloud. While expecting the droplets to grow with height, a decrease of the droplets size suggests either rainout of the larger droplets from the cloud; or creation of small new droplets at much greater heights than the cloud base, where they are usually created. Those processes can be related to the extent of decrease of cloud drop size with height. The detection of new droplet formation above the cloud base with satellite is introduced here for the first time and it allows to understand this microphysical process and its effect of the development of the clouds and precipitation.

## 1. Introduction

### 1.1. Background

The effects of aerosols on clouds and climate are among the largest uncertainties in the assessment and modeling of climate change (IPCC, 2013, 2021). The investigation of aerosols and clouds using remote sensing techniques allows global monitoring and analysis of aerosol-cloud interactions (e.g., Bréon et al., 2002; Grosvenor et al., 2018; Rosenfeld et al., 2016; Rosenfeld and Lensky, 1998). A main challenge, however, is to investigate aerosol and cloud properties at an accuracy that suffices to quantify the response of cloud microstructure and the corresponding radiative effects to aerosol perturbations. Aerosols serving as cloud condensation nuclei (CCN) determine the droplet number concentration ( $N_d$ ) at cloud base for a given cloud base updraft speed ( $w_b$ ), and

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subsequently dominate the cloud microphysical and precipitation processes. The impact of fine ( $d < 1 \mu\text{m}$ ) and ultrafine ( $d < 0.1 \mu\text{m}$ ) aerosol particles (UAP) on deep convective clouds can be substantial through the following two mechanisms:

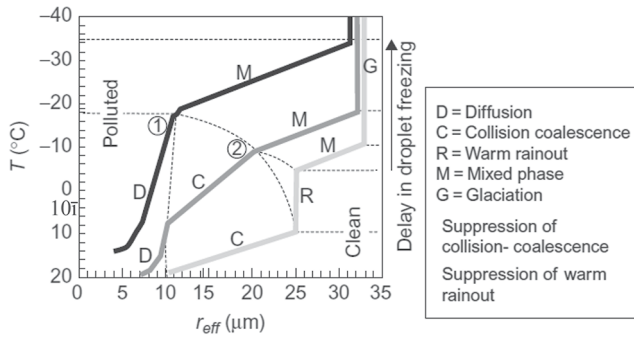
1. **Suppression of warm rain:** For a given updraft speed, polluted conditions with high CCN number concentrations ( $N_{\text{CCN}}$ ) lead to the formation of clouds with high  $N_d$  and small  $r_e$  at cloud base (Twomey, 1977). In such conditions, collision and coalescence processes are inhibited, thus delaying the droplets' growth into rain drops (Gunn and Phillips, 1957; Squires, 1958). Hence, extra latent heat is caused by the air parcels' ability to hold more liquid water content (LWC). This increases the updraft speed and maintains the clouds at a higher altitude. When postponing the rain initiation while the cloud continues to grow vertically, the cloud droplets are transported to above the zero-isotherm and freeze into ice hydrometeors (Braga et al., 2017). As a result, latent heat of freezing is released and increases the buoyancy, causing an invigoration of the cloud toward the tropopause, where it spreads aside into wide anvils (Rosenfeld, Lohmann, et al., 2008).
2. **Invigoration by activation of droplets above cloud base:** This is typically found in clean air masses with relatively low  $N_{\text{CCN}}$ , which result in large  $r_e$  and a corresponding high coalescence efficiency. The drop coalescence reduces the integrated droplets' surface area, which reduces the condensation rate. As a result, the supersaturation ( $S$ ) increases to the point where it exceeds the  $S$  at cloud base allowing activation of CCN with high critical  $S$  (typically UAP) that were not activated at cloud base. The activation of particles into cloud droplets high above the cloud base releases additional latent heat, increases the buoyancy, and invigorates the clouds (Fan et al., 2018; Khain et al., 2012).

The invigoration caused by these mechanisms creates larger and colder anvils (Fan et al., 2013; Pan et al., 2021), which may emit less longwave radiation to space and thus creates a positive radiative forcing and warms the Earth (Cotton et al., 2011; Slingo & Slingo, 1988). In addition, the evaporation of tiny ice crystals from these high anvils is likely to enrich the upper troposphere with water vapor, which is a strong greenhouse gas and thus likely to increase the positive radiative forcing.

Aerosol optical depth (AOD) has been extensively used as a proxy for  $N_{\text{CCN}}$  (Feingold et al., 2001; Quaas et al., 2009). The Aerosol Index (AI), which is the product of AOD and Angström exponent (Nakajima et al., 2001), provides a better estimation for  $N_{\text{CCN}}$ . However, Shinozuka et al. (2015), showed that such optically retrieved  $N_{\text{CCN}}$  could vary by an order of magnitude compared to the directly measured  $N_{\text{CCN}}$ . Furthermore, the satellite-retrieved AOD and AI become insensitive to  $N_{\text{CCN}} < 150 \text{ cm}^{-3}$  due to the small signal-to-noise ratio. This threshold varies greatly depending on the size of the CCN and their activation supersaturation (Shinozuka et al., 2015).

The  $N_d$  can also be retrieved based on the vertical evolution of  $r_e$  in the growing convective towers. This approach relies on the remarkable aircraft observations, which show that  $r_e$  increases with height above cloud base almost as if the clouds were adiabatic, despite the fact that the cloud LWC is mostly smaller than adiabatic (Freud et al., 2011). This pattern is caused by the nearly inhomogeneous mixing behavior of the convective clouds with the ambient air. Based on in-cloud aircraft measurements, Freud et al. (2011) showed that the cloud base  $N_d$  could be approximated by dividing the adiabatic LWC by the mass of a single droplet having an adiabatic  $r_e$ . Since both,  $r_e$  and cloud top temperature, can be retrieved by satellites,  $N_d$  could be calculated based on the assumption of adiabatic  $r_e$  and LWC (Rosenfeld, Fischman, et al., 2014; Rosenfeld, Liu, et al., 2014). The adiabatic LWC as a function of temperature can be calculated based on the retrieved cloud base temperature and pressure (Freud et al., 2011). The knowledge of the cloud base height above the surface,  $H_b$ , also allows estimating  $w_b$ , which can be obtained as  $w_b = 0.9H_b$ , with  $H_b$  in km, and  $w_b$  in  $\text{ms}^{-1}$  (Zheng and Rosenfeld, 2015). The cloud base maximum  $S$  ( $S_{\text{max}}$ ) can be calculated based on  $N_d$  and  $w_b$  (Pinsky et al., 2012). The  $N_d$  is then by definition  $N_{\text{CCN}}(S)$ . All these considerations were combined by Rosenfeld et al. (2016) into a satellite methodology to retrieve cloud base  $N_d$  and  $S$ , which further yields  $N_{\text{CCN}}(S)$  with an accuracy of  $\pm 30\%$ .

However,  $N_d$  and  $N_{\text{CCN}}$  can be retrieved from vertical  $T$ - $r_e$  profiles only as long as condensation dominates the cloud droplets' growth (i.e., for  $r_e < 14 \mu\text{m}$ ). Drop coalescence increases  $dr_e/dT$  beyond the condensational growth rate. Rosenfeld and Lensky (1998) provided a conceptual model of the vertical evolution of  $T$ - $r_e$  in deep convective clouds. They hypothesized that a stabilization or moderation of  $r_e$ , that is,  $dr_e/dT \approx 0$ , after exceeding a threshold of approximately  $14 \mu\text{m}$  was considered as resulting from rainout of the larger drops from the cloud. The appearance of ice particles increases the  $r_e$  further because ice particles are typically much larger than



**Figure 1.** Conceptual diagram of the microphysical zones in convective clouds based on the vertical evolution of the  $T$ - $r_e$  relationships as developed in (Rosenfeld & Lensky, 1998), here adopted from (Rosenfeld, 2018).

water drops in the same mixed-phase clouds. Eventually the cloud reaches a full glaciation, where the indicated  $r_e$  reaches a maximum. Rosenfeld and Lensky (1998) used the  $T$ - $r_e$  profile to partition the clouds into five vertical microphysical zones as shown in Figure 1: (a) Diffusion and condensational growth (in short “condensation”); (b) coalescence growth; (c) rainout; (d) mixed phase; (e) glaciated cloud.

## 1.2. Secondary Activation Zone

The  $T$ - $r_e$  conceptual model from Rosenfeld and Lensky (1998) in Figure 1, is based on the assumption of CCN activation into cloud droplets at cloud base. It does not consider the new activation of droplets in the course of cloud growth. The activation of particles into cloud droplets high above cloud base is referred to as secondary activation, as opposed to the primary activation of droplets at cloud base. Secondary activation can occur above cloud base in two ways: (a) The droplet coalescence reduces the integrated droplets'

surface area available for condensation, which reduces the condensation rate. As a result,  $S$  might exceed the  $S_{\max}$  at cloud base allowing inactivated UAP from the cloud base to nucleate into cloud droplets; (b) Additional entrained particles from the lateral and/or top boundaries of the cloud that can be activated even at the existing  $S$  at that level. Both ways will result in a reduction of the observed  $r_e$  with decreasing  $T$  due to the nucleation of small new droplets in the environment of coalesced droplets. This has been validated with several in-cloud penetration aircraft campaigns (Bera et al., 2022; Braga et al., 2017; Paluch and Baumgardner, 1989; Paluch and Knight, 1984; Prabha et al., 2011; Warner, 1969). The vertical microphysical zone where these processes occur is defined here as the “secondary activation zone” (SAZ). Clouds with enough latent heat left after warm rainout to keep the air parcel rising and which include a sufficient number of CCN can form a SAZ. The new droplets' formation releases enough latent heat to keep the cloud growing, increasing its longevity and height.

In this study, we describe an updated method to identify the microphysical zones of convective clouds based on  $T$ - $r_e$  profiles, taking the SAZ into account. We used a spectral bin microphysics adiabatic parcel model to simulate the  $N_d$ ,  $S$ , and the rain water fraction (RWF) profiles for four aerosol size distribution (ASD) and aerosol number concentrations below cloud base. The ASD and the simulated cloud properties are described in Section 3.1. Then, an application of the modeled cloud properties and their effect on the shape of the  $T$ - $r_e$  profile is described in Section 3.2, and the effect of real measured ASD on two satellite retrieved  $T$ - $r_e$  case studies is described in Section 3.3.

## 2. Methods

### 2.1. Model Description

In this study, we used an adiabatic cloud parcel model to describe the properties of convective clouds. The modeling was possible due to the inherent characteristic of  $r_e$  in convective clouds caused by the inhomogeneous nature of turbulent mixing at the cloud edges (Rosenfeld et al., 2016). The model is a spectral bin adiabatic parcel model developed by Pinsky and Khain (2002). This parcel model explicitly describes the cloud particle spectrum evolution using a moving mass grid containing 2000 mass bins, allowing effects on droplet collisions. The model describes the processes of aerosol particle growth, particle activation into cloud droplets, the vertical evolution of the cloud microstructure and precipitation-forming processes in the water phase ( $T > -10^\circ\text{C}$ ). The aerosol particles are considered as spherical NaCl particles to ensure compatibility with the model construction. The effect of different chemical compositions and hygroscopicity on cloud microphysics and ASD is beyond the scope of this study and has been addressed instead in Braga et al. (2021) and Pöhlker et al. (2021). Braga et al. (2017) have shown that the vertical evolution of the  $r_e$  measured by cloud probes generally agrees well with the theoretical  $r_e$  driven by the adiabatic growth of droplets through water condensation. When  $r_e$  exceeds  $14\ \mu\text{m}$ , coagulation and coalescence processes start to dominate. The coagulation process of cloud drops is calculated as proposed by Pinsky et al. (2001), who describe this process by solving stochastic collision equations according to Bott (1997). The coagulation of droplets is calculated after the initial condensational growth of particles. The drop-drop collision efficiency is calculated for each set of droplet masses allocated to the same parcel. Although

the model is very detailed, and strives to describe the basic physical processes, by considering only basic cloud properties while simplifying or ignoring additional environmental factors, the resulting microphysical processes may vary from real cloud behavior. This is discussed in Section 4.

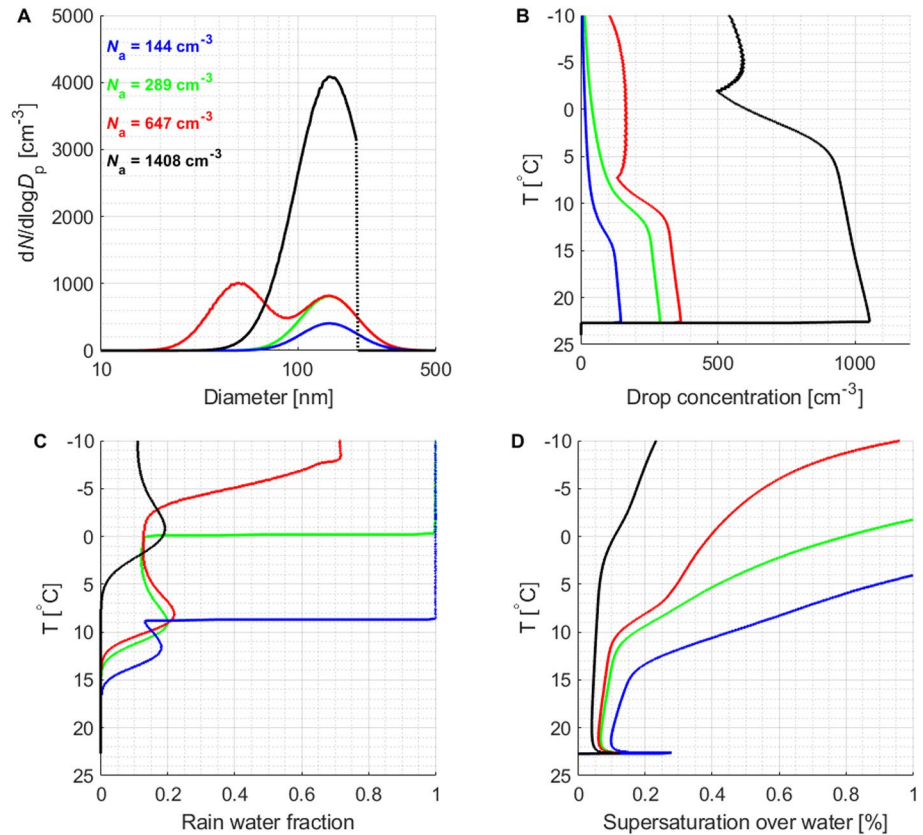
## 2.2. Satellite $T$ - $r_e$ Retrieval

The satellite retrievals of cloud properties for the case studies described in this study were performed using the Visible Infrared Imaging Radiometer Suite (VIIRS) passive sensor, onboard the Suomi National Polar-orbiting Partnership (NPP) satellite. The VIIRS sensor has a very high spatial resolution of 375 m at the nadir in five optical channels for detailed spectral information (Rosenfeld, Fischman, et al., 2014; Rosenfeld, Liu, et al., 2014). High spatial resolution is required to resolve the vertical structure of convective clouds. At lower resolution only the largest individual convective clouds can be resolved, thus missing the fine structure required to detect the SAZ. A snapshot of the  $T$ - $r_e$  profile of a cloud cluster provides the same information as obtained by tracking a single convective cloud throughout its vertical evolution (Arakawa and Schubert., 1974; Lensky and Rosenfeld, 2006). It is also assumed that  $r_e$  in the cloud top pixel has the same value as in more developed clouds at the same height, as long as precipitation does not fall through that level (Freud et al., 2008; Rosenfeld and Lensky, 1998). These assumptions enable to infer vertical  $T$ - $r_e$  development of clouds in a two-dimensional spatial image. Moreover, to reduce the data distortion, the satellite zenith angle of the clouds has to be between  $30^\circ$  and  $-20^\circ$ , and the solar zenith angle of the clouds has to be below  $65^\circ$ . The wavelengths used to retrieve the  $T$ - $r_e$  profile include 3.7 and 10.8  $\mu\text{m}$  (Rosenfeld, Fischman, et al., 2014; Rosenfeld, Liu, et al., 2014). This method assumes retrieval of water drops. However, ice particles are much larger than the droplets in the clouds, in which they are formed. They absorb more strongly the solar radiation at the 3.7  $\mu\text{m}$  band. Thus, the satellite retrieved  $r_e$  of ice particles appear larger than the same size water drops. Therefore, a very large  $r_e$  at supercooled temperature usually indicates the existence of ice in the cloud.

A designated software produces an RGB display that highlights the cloud microphysics and  $T$  and  $r_e$  retrievals in each pixel. On the basis of this display, cloud scenes of interest are manually sampled to avoid or minimize obstruction by multilayer or semi-transparent clouds that can disrupt the retrieved  $T$ - $r_e$ . A cloud mask algorithm, introduced by Zhu et al. (2014), filters out the partially filled cloudy pixels due to measurement errors in such pixels. The specific  $r_e$  values for each pixel are sorted by the detected temperature for each pixel into  $1^\circ\text{C}$  sized bins (i.e., 9.5–10.5 $^\circ\text{C}$  in the 10 $^\circ\text{C}$  bin etc.). Then, the 30-percentile median for all  $r_e$  values within one bin is calculated resulting in a continuous vertical cloud profile of  $r_e$  as a function of  $T$  (Rosenfeld et al., 2016). The effect of aerosols on the  $T$ - $r_e$  profile is mostly reflects in the growing clouds, rather than mature dissipating clouds. Thus, the use of the 30-percentile median is meant to capture the clouds in their growing phase (Lensky and Rosenfeld, 2006; Rosenfeld, Woodley, et al., 2008). Multilayer clouds are identified by a lack of continuity in the  $T$ - $r_e$  profiles. Therefore, this methodology is not valid for non-convective clouds or scenes obscured by multilayer clouds. The uncertainties of satellite-retrieved  $r_e$  and  $T$  are approximately 8% and 0.2–1.1 $^\circ\text{C}$ , respectively (Rosenfeld et al., 2016). Although, the importance of those uncertainties is reduced because the method described in this study is examining the relative change of  $r_e$  with  $T$  and not the absolute values.  $T$ - $r_e$  profiles retrieved in various locations show similar results for similar environments (e.g., Huang et al., 2022; Lensky and Rosenfeld, 2006; Rosenfeld, 2007; Rosenfeld, Woodley, et al., 2008) implying the robustness of  $T$ - $r_e$  retrievals and the fact that this method can be applied worldwide, in any meteorology and for any convective clouds as long as a valid unperturbed  $T$ - $r_e$  profile is retrieved. This study analyzes the satellite retrieved  $T$ - $r_e$  profiles by an automatic algorithm to detect the different microphysical zones. The entire description of the algorithm is described in Section 3.2.

## 2.3. Aerosol Observations at the Amazon Tall Tower Observatory

The Amazon Tall Tower Observatory (ATTO, 2020) is located in the central Amazon Basin (2.145 $^\circ\text{S}$ , 59.004 $^\circ\text{W}$  at 130 m above sea level). It monitors fundamental climatic and atmospheric parameters in a vast region of pristine rainforest. Some of the meteorological measurements at ATTO include temperatures, wind profiles, and precipitation (Andreae et al., 2015). These parameters were used to determine the coupling state of the clouds scenes and to choose case studies where the air parcel that was sampled at ATTO fed the cloud that was sampled by the satellite. For each of the case studies described in Section 3.3, an ASD was measured by a Scanning Mobility Particle Sizer (SMPS, TSI Inc.) connected to an inlet line with 60 m height at an 80 m tall mast (Franco et al., 2022). An overview of the atmospheric, geographic, and ecological conditions, as well as



**Figure 2.** Simulations of cloud properties for four different aerosol size distributions (ASD) and total particle number concentrations ( $N_a$ ). Panel A shows the four ASD including the cases (1) very clean unimodal (blue) with  $N_a = 144 \text{ cm}^{-3}$ , monomodal size distribution, with accumulation mode only; (2) clean unimodal (green) with  $N_a = 289 \text{ cm}^{-3}$ , monomodal size distribution, with accumulation mode only; (3) clean bimodal (red) with  $N_a = 647 \text{ cm}^{-3}$ , bimodal size distribution, the same number of accumulation mode particles as clean unimodal and additional Aitken mode particles; (4) polluted (black) with  $N_a = 1,408 \text{ cm}^{-3}$ , wide monomodal size distribution, with mostly accumulation mode reaching into the Aitken mode. The polluted ASD is truncated at 200 nm to prevent early warm rain initiation by enhanced coalescence. Panel B shows the simulated  $N_d$  with decreasing  $T$ . Panel C shows the vertical rain water fraction (RWF) profile. RWF is defined as the ratio between rain water content to total water content, where drops with a diameter  $>50 \mu\text{m}$  are considered raindrops. Due to the lack of significant activation of new droplet in the very clean unimodal and clean unimodal cases (shown in Figure 3), RWF jumps to unity when the smallest drops grow beyond the threshold of  $50 \mu\text{m}$ . In reality, the drop size distribution is wider, and therefore, there would not be such a sharp jump. In addition, clouds that completely rainout cannot develop due to a lack of condensation and buoyancy. The model forces the cloud to grow beyond this point in order to compare clouds profiles with and without significant new activation. Panel D shows the vertical profile of supersaturation over water.

commonly encountered air masses, can be found, for example, in Artaxo et al. (2013), Andreae et al. (2015), and Pöhlker et al. (2019). Characteristic aerosol and CCN conditions at ATTO have been described in recent studies (e.g., Moran-Zuloaga et al., 2018; Pöhlker et al., 2016, 2018).

### 3. Results

#### 3.1. Simulations of Cloud Microphysical Properties

This section discusses the vertical  $T$ - $r_e$  profiles of convective clouds using the parcel model with various initial-ization values. The simulations were performed for four ASD measured in subsaturated conditions (Braga et al., 2021; Pöhlker et al., 2021) as shown in Figure 2a. The data and initial thermodynamic conditions for the simulations, such as pressure, temperature, relative humidity, and height were collected below cloud bases of convective clouds during flight AC19 of the ACRIDICON-CHUVA campaign (Wendisch et al., 2016). The ASD include three monomodal ASD and one bimodal ASD that simulates a Hoppel minimum. The Hoppel minimum is a size separation gap between two size modes, which usually indicates the size of the smallest particles that can

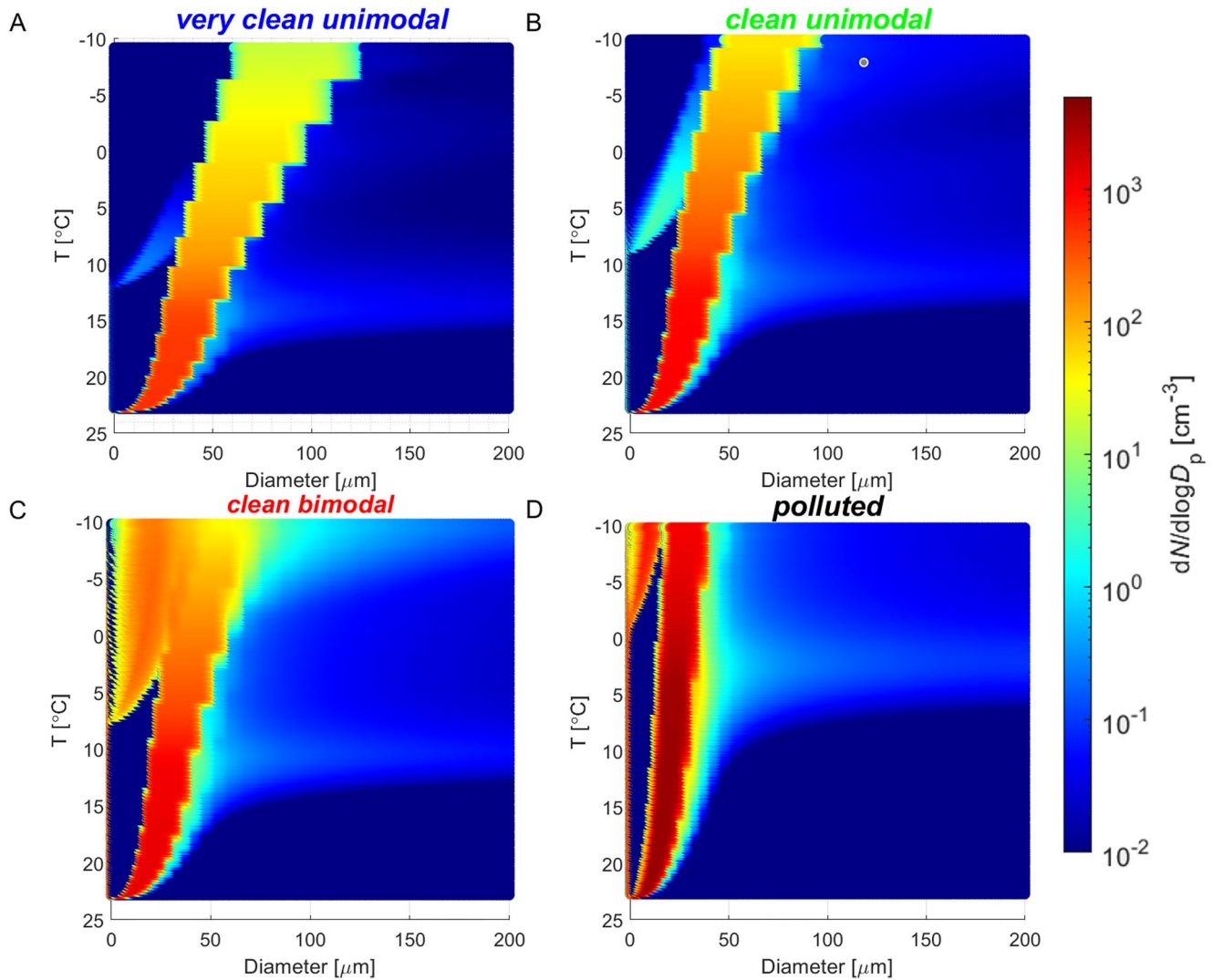
serve as CCN for the given thermodynamical conditions due to cloud processing (Hoppel and Frick, 1986; Hudson et al., 2015). The two size modes are the Aitken mode, with particles smaller than the Hoppel minimum (from approximately 10 to 70 nm), and the accumulation mode, which includes the larger ones (from approximately 70 to 200 nm) (Seinfeld and Pandis, 1998). The three monomodal ASD have only an accumulation mode for different aerosol number concentrations (cases: very clean unimodal, clean unimodal, and polluted) and the bimodal ASD has Aitken and accumulation modes (case: clean bimodal). Even though observations show no warm rain for similarly polluted cases (Braga et al., 2017), the model still simulates rain due to its construction. Therefore, to prevent or at least delay the initiation of warm rain by nucleation of large droplets and enhanced coalescence, the polluted ASD is truncated at 200 nm and lacks particles larger than this threshold. The comparison of the polluted ASD with and without the truncation is illustrated in the supplements (Figure S1). The increase of updraft speed with height is assumed to be linear with a constant acceleration of  $0.9 \text{ ms}^{-1} \text{ km}^{-1}$ . These are typical velocities found near the base of convective clouds (Zheng and Rosenfeld, 2015). The aerosol spectra and the respective vertical profiles of cloud properties for each ASD case are presented in Figure 2, the droplet size distribution (DSD) for each  $T$  bin is presented in Figure 3 and the corresponding  $T-r_e$  profiles from the model runs are shown in Figure 4.

The number of activated drops at cloud base is mostly influenced by accumulation mode particles, which are activated first (Pöhlker et al., 2016, 2018). As can be seen in Figure 2b,  $N_d$  at cloud base is almost equal to the  $N_a$  for the very clean unimodal and clean unimodal cases, while for the clean bimodal and polluted cases,  $N_d$  at cloud base is much smaller than the total  $N_a$ , meaning that a large fraction of the particles was not activated at cloud base. The inactivated particles are activated aloft when  $S$  exceeds  $S_{\text{max}}$  at a higher altitude (Figure 2d). Figure 3 shows the DSD in each  $T$  bin, colored by the number concentration of each size bin. For all four ASDs an increase of droplet concentration appears above the cloud base. However, the number concentration of activated droplets above cloud base in the very clean unimodal and clean unimodal cases is of at least two orders of magnitude lower than in the clean bimodal and polluted cases. The ratio of the number of activated droplets with respect to the number of eliminated large drops in the clean bimodal and polluted cases is much larger, which means a much more significant SAZ. This is noticeable in the increasing drop concentration with decreasing  $T$  for those cases (Figure 2b). In addition, Figure 3 shows the conversion of cloud droplets into raindrops which populate the larger size bins.

RWF is defined as the ratio between the rain water content ( $d > 50 \mu\text{m}$ ) and the total water content (Figure 2c). RWF starts to increase as coalescence starts and to decrease as the rain drops precipitate from the cloud parcel. This process takes place when the terminal velocity of raindrops is larger than the updraft speed within the cloud parcel. When the SAZ is negligible (very clean unimodal and clean unimodal cases) RWF increases with decreasing  $T$  up to unity when the smallest drops grow beyond the threshold of  $50 \mu\text{m}$ . In reality, the drop size distribution would be wider, and therefore, there would not be such a sharp jump. In addition, clouds that completely rain out cannot develop higher due to a lack of condensation and buoyancy. In this study, the model forces the cloud to grow beyond this point to illustrate how secondary activation changes the  $T-r_e$  profile compared to the alternative that is rarely realized in nature. The height for the onset of rain formation above the cloud base (RWF > 0) increases for increasing  $N_d$  at cloud base (Figures 2c and 3). This is consistent with the delayed coalescence at greater  $N_d$  due to slower droplet growth with height. The  $N_d$  starts to decrease sharply at the temperature of rain initiation, as the coalescence merges many cloud droplets into fewer raindrops. For simulations in which  $N_d$  starts to increase above cloud base (clean bimodal and polluted) the existence of a significant SAZ is identified. This occurs at the height where  $S$  exceeds the  $S_{\text{max}}$  at the cloud base. In these cases, a considerable amount of aerosol particles that were not activated into droplets near the cloud base are activated at the larger  $S$  aloft. This is valid for these model runs since the model does not include entrainment of additional aerosol particles. However, in nature, the origin of these particles is not necessarily inactivated aerosols from the cloud base, but they could also be from the surroundings, for which  $S$  does not have to exceed  $S_{\text{max}}$ . The significant activation of the new droplets in these cases replenishes the cloud water content; therefore, RWF values do not go up to 1.

### 3.2. Detection of Microphysical Zones Based on Modeled $T-r_e$ Profiles

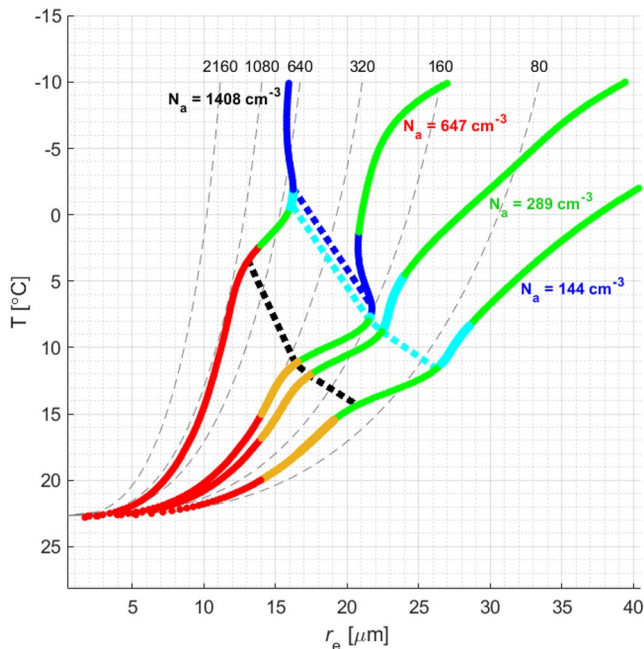
Different microphysical processes determine the shape of the  $T-r_e$  profiles of convective clouds. These processes were investigated using the model results described in Section 3.1. Figure 4 shows the  $T-r_e$  profiles of the four simulated clouds with varying shapes and colors representing the different microphysical zones. The gray lines in the background of the plot are the theoretical adiabatic  $r_e$  curves under the assumption of condensational growth for different cloud base  $N_d$ . The adiabatic curves are calculated based on the adiabatic LWC and the correspond-



**Figure 3.** The droplet size distribution (DSD) in each  $T$  bin, colored by the number concentration of each size bin, for the four model runs presented in Figure 2. It is inferred from this figure that SAZ occurs in all four ASD. However, in Panel A and B (very clean unimodal and clean unimodal cases) the number concentration of activated droplets above cloud base is at least two orders of magnitude lower than the clean bimodal and polluted cases shown in panels C and D. Therefore, the ratio of droplets activation to large drops elimination in the clean bimodal and polluted cases is larger, which means a much more significant SAZ. This causes an increase in drop concentration with decreasing  $T$  (Figure 2b). It is also shown here that the rain initiation is higher for cases with higher initial  $N_d$ .

ing development of  $r_e$ . The adiabatic LWC is computed using an adiabatic parcel model that rises from the cloud base temperature and pressure obtained from reanalysis (Rosenfeld et al., 2016). The deviations between the actual and the theoretic adiabatic  $T$ - $r_e$  profiles define the dominant microphysical zones, as follows:

1. **Condensation:** The first zone after droplets nucleation at cloud base (red lines in Figure 4). The  $r_e$  growth is similar to an adiabatic  $r_e$  curve originating at cloud base. Therefore, the  $T$ - $r_e$  line is nearly parallel to the background  $T$ - $r_e$  curve with the same cloud base  $N_d$ . During the condensational growth phase, no rain is expected, and thus, no rain was modeled.
2. **Adiabatic growth:** There are cases where the  $T$ - $r_e$  growth continues adiabatically well beyond the  $r_e$  threshold of rain initiation, which is commonly recognized near 14  $\mu\text{m}$  (e.g., Rosenfeld et al., 2006). Beyond this threshold, the microphysical zone can no longer be defined as dominantly condensational since coalescence has become very efficient and likely plays a significant role here as well. Accordingly, the continuation of the adiabatic increase in  $r_e$  beyond 14  $\mu\text{m}$  is defined as adiabatic growth here (orange lines in Figure 4) to be differentiated from the condensational growth. Such behavior may result from strong updrafts, where there is not enough time to form rain in the fast-rising parcels.



**Figure 4.**  $T$ - $r_e$  profiles of the four cloud simulations presented Figure 2, color-coded by different microphysical zones based on the growth rate of  $r_e$  with decreasing  $T$ . The thin gray dashed lines in the background are the theoretical adiabatic curves. The red lines are the condensational growth phase. The orange lines mark the adiabatic growth. The green lines mark the coalescence growth. The cyan lines mark the rainout, and the blue lines mark the secondary activation zone (SAZ). The thick dashed black, cyan and blue lines indicate the rain onset, maximum rainout, and the SAZ onset, respectively.

3. **Coalescence:** Drop coalescence increases the cloud drop size due to the drops merging. Therefore, a growth rate of  $r_e$  that exceeds the adiabatic growth rate marks the coalescence zone (green lines in Figure 4). The model shows the initiation of rain (Figure 2c) at the bottom of the coalescence zone. Consequently,  $N_d$  sharply reduces.
4. **Rainout:** Moderation of the growth rate of  $r_e$  with decreasing  $T$  above the coalescence zone indicates the rainout of the larger raindrops from the cloud parcel/top (cyan lines in Figure 4).
5. **Secondary activation zone:** A significant activation of new droplets after the occurrence of coalescence and/or rainout is indicated by negative  $dr_e/dT$  (blue lines in Figure 4).

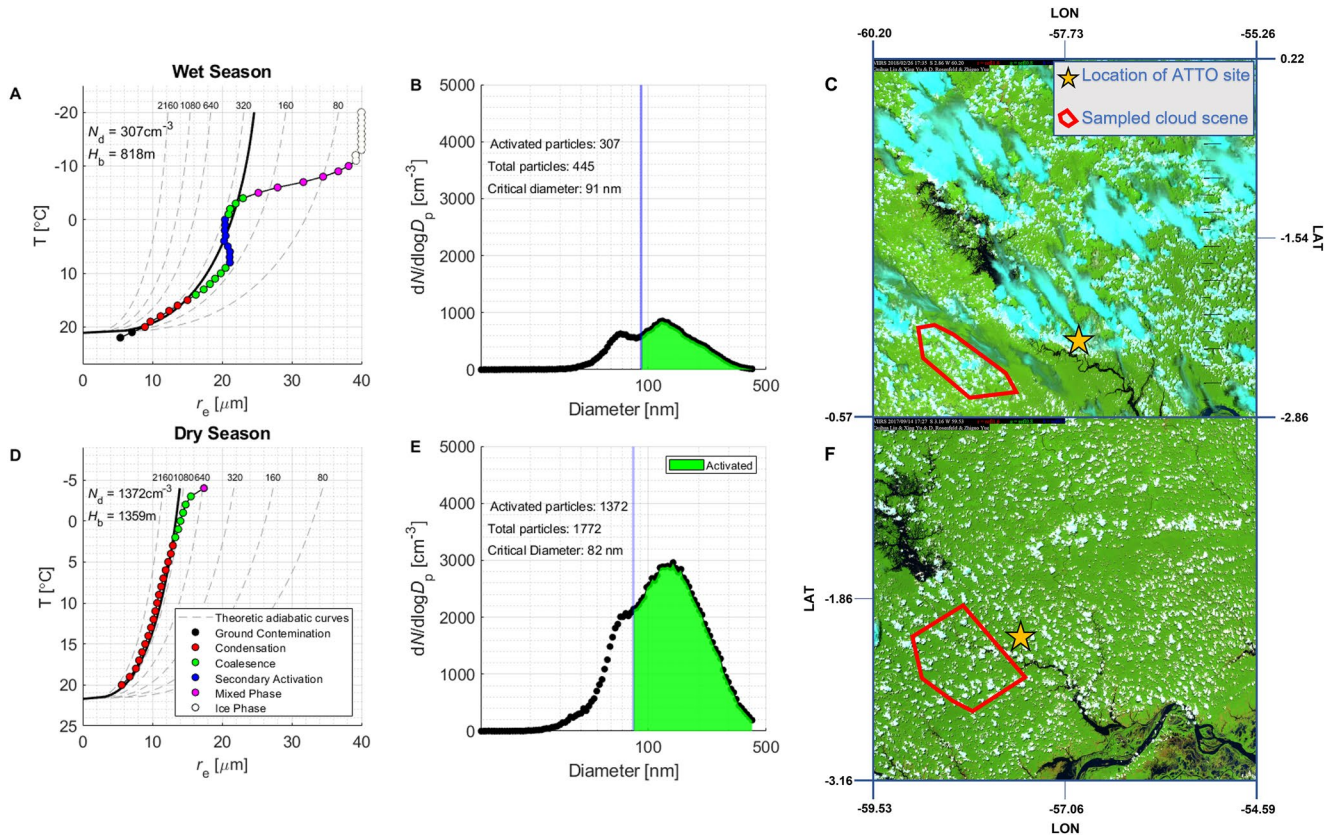
Figure 4 shows that the microphysical zones are distinguished from each other by the shape of the  $T$ - $r_e$  profiles. In all the four cloud simulations, the condensational part of the cloud (marked by red lines) behaves adiabatically and follows the theoretical adiabatic curves, depending on the  $N_a$  of each case. Then, the coalescence that comes after the adiabatic part indicates the cloud droplets' transition to raindrops. The depth of rain initiation ( $D_r$ , in meters above cloud base) is linearly correlated with  $N_d$  at cloud base, according to  $D_r \sim 4 \cdot N_d$  (Freud & Rosenfeld, 2012; Konwar et al., 2012). For example, in the very clean unimodal case,  $r_e$  increases fastest with decreasing  $T$  and reaches rain initiation at the warmest  $T$ . The maximum conversion rate of cloud into raindrops (maximum rainout, marked by the thick cyan dashed line) occurs at approximately 12°C. For cases with higher initial  $N_d$ , the temperature of rain initiation (marked by the dashed black line) and maximum rainout occur at colder temperatures (see Figures 2–4). Coalescence and precipitation lead to a strong reduction in the droplets' integrated surface area, causing limitations to the condensation rate. This increases  $S$  with decreasing  $T$ . In the clean and very clean unimodal cases, which have an accumulation mode only, all the aerosols get activated at cloud base, as seen in Figure 2b. The cloud droplets are

larger, and their precipitation causes only a slight decrease of  $dr_e/dT$ , which remains positive. However, in the clean bimodal and polluted cases, the significant amount of Aitken mode particles that were not activated at cloud base get activated above cloud base because of the increase of  $S$  beyond  $S_{\max}$  at cloud base. This results in a decrease of  $r_e$  with decreasing  $T$  (SAZ, marked by the blue lines). Such a pattern is observed due to the relative increase of the newly activated small droplet concentrations in the cloud parcel. This signature in  $T$ - $r_e$  profiles is not observed in the negligible secondary droplet activation above cloud base (see very clean unimodal and clean unimodal cases), because very little to no aerosols are left available for nucleation even at the very large  $S$  values aloft. Since the updraft in the model is forced to a linear increase of  $0.9 \text{ ms}^{-1} \text{ km}^{-1}$ , the cloud continues to grow vertically, leading to regrowth of  $r_e$  due to the ongoing coalescence of the remaining cloud droplets. As mentioned above, in nature, clouds with no new activation will rainout almost completely. They will reach very high  $S$ , but due to a lack of condensational heating, they will not develop any further. The practical impact of these results is to document the occurrence of secondary droplet activation well above the base of convective clouds with warm rain. The negative  $dr_e/dT$  in the modeled  $T$ - $r_e$  profile indicates the secondary activation of aerosols that were not activated at cloud base and implies that the SAZ can be detected by remote sensing. Nevertheless, the appearance of SAZ in nature is not necessarily connected to inactivated Aitken mode particles. It could also be caused by the entrainment of aerosols of any size that can be activated even at  $S$  below  $S_{\max}$ .

### 3.3. Detection of Microphysical Zones Based on Satellite $T$ - $r_e$ Profiles

There are unique relationships between the microphysical zones and their manifestation in the  $T$ - $r_e$  profiles. Accordingly, satellite-retrieved  $T$ - $r_e$  profiles can be used to infer microphysical zones, including the newly defined





**Figure 5.** Satellite retrieved  $T-r_e$  profiles of exemplary cloud scenes during the Amazonian wet and dry seasons close to the ATTO region, color-coded by different microphysical zones (panels A and D); and the 1-hr average ASD before the satellite overpass of each case study as measured by the SMPS on the ATTO mast (Panels B and E). Panels C and F show the true-color satellite image of each case. The ATTO location is marked by a star and the sampled cloud scene by the red polygon.

SAZ. Two cloud scenes close to ATTO were used as case studies to investigate and demonstrate the applicability of satellite  $T-r_e$  profiles. Cloud scene A (red polygon in Figure 5a) is a cumulonimbus cloud that was taken within approximately 69 km from the ATTO site during the wet season (26 February 2018 at 17:22 UTC). This season is characterized by relatively low CCN concentrations (Pöhlker et al., 2016). The retrieved  $N_d$  is  $307 \text{ cm}^{-3}$  and cloud base height is 818 m, as estimated by the lifting condensation level. The first two points of this  $T-r_e$  profile are semi-transparent cloud pixels that might cause a distortion of the retrieval and are considered as ground contamination. The 1-hr average ASD before the satellite overpass (Figure 5b) clearly shows a Hoppel minimum at 88 nm. The number of activated particles at cloud base (Figure 5b) is equal to  $N_d$ . Accumulating the  $N_d$  largest particles on the ASD yields the smallest diameter to be activated, that is, the critical diameter. In this case the critical diameter is 91 nm. This result suggests that the activated particles at cloud base are accumulation mode particles down to that threshold. The SAZ is clearly identified above the coalescence zone at the  $8^\circ\text{C}$  isotherm. The activation of either remaining Aitken mode particles from cloud base and/or entrained aerosols at higher levels, causes a reduction of  $r_e$  with decreasing  $T$ . The new droplets are added to the preexisting large cloud drops undergoing coalescence, which leads to a renewed growth of  $r_e$  with decreasing  $T$  above the  $0^\circ\text{C}$  isotherm. This means that the coalescence microphysical zone is again dominant there. The rapid growth of  $r_e$  toward colder subzero temperatures can also be interpreted as a mixed phase of liquid and ice and eventually full glaciation of the cloud (Rosenfeld and Lensky, 1998). Cloud scene B (red polygon Figure 5d) is a cluster of cumulus congestus clouds that was taken within approximately 26 km from ATTO during the dry season (14 September 2017 at 17:27 UTC), which is characterized by relatively high concentrations of CCN (Pöhlker et al., 2016). The retrieved  $N_d$  for the dry season case is  $1,372 \text{ cm}^{-3}$  and the cloud base height is 1,359 m. The 1-hr average ASD (Figure 5e) of this case is unimodal with an inflection point instead of a clear Hoppel minimum. Thus, the separation of modes is less clear in comparison with Cloud scene A. Higher  $N_d$  at cloud base with smaller  $r_e$  (compared to Cloud scene A), decrease the condensational growth rate of droplets with decreasing  $T$  and strongly suppress the coalescence

processes, therefore a SAZ signature is not identified. At the 2°C isotherm, the  $r_e$  becomes sufficiently large for the coalescence process to dominate up to the cloud's top.

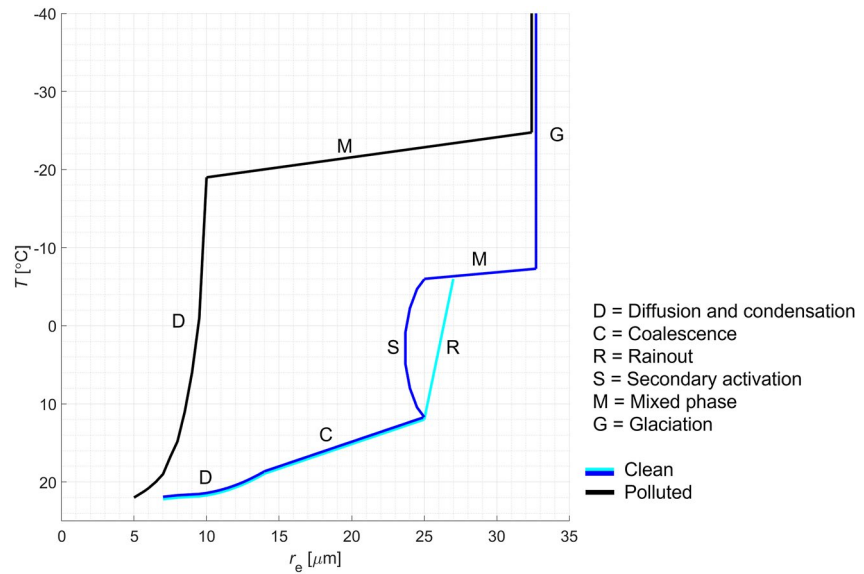
#### 4. Discussion

The SAZ signature allows us to identify a significant occurrence of new droplets nucleation above the cloud base. The origins of the particles on which the new droplets are formed can either be Aitken mode particles that were not activated at cloud base or entrained particles of all size aloft. In the first case,  $S$  aloft has to exceed the  $S_{\max}$  near cloud base. A stronger SAZ, which means lower  $-dr_e/dT$ , can imply an increased nucleation of smaller droplets, hence a greater latent heat release that increases the buoyancy. A deeper SAZ means that the top of the SAZ reaches lower temperatures. It can imply a deeper layer of aerosol activation above cloud base; thus the latent heat release occurs higher in the cloud. Increased buoyancy due to enhanced and/or higher occurrence of latent heat release leads to cloud invigoration. The invigoration of deep convective clouds may also lead to enhanced lightning activity (Thornton et al., 2017; Yuan et al., 2011).

As mentioned in Section 2.2, the parcel model is limited to cloud processes in the water phase only ( $T > -10^\circ\text{C}$ ). However, satellite retrieved cloud scenes can be more vertically developed than represented in the model, growing well above the  $-10^\circ\text{C}$  isotherm (see Cloud scene A). The rapid growth of  $r_e$  above the subzero temperatures indicates a mixed phase of ice and water particles. At a maximum threshold  $r_e$  of 40  $\mu\text{m}$  and/or temperatures colder than  $-38^\circ\text{C}$ , the ice phase is dominant, and the cloud is considered fully glaciated (Rosenfeld and Lensky, 1998). As mentioned in Section 3.1, another limitation of the model compared to observations is manifested in the polluted case. For example, cloud scene B compared to the modeled polluted case: Both cases represent highly polluted situations with high CCN concentration at cloud base. The parcel model initiates coalescence and warm rain at the 5°C isotherm and later on allows secondary activation, even with the truncation of the larger particles. However, in nature this does not happen in such polluted situations. Andreae et al. (2004), Rosenfeld, Fischman, et al. (2014), Rosenfeld, Liu, et al. (2014) and Braga et al. (2017) showed in different aircraft campaigns in deep tropical convective clouds that, when the air is highly polluted, the coalescence zone is mostly delayed to the mixed-phase zone or suppressed altogether. As can be seen in Cloud scene B, the coalescence is much weaker and leads into the beginning of the mixed phased zone with no secondary activation. Whether the origins of the particles are from the cloud base or entrainment aloft, the SAZ described in this study requires a large  $S$ , which is caused by the preceding coalescence. With that being said, in some cases secondary nucleation of droplets might also occur before coalescence dominates, during the condensational growth of the cloud. It might happen either due to entrainment, or in highly polluted clouds that are updraft limited. In those clouds the  $S_{\max}$  at the cloud base is rather small due to the activation of many particles. When the updraft increases with height to the extent that  $S$  exceeds the  $S_{\max}$  at cloud base, a SAZ might occur even during the condensational growth phase. This type of SAZ differs from the SAZ discussed in this paper and will be addressed in a subsequent study.

Figure 6 shows the updated conceptual  $T$ - $r_e$  model with the addition of the line of SAZ for clean and polluted clouds. The clean state (blue and cyan lines) is characterized by low aerosol concentration, thus the  $r_e$  increases rapidly above cloud base. The coalescence process is enhanced and leads to early warm rain. In the case where the cloud approaches supercooled temperatures and ice particles start to form, the observed  $r_e$  increases sharply with height until all droplets have frozen and  $r_e$  is stabilized. The cloud is then considered fully glaciated. The coalescence and rainout decrease the integrated surface area available for condensation, thus  $S$  increases. High  $S$  allows activation of additional cloud droplets if additional CCN are available, either from cloud base or entrained from aloft. In nature, clouds that undergo significant coalescence and/or warm rainout almost always experience SAZ (blue line). Otherwise, clouds would have difficulty developing further, due to the high  $S$  and a lack of sufficient condensational latent heating (cyan line).

In the polluted case (black line), there is a larger number of CCN at cloud base. That leads to a small rate of increase of  $r_e$  with height and suppression of the coalescence process. If the coalescence is suppressed up to the freezing level, rainout and SAZ will not appear on the  $T$ - $r_e$  profile.



**Figure 6.** An updated conceptual diagram of the microphysical zones for clean and polluted clouds based on the vertical evolution of the  $T$ - $r_e$  relationships including the secondary activation of droplets above cloud base. The clean cloud (blue and cyan lines) has low  $N_d$  and large  $r_e$  at cloud base. The enhanced coalescence leads to early warm rain. When the cloud reaches the subzero temperatures and start to glaciates,  $r_e$  is increased sharply with height. The coalescence and rainout decrease the integrated surface area available for condensation. The increased  $S$  allows activation of additional available CCN into cloud droplets. In nature, clouds with significant coalescence and/or warm rainout usually experience SAZ (blue line). Otherwise, clouds would have difficulty developing further, because of the high  $S$  and a lack of sufficient condensational latent heating (cyan line). The polluted case (black line) is characterized by larger  $N_d$  and a smaller  $r_e$  at cloud base. Suppression of the coalescence process occurs. If the suppression is all the way up to the freezing level, rainout and SAZ will not appear on the  $T$ - $r_e$  profile.

## 5. Summary

A method to identify the occurrence of secondary activation of droplets above deep convective cloud bases by satellite measurements was introduced in this paper. The earlier  $T$ - $r_e$  conceptual model by Rosenfeld and Lensky (1998) assumes CCN activation into cloud droplets at cloud base. Thus, a stabilization or moderation of  $r_e$  after exceeding a threshold of approximately 14  $\mu\text{m}$  with decreasing  $T$  was considered as resulting from rainout of the larger drops from the cloud. However, considering additional nucleation of droplets above cloud base, in the present study the moderation or inversion of  $dr_e/dT$  is ascribed to the existence of a significant secondary drop activation zone (SAZ). When pronounced coalescence and warm rain forming processes dominate during the cloud vertical growth, the droplets' integrated surface area is reduced. The reduction of the water surface area available for condensation decreases the condensation rate and increases the supersaturation to high values beyond the maximum supersaturation at cloud base. As a result, particles within the cloud parcel that could not be activated at cloud base due to their lower critical supersaturation are activated aloft. If the rate of large drops elimination is lower than the rate of new droplets activation, there is a decrease of  $r_e$  with decreasing  $T$  and increasing height. The origins of the newly activated particles can be either UAP from the cloud base and/or entrained particles aloft from the boundaries of the clouds.

The method presented in this study is an updated version of the previous conceptual  $T$ - $r_e$  model. The detection of SAZ signature on a  $T$ - $r_e$  profile, is introduced here for the first time. This conceptual model was demonstrated with adiabatic parcel model simulations for different scenarios of aerosol concentrations and size distributions. The SAZ signature was also observed in satellite-retrieved  $T$ - $r_e$  profiles. The possibility to apply this methodology to satellite retrieved  $T$ - $r_e$  profiles allows a detailed analysis of cloud microphysics in large areas and over long periods of time. This will increase the understanding of cloud-aerosol interactions at the micro and macro scale. The occurrence of the SAZ signature has climatic importance in shaping cloud microstructure aloft and a possible propagation to cloud radiative forcing.

## Disclaimer

This paper contains results of research conducted under the Technical/Scientific Cooperation Agreement between the National Institute for Amazonian Research, the State University of Amazonas, and the Max-Planck-Gesellschaft e.V.; the opinions expressed are the entire responsibility of the authors and not of the participating institutions.

## Conflict of Interest

The authors declare no conflicts of interest relevant to this study.

## Data Availability Statement

The satellite data of the case studies presented in this study is available on the VIIRS data set found at the NOAA's Comprehensive Large Array-data Stewardship System (2022; <https://www.avl.class.noaa.gov/>). The ATTO data used in this study are available via the ATTO data portal through <https://www.attoproject.org/>. The output data from the IDL program of the sampled cloud scenes have been deposited in supplementary data files for use in follow-up studies (Efraim et al., 2022; <https://doi.org/10.5061/dryad.h44j0zpnr>). For data requests beyond the available data, please refer to the corresponding author.

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