Vol. 6 (2013) *Acta Physica Polonica B Proceedings Supplement* No 3

GEOMETRICAL SCALING IN HIGH ENERGY COLLISIONS AND ITS BREAKING*

Michal Praszalowicz

The M. Smoluchowski Institute of Physics, Jagiellonian University
Reymonta 4, 30-059 Kraków, Poland

(Received July 11, 2013)

We analyze geometrical scaling (GS) in Deep Inelstic Scattering at HERA and in pp collisions at the LHC energies and in NA61/SHINE experiment. We argue that GS is working up to relatively large Bjorken $x \sim 0.1$. This allows to study GS in negative pion multiplicity $p_{\rm T}$ distributions at NA61/SHINE energies where clear sign of scaling violations is seen with growing rapidity when one of the colliding partons has Bjorken $x \geq 0.1$.

DOI:10.5506/APhysPolBSupp.6.809
PACS numbers: 13.85.Ni, 12.38.Lg

1. Introduction

In this short note, following Refs. [1–5] where also an extensive list of references can be found, we will focus on the scaling law, called geometrical scaling (GS), which has been introduced in the context of DIS [6]. Recently, it has been shown that GS is also exhibited by the $p_{\rm T}$ spectra at the LHC [1–3]. An onset of GS in heavy ion collisions at RHIC energies has been reported in Ref. [3]. At low Bjorken $x < x_{\rm max}$, proton is characterized by an intermediate energy scale $Q_{\rm s}(x)$ — called saturation scale [7, 8] — defined as the border line between dense and dilute gluonic systems within a proton (for review, see *e.g.* Refs. [9, 10]). For the present study, however, the details of saturation are not of primary interest, it is the very existence of $Q_{\rm s}(x)$ which is of importance.

Here, we present analysis of three different pieces of data which exhibit both emergence and violation of geometrical scaling. In Sect. 2 we briefly describe the method used to assess the existence of GS. Secondly, in Sect. 3 we describe our recent analysis [4] of combined HERA data [11] where it has

* Presented at the Workshop "Excited QCD 2013", Bjelašnica Mountain, Sarajevo, Bosnia–Herzegovina, February 3–9, 2013.

been shown that GS in DIS works very well up to relatively large $x_{\max} \sim 0.1$ (see also [12]). Next, in Sect. 4, on the example of the CMS $p_{\rm T}$ spectra in central rapidity [13], we show that GS can be extended to hadronic collisions. For particles produced at non-zero rapidities, one (larger) Bjorken $x = x_1$ may leave the domain of GS, *i.e.* $x_1 > x_{\max}$, and violation of GS should appear. In Sect. 5 we present analysis of very recent *pp* data from NA61/SHINE experiment at CERN [14] and show that GS is indeed violated once rapidity is increased. We conclude in Sect. 6.

2. Method of ratios

Geometrical scaling hypothesis means that some observable σ that, in principle, depends on two independent kinematical variables, say x and Q^2, in fact, depends only on a specific combination of them denoted as τ

$$\sigma\left(x, Q^2\right) = F(\tau)/Q_0^2 \,. \tag{1}$$

Here, function F in Eq. (1) is a dimensionless function of scaling variable

$$\tau = Q^2/Q_{\rm s}^2(x) \tag{2}$$

and

$$Q_{\rm s}^2(x) = Q_0^2 \left(x/x_0\right)^{-\lambda} \tag{3}$$

is the saturation scale. Here, Q_0 and x_0 are free parameters which can be extracted from the data within some specific model for σ, and exponent λ is a dynamical quantity of the order of $\lambda \sim 0.3$. Throughout this paper, we shall test the hypothesis whether different pieces of data can be described by formula (1) with *constant* λ, and what is the kinematical range where GS is working satisfactorily.

In view of Eq. (1), observables $\sigma(x_i, Q^2)$ for different x_i's should fall on one universal curve, if evaluated not in terms of Q^2 but in terms of τ. This means, in turn, that ratios

$$R_{x_i,x_{\rm ref}}(\lambda;\tau_k) = \frac{\sigma\left(x_i, \tau\left(x_i, Q_k^2;\lambda\right)\right)}{\sigma\left(x_{\rm ref}, \tau\left(x_{\rm ref}, Q_{k,{\rm ref}}^2;\lambda\right)\right)} \tag{4}$$

should be equal to unity independently of τ. Here, for some $x_{\rm ref}$, we pick up all $x_i < x_{\rm ref}$ which have at least two overlapping points in Q^2.

For $\lambda \neq 0$, points of the same Q^2 but different x's correspond, in general, to different τ's. Therefore, one has to interpolate $\sigma(x_{\rm ref}, \tau(x_{\rm ref}, Q^2;\lambda))$ to $Q_{k,{\rm ref}}^2$ such that $\tau(x_{\rm ref}, Q_{k,{\rm ref}}^2;\lambda) = \tau_k$. This procedure is described in detail in Refs. [4].

By tuning λ, one can make $R_{x_i,x_{\rm ref}}(\lambda;\tau_k) \to 1$ for all τ_k. In order to find optimal value $\lambda_{\min}$ that minimizes deviations of ratios (4) from unity, we form the chi-square measure

$$\chi^2_{x_i,x_{\rm ref}}(\lambda) = \frac{1}{N_{x_i,x_{\rm ref}} - 1} \sum_{k \in x_i} \frac{(R_{x_i,x_{\rm ref}}(\lambda;\tau_k) - 1)^2}{\Delta R_{x_i,x_{\rm ref}}(\lambda;\tau_k)^2}\,, \tag{5}$$

where the sum over k extends over all points of given x_i that have overlap with $x_{\rm ref}$, and $N_{x_i,x_{\rm ref}}$ is a number of such points.

3. Deep Inelastic Scattering at HERA

In the case of DIS, the relevant scaling observable is $\gamma^* p$ cross section and variable x is simply Bjorken x. In Fig. 1 we present 3d plot of $\lambda_{\min}(x, x_{\rm ref})$ which has been found by minimizing (5).

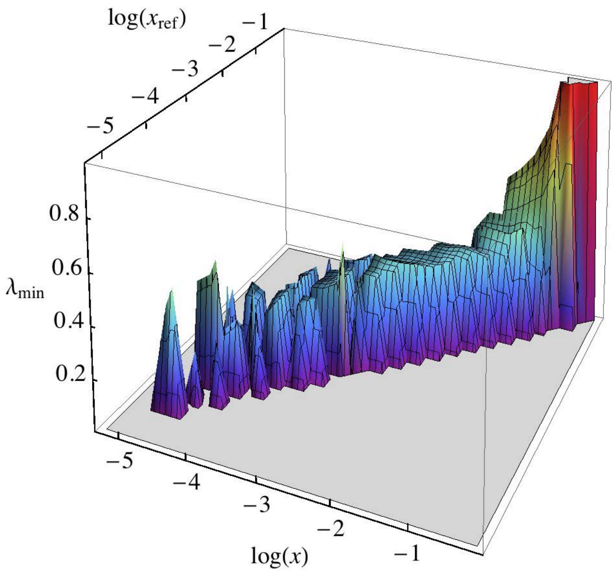

Fig. 1. Three dimensional plot of $\lambda_{\min}(x, x_{\rm ref})$ obtained by minimization of Eq. (5).

Qualitatively, GS is given by the independence of $\lambda_{\min}$ on Bjorken x and by the requirement that the pertinent value of $\chi^2_{x,x_{\rm ref}}(\lambda_{\min})$ should be small (for the discussion of the latter, see Refs. [4]). We see from Fig. 1 that the stability corner of $\lambda_{\min}$ extends up to $x_{\rm ref} \lesssim 0.1$, which is well above the original expectations. In Refs. [4] we have shown that

$$\lambda = 0.32 - 0.34 \qquad \text{for} \qquad x \leq 0.08\,. \tag{6}$$

4. Central rapidity $p_{\rm T}$ spectra at the LHC

In hadronic collisions at c.m. energy $W = \sqrt{s}$ particles are produced in the scattering process of two patrons carrying Bjorken x's

$$x_{1,2} = e^{\pm y}\, p_{\rm T}/W\,. \tag{7}$$

For central rapidities, $x = x_1 \sim x_2$. It has been shown that in this case charged particle multiplicity spectra exhibit GS [1]

$$\left.\frac{dN}{dy d^2 p_{\rm T}}\right|_{y\simeq 0} = \frac{1}{Q_0^2} F(\tau)\,, \tag{8}$$

where F is a universal dimensionless function of the scaling variable

$$\tau = p_{\rm T}^2/Q_{\rm s}^2(x) = p_{\rm T}^2/Q_0^2\, \left(p_{\rm T}/\left(x_0\sqrt{s}\right)\right)^{\lambda}\,. \tag{9}$$

Therefore, the scaling observable is $\sigma(W, p_{\rm T}^2) = dN/dyd^2p_{\rm T}$ and the method of ratios is applied to the multiplicity distributions at different energies (W_i taking over the role of x_i in Eq. (4)). For $W_{\rm ref}$, we take the highest LHC energy of 7 TeV. Therefore, one can form two ratios $R_{W_i,W_{\rm ref}}$ with $W_1 = 2.36$ and $W_2 = 0.9$ TeV. These ratios are plotted in Fig. 2 for the CMS single non-diffractive spectra for $\lambda = 0$ and for $\lambda = 0.27$, which minimizes (5) in this case. We see that original ratios plotted in terms of $p_{\rm T}$ range from 1.5 to 7, whereas plotted in terms of $\sqrt{\tau}$ they are well concentrated around unity. The optimal exponent λ is, however, smaller than in the case of DIS. Why this is so, remains to be understood.

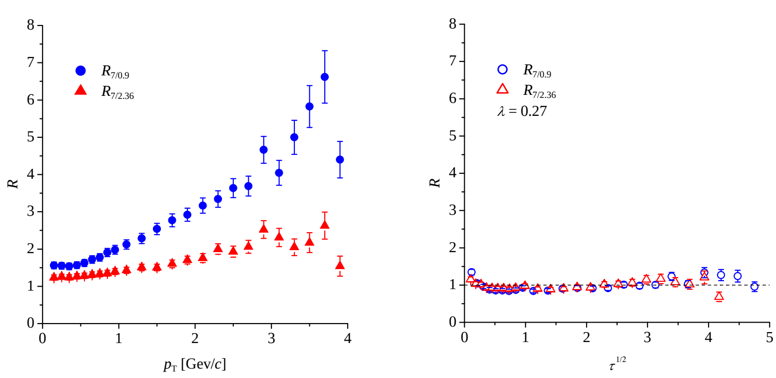

Fig. 2. Ratios of CMS $p_{\rm T}$ spectra [13] at 7 TeV to 0.9 (blue circles) and 2.36 TeV (red triangles) plotted as functions of $p_{\rm T}$ (left) and scaling variable $\sqrt{\tau}$ (right) for $\lambda = 0.27$.

5. Violation of geometrical scaling in forward rapidity region

For $y > 0$, two Bjorken x's can be quite different: $x_1 > x_2$. Therefore, looking at the spectra with increasing y one can eventually reach $x_1 > x_{\rm max}$ and GS violation should be seen. To this end, we shall use *pp* data from NA61/SHINE experiment at CERN [14] at different rapidities $y = 0.1$–3.5 and at five different energies $W_{1,\ldots,5} = 17.28, 12.36, 8.77, 7.75$, and 6.28 GeV.

In Fig. 3 we plot ratios $R_{1i} = R_{W_1,W_i}$ (4) for π^- spectra in central rapidity for $\lambda = 0$ and 0.27. For $y = 0.1$, the GS region extends towards the smallest energy because $x_{\rm max}$ is as large as 0.08. However, the quality of GS is the worst for the lowest energy W_5. By increasing y, some points fall outside the GS window because $x_1 \geq x_{\rm max}$, and finally for $y \geq 1.7$ no GS should be present in NA61/SHINE data. This is illustrated nicely in Fig. 4.

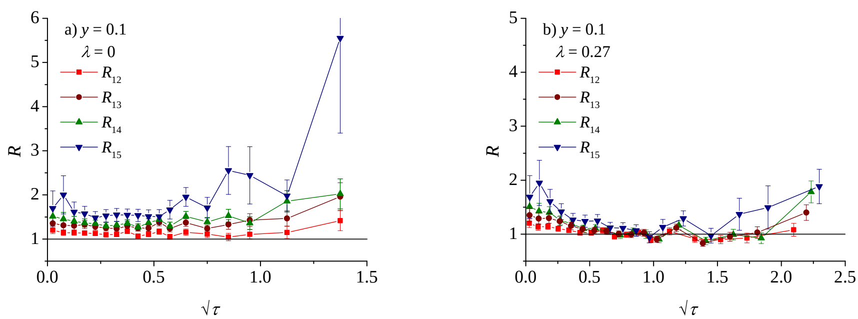

Fig. 3. Ratios R_{1k} as functions of $\sqrt{\tau}$ for the lowest rapidity $y = 0.1$: (a) for $\lambda = 0$ when $\sqrt{\tau} = p_{\rm T}$ and (b) for $\lambda = 0.27$ which corresponds to GS.

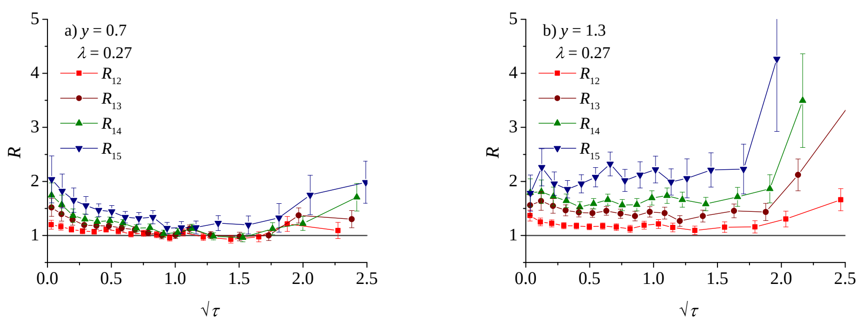

Fig. 4. Ratios R_{1k} as functions of $\sqrt{\tau}$ for $\lambda = 0.27$ and for different rapidities (a) $y = 0.7$ and (b) $y = 1.3$. With an increase of rapidity, gradual closure of the GS window can be seen.

6. Conclusions

We have shown that GS in DIS works well up to rather large Bjorken x's with exponent $\lambda = 0.32$–0.34. In *pp* collisions at the LHC energies in central rapidity GS is seen in the charged particle multiplicity spectra, however, $\lambda = 0.27$ in this case. By changing rapidity, one can force one of the Bjorken x's of colliding patrons to exceed $x_{\rm max}$ and GS violation is expected. Such behavior is indeed observed in the NA61/SHINE *pp* data.

The author wants to thank M. Gaździcki and Sz. Puławski for the access to the NA61/SHINE data and to T. Stebel for collaboration and remarks. Many thanks are due to the Organizers of this successful series of conferences. This work was supported by the Polish National Science Center (NCN) grant 2011/01/B/ST2/00492.

REFERENCES

[1] L. McLerran, M. Praszalowicz, *Acta Phys. Pol. B* **41**, 1917 (2010); *Acta Phys. Pol. B* **42**, 99 (2011).
[2] M. Praszalowicz, *Phys. Rev. Lett.* **106**, 142002 (2011).
[3] M. Praszalowicz, *Acta Phys. Pol. B* **42**, 1557 (2011); arXiv:1205.4538 [hep-ph].
[4] M. Praszalowicz, T. Stebel, *J. High Energy Phys.* **1303**, 090 (2013); *J. High Energy Phys.* **1304**, 169 (2013) [arXiv:1302.4227 [hep-ph]].
[5] M. Praszalowicz, *Phys. Rev.* **D87**, 071502(R) (2013) [arXiv:1301.4647 [hep-ph]].
[6] A.M. Stasto, K.J. Golec-Biernat, J. Kwiecinski, *Phys. Rev. Lett.* **86**, 596 (2001).
[7] L.V. Gribov, E.M. Levin, M.G. Ryskin, *Phys. Rep.* **100**, 1 (1983); A.H. Mueller, J.-W. Qiu, *Nucl. Phys.* **268**, 427 (1986); A.H. Mueller, *Nucl. Phys.* **B558**, 285 (1999).
[8] K.J. Golec-Biernat, M. Wüsthoff, *Phys. Rev. D* **59**, 014017 (1998).
[9] A.H. Mueller, arXiv:hep-ph/0111244.
[10] L. McLerran, *Acta Phys. Pol. B* **41**, 2799 (2010).
[11] C. Adloff *et al.* [H1 Collaboration], *Eur. Phys. J.* **C21**, 33 (2001); S. Chekanov *et al.* [ZEUS Collaboration], *Eur. Phys. J.* **C21**, 443 (2001).
[12] F. Caola, S. Forte, J. Rojo, *Nucl. Phys.* **A854**, 32 (2011).
[13] V. Khachatryan *et al.* [CMS Collaboration], *J. High Energy Phys.* **1002**, 041 (2010); *Phys. Rev. Lett.* **105**, 022002 (2010); *J. High Energy Phys.* **1101**, 079 (2011).
[14] N. Abgrall *et al.* [NA61/SHINE Collaboration], Report from the NA61/SHINE experiment at the CERN SPS, CERN-SPSC-2012-029, SPSC-SR-107; A. Aduszkiewicz, Ph.D. Thesis, University of Warsaw, 2013, in prepartation; Sz. Pulawski, talk at 9th Polish Workshop on Relativistic Heavy-Ion Collisions, Kraków, November 2012 and private communication.