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Subdifferential inclusions and quasi-static hemivariational inequalities for frictional viscoelastic contact problems

Review Article

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Abstract: We survey recent results on the mathematical modeling of nonconvex and nonsmooth contact problems arising in mechanics and engineering. The approach to such problems is based on the notions of an operator subdifferential inclusion and a hemivariational inequality, and focuses on three aspects. First we report on results on the existence and uniqueness of solutions to subdifferential inclusions. Then we discuss two classes of quasi-static hemivariational ineqaulities, and finally, we present ideas leading to inequality problems with multivalued and nonmonotone boundary conditions encountered in mechanics.

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1. Introduction

The paper is a shortened version of the original parts of an invited talk presented by the author at the Sixth Symposium on Nonlinear Analysis (SNA2011) held in Toruń, Poland, September 7–9, 2011. The goal is to review and report on the results of recent studies on problems involving operator subdifferential inclusions in Banach spaces, quasi-static hemivariational inequalities, and their applications.

The notion of hemivariational inequality was introduced by P.D. Panagiotopoulos in the early 1980s, cf. [17, 18], as a generalization of the variational inequality. The hemivariational inequality formulation exploits the notion of the generalized gradient of Clarke–Rockafellar, cf. [1]. The latter was introduced for a class of locally Lipschitz functions and

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allows to give variational formulations of several mechanical phenomena involving the nonconvex and nondifferentiable energy functionals. Thus the theory of hemivariational inequalities provides mathematical results which arise and are applicable to variational theory of engineering problems involving nonmonotone multivalued relations. In the stationary case the hemivariational inequalities can be formulated as the substationary point problems for the corresponding non-differentiable, nonsmooth and nonconvex energy functionals. The formulations of mechanical problems as hemivariational inequalities allow to give positive answers to unsolved or partially unsolved problems, cf. [9–12, 15, 16]. More information on modeling and analysis of contact phenomena can be found in [5, 7, 8, 20].

The goal of this paper is to give a survey on modeling of contact problems with nonsmooth potentials and on analysis of quasi-static hemivariational inequalities. We have tried to keep the presentation as simple as possible and this makes it accessible also to the newcomer in this field.

The content of the paper is as follows. We consider two classes of subdifferential inclusions in a framework of evolution triple of spaces. The first class involves inclusions with a history-dependent term for which we provide an existence and uniqueness result. The proof is based on arguments on pseudomonotone operators and fixed point. Then we specialize this result in the study of a class of history-dependent hemivariational inequalities. Such kind of problems arises in a large number of mathematical models which describe quasistatic processes of contact between a deformable body and an obstacle. To provide an example we consider a viscoelastic problem in which the frictional contact is modeled with subdifferential boundary conditions. We prove that this contact problem leads to a history-dependent hemivariational inequality in which the unknown is the velocity field. Then we apply our result in order to prove the unique weak solvability of the corresponding contact problem. In the second class of inclusions, we consider time-dependent possibly nonconvex nonsmooth functions and their Clarke subdifferentials operating on the unknown function. First we prove the existence of a weak solution. Then we study the asymptotic behavior of a sequence of solutions when a small parameter in the inertial term tends to zero. We prove that the limit function is a solution of a parabolic hemivariational inequality. Finally, we give applications to quasi-static viscoelastic frictional contact problems.

2. Preliminaries

In this section we recall some definitions we shall use in this paper, see e.g. Zeidler [21] and Denkowski et al. [2, 3]. Given a Banach space X, we denote its norm by $\|\cdot\|_X$. The dual space is denoted by X^* and $\langle \cdot, \cdot \rangle_{X^* \times X}$, denoted also by $\langle \cdot, \cdot \rangle$, is the duality pairing between X and X^* . The symbol w-X is used for the space X endowed with its weak topology. The notation $\mathcal{L}(X, Y)$ stands for the space of linear bounded operators from a Banach space X to a Banach space Y. Also, we denote by 2^X the collection of all subsets of X. For $U \subset X$, we also write $\|U\|_X = \sup\{\|u\|_X : u \in U\}$. Let $\theta: X \to \mathbb{R}$ be a locally Lipschitz function. The generalized directional derivative of θ at $x \in X$ in the direction $v \in X$ is defined by

$$\theta^{0}(x; v) = \limsup_{y \to x, \lambda \downarrow 0} \frac{\theta(y + \lambda v) - \theta(y)}{\lambda}$$

The generalized gradient of θ at x is the subset of X^* given by

$$\partial \theta(x) = \{ \zeta \in X^* : \theta^0(x; v) \ge \langle \zeta, v \rangle, v \in X \}.$$

A locally Lipschitz function θ is called regular (in the sense of Clarke, cf. [1]) at $x \in X$ if for all $v \in X$ the one-sided directional derivative $\theta'(x; v)$ exists and satisfies $\theta^0(x; v) = \theta'(x; v)$ for all $v \in X$.

- Let $T: X \to 2^{X^*}$. An operator T is said to be *pseudomonotone* if it satisfies the following conditions:
- (a) for every $x \in X$, Tx is a nonempty, convex and weakly compact set in X^* ;
- (b) T is u.s.c. on every finite dimensional subspace of X into X^* endowed with the weak topology; and
- (c) if $x_n \to x$ weakly in X, $x_n^* \in Tx_n$ and $\limsup \langle x_n^*, x_n x \rangle \leq 0$, then for each $z \in X$ there exists $x^*(z) \in Tx$ such that $\langle x^*(z), x z \rangle \leq \liminf \langle x_n^*, x_n z \rangle$.

Let $L: D(L) \subset X \to X^*$ be a linear densely defined maximal monotone operator. An operator T is said to be *pseu*domonotone with respect to D(L) (shortly *L-pseudomonotone*) if and only if (a) and (b) hold and

(d) if $\{x_n\} \subset D(L)$ is such that $x_n \to x$ weakly in $X, x \in D(L)$, $Lx_n \to Lx$ weakly in $X^*, x_n^* \in Tx_n, x_n^* \to x^*$ weakly in X^* and $\limsup \langle x_n^*, x_n \rangle \leq \langle x^*, x \rangle$, then $x^* \in Tx$ and $\langle x_n^*, x_n \rangle \to \langle x^*, x \rangle$.

An operator T is said to be *coercive* if there exists a function $c: \mathbb{R}^+ \to \mathbb{R}$ with $c(r) \to \infty$ as $r \to \infty$ such that $\langle x^*, x \rangle \ge c(\|x\|_X) \|x\|_X$ for every $(x, x^*) \in \text{Graph}(T)$.

A single-valued operator $T: X \to X^*$ is said to be *hemicontinuous*, if for all $u, v, w \in X$ the functional $t \mapsto \langle T(u+tv), w \rangle$ is continuous on [0, 1]. $T: X \to X^*$ is *pseudomonotone* if for each sequence $\{x_n\} \subseteq X$ such that it converges weakly to $x_0 \in X$ and $\limsup \langle Tx_n, x_n - x_0 \rangle \leq 0$, we have $\langle Tx_0, x_0 - x \rangle \leq \liminf \langle Tx_n, x_n - x \rangle$ for all $x \in X$.

Finally, we recall the following surjectivity result, cf. [2, Theorem 1.3.73], for multivalued operators which are pseudomonotone with respect to D(L).

Proposition 2.1.

If X is a reflexive, strictly convex Banach space, $L: D(L) \subset X \to X^*$ is a linear densely defined maximal monotone operator and $T: X \to 2^{X^*} \setminus \{\emptyset\}$ is bounded, coercive and pseudomonotone with respect to D(L), then the operator L + T is surjective.

3. History-dependent subdifferential inclusions

Let $\Omega \subset \mathbb{R}^d$ be an open bounded subset of \mathbb{R}^d with a Lipschitz continuous boundary $\partial\Omega$ and $\Gamma \subseteq \partial\Omega$, d = 1, 2, 3. Let V be a closed subspace of $H^1(\Omega; \mathbb{R}^s)$, $s \ge 1$, $H = L^2(\Omega; \mathbb{R}^s)$ and $Z = H^{\delta}(\Omega; \mathbb{R}^s)$ with a fixed $\delta \in (1/2, 1)$. Denoting by $i: V \to Z$ the embedding, by $\gamma: Z \to L^2(\Gamma; \mathbb{R}^s)$ and $\gamma_0: H^1(\Omega; \mathbb{R}^s) \to H^{1/2}(\Gamma; \mathbb{R}^s) \subset L^2(\Gamma; \mathbb{R}^s)$ the trace operators, we get $\gamma_0 v = \gamma(iv)$ for all $v \in V$. For simplicity, in what follows, we omit the notation of the embedding i and we write $\gamma_0 v = \gamma v$ for all $v \in V$. It is well known from the theory of Sobolev spaces, cf. [2, 3, 21], that (V, H, V^*) and (Z, H, Z^*) form evolution triples of spaces and the embedding $V \subset Z$ is compact. We denote by c_e the embedding constant of V into Z, by $\|\gamma\|$ the norm of the trace in $\mathcal{L}(Z, L^2(\Gamma; \mathbb{R}^s))$ and by $\gamma^*: L^2(\Gamma; \mathbb{R}^s) \to Z^*$ the adjoint operator to γ . We also introduce the following spaces:

$$\mathcal{V} = L^2(0, T; V), \qquad \mathcal{Z} = L^2(0, T; Z) \quad \text{and} \quad \widehat{\mathcal{H}} = L^2(0, T; H),$$

where $0 < T < +\infty$. Since the embeddings $V \subseteq Z \subseteq H \subseteq Z^* \subseteq V^*$ are continuous, it is known that the embeddings $\mathcal{V} \subseteq \mathcal{Z} \subseteq \widehat{\mathcal{H}} \subseteq \mathcal{Z}^* \subseteq \mathcal{V}^*$ are also continuous, where $\mathcal{Z}^* = L^2(0, T; Z^*)$ and $\mathcal{V}^* = L^2(0, T; V^*)$.

Let $A: (0, T) \times V \to V^*$, $S: \mathcal{V} \to \mathcal{V}^*$, $J: (0, T) \times L^2(\Gamma; \mathbb{R}^s) \to \mathbb{R}$ and $f: (0, T) \to V^*$ be given. We consider the following inclusion of subdifferential form.

Problem 3.1.

Find $u \in \mathcal{V}$ such that

$$A(t, u(t)) + Su(t) + \gamma^* \partial J(t, \gamma u(t)) \supseteq f(t) \qquad \text{for a.e.} \quad t \in (0, T).$$

$$\tag{1}$$

We note that in (1) the notation Su(t) stands for (Su)(t), i.e. Su(t) = (Su)(t) for all $u \in \mathcal{V}$ and a.e. $t \in (0, T)$. The symbol ∂J denotes the Clarke subdifferential of a locally Lipschitz function $J(t, \cdot)$. We adopt the following definition.

Definition 3.2.

A function $u \in \mathcal{V}$ is called a *solution to Problem 3.1* if and only if there exists $\zeta \in \mathcal{I}^*$ such that for a.e. $t \in (0, T)$,

$$A(t, u(t)) + Su(t) + \zeta(t) = f(t)$$
 and $\zeta(t) \in \gamma^* \partial J(t, \gamma u(t)).$

In order to state a result on the solvability of Problem 3.1, we need the following hypotheses.

- I. A: $(0, T) \times V \rightarrow V^*$ is such that
 - (I.a) $A(\cdot, v)$ is measurable on (0, T) for all $v \in V$;
 - (I.b) $A(t, \cdot)$ is hemicontinuous and strongly monotone for a.e. $t \in (0, T)$, i.e., $\langle A(t, v_1) A(t, v_2), v_1 v_2 \rangle_{V^* \times V} \ge m_1 \|v_1 v_2\|_V^2$ for all $v_1, v_2 \in V$ with $m_1 > 0$;
 - (I.c) $||A(t, v)||_{V^*} \le a_0(t) + a_1 ||v||_V$ for all $v \in V$, a.e. $t \in (0, T)$ with $a_0 \in L^2(0, T)$, $a_0 \ge 0$ and $a_1 > 0$;
 - (I.d) A(t, 0) = 0 for a.e. $t \in (0, T)$.
- II. $S: \mathcal{V} \to \mathcal{V}^*$ is such that

$$\|Su_1(t) - Su_2(t)\|_{V^*} \le L_S \int_0^t \|u_1(s) - u_2(s)\|_V ds$$

for all u_1 , $u_2 \in \mathcal{V}$, a.e. $t \in (0, T)$ with $L_S > 0$.

- III. $J: (0, T) \times L^2(\Gamma; \mathbb{R}^s) \to \mathbb{R}$ is such that
 - (III.a) $J(\cdot, u)$ is measurable on (0, T) for all $u \in L^2(\Gamma; \mathbb{R}^s)$
 - (III.b) $J(t, \cdot)$ is locally Lipschitz on $L^2(\Gamma; \mathbb{R}^s)$ for a.e. $t \in (0, T)$;
 - $(III.c) \ \|\partial J(t, u)\|_{L^{2}(\Gamma; \mathbb{R}^{s})} \leq c_{0} + c_{1} \|u\|_{L^{2}(\Gamma; \mathbb{R}^{s})} \text{ for all } u \in L^{2}(\Gamma; \mathbb{R}^{s}), \text{ a.e. } t \in (0, T) \text{ with } c_{0}, c_{1} \geq 0;$
 - (III.d) $(z_1 z_2, u_1 u_2)_{L^2(\Gamma; \mathbb{R}^s)} \ge -m_2 \|u_1 u_2\|_{L^2(\Gamma; \mathbb{R}^s)}^2$ for all $z_i \in \partial J(t, u_i)$, $u_i, z_i \in L^2(\Gamma; \mathbb{R}^s)$, i = 1, 2, a.e. $t \in (0, T)$, $m_2 \ge 0$;
 - (III.e) $J^{0}(t, u; -u) \leq d_{0}(1 + ||u||_{L^{2}(\Gamma; \mathbb{R}^{s})})$ for all $u \in L^{2}(\Gamma; \mathbb{R}^{s})$, a.e. $t \in (0, T)$ with $d_{0} \geq 0$.

We remark that condition II is satisfied for the operator $S: \mathcal{V} \to \mathcal{V}^*$ given by

$$Sv(t) = R\left(t, \int_0^t v(s) \, ds + v_0\right) \qquad \text{for all} \quad v \in \mathcal{V}, \qquad \text{a.e.} \quad t \in (0, T), \tag{2}$$

where $R: (0, T) \times V \to V^*$ is such that $R(\cdot, v)$ is measurable on (0, T) for all $v \in V$, $R(t, \cdot)$ is a Lipschitz continuous operator for a.e. $t \in (0, T)$ and $v_0 \in V$. It is also satisfied for the *Volterra operator* $S: \mathcal{V} \to \mathcal{V}^*$ given by

$$Sv(t) = \int_0^t C(t-s)v(s) \, ds \qquad \text{for all} \quad v \in \mathcal{V}, \qquad \text{a.e.} \quad t \in (0, T), \tag{3}$$

where $C \in L^{\infty}(0, T; \mathcal{L}(V, V^*))$. Clearly, in the case of the operators (2) and (3) the current value $\delta v(t)$ at the moment t depends on the history of the values of v at the moments $0 \leq s \leq t$ and, therefore, we refer to the operators of the form (2) or (3) as *history-dependent*. We extend this definition to all the operators $\mathcal{S}: \mathcal{V} \to \mathcal{V}^*$ which satisfy condition II and, for this reason, we say that the subdifferential inclusions of the form (1) are *history-dependent subdifferential inclusions*. The main feature of such inclusions consists in the fact that they contain operators which, at any moment $t \in (0, T)$ depend on the *history* of the solution up to the moment t. This feature makes the difference with respect to the time-dependent subdifferential inclusions studied in literature in which, usually, the operators involved are assumed to depend on the *current* value of the solution u(t).

In order to prove the existence and uniqueness for Problem 3.1, we first state a result on the unique solvability of a subdifferential inclusion in which the time variable plays the role of a parameter.

Lemma 3.3.

Assume that I holds and $f \in \mathcal{V}^*$. If one of the following hypotheses:

- (i) (III.a)–(III.d) and $m_1 > \max\{c_1, m_2\}c_e^2 \|\gamma\|^2$,
- (ii) III and $m_1 > m_2 c_e^2 \| \boldsymbol{\gamma} \|^2$

is satisfied, then the problem

$$A(t, u(t)) + \gamma^* \partial J(t, \gamma u(t)) \ni f(t) \qquad \text{for a.e.} \quad t \in (0, T)$$
(4)

has a unique solution $u \in \mathcal{V}$.

Proof. We provide main steps of the proof without details. First, since the operator $A(t, \cdot)$ is strongly monotone and A(t, 0) = 0 for a.e. $t \in (0, T)$, it follows that $A(t, \cdot)$ is coercive with $m_1 > 0$, i.e. $\langle A(t, v), v \rangle_{V^* \times V} \ge m_1 \|v\|_V^2$ for all $v \in V$, a.e. $t \in (0, T)$. Moreover, since the operator $A(t, \cdot)$ satisfies (1.b)–(1.c), it is pseudomonotone for a.e. $t \in (0, T)$. This follows from the facts that every strongly monotone operator is monotone and every bounded, hemicontinuous and monotone operator is pseudomonotone, cf. [21, Proposition 27.6]. Next, we define the operator $B: (0, T) \times V \to 2^{Z^*}$ by $B(t, v) = \gamma^* \partial J(t, \gamma v)$ for all $v \in V$, a.e. $t \in (0, T)$. Under either the hypothesis (i) or (ii), we can establish the following properties of the operator B.

- (IV.a) $B(\cdot, v)$ is measurable for all $v \in V$;
- (IV.b) $||B(t, v)||_{Z^*} \le b_0(1 + ||v||_V)$ for all $v \in V$, a.e. $t \in (0, T)$ with $b_0 > 0$;
- (IV.c) for all $v \in V$ and a.e. $t \in (0, T)$, B(t, v) is nonempty, convex, weakly compact subset of Z^* ;
- (IV.d) $\langle B(t,v), v \rangle_{V^* \times V} \ge -b_1 \|v\|_V^2 b_2 \|v\|_V b_3$ for all $v \in V$, a.e. $t \in (0, T)$ with $b_1, b_2, b_3 \ge 0$;
- (IV.e) the graph of $B(t, \cdot)$ is closed in $Z \times (w Z^*)$ topology for a.e. $t \in (0, T)$, i.e. if $\zeta_n(t) \in B(t, v_n)$ with $v_n, v \in V$, $v_n \rightarrow v$ in Z and $\zeta_n(t), \zeta(t) \in Z^*, \zeta_n(t) \rightarrow \zeta(t)$ weakly in Z^* , then $\zeta(t) \in B(t, v)$.

Subsequently, we define the multivalued map $\mathcal{F}: (0, T) \times V \to 2^{V^*}$ by $\mathcal{F}(t, v) = A(t, v) + B(t, v)$ for all $v \in V$ and a.e. $t \in (0, T)$. From (I.a) and (IV.a), it is clear that $\mathcal{F}(\cdot, v)$ is a measurable multifunction for all $v \in V$. Exploiting [3, Proposition 6.3.66], we show that $\mathcal{F}(t, \cdot)$ is pseudomonotone and coercive for a.e. $t \in (0, T)$. Therefore, applying the fundamental surjectivity result, cf. e.g. [3, Theorem 6.3.70], it follows that $\mathcal{F}(t, \cdot)$ is surjective. This implies that for a.e. $t \in (0, T)$ there exists a solution $u(t) \in V$ of the problem (4). Furthermore, owing to the coercivity of $\mathcal{F}(t, \cdot)$, we deduce the following estimate:

$$\|u(t)\|_{V} \le c(1 + \|f(t)\|_{V^{*}}) \quad \text{for a.e.} \quad t \in (0, T) \quad \text{with} \quad c > 0.$$
(5)

Using the strong monotonicity of $A(t, \cdot)$, (III.d) and the hypothesis $m_1 > m_2 c_e^2 \|\gamma\|^2$, we prove now that the solution to the problem (4) is unique. Also, we are able to prove that the solution of the problem (4) is a measurable function on (0, *T*). Since $f \in \mathcal{V}^*$, from the estimate (5), we conclude that $u \in \mathcal{V}$ and (4) hold, which completes the proof of the lemma.

The existence and uniqueness result for Problem 3.1 reads as follows.

Theorem 3.4.

Assume I, II and $f \in \mathcal{V}^*$. If either (i) or (ii) of the hypothesis of Lemma 3.3 holds, then Problem 3.1 has a unique solution.

Proof. We use a fixed point argument. Let $\eta \in \mathcal{V}^*$. We denote by $u_\eta \in \mathcal{V}$ the solution of the following problem:

$$A(t, u_n(t)) + \gamma^* \partial J(t, \gamma u_n(t)) \ge f(t) - \eta(t) \quad \text{for a.e.} \quad t \in (0, T).$$
(6)

By Lemma 3.3 we know that $u_{\eta} \in \mathcal{V}$ exists and is unique. Next, we consider the operator $\Lambda: \mathcal{V}^* \to \mathcal{V}^*$ defined by

$$\Lambda \eta(t) = Su_n(t)$$
 for all $\eta \in \mathcal{V}^*$, a.e. $t \in (0, T)$.

We show by using the Banach contraction principle that the operator Λ has a unique fixed point $\eta^* \in \mathcal{V}^*$. Then u_{η^*} is a solution to Problem 3.1, which concludes the existence part of the theorem. The uniqueness part follows from the uniqueness of the fixed point of Λ . Namely, let $u \in \mathcal{V}$ be a solution to Problem 3.1 and define the element $\eta \in \mathcal{V}^*$ by $\eta(t) = Su(t)$ for a.e. $t \in (0, T)$. It follows that u is a solution to the problem (6) and, by the uniqueness of solutions to (6), we obtain $u = u_{\eta}$. This implies $\Lambda \eta = Su_{\eta} = Su = \eta$ and by the uniqueness of the fixed point of Λ we have $\eta = \eta^*$, so $u = u_{\eta^*}$, which completes the proof.

4. Quasistatic hemivariational inequalities

Let *V* and *Z* be reflexive, separable Banach spaces, and let *H* be a Hilbert space. Suppose that $V \,\subset Z \,\subset H \approx H^* \,\subset Z^* \,\subset V^*$, where H^*, Z^* and V^* denote dual spaces to *H*, *Z* and *V*, respectively, all embeddings are continuous and *V* is compactly embedded in *Z*. We denote by $\langle \cdot, \cdot \rangle$ the duality of *V* and *V**, and the pairing between *Z* and *Z** as well. Given a fixed number $0 < T < \infty$ and $2 \leq p < \infty$, we introduce the following spaces: $\mathcal{V} = L^p(0, T; V)$, $\mathcal{Z} = L^p(0, T; Z)$, $\widehat{\mathcal{H}} = L^2(0, T; H)$, $\mathcal{Z}^* = L^q(0, T; Z^*)$, $\mathcal{V}^* = L^q(0, T; V^*)$ with 1/p + 1/q = 1 and $\mathcal{W} = \{w \in \mathcal{V} : w' \in \mathcal{V}^*\}$, where the time derivative is understood in the sense of vector-valued distributions. The latter is a separable, reflexive Banach space with the norm $\|w\|_{\mathcal{W}} = \|w\|_{\mathcal{V}} + \|w'\|_{\mathcal{V}^*}$. We have $\mathcal{W} \subset \mathcal{V} \subset \mathcal{Z} \subset \widehat{\mathcal{H}} \subset \mathcal{Z}^* \subset \mathcal{V}^*$ with continuous embeddings. It is well known, cf. e.g. [3, Proposition 8.4.14], that the space \mathcal{W} is embedded continuously in C(0, T; H) (the space of continuous functions on [0, T] with values in *H*), i.e. every element of \mathcal{W} , after a possible modification on a set of measure zero, has a unique continuous representative in C(0, T; H). Moreover, since *V* is embedded compactly in *Z*, then so is \mathcal{W} in \mathcal{Z} , cf. [3, Theorem 8.4.13].

Consider the following evolutionary inclusion of the form:

$$\begin{cases} A(t)u'(t) + Bu(t) + M^* \partial J(t, Mu(t)) \ni f(t) & \text{ for a.e. } t \in (0, T), \\ u(0) = u_0. \end{cases}$$
(7)

We remark that, by the definition of the Clarke subdifferential, problem (7) is equivalent to the following inequality:

$$\begin{cases} \left\langle A(t)u'(t) + Bu(t) - f(t), v \right\rangle + J^0(t, Mu(t); Mv) \ge 0 & \text{for all } v \in V, \quad \text{a.e. } t \in (0, T), \\ u(0) = u_0, \end{cases}$$

where J^0 stands for the generalized directional derivative of $J(t, \cdot)$. For this reason problem (7) is called a hemivariational inequality.

Definition 4.1.

A function $u \in L^{\infty}(0, T; V)$ is called a *solution to* (7) if and only if $u' \in \mathcal{V}$ and there exists $\eta \in \mathbb{Z}^*$ such that

$$\begin{cases} A(t)u'(t) + Bu(t) + \eta(t) = f(t) & \text{ for a.e. } t \in (0, T), \\ \eta(t) \in M^* \partial J(t, Mu(t)) & \text{ for a.e. } t \in (0, T), \\ u(0) = u_0. \end{cases}$$

We note that if u is a solution to (7), then $u \in L^{\infty}(0, T; V) \subset \mathcal{V}$, i.e. $u \in W^{1,p}(0, T; V)$. Since $W^{1,p}(0, T; V) \subset C(0, T; V)$ continuously, the initial condition u(0) has a meaning in V.

In addition to our notation, let X be a Banach space. We assume the following hypotheses.

- H(A) $A \in L^{\infty}(0, T; \mathcal{L}(V, V^*))$ is an operator such that A(t) is coercive, i.e. there is a constant $\alpha > 0$ such that for a.e. $t \in (0, T), \langle A(t)v, v \rangle \ge \alpha \|v\|_{V}^{p}$ for all $v \in V$.
- H(B) $B \in \mathcal{L}(V, V^*)$ is monotone (nonnegative) and symmetric.
- H(J) $J: (0, T) \times X \rightarrow \mathbb{R}$ is a function such that
 - (i) $J(\cdot, x)$ is measurable on (0, T) for all $x \in X$;
 - (ii) $J(t, \cdot)$ is locally Lipschitz on X for a.e. $t \in (0, T)$;
 - (iii) $\|\partial J(t, x)\|_{X^*} \le c(1 + \|x\|_X^{2/q})$ for all $x \in X$, a.e. $t \in (0, T)$ with c > 0.

 $H(M) \quad M \in \mathcal{L}(Z, X).$

- $\mathsf{H}_0 \quad f \in \mathcal{V}^*, \, u_0 \in V, \, u_1 \in H.$
- H₁ If p = 2, then $\alpha > 2cTc_e^2 ||M|| \max\{1, ||M||\}$, where $c_e > 0$ is an embedding constant of V into Z and $||M|| = ||M||_{\mathcal{L}(Z,X)}$.

The following is the main existence result on the evolution inclusion (7).

Theorem 4.2.

Under hypotheses H(A), H(B), H(J), H(M), $f \in \mathcal{V}^*$, $u_0 \in V$ and H_1 , the inclusion (7) admits at least one solution.

The proof of Theorem 4.2 is based on the so-called vanishing acceleration method which we describe below. To this end, we consider an evolution second order inclusion of the form

$$\begin{cases} \varepsilon u''(t) + A(t)u'(t) + Bu(t) + M^* \partial J(t, Mu(t)) \ni f(t) & \text{for a.e. } t \in (0, T), \\ u(0) = u_0, & \sqrt{\varepsilon} u'(0) = u_1, \end{cases}$$
(8)

where $\varepsilon > 0$. For ε fixed, we write for simplicity *u* for the solution u_{ε} of (8).

Definition 4.3.

A function $u \in \mathcal{V}$ is called a *solution to* (8) if and only if $u' \in \mathcal{W}$ and there exists $\eta \in \mathcal{I}^*$ such that

$$\begin{cases} \varepsilon u''(t) + A(t)u'(t) + Bu(t) + \eta(t) = f(t) & \text{for a.e. } t \in (0, T), \\ \eta(t) \in M^* \partial J(t, Mu(t)) & \text{for a.e. } t \in (0, T), \\ u(0) = u_0, & \sqrt{\varepsilon} u'(0) = u_1. \end{cases}$$

Similarly, as before, we note that if u is a solution to (8), then $u \in W^{1,p}(0, T; V)$. Since the embeddings $W^{1,p}(0, T; V) \subset C(0, T; V)$ and $W \subset C(0, T; H)$ are continuous, the initial conditions u(0) and u'(0) have a meaning in V and H, respectively.

First, we comment on the existence result for (8). The important step is to derive the following uniform estimate.

Lemma 4.4.

Let $\varepsilon > 0$ be fixed. Assume hypotheses H(A), H(B), H(J), H(M) and H₀, and let u be a solution to (8). If p > 2, then there exists a constant C > 0 independent of ε such that

$$\|u\|_{C(0,T;V)} + \|u'\|_{\mathcal{V}} + \sqrt{\varepsilon} \|u'\|_{L^{\infty}(0,T;H)} + \varepsilon \|u''\|_{\mathcal{V}^*} \le C \left(1 + \|u_0\|_{V}^{2/q} + \|u_1\|_{H}^{2/q} + \|f\|_{\mathcal{V}^*}^{2/q}\right)$$

Moreover, this estimate still holds for p = 2 provided H_1 is satisfied.

From [13, Theorem 8], we obtain the following result.

Theorem 4.5.

If hypotheses H(A), H(B), H(J), H(M), H₀ and H₁ hold, then for every fixed $\varepsilon > 0$ the problem (8) admits at least one solution.

The following is the existence result for the operator inclusion (7).

Theorem 4.6.

Let the hypotheses H(A), H(B), H(J), H(M), H_0 and H_1 hold. For every ε fixed, let u_{ε} be a solution to (8) given by Theorem 4.5. Then there exists u fulfilling $u \in L^{\infty}(0, T; V)$, $u' \in W$, such that the following convergences hold:

$$\begin{split} u_{\varepsilon} &\to u \quad weakly^* \ in \ L^{\infty}(0, T; V), \qquad u'_{\varepsilon} \to u' \quad weakly \ in \ \mathcal{V}, \\ \sqrt{\varepsilon} u'_{\varepsilon} &\to 0 \quad weakly^* \ in \ L^{\infty}(0, T; H), \qquad \varepsilon u''_{\varepsilon} \to 0 \quad weakly \ in \ \mathcal{V}^*, \end{split}$$

as $\varepsilon \to 0$. Moreover, the limit function u is a solution to the problem (7).

The study of inclusions of type (7) is motivated by several contact problems of solid mechanics. It is well known that the dynamic equations of motion, representing momentum conservation, that govern the evolution of the state of the body, are of the form $u''_i - \sigma_{ij,j} = f_i$, where u is a displacement, σ is the stress tensor and f is the density (per unit volume) of applied forces. In many cases we are interested in situations in which the system configuration and the external forces and tractions evolve slowly in time in such a way that the accelerations in the system are rather small and negligible, so that the inertial (the second order time derivative) terms can be neglected. In such a way, we obtain the quasistatic approximation (equilibrium equations) for the equations of motion, $-\text{Div }\sigma = f$, where Div is the divergence operator. In this approximation, at each time instant the system is in equilibrium, and the external forces are balanced by the internal stresses. Rigorous mathematical treatment of quasistatic problems is recent. The rapid growth of the theory of quasistatic contact models can be seen from several monographs and many papers dedicated to such phenomena, cf. Duvaut and Lions [4], Han and Sofonea [6], Shillor et al. [19] and references therein. When we assume that the process is slow and the accelerations are negligible, mathematically it means that the system changes character, from being of a hyperbolic type to an elliptic or a parabolic type. To our knowledge quasistatic hemivariational inequalities have been studied in the literature only recently and a recent existence result for such problems is provided in Migórski and Ochal [13]. We mention that quasistatic problems modeled by variational inequalities with strongly monotone and Lipschitz operators were considered in [6, 19]. We also point out that the result of Theorem 4.6 can be applicable to several important and challenging quasistatic models of contact phenomena. Finally, we remark that the question of uniqueness of solutions to both dynamical and guasistatic hemivariational inequality is left open. It would be also interesting to extend the result of Theorem 4.6 to a class of problems with nonlinear operators A(t) and B.

5. History-dependent hemivariational inequality

In this section we provide a result on existence and uniqueness of a solution to a class of hemivariational inequalities associated with Problem 3.1. We adopt the notation of Section 3. The problem under consideration reads as follows.

Problem 5.1.

Find $u \in \mathcal{V}$ such that

$$\langle A(t, u(t)), v \rangle_{V^* \times V} + \langle Su(t), v \rangle_{V^* \times V} + \int_{\Gamma} j^0(t, \gamma u(t); \gamma v) \, d\Gamma \geq \langle f(t), v \rangle_{V^* \times V} \tag{9}$$

for all $v \in V$ and a.e. $t \in (0, T)$.

In the study of the hemivariational inequality (9), in addition to the previous assumptions, we need the following hypothesis.

V. $j: \Gamma \times (0, T) \times \mathbb{R}^s \to \mathbb{R}$ is such that

(V.a) $j(\cdot, \cdot, \xi)$ is measurable on $\Gamma \times (0, T)$ for all $\xi \in \mathbb{R}^s$ and $j(\cdot, \cdot, 0) \in L^1(\Gamma \times (0, T))$; (V.b) $j(x, t, \cdot)$ is locally Lipschitz on \mathbb{R}^s for a.e. $(x, t) \in \Gamma \times (0, T)$;

- (V.c) $\|\partial j(x, t, \xi)\|_{\mathbb{R}^s} \leq \overline{c}_0 + \overline{c}_1 \|\xi\|_{\mathbb{R}^s}$ for a.e. $(x, t) \in \Gamma \times (0, T)$, all $\xi \in \mathbb{R}^s$ with $\overline{c}_0, \overline{c}_1 \geq 0$;
- (V.d) $(\zeta_1 \zeta_2) \cdot (\xi_1 \xi_2) \ge -m_2 \|\xi_1 \xi_2\|_{\mathbb{R}^5}^2$ for all $\zeta_i, \xi_i \in \mathbb{R}^s, \zeta_i \in \partial j(x, t, \xi_i), i = 1, 2, \text{ a.e. } (x, t) \in \Gamma \times (0, T)$ with $m_2 \ge 0$;
- (V.e) $j^0(x, t, \xi; -\xi) \leq \overline{d}_0(1 + \|\xi\|_{\mathbb{R}^s})$ for a.e. $(x, t) \in \Gamma \times (0, T)$, all $\xi \in \mathbb{R}^s$ with $\overline{d}_0 \geq 0$.

Note that in the condition (V.d) the dot denotes the inner product in \mathbb{R}^{s} . From Theorem 3.4, we deduce the following existence and uniqueness result for the hemivariational inequality (9). The details on its proof can be found in [14].

Theorem 5.2.

Assume that | and || hold and $f \in \mathcal{V}^*$. If one of the following hypotheses:

- (i) (V.a)–(V.d) and $m_1 > \max{\{\sqrt{3}\overline{c}_1, m_2\}} c_e^2 \|\gamma\|^2$,
- (ii) V and $m_1 > m_2 c_e^2 \| \mathbf{y} \|^2$

is satisfied, then Problem 5.1 has a solution $u \in \mathcal{V}$. If, in addition, the regularity condition:

either $j(x, t, \cdot)$ or $-j(x, t, \cdot)$ is regular on \mathbb{R}^s for a.e. $(x, t) \in \Gamma \times (0, T)$

holds, then the solution of Problem 5.1 is unique.

6. Quasi-static frictional contact problems

A large number of quasistatic contact problems with elastic or viscoelastic materials leads to a hemivariational inequality of the form (9) in which the unknown is either the velocity or the displacement field. In both cases the abstract results of Sections 3, 4 and 5 work and can be used to provide the unique weak solvability of the corresponding contact problems. In this section we illustrate the use of these results in the study of two contact problems.

The physical setting of the first contact problem is as follows. A viscoelastic body occupies a domain Ω of \mathbb{R}^d , d = 2, 3, with surface $\partial\Omega$ which is partitioned into three disjoint measurable parts Γ_D , Γ_N and Γ_C such that $m(\Gamma_D) > 0$. We are interested in the evolution process of the mechanical state of the body in the bounded interval of time [0, T], where T > 0. The body is clamped on Γ_D and so the displacement field vanishes there. Surface tractions of density \mathbf{f}_N act on Γ_N and volume forces of density \mathbf{f}_0 act in Ω . We assume that the forces and tractions change slowly in time so that the acceleration of the system is negligible and, therefore, the process is quasistatic. Moreover, the body is in contact with an obstacle on Γ_C , the so-called foundation. The contact is frictional and is modeled with subdifferential boundary conditions.

We use the notation $\mathbf{v} = (\mathbf{v}_i)$ for the outward unit normal at $\partial\Omega$. We denote by $\mathbf{u} = (u_i)$, $\mathbf{\sigma} = (\sigma_{ij})$, and $\mathbf{\varepsilon}(\mathbf{u}) = (\varepsilon_{ij}(\mathbf{u}))$ the displacement vector, the stress tensor, and the linearized strain tensor, respectively. We recall that the components of the linearized strain tensor $\mathbf{\varepsilon}(\mathbf{u})$ are given by $\varepsilon_{ij}(\mathbf{u}) = (u_{i,j} + u_{j,i})/2$. We denote by \mathbb{S}^d the space of second order symmetric tensors on \mathbb{R}^d or, equivalently, the space of symmetric matrices of order d. The canonical inner products and the corresponding norms on \mathbb{R}^d and \mathbb{S}^d are given by

$$\begin{split} \boldsymbol{u} \cdot \boldsymbol{v} &= u_i v_i, \quad \|\boldsymbol{v}\|_{\mathbb{R}^d} = (\boldsymbol{v} \cdot \boldsymbol{v})^{1/2} \quad \text{ for all } \boldsymbol{u} = (u_i), \quad \boldsymbol{v} = (v_i) \in \mathbb{R}^d, \\ \boldsymbol{\sigma} &: \boldsymbol{\tau} = \sigma_{ij} \tau_{ij}, \quad \|\boldsymbol{\tau}\|_{\mathbb{S}^d} = (\boldsymbol{\tau} : \boldsymbol{\tau})^{1/2} \quad \text{ for all } \boldsymbol{\sigma} = (\sigma_{ij}), \quad \boldsymbol{\tau} = (\tau_{ij}) \in \mathbb{S}^d, \end{split}$$

respectively.

With these preliminaries, the classical formulation of the quasistatic contact problem we consider in this section is the following.

Problem 6.1.

Find the displacement field $\boldsymbol{u}: \Omega \times [0, T] \to \mathbb{R}^d$ and the stress field $\boldsymbol{\sigma}: \Omega \times [0, T] \to \mathbb{S}^d$ such that, for all $t \in (0, T)$,

 $\sigma(t) = \mathcal{A}(t, \varepsilon(u'(t))) + \mathcal{B}(t, \varepsilon(u(t))) \quad \text{in } \Omega,$ $\text{Div } \sigma(t) + f_0(t) = \mathbf{0} \quad \text{in } \Omega,$ (10a)
(10b)

$$\boldsymbol{u}(t) = \boldsymbol{0} \qquad \text{on} \quad \boldsymbol{\Gamma}_{D}, \tag{10c}$$

$$\sigma(t)\mathbf{v} = f_N(t) \quad \text{on} \quad \Gamma_N, \tag{10d}$$

$$-\sigma_{v}(t) \in \partial j_{v}(t, u'_{v}(t)) \quad \text{on } \mid_{C}, \tag{10e}$$
$$-\sigma_{v}(t) \in \partial i_{v}(t, u'_{v}(t)) \quad \text{on } \mid_{C} \tag{10f}$$

$$\boldsymbol{\sigma}_{\tau}(t) \in \mathcal{O}_{J_{\tau}}(t, \boldsymbol{u}_{\tau}(t)) \qquad \text{on} \quad \boldsymbol{\Gamma}_{C}, \tag{10f}$$

$$\boldsymbol{u}(0) = \boldsymbol{u}_0 \qquad \text{in} \quad \boldsymbol{\Omega}. \tag{10g}$$

We present a short description of the equations and conditions in Problem 6.1 and we refer the reader to [6, 15, 19] for more details and mechanical interpretation.

The equation (10a) represents the viscoelastic constitutive law in which A and B are given nonlinear operators, called the viscosity operator and elasticity operator, respectively. The explicit dependence of the viscosity and elasticity operators A and B with respect to the time variable allows to model situations when the properties of the material depend on the temperature, which plays the role of a parameter, i.e. its evolution in time is prescribed. Equality (10b) is the equilibrium equation, where Div represents the divergence operator, i.e. Div $\sigma = (\sigma_{ij,j})$. Conditions (10c) and (10d) are the displacement and traction boundary conditions, respectively, and (10g) is the initial condition in which the function u_0 denotes the initial displacement field.

Conditions (10e) and (10f) represent the frictional contact conditions in which j_v and j_τ are given functions. The subscripts v and τ for σ and u' indicate normal and tangential components of tensors and vectors. The symbol ∂j denotes the Clarke subdifferential of j with respect to the last variable. Concrete examples of frictional models which lead to subdifferential boundary conditions of the form (10e), (10f) with the functions j_v and j_τ satisfying assumptions VIII and IX below can be found in [12, 15]. Here, we restrict ourselves to remark that these examples include the viscous contact and the contact with nonmonotone normal damped response, associated to a nonmonotone friction law, to Tresca's friction law or to a power-law friction. Our results below are valid for the corresponding quasistatic frictional contact problems.

In the study of the contact problem (10a)–(10g) we use standard notation $Q = \Omega \times (0, T)$, $\Sigma_D = \Gamma_D \times (0, T)$, $\Sigma_N = \Gamma_N \times (0, T)$ and $\Sigma_C = \Gamma_C \times (0, T)$. For all $\mathbf{v} \in H^1(\Omega; \mathbb{R}^d)$ we still denote by \mathbf{v} the trace of \mathbf{v} on Γ and we use the notation v_v and v_τ for the normal and tangential components of \mathbf{v} on $\partial\Omega$ given by $v_v = \mathbf{v} \cdot \mathbf{v}$ and $v_\tau = \mathbf{v} - v_v \mathbf{v}$. We recall that the normal and tangential components of the stress field $\boldsymbol{\sigma}$ on the boundary are defined by $\sigma_v = (\boldsymbol{\sigma} \mathbf{v}) \cdot \mathbf{v}$ and $\boldsymbol{\sigma}_\tau = \boldsymbol{\sigma} \mathbf{v} - \sigma_v \mathbf{v}$. Next, we introduce the spaces V and \mathcal{H} , defined by $V = \{\mathbf{v} = (v_i) \in H^1(\Omega; \mathbb{R}^d) : \mathbf{v} = \mathbf{0} \text{ a.e. on } \Gamma_D\}$ and $\mathcal{H} = L^2(\Omega; \mathbb{S}^d)$.

We assume that the viscosity operator $\mathcal A$ and the elasticity operator $\mathcal B$ satisfy

- VI. $\mathcal{A}: Q \times \mathbb{S}^d \to \mathbb{S}^d$ is such that
 - (VI.a) $\mathcal{A}(\cdot, \cdot, \varepsilon)$ is measurable on Q for all $\varepsilon \in \mathbb{S}^d$;
 - (VI.b) $\mathcal{A}(\mathbf{x}, t, \cdot)$ is continuous on \mathbb{S}^d for a.e. $(\mathbf{x}, t) \in Q$;
 - (VI.c) $(\mathcal{A}(\mathbf{x}, t, \varepsilon_1) \mathcal{A}(\mathbf{x}, t, \varepsilon_2)): (\varepsilon_1 \varepsilon_2) \ge m_{\mathcal{A}} \|\varepsilon_1 \varepsilon_2\|_{\mathbb{S}^d}^2$ for all $\varepsilon_1, \varepsilon_2 \in \mathbb{S}^d$, a.e. $(\mathbf{x}, t) \in Q$ with $m_{\mathcal{A}} > 0$;
 - $(\text{VI.d}) \|\mathcal{A}(\mathbf{x}, t, \boldsymbol{\varepsilon})\|_{\mathbb{S}^d} \leq \overline{a}_0(\mathbf{x}, t) + \overline{a}_1 \|\boldsymbol{\varepsilon}\|_{\mathbb{S}^d} \text{ for all } \boldsymbol{\varepsilon} \in \mathbb{S}^d, \text{ a.e. } (\mathbf{x}, t) \in Q \text{ with } \overline{a}_0 \in L^2(Q), \ \overline{a}_0 \geq 0 \text{ and } \overline{a}_1 > 0;$
 - (VI.e) $\mathcal{A}(\mathbf{x}, t, \mathbf{0}) = \mathbf{0}$ for a.e. $(\mathbf{x}, t) \in Q$.

VII. $\mathcal{B}: Q \times \mathbb{S}^d \to \mathbb{S}^d$ is such that

- (VII.a) $\mathcal{B}(\cdot, \cdot, \varepsilon)$ is measurable on Q for all $\varepsilon \in \mathbb{S}^d$;
- (VII.b) $\|\mathcal{B}(\mathbf{x}, t, \varepsilon_1) \mathcal{B}(\mathbf{x}, t, \varepsilon_2)\|_{\mathbb{S}^d} \le L_{\mathcal{B}} \|\varepsilon_1 \varepsilon_2\|_{\mathbb{S}^d}$ for all $\varepsilon_1, \varepsilon_2 \in \mathbb{S}^d$, a.e. $(\mathbf{x}, t) \in Q$ with $L_{\mathcal{B}} > 0$;
- (VII.c) $\mathcal{B}(\cdot, \cdot, \mathbf{0}) \in L^2(Q; \mathbb{S}^d)$.

The contact potentials j_{ν} and j_{τ} satisfy the following hypotheses.

VIII. $j_{\nu} : \Sigma_C \times \mathbb{R} \to \mathbb{R}$ is such that

(VIII.a) $j_{\nu}(\cdot, \cdot, r)$ is measurable on Σ_{C} for all $r \in \mathbb{R}$ and $j_{\nu}(\cdot, \cdot, 0) \in L^{1}(\Sigma_{C})$;

- (VIII.b) $j_{\nu}(x, t, \cdot)$ is locally Lipschitz on \mathbb{R} for a.e. $(x, t) \in \Sigma_C$;
- (VIII.c) $|\partial j_{\nu}(x, t, r)| \leq c_{0\nu} + c_{1\nu}|r|$ for all $r \in \mathbb{R}$, a.e. $(x, t) \in \Sigma_C$ with $c_{0\nu}, c_{1\nu} \geq 0$;
- (VIII.d) $(\zeta_1 \zeta_2)(r_1 r_2) \ge -m_v |r_1 r_2|^2$ for all $\zeta_i \in \partial j_v(x, t, r_i), r_i \in \mathbb{R}, i = 1, 2, \text{ a.e. } (x, t) \in \Sigma_C$ with $m_v \ge 0$;
- (VIII.e) $j_{\nu}^{0}(x, t, r; -r) \leq d_{\nu}(1+|r|)$ for all $r \in \mathbb{R}$, a.e. $(x, t) \in \Sigma_{C}$ with $d_{\nu} \geq 0$.
- IX. $j_{\tau} \colon \Sigma_C \times \mathbb{R}^d \to \mathbb{R}$ is such that
 - (IX.a) $j_{\tau}(\cdot, \cdot, \xi)$ is measurable on $\Sigma_{\mathcal{C}}$ for all $\xi \in \mathbb{R}^d$ and $j_{\tau}(\cdot, \cdot, \mathbf{0}) \in L^1(\Sigma_{\mathcal{C}})$;
 - (IX.b) $j_{\tau}(x, t, \cdot)$ is locally Lipschitz on \mathbb{R}^d for a.e. $(x, t) \in \Sigma_C$;
 - (IX.c) $\|\partial j_{\tau}(\mathbf{x}, t, \boldsymbol{\xi})\|_{\mathbb{R}^d} \leq c_{0\tau} + c_{1\tau} \|\boldsymbol{\xi}\|_{\mathbb{R}^d}$ for all $\boldsymbol{\xi} \in \mathbb{R}^d$, a.e. $(\mathbf{x}, t) \in \Sigma_C$ with $c_{0\tau}, c_{1\tau} \geq 0$;
 - $(\mathsf{IX.d}) \ (\boldsymbol{\zeta}_1 \boldsymbol{\zeta}_2) \cdot (\boldsymbol{\xi}_1 \boldsymbol{\xi}_2) \geq -m_\tau \| \boldsymbol{\xi}_1 \boldsymbol{\xi}_2 \|_{\mathbb{R}^d}^2 \text{ for all } \boldsymbol{\zeta}_i \in \partial j_\tau(\boldsymbol{x}, t, \boldsymbol{\xi}_i), \ \boldsymbol{\xi}_i \in \mathbb{R}^d, \ i = 1, 2, \text{ a.e. } (\boldsymbol{x}, t) \in \Sigma_C \text{ with } m_\tau \geq 0;$
 - (IX.e) $j^0_{\tau}(\mathbf{x}, t, \boldsymbol{\xi}; -\boldsymbol{\xi}) \leq d_{\tau}(1 + \|\boldsymbol{\xi}\|_{\mathbb{R}^d})$ for all $\boldsymbol{\xi} \in \mathbb{R}^d$, a.e. $(\mathbf{x}, t) \in \Sigma_C$ with $d_{\tau} \geq 0$.

The volume force and traction densities satisfy

$$\boldsymbol{f}_{0} \in L^{2}(0, T; L^{2}(\Omega; \mathbb{R}^{d})), \qquad \boldsymbol{f}_{N} \in L^{2}(0, T; L^{2}(\Gamma_{N}; \mathbb{R}^{d}))$$

$$(11)$$

and, finally, the initial displacement is such that

$$\boldsymbol{u}_0 \in \boldsymbol{V}. \tag{12}$$

We turn now to the variational formulation of Problem 6.1. Suppose that (u, σ) is a couple of sufficiently smooth functions which solve (10a)-(10g). Let $v \in V$. Then, using (10b), we have

$$(\boldsymbol{\sigma}(t),\boldsymbol{\varepsilon}(\boldsymbol{v}))_{\mathcal{H}} = (\boldsymbol{f}_0(t),\boldsymbol{v})_{L^2(\Omega;\mathbb{R}^d)} + \int_{\partial\Omega} \boldsymbol{\sigma}(t) \boldsymbol{v} \cdot \boldsymbol{v} \, d\Gamma \qquad \text{for a.e.} \quad t \in (0,T).$$
(13)

We take into account the boundary conditions (10c) and (10d) to see that

$$\int_{\partial\Omega} \boldsymbol{\sigma}(t) \, \boldsymbol{v} \cdot \boldsymbol{v} \, d\Gamma = \int_{\Gamma_N} \boldsymbol{f}_N(t) \cdot \boldsymbol{v} \, d\Gamma + \int_{\Gamma_C} (\boldsymbol{\sigma}_v(t) v_v + \boldsymbol{\sigma}_\tau(t) \cdot \boldsymbol{v}_\tau) \, d\Gamma \qquad \text{for a.e.} \quad t \in (0, T).$$
(14)

On the other hand, from the definition of the Clarke subdifferential, (10e) and (10f), we have

$$-\sigma_{\nu}(t)v_{\nu} \leq j_{\nu}^{0}(t, u_{\nu}'(t); v_{\nu}), \qquad -\sigma_{\tau}(t) \cdot v_{\tau} \leq j_{\tau}^{0}(t, u_{\tau}'(t); v_{\tau}) \qquad \text{on} \quad \Sigma_{C},$$

which imply that

$$\int_{\Gamma_{C}} \left(\sigma_{\nu}(t) v_{\nu} + \boldsymbol{\sigma}_{\tau}(t) \cdot \boldsymbol{v}_{\tau} \right) d\Gamma \geq - \int_{\Gamma_{C}} \left(j_{\nu}^{0}(t, \boldsymbol{u}_{\nu}'(t); \boldsymbol{v}_{\nu}) + j_{\tau}^{0}(t, \boldsymbol{u}_{\tau}'(t); \boldsymbol{v}_{\tau}) \right) d\Gamma$$
(15)

for a.e. $t \in (0, T)$. Consider the function $f: (0, T) \to V^*$ given by

$$\langle f(t), v \rangle_{V^* \times V} = (f_0(t), v)_{L^2(\Omega; \mathbb{R}^d)} + (f_N(t), v)_{L^2(\Gamma_N; \mathbb{R}^d)}$$
(16)

for all $v \in V$ and a.e. $t \in (0, T)$. We combine (13)–(16) to obtain

$$(\boldsymbol{\sigma}(t),\boldsymbol{\varepsilon}(\boldsymbol{v}))_{\mathcal{H}} + \int_{\Gamma_{C}} (j_{\boldsymbol{v}}^{0}(t,\boldsymbol{u}_{\boldsymbol{v}}'(t);\boldsymbol{v}_{\boldsymbol{v}}) + j_{\tau}^{0}(t,\boldsymbol{u}_{\tau}'(t);\boldsymbol{v}_{\tau})) d\Gamma \geq \langle \boldsymbol{f}(t),\boldsymbol{v} \rangle_{V^{*} \times V}$$

and, using the constitutive law (10a), it follows

$$\left(\mathcal{A}(t,\boldsymbol{\varepsilon}(\boldsymbol{u}'(t))),\boldsymbol{\varepsilon}(\boldsymbol{v})\right)_{\mathcal{H}} + \left(\mathcal{B}(t,\boldsymbol{\varepsilon}(\boldsymbol{u}(t))),\boldsymbol{\varepsilon}(\boldsymbol{v})\right)_{\mathcal{H}} + \int_{\Gamma_{C}} (j_{\nu}^{0}(t,\boldsymbol{u}_{\nu}'(t);\boldsymbol{v}_{\nu}) + j_{\tau}^{0}(t,\boldsymbol{u}_{\tau}'(t);\boldsymbol{v}_{\tau})) d\Gamma \geq \langle \boldsymbol{f}(t),\boldsymbol{v} \rangle_{V^{*}\times V}$$
(17)

for all $v \in V$ and a.e. $t \in (0, T)$. Let w = u' denote the velocity field. Then, by using the initial condition (10g), it follows that

$$u(t) = \int_0^t w(s) \, ds + u_0$$
 for all $t \in [0, T]$. (18)

Therefore, (17) and (18) lead to the following variational formulation of problem (10a)-(10g), in terms of velocity.

Problem 6.2.

Find a velocity field $w \in \mathcal{V}$ such that

$$\left(\mathcal{A}(t,\boldsymbol{\varepsilon}(\boldsymbol{w}(t))),\boldsymbol{\varepsilon}(\boldsymbol{v})\right)_{\mathcal{H}}+\left(\mathcal{B}\left(t,\boldsymbol{\varepsilon}\left(\int_{0}^{t}\boldsymbol{w}(s)\,ds+\boldsymbol{u}_{0}\right)\right),\boldsymbol{\varepsilon}(\boldsymbol{v})\right)_{\mathcal{H}}+\int_{\Gamma_{C}}(j_{\nu}^{0}(t,w_{\nu}(t);\nu_{\nu})+j_{\tau}^{0}(t,w_{\tau}(t);\nu_{\tau}))\,d\Gamma\geq\langle\boldsymbol{f}(t),\boldsymbol{v}\rangle_{V^{*}\times V}\right)_{\mathcal{H}}$$

for all $v \in V$ and a.e. $t \in (0, T)$.

The hemivariational inequality in Problem 6.2 is of the form of the inequality in Problem 5.1. From Theorem 5.2, we obtain the following result on Problem 6.2.

Theorem 6.3.

Assume that (VI.a), VII, (11) and (12) hold. If one of the following hypotheses:

- (i) (VIII.a)–(VIII.d), (IX.a)–(IX.d) and $m_1 > \max\{\sqrt{3}(c_{1\nu}+c_{1\tau}), m_{\nu}, m_{\tau}\}c_{\rho}^2 \|\gamma\|^2$,
- (ii) VIII, IX and $m_1 > \max\{m_v, m_\tau\} c_e^2 \|\gamma\|^2$

is satisfied, then Problem 6.2 has at least one solution. If, in addition,

either
$$j_{y}(x, t, \cdot), j_{\tau}(x, t, \cdot)$$
 are regular or $-j_{y}(x, t, \cdot), -j_{\tau}(x, t, \cdot)$ are regular for a.e. $(x, t) \in \Sigma_{C}$, (19)

then the solution of Problem 6.2 is unique.

Let *w* be a solution of Problem 6.2 and denote by *u* and σ the functions defined by (18) and (10a). Then, the couple (u, σ) is called a *weak solution* of the frictional contact problem (10a)–(10g). We conclude, under the hypotheses of Theorem 6.3, that the frictional contact problem (10a)–(10g) has at least one weak solution with the following regularity:

 $\boldsymbol{u} \in W^{1,2}(0, T; V)$ and $\boldsymbol{\sigma} \in L^2(0, T; \mathcal{H})$ with $\text{Div}\,\boldsymbol{\sigma} \in L^2(0, T; V^*)$.

If, in addition, the regularity condition (19) holds, then the weak solution of Problem 6.1 is unique.

We now pass to the second problem of this section. We consider the quasistatic viscoelastic contact with nonmonotone normal compliance and friction. As in the first problem, we assume that the volume forces and surface tractions change slowly in time so that the acceleration in the system is negligible. We show that the quasistatic model can be formulated as a time dependent hemivariational inequality of the form (7) and the abstract result of Theorem 3.4 is applicable in this case. For the mechanical formulation of the process we use the notation introduced above. We set $Q = \Omega \times (0, T)$. Again we denote by $u: Q \to \mathbb{R}^d$ the displacement field, by $\sigma: Q \to \mathbb{S}^d$ the stress tensor and by $\varepsilon(u) = {\varepsilon_{ij}(u)}, \varepsilon_{ij}(u) = (\partial_j u_i + \partial_i u_j)/2$ the strain tensor. We assume a linear viscoelastic model with the constitutive law of the Kelvin–Voigt type

$$\sigma_{ij} = a_{ijkl} \varepsilon_{kl}(u') + b_{ijkl} \varepsilon_{kl}(u) \quad \text{in } Q$$

where $\mathcal{C}(t) = \{a_{ijkl}(t)\}$ and $\mathcal{G} = \{b_{ijkl}\}, i, j, k, l = 1, ..., d$, are the viscosity and the elasticity tensors, respectively. Concerning the contact conditions, we consider the following subdifferential relations $-\sigma_v \in \partial j_v(x, t, u_v)$ and $-\sigma_\tau \in \partial j_\tau(x, t, u_\tau)$. The functions $j_v: \Gamma_C \times (0, T) \times \mathbb{R} \to \mathbb{R}$ and $j_\tau: \Gamma_C \times (0, T) \times \mathbb{R}^d \to \mathbb{R}$ are locally Lipschitz in their last variables and $\partial j_v, \partial j_\tau$ denote their Clarke subdifferentials. The initial displacement is denoted by u_0 . The classical formulation of the mechanical problem is as follows. Find the displacement field $u: \Omega \times [0, T] \to \mathbb{R}^d$ and the stress field $\sigma: \Omega \times [0, T] \to \mathbb{S}^d$ such that, for all $t \in (0, T)$,

$$\boldsymbol{\sigma}(t) = \mathbb{C}(t)\boldsymbol{\varepsilon}(\boldsymbol{u}'(t)) + \mathcal{G}\boldsymbol{\varepsilon}(\boldsymbol{u}(t)) \quad \text{in } \Omega,$$
(20a)

Div
$$\boldsymbol{\sigma}(t) + \boldsymbol{f}_0(t) = \boldsymbol{0}$$
 in Ω , (20b)

$$\boldsymbol{u}(t) = \boldsymbol{0} \qquad \text{on} \quad \boldsymbol{\Gamma}_D, \tag{20c}$$

$$\boldsymbol{\sigma}(t)\,\boldsymbol{\nu}=\boldsymbol{f}_N(t)\qquad\text{on}\quad\boldsymbol{\Gamma}_N,\tag{20d}$$

$$-\sigma_{\nu}(t) \in dj_{\nu}(t, u_{\nu}(t)) \quad \text{on } \Gamma_{C},$$
(20e)

$$\boldsymbol{\sigma}_{\tau}(t) \in \partial j_{\tau}(t, \boldsymbol{u}_{\tau}(t)) \quad \text{on } \boldsymbol{\Gamma}_{C}, \tag{201}$$

$$\boldsymbol{u}(0) = \boldsymbol{u}_0 \qquad \text{in} \quad \Omega. \tag{20g}$$

For concrete examples of boundary conditions (20e) and (20f), we refer to [4, 6, 16-18].

In order to obtain a variational formulation of the problem (20a)–(20g) we need additional notation. Let $V = \{v = (v_i) \in H^1(\Omega; \mathbb{R}^d) : v = 0$ a.e. on $\Gamma_D\}$, $Z = H^{\delta}(\Omega; \mathbb{R}^d)$ with a fixed $\delta \in (1/2, 1)$, $H = L^2(\Omega; \mathbb{R}^d)$ and $X = L^2(\Gamma_C; \mathbb{R}^d)$. On V we consider the inner product and the corresponding norm given by $\langle u, v \rangle_V = \langle \varepsilon(u), \varepsilon(v) \rangle_{L^2(\Omega; \mathbb{S}^d)}$, $\|v\|_V = \|\varepsilon(v)\|_{L^2(\Omega; \mathbb{S}^d)}$ for $u, v \in V$. Since $m(\Gamma_D) > 0$, the Korn inequality $\|v\|_{H^1(\Omega; \mathbb{R}^d)} \le c \|\varepsilon(v)\|_{L^2(\Omega; \mathbb{S}^d)}$ for $v \in V$ with c > 0 holds, and therefore the norms $\|\cdot\|_{H^1(\Omega; \mathbb{R}^d)}$ and $\|\cdot\|_V$ are equivalent on V. Moreover, we put $M = \gamma \in \mathcal{L}(Z, X)$, γ is the trace operator, and, for simplicity, p = 2.

We also assume that the viscosity and elasticity tensors have the usual properties of ellipticity, symmetry and positivity.

- H(C) C: $Q \times \mathbb{S}^d \to \mathbb{S}^d$ is a viscosity tensor, $\mathbb{C}(t) = \{a_{ijkl}(t)\}$, such that $a_{ijkl} = a_{klij} = a_{ijlk} \in L^{\infty}(Q)$ and there exists $m_1 > 0$ satisfying $\mathbb{C}(t)\tau:\tau \ge m_1 \|\tau\|_{\mathbb{S}^d}^2$ for all $\tau \in \mathbb{S}^d$, a.e. in Q.
- H(G) $\mathfrak{G}: \Omega \times \mathbb{S}^d \to \mathbb{S}^d$ is an elasticity tensor, $\mathfrak{G} = \{b_{ijkl}\}$, such that $b_{ijkl} = b_{klij} = b_{ijlk} \in L^{\infty}(\Omega)$ and $\mathfrak{G}\tau: \tau \ge 0$ for all $\tau \in \mathbb{S}^d$, a.e. in Ω .

Next, we define the operators A(t), $B \in \mathcal{L}(V, V^*)$ by

$$\langle A(t)u, v \rangle = \langle \mathcal{C}(t)\boldsymbol{\varepsilon}(u), \boldsymbol{\varepsilon}(v) \rangle_{L^2(\Omega; \mathbb{S}^d)}, \qquad \langle Bu, v \rangle = \langle \boldsymbol{\Im}\boldsymbol{\varepsilon}(u), \boldsymbol{\varepsilon}(v) \rangle_{L^2(\Omega; \mathbb{S}^d)}$$
(21)

for $t \in (0, T)$ and $u, v \in V$, and $f \in \mathcal{V}^*$ by (16). Analogously as in the first contact problem, we derive the following hemivariational inequality which is a weak form of (20a)–(20g): find $u: (0, T) \to V$ such that $u \in L^{\infty}(0, T; V)$, $u' \in \mathcal{V}$ and

$$\begin{cases} \langle A(t)u'(t) + Bu(t), v \rangle + \int_{\Gamma_C} (j_v^0(t, u_v; v_v) + j_\tau^0(t, u_\tau; v_\tau)) \, d\Gamma \geq \langle f(t), v \rangle & \text{for a.e. } t \in (0, T) \text{ and all } v \in V, \\ u(0) = u_0. \end{cases}$$

$$(22)$$

Let us consider the functional $J: (0, T) \times X \to \mathbb{R}$ defined by

$$J(t,v) = \int_{\Gamma_C} (j_v(x,t,v_v(x)) + j_\tau(x,t,v_\tau(x))) d\Gamma \quad \text{for a.e.} \quad t \in (0,T) \quad \text{and all} \quad v \in X,$$
(23)

where $X = L^2(\Gamma_C; \mathbb{R}^d)$.

Lemma 6.4.

Under the hypotheses VIII and IX, the functional J defined by (23) satisfies

- (i) $J(\cdot, v)$ is measurable for all $v \in X$ and $J(\cdot, 0) \in L^{1}(0, T)$;
- (ii) $J(t, \cdot)$ is Lipschitz on bounded subsets of X;
- (iii) $\|\partial J(t,v)\|_{X^*} \leq \overline{c}(1+\|v\|_X)$ for all $v \in X$, a.e. $t \in (0,T)$ with $\overline{c} = 2\sqrt{2} \max\{c_v, c_\tau\} \max\{1, \sqrt{m(\Gamma_C)}\}$;
- (iv) for all $v, w \in X$, we have

$$J^{0}(t, v; w) \leq \int_{\Gamma_{C}} (j_{v}^{0}(t, v_{v}; w_{v}) + j_{\tau}^{0}(t, v_{\tau}; w_{\tau})) d\Gamma, \qquad (24)$$

where $J^0(t, v; w)$ denotes the directional derivative of $J(t, \cdot)$ at a point $v \in X$ in the direction $w \in X$. If, in addition,

either
$$j_{v}(x, t, \cdot), j_{\tau}(x, t, \cdot)$$
 are regular or $-j_{v}(x, t, \cdot), -j_{\tau}(x, t, \cdot)$ are regular, (25)

then either $J(t, \cdot)$ or $-J(t, \cdot)$ is regular, respectively and (24) holds with equality.

Under our notation we associate with the hemivariational inequality (22) the following inclusion of type (7): find $u \in L^{\infty}(0, T; V)$ with $u' \in \mathcal{V}$ such that

$$\begin{cases} A(t)u'(t) + Bu(t) + \gamma^* \partial J(t, \gamma u(t)) \ni f(t) & \text{for a.e. } t \in (0, T), \\ u(0) = u_0. \end{cases}$$
(26)

Remark 6.5.

We notice that if the hypotheses VIII and IX hold, then every solution to (26) is a solution to (22). The converse holds provided j_v and j_τ satisfy the regularity condition (25). These facts follow from the definition of the Clarke subdifferential and Lemma 6.4 (iv).

The existence and regularity result for the hemivariational inequality (22) reads as follows.

Theorem 6.6.

If the hypotheses H(C), H(G), (VIII.a)–(VIII.c), (IX.a)–(IX.c), (11), (12) hold and $m_1 > 2\overline{c}Tc_e^2 \|\gamma\| \max\{1, \|\gamma\|\}$, where \overline{c} is given in Lemma 6.4 (iii), $c_e > 0$ is the embedding constant of V into Z and $\|\gamma\| = \|\gamma\|_{\mathcal{L}(Z;X)}$, then the problem (22) admits at least one solution $u \in L^{\infty}(0, T; V)$ with $u' \in \mathcal{V}$.

Proof. It follows from H(\mathcal{C}) and H(\mathcal{G}) that the operators A(t) and B defined by (21) satisfy H(A) with $\alpha = m_1$ and H(B), respectively. It is a consequence of Lemma 6.4 (i)–(iii) that the functional J given by (23) satisfies H(J). Also H(M) follows easily by the properties of the trace operator. The conclusion follows from Theorem 4.2 and Remark 6.5.

We conclude this section with short comments on multivalued boundary conditions (20e) and (20f) which are met in solid mechanics. The condition (20e) is a generalization of a normal compliance condition to the nonmonotone setting. Let the function $j_{\nu} : \mathbb{R} \to \mathbb{R}$ be defined by $j_{\nu}(r) = \int_{0}^{r} p_{\nu}(s) ds$, where for simplicity we drop the (x, t)-dependence and the function $p_{\nu} \in L^{\infty}_{loc}(\mathbb{R})$ is such that $|p_{\nu}(s)| \leq p_{1}(1+|s|)$ for $s \in \mathbb{R}$ with $p_{1} > 0$ and limits $\lim_{r \to s\pm} p_{\nu}(r)$ exist for all $s \in \mathbb{R}$. It is well known, see e.g. [16, 18], that $\partial j_{\nu}(r) = \hat{p}_{\nu}(r)$ for $r \in \mathbb{R}$, where the multivalued function $\hat{p}_{\nu} : \mathbb{R} \to 2^{\mathbb{R}}$ is given by $\hat{p}_{\nu}(r) = [p_{\nu}^{(1)}(r), p_{\nu}^{(2)}(r)], [\cdot, \cdot]$ denotes an interval in \mathbb{R} and

$$p_{\nu}^{(1)}(r) = \lim_{\varepsilon \to 0^+} \operatorname{ess\,inf}_{|\tau - r| \leq \varepsilon} p_{\nu}(\tau), \qquad p_{\nu}^{(2)}(r) = \lim_{\varepsilon \to 0^+} \operatorname{ess\,sup}_{|\tau - r| \leq \varepsilon} p_{\nu}(\tau).$$

In this case VIII holds and (20e) has the form $-\sigma_v \in \hat{p}_v(u_v)$ on $\Gamma_C \times (0, T)$. Modifications of this example describe the nonmonotone zigzag relations in contact laws for a granular material and a reinforced concrete, cf. [18, Sections 2.4 and 7.2], [16, Section 4.6] and [12]. We also observe that if, in addition, p_v is a continuous function, then (20e) reduces to $-\sigma_v = p_v(u_v)$ on $\Gamma_C \times (0, T)$ which is a well-known classical normal compliance condition. The normal compliance contact condition describes a reactive foundation. It assigns a reactive normal pressure which depends on the interpenetration of the asperities on the boundary surface and those of the foundation. We mention that the normal compliance condition is considered in engineering and certain mathematical publications as an approximation of the Signorini nonpenetration condition, and this idealization makes sense for practical reasons, cf. [19].

In the friction condition (20f), if $j_{\tau} = 0$, we obtain the frictionless contact. If $j_{\tau}(x, t, \xi) = S(x, t) \|\xi\|_{\mathbb{R}^d}$, where $S \in L^{\infty}(\Gamma_C \times (0, T))$ and S > 0 a.e., then we get a version of the static Tresca friction law where the friction bound depends on time. In the friction condition (20f) we suppose that the tangential shear on the contact surface is given as a nonmonotone, time dependent and possibly multivalued function of the tangential displacement. For other examples of functions j_v and j_{τ} , we refer to [12]. Other two- and three-dimensional nonconvex zigzag superpotential laws are detailed in [15, 16].

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