

Proper Orthogonal Decomposition for Reduced Order Modeling: 2D Heat Flow

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Abstract— Modeling issues of infinite dimensional systems is studied in this paper. Although the modeling problem has been solved to some extent, use of decomposition techniques still pose several difficulties. A prime one of this is the amount of data to be processed. Method of snapshots integrated with POD is a remedy. The second difficulty is the fact that the decomposition followed by a projection yields an autonomous set of finite dimensional ODEs that is not useful for developing a concise understanding of the input operator of the system. A numerical approach to handle this issue is presented in this paper. As the example, we study 2D heat flow problem. The results obtained confirm the theoretical claims of the paper and emphasize that the technique presented here is not only applicable to infinite dimensional linear systems but also to nonlinear ones.

I. INTRODUCTION

Although the applications of Proper Orthogonal Decomposition (POD) particularly focus on the extraction of coherent and dominant modes available in aerodynamic flows, the problem of tackling with huge amounts of data and technical difficulties in the obtained model constitute barriers between the stipulated efforts and the thorough understanding of the process, [1-6]. For this reason, we study the modeling problem that addresses the above mentioned difficulties appropriately on a 2D, simple, and a linear dynamics, namely the heat flow. The goal of the paper is to introduce how the data is obtained, how the external stimuli enter into the dynamics, and how the data set is processed so as to obtain a set of orthogonal basis, and finally, how the external stimuli is made explicit in an autonomous set of ODEs.

Modeling of systems displaying spatial continuum requires a careful consideration since the physical process under investigation is an infinite dimensional one due to the spatial continuity. Efforts in understanding the behavior of such systems have particularly focused on the low dimensional models capturing the essential behavioral properties with a few Ordinary Differential Equations (ODEs). This has been done by using modal decompositions such as Proper Orthogonal Decomposition (POD) and Singular Value Decomposition (SVD). Although neither the decomposition techniques nor the infinite dimensionality are new issues in this field, obtaining a model having the boundary conditions as external inputs is a major problem in the POD and SVD methods. More

explicitly, these approaches result in models where external control input appears in the dynamical equations implicitly, and this is not very useful for controller design. Another difficulty is the presence of modeling uncertainties, which stem from varying internal parameters or hypotheses that are not thoroughly valid. For the heat transport process, imprecise knowledge on thermal diffusivity parameter is a good example to study uncertainties.

The use of decomposition techniques in modeling of spatially continuous systems has extensively been studied in the field of aerodynamic flow control problems, [1-4]. Since the dynamics of the process under investigation is governed by Navier-Stokes equations, obtaining closed form solutions are very difficult and the modeling studies particularly focus on the real time observations from the process. For systems having two or more spatial dimensions, the POD technique has been utilized with the aid of snapshots method, [1-2]. Alternatively, for single dimensional processes, the same modeling procedure can be followed by exploiting the SVD technique.

Procedurally, in both of them, if the numerical data contains coherent modes, the expansion accurately describes the temporal modes and the spatial components distributing them over the physical domain of the process. Furthermore, the orthogonality of the basis functions, which describe the spatial properties, helps in finding a set of ODEs synthesizing the temporal modes. Although the algorithmic part seems straightforward, the final form of the ODEs depicts an autonomous system having no external input. At this point, several modifications are needed to separate the effect of boundary conditions, which constitute the inputs exciting the process. 2D heat transport problem is therefore a good candidate to study how modeling issues are addressed.

A number of variations of this problem has been taken into consideration in former studies, [5-7]. Atwell et al [5-6], have considered 2D heat transport problem with control input explicitly available in the Partial Differential Equation (PDE). The thermal diffusivity parameter has been taken as a known constant and several control strategies have been assessed with the modeling results of POD approach. In [6], the design has been discussed from the computational point of view.

Another work focusing on 1D heat transport problem re-

ports the design of time-optimal boundary control, [7]. It is emphasized in [7] that the time-optimal control has the bang-bang property, and the solution has been postulated by the techniques of Hilbert spaces. Röscher [8], views the characterization of boundary condition as an identification problem, and presents an iterative approach to meet the conditions of optimality.

Although the techniques of functional analysis suggest several solutions to the problem at hand [9], the presented technique here can be generalized to a variety of systems displaying arbitrarily complicated behavior. This fact is intimately related to the observation-based nature of the approach.

The way this paper differs from what have been appeared in the literature is to demonstrate how a low-order model having an external control input is derived. In the sequel, the details will be presented in the following organization: The second section presents the POD technique and its relevance to the modeling strategy. In the third section, development of the reduced order model for the 2D heat flow is analyzed. The fourth section presents the results and the concluding remarks are given at the end of the paper.

II. PROPER ORTHOGONAL DECOMPOSITION

Consider the ensemble $U_i(x, y)$ $i = 1, 2, \dots, N$. Apparently, every element of this set corresponds to a snapshot observed from a process, say for example, the 2D heat equation $u_t(x, y, t) = c^2(u_{xx}(x, y, t) + u_{yy}(x, y, t))$, where c is constant.

The goal is to find an orthogonal basis set letting us to write the solution as

$$u(x, y, t) \approx \sum_{i=1}^N \alpha_i(t) \phi_i(x, y), \quad (1)$$

where $\alpha_i(t)$ is the temporal part, and $\phi_i(x, y)$ is the spatial part. It will later be clear that if the basis set $\{\phi_i(x, y)\}_{i=1}^N$ is an orthogonal set, then the modeling task can exploit Galerkin projection technique.

Let us summarize the POD procedure.

Step 1. Start calculating the $N \times N$ correlation matrix L , the $(ij)^{th}$ entry of which is $L_{ij} := \langle U_i, U_j \rangle_\Omega$, where $\langle \cdot, \cdot \rangle_\Omega$ is the inner product operator defined over the spatial domain (Ω) of the process.

Step 2. Find the eigenvectors (denoted by v_i) and the associated eigenvalues (λ_i). Sort them in a descending order in terms of the magnitudes of λ_i . Note that every v_i is an $N \times 1$ vector satisfying $v_i^T v_i = \frac{1}{\lambda_i}$, here, for simplicity of the exposition we assume that the eigenvalues are distinct.

Step 3. Construct the basis set by using

$$\phi_i(x, y) = \sum_{j=1}^N v_{ij} U_j(x, y), \quad (2)$$

where v_{ij} is the j^{th} entry of the eigenvector v_i , and $i = 1, 2, \dots, \text{rank}(L)$. It can be shown that $\langle \phi_i(x, y), \phi_j(x, y) \rangle_\Omega = \delta_{ij}$ with δ_{ij} being the Kronecker delta function.

Remark 1. Notice that the basis functions are admixtures of the snapshots.

Step 4. Calculate the temporal coefficients. Taking the inner product of both sides of (1) with $\phi_i(x, y)$, the orthogonality lets us have

$$\begin{aligned} \alpha_i(t_0) &= \langle \phi_i(x, y), u(x, y, t_0) \rangle_\Omega \\ &= \langle \phi_i(x, y), U_{t_0} \rangle_\Omega, \end{aligned} \quad (3)$$

Without loss of generality, an element of the ensemble $\{U_i(x, y)\}_{i=1}^N$ may be $U(x, y, t_0)$. Therefore, to generate the temporal gain ($\alpha_i(t)$) of the spatial basis $\phi_i(x, y)$, one would take the inner product with the elements of the ensemble with sampled forms of the basis functions.

Remark 2. Notice that the following hold true:

$$\begin{aligned} \langle U_1(x, y), \phi_k(x, y) \rangle_\Omega &= \alpha_k(t_1) \\ \langle U_2(x, y), \phi_k(x, y) \rangle_\Omega &= \alpha_k(t_2) \\ &\vdots \\ \langle U_N(x, y), \phi_k(x, y) \rangle_\Omega &= \alpha_k(t_N) \end{aligned}$$

$$\begin{aligned} \sum_{i=1}^N \langle U_i(x, y), \phi_k(x, y) \rangle_\Omega^2 &= \sum_{i=1}^N \alpha_i^2(t_i) \\ &= \lambda_k \end{aligned}$$

Apparently, the temporal coefficients satisfy orthogonality properties over $t \in \{t_1, t_2, \dots, t_N\}$

In the next section, we demonstrate how the boundary condition is transformed to an explicit control input in the ODEs.

III. REDUCED ORDER MODELING

Consider the heat conduction problem over the spatial domain $\Omega := (x, y) \in [0, 1] \times [0, 1]$. Truncate the sum in (1) after some certain integer, say M , and assume that the modes from $M + 1$ to ∞ are reasonably small in magnitude such that we can write the following:

$$u(x, y, t) = \sum_{i=1}^M \alpha_i(t) \phi_i(x, y). \quad (4)$$

The solution above must satisfy the PDE, i.e. we get

$$\sum_{i=1}^M \dot{\alpha}_i(t) \phi_i(x, y) = c^2 \sum_{i=1}^M \alpha_i(t) \zeta_i(x, y) \quad (5)$$

where $\zeta_i(x, y) = \phi_i(x, y)_{xx} + \phi_i(x, y)_{yy}$. Clearly taking the inner product of both sides with $\phi_k(x, y)$ yields

$$\dot{\alpha}_k = c^2 \sum_{i=1}^M \alpha_i(t) \langle \phi_k(x, y), \zeta_i(x, y) \rangle_\Omega, \quad (6)$$

where $k = 1, 2, \dots, M$. Apparently no matter what the boundary conditions, the above representation contains the

effect of them implicitly. To overcome this problem, denote the set of points, at which boundary excitations are independently specified, by $\partial\Omega$. The underlying idea in developing the dynamic model is the following fact:

$$\langle \phi_k(x, y), \zeta_i(x, y) \rangle_{\Omega} = \langle \phi_k(x, y), \zeta_i(x, y) \rangle_{\partial\Omega} + \langle \phi_k(x, y), \zeta_i(x, y) \rangle_{\Omega \setminus \partial\Omega}, \quad (7)$$

which basically corresponds to the repartitioning of the domain by changing limits of a Riemann integral computed over a non-intersecting subdomains embodying the domain of the original integral when they are united.

Define $\partial\Omega := \{(0, 0), (0, 1), (1, 0), (1, 1)\}$. This means that the external stimuli enter from the corners of the domain. In other words, the dynamic input will have four inputs denoted by $\gamma_{00}, \gamma_{01}, \gamma_{10}$ and γ_{11} . Clearly, in this notation the first index stands for x and the second stands for y .

Denote the solution observed when $\gamma_{00}(t) = g_{00}(t)$ with other inputs are equal to zero as $u^{00}(x, y, t)$ with $g_{00}(t)$ being an arbitrarily chosen test signal. Perform the POD procedure to obtain a solution in the form of (4). Repeating the same procedure till every element of the set $\partial\Omega$ is contained, one would end up with the following solution when each input assumes arbitrarily chosen values,

$$u(x, y, t) = \sum_{p=0}^1 \sum_{q=0}^1 \sum_{i=1}^M \alpha_i^{pq}(t) \phi_i^{pq}(x, y), \quad (8)$$

which is the result of the superposition principle. Note that for nonlinear PDEs, this expansion would not be valid.

Now assume that the system is excited only from $(x, y) = (0, 0)$ entry. Since the solution in (4) must be satisfied also on $\partial\Omega$, this lets us write the following,

$$\begin{aligned} u(0, 0, t) &= u^{00}(0, 0, t) \\ &= \gamma_{00}(t) \\ &= \sum_{i=1}^M \alpha_i^{00}(t) \phi_i^{00}(0, 0) \\ &= \alpha_k^{00}(t) \phi_k^{00}(0, 0) \\ &\quad + \sum_{i=1}^M (1 - \delta_{ik}) \alpha_i^{00}(t) \phi_i^{00}(0, 0), \end{aligned} \quad (9)$$

$$\begin{aligned} \alpha_k^{00}(t) \phi_k^{00}(0, 0) \zeta_k^{00}(0, 0) &= \gamma_{00}(t) \zeta_k^{00}(0, 0) \\ &\quad - \sum_{i=1}^M (1 - \delta_{ik}) \alpha_i^{00}(t) \phi_i^{00}(0, 0) \zeta_k^{00}(0, 0) \end{aligned} \quad (10)$$

In a similar fashion, one should investigate the spatial gain associated with other inputs, i.e. $(x, y) \in \{(0, 1), (1, 0), (1, 1)\}$. Performing this would simply let us generalize (10) as follows:

$$\begin{aligned} \alpha_k^{00}(t) \phi_k^{00}(p, q) \zeta_k^{00}(p, q) &= \gamma_{pq}(t) \zeta_k^{00}(p, q) \\ &\quad - \sum_{i=1}^M (1 - \delta_{ik}) \alpha_i^{00}(t) \phi_i^{00}(p, q) \zeta_k^{00}(p, q), \end{aligned} \quad (11)$$

where p and q assume values either zero or one. Notice that the manipulations in (9)-(11) are for the single excitation from $(x, y) = (0, 0)$. Therefore if $p = 0, q = 0$, the expressions (10) and (11) become identical. However, other combinations of p and q stipulate that $\gamma_{pq}(t)$ may have a nonzero spatial gain (See (11)) although the $\gamma_{pq}(t)$ is identically equal to zero.

Considering (7) in (6) and using the above equation for every p and q would result in the dynamical system corresponding to the described excitation configuration.

$$\dot{\alpha}_k^{00}(t) = \sum_{i=1}^M A_{ki}^{00} \alpha_i^{00}(t) + c^2 \sum_{p=0}^1 \sum_{q=0}^1 \gamma_{pq}(t) \zeta_k^{00}(p, q), \quad (12)$$

where

$$\begin{aligned} A_{ki}^{00} &= c^2 \langle \phi_k^{00}(x, y), \zeta_i^{00}(x, y) \rangle_{\Omega} \\ &\quad - c^2 \sum_{p=0}^1 \sum_{q=0}^1 \phi_i^{00}(p, q) \zeta_k^{00}(p, q). \end{aligned}$$

The model above can be written in the well known state space form,

$$\begin{aligned} \dot{\underline{\alpha}}^{00}(t) &= A^{00} \underline{\alpha}^{00}(t) + B^{00} \underline{\gamma}(t) \\ u^{00}(x, y, t) &= \underline{C}^{00}(x, y) \underline{\alpha}^{00}(t), \end{aligned} \quad (13)$$

where $\underline{\gamma} = (\gamma_{00} \ \gamma_{01} \ \gamma_{10} \ \gamma_{11})^T$, and $B_{ki}^{00} = c^2 \zeta_k^{00}(p, q)$ with $i = 2p + q + 1$. The row vector $\underline{C}^{00}(x, y)$ is apparently the vector of 2D basis functions calculated at given points x and y and for given excitation conditions, i.e. 00 for this case.

Clearly, the above model has M states and four inputs. In order to obtain the full representation, the procedure described in (9)-(13) must be repeated for every $\gamma_{pq}(t)$ to obtain the individual components of (8).

Remark 3. Considering (12) and (13), it is apparent how the dynamical model is affected from the uncertainties on the thermal diffusivity parameter c .

Let $U(x, y, s) := \mathcal{L} \{u(x, y, s)\}$, $U^{pq}(x, y, s) := \mathcal{L} \{u^{pq}(x, y, s)\}$ and $\Gamma_{pq}(s) := \mathcal{L} \{\gamma_{pq}(s)\}$ with \mathcal{L} being the Laplace transform operator. In the view of the analysis presented, the total system can be expressed as follows

$$U(x, y, s) = \sum_{m=0}^1 \sum_{n=0}^1 U^{mn}(x, y, s), \quad (14)$$

where

$$U^{mn}(x, y, s) = \sum_{p=0}^1 \sum_{q=0}^1 G_{pq}^{mn}(x, y, s) \Gamma_{pq}(s). \quad (15)$$

IV. MODELING RESULTS

According to the described procedure, several tests have been done. The PDE has been solved by using Crank-Nicholson method (See [10] for details), with a step size of 5 msec. The initial thermal distribution is taken as zero everywhere and the thermal diffusivity constant is set as $c = 0.4$. In order to form the solution, a linear grid having $N_x = N_y = 25$ points in x -direction and y -direction respectively. According to the above parameter values, a total of 201 snapshots embody the entire numerical solution, among which a linearly sampled $N = 21$ snapshots have been used for the POD scheme. Although one may use the entire set of snapshots, it has been shown that a reasonably descriptive subset of them can be used. In the literature, this approach is called *method of snapshots*, which significantly reduce the computational intensity of the overall scheme, [1-2]. Once the modes have been obtained, we truncated the solution at $M = 5$, which represents (in average)

According to the above settings, the inner product seen in (6) is calculated as follows:

$$\langle \phi_k(x, y), \zeta_i(x, y) \rangle_{\Omega} = \frac{1}{N} \sum_{l=1}^{N_x} \sum_{j=1}^{N_y} \phi_k(x_l, y_j) \zeta_i(x_l, y_j) \quad (16)$$

in which we calculate the derivatives numerically.

The system has been stimulated from every corner, with one-at-a-time basis. For collecting data, we used several test signals given as

$$\begin{aligned} g_{pq}(t) &= 0.1 + 2 \exp(-6t) \sin(8\pi t) \\ g_{pq}(t) &= \sin(2\pi t) \\ g_{pq}(t) &= 1 - \exp(-6t) \\ g_{pq}(t) &= t, \end{aligned} \quad (17)$$

where subscript pq indicates the chosen corner. For each case, the basis surfaces change very slightly and this influences the content of the matrices in (13). Checking the eigenvalues of the matrix A is an indicator of the similarity between each case. This is illustrated in Figure 1, in which the eigenvalues form individual clusters.

In order to construct the dynamic model, the matrices obtained for different test conditions at a chosen corner have been averaged. To compare the overall solution with the numerically obtained data, we choose the following boundary excitations, which are applied simultaneously:

$$\begin{aligned} \gamma_{00}(t) &= \sin(2\pi t) \\ \gamma_{01}(t) &= -\sin(3\pi t) \\ \gamma_{10}(t) &= \text{sgn}(\sin(2\pi t)) \\ \gamma_{11}(t) &= 1 - \exp(-6t). \end{aligned} \quad (18)$$

The results for this case are depicted in Figure 2. Every pair in this figure represent the numerical solution (on the left) and

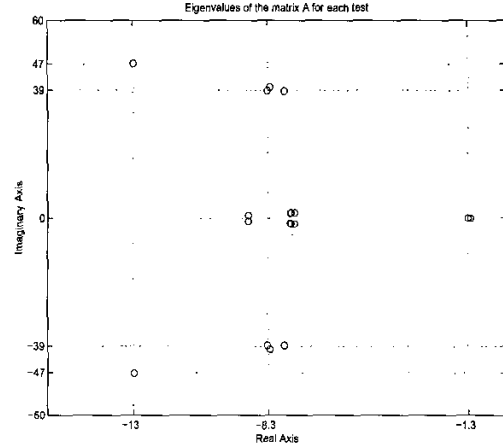


Fig. 1. Clusters of eigenvalues appear as the procedure is repeated for different test conditions. A test signal is chosen and an input point is chosen to complete one experiment. The figure depicts 16 experiments and 80 data points.

the reconstructed solution (on the right). The pairs in the first line stand for $t = 0.1$ sec. (on the left) and $t = 0.2$ sec. (on the right) and the rest includes solutions sampled at the same rate till $t = 1$ sec.

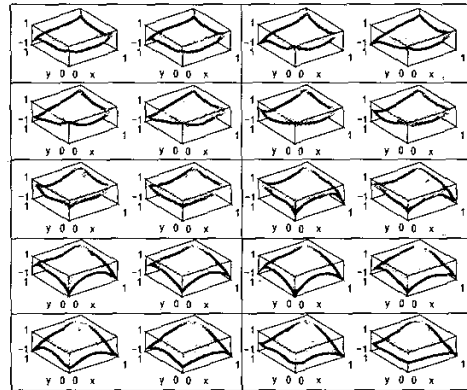


Fig. 2. A comparison of the solutions generated by the numerical solver and the low-order dynamic model. The frequency ranges of the test conditions and the validation are similar.

It is clear from Figure 2 that the surfaces at every cell are reasonably similar to each other. This result confirms the analytical claims and validates the method described to separate the effect of control terms.

Remark 4. It should be noted that the modes of the infinite dimensional system have been excited by a set of test signals (17), and the result has been validated on a similar set of boundary conditions (19). The answer to the immediate question "why" is as follows: The model is valid dominantly for the frequency range covered by the test signals. If one wants to cover a wide range of frequency spectrum, then this would first require smaller step size for the PDE solver and much longer time to obtain a representative set of snapshots.

On the other hand, the POD procedure will filter out some of the high frequency content. The reader should notice that the underlying idea of the POD procedure is to reconstruct the solution from its samples (snapshots), however, a good reconstruction can be performed only when the solution is dominated by coherent modes, which implies relatively low frequency content of the entire solution. A good comparison is presented in Figure 4 for the following boundary conditions:

$$\begin{aligned}\gamma_{00}(t) &= \sin(20\pi t) \\ \gamma_{01}(t) &= \sin(2\pi t) \\ \gamma_{10}(t) &= \text{sgn}(\sin(14\pi t)) \\ \gamma_{11}(t) &= -\sin(30\pi t),\end{aligned}\quad (19)$$

in which the frequency range is increased approximately about ten times for all corners except $\gamma_{01}(t)$, which is kept in the frequency range of the modeling stage so that the difference will be clear. Referring to Figure 3, it is apparent that the evolution of the corner signals are so fast that the reduced order model cannot reconstruct the solution around them, however around the corner $\gamma_{10}(t)$, the reconstruction is unsurprisingly good. This result simply validates our remarks.

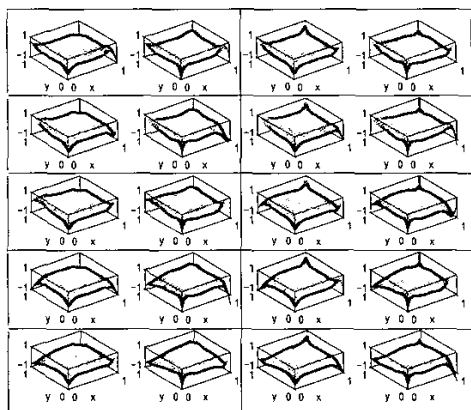


Fig. 3. A comparison of the solutions generated by the numerical solver and the low-order dynamic model. The frequency ranges of the test conditions and the validation are not similar.

Remark 5. The dissimilarity observed in the second case raises the following question: Does the percent energy described earlier depend upon the frequency content of the test signals? The answer is apparently yes. Basically, a model is valid in the frequency range of data that leads to the model. The degree of confidence decreases as the overlapping between the model derivation conditions and the operating conditions becomes dissimilar. Alternatively, one can interpret this as follows: Fixing the number of modes (M) (or equivalently fixing the energy level) cannot lead to the reconstruction of the information hidden in the modes $M + 1$ to N .

V. CONCLUSIONS

Many physical phenomena, such as aerodynamic flows, flexible object dynamics, and heat transport have infinitely many

dynamical modes describing the entire behavior when they are considered collectively. In this paper, the reduced order modeling of 2D heat equation is considered. The physical domain is a square plate, and the excitation enters into the system from the corners of the plate. The algorithmic approach has been shown to be capable of capturing the dynamically significant modes of the solution. Once the number of modes is decided, the procedure yields a set of autonomous ODEs. We present a method to separate the external stimuli from these ODEs. The described approach has been shown to be successful in terms of the similarity between the numerical solution and the solution based on the low-dimensional model. We observed that the widely accepted form of the realization performance, which is the percent energy captured is a frequency-variable quantity, and it does not constitute a globally valid measure of realization performance. As long as the data is a representative for the operating conditions, the results are highly promising in the sense of extending the modeling and separation techniques for more difficult and nonlinear cases, such as for aerodynamic flows, which constitute the ultimate goal of this research.

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