

near the center (where there is a high probability of finding the electron and hole at the same relative position). Next ($t = 250$ fs), interference between the excitonic and free-carrier WP is seen. Because the WP is highly anisotropic, there is a net dipole moment that results in the THz spectrum shown in Figure 2. At later times ($t = 400, 640$ fs), side lobes can be seen in the WP; these are formed by the combination of slow transverse spreading, the relatively fast field driven motion in the polarization direction, and the excitonic attraction (Coulombic rescattering). The resulting WPs for the $\phi = \pi/2$ case are significantly different.⁵

In conclusion, we have theoretically investigated the simultaneous exposure of an $\text{In}_x\text{Ga}_{1-x}\text{As}$ QW to a broadband optical pulse and a 1 THz driving field. The scenario produces interfering, electron-hole WPs whose anisotropic beating structure can be manipulated to achieve harmonic-emission in the THz regime. This technique may prove to be a useful means of generating tunable THz transients in the laboratory. Moreover, enhanced tunability has been very recently demonstrated via the selective excitation of spatially chosen WPs using two-pulse nondegenerate optical pulses.⁷

Acknowledgments

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Quarter Wave Fresnel Rhomb with Colinear Output Beam

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Abstract

A standard quarter wave ($\lambda/4$) Fresnel rhomb laterally displaces the output beam with respect to the input beam. This lateral displacement can cause problems when inserting the rhomb in an already aligned optical setup. This note describes the use of two $\lambda/8$ rhombs to construct a $\lambda/4$ rhomb. The output of the proposed rhomb is colinear with the input beam, and allows a much more flexible optical alignment.

Introduction

Achromatic quarter wave ($\lambda/4$) or half wave ($\lambda/2$) retarders are manufactured using Fresnel rhombs.¹ The standard $\lambda/4$ rhomb creates a phase shift (δ) of 90° between the parallel ($//$) and the perpendicular (\perp) components of an incident polarized beam upon total internal reflection (TIR). If properly aligned, it can convert a linearly polarized beam to a circularly polarized beam and vice versa. The output of the $\lambda/4$ rhomb is displaced laterally with respect to the input beam. This lateral displacement causes alignment problems, when inserting the rhomb in an already aligned optical setup.

The standard $\lambda/2$ rhomb creates a phase shift (δ) of 180° between the $//$ and \perp components of an incident polarized beam. If properly aligned, it can rotate a linearly polarized beam by 90° . The standard $\lambda/2$ rhomb is composed of two $\lambda/4$ rhombs. The output of the $\lambda/2$ rhomb is colinear with the input beam, making the optical alignment much more flexible.

This note proposes a $\lambda/8$ Fresnel rhomb that creates a phase shift (δ) of 45° between the $//$ and \perp components of an incident polarized beam. The proposed $\lambda/4$ Fresnel rhomb is composed of two $\lambda/8$ rhombs. The output of the proposed $\lambda/4$ rhomb is colinear with the input beam, making the optical alignment much more flexible.

Calculation of the phase shift

The working principle of the Fresnel rhomb is based on the phase shift between the $//$ and \perp components of an incident polarized beam upon TIR. This phase shift can be calculated using

$$\delta = 2 \tan^{-1} \left\{ \frac{\cos(\theta_i) [\sin^2(\theta_i) - n^2]^{1/2}}{\sin^2(\theta_i)} \right\}, \quad (1)$$

where θ_i is the angle of incidence and n is the refractive index of the rhomb.

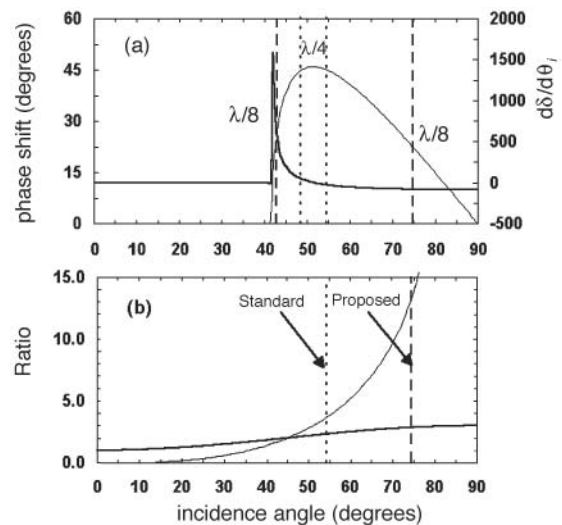


Figure 1. (a) Calculated phase shift (thin line) and its derivative (thick line) as a function of the incidence angle for a glass rhomb ($n = 1.51$). (b) Calculated L/A (thin line) and H/A (thick line) ratios as a function of the incidence angle.

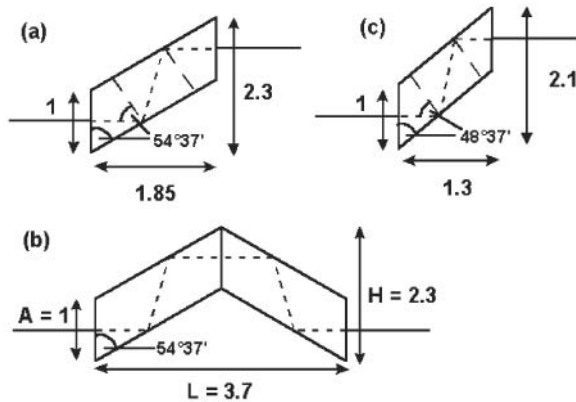


Figure 2. (a) Optical schematic of the (a) standard $\lambda/4$, (b) standard $\lambda/2$, and (c) unrealized $\lambda/4$ rhombs.

Figure 1a (previous page) shows the phase shift, δ , between the \parallel and \perp polarized waves as a function of the incidence angle for glass ($n = 1.51$) calculated using Eq. 1. Figure 1a also shows the first derivative of δ with respect to the incidence angle ($d\delta/d\theta_i$) to estimate the variation of δ for small angular misalignments, with larger variations corresponding to smaller angular tolerances.

Calculation of the physical dimensions

Figure 1b shows the length/aperture (L/A) and the height/aperture (H/A) ratios of a double rhomb calculated using

$$L/A = 4\sin^2(\theta_i)\tan(\theta_i), \tag{2a}$$

$$H/A = 1 + 2\sin^2(\theta_i). \tag{2b}$$

The standard $\lambda/4$ and $\lambda/2$ rhombs

As indicated by the dotted lines labeled $\lambda/4$ in Figure 1a, a phase shift of $\lambda = 45^\circ$ occurs for both $(\theta_i) = 48^\circ37'$ (with $d\delta/d\theta_i = 53$) and $\theta_i = 54^\circ37'$ (with $d\delta/d\theta_i = -27$). Since the angular tolerance of $\theta_i = 54^\circ37'$ is higher than $\theta_i = 48^\circ37'$, the standard $\lambda/4$ rhomb uses two $\delta = 45^\circ$ TIR at $\theta_i = 54^\circ37'$. The standard $\lambda/4$ rhomb is shown in Figure 2a. The rhomb angle is equal to the incidence angle of $\theta_i = 54^\circ37'$. The output beam is displaced laterally, parallel to the input beam. As mentioned previously, this lateral displacement creates difficulties in optical alignment. For a standard $\lambda/4$ rhomb, a typical aperture size of 1 cm requires a length of 1.85 cm

and a height of 2.3 cm, as indicated by the dotted “standard” line at $\theta_i = 54^\circ37'$ (Fig. 1b).

Figure 2c describes the steeper $\lambda/4$ rhomb with $\theta_i = 48^\circ37'$, a typical aperture of 1 cm, shorter length of 1.3 cm, and height of 2.1 cm, which would fit into a crowded system easier. But since the angular tolerance of $\theta_i = 48^\circ37'$ (with $d\delta/d\theta_i = 53$) is smaller than $\theta_i = 54^\circ37'$ (with $d\delta/d\theta_i = -27$), the $\theta_i = 48^\circ37'$ configuration is not realized.

The standard $\lambda/2$ rhomb is shown in Figure 2b and is composed of two $\lambda/4$ rhombs. Four reflections within a pair of $\lambda/4$ rhombs produces a $\lambda/2$ retarder. The output beam is colinear with the input beam, and the optical alignment is very flexible. For a standard $\lambda/2$ rhomb, a typical aperture size of 1 cm requires a length of 3.7 cm and a height of 2.3 cm, as indicated by the dotted “standard” line at $\theta_i = 54^\circ37'$ (see Fig. 1b).

The proposed $\lambda/8$ and $\lambda/4$ rhombs

As indicated by the dashed lines labeled $\lambda/8$ in Figure 1a, a phase shift of $\delta = 22.5^\circ$ occurs at $\theta_i = 42^\circ20'$ (with $d\delta/d\theta_i = 783$) and $\theta_i = 74^\circ42'$ (with $d\delta/d\theta_i = -81$). Since the angular tolerance of $\theta_i = 74^\circ42'$ is higher than $\theta_i = 42^\circ20'$, the proposed $\lambda/8$ rhomb uses two $\delta = 22.5^\circ$ TIR at $\theta_i = 74^\circ42'$ (see Fig. 3a). The rhomb angle is equal to the incidence angle of $\theta_i = 74^\circ42'$. The output beam is displaced laterally, parallel to the input beam. For the proposed $\lambda/8$ rhomb, a typical aperture size of 1 cm requires a length of 6.8 cm and a height of 2.9 cm, as indicated by the dashed “proposed” line at $\theta_i = 74^\circ42'$ (see Fig. 1b).

Figure 3c describes the steeper $\lambda/8$ rhomb with $\theta_i = 42^\circ20'$, a typical aperture of 1 cm, shorter length of 0.8 cm, and a height of 1.9 cm, which would fit into a crowded system easier. But since the angular tolerance of the $\theta_i = 42^\circ20'$ (with $d\delta/d\theta_i = 783$) is smaller than the angular tolerance of $\theta_i = 74^\circ42'$ (with $d\delta/d\theta_i = -81$), the $\theta_i = 42^\circ20'$ configuration is not explored. Also since $42^\circ20'$ is smaller than 45° , the $\theta_i = 42^\circ20'$ configuration allows the transmission of direct rays as indicated in Figure 3c.

The proposed $\lambda/4$ rhomb is shown in Figure 3b and is composed of two $\lambda/8$ rhombs. The output beam is colinear with the input beam (similar to the case of the standard $\lambda/2$ rhomb) making optical alignment very flexible. For the proposed $\lambda/4$ rhomb, a typical aperture size of 1 cm requires a length of 13.6 cm and a height of 2.9 cm, as indicated by the dashed “proposed” line at $\theta_i = 74^\circ42'$ (see Fig. 1b). For the same aperture size the proposed $\lambda/4$ rhomb with colinear output beam has 7.4 times the length and 1.3x the height of the standard $\lambda/4$ rhomb. The angular tolerance of the proposed $\lambda/4$ rhomb (with $d\delta/d\theta_i = -81$) is only 3x lower than the standard $\lambda/4$ rhomb (with $d\delta/d\theta_i = -27$). Therefore, for applications where colinearity is an important factor, the proposed $\lambda/4$ rhomb is a good candidate for replacing the standard $\lambda/4$ rhomb.

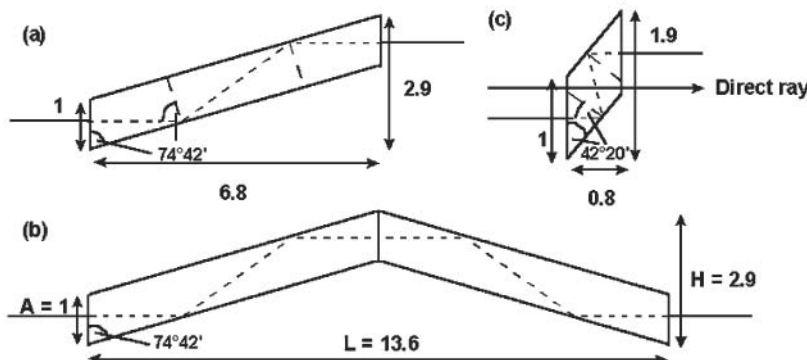


Figure 3. (a) Optical schematic of the (a) proposed $\lambda/8$, (b) proposed $\lambda/4$, and (c) unexplored $\lambda/8$ rhombs.

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